A NEW PERSPECTIVE ON COASTALLY TRAPPED DISTURBANCES USING DATA FROM THE SATELLITE ERA

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Satellite observations reveal that large-scale changes in the Pacific high may yield polluted coastal marine clouds and fogs associated with coastally trapped disturbances.

Coastal marine environments are regions of significant meteorological interest as a result of the direct impact of atmospheric processes on, for example, commerce and trade, naval operations, and civilian activities. Low clouds and fogs, in particular, undoubtedly affect these sectors. Forecasting the life cycle (initiation, maintenance, and dissipation) and spatial and vertical coverage of these ubiquitous coastal clouds remains an outstanding challenge due to the inherently complex relationship between the ocean–land–atmosphere system (Boucher et al. 2013; Koraĉin and Dorman 2017).

Considerable effort within the atmospheric science community focuses on marine boundary layer (MBL) stratiform clouds because their radiative response to changes in the climate system is substantial yet inconsistent between various global circulation models (e.g., Palmer and Anderson 1994; Delecluse et al. 1998; Bachiochi and Krishnamurti 2000; Bony and Dufresne 2005; Webb et al. 2006). Because their albedo (30%–70%; Randall et al. 1984; Koraĉin and Dorman 2017) is relatively high compared to the ocean surface (10%) and their temperature is nearly equal to that of the ocean surface, these cloud types exert an average shortwave (SW) radiative forcing of approximately 60–120 W m^{-2} (cooling effect; e.g., Hartmann et al. 1992; Zhang and Li 2013), although the magnitude is dependent on both macrophysical and microphysical properties (e.g., Twomey 1977; Albrecht 1989; Twohy et al. 1995, 2005; Wood 2012).

Low, marine clouds are expansive, as their global coverage is about one-third at any time and up to ~60% over the subtropic and midlatitude oceanic

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The abstract for this article can be found in this issue, following the table of contents.
DOI:10.1175/BAMS-D-18-0002.1

In final form 30 August 2018

regions near California, Peru, and Namibia during summer months (e.g., Klein and Hartmann 1993; Jiang et al. 2014). In the present study, we focus on the northeastern Pacific Ocean.

During the boreal warm season, the combination of the Pacific high and thermal low in the desert southwest of the United States leads to an MBL that slopes downward toward the east and supports northerly flow (e.g., Beardsley et al. 1987; Zemba and Friehe 1987; Parish 2000). A coastal jet sits at the base of a strong subsidence inversion (oftentimes 10°–20°C) that separates the cool, moist MBL from the warm, dry free troposphere above. Several times during each summer month, the strong, northerly winds weaken (also termed "relax"; e.g., Winant et al. 1987; Melton et al. 2009) and even reverse (e.g., Bond et al. 1996; Mass and Bond 1996).

The synoptic-scale conditions that support reversal events are marked by a migration and elongation of the Pacific high toward the northwestern United States (e.g., Mass and Bond 1996; Nuss et al. 2000). As a result, the offshore flow aloft strengthens, the MBL thins, and part of the cloud deck clears [see, e.g., Kloesel (1992) and Crosbie et al. (2016) for discussion of cloud clearings]. A reversal of the alongshore pressure gradient along the Northern California coastline occurs for two reasons: i) air descends from higher terrain over land to sea level and warms adiabatically and ii) warm air over land advects toward the MBL. Hence, the resultant pressure field preconditions the environment and becomes conducive for southerly flow. Bond et al. (1996) and Mass and Bond (1996) emphasize that the alongshore synoptic pressure gradient must deviate from climatology for a reversal to exist; that is, a southward-directed pressure gradient is required.

In the literature, these events are called wind reversals, southerly surges, coastally trapped wind reversals, or coastally trapped disturbances (CTDs). Here, we refer to them as CTDs. In the summer, CTDs occur along not only the western coast of the United States, but also the western coasts of South America (e.g., Garreaud et al. 2002; Garreaud and Rutllant 2003) and southern Africa (e.g., Gill 1977; Reason and Jury 1990), in addition to the southeastern coast of Australia (e.g., Holland and Leslie 1986; Reid and Leslie 1999; Reason et al. 1999). These phenomena are considered coastally trapped within $O(100)$ km of the coastline as determined by the Rossby radius of deformation (e.g., Bond et al. 1996; Ralph et al. 1998).

The dynamics of CTDs have been debated in the literature over the last several decades, with a thorough review given by Nuss et al. (2000). Interpretations include a Kelvin wave (e.g., Dorman 1985), a topographic Rossby wave (Skamarock et al. 1999), a topographically trapped density current (e.g., Mass and Albright 1987), a combination Kelvin wave–bore (Ralph et al. 2000), and an ageostrophic response to the alongshore pressure gradient (e.g., Mass and Bond 1996). Regardless of the primary dynamical mechanism, observations and modeling results show that a broad region of stratus clouds and fogs often accompany these disturbances (e.g., Bond et al. 1996; Thompson et al. 2005; Rahn and Parish 2008). The microphysical properties of these clouds and fogs have received little attention in the literature and form the focus of the current study.

Figure 1 shows the temporal evolution of the MBL cloud field during the 22–25 June 2006 CTD case. Radar observations from this particular event are presented in Parish (2008) and suggest that drizzle processes (accretion and self-collection) may dominate over condensational growth of cloud droplets

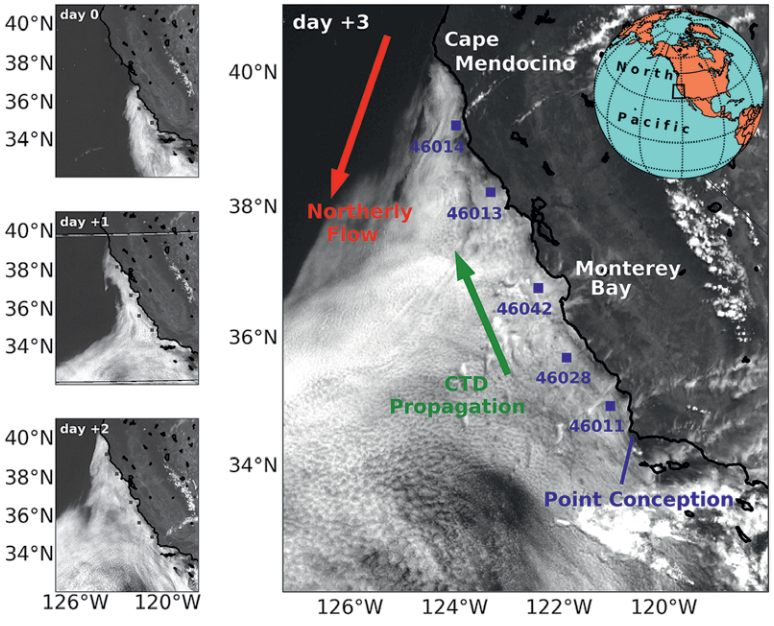

Fig. 1. Visible satellite imagery from *GOES-11* showing the temporal evolution of an example CTD cloud deck from 22 to 25 Jun 2006 (day 0 to day +3). The day +3 panel displays select topographic locations in addition to the locations of buoy sites used in the analysis. Buoy locations are also shown for days 0, +1, and +2.

in localized cells; marine stratocumulus that are not associated with CTDs are known to show similar structure (e.g., Wood 2005; vanZanten et al. 2005).

A modeling investigation of the 15–16 June 2000 CTD case conducted by Thompson et al. (2005) finds that MBL depth and cloud extent are a strong function of buoyancy generation resulting from cloud-top radiative flux divergence and upward sea surface heat fluxes. Moreover, the wind shift at the leading edge of the CTD forces convergence and rising motion while preceding the stratus tongue and playing an important role in the distribution of cloud liquid water. The cloud deck to the south is found to lift off the surface and decouple due to entrainment of dry air resulting from Kelvin–Helmholtz instability at cloud top.

To the authors' best knowledge, however, the literature does not explicitly report comprehensively on aerosol–cloud and air–sea interactions in the context of CTDs. Therefore, the results from this study could provide a valuable test bed for better understanding and predicting aerosol–cloud–precipitation interactions in marine stratocumulus for both small- and large-scale modeling efforts.

Table 1. List of CTD events with corresponding buoy data availability. For each event, the first day refers to the initiation day (day 0). The data at each NDBC site are valid for a particular event (denoted by X) if the wind has a positive alongshore component for at least six consecutive hours on day 0. (When the threshold is increased to 10 consecutive hours, only buoys 46011 and 46014 are sensitive, most likely due to their geographical positions along the coastline.) Furthermore, if a buoy is missing data for one hour, then the meteorological variables are linearly interpolated. However, if a buoy is missing data for two or more consecutive hours, then the data for that particular event are considered invalid. These criteria are similar to those used in Bond et al. (1996).

Year	Dates	46011	46028	46042	46013	46014
2004	16–18 Jun	X	X	X	X	X
2006	22–25 Jun	X	X	X	X	
2007	6–7 Sep	X		X	X	X
2008	12–13 Jun	X	X	X	X	X
	27–28 Jun	X	X	X	X	X
	8–13 Jul	X	X	X	X	X
2009	29 Jun–3 Jul	X	X	X	X	X
2010	4–8 Jul	X	X	X	X	X
2011	27–28 Jul	X	X	X	X	X
	1–5 Sep	X	X	X	X	X
	10–11 Sep	X	X	X	X	X
	28–29 Sep	X	X	X	X	X
2012	11–13 Sep			X	X	X
2013	3–5 Jun		X	X	X	X
2014	9–10 Jun		X	X	X	
	30 Jun–2 Jul		X	X	X	X
	3–7 Sep		X	X	X	X
2015	16–18 Jul	X			X	X
	28–30 Jul			X	X	
	17–19 Aug			X	X	
2016	15–16 Jul	X	X	X		
	17–19 Aug		X	X	X	X
	7–10 Sep	X	X	X		

EVENT SELECTION. The research presented here builds upon the 1981 to 1991 CTD climatology outlined in Bond et al. (1996) and Mass and Bond (1996) by considering the warm season (June–September) from 2004 to 2016. Potential CTD cases are first identified using wind data from the National Data Buoy Center (NDBC)—specifically, from buoy 46013 (Bodega Bay; see Fig. 1). The list of events is then filtered to remove any landfalling cyclones, consistent with Bond et al. (1996). Visible satellite images are used in this filtering step. The *Aqua* Moderate Resolution Imaging Spectroradiometer (MODIS) imagery is interrogated qualitatively for the remaining ~60 events to ensure that the instrument could clearly observe each CTD cloud deck propagate northward over a period lasting at least 2 days with little to no contamination from clouds originating from other regions. After employing these techniques, a total of 23 cases are identified as clear CTDs with associated cloud fields.

The set of 23 cases is listed in Table 1. Five NDBC stations along the California coastline are chosen to

confirm the presence of a CTD and to characterize the surface meteorology during the reversal. These buoys, from south to north, are 46011 (Santa Maria), 46028 (Cape San Martin), 46042 (Monterey), 46013 (Bodega Bay), and 46014 (Point Arena). The total number of valid events (see Table 1) at buoys 46011, 46028, 46042, 46013, and 46014 are 15, 18, 22, 21, and 18, respectively. Each event has valid data from at least two buoys: two events have valid data from two buoys, five events from three buoys, six events from four buoys, and 10 events from all five buoys. Using the buoy data, we identify day 0 as the day on which the wind reversal begins at buoy 46011 (see Table 1 for additional details).

REMOTE SENSING MEASUREMENTS. Data from 1-km level 2 *Aqua* MODIS retrievals allow for a new, comprehensive, and quantitative perspective on the cloud physical properties of the 23 CTD cases. Satellite passes over the region of interest occur in the early afternoon hours, typically between ~2030 and 2200 UTC (~1330 and 1500 local time). Any two MODIS swath scenes that are consecutively retrieved are stitched together. The MODIS images are filtered to eliminate regions with missing data, overlapping clouds, and cloud-top temperature less than or equal to 0°C.

Table 2. List of non-CTD events analyzed using MODIS data and compared with CTD events.

Year	Dates
2004	9–11 Jun
2005	23 Jun–6 Jul
	27–31 Jul
	16–21 Aug
2006	5–9 Jun
	2–4 Jul
2007	23–26 Jul
2008	17–21 Jul
	3–9 Aug
	22–24 Aug
2009	8–13 Jun
	20–30 Jul
2010	1–12 Aug
2011	7–13 Jun
	11–16 Jul
	3–7 Aug
	19–22 Aug
2012	5–15 Jul
2013	6–9 Jul
2014	2–4 Jun
	17–20 Aug
2015	31 Aug–4 Sep
2016	1–8 Jul

The number of MODIS retrievals considered for each CTD case is dependent upon the duration of the case and the amount of contamination by clouds external to the CTD. Therefore, a total of 75 retrievals comprise the CTD cases. To compare with these cases, we use visible satellite images to compile a set of 23 non-CTD stratus and stratocumulus cases[1] (Table 2). A total of 141 retrievals comprise the non-CTD cases. For both the CTD and non-CTD composites, exactly one retrieval is used to represent the cloud field for each day.

Two key variables, cloud droplet effective radius[2] r_e and optical thickness[3] τ, are retrieved by MODIS utilizing a bispectral solar reflectance method (Nakajima and King 1990). The 3.7-μm channel is used for the r_e and τ retrievals because it best represents the actual value of r_e at cloud top (Platnick 2000; Rausch et al. 2017).

The liquid water path (LWP) is then inferred from r_e and τ and given by the relationship LWP = $C\rho_l r_e \tau$, where C is a function of the assumed vertical distribution of cloud liquid water and ρ_l is the density of liquid water (Miller et al. 2016). We assume an approximately adiabatic (C = 5/9) cloud vertical profile whereby the cloud liquid water content is expected to increase approximately linearly with height (e.g., Wood and Hartmann 2006). Also, the cloud droplet number concentration is assumed to be roughly constant with height. Observational studies suggest that assuming an adiabatic or subadiabatic profile yields more realistic LWP results than when assuming a vertically homogeneous (C = 6/9) cloud profile (e.g., Brenguier et al. 2000; Wood 2005).

Data from the Advanced Scatterometer (ASCAT) aboard the European Organisation for the Exploitation of Meteorological Satellites (EUMETSAT) *MetOp-A* satellite are obtained from Remote Sensing Systems (RSS) and provide wind measurements over the ocean at 12.5-km resolution. The microwave scatterometer operates at 5.255 GHz (~5.705 cm; C band) and relates backscattered power to sea surface roughness, which correlates well with near-surface wind speed and direction. Optimally interpolated sea surface temperature (SST) data are also sourced from RSS and use the scheme developed in Reynolds and Smith (1994). The interpolated dataset utilized in this study combines retrieval methods at microwave and infrared wavelengths to yield high-resolution (~9 km) SSTs under both clear and cloudy sky conditions.

[1] These cases are carefully selected to ensure the presence of a northerly flow regime under relatively quiescent synoptic conditions. We choose time periods with fairly persistent and continuous cloud decks because these qualities also characterize CTD clouds.

[2] Cloud droplet effective radius r_e is defined as the ratio of the third moment to the second moment of the cloud droplet size distribution: $r_e = \int_0^\infty n(r) r^3 dr / \int_0^\infty n(r) r^2 dr$.

[3] Optical thickness τ is defined as the line integral of the extinction (absorption plus scattering) coefficient between levels z_1 and z_2: $\tau = \int_{z_1}^{z_2} \beta_e(z) dz$.

SYNOPTIC CONDITIONS. We now examine changes in the large-scale dynamics and thermodynamics that occur when the Pacific high shifts toward the northeast prior to a CTD. Grids from the North American Mesoscale Forecast System (NAM) 12-km analysis at 0000 UTC are analyzed for all 23 CTD cases. First, we average all 0000 UTC NAM grids (97.7% availability) for June–September 2004–16 to produce a climatology. Then, we construct anomalies for each of the four days prior to a CTD and for the day of the initiation (days –4 to 0) using the 0000 UTC NAM grids (Fig. 2).

During the warm season, the Pacific high typically resides hundreds of kilometers west of California and generates low-level, northerly flow along the shore; winds at 850 hPa are approximately 5–8 m s^{-1}. The isobaric surface slopes downward toward the coast as the 850-hPa temperatures increase to the east. At 850 hPa, the cross-coast height gradient is strongest (~5 m per 100 km) and directed away from the coastline offshore Northern California; the height gradient decreases in magnitude with decreasing latitude. Average conditions at 500 hPa reveal a mostly zonal (westerly) flow with a weak trough offshore. Three to four days before a CTD initiates, the large-scale conditions are similar to climatology, although the position of the anticyclone is slightly south, as indicated by the anomalously low sea level pressure (SLP) values and 850-hPa heights to the north. The 850-hPa temperatures are slightly below (above) average north (south) of central California along the coastline, while at 500 hPa they are generally above average offshore California. Diffluence occurs at 850 hPa near 130°W, and the zonal flow at 500 hPa is greater than climatology.

Over the next several days, the Pacific high shifts northeastward and strengthens noticeably, and the

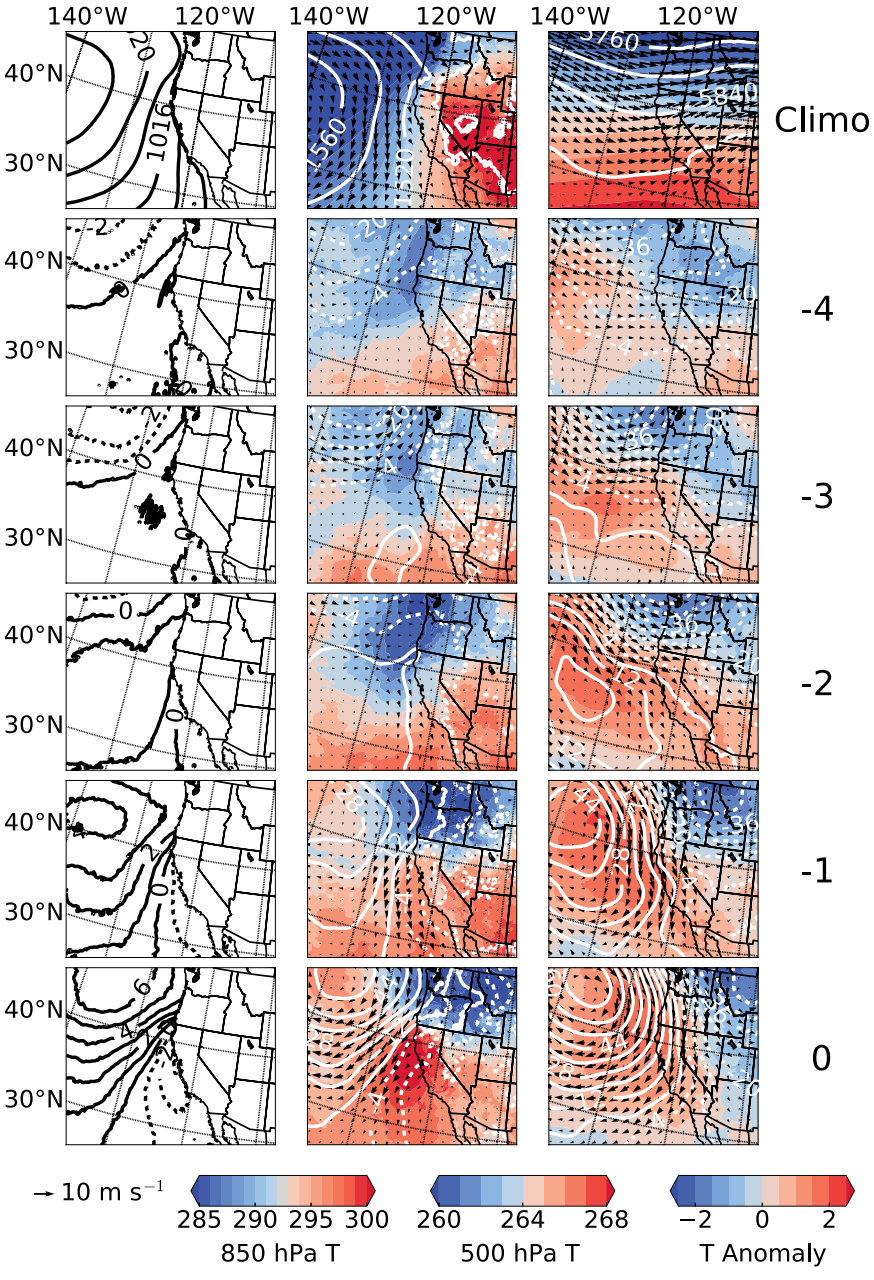

Fig. 2. Composites of 0000 UTC 12-km NAM grids: (left) SLP (black contours; units in hPa), and (middle) 850-hPa and (right) 500-hPa height (white contours; m), wind arrows, and temperature (color contoured with color bar; K). The topmost panel shows the climatological average, while the following five panels show the anomalies relative to climatology for days –4, –3, –2, –1, and 0. For the anomaly plots, solid (dashed) contours indicate positive (negative) values. A reference wind vector is also plotted.

thermal low pressure extends offshore Southern California, causing a tighter pressure gradient. An enhancement in the northerly flow is observed at 850 hPa; wind speeds exceed 10 m s^{-1} in many areas. Troughing along the Northern California coast is evident on day –1 as offshore continental flow begins and 850-hPa temperatures increase. On the day that a typical CTD initiates (day 0), SLP values are greater than 1028 hPa—more than 6 hPa (8 hPa) greater than climatology (on day –3)—around 42°N, 135°W. The SLP anomaly near central California on day 0 is about –3 hPa. A slight rotation in the 850-hPa height contours fosters additional offshore flow; the 850-hPa height gradient, which increases to ~15 m per 100 km offshore Northern California, reverses sign near Monterey Bay to become directed toward the coastline (not shown). This additional tightening of the height gradient leads to even stronger winds offshore that are blowing from the north-northeast. Adiabatic descent and warm air advection induces troughing along the California coast as far north as Cape Mendocino. At 500 hPa, a ridge–trough pattern develops offshore Northern California. A similar evolution in both the lower- and upper-level conditions are reported in the Mass and Bond (1996) analysis.

The synoptic-scale composites in Fig. 2 are representative of all 23 CTD events as determined by a qualitative analysis of the NAM grids for each event. This is confirmed through a statistical evaluation whereby we consider each variable for day –4 through day 0. The percentage area over the oceanic domain illustrated in Fig. 2 that is statistically significant at the 95% confidence level is listed for all variables in Table 3. The number of statistically significant grid boxes is smallest for all variables on day –4, reiterating that conditions deviate little from climatology. Unsurprisingly, the percentage of statistically significant areas increases toward day 0 and is greatest on day –2, –1, or 0 for all variables. Greater than 60% of the oceanic region is statistically significant for seven of eight variables on at least one day, which justifies the claim that the large-scale meteorological conditions are perturbed substantially preceding the inception of a CTD.

AIR–SEA INTERACTIONS. Changes in SSTs over the 14-day period surrounding the initiation of a CTD are shown in Fig. 3. The actual SST and surface wind fields are plotted for day 0, while the anomalies relative to day 0 are shown for all other days. About a week prior to the CTD, low-level, offshore winds weaken to decrease ocean upwelling and increase SSTs. A large swath of SST anomalies greater than +0.5°C are observed. Maximum warming is seen off the Northern California coast around days –3 and –2.

The offshore winds then begin to increase in accordance with the strengthening SLP gradient; the strongest northerly winds are observed on the day that a typical CTD commences (day 0). Because upwelling-favorable winds are enhanced, SSTs decrease from day –1 to +3 offshore Northern California. Anomalies approach –0.5°C and lead to a net decrease in SST of more than 1.0°C over 4 days across ~50,000 km^2. A lag of 2–3 days is observed between the changing sea surface wind stress and subsequent SST anomaly. A lag in cooling SSTs following the reestablishment of northerly winds after a relaxation event is described in Goela et al. (2016).

As expected, the southerly flow associated with the CTD extends northward over time and disrupts the upwelling regime beginning at day 0. After day +2, the SST anomalies increase significantly, and by day +7, anomalies reach a maximum in magnitude and spatial extent: warm anomalies exceed +1.0°C in many areas and are statistically significant at the 90% confidence level over ~350,000 km^2 (hatched region in Fig. 3).

The three-stage cycle in near-surface winds and SSTs here is similar to that identified by Fewings et al. (2016) and Flynn et al. (2017) during wind relaxation events. The authors note that, on average, the

Table 3. Percentage of area where anomaly for each listed meteorological variable is statistically significant at the 95% confidence level. We use a *t* test and let the null hypothesis be the climatological value. For each variable, the boldface value represents the largest percentage.

	Day –4	Day –3	Day –2	Day –1	Day 0
SLP	0.1	2.0	35.7	**80.1**	77.1
850-hPa *u*	1.2	24.0	18.1	25.3	**63.8**
500-hPa *u*	1.1	22.3	30.2	25.1	**64.9**
850-hPa *v*	4.3	14.7	24.2	48.6	**62.1**
500-hPa *v*	4.1	15.8	**60.0**	58.8	51.6
850-hPa *Z*	0.0	4.7	28.3	70.3	**71.9**
500-hPa *Z*	19.5	22.2	55.6	**67.0**	65.2
850-hPa *T*	0.1	11.3	38.0	33.4	**46.0**
500-hPa *T*	3.5	33.8	65.2	**72.1**	43.0

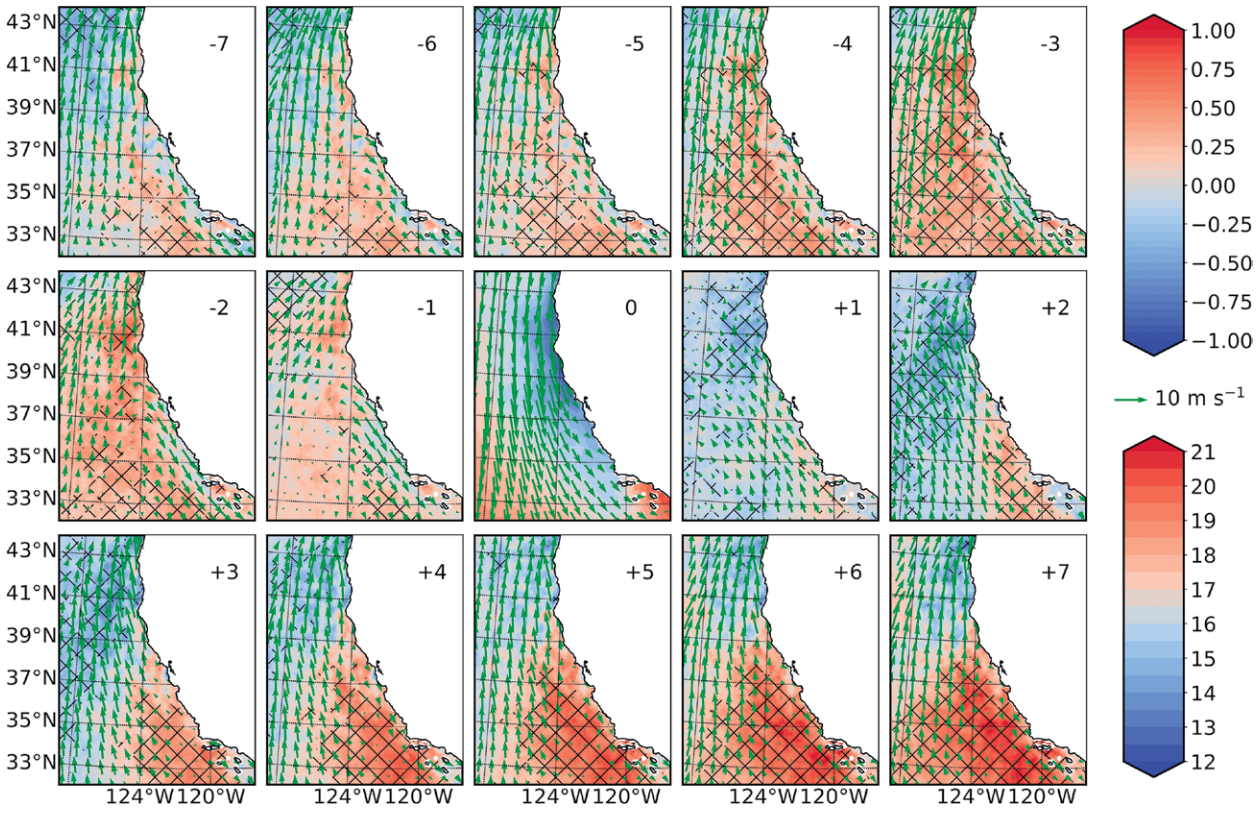

Fig. 3. Composites of SST and surface wind from satellite observations. Numbers in the top-right corner of each panel refer to the day relative to the initiation of a CTD (day 0). The actual SST and wind fields are plotted for day 0, while the anomalies relative to day 0 are shown for all other days. The hatching indicates regions of SST anomalies that are statistically significant at the 90% confidence level. The top and bottom color bars refer to SST anomalies and the actual SST field (units in °C), respectively. A reference wind vector is also shown. Wind data are not available for the 2004 and 2006 CTD cases.

wind stress does not reverse sign. While this suggests that downwelling-favorable conditions are not common in their study, such conditions are commonplace during a CTD. We suspect that the underlying cause for the difference in downwelling frequency is due to the strength of the Pacific high: anomalies in SLP one day before a typical wind relaxation event began are only ~2 hPa (Fewings et al. 2016).

SURFACE TRANSITION. Hourly observations at each of the five buoy stations for all valid events are analyzed for the 24-h period across each CTD (Fig. 4). Each event is normalized such that the time when the buoy first reports northward flow is at the center of the plot (0 h). Similar to the Bond et al. (1996) study, most of the CTDs begin propagating overnight or in the early morning hours (typically between 0000 and 1200 UTC on day 0). The initiation and surging of a CTD overnight are tied to diabatic processes. Daytime heating enhances the cross-coast pressure gradient that favors northerly wind near the coast (Beardsley et al. 1987), which inhibits the formation of a CTD during the day. At night, the temperature in the cloud-free region north of the CTD changes little compared to the cloudy region of the incipient CTD, where the temperature decreases rapidly and the MBL deepens due to cloud-top longwave radiative cooling. The increased density contrast produces a northward acceleration along the coast and is responsible for the characteristic surge overnight and into the morning (Rahn and Parish 2008).

A stark contrast in the wind speed and direction is evident at nearly all buoys. Twelve hours prior to a CTD, the surface wind speed is typically around 5 m s^{-1} except for buoy 46011, where the wind speed is slightly weaker. At all buoy sites, wind speeds decrease as the CTD approaches. A weak, positive cross-coast component—indicating onshore flow—is evident at all buoy sites except buoy 46014, where values are slightly negative. Leading up to the reversal, strong negative alongshore (southward) winds are apparent. At the time of the reversal, a minimum in the wind speed and magnitude of the alongshore component is observed. After the leading edge of

the CTD passes, wind speeds begin to increase and reach a maximum by about +3 to +6 h. The signal in the cross-coast component is not consistent between the buoys; only buoy 46014 shows a clear signal with offshore (onshore) flow ahead of (behind) the CTD. As expected, however, the alongshore wind turns positive and is approximately 3–4 m s^{-1} on average.

The number of strong CTD events (alongshore component of at least 5 m s^{-1} for at least one report) at buoys 46011, 46028, 46042, 46013, and 46014 is 3 (20% of all valid events), 8 (44%), 9 (41%), 12 (57%), and 11 (61%), respectively. These frequencies suggest that the northern sites are more likely to experience stronger alongshore winds within the CTD. We hypothesize that CTDs usually initiate near the southern locations, so they tend to be less developed and weaker, while the northern locations are more representative of well-developed and stronger CTDs that surge northward along the coast. Overall, we find that the probability of observing a strong CTD event is much greater than previously reported (~10%–15%; Bond et al. 1996).

Pressure and air temperature perturbations are also evaluated during the transition from northerly to southerly flow. For each event, the pressure and temperature at the time of the CTD (0 h) are considered the mean values. Pressure generally reaches a minimum about 3–6 h before the reversal. Air temperature is relatively high prior to the CTD, although buoy 46042 shows slightly lower temperature during these times. During the 12 h following a typical CTD, the pressure at each buoy rises. This pressure increase after the passage of a CTD has been previously noted (e.g., Bond et al. 1996; Parish et al. 2008). Temperature observations also show an increase after the passage, with a minimum typically occurring at or just after the wind shift. The temperature for buoy 46042 does not show this coherent trend. Monterey Bay influences the buoy 46042 measurements because the change in coastline geometry and the presence of complex terrain may significantly alter the CTD propagation and the land–sea breeze circulation (Banta et al. 1993). Satellite imagery regularly shows that the flow at the head of the CTD

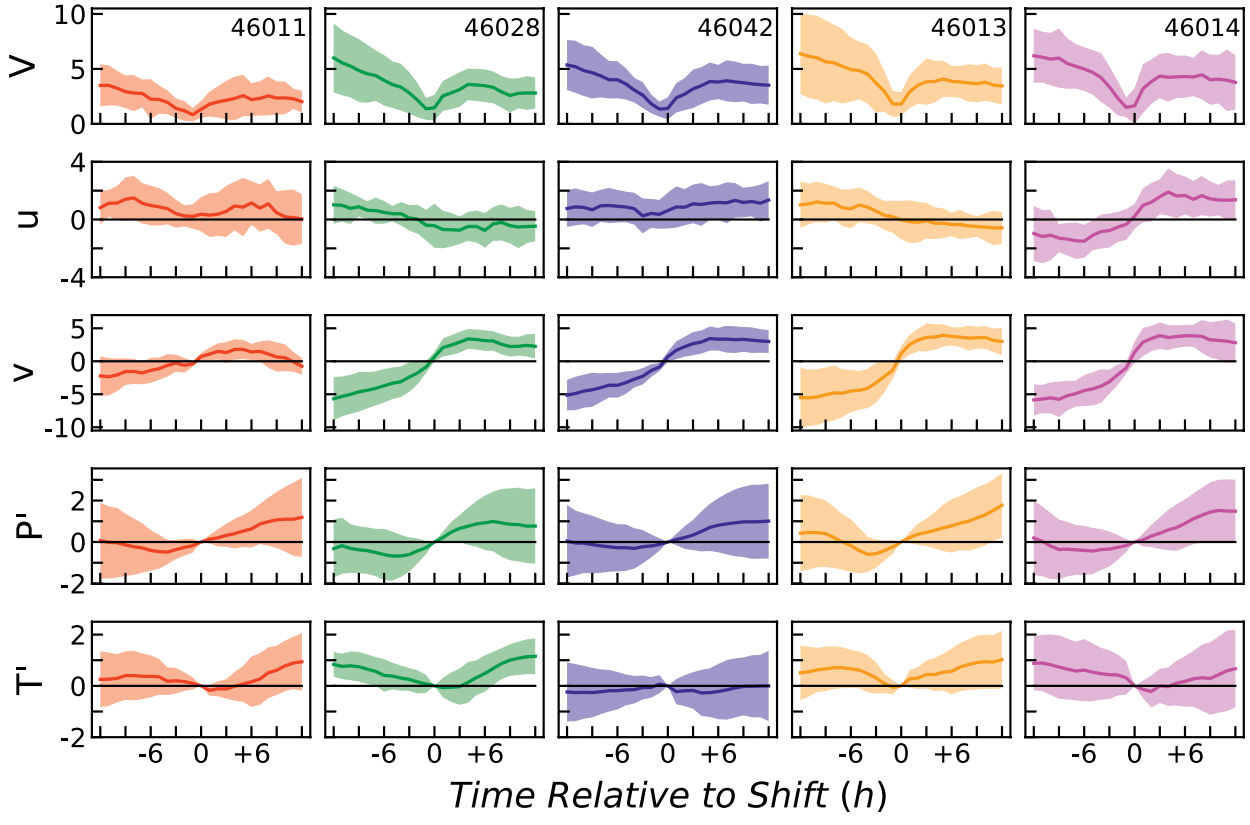

Fig. 4. Time series of horizontal wind speed (*V*; units in m s^{-1}), cross-coast wind speed (*u*; m s^{-1}), alongshore wind speed (*v*; m s^{-1}), pressure perturbation (*P′*; hPa), and temperature perturbation (*T′*; °C). The coordinate system is rotated 30° to align approximately with the coastline such that the cross-coast (*u*) and alongshore (*v*) wind components are defined as 330° and 60°, respectively, in accordance with Ralph et al. (2000) and Parish et al. (2008). Buoy observations are arranged from south on the left to north on the right.

follows the coastline into Monterey Bay and the low cloud is often seen entering the Salinas Valley south of the bay before the CTD continues northward along the coast. Overall, the air temperature analysis here differs from the CTD climatology by Bond et al. (1996), where buoy data indicated little trend.

The discrepancies in the i) number of strong CTD events and ii) alongshore temperature gradient between the Bond et al. (1996) study and present study may be attributed to the case selection process. Because more than half of our cases propagate north of Cape Mendocino, we hypothesize that the average synoptic-scale, alongshore pressure gradient on the order of 1 hPa per 100 km driving the CTDs considered here, which is stronger than that found in the Bond et al. (1996) results, increases the in-cloud wind speed. A strong alongshore pressure gradient is due to strong offshore flow that enhances subsidence and inversion strength, which act to decrease cloud-top entrainment and maintain the cloud layer (e.g., Wood and Bretherton 2006; Iacobellis and Cayan 2013). Because a well-developed cloud deck is present in all 23 cases here, the average CTD may be characterized by an alongshore thermal gradient that is stronger than that found for the typical CTD in Bond et al. (1996).

CONTRASTING CLOUD REGIMES. We now compare CTD cloud characteristics with those for

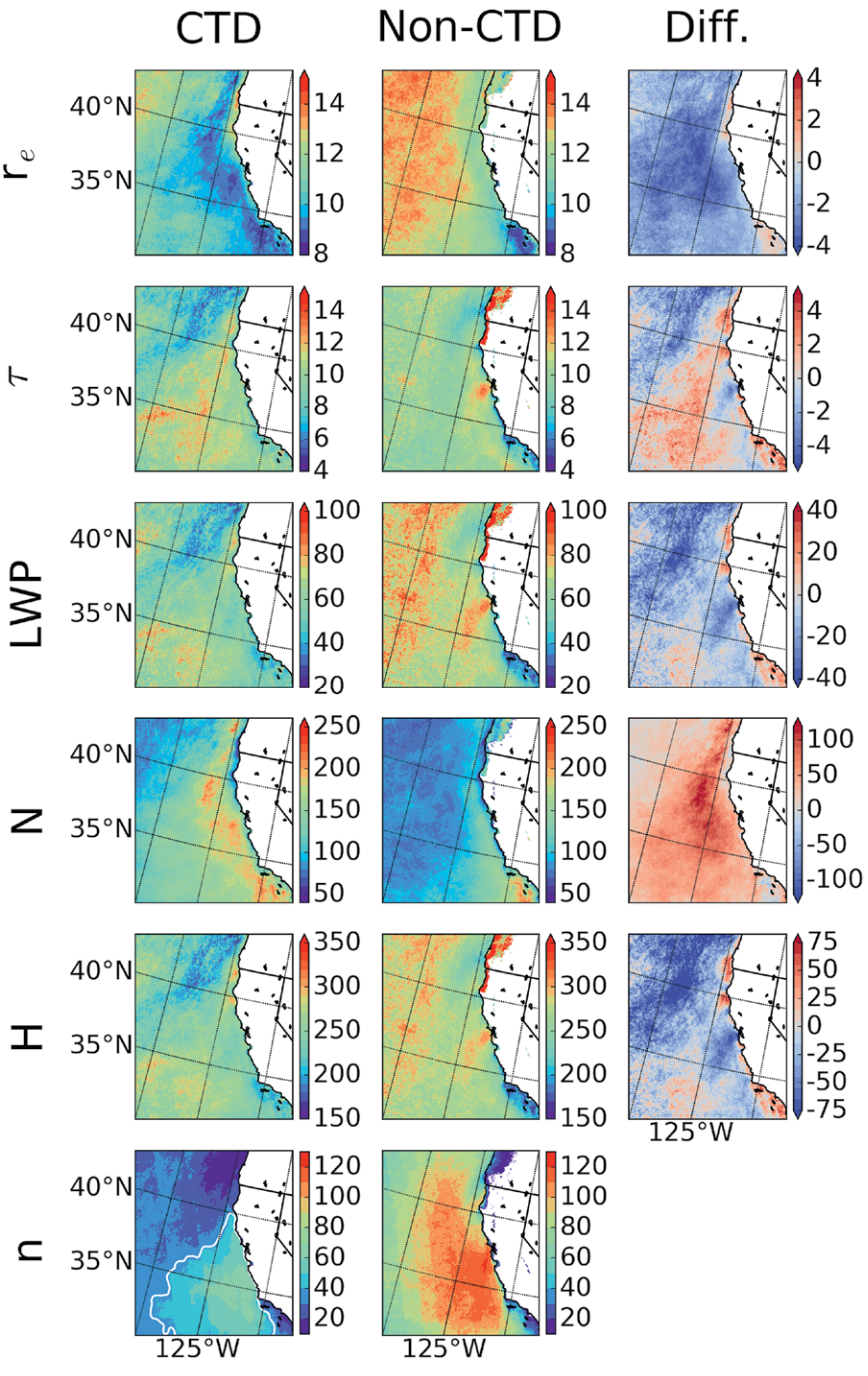

FIG. 5. Mean microphysical parameters contoured for (left) CTD and (middle) non-CTD events along with the (right) difference (CTD minus non-CTD). Positive (negative) differences indicate larger (smaller) values during CTD events. Properties shown include r_e (units in μm), τ (unitless), LWP (g m^{-2}), N (cm^{-3}), and H (m). The number of samples n at each location is also color contoured. We contour the n = 37 line (~54% of n_{max} for CTD cases) in bold white to provide a visual for the approximate mean cloud deck extent during a mature CTD. We consider all days listed in Tables 1 and 2. The 1-km MODIS retrievals are interpolated to a regular 1/20° × 1/20° (~5 km × 5 km) grid using inverse distance weighting.

marine clouds during non-CTD events (Fig. 5; also see sidebar "Estimates of cloud depth and cloud droplet number concentration"). Several key distinctions between cloud regimes are evident. Minimum values of r_e (~8–10 μm) are observed in a large swath west of the U.S. coastline during the CTD cases. During the non-CTD events, however, r_e is much greater offshore; in many regions, differences in r_e between the CTD and non-CTD events exceed 3 μm (~20%–40%). Accompanying the relatively small values of r_e are large values of N. Values approach 250 cm^{-3} during CTDs and are nearly a factor of 2 greater compared to non-CTDs in many areas. For nearly all points in the domain, N is, on average, greater during CTD than non-CTD cases.

Capes and points along the coastline appear to modify the cloud field during CTD events. An increase in r_e (~10–14 μm) is observed close to the shoreline with maximum values adjacent to the coast found near the California–Oregon border. However, the sample size near the coast north of Cape Mendocino is small because not all of the CTDs propagate to this region. Pockets of diminished N (~50–150 cm^{-3}) are also present in these areas. Overall, differences in r_e, τ, LWP, and H between the CTD and non-CTD cases display a dipole pattern on the northern and southern locations of capes and points. Positive (negative) differences in these variables are observed on the south (north) side of terrain features. We hypothesize that because points and capes act as barriers to and induce hydraulic features within the marine flow, they modify local-scale dynamics, cloud layer depth, and microphysical processes (e.g., condensational growth, entrainment, and drizzle production). For instance, flow convergence upstream of a barrier may promote cloud layer deepening (e.g., due to an oblique hydraulic jump; see Haack et al. 2001), while flow divergence downstream of a barrier may promote cloud layer thinning (e.g., due to an expansion fan; see Juliano et al. 2017). Future high-resolution modeling studies should be performed to test this hypothesis.

Frequency distributions plotted in Fig. 6 confirm the disparity in r_e and N between CTD and non-CTD events. The r_e distribution for CTD cases is relatively narrow and has a peak around 8 μm. The peak is also around 8 μm for non-CTD cases, but the distribution is much broader with a tail that falls off much slower. Both distributions are right-tailed. Due to the inverse relationship between r_e and N, the peak of the N distribution occurs at ~40 cm^{-3} for non-CTD cases and ~110 cm^{-3} for CTD cases. Similar to the r_e distribution, both N distributions are positively skewed, and the one for non-CTDs falls off quickly, while there are many more instances of large N (>150 cm^{-3}) during CTDs. Differences in τ, LWP, and H between CTD and non-CTD cloud regimes are not as large, although there are generally more instances of optically thicker (larger τ) and shallower (smaller H and LWP) clouds associated with CTDs.

Normalized joint probability density functions are used to highlight CTD and non-CTD differences in the cloud physics (Fig. 7). As has been the common theme, the likelihood of observing clouds with more numerous (N) and smaller (r_e) droplets is higher in CTD cloud

ESTIMATES OF CLOUD DEPTH AND CLOUD DROPLET NUMBER CONCENTRATION

Cloud depth H and cloud droplet number concentration N may be derived from observations of τ and LWP after assuming an adiabatic cloud model (Bennartz 2007). These relationships [Eqs. (8) and (9) in Bennartz 2007] are expressed as follows:

$$H = \left(2\frac{\mathrm{LWP}}{c_w C_F} \right)^{1/2}, \qquad \text{(SB1)}$$

$$N = \frac{2^{-5/2}}{k}\tau^3 \left(\frac{\mathrm{LWP}}{C_F} \right)^{-5/2} \left(\frac{3}{5}\pi Q \right)^{-3} \left(\frac{3}{4\pi\rho_l} \right)^{-2} c_w^{1/2}, \qquad \text{(SB2)}$$

where c_w is the condensation rate (often called the liquid water content lapse rate) in kg m^{-4}, C_F is the fraction of warm clouds, k is the ratio between the volume mean radius and the effective radius to the third power, and Q is the scattering efficiency. We calculate the condensation rate, which is a function of cloud base temperature and pressure, for each 1 × 1 km^2 grid box containing cloud. Cloud-top temperature and pressure are derived from the MODIS retrievals. Due to the shallow nature of MBL clouds, we assume that i) cloud base temperature is equal to cloud-top temperature and ii) cloud-base pressure is 25 hPa greater than cloud-top pressure. Furthermore, Lu and Seinfeld (2006) report that k varies between 0.5 and 0.9, and so we choose a value of 0.8 following Bennartz and Rausch (2017). Finally, Q = 2 is consistent with Mie theory because cloud droplets are much larger than the incident solar radiation (size parameter $\gamma \geq 30$; Bennartz 2007). For a review of limitations and sensitivities relating to estimates of N from *Aqua* MODIS retrievals, the reader is referred to the recent discussion by Bennartz and Rausch (2017).

decks than in typical marine stratus and stratocumulus decks. These CTD clouds with high N may be relatively shallow or deep (and therefore characterized by low or high LWP) and optically thin or thick (small or large τ). Once the number of cloud droplets approaches ~250 cm^{-3}, however, H, LWP, and τ converge toward ~175 m, ~40 g m^{-2}, and ~7.5, respectively. Moreover, when relatively shallow MBL clouds are present, there is a higher probability of observing optically thicker (thinner) clouds in CTD (non-CTD) cases because N is typically larger during CTD events.

CLOUD DROPLET NUMBER CONCENTRATION ENHANCEMENT. Through MODIS retrievals, we show that marine stratiform clouds accompanying CTDs exhibit, on average, a clear increase in cloud droplet number concentration N when compared to N in marine clouds that develop during the typical northerly wind regime (non-CTDs). We hypothesize that the fundamental cause for this substantial increase in the number of cloud droplets is threefold:

i) The observed wind stress–SST cycle promotes mixing of sea salt aerosol into the MBL.
ii) Before entering the cloud deck, CTD air parcels spend a considerable amount of time passing through major shipping lanes.
iii) Offshore continental flow, which is a requirement for the initiation of a CTD, transports continental aerosol into the marine layer.

We depict these three processes using schematics (Fig. 8) and elaborate with the following discussion.

Sea salt aerosol concentration is known to depend strongly on surface wind speed (e.g., Gong 2003; Pierce and Adams 2006; Feng et al. 2017). We suspect that as northerly winds offshore of California ramp up immediately preceding a CTD, elevated ocean stress leads to an increase in sea spray and sea salt particles mixing into the MBL (Fig. 8, top panel). Given that their residence time, which is inversely proportional to particle size, may range from 30 min to over 2 days (e.g., Lovett 1978; Gong et al. 1997; Jaeglé et al. 2011), sea salt aerosol may efficiently mix into CTD clouds due to strong convergence (e.g., Thompson et al. 2005; Parish et al. 2008; Rahn and Parish 2008) and act as nucleation sites for cloud droplets. These aerosols may play an important role in broadening the MBL cloud droplet distribution, as only a few

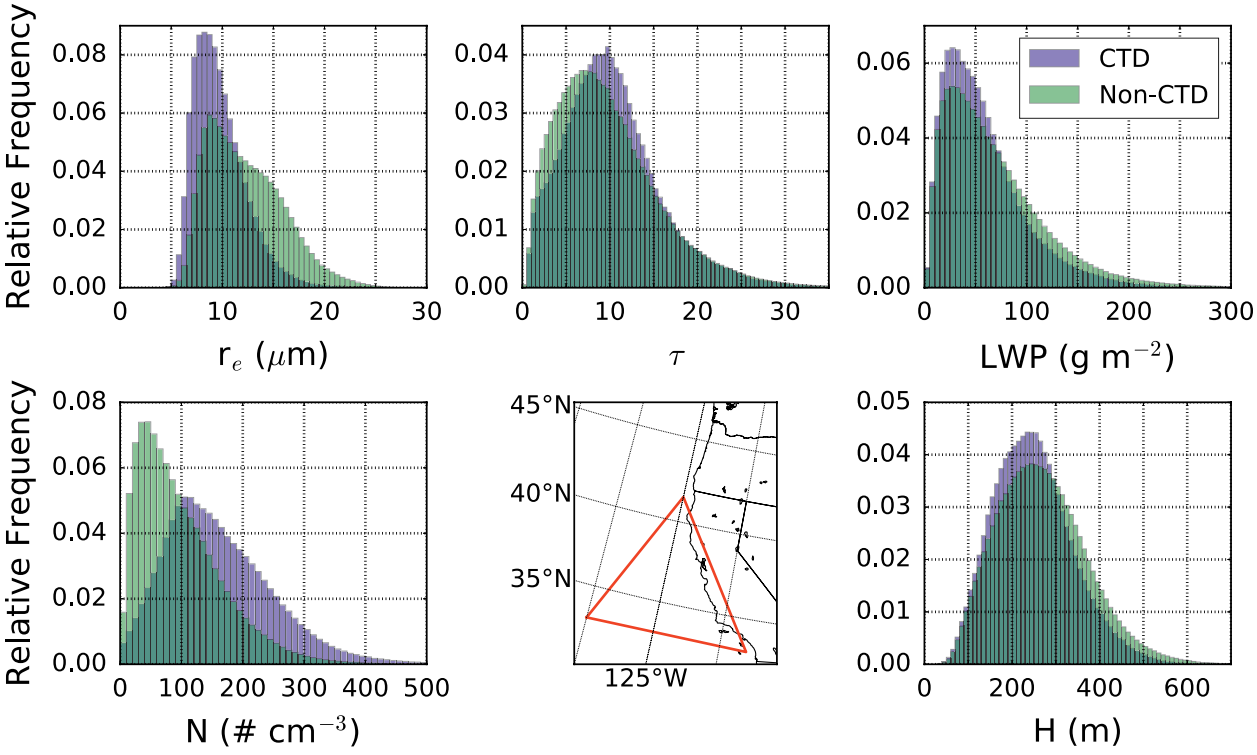

Fig. 6. Histograms comparing relative frequency of occurrence for the microphysical variables presented in Fig. 5 between CTD and non-CTD cases. The binned data are constrained to the outlined polygon (bottom middle panel) because this region approximates the spatial extent of a typical CTD. The histograms are mostly insensitive to the location, size, and shape of the polygon.

giant cloud condensation nuclei (CCN) are needed to initiate drizzle even in clouds with relatively high CCN concentrations (Feingold et al. 1999; Jensen and Nugent 2017).

Dimethyl sulfide and, predominantly, ship exhaust contribute to the marine sulfate budget (e.g., Langley et al. 2010). While we have no reason to believe that shipping emissions differ between CTD and non-CTD cases, it is possible that CTD air parcels are more likely to interact with exhaust plumes prior to entering the cloud deck because they tend to follow shipping lanes (Fig. 8, bottom panel; e.g., see Fig. 9 in Coggon et al. 2012). In situ cloud samples from a CTD case considered here (27–28 July 2011; Coggon et al. 2012) reveal enhanced N (>200 cm^{-3}) and high concentrations of vanadium, which is a primary by-product of ship exhaust (e.g., Agrawal et al. 2008).

The California coastal zone is home to a variety of aerosol that may be advected offshore (Fig. 8, bottom panel). Aircraft measurements suggest that biogenic volatile organic compounds, which are emitted by the heavily forested region in Northern California and southern Oregon, may make up a large portion of above-cloud aerosol (Coggon et al. 2014). Observations by Hegg et al. (2009) and Hegg et al. (2010) show that anywhere from about 50% to 67% of CCN offshore of California originate from anthropogenic sources. Due to agricultural activity, central and Southern California are large producers of mineral dust (e.g., Clausnitzer and Singer 1997; Ngo et al. 2010; Hand et al. 2017), which may act as CCN when coated by sea salt or sulfate (e.g., Levin et al. 2005; Gibson et al. 2007). Additionally, as drought risk in California continues to increase (e.g., Cook et al. 2015; Williams et al. 2015), the influence of

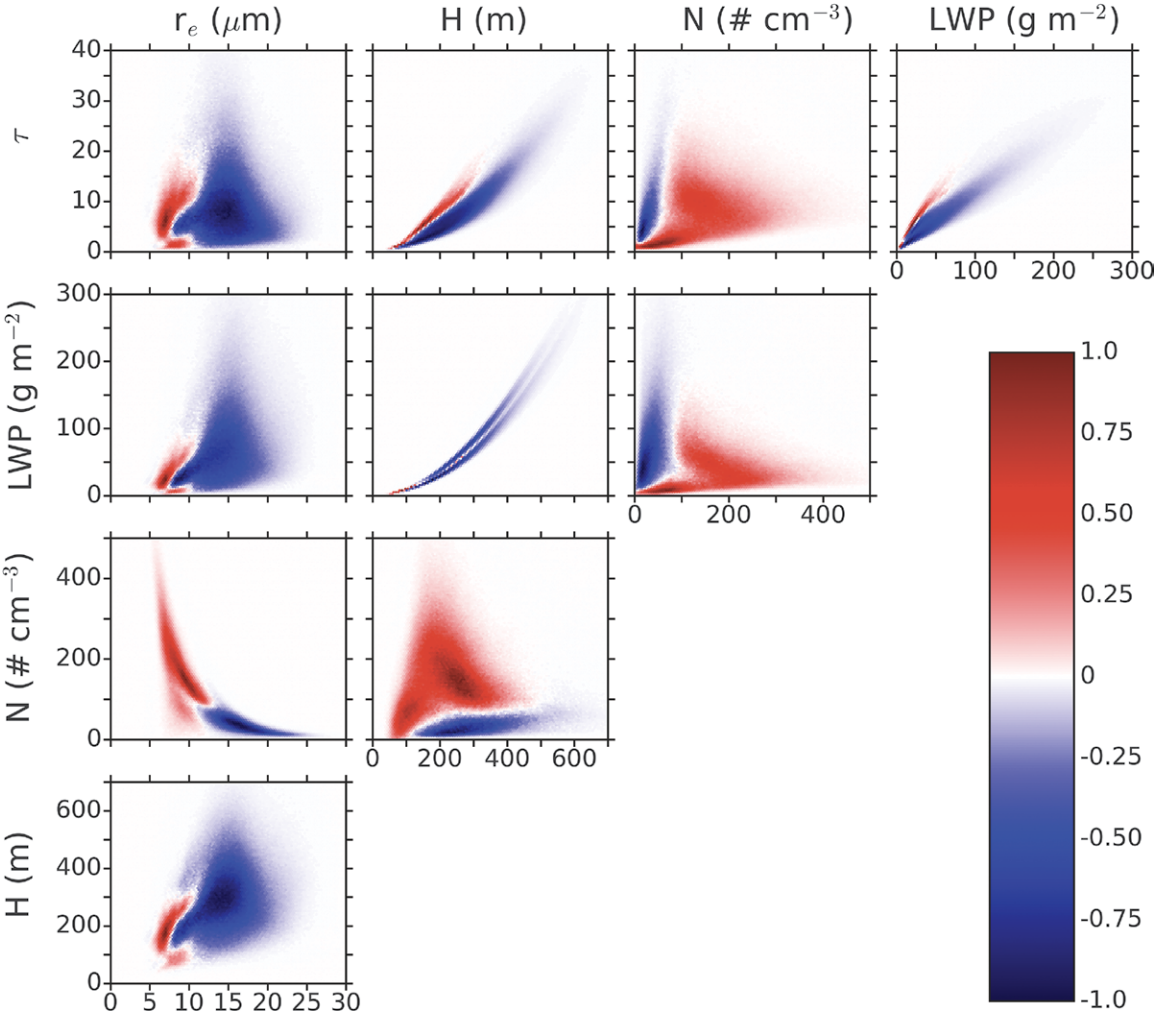

Fig. 7. Normalized joint probability density functions for the various microphysical parameter spaces presented in Fig. 5. Red (blue) color contours represent regions in the specified parameter space where the probability of finding a sample is higher (lower) in CTD cases compared to non-CTD cases.

biomass burning aerosol on marine stratocumulus warrants additional research, as their radiative impact is profound (e.g., Brioude et al. 2009; Lu et al. 2018).

To address the extent to which CTDs may be influenced by these various anthropogenic sources, we generate composite backward trajectories for CTD and non-CTD cases (Fig. 9). The trajectories suggest that for both nearshore and offshore locations, CTD cloud decks are more susceptible to shipping and continental emission sources than non-CTD cloud decks. This finding is expected considering the large-scale meteorological changes that occur during CTD events to promote alongshore flow near the surface and offshore flow above the MBL. Any lateral or vertical cloud entrainment will therefore allow polluted air parcels to mix into the cloud layer and alter cloud properties.

Using backward trajectories and data from the Cloud–Aerosol Lidar with Orthogonal Polarization (CALIOP) for an individual case study

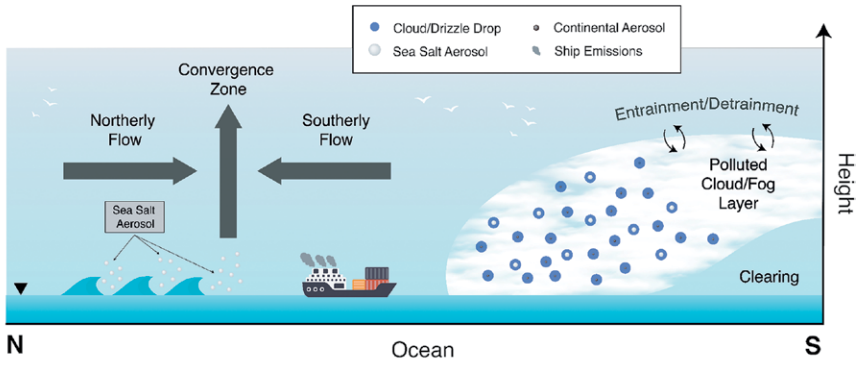

Fig. 8. Schematic diagrams of CTD cloud processes in (top) vertical cross-sectional framework and (bottom) horizontal plan view framework. See text for discussion. Abbreviations are also shown for reference: Washington (WA), Oregon (OR), Nevada (NV), California (CA), Cape Mendocino (CM), Sacramento (SAC), San Francisco (SF), Monterey Bay (MB), Point Conception (PC), Los Angeles (LA), and San Diego (SD).

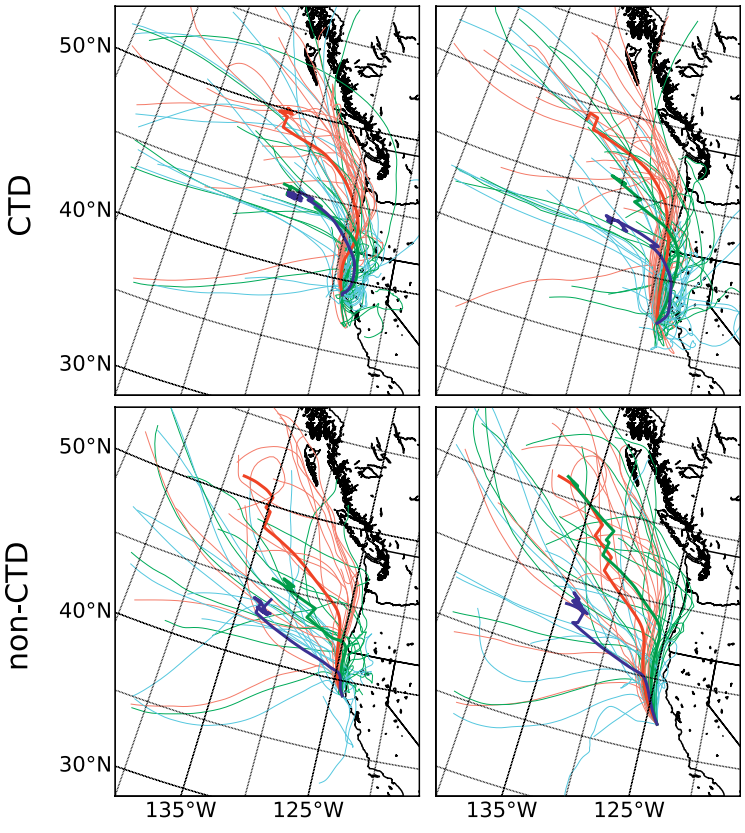

Fig. 9. Composites of 72-h backward trajectories using 12-km NAM grids in the Hybrid Single-Particle Lagrangian Integrated Trajectory (HYSPLIT) model beginning at (left) 39°N, 124°W and (right) 37°N, 124°W for (top) CTD and (bottom) non-CTD cases. Red, green, and blue lines represent trajectories ending at 100, 500, and 1,500 m AGL, respectively. Light, thin lines are shown for all trajectories while dark, bold lines represent mean trajectories. For each event, trajectories begin at 0000 UTC one day after the first dates listed in Tables 1 and 2. Grids are not available for two and seven of the CTD and non-CTD cases, respectively.

(9 July 2008), we show the numerous aerosol sources that may affect CTD clouds (Fig. 10). According to the CALIOP data, marine (sea salt), polluted continental (sulfate, nitrate, organic carbon, and ammonium mixed with urban pollution), and biomass burning (soot and organic carbon) aerosol are all present in the vicinity of the cloud deck. Any regions labeled as "other" indicate that a mixture of at least two aerosol types is present; therefore, it is likely that over the ocean the omnipresent sea salt aerosol often coexists with ship exhaust and/or various types of continental aerosol. A daytime *Aqua* MODIS "true color" image confirms that the smoke plumes on this day are from multiple wildfires in California. Examination of additional images from CALIOP (not shown) suggests that aerosol from both marine and continental sources are present in the vicinity of the cloud deck during each of the 23 CTD events.

IMPLICATIONS FOR RADIATION BALANCE. Due to their anomalous properties, we expect marine clouds that accompany CTDs to affect the radiation balance differently than non-CTD marine clouds over the northeast Pacific Ocean. We compare fractional cloud albedo α_c for CTD and non-CTD cases through frequency distributions in Fig. 11. Unsurprisingly, due to the large difference between cloud and ocean albedos, both regimes exhibit a strong SW cloud radiative forcing (CRF_{SW}), or cooling effect (see sidebar "Shortwave cloud radiative forcing"). The mean values for α_c and CRF_{SW} are 53.5% and –148.9 W m^{-2} (51.6% and –142.5 W m^{-2}) for CTD (non-CTD) cases, respectively. The histograms show more cases of high α_c during CTDs; thus, on average, these clouds reflect more incoming shortwave radiation than their non-CTD counterparts. We attribute the ~2% albedo increase to the enhancement of droplet concentration and subsequent increase (decrease) in CTD optical depth (droplet size). In situ observations that compare a CTD event to two non-CTD events also show that CTD clouds have relatively high τ (14.9 vs 10.2 and 7.6), low r_e (8.9 μm vs 9.2 and 14.3 μm), and high α_c (63% vs 59% and 48%) values (Crosbie et al. 2016).

Assuming the longwave radiative forcing of the two cloud regimes is equal, results suggest that, on average, marine stratiform CTD clouds off the coast of California exert a top-of-the-atmosphere (TOA) SW radiative forcing of ~4% (~6.4 W m^{-2}) more than non-CTD clouds. We emphasize that CTD clouds are present less frequently than marine stratiform clouds that form under typical summertime conditions. Moreover, the cloud clearing that usually occurs prior to the arrival of a CTD also has implications for the

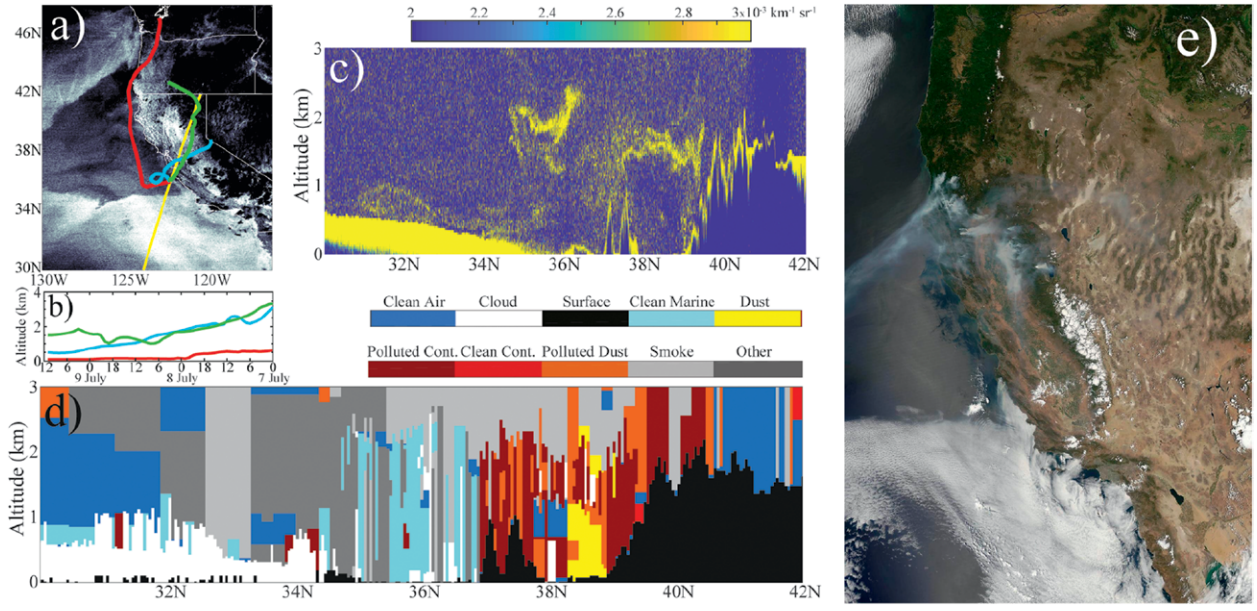

Fig. 10. (a) The 1000 UTC 9 Jul 2008 GOES fog product using band 4 (10.7 μm) minus band 2 (3.9 μm), the 1013–1020 UTC *Cloud–Aerosol Lidar and Infrared Pathfinder Satellite Observations (CALIPSO)* overpass used for cross sections (yellow), and the 1000 UTC HYSPLIT backward trajectories at 100 m (red), 500 m (blue), and 1,500 m (green) for 36°N, 122.2°W. (b) Time–height plot of backward trajectories. (c) Cross section of 532-nm total attenuated backscatter (units in km^{-1} sr^{-1}). (d) Cross section of CALIOP vertical feature mask including aerosol feature subtypes and corresponding legend. (e) A 9 Jul 2008 daytime *Aqua* MODIS "true color" image.

radiation balance (i.e., relatively less reflected SW).

SUMMARY AND FUTURE WORK. Through the use of various observational tools, we provide a new perspective on the disruption of the familiar summertime marine stratiform deck offshore California. It is shown that MBL clouds which form under the typical northerly flow conditions (non-CTD) differ notably compared to those which form under the anomalous southerly flow conditions (CTD). The main characteristics differentiating these two cloud regimes are summarized in Table 4.

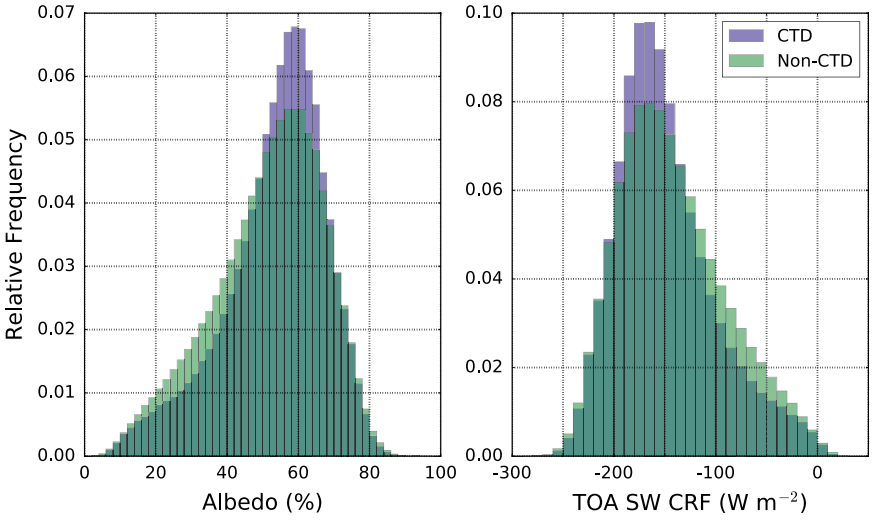

Fig. 11. Histograms comparing relative frequency of occurrence for albedo (units in %) and TOA CRF_{sw} (W m^{-2}) between CTD and non-CTD cases. As in the previous histograms, the binned data here are constrained to the outlined polygon in Fig. 6.

In the context of CTDs, the interplay between aerosol and clouds is clearly a complex issue that must be further addressed. Research pertaining to individual aspects of the aforementioned aerosol sources is ongoing; untangling the symbiotic relationships between these aerosols is not trivial. Although we do not present here any direct observational evidence regarding the composition of particles acting as CCN, we confidently hypothesize that the enhancement of N is primarily due to an abundance of hygroscopic aerosols originating from both continental and oceanic sources. The degree to which these aerosols are successful in their competition to become activation sites and contribute toward subsequent cloud processes (e.g., diffusional growth, autoconversion, collision–coalescence) is a question that remains unanswered.

It is evident that the large-scale meteorological conditions strongly influence the occurrence of a CTD. A first step is to understand better the frequency of CTD events, in addition to the specific synoptic-scale pattern(s) that affect particular aspects of the CTD life cycle. For instance, how do global cycles, such as the El Niño–Southern Oscillation and Pacific decadal oscillation, affect the shift in the Pacific high during the boreal summer? Moreover, it may be informative to identify any clear differences in the large-scale conditions (e.g., strength and location of offshore continental flow) that support weak versus strong CTDs because they may differ in terms of propagation speed, cloud properties, and lifetime. Unsupervised machine learning techniques, such as self-organizing maps, are sometimes applied to extensive datasets in atmospheric science (e.g., Liu and Weisberg 2011). Employing

SHORTWAVE CLOUD RADIATIVE FORCING

If Rayleigh scattering is neglected for a nonabsorbing, plane-parallel atmosphere (homogeneous scattering layer) with an asymmetry factor, $g \approx 0.85$ (Lacis and Hansen 1974), then scattering of solar radiation by clouds may be described by

$$\alpha_c \approx \frac{0.13\tau}{1+0.13\tau}, \tag{SB3}$$

where α_c is fractional cloud albedo. Upon calculating α_c, the TOA CRF_{SW} may be estimated as

$$\mathrm{CRF_{SW}} = \left(\frac{S_o}{4}\right)(\alpha_o - \alpha_c), \tag{SB4}$$

where S_o is the solar constant, and α_o is ocean albedo. We assume S_o = 1,370 W m^{-2} and α_o = 0.10 (10%). Because the albedo of marine stratiform clouds is much greater than 10%, we expect CRF_{SW} to be large and negative.

Table 4. Summary of main differences between non-CTD and CTD features.

Feature	Non-CTD	CTD
Pacific high	Relatively weak and located far offshore	Relatively strong and located closer to Pacific Northwest
Alongshore pressure gradient force	Directed southward everywhere	Directed southward (northward) north (south) of central California
Near-surface winds	Northerly flow everywhere	Northerly (southerly) flow outside (inside) CTD
SSTs	Persistent upwelling regime leading to relatively cool SSTs	Three-stage cycle of relatively weak–strong–weak winds leading to relatively warm–cool–warm SSTs
Microphysical properties	Relatively low N and large r_e	Relatively high N and small r_e
Aerosol influences	Primarily marine sources (sea salt and ship exhaust)	Both marine sources and continental sources (forested, wildfire, agricultural, and urbanized regions)
Radiative impact	Relatively low albedo and less reflected SW	Relatively high albedo and more reflected SW

these techniques to CTDs may yield important details regarding the physical connection between predictive variables and ultimately lead to higher-accuracy forecasts across a range of temporal and spatial scales.

Observations (e.g., Minnis et al. 1992; Painemal et al. 2013; Burleyson et al. 2013) and idealized large-eddy simulations (e.g., Caldwell et al. 2005; Caldwell and Bretherton 2009; Kazil et al. 2016) show that marine stratocumulus display notable diurnal variability in, for example, H, LWP, entrainment rate, and drizzle production. Capturing the midmorning and early evening transition periods in a numerical model is critical for the cloud field evolution (e.g., horizontal and vertical distribution of liquid water and the subsequent impact on drizzle processes) and continues to challenge both the regional forecasting and climate prediction communities. One way to potentially ameliorate forecasting errors associated with coastal marine cloud–radiation interactions is to collocate aircraft (top-down) and surface-based (bottom-up) measurements (e.g., microwave radiometer).

In general, the diurnal variability in MBL cloud properties is difficult to analyze using satellite retrievals because most techniques utilize solar wavelengths and are limited to daylight hours. The Visible Infrared Imaging Radiometer Suite (VIIRS) Day/Night Band, however, is a unique sensor that can be utilized to detect marine stratocumulus clouds at night (e.g., Miller et al. 2013; Seaman and Miller 2015). Additionally, future efforts should exploit the new channels aboard the Geostationary Operational Environmental Satellite system (GOES-R) that can retrieve aerosol and cloud properties while coordinating field campaigns to validate these remote sensing capabilities.

Additional observations, specifically in situ, are needed to confirm the remote sensing *Aqua* MODIS observations shown here that suggest CTD clouds and fogs are composed of more cloud droplets than typical marine stratiform clouds. In situ measurements should include instruments that are capable of characterizing the chemical composition of aerosol. Aerosol size distribution measurements focusing on sea salt are necessary to determine the potential impact on drizzle processes. Because CTD clouds typically extend to near the ocean surface, accurately measuring aerosol below cloud base is very challenging. Therefore, utilizing platforms such as unmanned aerial vehicles seems appealing, especially as momentous advances continue in their deployment for atmospheric, and specifically air quality, observations (e.g., Villa et al. 2016).

CTDs should also be probed through high-resolution numerical modeling methods. Sensitivity studies using a framework such as the Weather Research and Forecasting Model (Skamarock et al. 2008) may allow for a deeper understanding of cloud processes relevant to CTDs. Once the model accurately represents the horizontal and vertical distributions of cloud liquid water using a sensible combination of physical parameterizations, then various components may be perturbed to yield additional insight into small-scale atmospheric processes. Furthermore, recent modeling efforts targeting more detailed microphysical treatments [e.g., prognostic supersaturation using the predicted particle properties (P3) scheme (Morrison and Milbrandt 2015) and prognostic aerosol using the Thompson aerosol-aware scheme (Thompson and Eidhammer 2014)] may prove useful when examining CTDs. Ultimately, because the real-life problem is a complicated interaction between multiple systems (ocean, land, and atmosphere), running a fully coupled model is the most ideal technique to represent CTD systems.

ACKNOWLEDGMENTS. Author TWJ is supported in part by the National Science Foundation through Grant AGS-1439594 and by the State of Wyoming. Author ZJL thanks support through the Department of Energy from Grant DE-SC0016354. Author DAR is supported in part by the National Science Foundation through Grant AGS-1439515. We gratefully thank Justine Sulia for creating the schematics. We also thank the three anonymous reviewers and the editor whose comments have improved the manuscript. Steven Platnick provided helpful information regarding the *Aqua* MODIS data processing. The MODIS "true color" image (Fig. 10e) is courtesy of the Space Science and Engineering Center at the University of Wisconsin–Madison.

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