

**Methane dissociation on the steps and terraces of Pt(211)  
resolved by quantum state and impact site**

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Methane dissociation on the step and terrace sites of a Pt(211) single crystal was studied by reflection absorption infrared spectroscopy (RAIRS) at a surface temperature of 120 K. The C-H stretch RAIRS signal of the chemisorbed methyl product species was used to distinguish between absorption on step and terrace sites allowing methyl uptake to be monitored as a function of incident kinetic energy for both sites. Our results indicate a direct dissociation mechanism on both sites with higher reactivity on steps than on terraces consistent with a difference in activation barrier height of at least 30 kJ/mol. State-specific preparation of the incident CH<sub>4</sub> with one quantum of antisymmetric ( $\nu_3$ ) stretch vibration further increases the CH<sub>4</sub> reactivity enabling comparison between translational and vibrational activation on both steps and terraces. The reaction is modeled with a first principles quantum theory that accurately describes dissociative chemisorption at different sites on the surface.

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## I. INTRODUCTION

The dissociation of small polyatomic molecules on transition metal surfaces often represents the rate limiting step in the heterogeneously catalyzed processes used to produce chemicals on an industrial scale<sup>1</sup>. Developing a predictive understanding of these dissociation reactions is not only of fundamental interest, but also of practical importance.

In steam reforming, used commercially to produce hydrogen, the dissociation of methane on the transition metal catalyst is the rate determining step<sup>2</sup>. Experiments have shown this reaction is non-statistical<sup>3</sup> with mode-specificity<sup>4-6</sup>, bond-selectivity<sup>5,7,8</sup> and stereo-specificity<sup>9,10</sup> all being observed, making obtaining an accurate theoretical description challenging. Only very recently has quantitative agreement (within chemical accuracy, 4.2 kJ/mol) between experiment and theory been achieved by Kroes and co-workers using specific reaction parameter (SRP) density functional theory<sup>11,12</sup>. By using the initial reaction probability ( $S_0$ ) for the dissociation of CHD<sub>3</sub> on Ni(111) from one experiment to optimize a single parameter in their calculations, they were able to reproduce  $S_0$  obtained experimentally for CHD<sub>3</sub> dissociation both under laser-off conditions and prepared with a quantum of C-H stretch vibration on Ni(111), Pt(111) and Pt(211).

In previous quantum state-resolved studies of methane dissociation on transition metal surfaces<sup>13-17</sup>, only the total reactivity averaged over all impact sites on the surface has been determined. By combining quantum state-specific reactant preparation in a molecular beam with reflection absorption infrared spectroscopy (RAIRS) product detection, we demonstrate here that we can additionally resolve whether methane dissociates on the steps (red) or terraces (black) of a Pt(211) crystal (shown schematically in Fig. 1), allowing both quantum state-specific and site-resolved reactivity to be determined simultaneously. This additional level of resolution provides a further stringent test of theory, requiring that the reactivity determined from calculations can also be resolved by impact site on the surface. Here, we use

15 dimensional quantum reactive scattering methods based on the Reaction Path Hamiltonian (RPH) approach<sup>18–24</sup> to calculate site-resolved values of  $S_0$  which we compare to the values obtained experimentally. While traditional quantum methodologies for gas-surface scattering have advanced significantly in recent years<sup>25</sup>, our RPH approach remains the only method that can accurately include and describe all of the vibrational modes of methane.

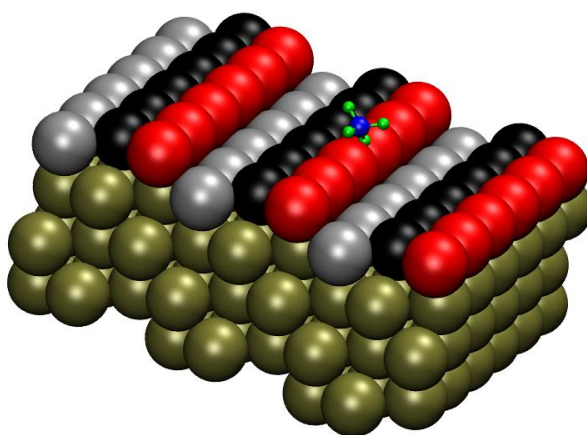


FIG 1. The Pt(211) surface has three atom wide (111) terraces and one atom high (100) steps. We label the row of atoms on the step edge as step (red), the middle row terrace (black) and the final row corner (gray).

## II. EXPERIMENTAL METHODS

The experimental apparatus has been described in detail previously<sup>26</sup>, and only the most relevant details will be presented here. In brief, the experiments were performed in a surface science machine consisting of a triply differentially pumped molecular beam source chamber attached to an ultrahigh vacuum (UHV) chamber with a base pressure of  $5 \times 10^{-11}$  mbar. The single crystal Pt(211) surface, cut to within  $0.1^\circ$  of the (211) plane, was mounted in the UHV chamber in a tantalum support between two tungsten wires. A K-type thermocouple was spot welded to the tantalum mount to monitor the temperature, which was stabilized using a PID controller. From temperature programmed desorption measurements, we estimate that there was a maximum of 10 K difference between the temperature of the

tantalum mount and the Pt(211) surface. The surface could be heated to over 1100 K by passing a current along the tungsten wires, and cooled to 100 K through thermal contact with a liquid nitrogen reservoir. Between deposition measurements, the Pt(211) crystal was cleaned by exposure to  $5 \times 10^{-8}$  mbar of  $O_2$  at 700 K before annealing at 1100 K. At the end of each day cleaning by  $Ar^+$  sputtering and annealing at 1100 K was used. The surface cleanliness was confirmed using Auger electron spectroscopy.

The molecular beam was formed by expanding 2 bar of a 3%  $CH_4$  seeded in He gas mix through a stainless steel nozzle with a 50  $\mu m$  diameter hole before passing through a 2 mm diameter skimmer. To change the incident kinetic energy ( $E_k$ ) of the molecules the nozzle was resistively heated to temperatures of up to 800 K. The nozzle temperature was monitored by a K-type thermocouple spot welded to the end of the nozzle and stabilized to within 1 K using a PID controller. The velocity distribution of the molecular beam was measured by the time of flight method using a fast chopper wheel in conjunction with an on-axis quadrupole mass spectrometer (QMS). The experimental distribution of  $CH_4$  arrival times was fit by a flux weighted Maxwell Boltzmann distribution<sup>27</sup>. During the molecular beam deposition  $CH_4$  was incident perpendicular to the plane of the Pt(211) surface. The  $CH_4$  flux incident on the Pt(211) surface was monitored using an off-axis QMS calibrated against a cold-cathode ion gauge.

For the laser-on measurements, a fraction of the incident methane was prepared quantum state selectively in the antisymmetric C-H stretch fundamental vibration  $v_3 = 1, J = 2$  via the R(1) transition using a continuous wave optical parametric oscillator. The laser frequency was stabilized to the ( $v_3 = 1, J = 2 \leftarrow v = 0, J = 1$ ) transition frequency of 3038.490  $cm^{-1}$  by locking to a Lamb-dip<sup>10</sup> created in a static absorption cell containing 30  $\mu bar$  methane. We excite the molecular beam using rapid adiabatic passage<sup>28</sup> (RAP) by focusing the laser with a 254 mm focal length cylindrical lens to create curved wavefronts in the

region where the molecular beam traverses the laser beam. CH<sub>4</sub> molecules passing through the focused laser beam experience a frequency sweep due to the Doppler Effect which, providing certain conditions are met, can completely transfer the population of the initial rovibrational state to the final rovibrational state. A room temperature pyroelectric detector was used to determine the flux of state prepared molecules created by RAP.

We recorded the methyl uptake resulting from exposure of the Pt(211) surface to an incident molecular beam of CH<sub>4</sub> at a surface temperature ( $T_s$ ) of 120 K, where the nascent methane dissociation products CH<sub>3</sub>(ads) and H(ads) are stable. CH<sub>3</sub> adsorbed on Pt(211) was detected by the RAIRS technique using a Bruker Vertex V70 Fourier Transform Infrared Spectrometer with external liquid nitrogen cooled InSb detector. Each RAIRS spectrum was recorded with 4 cm<sup>-1</sup> resolution averaging over either 256 scans or 512 scans, corresponding to an acquisition time of approximately 30 seconds or one minute. To convert the spectra to uptake curves, we fit the spectra to the sum of two Gaussians to obtain the area of the absorption peaks for CH<sub>3</sub>(ads) on the step and terrace. The full width at half maximum (FWHM) of both peaks were on the order of 5 cm<sup>-1</sup> to 8 cm<sup>-1</sup>. As the RAIRS spectra were recorded with a resolution of 4 cm<sup>-1</sup>, the FWHM of each peak showed significant scatter around the general trend, with the problem then transferred to the calculated area. To reduce the noise in the area of the spectra the FWHM of the spectra were fit to a logarithm function, and the FWHM from this fit used with the peak height obtained from the Gaussian fit to obtain the absorption area which is presented as the CH<sub>3</sub> RAIRS absorption signal.

The flux of methane molecules impinging on the surface was calculated using<sup>29</sup>

$$\text{Flux} = \frac{\Delta P S}{A k_B T_g} \quad (1)$$

where  $\Delta P$  is the pressure rise when the molecular beam enters the UHV chamber obtained from the calibrated off-axis QMS,  $S$  is the experimentally determined pumping speed of the UHV chamber for CH<sub>4</sub> (520 l/s),  $A$  is the area of the molecular beam spot on the surface

determined by Auger electron spectroscopy,  $k_B$  is Boltzmann's constant and  $T_g$  the temperature of the static gas in the UHV chamber (298 K). The dose is the product of incident flux multiplied by the deposition time. For the laser-on measurements, we take into account the excited fraction ( $f_{\text{exc}}$ ) of the incident  $\text{CH}_4(v_3 = 1, J = 2)$  prepared by RAP.

### III. THEORETICAL METHODS

Electronic structure calculations are performed using the Density Functional Theory (DFT)-based Vienna *ab-initio* simulation package (VASP), developed at the Institut für Materialphysik of the Universität Wien<sup>30–34</sup>. A four-layer  $3 \times 2$  supercell with periodic boundary conditions is used to represent the metal as a series of infinite slabs, with a large vacuum space of 16 Å between the slabs. The interactions between the ionic cores and the electrons are described by fully nonlocal optimized projected augmented-wave (PAW) potentials<sup>34,35</sup>. The exchange-correlation effects are treated within the generalized gradient approximation using the Perdew-Burke-Ernzerhof (PBE) functional<sup>36,37</sup>. We also use the SRP functional developed by the Kroes group in Leiden<sup>11,12</sup>, which has been shown to accurately describe the dissociative chemisorption of methane on both Ni and Pt surfaces<sup>11,12,38</sup>. The SRP functional consists of a combination of PBE<sup>36,37</sup> and RPBE<sup>39</sup> for exchange, and vdW-DFT<sup>40</sup> for correlation.

Our Hamiltonian for dissociation over a rigid metal lattice is:

$$H = K + V = -\frac{\hbar^2}{2} \sum_{i=1}^{15} \frac{\partial^2}{\partial x_i^2} + V(x_1, x_2, \dots, x_{15}), \quad (2)$$

where the  $x_i$  are the 15 mass-weighted Cartesian coordinates of the  $\text{CH}_4$  nuclei. To construct our potential energy surface (PES) we first locate all relevant transition states (TSs) and the corresponding minimum energy paths (MEPs) for reaction. We do not relax the metal lattice in the presence of the  $\text{CH}_4$  because we model a high-energy collision where the lattice does

not have time to relax during the collision. We have characterized four transition states for dissociative chemisorption: dissociation along the step edge (path N), dissociation across the edge (paths M and L), and dissociation over a terrace atom. Details are presented in the next Section. In Fig. 2 we plot the total energy,  $V_0(s)$ , along the L and N MEPs, as a function of the distance along the path,  $s$ , where  $(ds)^2 = \sum_{i=1}^{15} (dx_i)^2$ ;  $s = 0$  at the TS. For comparison we also plot the MEP for dissociation on Pt(111), which is similar to that on the Pt(211) terrace, except for a change in barrier height.

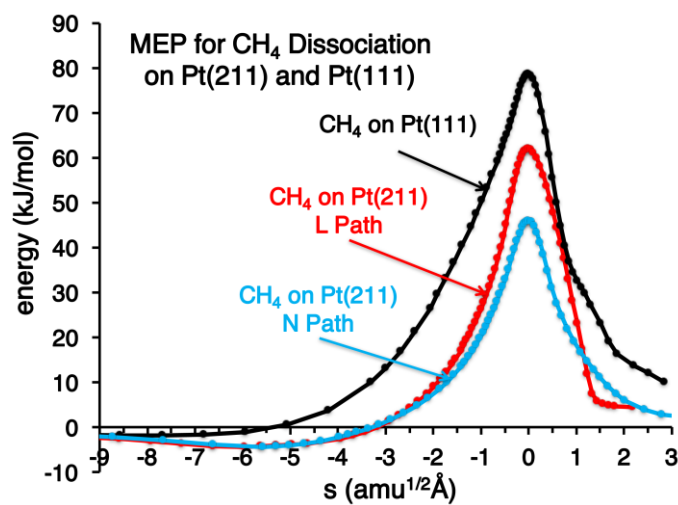


FIG. 2. Reaction paths for methane dissociation on Pt(211) and Pt(111).

At several dozen images along each reaction path we compute and diagonalize the force-projected Hessian to get the fourteen normal vibrational coordinates,  $Q_k$ , and corresponding frequencies,  $\omega_k(s)$ , describing harmonic motion orthogonal to the reaction path. Ignoring higher order (anharmonic) terms, our PES can be written in the reaction path coordinate  $s$  and  $\{Q_k\}$ :

$$V = V_0(s) + \sum_{k=1}^{14} \frac{1}{2} \omega_k^2(s) Q_k^2. \quad (3)$$

We write our molecular wave function, in reaction path coordinates, as:

$$Y(t) = \hat{a} \sum_n c_n(s;t) F_n(\{Q_k\};s), \quad (4)$$

where the  $F_n$  are products of harmonic oscillator eigenfunctions that depend parametrically on  $s$ , and the vector  $\mathbf{n}$  labeling the vibrational states corresponds to a set of quantum numbers  $n_k$ . The equations of motion for the wave packets,  $\chi_n(s; t)$ , are<sup>20</sup>:

$$i\hbar \frac{\partial \chi_n(s; t)}{\partial t} = \left[ \frac{1}{2} p_s^2 + V_0(s) + \sum_{k=1}^{14} \hbar \omega_k(s) \left( n_k + \frac{1}{2} \right) \right] \chi_n(s; t) + \sum_{n'} F_{nn'} \chi_{n'}(s; t). \quad (5)$$

The wave packets evolve on vibrationally adiabatic potential energy surfaces for each vibrational state  $\mathbf{n}$ , and the operators  $F_{nn'}$  couple states of the same symmetry. The  $F_{nn'}$  are proportional to the vibrationally nonadiabatic couplings computed from the normal mode eigenvectors. Because of the parametric dependence of the  $\Phi_n$  on  $s$ , the  $F_{nn'}$  also contain momentum operators,  $p_s$ , and curve crossing (transitions between vibrationally adiabatic states) becomes more probable at higher velocities, as well as for larger values of the coupling. The sums over  $\mathbf{n}$  in Eqs. 4 and 5 include the vibrationally adiabatic ground state and all states with either one or two vibrational quanta in the nine normal modes of the incident molecule. A detailed derivation can be found in a recent publication<sup>22</sup>. For a given initial vibrational state,  $\mathbf{n}_0$ , standard techniques are used to evolve the wave packets<sup>19,20</sup> and energy-analyze the reactive flux<sup>41,42</sup>. The result is the rigid-lattice reaction probability,  $P_0(E_k, \mathbf{n}_0)$ , for collision at a surface impact site corresponding to one of the 4 MEPs. The rotational orientation of the molecule has been allowed to evolve adiabatically.

To compute  $S_0$  we first average  $P_0$  over all surface impact sites assuming sudden behavior for  $X$  and  $Y$ , the location of the molecular center of mass over the surface unit cell. Motion along  $X$  and  $Y$  should be slow on collisional time scales, given the normal incident conditions. *Ab-initio* molecular dynamics (AIMD) studies confirm that there is little translational steering<sup>21</sup>. We divide the Pt(211) surface into three regions: the step region, the terrace region and the corner region, as shown schematically in Fig. 1. We assume that the corner region contributes nothing to  $S_0$ , given the large barrier. We average  $P_0$  for the terrace



over all impact sites within that region, using the following to approximate  $P_0$  at impact sites  $(X, Y)$  close to the minimum barrier site at  $X = 0, Y = 0$ :

$$P_0(E_i, \mathbf{n}_0; X, Y) \approx P_0(E_i - DV, \mathbf{n}_0; 0, 0). \quad (6)$$

$\Delta V(X, Y)$  is the increase in barrier height relative to the minimum barrier site. We can approximate this using the normal modes, or compute it directly using DFT for several molecular configurations near the TS. We take a similar approach for the step, though that region is broken up into sub-regions corresponding to the M, N and L reaction paths.

AIMD also suggests that the rotational behavior might be closer to sudden at high incident  $E_k$ <sup>21,43</sup>. To compute the sticking in the rotationally sudden limit, we use a similar approach to Eq. (6), averaging over orientations of the molecule at the TS, as described in a recent work<sup>24</sup>. As in that work, we then use a switching function to combine these two results so that the final  $S_0(E_k)$  is 90% rotationally adiabatic at  $E_k = 0.2$  eV and 90% rotationally sudden at  $E_k = 0.9$  eV.

We treat lattice motion using a sudden model described elsewhere<sup>44,45</sup>. Given the metal atom mass and the short collision times, we assume the lattice atom is stationary during the collision. On Ni and Pt surfaces, as the lattice atoms vibrate the height of the barrier to methane dissociation changes. On the (100) and (111) surfaces of Ni and Pt, and for the Pt(211) terrace atoms, this is mostly confined to the motion, normal to the surface, of the metal atom over which the methane dissociates<sup>44,45</sup>. If  $Q$  is the displacement of this lattice atom in and out of the plane of the surface, and  $Q > 0$  for motion away from the bulk, the barrier height changes by an amount  $-\beta Q$ , where  $\beta = 91.9$  kJ/mol/Å on Pt(111)<sup>46</sup>. We average the sticking probability over all values of  $Q$ , using a Debye model for the probability of having a value of  $Q$  for a given  $T_s$ . To do this we need to know the root mean square value for  $Q$ ,  $Q_{\text{rms}}$ , at some  $T_s$ . We have typically gotten this from surface scattering experiments, but such data is not available for Pt(211), nor is such data site specific. We

have thus used AIMD with our Pt(211) slab to compute  $Q_{\text{rms}}$  at 300 K. We find that for the terrace atom,  $Q_{\text{rms}} = 0.109 \text{ \AA}$  in the direction normal to the (111) terrace, very similar to the  $Q_{\text{rms}}$  value of  $0.104 \text{ \AA}$  on Pt(111). On the step edge, the behavior is more complex. We find  $\beta = 96.3$  and  $119.7 \text{ kJ/mol/ \AA}$  for the L and N paths, respectively, but  $Q$  has components both normal and lateral to the surface, away from the edge. The motion of other atoms can also modify the barrier, but their contributions are small. Our AIMD calculations give  $Q_{\text{rms}} = 0.099 \text{ \AA}$  for the motion of step atoms in the direction perpendicular to the surface at 300 K. For motion lateral to the surface,  $Q_{\text{rms}}$  is not significantly different from this. Lattice vibration also changes the location of the TS along  $Z$ . We find that as the lattice is displaced normal to the (211) surface by  $Q_Z$ , the location of the TS changes by  $\alpha Q_Z$ , where  $\alpha = 0.8379$  and  $0.8993$  for the L and N paths, respectively. We use the modified Surface Mass Model<sup>47</sup> to include this effect.

#### IV. RESULTS AND DISCUSSION

It is well known that  $\text{CH}_4$  dissociates on Pt(111) by a direct reaction and that a single C-H bond is broken, leading to chemisorbed products  $\text{CH}_3(\text{ads})$  and  $\text{H}(\text{ads})$ <sup>8,26,48–50</sup>.  $\text{CH}_3(\text{ads})$  on Pt(111) adsorbs on top of a Pt atom and is stable at  $T_{\text{S}}$  below  $200 \text{ K}$ <sup>49</sup> and can be detected by RAIRS *via* absorption peaks at  $2883 \text{ cm}^{-1}$  and  $2755 \text{ cm}^{-1}$  (Fig. 3A top) assigned to the symmetric  $\text{CH}_3(\text{ads})$  stretch<sup>26,50,51</sup> and the first overtone of the antisymmetric bend, respectively. For dissociation of  $\text{CH}_4$  on Pt(211), we observe two peaks in the  $\text{CH}_3(\text{ads})$  stretch region at  $2886 \text{ cm}^{-1}$  and  $2903 \text{ cm}^{-1}$  (Fig. 3A bottom) which we assign by comparison with the spectrum for Pt(111) to the symmetric stretch of  $\text{CH}_3(\text{ads})$  on terrace and step sites respectively. The assignment of the  $2903 \text{ cm}^{-1}$  peak to  $\text{CH}_3(\text{ads})$  on step sites is supported by the observation of much higher reactivity for this site (see Fig. 3) consistent with the calculated dissociation barrier heights for step, terrace and corner sites (Table I.)

We attribute the additional peak at  $2752\text{ cm}^{-1}$  to the bend overtone of the  $\text{CH}_3(\text{ads})$  on both sites. The detection of two peaks in the C-H stretch region of the RAIRS spectrum rather than three suggests that the corner sites are not occupied by  $\text{CH}_3(\text{ads})$  under the conditions of our experiments.

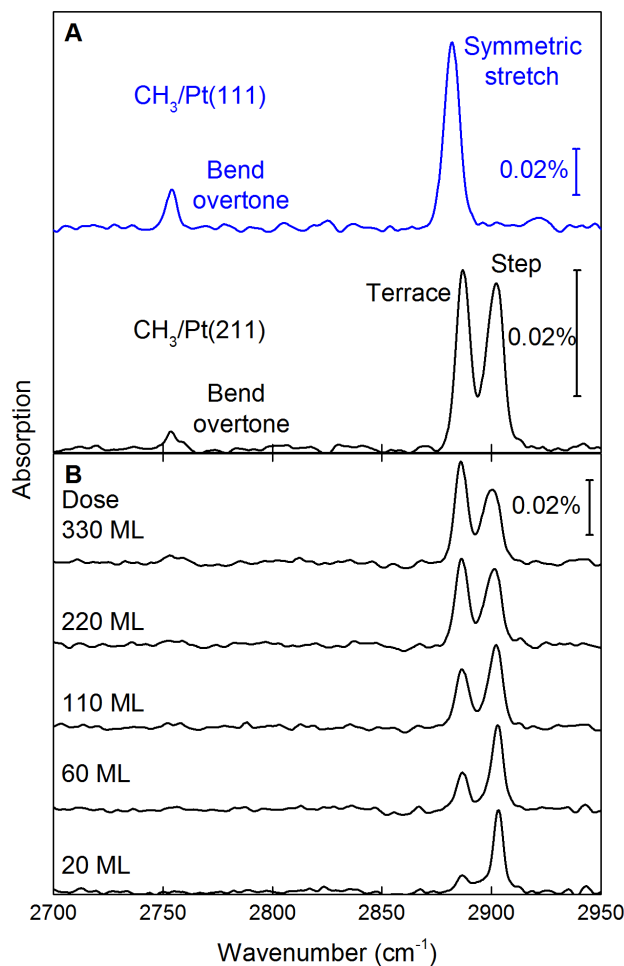


FIG. 3. A. Comparison of the RAIRS spectra of  $\text{CH}_3(\text{ads})$  on Pt(111) (top, blue) and Pt(211) (bottom, black) at  $T_S = 120\text{ K}$ . B. Variation of the RAIRS spectra for  $\text{CH}_3(\text{ads})$  with incident methane dose at  $T_S = 120\text{ K}$  and  $E_k = 65\text{ kJ/mol}$ .

These observations are consistent with our DFT calculations.  $\text{CH}_3$  prefers to bind on the top sites of Pt(111) and Pt(100)<sup>52,53</sup> and we find similar behavior on Pt(211). The adsorption energies for the three Pt(211) top sites are listed in Table I. Adsorption at the step edge is the most stable. We find that adsorption of  $\text{CH}_3$  on the top sites of Pt(111) is very similar to that on the Pt(211) terrace atoms. The most stable H adsorption is on the bridge site along the step edge, with an adsorption energy of  $275\text{ kJ/mol}$ ,  $13\text{ kJ/mol}$  larger than the

second most stable state. In Table II we summarize the properties of the four most important transition states for the dissociation of CH<sub>4</sub> on Pt(211). Consistent with the strong product binding on the step edge, the lowest barrier is along path N, where the activation energy is only 42 kJ/mol. For all transition states, the carbon atom is approximately above the top site with the reactive C-H bond angled towards the metal. The N transition state geometry, with the carbon over an edge atom and the breaking C-H bond oriented parallel to the edge, is consistent with that reported in other studies on Pt(211)<sup>12</sup>. For the M and L transition states, the carbon is also over an edge atom, but the dissociating bond is oriented perpendicular to the step edge, with the reactive H towards the (111) terrace for the M path and towards the (100) step for the L path. This geometry allows the carbon to get closer to the Pt at the TS, and the bond is stretched a bit more. The calculated activation energy for dissociation over a Pt(211) terrace atom is 84 kJ/mol, larger than that on Pt(111) (68 kJ/mol) or any of the edge sites. The terrace TS structure is very similar to that on Pt(111), taking into account that the (111) terrace of Pt(211) is not perpendicular to the surface normal. For the Pt(211) corner atoms, the CH<sub>3</sub>(ads) binding energy is only 123 kJ/mol, and we estimate that the barrier for dissociation is 183 kJ/mol. Thus, dissociation is not likely to occur at the corner sites at the molecular energies considered in this study.

Table I. The adsorption energy for CH<sub>3</sub> at step, terrace and corner top sites on Pt(211), and the top site on Pt(111).  $E_{ads}^{PBE}$  and  $E_{ads}^{SRP}$  are adsorption energies with lattice atoms unrelaxed, using PBE and SRP functionals, respectively.  $E_{ads}^{SRP,relax}$  is the adsorption energy when the top two layers of the metal lattice are allowed to relax.

	$E_{ads}^{PBE}$ (kJ/mol)	$E_{ads}^{SRP}$ (kJ/mol)	$E_{ads}^{SRP,relax}$ (kJ/mol)
Pt(211) Step	-204	-210	-215
Pt(211) Terrace	-178	-194	-209
Pt(211) Corner	-112	-123	-159
Pt(111) Top	-189	-194	-206

Table II. Transition state data for methane dissociation on Pt(211) and Pt(111).  $Z_C^\ddagger$  is the distance from the C atom to the Pt atom directly below it.  $r^\ddagger$  and  $\theta^\ddagger$  are the length of the dissociating bond and its angle relative to the surface normal, respectively.  $E_b$  is the barrier height relative to the methane infinitely far from the surface, and  $E_a$ , the activation energy, is  $E_b$  with zero point energy (ZPE) corrections. All results are in kJ/mol and are for the PBE functional, except for  $E_a^{SRP}$ , the activation energy using the SRP functional.

	$Z_C^\ddagger$ (Å)	$r^\ddagger$ (Å)	$\theta^\ddagger$ (Å)	$E_b$	ZPE	$E_a$	$E_a^{SRP}$
Pt(211), N	2.242	1.480	133.6	46	-10	36	42
Pt(211), M	2.157	1.608	122.0	54	-7	47	47
Pt(211), L	2.175	1.654	124.3	62	-9	53	55
Pt(211), terr	2.254	1.525	129.4	101	-11	90	84
Pt(111)	2.241	1.521	131.7	79	-11	68	68

We find that our peak assignment for  $\text{CH}_3(\text{ads})$  is reversed compared to those previously reported for CO on the steps and terraces of Pt(211)<sup>54,55</sup> and on curved Pt crystals<sup>56</sup>. For CO(ads) on steps, the C=O stretch is lower in frequency than for adsorption on the terrace site, whereas for  $\text{CH}_3(\text{ads})$  the step peak is  $17 \text{ cm}^{-1}$  higher in frequency than the terrace peak. The redshift for the CO stretch on the steps is rationalized by weakening of the C=O bond due to electron transfer from the metal to  $2\pi^*$  orbitals of CO leading to a stronger bond softening for CO on the step edges. For the  $\text{CH}_3(\text{ads})$ , our RAIRS data indicate the opposite trend: stronger C-H bond softening occurs on the terrace site consistent with calculations reported by Michaelides and Hu<sup>57</sup>. These authors examined the C-H stretch frequency of  $\text{CH}_3(\text{ads})$  on top and threefold hollow sites on Pt(111) and Ni(111) and calculated a stronger redshift for the hollow sites due to the C-H bonds being closer to the metal atoms than for the top sites, consistent with experimental values for  $\text{CH}_3(\text{ads})$  on Ni(111) and Pt(111). We do not observe the same splitting in the overtone of the bend, showing this vibration is affected less by this bond softening.

The non-invasive nature of RAIRS allows us to monitor the simultaneous, site-selective uptake of  $\text{CH}_3(\text{ads})$  on step and terrace sites by recording spectra throughout a  $\text{CH}_4$  deposition experiment. Fig. 3B presents RAIRS spectra taken during a measurement at  $E_k =$

65 kJ/mol,  $T_S = 120$  K without laser excitation. The higher frequency peak appears first, consistent with  $\text{CH}_4$  dissociation on the more reactive steps. As the step peak saturates, the peak at  $2886\text{ cm}^{-1}$  grows more slowly in intensity indicating  $\text{CH}_3(\text{ads})$  on the terraces. With increasing  $\text{CH}_3(\text{ads})$  coverage on the terraces, the higher frequency step peak broadens, most likely due to interaction between  $\text{CH}_3(\text{ads})$  on steps and on terraces.

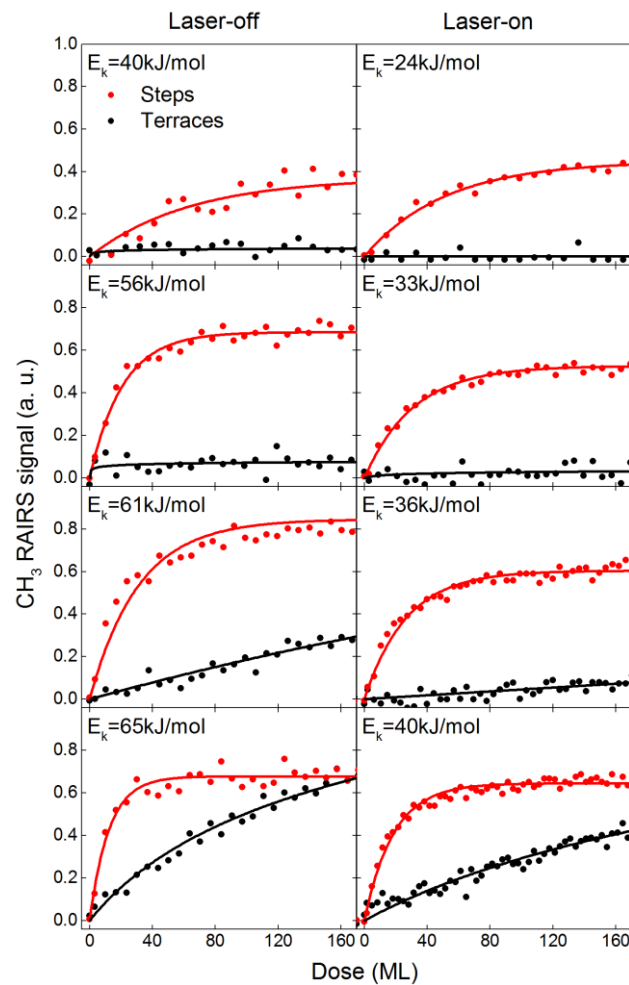


FIG. 4. Uptake curves measured under laser-off (left panels) and laser-on (right panels) conditions for dissociation of methane on the steps (red points) and terraces (black points) on Pt(211) at  $T_S = 120$  K. The solid lines are fits to the data obtained using Eq. (7).

Employing the analysis described in Section II, the RAIRS spectra are converted into the uptake curves shown in Fig. 4, where the  $\text{CH}_3$  RAIRS signal is plotted as a function of incident methane dose for uptake on the steps (red points) and terraces (black points) for measurements without laser excitation (left hand column) and laser-on depositions where

between 10% and 45% of the incident methane is prepared with a single quantum of antisymmetric stretch ( $\nu_3$ ) vibration (right hand column). At  $E_k = 40$  kJ/mol without laser excitation, uptake is only observed on the steps. Increasing  $E_k$  or adding a quantum of  $\nu_3$  vibration leads to  $\text{CH}_3(\text{ads})$  uptake on both steps and on the terraces, demonstrating that methane dissociation on Pt(211) is promoted by both translational and vibrational energy.

Without laser excitation,  $\text{CH}_3(\text{ads})$  uptake is first seen on the steps of Pt(211) at  $E_k = 33$  kJ/mol (data not shown) and on the terraces at  $E_k = 61$  kJ/mol. From these values, we estimate the activation barrier on the step to be at least 30 kJ/mol lower than the activation barrier on the terrace. This is consistent with calculated values of 42 and 37 kJ/mol for the difference between the activation energy on the terrace and on the two lowest barrier sites on the step edge. The experimental value establishes only a lower limit for the true difference in activation barriers as the increase in  $E_k$  was achieved by nozzle heating, which will thermally populate low lying bending vibrations of  $\text{CH}_4$ . With increasing nozzle temperature more molecules are in vibrationally excited states, leading to a stronger increase in reactivity than would be expected from increasing only  $E_k$ .

The theoretical activation barriers were calculated for methane dissociation on a clean Pt(211) surface, whereas the RAIRS uptake curves above show that there is significant  $\text{CH}_3(\text{ads})$  on the steps when methane dissociation is observed on the terraces. As  $E_k$  increases, the maximum (saturation)  $\text{CH}_3(\text{ads})$  coverage on the Pt(211) surface increases, as shown by the increasing asymptotes of the uptake curves for both the laser-off and laser-on data. The saturation coverage of  $\text{CH}_3(\text{ads})$  on the steps also decreases slightly when there is significant  $\text{CH}_3(\text{ads})$  uptake on the terrace. Both these observations suggest that  $\text{CH}_3(\text{ads})$  increases the effective activation barrier height for methane dissociation compared to that on a clean Pt(211) surface, as has been observed previously on Pt(111)<sup>58</sup>. DFT calculations in that work

showed that this was due to substantial charge transfer between the metal and the adsorbed molecules, leading to repulsive forces between the dissociating CH<sub>4</sub> and CH<sub>3</sub>(ads)<sup>58</sup>.

A previous study on the dissociative chemisorption of methane on the stepped Pt(322) and Pt(355) surfaces performed at T<sub>S</sub> ≈ 120 K suggested that the CH<sub>3</sub>(ads) formed on the terraces migrated to the steps<sup>59</sup>. The uptake curves presented in Fig. 4 show no evidence of CH<sub>3</sub>(ads) migration from step to terrace sites on Pt(211) at T<sub>S</sub> = 120 K. If the methane dissociated on the terraces and the resultant CH<sub>3</sub>(ads) migrated to the steps, then uptake on the terraces should only be detected once the steps are saturated, which is inconsistent with the data shown in Fig. 4. Furthermore, we detect no change in the RAIRS signal for the steps and terraces due to CH<sub>3</sub>(ads) migration from the terraces to steps when the CH<sub>4</sub> deposition is stopped before the steps are saturated. These observations are fully consistent with our DFT calculations. Using the SRP functional, and allowing the top two layers of the metal to fully relax, the computed activation energy for CH<sub>3</sub>(ads) diffusion from the terrace to the more stable step site is 64.4 kJ/mol, and the corresponding rate constant is only about 10<sup>-15</sup> s<sup>-1</sup> at 120 K.

The reactivity of the methane on the step and terrace sites can be quantified by the initial sticking coefficient which is the probability that the molecule dissociates on either the step ( $S_0(\text{step})$ ) or the terrace ( $S_0(\text{terr})$ ) on the clean Pt(211) surface. This value is proportional to the initial gradient of the uptake curves presented in Fig. 4. Increasing  $E_k$  or adding a quantum of  $\nu_3$  leads to an increase in both  $S_0(\text{step})$  and  $S_0(\text{terr})$  as more incident energy is available to overcome the activation barrier for the dissociation. To be able to obtain absolute values of  $S_0$ , the data were fit using a Langmuir type uptake model:

$$\frac{d\theta}{d\varepsilon} = S_0(1 - n_s\theta), \quad (7)$$



where  $\Theta$  is the coverage,  $\varepsilon$  is the dose and  $n_s$  is the number of sites that each adsorbate blocks on the surface. Each uptake was fit independently, i.e. assuming that the uptake of  $\text{CH}_3(\text{ads})$  on the terrace was not influenced by  $\text{CH}_3(\text{ads})$  uptake on the steps, and vice versa.

In order to be able to fit the data, it is necessary to convert the  $\text{CH}_3$  RAIRS signal to  $\text{CH}_3(\text{ads})$  coverage. King and Wells (K&W) measurements<sup>60,61</sup> were used to determine the conversion for the step sites. Deposition measurements were done at  $E_k = 65$  kJ/mol and were monitored using K&W and RAIRS (64 scans/spectra) simultaneously. We assumed that under these conditions  $S_0(\text{terr})$  made a negligible contribution to the sticking coefficient obtained from the K&W measurement. The RAIRS data were analyzed as above to obtain the uptake curve of  $\text{CH}_3$  RAIRS signal against dose. The y-axis was then scaled so that the initial gradient of the uptake curve matched the value of  $S_0$  obtained from the K&W measurement. This procedure was repeated five times, and the average result was used as the conversion from  $\text{CH}_3$  RAIRS signal to coverage for all the uptake curves on the steps. The conversion factor for the terraces was taken to be the same as that for  $\text{CH}_3(\text{ads})$  on a Pt(111) surface.

The state-resolved sticking coefficients,  $S_0^{v_3}$  were found from the laser-off ( $S_0(\text{laser-off})$ ) and laser-on ( $S_0(\text{laser-on})$ ) sticking coefficients using<sup>14</sup>:

$$S_0^{v_3} = \frac{S_0(\text{laser-on}) - S_0(\text{laser-off})}{f_{\text{exc}}} + S_0(\text{laser-off}). \quad (8)$$

In Fig. 5, the values of  $S_0(\text{step})$  (filled points, panel A) and  $S_0(\text{terr})$  (filled points, panel B) are compared with those calculated using the RPH model (solid lines and open points) described in Section III<sup>19-24</sup>. Good agreement between theory and experiment is observed for dissociation on the step sites, where the experimental  $S_0(\text{step})$  has been calibrated using K&W measurements<sup>60,61</sup>. Theory also confirms that contributions to the total sticking from reactions on the terrace sites is negligible at  $E_k = 65$  kJ/mol, as was assumed in the calibration of  $S_0(\text{step})$ . The theoretical laser-off values were estimated by including contributions from

vibrationally excited states at the experimental nozzle temperatures. Given the low activation barriers and large  $S_0$  for ground state reaction at the step edge, these effects are relatively minor in panel 5A. Finally, we note that  $S_0$  computed using the PBE functional gives a reasonable agreement with experiment (not shown). However, rescaling our MEP barrier heights to the SRP values leads to the improved agreement in Fig. 5A.

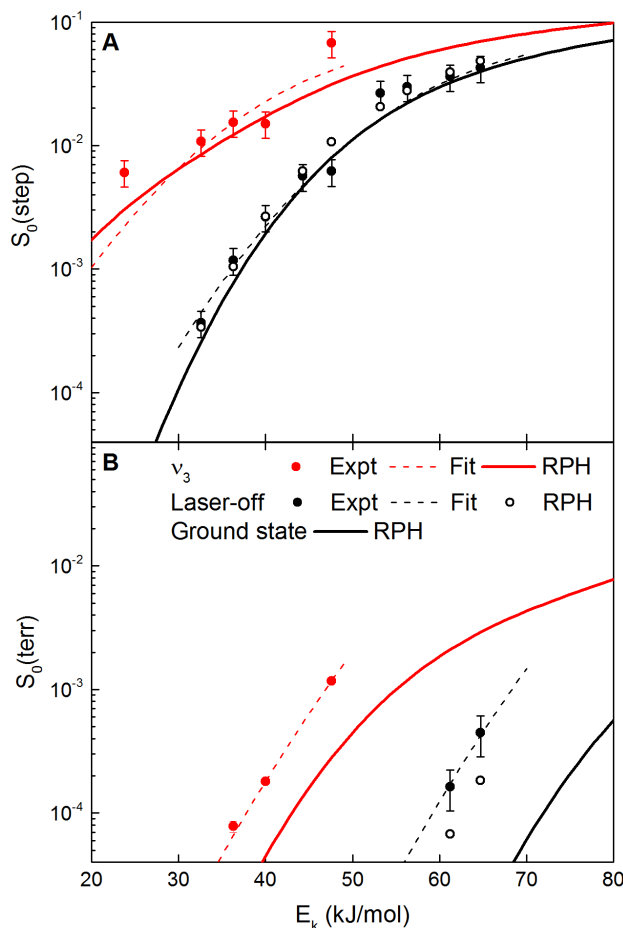


FIG. 5. A. A comparison of experimentally determined state-resolved sticking coefficients for  $\text{CH}_4$  prepared with a single quantum of  $v_3$  (red symbols) and under laser-off conditions (black symbols) for dissociation on the steps of Pt(211) at  $T_S = 120$  K with those obtained from RPH calculations using the SRP functional (solid lines). The dashed lines are S-shape curve fits to the experimental data using Eq. (9). B. As panel A, but for dissociation on the terraces.

The agreement between theory and experiment for dissociative sticking on the terrace sites is less good. However, the experimental  $S_0(\text{terr})$  were estimated using the same conversion factor obtained for Pt(111)<sup>26</sup>. It is clear from Table II that the energetics of methane chemisorption on the terraces of Pt(111) and Pt(211) are different, and the

magnitude of the derivative dipole coupling for the symmetric stretch is likely to be different for  $\text{CH}_3$  adsorbed on these two sites. If we rescale the experimental  $S_0(\text{terr})$  by 0.3, the agreement improves significantly. Note that for the higher activation barriers and smaller values of sticking at the terrace sites, the laser-off corrections are much more important.

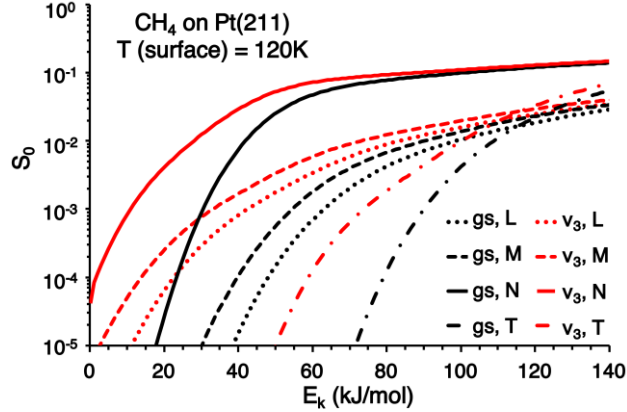


FIG. 6. Computed contributions to the dissociative sticking probability of methane from the three step (L, M, N) and one terrace (T) transition states. Results are shown for methane initially in the ground state (gs) and the  $1\nu_3$  excited state, at  $T_S = 120$  K.

It is interesting to compare the contributions to the total sticking from the 4 primary reaction pathways. In Fig. 6, we plot the computed contributions to the total dissociative sticking probability from the four reaction paths, at  $T_S = 120$  K, for the PBE functional. For molecules initially in the ground state, the contribution to the reactive sticking from the N MEP is an order of magnitude larger than that from any other path, except at the highest energies, where other paths can contribute several percent. Excitation of the antisymmetric stretch ( $1\nu_3$ ) promotes reaction at all sites. The vibrational enhancement is largest for the L and M paths and for dissociation on the terrace sites. This is consistent with the elongation of the reactive bond at the transition state being larger for these paths (see Table II).

To compare the effect of adding kinetic or vibrational energy, the experimental data were fit using S-shape curves (dashed lines, Fig. 5) given by<sup>14,62</sup>:

$$S_0(E_k) = \frac{A}{2} \left[ 1 + \operatorname{erf} \left( \frac{E_k - E_0}{W} \right) \right], \quad (9)$$

where  $A$  is the asymptotic value of  $S_0$  at infinitely high incident energy,  $E_0$  is the average activation barrier height and  $W$  the width of the distribution of activation barriers. The laser-off and  $v_3$  data were fit simultaneously with  $A$  and  $W$  restricted to be the same for both sets of data. In the fit, each value of  $S_0$  was weighted by  $1/\sigma$ , where  $\sigma$  is the error bar on the experimental data representing 68% confidence limits. The vibrational efficacy ( $\eta^{v_3}$ ) for each site on the surface was then found using<sup>62</sup>:

$$\eta^{v_3} = \frac{(E_0(\text{laser-off}) - E_0(v_3))}{E_{\text{vib}}}, \quad (10)$$

where  $E_{\text{vib}}$  is the vibrational energy of a single quantum of antisymmetric stretch vibration (36 kJ/mol). The efficacy quantifies the relative efficiency with which vibrational energy promotes reactivity compared to the same amount of kinetic energy.  $\eta^{v_3}(\text{step}) = 0.45$  and  $\eta^{v_3}(\text{terr}) = 0.59$ , showing adding a quantum of  $v_3$  promotes the reactivity more on the terraces of Pt(211) than the steps but that kinetic energy is more efficient on both sites. The efficacies extracted from the theoretical sticking curves typically vary with  $S_0$ . We find that the computed step efficacies vary over the range of the data from  $\eta^{v_3}(\text{step}) = 0.50$  at  $S_0 = 3 \times 10^{-3}$  to  $\eta^{v_3}(\text{step}) = 0.37$  at  $S_0 = 10^{-1}$ , in good agreement with experiment. On the other hand, for the terrace site we find that  $\eta^{v_3}(\text{terr}) = 0.79$  at  $S_0 = 10^{-4}$  and  $\eta^{v_3}(\text{terr}) = 0.76$  at  $S_0 = 10^{-3}$ . However, we note that the experimental efficacies are computed relative to the laser off curves, while the theoretical efficacies are relative to the true ground state. While this doesn't matter for reaction on the steps, we have shown that vibrationally excited molecules make a significant contribution to  $S_0(\text{terr})$ . If we compute the  $v_3$  efficacies relative to the two (theory) laser off data points in Fig. 5B, we find  $\eta^{v_3}(\text{terr}) = 0.53$  and  $0.55$ , much closer to the experimental result. All of the above results are consistent with the reactive C-H bond being more stretched at the transition state for dissociation on the terrace site, 1.525 Å, than at the transition state for the dominant reaction path along the step edge, 1.480 Å (see Table II). As

discussed in several recent papers<sup>19–21,23,24,63,64</sup>, this vibrational enhancement in  $S_0$  arises from a coupling of one of the antisymmetric stretch modes to the symmetric stretch, which correlates adiabatically with the dissociating bond. The sudden vector projection model proposed by Guo and co-workers provides another qualitative way to compare efficacies<sup>65</sup>, but we have not done such calculations here.

## V. SUMMARY

It has long been proposed that reactions on most catalysts occur preferentially on step edges and other defect sites. We report here the results of a collaboration between experiment and theory that examines the dissociative chemisorption of methane on the stepped Pt(211) surface, where both the experimental and theoretical methods are resolved with respect to reaction site and the quantum state of the incident molecule. The dissociative sticking of  $\text{CH}_4$  produces  $\text{H}(\text{ads})$  and  $\text{CH}_3(\text{ads})$  fragments, and the methyl product is stable at our experimental surface temperature of 120 K. We demonstrate that it is possible to use RAIRS detection of the methyl fragment to distinguish between  $\text{CH}_3(\text{ads})$  on the step and terrace atoms of Pt(211). The dissociative sticking of  $\text{CH}_4$  is measured over a range of incident energies by using this site-specific RAIRS detection to monitor the formation of  $\text{CH}_3(\text{ads})$  on the step and terrace sites. We find that the dissociation on Pt(211) is a direct reaction on both sites and that diffusion of the  $\text{CH}_3(\text{ads})$  product does not occur at  $T_s=120$  K.  $\text{CH}_4$  reactivity on Pt(211) is dominated by dissociation on step sites due to a lower barrier of at least 30 kJ/mol compared to the terrace sites. There is no evidence for dissociation at the corner sites on Pt(211).

DFT studies are consistent with these observations. Methyl groups prefer the top sites on Pt(211), and the binding is strongest over the atoms along the step. Consequently, the barriers are lowest there. We find that the activation energy for dissociation on the terrace is

84 kJ/mol, while on the step edge we find three transition states with activation energies of 42, 47 and 55 kJ/mol. The barrier for dissociation over the corner atoms is very large, 183 kJ/mol, and reaction is not likely to occur there at the energies considered in this study. We also find that the barrier to methyl diffusion on the surface is large, 64 kJ/mol, and that diffusion should not occur at 120 K.

In the experimental sticking results methane dissociation on both sites is promoted by adding kinetic energy or a single quantum of antisymmetric stretch vibration. First principles quantum scattering calculations of the dissociative sticking probability are in very good agreement with experiment for reaction at the step sites. For reaction at the terrace sites the agreement is not as good, with experiment and theory differing by a factor of three. Unlike sticking on the step sites, it was not possible to independently calibrate sticking on the terrace. A factor of 3 is reasonable given variation in the RAIRS conversion factor and errors in the theory. Theory finds that reaction at the lowest barrier step site, where the molecule dissociates along the edge, dominates the sticking at all but the highest energies. The measured vibrational efficacy is found to be less than one on both the step and the terrace, but is larger for dissociation on the terrace than on the step. This is consistent with our DFT studies, which find that the elongation of the dissociating bond is larger at the transition state for dissociation on the terrace than over the lowest barrier step site. Our calculations of sticking reproduce the experimentally measured efficacies for both the step and terrace sites.

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