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Prediction and Diagnosis of Tropical Cyclone Formation in an NWP System. Part I: The Critical Role of Vortex Enhancement in Deep Convection

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ABSTRACT

This is the first of a three-part investigation into tropical cyclone (TC) genesis in the Australian Bureau of Meteorology's Tropical Cyclone Limited Area Prediction System (TC-LAPS), an operational numerical weather prediction (NWP) forecast model. The primary TC-LAPS vortex enhancement mechanism is presented in Part I, the entire genesis process is illustrated in Part II using a single TC-LAPS simulation, and in Part III a number of simulations are presented exploring the sensitivity and variability of genesis forecasts in TC-LAPS.

The primary vortex enhancement mechanism in TC-LAPS is found to be convergence/stretching and vertical advection of absolute vorticity in deep intense updrafts, which result in deep vortex cores of 60–100 km in diameter (the minimum resolvable scale is limited by the 0.15° horizontal grid spacing). On the basis of the results presented, it is hypothesized that updrafts of this scale adequately represent mean vertical motions in real TC genesis convective regions, and perhaps that explicitly resolving the individual convective processes may not be necessary for qualitative TC genesis forecasts. Although observations of sufficient spatial and temporal resolution do not currently exist to support or refute this proposition, relatively large-scale (30 km and greater), lower- to midlevel tropospheric convergent regions have been observed in tropical oceanic environments during the Global Atmospheric Research Programme (GARP) Atlantic Tropical Experiment (GATE), the Equatorial Mesoscale Experiment (EMEX), and the Tropical Ocean Global Atmosphere Coupled Ocean–Atmosphere Response Experiment (TOGA COARE), and regions of extreme convection of the order of 50 km are often (remotely) observed in TC genesis environments. These vortex cores are fundamental for genesis in TC-LAPS. They interact to form larger cores, and provide net heating that drives the system-scale secondary circulation, which enhances vorticity on the system scale akin to the classical Eliassen problem of a balanced vortex driven by heat sources. These secondary vortex enhancement mechanisms are documented in Part II.

In some recent TC genesis theories featured in the literature, vortex enhancement in deep convective regions of mesoscale convective systems (MCSs) has largely been ignored. Instead, they focus on the stratiform regions. While it is recognized that vortex enhancement through midlevel convergence into the stratiform precipitation deck can greatly enhance midtropospheric cyclonic vorticity, it is suggested here that this mechanism only increases the potential for genesis, whereas vortex enhancement through low- to midlevel convergence into deep convective regions is necessary for genesis.

1. Introduction

As part of the ongoing development of the Australian Bureau of Meteorology's tropical cyclone (TC)

version of the Limited Area Prediction System (LAPS), a detailed investigation of the TC life cycle in TC-LAPS is in progress. Clearly one of the most important aspects of the TC forecast is genesis; that is, whether a TC will develop in the forecast area or not. Before improvements can be made to the prediction system it is necessary to develop a basic understanding of genesis in the model, and an understanding of how realistic the model depiction of genesis is. As is frequently noted in

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the literature regarding genesis and intensification (e.g., Gray 1998) there are many wrong ways to the right answer. Thus, if the model is performing well for the incorrect reasons, improvements in the realism of certain aspects of the modeling system could in fact reduce the performance of the system. With this in mind, a detailed examination of TC genesis in TC-LAPS is underway and the first round of results is presented in this study and two companion papers, Tory et al. (2006, hereafter Part II; 2005, manuscript submitted to *J. Atmos. Sci.*, hereafter Part III). In this paper (Part I), the TC-LAPS genesis processes are introduced, and the primary vortex enhancement mechanism illustrated. Part II provides a detailed diagnostic analysis of the genesis processes and includes forecast verification for a simulation of TC Chris (February 2002, off the western Australian coast). Sensitivity and variability of the genesis processes are examined in Part III, where a number of developing and nondeveloping TC genesis simulations are presented.

Tropical cyclone genesis has been described as the process that leads to the development of a self-sustaining surface-concentrated vortex, in which the flux of energy from the sea surface to the vortex, governed partly by the intensity of the vortex, is sufficient to maintain and amplify the vortex; that is, the hurricane heat engine [e.g., Wind Induced Surface Heat Exchange (WISHE); Emanuel 1986; Rotunno and Emanuel 1987]. Alternatively, Saunders and Montgomery (2004), Hendricks et al. (2004, hereafter H04), and Montgomery et al. (2006, hereafter M06), each described the genesis process as the development of a warm core vortex that extends from the surface to at least the midtroposphere. This is independent of the positive feedback between the surface fluxes and wind speed as described by WISHE (i.e., it is pre-WISHE). Both definitions require the generation of finite-amplitude cyclonic surface vorticity.

Over the last decade or so the search for the process that generates such a finite amplitude surface vortex has focused on the observation that TC formation is associated with mesoscale convective systems (MCSs) and their accompanying mesoscale convective vortices (MCVs). Although the dynamics and thermodynamics of midlatitude terrestrial MCSs and their accompanying MCVs has been extensively studied and reasonably well documented (e.g., Fritsch et al. 1994; Houze 2004), comparatively little has been published about their tropical oceanic counterparts. Mapes and Houze (1995, hereafter MH95) used airborne Doppler radar to construct vertical profiles of mean horizontal divergence in tropical oceanic MCSs. These profiles can provide insight into the type of vorticity structure one might ex-

pect in MCSs. They found the profiles to consist mostly of just the two deepest internal modes spanning the troposphere; these being lower tropospheric convergence with upper tropospheric divergence, typically associated with deep convective precipitation, and midtropospheric convergence combined with upper and lower tropospheric divergence, typically associated with stratiform precipitation. Hereafter we will refer to these two divergence profiles as convective divergence profiles (CDP) and stratiform divergence profiles (SDP) respectively. This result is consistent with Raymond et al. (1998) who investigated the mean vorticity, divergence, and vertical mass flux within MCSs associated with developing TCs observed during TEXMEX. Often in the genesis process SDP was significant early on, but later CDP dominated.

The well-documented midlatitude terrestrial MCS is dominated by stratiform precipitation at maturity, and consequently an SDP. Because convergence (divergence) increases (decreases) the absolute vorticity magnitude in a rotating environment, the MCV consists of a cyclonic vortex maximized in the middle troposphere. Evaporative cooling below the MCS is often responsible for a surface anticyclone (Johnson et al. 1989). Hereafter we label vortex intensification in a stratiform midlevel convergence region stratiform vortex enhancement (SVE), and vortex intensification in a convective low- to midlevel convergence region convective vortex enhancement (CVE). MH95 identified examples of almost pure SDP and CDP, averaged over diameters of 30–60 km, although most were comprised of combinations of both (see also Houze 1997). The absolute vorticity magnitude in a rotating environment dominated by CDP would be expected to increase throughout the low- to midtroposphere (i.e., CVE) and decrease above. The majority of profiles presented in MH95 included nontrivial convergence from the surface to mid-troposphere, which suggests CDP often dominated over SDP at low levels. Such mean divergence profiles within a TC genesis environment (nontrivial low- to midlevel cyclonic absolute vorticity) should enhance cyclonic vorticity from the low to midlevels and lead to the generation of cyclonic vortex cores of equivalent depth.

The TC genesis theory proposed in Simpson et al. (1997) and Ritchie and Holland (1997), did not consider CVE in the genesis process. Although they acknowledged that it could exist, they believed it to act only as a TC enhancement mechanism after genesis was complete. The TC genesis theory of Bister and Emanuel (1997) considered CVE only in the final stages of genesis. Both theories were based on the midlatitude terrestrial MCV conceptual model that includes a sur-

face anticyclone (as mentioned above) and a stratiform precipitation region of considerably greater horizontal scale than the convective region. Thus these theories first required a mechanism to replace the surface anticyclone with cyclonic vorticity from above. The main focus of both theories was on mechanisms that bring midlevel cyclonic vorticity to the surface. Using a mid-latitude terrestrial MCV conceptual model as a basis for their genesis theories may have been misleading, however, because they did not consider the potential for CVE.

In this paper we advance the hypothesis that CVE is critical for TC genesis. Since, as of yet, there is insufficient observational data to prove or refute this proposition, it must remain a hypothesis until tested further. We acknowledge that SVE is likely to play a role in TC genesis by enhancing large areas of midlevel vorticity, particularly early in the genesis process, but we suggest genesis will not proceed without CVE. Furthermore, in some instances there might be a system-scale transition from vortex intensification dominated by SVE to that dominated by CVE. (Note, this transition is opposite to the typical individual MCS life cycle, where the young MCS is dominated by convective precipitation, before the stratiform region has time to grow.) If such a transition does exist it may vary from basin to basin. For example the dry Saharan air layer in the Atlantic basin is likely to favor SVE by increasing low-level evaporation, which both strengthens the stratiform circulation, and inhibits the growth of large areas of near downdraft free convection.

There are a number of modeling studies in support of the hypothesis that CVE is critical for genesis. Montgomery and Enagonio (1998) considered CVE when they looked at the interaction of an MCV with a low- to midlevel vortex core that would likely be generated by vorticity convergence into a convective hot tower complex. They found a single upright cyclonic vortex core resulted, and proposed such an interaction may provide the necessary low-level vortex enhancement for TC genesis. More recent studies by H04 and M06 found CVE to be critical for TC genesis in both realistic and idealized models. They found intense vortices were generated in the low- to midtroposphere through vorticity convergence into convective hot towers. These vortices interacted to form larger low- to midlevel vortices. They also found the sum of diabatic heating from condensation and adiabatic cooling from expansion in the hot towers was slightly positive, and the net heating from all hot towers enhanced the system-scale secondary circulation in a process akin to the Eliassen balanced vortex model forced by heat sources (Eliassen 1951). They termed these two processes the vortex up-

scale cascade, and the system scale intensification (SSI), respectively.

It is interesting to note that CVE was responsible for the intensification in Kurihara and Tuleya's (1981) pioneering simulation of a tropical storm (TS) that developed in an easterly wave. Furthermore, their temperature budgets were remarkably similar to H04 and M06 (and the TC-LAPS budgets presented in Part II), and suggest the SSI process played an important role.

More recently, an observational study by Reasor et al. (2005) identified low-level vortex enhancement in the vicinity of a convective hot tower (consistent with H04 and M06) in Doppler radar observations during the genesis of Hurricane Dolly. They also identified midlevel vortex enhancement in a stratiform precipitation region. The spatial resolution of their observational data was sufficient to identify both CVE and SVE as separate entities within regions of similar scale to MH95's averaging domains. Although they did not carry out a quantitative partitioning of the two, their results suggest that both CVE and SVE act together to enhance cyclonic vorticity in a TC genesis environment.

In Part I of this series we document the primary vortex enhancement mechanism in TC-LAPS, and show that CVE is responsible for the construction of a low- to midtropospheric cyclonic vorticity core. We show in Part II how these cores contribute to the construction of a cyclonic potential vorticity (PV) monolith through the vortex upscale cascade and SSI processes of H04 and M06. In Part III we begin to identify common TC genesis features, and necessary conditions for TC genesis in TC-LAPS, using simulations of a number of developing and nondeveloping TCs.

An outline of Part I is as follows. The modeling system used, TC-LAPS, is described in section 2. The primary low-level vortex enhancement mechanism active in all TC-LAPS genesis simulations is illustrated in section 3. Observational evidence in support of this mechanism and the ability of TC-LAPS to adequately represent the mechanism is presented in section 4. The work is summarized in section 5.

2. TC-LAPS: Model description

LAPS is an operational Numerical Weather Prediction (NWP) forecast model. A number of domains of varying sizes and horizontal resolution are run twice daily, together with a global model, to make up the suite of NWP forecasts produced by the Australian Bureau of Meteorology. The largest LAPS domain is illustrated in Fig. 1. It encompasses the Australian continent and much of the surrounding oceans and seas, and neighboring islands (longitude 65.0° to 184.625°, latitude -65.0° to 17.125°). The horizontal grid resolu-

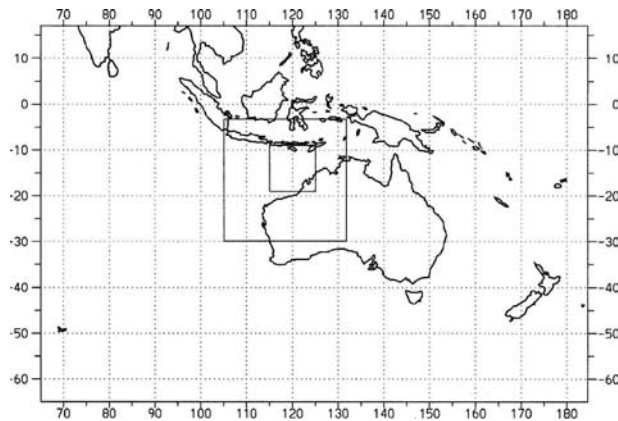


FIG. 1. LAPS375 domain, with the TC Chris TC-LAPS domain embedded. The innermost rectangle represents the subdomain featured in Fig. 2.

tion of this domain is 0.375° (LAPS375). LAPS375 is nested in the Bureau of Meteorology's global model Global Assimilation Prediction System (GASP). Three mesoscale versions of LAPS are nested (one-way) in LAPS375, as well as a TC version (TC-LAPS) for monitoring tropical systems of interest that develop in the LAPS375 domain.¹ By its nature the domain of TC-LAPS is not fixed. An example of the TC-LAPS domain embedded in LAPS375 is shown in Fig. 1 (large rectangle). It is the TC-LAPS domain used in the simulation of TC Chris that is discussed in the next section, and in Part II. The operational resolution is 0.15° on 29σ levels and covers a $27^\circ \times 27^\circ$ area centered on the system of interest.

A detailed description of LAPS is provided by Puri et al. (1998), and summary of the LAPS components, including a number of additional options now available and used in TC-LAPS, is outlined here. The Miller–Pearce time-stepping scheme (Miller and Pearce 1974) is combined with a third-order upwinding advection scheme, and implemented on an Arakawa A-grid (non-staggered). The European Centre for Medium-Range Weather Forecasts (ECMWF) boundary layer parameterization and land surface schemes (ECMWF 1995) are now available and are used in all operational LAPS models. The vertical diffusion parameterization employs a Monin–Obukhov surface layer with first-order closure schemes for determining the exchange coefficients (eddy diffusivities) above the surface. Under stable conditions the exchange coefficients are determined from mixing length theory using Richardson

number-dependent stability functions. For unstable conditions they are scaled by the boundary layer height, based on Troen and Mahrt (1986). Each version of LAPS now includes 29 levels in the vertical, including nine in the lowest (approximately) 1000 m.

The Tiedke mass flux convection scheme (Tiedke 1989) is used to parameterize subgrid-scale convection. The penetrative convection component is triggered where subcloud-layer moisture convergence is positive. The cloud-base mass flux is determined from the subcloud-layer moisture convergence, which together with the mass convergence in the lower half of the cloud due to organized entrainment provides the basic closure for the parameterization scheme. The scheme takes into account: turbulent entrainment and detrainment; cloud-top detrainment at and above the zero-buoyancy level; downdrafts, when sufficient dry environmental air mixed into the cloud evaporates cloud water and cools the air until it becomes negatively buoyant; precipitation; and evaporation of precipitation in the cloud and in cloud air detrained into the environment. This convection parameterization influences only the grid-scale temperature and humidity tendencies. For grid-resolved updrafts, any moisture in excess of 100% humidity is removed as rain.

The radiation scheme (Fels and Schwartzkopf 1975) has not changed since Puri et al. (1998), although the radiation tendencies are now updated hourly instead of 3-hourly to better resolve the diurnal cycle. The analysis system (multivariate statistical interpolation; MVSII) is also largely unchanged. An option to use sea surface temperature (SST) analyzed daily has been added, but not used here (weekly averaged SST is used). Currently, there is no feedback between the ocean and atmosphere in any of the LAPS systems.

Dynamic initialization

The initialization procedure follows Davidson and Puri (1992) with only a few changes. A summary of the current procedure follows. For genesis simulations, no bogus vortices are included in the initialization. Initialization is performed in two steps: 1) the analysis step, where high-resolution analyses (called target analyses) are obtained at 6-hourly intervals during the initialization period; and 2) the dynamical nudging step. Dynamical nudging is used to “grow” a numerically consistent and balanced initial state that best represents the real atmosphere, by nudging a numerical forecast, during the simulation, toward an analyzed state of the atmosphere (the target analyses generated in step 1). In the current study and in the operational TC-LAPS the dynamical nudging is performed over the 24-h period prior to the initial forecast time.

¹ An additional LAPS domain is centered on the equator (T-LAPS, longitude 70.0° to 189.625° , latitude -45.0° to 44.625°). It provides nesting for TC-LAPS in tropical regions north of the LAPS375 domain.

1) ANALYSIS STEP

The target analyses are produced using the objective analysis system described in Puri et al. (1998), and an assimilation scheme involving four 6-h model integrations. LAPS375 or tropical LAPS (T-LAPS) analyses are used to initialize the first 6-h model integration and for boundary conditions throughout the assimilation. Objective analyses are performed at the end of each 6-h model integration, and the resulting fields provide both the initial conditions for the next 6-h model integration, and the target analyses used in the dynamic nudging step. The observation base used by the objective analysis includes all standard observational data, together with the additional scatterometer and surface wind observations included in the research version of TC-LAPS used here.

2) DYNAMICAL NUDGING STEP

The $T = -24$ h target analysis is used to initialize TC-LAPS during the 24-h period of dynamical nudging. Linear interpolation in time between the neighboring target analyses provides a target analysis at every nudging time step during the procedure. Throughout the 24-h period the numerical solution is nudged toward the analyzed vorticity, while allowing the model to develop its own divergence. This leads to the preservation of the observationally reliable rotational wind component in the target analyses, while allowing the model-generated divergent wind field to dominate over the less reliable analyzed divergent winds. The surface pressure and temperature fields are also nudged toward the target analyses to preserve the mass field.

As mentioned in Davidson and Puri (1992), this form of dynamical nudging does not guarantee the convection will develop in the right place or time. To overcome this problem cloud-top temperature nudging is applied, where artificial heat sources are used to force model convection in regions of deep convection identified by satellite observations. This forcing is applied during the assimilation process as well as the dynamical nudging. The method is well described in Davidson and Puri (1992) and will not be elaborated on here except where changes to the procedure have been made. These changes include the reduction of the triggering temperature (the maximum cloud-top temperature at which the heat source is applied) from 273 to 233 K, to ensure forced convection is only applied where the observed convection is deep. Another change, although minor, involves nudging of the numerically predicted atmosphere toward the imposed heating profile, rather than the direct replacement of the convective param-

eterization with the heating profile, as mentioned in Davidson and Puri (1992).

3. Primary TC-LAPS vortex intensification mechanism

Here we document the primary vortex intensification mechanism present in all TC-LAPS genesis simulations: horizontal convergence (or stretching) of absolute vorticity in deep convective cores (i.e., CVE). The mechanism in itself is not responsible for vortex amplification to TC intensity. Instead the mechanism builds vortex cores that interact with one another (H04 and M06's diabatic upscale vortex cascade) and provide heating that fuels the SSI process (H04; M06). (Hereafter, we label the diabatic upscale vortex cascade and SSI processes as secondary vortex enhancement mechanisms.) In this way the primary mechanism indirectly contributes to the construction of a monolithic cyclonic vortex of sufficient intensity and scale for the self-sustaining and self-amplification process (e.g., WISHE amplification) to set in.² The construction of the vortex monolith in a simulation of TC Chris is analyzed and described in detail in Part II.

The horizontal convergence/stretching of absolute vorticity in deep convective cores was described by H04 and M06. They labeled the resulting vortices vortical hot towers (VHTs). The main difference between the TC-LAPS convective vortex cores and VHTs is the scale of the updrafts and resulting vortices. As mentioned in section 2, the TC-LAPS horizontal grid spacing is 0.15° (about 15 km), giving a minimum resolvable convective scale of about 60 km. Typically the resolved updrafts are of the order of 60–100 km in diameter,³ about 4–5 times greater than those documented in H04 and M06. It is worth noting that even on scales of 300 km Kurihara and Tuleya (1981) showed essentially the same vortex enhancement processes were active in their convective updrafts, forced onto such scales by a 0.625° grid spacing.

The convection in the TC-LAPS updrafts, which extend to depths greater than 14 km, is both explicit and parameterized. At 0.15° grid spacing the individual convective cells are not resolved and thus a parameterization scheme is required. The parameterization scheme,

² To test when WISHE might become active a sensitivity simulation was run in which the wind speed influence on surface fluxes was capped at 5 m s^{-1} . A significant difference in growth rate was not evident until after 18 h of model time, which suggests that WISHE is likely to become active at about this time.

³ Horizontal smoothing has been applied to the vertical motion plots featured in this paper, which tend to give the appearance of greater horizontal scale to the updrafts.

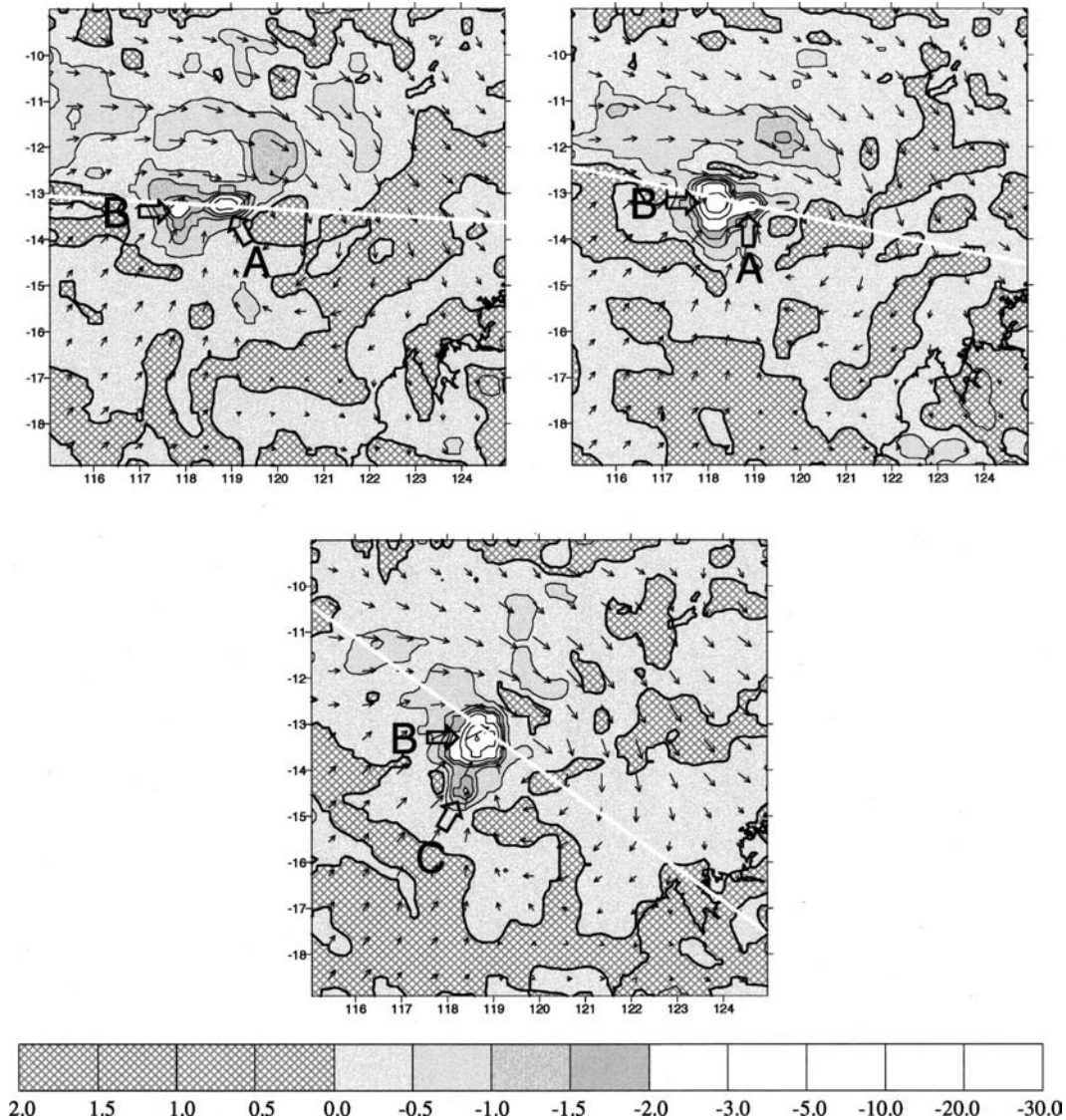


FIG. 2. Vertical motion on the $\sigma = 0.25$ level (Pa s^{-1} , shaded) with horizontal wind on the $\sigma = 0.9943$ level overlaid (vectors), 4 (left), 6 (right), and 8 (lower) hours into the simulation of TC Chris initialized at 1100 UTC 1 Feb 2002. The white lines indicate the positions of cross sections presented in Fig. 3, and convective updrafts/cores referenced in the text are labeled A, B, and C.

in theory, adjusts the temperature and moisture fields to represent the effect of the unresolved convection. It does not transport mass; instead the explicit, or resolved, updrafts represent the mean vertical motion within the larger convective region. It is not clear how realistic this combination is, but it is clear that some form of convective parameterization is required and that the parameterization alone cannot represent the relatively large convective divergence profiles observed by MH95 in tropical oceanic MCSs. Below the modeled mature updrafts, boundary layer subsidence (up to 400 m deep) is often present with associated cooling of up to 2 K (not shown). This subsidence is likely to have

developed in response to the temperature and moisture field adjustment associated with the parameterization of downdrafts and evaporative cooling.

Examples of the TC-LAPS updrafts are shown in Fig. 2. The images in this figure come from the simulation of TC Chris featured in Part II, and show a cycle of updraft development and decay between 4, 6, and 8 h into the simulation. Vertical motion on the $\sigma = 0.25$ level (approximately 10 km) is combined with the low-level horizontal winds to identify regions of deep intense updrafts embedded in low-level cyclonic flow. In the left panel there are two updraft cores labeled A and B. Core A is an old updraft undergoing decay, while core

B is relatively young. The right panel shows core B in its mature stage, with only a small remnant of core A remaining. The lower panel shows the emergence of a third updraft labeled C just prior to the decay of core B. Vertical cross sections of vertical velocity, absolute vorticity, and the contributions to absolute vorticity tendency from vertical advection and stretching, are shown in Fig. 3, for core B during the developing, mature, and decaying stages (same times as Fig. 2). The white lines in Fig. 2 indicate the locations of the cross sections. The second row of Fig. 3 shows the intensification of absolute vorticity at low to midlevels in the vicinity of core B. During this 4-h period the maximum cyclonic vortex intensity increased from about $-3 \times 10^{-4} \text{ s}^{-1}$ to about $-8 \times 10^{-4} \text{ s}^{-1}$ (Southern Hemisphere), with the maximum value located in the lowest 2 km.

In pressure coordinates the tendency equation for the vertical component of absolute vorticity (s_a) can be expressed as (Holton 2004)

$$\frac{\partial s_a}{\partial t} = -\mathbf{u}_h \cdot (\nabla_h s_a) - \omega \frac{\partial s_a}{\partial p} - (\nabla \cdot \mathbf{u}_h) s_a - \left(\frac{\partial \omega}{\partial x} \frac{\partial v}{\partial p} - \frac{\partial \omega}{\partial y} \frac{\partial u}{\partial p} \right),$$

where the subscript h refers to the horizontal components of the wind and the other variables have their usual meaning. The terms on the right-hand side represent horizontal advection, vertical advection, stretching/convergence, and tilting of absolute vorticity. All the contributions to the vertical vorticity tendency from these terms were calculated and the dominant vortex enhancement terms, vertical advection, and stretching/convergence, are illustrated in Fig. 3.

The contributions from the horizontal advection and tilting terms do not contribute to cyclonic vortex enhancement in the vicinity of the updraft core. Tilting opposed the contribution from vertical advection at the updraft edge. The contribution from horizontal advection is more difficult to assess because it is highly dependent on the frame of reference in which the term is calculated. It represents both the advection of the vortex core within the background monsoon circulation, and the advection of vorticity into and out of the vortex core. The former, which is of no interest to understanding the development of the vortex core, can be removed by subtracting the background flow. However, because it is not uniform it is impractical to subtract the background flow at every point of interest. Instead we investigated the effects of horizontal advection into and out of the vortex cores at only a few points (not shown). We found anticyclonic (cyclonic) tendencies in convergent (divergent) regions due to the inward (outward) flow in an environment of increasing cyclonic vorticity toward the center of the vortex core; that is, they op-

posed the tendencies from the convergence/stretching term. At low levels where the convergence was maximized this opposition was up to 25% of the convergence/stretching term.

Figure 3 shows the vortex is enhanced at low to midlevels by convergence/stretching of absolute vorticity in the deep intense updraft, and at mid- to upper levels by vertical advection of absolute vorticity. Thus vorticity is being concentrated and stretched from below and advected upward by the model updraft. Figure 3 also shows how the tendency terms evolve with the updraft life cycle. During the developing phase, when the updraft maximum is located at relatively low levels, the contribution from the stretching/convergence term is confined to low levels, and it grows deeper with time as the location of the updraft maximum moves upward with time (associated with the evacuation of mass from the horizontal plane where the upward flow is accelerating).⁴ Also evident in Fig. 3 is anticyclonic growth on the left edge of the updraft above about 7000 m (cf. the center and right panels, second row). This is due to tilting and appears to indicate that the net change in absolute vorticity is nearly zero at this level. This would be consistent with Haynes and McIntyre's (1987) demonstration that there can be no net change in absolute vorticity on a pressure surface. The PV associated with the same convective burst is presented in Fig. 7 of Part II, where we show in another cross section two hours later that the cyclonic anomaly does become considerably more intense at these levels. With time vortex advection effects lead to the ejection of the anticyclonic anomaly from the cyclonic core region.

To illustrate that this primary vortex enhancement mechanism is a very common feature of TC-LAPS simulations a similar analysis has been performed on another three updrafts from three additional TC forecasts. These are: the Elcho Island storm (Arafura Sea, January 2003, it reached TC intensity just prior to landfall and as a consequence was not named), TC Fiona (Arafura Sea, February 2003) and TC Erica (Coral Sea, southwest Pacific, March 2003). Part III contains a more detailed analysis of these events. The vertical velocity, absolute vorticity, and the two dominant abso-

⁴ The anticyclonic convergence tendency between 4000 and 8000 m in the decaying core A (Fig. 3, lower left panel) appears to contradict the combination of absolute vorticity and implied divergence from the panels above. Horizontal smoothing has merged two side-by-side anticyclonic tendency anomalies (not shown). These tendency anomalies are anticyclonic because on the left, the absolute vorticity is anticyclonic and the flow is horizontally convergent, and on the right, the flow is divergent and absolute vorticity cyclonic.

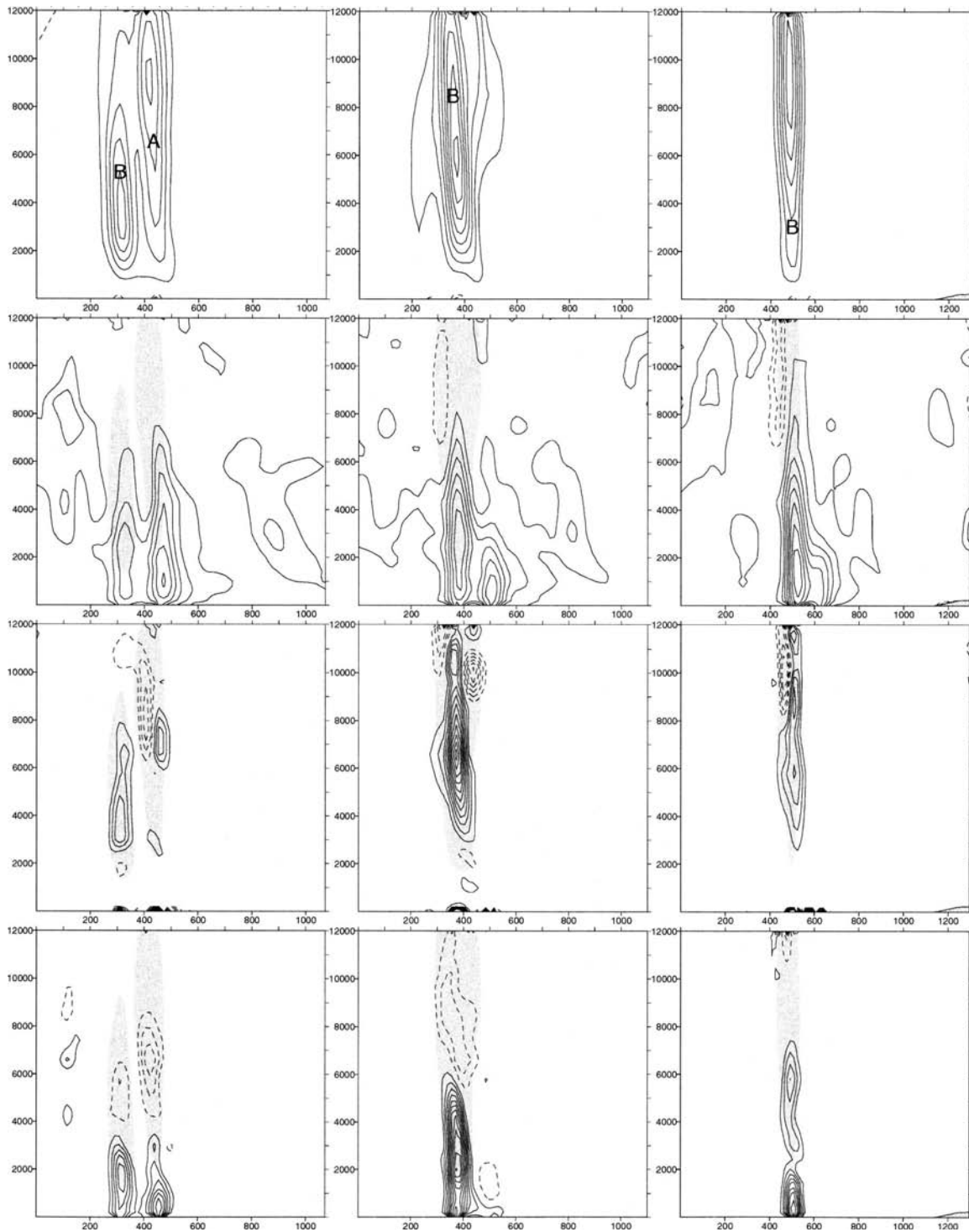


FIG. 3. Vertical cross sections of (top) vertical motion (contour interval = 1 Pa s^{-1}), (second row) absolute vorticity (contour interval = $1 \times 10^{-4} \text{ s}^{-1}$), (third row) contributions to the absolute vorticity tendency from vertical advection, and (bottom) stretching (contour interval = $2.5 \times 10^{-8} \text{ s}^{-2}$), at (left) 4, (center) 6, and (right) 8 hours into the simulation of TC Chris initialized at 1100 UTC 1 Feb 2002. All zero contours have been omitted. Dashed lines represent downward flow (top), anticyclonic vorticity (second row), and anticyclonic vorticity tendency (third and bottom rows). White lines in Fig. 2 indicate the cross-section locations, and convective updrafts/cores referenced in the text are labeled A and B. Shading has been added to the lower three rows to provide an indication of the updraft location (updraft intensity shaded, $\omega > 2.5 \text{ Pa s}^{-1}$). Horizontal and vertical axis scales are km and m, respectively.

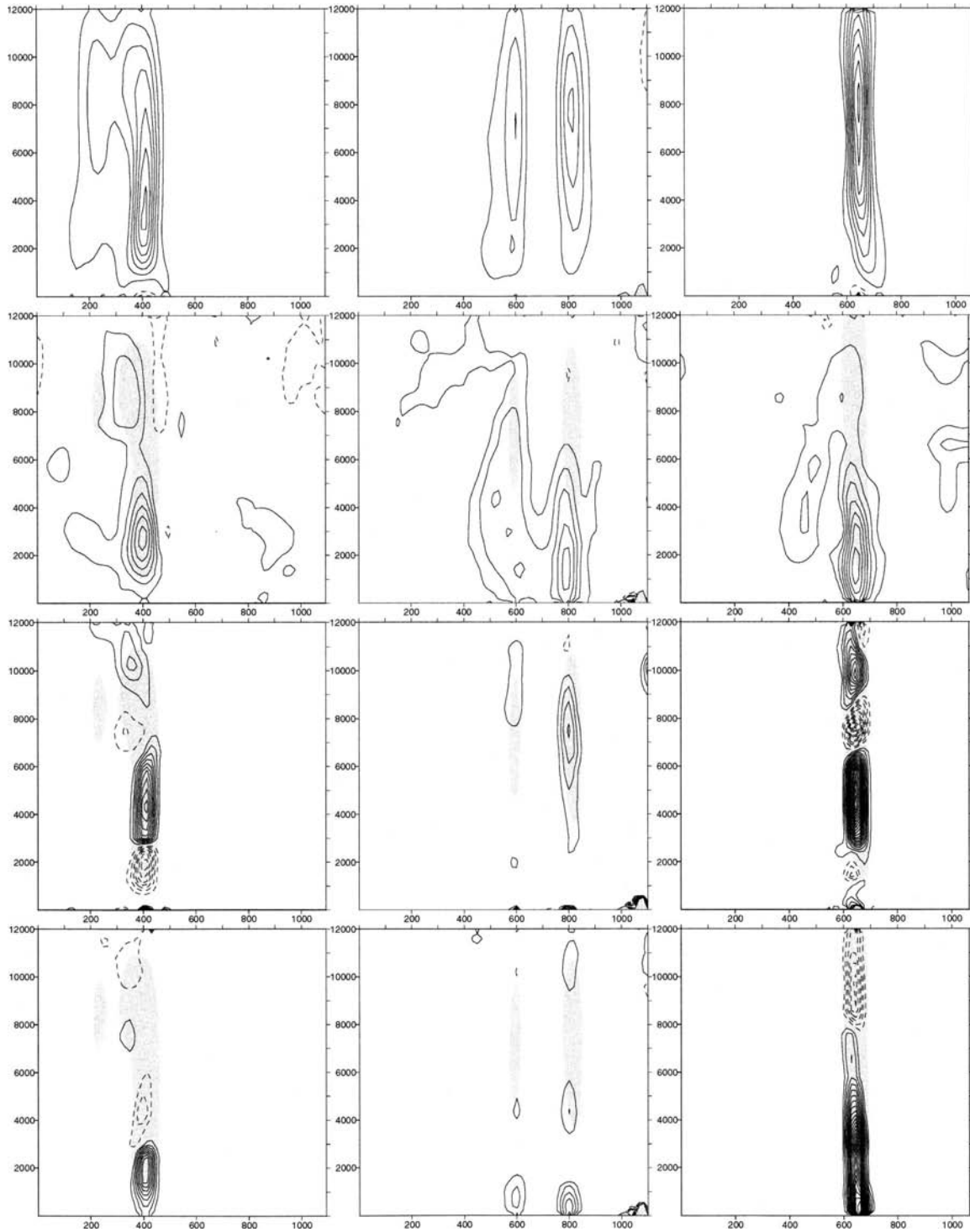


FIG. 4. As in Fig. 3 except for three additional TC simulations: (left) Elcho Island Storm, (center) TC Fiona, and (right) TC Erica. These simulations are presented in Part III.

lute vorticity tendency terms (vertical advection and convergence/stretching) for the three updrafts are presented in Fig. 4. This figure provides an indication of the variability of the process from case to case.

The size and intensity of the updrafts from the Elcho Island (left column, Fig. 4) and TC Erica simulations (right column) are examples of strong TC-LAPS updrafts, whereas the two updrafts evident in the TC

Fiona simulation (center column) are particularly weak. The Elcho Island “snapshot” was taken during a time of rapid development on both the updraft and system scales (5 h into the forecast). The rapid development on the updraft scale is evident in the two tendency terms in Fig. 4, which are of similar intensity to the mature updraft from TC Chris mentioned above (center column, Fig. 3). One obvious difference is the negative contribution to intensification from the vertical advection term below 2800 m, which results from the positive vertical gradient of cyclonic absolute vorticity in this layer (i.e., the absolute vorticity maximum is located at 2800 m). This low-level negative contribution is more than compensated for by the strong positive contribution from the stretching/convergence term, which leads, with time, to a downward migration of the maximum in cyclonic absolute vorticity. It is interesting to note that such a downward migration would give the appearance of the vortex growing down from midlevels, when clearly the dynamics describe intensification from below. This phenomenon has been observed in a number of TC-LAPS simulations, and is also evident in Chen and Frank (1993).

In Chen and Frank (1993) a 3D hydrostatic meso-scale model, the Pennsylvania State University–National Center for Atmospheric Research (PSU–NCAR), was used to simulate the growth of a midlatitude MCS and associated MCV. Between 4 and 8 h into the simulation a TC-LAPS-like updraft core developed and the subsequent vortex development was almost identical to that shown in Fig. 4 for the Elcho Island storm, except the vortex maximum was deeper, and as a consequence the tendency terms extended over a deeper layer of the lower troposphere. (Chen and Frank associated the vortex enhancement with stratiform dynamics, although the divergence profile appears to be convective.)

The TC Fiona simulation was characterized by very slow development on the system scale. This was due in part to relatively weak and short-lived updrafts, and relatively long periods between convective outbreaks that fuelled the genesis process. The weak contributions to vortex intensification on the updraft scale are evident in Fig. 4 (cf. with the other two examples). Clearly, as Eq. (1) would suggest, the stronger the updraft and the greater the absolute vorticity the updraft is embedded in, the greater the vortex intensification.

This point is very apparent in the case of TC Erica, where the two tendency terms are about double the magnitude of the Elcho Island and TC Chris terms. Despite this very significant vortex enhancement on the updraft scale, the TC Erica simulation failed to spin up a TC (see Part III). This particular updraft formed

about 100 km outside of a monsoon gyre, and about 300 km from the center of the gyre. It was also subjected to low-level vertical shear. The shear in itself may have been sufficient to inhibit development, however it is also possible that the cyclonic environment in the vicinity of the convection was not sufficiently large or intense to sustain the development (see Part III).

4. Discussion

In the previous section we have described a vortex intensification mechanism active in the TC-LAPS updrafts, and we note that it is the primary mechanism responsible for TC genesis in that it provides seed vortices and net heating that drive the secondary mechanisms of vortex upscale cascade and the SSI process (Parts II and III). This primary mechanism is essentially CVE.

A number of questions arise regarding the realism of the simulations.

- 1) Is the CVE process, as documented in the previous section, active in the real atmosphere?
- 2) If so, is the CVE process critical for TC genesis in the real atmosphere?
- 3) If so, are we adequately representing and resolving the process?

Given that two prominent genesis theories (Simpson et al. 1997; Ritchie and Holland 1997; Bister and Emanuel 1997) focus on SVE, which does not appear to play an obvious role in the TC-LAPS simulations, further questions arise.

- 4) Does the SVE process play an important role in TC genesis?
- 5) Does TC-LAPS adequately incorporate the SVE process?

We are unaware of observational evidence of sufficient temporal or spatial resolution to answer these questions with certainty. However, we believe the observations that do exist provide strong evidence in favor of the CVE process as a critical genesis mechanism. This is backed up by contemporary modeling studies (H04; M06). In presenting the argument for the CVE process being active and important for genesis in the real world we will address the five questions posed above.

a. Is the CVE process active in the real atmosphere?

Forecasters for decades have known that vortex intensification often follows periods of sustained deep convection in tropical systems ranging from depressions through to full-blown hurricanes. Zehr (1992) docu-

mented such observations focusing mostly on satellite imagery from two seasons (1983, 1984) of tropical storms and typhoons in the northwest Pacific (50 events in total). Low-level U.S. Air Force reconnaissance flight-level data (at approximately 500 m above sea level) were available for many of the systems, which provided confirmation of the role such convection played in intensifying low-level vorticity. He found that after periods of greatly enhanced deep cumulonimbus convection the low-level vorticity was significantly enhanced. We note that such observations do not rule out the genesis theories based on SVE, since it is conceivable that decaying convection contributed to the construction of a stratiform precipitation deck. However, we feel it is unlikely that SVE was responsible for the low-level vortex intensification reported by Zehr, given the response time required to build the stratiform precipitation deck (e.g., Houze 2004) and then somehow transfer the vorticity down from midlevels. Zehr suggests the intensification is evident in a matter of hours, which is consistent with direct low-level vortex intensification by CVE. Furthermore, in the TC genesis environment large regions of very cold cloud tops (significant overshooting) are often observed. Zehr showed one satellite image (Zehr's Fig. 7.5) of cold cloud-top temperatures covering nearly $3^\circ \times 3^\circ$ square region of clouds less than -65°C with an almost 1° diameter region less than -80°C embedded, suggesting a large contiguous region of deep intense convection was active. [We note in Part II that cloud-top temperatures less than -85°C correlate very well with convective rain rates obtained from Tropical Rainfall Measuring Mission (TRMM) data for TC Chris.]

MH95 provide evidence in support of CVE in tropical oceanic regions in nongeneration environments. They produced vertical profiles of mean horizontal divergence within 10 MCSs observed during TOGA COARE, and found that both SDP and CDP were present. On average neither was insignificant, which suggests vortex enhancement through convergence of absolute vorticity should be apparent from low to midlevels. They noted that some systems sampled were dominated by SDP and others CDP. Data collected from two of the systems dominated by CDP, included a TC rainband (TC Oliver), and a young vigorous, near-draft free, convective region.

Further evidence of mesoscale low- to midlevel convergence in large convective regions has been found in the Global Atmospheric Research Programme Atlantic Tropical Experiment (GATE) and the jointly conducted Equatorial Mesoscale Experiment (EMEX), and the Australian Monsoon Experiment (AMEX). Zipser and Gautier (1978) documented the intensifica-

tion of a tropical depression (TD) during GATE near Dakar on 15 July 1974. They noted, among other things, mesoscale cyclogenesis was preceded by mesoscale organization of deep convection, accompanied by strong mesoscale convergence at low levels; strong convective echoes were "within and downwind of a broad-zone of 10^{-4} s^{-1} convergence at 990 mb that covers an entire 1° square." Mapes and Houze (1992) commented that MCSs that developed in weak cyclonic depressions were characterized by very broad, deep bands of convective clouds, with large apparent upward mass flux. They illustrated this with an example from EMEX where an updraft exceeded 1 m s^{-1} continuously along a 40-km flight leg. Although the flight passed along a convective line, they noted that it was not a particularly narrow one. Finally, Reasor et al. (2005) provide evidence of low-level vortex enhancement associated with deep convection during the genesis of Hurricane Dolly.

b. Is the CVE process critical for TC genesis in the real atmosphere?

While we cannot answer this question with any certainty yet, we feel there is sufficient evidence to suggest it is quite likely. As noted above, (i) bursts of intense convection have long been associated with vortex intensification, and TC genesis; (ii) we suggested the CVE response time is likely to be more consistent with Zehr's observations than the indirect SVE; and (iii) we believe the relatively large areas of very cold cloud tops represent overshooting in deep convective regions where the mean divergence profile is likely to be consistent with MH's CDP. For these reasons we associate the relatively short-lived (order 6 h) but intense convective bursts that tend to occur sporadically throughout the genesis environment (e.g., Part II; Ritchie and Holland 1997; Ritchie et al. 2003; Harr et al. 1996), with CVE.

c. Are we adequately representing and resolving the CVE process?

To answer the question regarding the adequate representation of the mechanism we need to be sure that the scale and intensity of the TC-LAPS updrafts realistically represent the vertical profile of the mean horizontal divergence. The observations of Zipser and Gautier (1978) from GATE, the Mapes and Houze (1992) EMEX updraft observations, the MH observational study of convergence profiles taken during TOGA COARE, the large areas of very cold cloud-top temperatures observed by Zehr (1992), and finally the comments by Gray (1998) that pockets of extreme convection on the scale of 50 km "sometimes act as the focus from which the centers of tropical cyclones de-

velop," all suggest that relatively large areas (30 km or greater) of enhanced low- to midlevel convergence do exist in the oceanic tropical environment. A comment in Houze (1997, p. 2186) stated that deep convective precipitation areas in TOGA COARE typically reached 140 km in diameter. It follows that this large area, on average, would likely be dominated by a CDP. This latter scale is greater than the largest TC-LAPS updraft cores. The mean convergence in core B (Fig. 2 and 3) averaged over a diameter of 30 km yields values of $-4 \times 10^{-4} \text{ s}^{-1}$ above the boundary layer decreasing in magnitude to about $-2 \times 10^{-4} \text{ s}^{-1}$ near 3 km (not shown). This is about twice the intensity of the convective profiles measured by MH95, although it should be noted that they did not sample convective regions in an intensifying TC genesis environment. On the other hand, when averaged over a $1^\circ \times 1^\circ$ square, it is of similar intensity to that measured by Zipser and Gautier (1978) in an intensifying tropical depression. Essentially this means the vertical profiles of horizontal mean divergence observed in TC-LAPS are of similar type and spatial scale as those frequently observed in tropical oceanic environments, and of similar magnitude to one measurement in a TC genesis-like environment.

To answer the question regarding the adequate resolution we need to understand the importance of the finer details of the convection for TC genesis. Currently the TC-LAPS updrafts are consistent with an MH95 CDP. The associated vortex enhancement fuels the secondary intensification mechanisms that lead to TC genesis. It would appear from the many simulations we have performed using TC-LAPS (some reported in Parts II and III) that the TC-LAPS updrafts responsible for the primary vortex enhancement mechanism are qualitatively accurate to drive TC formation and nonformation. However, the finer details may be quantitatively important when it comes to accurately forecasting the location, timing and intensity of the developing TC.

d. Does the SVE process play an important role in TC genesis?

MH95 have shown the SDP is present in a number of nongensis oceanic tropical environments, and that it tends to dominate in long-lived MCSs. Vortex structures, consistent with the well-documented midlatitude terrestrial MCV that form in the stratiform precipitation deck, have been documented in the TC genesis environment (e.g., Harr and Elsberry 1996; Bister and Emanuel 1997; Raymond et al. 1998; Reasor et al. 2005). Clearly, SVE is favorable for TC genesis, in that it enhances the midlevel pool of vorticity, and it is likely

to favor convection (Raymond and Jiang 1990). However, the associated low-level divergence weakens the absolute vorticity magnitude at low levels, which is counter to the genesis process. Montgomery and Enagonio (1998) showed that the interaction of two vortices, one consistent with construction by SVE (traditional MCV) and the other consistent with construction by CVE (vortex core maximized at low levels), leads to the generation of a single upright vortex core. Thus SVE is likely to be important for TC genesis when it interacts with a vortex formed in a convective region. Therefore, we believe SVE may play a role in improving the likelihood of genesis, but ultimately genesis will not proceed without CVE.

e. Does TC-LAPS adequately incorporate SVE?

The MH95 vertical profiles of mean horizontal divergence have shown that MCSs can often contain SDP and CDP of similar intensity. It can be assumed in these cases the MCS being sampled consists of distinct convective and stratiform regions, with their typical respective divergence structures. However, averaged over the whole sampling region significant convergence extends from low to midlevels. This is evident in the TC-LAPS updrafts, particularly the mature updrafts (e.g., see the lower panels of Figs. 3 and 4). Whether this is just the nature of convective updrafts in the TC-LAPS TC genesis environment (i.e., maximum diabatic heating at 7 or 8 km, rather than 5 or 6 km), or whether TC-LAPS is in some way incorporating the SDP is yet to be determined. Regardless of the reason it is evidence that a divergence profile consistent with a combination of SDP and CDP, dominated by the latter, is being incorporated in the TC-LAPS updrafts.

5. Summary and conclusions

The search for an understanding of TC genesis over the last ten years or so has focused on the observation that TCs often develop in the vicinity of one or more MCSs. It was believed that the transition from MCS to a TC-like vortex required the generation of low-level cyclonic vorticity below the MCS, and the search for a TC genesis mechanism became the search for a mechanism that provided this sub-MCS low-level cyclonic vorticity. Two of the earlier genesis theories focused almost entirely on the vortex enhancement in the stratiform region of the MCS (Simpson et al. 1997; Bister and Emanuel 1997), and thus required mechanisms for bringing midlevel cyclonic vorticity down to the surface. Another genesis theory suggested the necessary low-level vorticity could be provided by vortex enhancement by low-level convergence into deep convective regions (Montgomery and Enagonio 1998). More

recent modeling studies have highlighted the importance of CVE in realistic (H04) and idealized environments (M06). CVE has been identified as the dominant vortex intensification mechanism in the TC-LAPS genesis simulations.

The vorticity tendency analysis of the TC-LAPS updrafts presented in section 3 showed that vortex cores were generated by convergence and vertical advection of vorticity in convective updrafts. We found this to be the TC-LAPS primary vortex enhancement mechanism. We show in Part II that the net heating generated in these cores drives the SSI process (H04; M06) and the vorticity generated in the cores contributes to the upscale vortex cascade (H04; M06), both secondary vortex enhancement mechanisms that lead to the construction of a cyclonic PV monolith.

The generation of deep vortex cores is fundamental to genesis in TC-LAPS, but it is not yet clear whether this mechanism is of equal importance in the real atmosphere. Evidence of deep convective regions of TC-LAPS scales do exist in tropical oceanic environments (e.g., Zipser and Gautier 1978; Mapes and Houze 1992; MH95), and more specifically are very commonly observed in genesis environments (Zehr 1992; Gray 1998). These latter observations are consistent with the hypothesis we propose that deep convective regions are critical for genesis in the real atmosphere. The TC-LAPS resolved updrafts ideally represent the mean flow in real convective regions. It may be that the interaction of these updrafts with the larger-scale environment is of sufficient accuracy to forecast formation with qualitative success, whereas greater resolution may be required to adequately represent smaller-scale features and interactions (and possibly separate CVE and SVE) before improvements in formation timing, rate, and intensity can be made.

SVE does not appear to play a significant role in the TC-LAPS simulations. Such midlevel vortex enhancement increases the likelihood of genesis, but as mentioned above genesis cannot proceed without some other mechanism to enhance low-level cyclonic vorticity (the SDP is divergent at low levels, thus can only act to reduce the low-level absolute vorticity magnitude). Of the two dominant divergence profiles observed by MH, only the CDP in a cyclonic environment can directly enhance the low-level vorticity. Thus SVE would not appear to be essential for genesis, but it may combine well with CVE to provide deep net convergent regions in MCSs, provided the CDP dominates at low levels. The vertical velocity profiles in the TC-LAPS updrafts are consistent with such deep net convergence, which suggests TC-LAPS may be adequately incorpo-

rating SVE at least during periods of strong convective activity.

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REFERENCES

- Bister, M., and K. A. Emanuel, 1997: The genesis of Hurricane Guillermo: TEXMEX analyses and a modeling study. *Mon. Wea. Rev.*, **125**, 2662–2682.
- Chen, S. S., and W. M. Frank, 1993: A numerical study of the genesis of extratropical convective mesovortices. Part I: Evolution and dynamics. *J. Atmos. Sci.*, **50**, 2401–2426.
- Davidson, N. E., and K. Puri, 1992: Tropical prediction using dynamical nudging, satellite defined convective heating and a cyclone bogus. *Mon. Wea. Rev.*, **120**, 2501–2522.
- ECMWF, 1995: ECMWF forecast model: Physical parameterization. 4th ed. European Centre for Medium-Range Weather Forecasts, Reading, United Kingdom, 142 pp.
- Eliassen, A., 1951: Slow thermally or frictionally controlled meridional circulation in a circular vortex. *Astrophys. Norv.*, **5**, 18–60.
- Emanuel, K. A., 1986: An air–sea interaction theory for tropical cyclones. Part I: Steady-state maintenance. *J. Atmos. Sci.*, **43**, 585–604.
- Fels, S. B., and M. D. Schwartzkopf, 1975: The simplified exchange approximation: A new method for radiative transfer calculations. *J. Atmos. Sci.*, **32**, 1475–1488.
- Fritsch, J. M., J. D. Murphy, and J. S. Kain, 1994: Warm core vortex amplification over land. *J. Atmos. Sci.*, **51**, 1780–1807.
- Gray, W. M., 1998: The formation of tropical cyclones. *Meteor. Atmos. Phys.*, **67**, 37–69.
- Harr, P. A., and R. L. Elsberry, 1996: Structure of a mesoscale convective system embedded in Typhoon Robyn during TCM-93. *Mon. Wea. Rev.*, **124**, 634–652.
- , —, and J. C. L. Chan, 1996: Transformation of a large monsoon depression to a tropical storm during TCM-93. *Mon. Wea. Rev.*, **124**, 2625–2643.
- Haynes, P. H., and M. E. McIntyre, 1987: On the evolution of vorticity and potential vorticity in the presence of diabatic heating and frictional or other forces. *J. Atmos. Sci.*, **44**, 828–841.
- Hendricks, E. A., M. T. Montgomery, and C. A. Davis, 2004: The

- role of "vortical" hot towers in the formation of Tropical Cyclone Diana (1984). *J. Atmos. Sci.*, **61**, 1209–1232.
- Holton, J. R., 2004: *An Introduction to Dynamic Meteorology*. 4th ed. Elsevier/Academic Press, 535 pp.
- Houze, R. A., 1997: Stratiform precipitation in regions of convection: A meteorological paradox? *Bull. Amer. Meteor. Soc.*, **78**, 2179–2196.
- , 2004: Mesoscale convective systems. *Rev. Geophys.*, **42**, RG4003, doi:10.1029/2004RG000150.
- Johnson, R. H. S., S. Chen, and J. J. Toth, 1989: Circulations associated with a mature-to-decaying midlatitude mesoscale convective system. Part I: Surface features—Heat bursts and mesoscale development. *Mon. Wea. Rev.*, **117**, 942–959.
- Kurihara, Y., and R. E. Tuleya, 1981: A numerical simulation study on the genesis of a tropical storm. *Mon. Wea. Rev.*, **109**, 1629–1653.
- Mapes, B. E., and R. A. Houze, 1992: An integrated view of the 1987 Australian monsoon and its mesoscale convective systems. I: Horizontal structure. *Quart. J. Roy. Meteor. Soc.*, **118**, 927–963.
- , and —, 1995: Diabatic divergent profiles in western Pacific mesoscale convective systems. *J. Atmos. Sci.*, **52**, 1807–1828.
- Miller, M. J., and R. P. Pearce, 1974: A three-dimensional primitive equation model of cumulonimbus convection. *Quart. J. Roy. Meteor. Soc.*, **100**, 133–154.
- Montgomery, M. T., and J. Enagonio, 1998: Tropical cyclogenesis via convectively forced vortex Rossby waves in a three-dimensional quasigeostrophic model. *J. Atmos. Sci.*, **55**, 3176–3207.
- , M. E. Nicholls, T. A. Cram, and A. B. Saunders, 2006: A vortical hot tower route to tropical cyclogenesis. *J. Atmos. Sci.*, **63**, 355–386.
- Puri, K., G. S. Dietachmayer, G. A. Mills, N. E. Davidson, R. A. Bowen, and L. W. Logan, 1998: The new BMRC Limited Area Prediction System, LAPS. *Aust. Meteor. Mag.*, **47**, 203–223.
- Raymond, D. J., and H. Jiang, 1990: A theory for long-lived mesoscale convective systems. *J. Atmos. Sci.*, **47**, 3067–3077.
- , C. Lopez-Carillo, and L. Lopez Cavazos, 1998: Case-studies of developing east Pacific easterly waves. *Quart. J. Roy. Meteor. Soc.*, **124**, 2005–2034.
- Reasor, P. D., M. T. Montgomery, and L. F. Bosart, 2005: Mesoscale observations of the genesis of Hurricane Dolly (1996). *J. Atmos. Sci.*, **62**, 3151–3171.
- Ritchie, E. A., and G. J. Holland, 1997: Scale interactions during the formation of Typhoon Irving. *Mon. Wea. Rev.*, **125**, 1377–1396.
- , J. Simpson, W. T. Liu, J. Halverson, C. S. Velden, K. F. Brueske, and H. Pierce, 2003: Present day satellite technology for hurricane research. *Hurricane: Coping with Disaster*, R. Simpson, Ed., Amer. Geophys. Union, 249–289.
- Rotunno, R., and K. A. Emanuel, 1987: An air–sea interaction theory for tropical cyclones. Part II: Evolutionary study using a nonhydrostatic axisymmetric numerical model. *J. Atmos. Sci.*, **44**, 542–561.
- Saunders, A., and M. T. Montgomery, 2004: A closer look at vortical hot towers within a tropical cyclogenesis environment. Department of Atmospheric Science Blue Book, Colorado State University, Vol. 752, 105 pp.
- Simpson, J., E. A. Ritchie, G. J. Holland, J. Halverson, and S. Stewart, 1997: Mesoscale interactions in tropical cyclone genesis. *Mon. Wea. Rev.*, **125**, 2643–2661.
- Tiedke, M., 1989: A comprehensive mass flux scheme for cumulus parameterization large-scale models. *Mon. Wea. Rev.*, **117**, 1779–1800.
- Tory, K. J., M. T. Montgomery, N. E. Davidson, and J. D. Kepert, 2006: Prediction and diagnosis of tropical cyclone formation in an NWP system. Part II: A diagnosis of Tropical Cyclone Chris formation. *J. Atmos. Sci.*, **63**, 3091–3113.
- Troen, I. B., and L. Mahrt, 1986: A simple model of the atmospheric boundary layer: sensitivity to surface evaporation. *Bound.-Layer Meteor.*, **37**, 129–148.
- Zehr, R., 1992: Tropical cyclogenesis in the western North Pacific. NOAA Tech. Rep. NESDIS 61, 181 pp.
- Zipser, E. J., and C. Gautier, 1978: Mesoscale events within a GATE tropical depression. *Mon. Wea. Rev.*, **106**, 789–805.