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THESIS

DATA REDUCTION FOR THE UNSTEADY AERODYNAMICS ON A CIRCULATION CONTROL AIRFOIL

by

Billy Murel Pickelsimer

March 1977

Thesis Advisor: Louis V. Schmidt

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Data Reduction for the Unsteady Aerodynamics on a Circulation Control Airfoil

by

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Submitted in partial fulfillment of the requirements for the degree of

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ABSTRACT

Calculating the lift, drag, and pitching moment coefficients for an airfoil from the static pressure distribution obtained from wind tunnel tests is routine task when steady flow is considered, but it is much more complicated when the airfoil is operating in an unsteady flow field, similar to that experienced by a helicopter rotor blade, produced by an oscillating wind tunnel. A data reduction routine capable of condensing the large numbers of data associated with the unsteady investigation, as well as a numerical integration algorithm for the unsteady aerodynamic coefficients, were developed; however, no unsteady data were collected due to hardware failures. The ability of the program was demonstrated on previously obtained steady and quasi-steady data and sample results were presented.

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I. INTRODUCTION

The primary objectives of this investigation were: 1) to report on the data reduction techniques used to obtain values for both steady and non-steady aerodynamic coefficients, resulting from the experimentally acquired static pressure distribution over the model when the tunnel was operating at a specified dynamic pressure, with Coanda blowing oscillated at various frequencies and amplitudes of momentum blowing coefficient; and 2) to present a sample of results. Removal of the model from the tunnel for modifications in November 1976 precluded any actual processing of non-steady data consistent with the stated goals.¹

However, the data reduction techniques are applicable to the pressure distribution obtained with steady Coanda blowing and will be demonstrated on such data obtained in April 1976. Prior to removal of the model, the tunnel was operated without blowing in an effort to generate a set of non-steady data to be used in tailoring the data reduction procedures to be used downstream in the on-going project. In that test, the model experienced self-excited harmonic pressures resulting from vortices being shed alternately

The modifications, to correct for surface delaminations near the upper jet nozzle, include a steel upper surface trailing edge and slot height adjustment to 0.015 inches. Additional pressure taps were installed on the Additional pressure taps were installed on the upper surface forward of the Coanda slot, and leaking tubes were repaired while the modifications were being implemented.

from the upper and lower surfaces. The manually recorded data consisted of the RMS and mean values of the unsteadypressure, and the phase angle in respect to a reference for each pressure orifice. A more complete description of the data acquisition system is given in Section II B.

 $\mathcal{L}^{\text{max}}_{\text{max}}$

II. EQUIPMENT AND INSTRUMENTATION

A. AIRFOIL

The airfoil used in this study was obtained from the Lockheed-California Company after completion of their design feasibility study on circulation control rotors (Ref.l). It had an elliptical cross-section with a chord of 10.215 inches, thickness ratio of 0.20, and camber of 3.3% (Figure 1). A slot to facilitate Coanda blowing was machined in the upper surface at chord position X/C=0.95, with the slot height set at 0.010 inches. Coanda air was supplied to the plenum cavity at both ends of the section to insure uniform flow along the span. Fifty-three pressure taps at mid-span with a chord-wise orientation were monitored through two scanivalve transducers for an electrical representation of the pressure distribution over the airfoil. The pressure tap locations are listed in Table I.

The 2.0 x 2.0 foot unsteady flow wind tunnel in the Department of Aeronautics Laboratories at the Naval Postgraduate School was used as it is capable of simulating the velocity flow field experienced by blade elements of a helicopter in forward flight. The model was mounted at an angle of attack of -5 degrees, which closely corresponds to the zero-lift angle of attack for the data reported herein.

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RIRFOIL CROSS-SECTION AITH THP LOCATIONS

| THBLE | |
|-------|--|
|-------|--|

AIRFOIL PRESSURE THP LOCATIONS

X IN INCHES HFT OF LEADING EDGE

Y IN INCHES FROM CHORD LINE

B. INSTRUMENTATION

The data acquisition system evolved from a classic manometer board arrangement with readings hand recorded (Ref. 2) and reduced to the present system, which employed electrical read-outs with an existing Digitec system with customized signal conditioning. The digital multiplexing circuitry in the Digitec system was interconnected through the scanivalve solenoid controllers for automatic scanivalve advancing as the output of the digital voltmeter was sequentially printed on a paper tape. Pressure sensing was accomplished by connecting each of the pressure taps to one of the two scanivalves (Figure 2) through tubing of uniform diameter and length. This was done to insure a constant transfer function for all pressure taps, when measuring unsteady pressures. Calibration tests using the method reported by Johnson (Ref. 3) showed an acceptable variation in dynamic gain in the frequency range from 0 to 300 Hz (Figure 3) , which included the range of practical interest.

For the unsteady aerodynamics analysis $(C_{ij} = 0)$, the voltage representation of the surface static pressures was of the form

$$
e_x = a_x + b_x(t). \tag{2.1}
$$

Full analog measuring techniques were employed in measuring the mean voltage, RMS voltage, and phase relationship to a reference for each pressure tap. The mean voltage was obtained by processing the signal through a low-pass filter circuit (Figure 4) with a two second time constant which

14

Figure 3

corresponds to a corner frequency of 0.08 Hz on a Bode plot. Therefore, the influence of the unsteady voltage was essentially negated and the output readings represented the steady components of the signal.

C. ADVANCED DATA ACQUISITION SYSTEM

To meet the primary objective mentioned in Section I, a more sophisticated data acquisition and reduction system will be required and is now being constructed by Englehardt² at the Naval Postgraduate School in the Department of Aeronautics. The system, based on the Intel MDS-80 microprocessor unit, will control the logic of the experiment, collect and store data, and handle the entire data reduction process. The algorithm for data reduction discussed in Section V was proven sufficiently simple and capable of performing all required data reduction and numerical integration.

 2 Master of Science thesis by Englehardt with a complete 2 description of the entire microprocessor system will be published in June 1977.

 $\ddot{}$

Figure 4 Low-Pass Filter Circuit

III. DATA REDUCTION

A. STEADY DATA PROCESSING

The aerodynamic coefficients are defined by integral equations but must be approximated by numerical integrations, since data are available only at a finite number of pressure tap locations. The approximate steady normal force, chord force, and pitching moment coefficients are

$$
C_{n} = \int_{0}^{1.0} (C_{p_{1}} - C_{p_{u}}) d({X}_{/C})
$$

\n
$$
C_{c} = \int_{Y/C \text{ (max)}}^{Y/C \text{ (max)}} (C_{p_{f}} - C_{p_{v}}) d({Y}_{/C})
$$

\n
$$
C_{m(TE)} = \int_{-Y/C \text{ (min)}}^{Y/C \text{ (max)}} (C_{p_{f}} - C_{p_{v}}) (Y_{/C}) d({Y}_{/C}) +
$$

\n
$$
\int_{0}^{1.0} (C_{p_{1}} - C_{p_{u}}) (1 - X_{/C}) d({X}_{/C})
$$

Then, with an angle of attack correction for the force coefficients and a moment transfer to the quarter-chord position, the aerodynamic coefficients are

$$
C_{1} = C_{n} \cos \alpha - C_{c} \sin \alpha
$$

\n
$$
C_{d} = C_{n} \sin \alpha + C_{c} \cos \alpha
$$
 (3.2)
\n
$$
C_{m} (c/4) = C_{m} (TE) - 0.75 C_{n}.
$$

The pressure data from the April 1976 tests and the November 1976 mean pressure data were transferred to a

Hewlett-Packard 9830 tape cassette and processed using the program listed in Appendix A, which employs the equivalent numerical relationship to the equations stated above.

B. UNSTEADY DATA PROCESSING

Before the unsteady data were suitable for numerical integration, a dynamic gain correction was necessary to account for the transfer function discussed in Section II. The frequency of the self-induced oscillations was constant at 262 Hz when the tunnel airspeed was approximately 100 ft/s. From Figure 3, the corresponding gain correction, $|G_{\text{f}}|$, was approximately 0.5, which implied that the amplitude of the surface static pressure was attenuated by a factor of one-half at the transducer. All phase angle values were relative to tube number 7, which was arbitrarily chosen as a clock reference. Phase angle errors resulting from the two different scanivalve circuits were faired out in the analysis (Figure 5).

With the above corrections, the numerical integration of the unsteady static pressure distribution was carried out in a manner similar to the steady pressure integrations. For the unsteady normal force coefficient

$$
\langle c_n^2 \rangle^{\frac{1}{2}} \sin \phi = \int_0^{1.0} (c_{p_1} \sin \phi - c_{p_u} \sin \phi) d\left(\frac{x}{c}\right)
$$

$$
\langle c_n^2 \rangle^{\frac{1}{2}} \cos \phi = \int_0^{1.0} (c_{p_1} \cos \phi - c_{p_u} \cos \phi) d\left(\frac{x}{c}\right)
$$
 (3.3)

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Then
$$
\langle C_n^2 \rangle^{\frac{1}{2}} = \sqrt{(\langle C_n^2 \rangle^{\frac{1}{2}} \sin \phi)^2 + (\langle C_n^2 \rangle^{\frac{1}{2}} \cos \phi)^2}
$$

and $\phi_T = \tan^{-1} \frac{(\langle C_n^2 \rangle^{\frac{1}{2}} \sin \phi)}{(\langle C_n^2 \rangle^{\frac{1}{2}} \cos \phi)}$

where $\langle c_n^2 \rangle^2$ is the RMS value of the normal force coefficient and $\begin{array}{c} \Phi_{\rm T} \end{array}$ is the phase angle relative to tube number 7. The chord force and pitching moment coefficients were calculated similarly.

$$
\langle C_{c}^{2} \rangle^{\frac{1}{2}} \sin \phi = \int_{-Y/C(max)}^{Y/C(max)} (C_{p_{f}} \sin \phi - C_{p_{r}} \sin \phi) d(Y_{c})
$$

$$
\langle C_{c}^{2} \rangle^{\frac{1}{2}} \cos \phi = \int_{-Y/C(min)}^{Y/C(max)} (C_{p_{f}} \cos \phi - C_{p_{r}} \cos \phi) d(Y_{c})
$$

$$
\langle C_{c}^{2} \rangle^{\frac{1}{2}} = \sqrt{(\langle C_{c}^{2} \rangle^{\frac{1}{2}} \sin \phi)^{2} + (\langle C_{c}^{2} \rangle^{\frac{1}{2}} \cos \phi)^{2}}, \phi_{T} = \frac{(\langle C_{c}^{2} \rangle^{\frac{1}{2}} \sin \phi)}{(\langle C_{c}^{2} \rangle \cos \phi)}
$$

$$
\langle c_m^2 \rangle_{\text{te}}^{\frac{1}{2}} \sin \phi = \int_{Y/C(\text{min})}^{Y/C(\text{max})} (c_{p_f} \sin \phi - c_{p_r} \sin \phi) \left(\frac{Y}{C} \right) d \left(\frac{Y}{C} \right) +
$$

$$
\int_{0}^{1.0} (C_{p_1} \sin \phi - C_{p_u} \sin \phi) (1 - \frac{x}{c}) d\left(\frac{x}{c}\right)
$$

$$
\langle C_m^2 \rangle_{te}^{\frac{1}{2}} \cos \phi = \int_{-Y/C(min)}^{Y/C(max)} (C_{p_f} \cos \phi - C_{p_r} \cos \phi) \left(\frac{Y}{C} \right) d \left(\frac{Y}{C} \right) +
$$

$$
\int_{0}^{1} (C_{p_1} \cos \phi - C_{p_u} \cos \phi) (1 - \frac{x}{c}) d\left(\frac{x}{c}\right)
$$

$$
\langle C_m^2 \rangle_{te}^{\frac{1}{2}} = \sqrt{(\langle C_m^2 \rangle^{\frac{1}{2}} \sin \phi)^2 + (\langle C_m^2 \rangle^{\frac{1}{2}} \cos \phi)^2}, \phi_T = \frac{(\langle C_m^2 \rangle^{\frac{1}{2}} \sin \phi)}{(\langle C_m^2 \rangle^{\frac{1}{2}} \cos \phi)}
$$

The unsteady aerodynamic coefficients are calculated as in Eqs . 3.2, with the aid of elementary trigonometric identities to account for the phase angles. The program listed in Appendix A has the option for unsteady numerical integration incorporated.

It is intended that the $C_\mu^{}$ =0 $\,$ oscillating pressure $\,$ measurements be repeated when the improved data acquisition system becomes available. The tests will be conducted at several tunnel dynamic pressures to determine the Reynolds number dependence of the data. The previous measurements were at approximately $R_e = 0.54 \times 10^6$.

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IV. RESULTS

A. STEADY DATA RESULTS

The validity of the concept of circulation control was demonstrated by Englar (Refs. 4 and 5) and is confirmed here with the steady momentum blowing coefficient results. The resulting aerodynamic coefficients are summarized in Table II. The ability to produce relatively large increases in lift coefficient with only a small increase in blowing coefficient, shown graphically in Figures 6a, b, and c, is one of the desirable characteristics of the system. Probably the best measure of the effectiveness is the aerodynamic lift amplification ratio, $3C_1/3C_1$, which was computed to be approximately 15.3, the low value being attributed to the geometric deficiencies of the trailing edge slot as alluded to in Section I. However, the results are encouraging.

Close scrutiny of the coefficient of pressure information plotted against X/C (Figures 7-12) reveals a pressure distribution over the model such that the center of pressure was at the mid-chord position, which is in agreement with the theory. The center of pressure position is also verified by the graph of $C_{m(c/4)}$ Vs. C_{u} .

ed by the graph of
$$
C_{m(c/4)}
$$
 Vs. C_{μ} .
\n
$$
(\frac{x}{C})_{C.P.} = \frac{C}{4} - (\frac{3C_{m}/3C_{\mu}}{(3C_{1}/3C_{\mu})}) = .25 - \frac{(-3.5)}{(15.5)} = .48
$$

 \mathcal{I}

TABLE. II

AERODYNAMIC COEFFICIENT SUMMARY

STEADY AERODYNAMIC COEFFICIENTS

Angle of Attack=-5 degrees U=100 ft/s

UNSTEADY AERODYNAMIC COEFFICIENTS

 $C_1 = -0.0765$ $C_1(RMS) = 2.1404$ $\phi = 225$ degrees C_{d} = 0.0451 C_{d} (RMS)=0.1629 ϕ = 39 degrees C_m =-0.1063 C_m (RMS)=0.5817 ϕ = 63 degrees

Figure 7

 $\bar{\bar{a}}$

 \bar{b}

Also, the low drag qualities of the airfoil are favorable when full thrust recovery of the momentum blowing coefficient is assumed. Pressure coefficient data are listed in Appendix B.

B. UNSTEADY DATA RESULTS

The unsteady aerodynamics of the elliptically shaped airfoil have been given a rather cavalier treatment in the past but their contribution to the total forces and moments must be realized. The implications of the unsteady aerodynamics is not generally well understood; therefore, a simple example is presented here to emphasize their importance before discussion of the unsteady results.

Consider an airfoil at zero angle of attack with the RMS pressure distribution equal on the upper and lower surfaces (Figure 13a) but with unequal relative phase angles (Figure 13b) taken from the quarter-chord and the three-quarter-chord positions from the actual unsteady data (Figure 5). The unsteady normal force coefficient, $\langle C_n^2 \rangle^2$, is calculated from the numerical approximation to Eq. 3.3, as :

$$
\langle C_n^2 \rangle^{\frac{1}{2}} \sin \phi = 0.5 (\sin(210^\circ) + \sin(230^\circ) - \sin(35^\circ) - \sin(55^\circ))
$$

$$
\langle C_n^2 \rangle^{\frac{1}{2}} \sin \phi = -1.33
$$

$$
\langle C_n^2 \rangle^{\frac{1}{2}} \cos \phi = 0.5 (\cos(210^\circ) + \cos(230^\circ) - \cos(35^\circ) - \cos(55^\circ))
$$

$$
\langle C_n^2 \rangle^{\frac{1}{2}} \cos \phi = -1.45
$$

$$
\langle C_n^2 \rangle^{\frac{1}{2}} = \sqrt{1.33^2 + 1.45^2} \qquad \phi = \frac{-1.33}{-1.45}
$$

$$
\langle C_n^2 \rangle^{\frac{1}{2}} = 1.97 \qquad \phi = 223^\circ
$$

Figure 13b Sample Pressure Distribution

Similarly, for the moment coefficient (referenced to the trailing edge) $\left\langle C_m^2 \right\rangle^2 \sin \phi = 0.5((\sin(210^{\circ}) - \sin(35^{\circ}))0.75 + (\sin(230^{\circ}) - \sin(55^{\circ}))0.25)$ $\langle C_m^2 \rangle^2 \sin \phi = -0.6$ $\langle C_{\rm m}^2 \rangle^2$ cos = 0.5((cos(210^o) - cos(35^o))0.75+(cos(230^o) - cos(55^o))0.25) *s $\langle C_m \rangle$ $2\sqrt{cos} = -0.78$ $\frac{1}{2}$ $\langle C_m^2 \rangle^{\frac{1}{2}} = \sqrt{0.6^2 + 0.78^2}$ $\phi = \frac{-0.6}{-0.78}$ $\left\langle C_{\infty} \right\rangle^{1/2} = 0.98$ $\phi = 217^\circ$

Clearly, the force and moment coefficients are of significant magnitude even if a uniform distribution (RMS) of pressure coefficient were assumed, because of the importance of considering relative phase relationships. The dynamic behavior of the unsteady aerodynamics of the circulation control airfoil, particularly the lift amplification ratio, in the unsteady environment caused by cyclic blowing will be investigated after the improved data acquisition system is operational.

As discussed in Section I, measurements were taken in November 1976, with the airfoil experiencing harmonic pressure variations. The unsteady data are listed in Appendix C. The experimental data were numerically integrated, similar to the previous example, for the lift, drag, and pitching moment coefficients (Figures 14 and 15). It should be recognized that the phase shift between the upper

35

and lower surfaces is approximately 180 degrees, the result being the relatively large amplitude of the aerodynamic coefficients. Highlights of the unsteady results, summarized in Table II, were that:

1. The time varying pressures over the model were predominantly harmonic, with a non-dimensional frequency (Strouhal no.) of

$$
S = \frac{f}{U} = 0.449.
$$

2. The RMS value of the normal force coefficient was approximately 2.14, which was an order of magnitude greater than the RMS drag coefficient.

3. The test Reynolds no. was approximately

$$
R_e = 0.54 \times 10^6.
$$

4. The location on the airfoil for the center of pressure due to unsteady loadings was approximately 0.491 chord and the RMS value of pitching moment at this reference point was

$$
\left\langle C_m^2 \right\rangle^{\frac{1}{2}} = 0.1694
$$

with a phase angle of 90 degrees lag relative to the normal force. This term may be considered as an apparent mass type effect.

CP(RM5)*5IN(PHI) V5. X/C

V. DATA REDUCTION FOR THE ADVANCED DATA REDUCTION SYSTEM

The advanced data acquisition system described in Section II C will allow a more comprehensive investigation of the unsteady aerodynamics of the model, but the data reduction techniques must be tailored to handle the large numbers of data to be collected. Oscillation of the plenum pressure as described by Bauman (Reference 6) induces harmonic blowing from the trailing edge slot at a specified frequency, W_1 . It is assumed that the unsteady static pressure at each station will have a steady value plus a periodic value with frequency W_{1} superimposed due to the oscillating momentum blowing coefficient.

Since the data are assumed periodic in form, the unsteady pressure at any point on the airfoil may be defined by a Fourier series (Ref. 7) as follows:

 $P(x,t) = A_0(x) + \sum_{n=1}^{\infty} [A_n(x) \cos(nW_1 t) + B_n(x) \sin(nW_1 t)]$ (5.1) where W_1 is the frequency of the oscillation. The standard formulae for obtaining the Fourier coefficients for the infinite series are:

$$
A_0(x) = \frac{1}{2T} \int_{-T}^{T} P(x) dx
$$

\n
$$
A_n(x) = \frac{1}{T} \int_{-T}^{T} P(x) \cos(nx) dx
$$
 (5.2)
\n
$$
B_n(x) = \frac{1}{T} \int_{-T}^{T} P(x) \sin(nx) dx
$$

Because the experimental data are in the form of a discritization of a truncated periodic signal, the above equations are slightly altered for estimating the Fourier coefficients. The modified form is:

$$
\hat{X}_{o}(x) = \frac{1}{K} \sum_{i=1}^{K} P_{i}(x)
$$

$$
\hat{A}_{n}(x) = \frac{2}{K} \sum_{i=1}^{K} P_{i}(x) \cos(nW_{1}t) ; n=1,...,n \qquad (5.3)
$$

$$
\hat{B}_{n}(x) = \frac{2}{K} \sum_{i=1}^{K} P_{i}(x) \sin(nW_{1}t)
$$

where k is the number of samples to be analyzed and n is the harmonic number. It is recognized that the physical meaning of $A_0(x)$ is the steady or mean value of the oscillating pressure. That is, $A_0(x) = \langle p_0(x) \rangle$ where

$$
\left\langle P_{\mathbf{O}}(x)\right\rangle \equiv \frac{\text{Lim}}{\text{T} + \infty} \quad \frac{1}{2\text{T}} \int_{-\text{T}}^{\text{T}} P(x) \, \mathrm{d}x \tag{5.4}
$$

By subtracting out the mean pressure value, the signal becomes biased to a pure oscillating pressure with a mean value equal to zero. The terms $\tilde{A}_n(x)$ and $\tilde{B}_n(x)$ are equivalent to the cosine and sine components, respectively, discussed in Section III B. Figure 16 illustrates the reconstruction of a discritized signal utilizing this algorithm programmed on the Hewlett-Packard 9830 calculator (program listing in Appendix A). The error was minimized by insuring that the number of samples analyzed was an integer multiple of the frequency; however, the truncation error was never greater than 2%.

UNETERDY PRESENEE

Having developed an effective method for handling the large numbers of data, calculation of the aerodynamic force and moment coefficients can be accomplished. Before the actual processing of data begins, the sample pressures are converted to non-dimensional coefficient form to expedite calculations. The unsteady pressure coefficients are

$$
C_{p_i}(t) = \frac{p_i - p}{p - H} \qquad i = 1, 2, 3, ..., 53 \qquad (5.5)
$$

where P and H are steady pressure values from a pitotstatic tube mounted on the floor of the wind tunnel. Then the oscillating signal at any pressure tap is;

$$
C_p(x, t) = \sum A_n \cos(nw_1 t) + B_n \sin(nw_1 t)
$$

 \mathbf{f} n \mathbf{f} \mathbf{h} \mathbf{f} \mathbf{n} \mathbf{f}

$$
C_p(x, t) = \sum T_n \sin(nw_1 t + \phi_n)
$$
 (5.6)

$$
T_n = \sqrt{A_n^2 + B_n^2}, \phi_n = B_{n/2}
$$

The unsteady $C_p(x,t)$ are related by the phase angles ϕ , however, the phase angles are meaningless unless they are measured relative to a "clock" or reference signal. Therefore, it is necessary to designate one pressure tap to be used as the reference. Then, the relative phase angle is

$$
\phi_{\text{rel}} = \phi_{\text{i}} - \phi_{\text{ref}} \tag{5.7}
$$

For simplicity, ϕ_i shall hereafter be interpreted as the relative phase angle.

With the steady pressure coefficient removed from the oscillating signal, separate aerodynamic coefficients can be calculated for the steady and unsteady signals. The

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steady force and moment coefficients are found exactly as in Section III A. The unsteady aerodynamic coefficients can be defined in a Fourier series form as:

$$
C_n(t) = \sum C_{n_k} \sin(kw_1 t + \phi_k)
$$

\n
$$
C_c(t) = \sum C_{c_k} \sin(kw_1 t + \phi_k)
$$
 (5.8)
\n
$$
C_m(t) = \sum C_{m_k} \sin(kw_1 t + \phi_k)
$$

The magnitudes and phase angles of the coefficients can be determined thusly:

for the first harmonic

$$
C_{n_1} \sin(w_1 t + \phi) = C_{n_1} \sin\phi \cos(w_1 t) + C_{n_1} \cos\phi \sin(w_1 t)
$$

where

$$
C_{n_1} \sin \phi = \int_0^1 \left[C_{p_1} \left(\frac{x}{c} \right) \sin \phi_1 \left(\frac{x}{c} \right) - C_{p_1} \left(\frac{x}{c} \right) \sin \phi_1 \left(\frac{x}{c} \right) \right] d \left(\frac{x}{c} \right)
$$

$$
C_{n_1} \cos \phi = \int_0^1 C_{p_1} (\frac{x}{c}) \cos \phi_1 (\frac{x}{c}) - C_{p_1} (\frac{x}{c}) \cos \phi_1 (\frac{x}{c}) d(\frac{x}{c})
$$
 (5.9)

then

$$
C_{n_1} = \sqrt{(C_{n_1} \cos \phi_1)^2 + (C_{n_1} \sin \phi_1)^2}, \quad \phi_1 = \tan^{-1} \left(\frac{C_{n_1} \sin \phi_1}{C_{n_1} \cos \phi_1} \right)
$$

Similar calculations are done to obtain values for the higher harmonics $(c_{n_2}, c_{n_3}, \cdots, c_{n_i})$ as well as the chord force and pitching moment coefficients using the higher harmonics of pressure coefficient information- as obtained during the initial data reduction previously described.

VI. CONCLUSIONS

As mentioned in Section I, the two-fold objective of this paper was to develop a computer program for unsteady aerodynamic analysis and to demonstrate its capabilities on experimentally generated unsteady data. The objective was partially satisfied; a computer program, well within the limitations of mini-calculators such as the Hewlett-Packard 9830, was developed which employed a straightforward numerical integration of pressure coefficients to obtain values for the aerodynamic coefficients. The program can be modified for unsteady data by properly taking into account the cosine and sine components of the pressure signal relative to a fixed harmonic (clock) signal. Also, the extraction of the Fourier coefficients of the truncated, time-varying signal appears to be within the capabilities of the updated data acquisition system.

APPENDIX A

COMPUTER PROGRAM LISTINGS AND OUTPUT

```
1 REM PROGRAM FOR DECOMPOSITION OF A FUNCTION
  2 REM INTO FOURIER COEFFICIENTS
 10 DIM EC2561
 30 W1=(20*PI)
 35 DISP "DIVISIONS PER CYCLE?";
 36 INPUT D
 42 P=2*PI/W1
 43 T=P/D
 50 REM COMPUTE SAMPLE FUNCTION
 60 FOR 1=1 TO 220
 65 X=W1*(I-1)*T70 E[I]=SIN(X)+0.5*SIN(2*X)+0.3*SIN(3*X)80 NEXT I 100 PRINT
105 PRINT "THE DRIVING FREQUENCY IS 10 HZ"
106 PRINT D; "DIVISIONS PER CYCLE"
107 PRINT
110 PRINT " COEFFICIENTS"
111 PRINT
113 PRINT TAB2"COS"TAB15"SIN"
114 FOR J=l TO 3
115 PRINT
117 PRINT TAB2"A";J, "B";J, "PHASE ANGLE"
119 A1=B1=0
120 FOR 1=1 TO 220
130 Bl=(B1+E[I] *SIN(J*W1*(I-1)*T))
140 Al=(A1+ELI^*COS(J^*W1^*(I-1)^*T))150 NEXT I 159 P=ATN(B1/A1)
160 WRITE (15,170) (Al*2/220) , (Bl*2/220) , (P*180/PI)
170 FORMAT F6.3,7X, F6.3, 13X, F5.1
175 NEXT J
180 END
            THE DRIVING FREQUENCY IS 10 HZ
                 10 DIVISIONS PER CYCLE
             COEFFICIENTS
            COS SIN
             A 1 B 1 PHASE ANGLE
            0.000 1.000 90.0
```



```
1 REM COMPUTE CL , CD, CM- -STEADY AND UNSTEADY
  2 REM FROM CP DATA ON TAPE CASSETTE
 10 COM R1,DC80J ,P[801 ,C£60]
 20 DIM AC1061
 30 PRINT
 40 LOAD DATA 4,
 60 C9=0
 70 X9=0.95
 80 Y9=0.048
 90 DISP "CP DATA FILE NO. IS";
100 INPUT Z 110 LOAD DATA Z 111 DC22 = 0.4112 DC36\overline{J} = 0.539113 Gl=l
115 PRINT "INPUT OPTION"
116 PRINT "1 FOR STEADY DATA"
117 PRINT "2 FOR NON-STEADY DATA"
118 INPUT G3
119 GOTO G3 OF 130,212
120 GOTO 130
130 PRINT
140 PRINT "RUN NO. IS"R1
150 PRINT TAB7"UPPER"TAB24"LOWER"
160 PRINT TAB2"N"TAB9"CP"TAB19"N"TAB26"CP"
170 FORMAT F3.0, 2X, F7.3, 5X, F3.0, 2X, F7.3
180 FOR N=l TO 27
190 M=N+26
200 WRITE (15,170)N,C[N] ,M,C[M]
201 NEXT N
202 FIXED 3
203 PRINT "CP1="C9"AT X/C="X9", Y/C="Y9
204 STANDARD
205 DISP "CORRECT THE CP VALUES"
206 STOP
210 GOTO 270
212 DISP "INPUT DYNAMIC GAIN";
213 INPUT G2
214 FIXED 3
216 PRTNT "RUN NO. IS"R1
217 C9 = C[21] + 0.33* (C[22] - C[21])
218 PRINT "TAP"TAB7"CP (MEAN) "TAB17"CP (RMS) "TAB2 7"PHI"
220 FORMAT F3.0,3X,F7.3,3X,F7.3,3X,F4.0<br>222 FOR I=1 TO 53
224 WRITE (15, 220)1, C[IJ ,D[I J ,P [I]
226 NEXT I 227 PRINT "C9="C9
228 DISP "CORRECT THE CP VALUES"
230 STOP
235 GOTO 270
240 FOR 1=1 TO 53
242 C[I]=G2*D[I]*SIN(P[IJ/57.296)
244 NEXT I 245 C9 = (D[21] + 0.33*(D[22] - D[21]))246 G1=2
247 GOTO 270
```


```
250 FOR I=1 TO 53
 252 C[I] = G2*D[I]*COS(P[I]/57.296)254 NEXT I
 256 C9 = (D[21]+0.33*(D[22]-D[21]))257 G1 = 3258 GOTO 270
 270 L1=0272 L2=0274 P1=0276 P2=0310 FOR I=1 TO 26
 320 IF I=21 THEN 490
 330 C1 = - C[I]
 340 C2 = -C[T+1]350 X1 = A[T]360 X2 = A[I + 1]370 GOSUB 1120
 380 L1=L1+N1
 390 P1=P1+M1
400 C1 = C [I + 27]410 C2 = C[I + 26]420 X1 = A[I + 27]430 X2 = A[I + 26]440 GOSUB 1120
450 L2=L2+N1460 P2 = P2 + M1470 NEXT I
480 GOTO 640
490 C1 = -C [21]
500 C2 = -C9510 X1 = A [21]520 \tX2 = X9530 GOSUB 1120
540 L1=L1+N1
550 P1=P1+M1
560 C1 = -C9570 C2 = -C[22]580 X1 = X9590 X2 = A 22
600 GOSUB 1120
610 L1=L1+N1
620 P1=P1+M1
630 GOTO 400
640 D1=0650 D2=0
660 P3=0670 P4=0680 FOR I=1 TO 13
690 C1+C[I]
700 C2 = C[I+1]710 Y1=A[I+53]
720 Y2 = A[I + 54]730 GOSUB 1160
740 D1=D1+N1
750 P3=P3+M1
760 NEXT I
```


```
770 FOR I=40 TO 52
 780 C1=C[I]
 790 C2=C[I+1]
 800 Y1=A[I+53]
 810 Y2 = A[I + 54]820 GOSUB 1160
 830 D1=D1+N1
 840 P3=P3+M1
 850 NEXT I
 860 FOR I=14 TO 39
 870 IF I=21 THEN 970
 880 C1 = - C[I+1]
 890 C2=C[I]
 900 Y1 = A[T + 54]910 Y2 = A[I + 53]920 GOSUB 1160
 930 D2 = D2 + N1940
    P4 = P4 + M1950 NEXT I
 960 GOTO 1200
 970 \text{ } CI = -C9980 C2=C 21
 990 Y1=Y9
1000 Y2 = A 74
1010 GOSUB 1160
1020 D2=D2+N11030
    P4 = P4 + M11040 \text{ CI} = -C[22]1050 C2 = -C91060
    Y1 = A[75]1070 Y2=Y9
1080 GOSUB 1160
1090 D2=D2+N11100 P4 = P4 + M11110 GOTO 950
1120 M1=C1* (X2-X1)* (1-0.5* (X2+X1)
1130 M1=M1+(C2-C1)*(X2-X1)*(0.5-(X2/3)-(X1/6))
1140 N1=0.5* (C2+C1)* (X2-X1)1150 RETURN
1160 N1=0.5* (C2+C1)* (Y2-Y1)1170 M1=0.5*C1* (Y2-Y1)* (Y2+Y1)1180 M1=M1+(C2-C1)*(Y2-Y1)*((Y1/6)+(Y2/3))
1190 RETURN
1200 PRINT
1205 GOTO G1 OF 1210, 1500, 1600
1210
    PRINT "RUN NO. IS"R1
1220 PRINT
1225
     FIXED 4
1230
    PRINT
           "COEFFICIENT BUILDUP IS"
1250
     PRINT
           TAB2"CNU="L1
1260
    PRINT TAB2"CNL="L2
1270
    PRINT
            TAB10"TOTAL CN="(L1+L2)
1280
     PRINT
           TAB2"CCF="D1
1290
    PRINT
            TAB2"CCR="D2
1300
    PRINT TAB10"TOTAL CC="(D1+D2)
```


```
1310 PRINT TAB2"CMU="P1
1320 PRINT TAB2"CML="P2 » 1330 PRINT TAB2"CMF="P3
1340 PRINT TAB2"CMR="P4
1350 PRINT TAB10"TOTAL CM=" (P1+P2+P3+P4)"AT T.E."
1360 M1=(P1+P2+P3+P4)
1370 N1=(L1+L2)
1380 C1=(D1+D2)
1390 M5=M1-(0.75*N1)
1400 DISP "INPUT ANGLE OF ATTACK (DEGS.)";
1410 INPUT A2
1415 A2=0.01745*A2
1416 A9=A2
1420 N5=N1*C0S(A2)-C1*SIN(A2)
1430 C5=C1*C0S(A2)+N1*SIN(A2)
1440 PRINT
1445 C2=A2*57.296
1450 PRINT TAB 2"CL (MEAN) ="N5 ;TAB21"CD(MEAN) ="C5
1460 PRINT TAB2"ALPHA ="C2"DEGS"
1470 STANDARD
1480 GOTO G3 OF 90,240,250
1490 END
1500 N3=L1+L2
1502 C3=D1+D2
1504 M3+P1+P2+P3+P4
1505 FIXED 4
1506 PRINT "CN(RMS)*SIN(PHI)="N3
1508 PRINT "CC(RMS)*SIN(PHI)="C3
1510 PRINT "CM(RMS)*SIN(PHI)="M3
1511 PRINT
1595 GOTO 250
1600 N4=L1+L2
1602 C4=D1+D2
1604 M4=P1+P2+P3+P4
1605 FIXED 4
1606 PRINT "CN(RMS)*COS(PHI)="N4
1608 PRINT "CC(RMS)*COS(PHI)="C4
1610 PRINT "CM(RMS)*COS(PHI)="M4
1611 PRINT
1612 Nl=SQR(N3+2+N4+2)
1613 P1=ATN(N3/N4)
1614 X=N3
1615 Y=N4
1616 F=P1
1617 GOSUB 1900
1618 P1=F
1620 PRINT "CN(RMS)="N1,"PHI="P1*57.296
1630 Cl=SQR(C3+2+C4+2)
1640 P2=ATN(C3/C4)
1641 X=C3
1642 Y=C4
1643 F=P2
1644 GOSUB 1900
1645 P2=F
1650 PRINT "CC(RMS)="C1,"PHI="P2*57.296
1660 Ml=SQR(M3+2+M4+2)
```


```
1670 P3=ATN(M3/M4)
1671 X=M3
1672 Y=M4
1673 F=P3
1674 GOSUB 1900
1675 P3=F
1680 PRINT "CM(RMS)="M1,"PHI="P3*57.296
1690 A1=N1*C0S(A9)*SIN(P1)-C1*SIN(A9)*SIN(P2)
1700 B1=N1*C0S(A9)*C0S(P1)-C1*SIN(A9)*C0S(P2)
1710 L9=SQR(Al+2+Bl+2)
1720 L8=ATN(A1/B1)
1721 X=A1
1722 Y=B1
1723 F=L8
1724 GOSUB 1900
1725 L8=F
1730 A2=N1*SIN(A9)*SIN(P1)+C1*C0S(A9)+SIN(P2)
1740 B2=N1*C0S(P1)*SIN(A9)+C1*C0S(A9)*C0S(P2)
1750 D9=SQR(A2+2+B2+2)
1760 D8=ATN(A2/B2)
1761 X=A2
1762 Y=B2
1763 F=D8
1764 GOSUB 1900
1765 D8=F
1770 A3=M1*SIN(P3)-0.75*N1*SIN(P1)
1780 B3=M1*COS(P3)-0.7 5*N1*COS(P1)
1790 M9=SQR(A3+2+B3+2)
1800 M8=ATN(A3/B3)
1801 X=A3
1802 Y=B3
1803 F=M8
1804 GOSUB 1900
1805 M8=F
1810 PRINT
1820 PRINT TAB10"AERODYNAMIC COEFFICIENT SUMMARY"
1830 PRINT
1840 PRINT "CL="N5" +"L9"SIN(W1*T+"L8*57 . 296")
1850 PRINT
1860 PRINT "CD="C5"+"D9"SIN(W1*T+"D8*57.296")"
1870 PRINT
1880 PRINT "CM(C/4)="M5"+"M9"SIN(W1*T+"M8*57.296")"
1890 GOTO 90
1900 IF X<0 AND Y<0 THEN 1920
1910 IF X>0 AND Y<0 THEN 1920
1915 GOTO 1930
1920 F=F+PI
19 30 RETURN
1940 STOP
```
\bar{I} ϵ

UNSTEADY DATA OUTPUT

 $CN(RMS)*SIN(PHI) = -1.5385$ $CC(RMS)^*SIN$ $(PHI)^*=-0.0312$ $CM(RMS) * SIN(PHI) = -0.6336$ $CN (RMS)*COS (PHI) = -1.4968$ $CC(RMS)^*COS(PHI) = -0.0042$ $CM(RMS)*COS(PHI) = -0.8581$ $CN(RMS) = 2.1465$
 $CC(RMS) = 0.0315$

PHI=262.2760 $CC(RMS) = 0.0315$
 $CM(RMS) = 1.0667$
 $PHI = 216.4451$ $CM(RMS) = 1.0667$

AERODYNAMIC COEFFICIENT SUMMARY

HPPFNDIX R

RPPENDIX C

UNSTERDY PRESSURE DRTR

UNSTERDY PRESSURE DATA

(CONTINUED)

HNGLE OF HTTHCK=-5 DEG

 $U = 127$ FT/SEC

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