

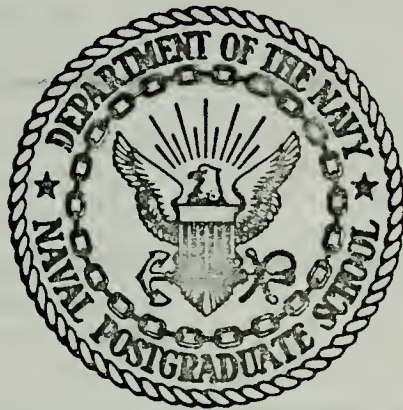
ELECTROEXCITATION OF GIANT RESONANCES  
BETWEEN 5 MeV AND 40 MeV  
EXCITATION ENERGY IN <sup>197</sup>Au

Kenneth Paul Ferlic

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# NAVAL POSTGRADUATE SCHOOL

## Monterey, California



# THESIS

ELECTROEXCITATION OF GIANT RESONANCES  
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by

Kenneth Paul Ferlic  
and  
Ronald Dallas Waddell

June 1974

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by

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June 1974



ABSTRACT

Inelastic electron scattering from  $^{197}\text{Au}$  between 5 MeV and 40 MeV excitation energy revealed giant multipole resonances. Angular distribution studies with incident 90 MeV electrons show evidence of previously unreported giant monopole states. Reduced matrix elements have been extracted and multipolarity assignments have been made. Seven states have been observed at excitations of 7.3 (undetermined magnetic), 9.2 (E0), 10.8 (E2), 14.0 (E1), 18.0 (undetermined), 22.5 (E2), and 33.5 MeV (E0 or E2).



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## I. INTRODUCTION

Giant resonance investigations of heavy nuclei have been actively pursued at the Naval Postgraduate School electron linear accelerator since 1973. Initial studies by Warshawsky and Webber [Ref. 1] demonstrated the feasibility of measuring inelastic electron spectra in the continuum region of  $^{197}\text{Au}$ . They adapted analysis techniques developed at Darmstadt and measured cross sections for several known resonances. Warshawsky and Webber further suggested the existence of previously unreported resonances at about 9 MeV and 23 MeV excitation energy, but their analysis did not conclusively identify the multipolarities.

The work described in this thesis completes the study of  $^{197}\text{Au}$  discussed above and extends the excitation energy range to 40 MeV. Seven giant multipole resonances have been observed. Besides the previously reported 7.3, 10.8, and 14.0 states [Refs. 1-4, 18 and 19], the existence of resonances at 9.2 and 22.5 MeV are required to explain the observed electron spectra.

Experimental values of the inelastic scattering form factors were measured at scattering angles of  $60^\circ$ ,  $75^\circ$ ,  $90^\circ$ ,  $105^\circ$ , and  $120^\circ$  with incident electrons of 90 MeV energy. Transition multipolarities have been deduced and reduced transition probabilities have been extracted. Newly identified are resonances at excitation energies of 9.2 MeV (isoscalar E0), 22.5 MeV (isovector E2), and 33.5 MeV (isovector E0 or E2). Previously reported multipolarities for the 10.8 MeV (isoscalar E2) and 14.0 MeV (isovector E1) [Ref. 18] resonances have been confirmed. The analysis contained here is unable to uniquely assign transition multipolarities to the 7.3 MeV (magnetic) and 18.0 MeV structures.



## II. THEORY

### A. ELECTRON SCATTERING

The Mott cross section for an extremely relativistic electron which scatters elastically from a nucleus of point charge  $Ze$  is given by [Ref. 6]

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} = \left(\frac{Ze^2}{2E_i}\right)^2 \frac{\cos^2 \frac{\Theta}{2}}{\sin^4 \frac{\Theta}{2}} \quad \text{II-1}$$

The electron is assumed to have spin while the nucleus has neither spin nor magnetic moment. Nuclear charge structure in a nucleus of finite extent modifies the scattering cross section to [Ref. 6]

$$\left(\frac{d\sigma}{d\Omega}\right)_{\text{measured}} = \left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}} \left| \int_{\text{nuclear volume}} \rho(\vec{r}) e^{i\vec{q}\cdot\vec{r}} dV \right|^2 \quad \text{II-2}$$

$\rho(\vec{r})$  is the charge density within the nucleus as a function of  $\vec{r}$ , the radius vector from the center of the nucleus, and  $\vec{q}$  is the momentum transfer vector. The Born approximation has been assumed. The term describing the nuclear structure is called the square of the form factor,

$$\left| F(q^2) \right|^2 = \frac{\left(\frac{d\sigma}{d\Omega}\right)_{\text{measured}}}{\left(\frac{d\sigma}{d\Omega}\right)_{\text{Mott}}} \quad \text{II-3}$$

For inelastic scattering, the cross section may also be written as a modification to the Mott scattering cross section, but the form is more complex, with both transverse and longitudinal contributions which are described elsewhere [Ref. 7]. For inelastic longitudinal scattering the form factor is related to the matrix element of the charge





operator evaluated between initial and final nuclear states. Although transverse components require modification of the simple expression II-2, it is conventional to call the experimental form factor squared the ratio of measured to Mott cross sections as in II-3. This value is compared to theoretical calculations based on particular nuclear models.

For inelastic scattering, the Born approximation may be used to compute the theoretical form factor based on an assumed nuclear transition charge distribution. First attempts at a solution were made by describing the perturbing fields in terms of a multipole expansion, and assuming the incident and scattered electron wave functions to be plane waves. This so called plane wave Born approximation (PWBA) is satisfactory for light nuclei, but becomes less accurate as the number of protons in the nucleus increases. For heavy nuclei, the Coulomb field surrounding the nucleus distorts both the incident and scattered electron wave functions sufficiently that a plane wave solution no longer gives the correct results. Consequently, the plane wave solutions of the PWBA have to be replaced by phase shifted spherical waves. This method, termed the distorted wave Born approximation (DWBA), in practice requires computer computations to extract the form factors [Ref. 8]. Form factors thus calculated are plotted as a function of scattering angle for each desired electric and magnetic transition. Comparison of these plots with experimental results allows for the assignment of transition modes to the observed resonances.

## B. GIANT RESONANCES

The term "giant resonance" has been used until recently to describe the giant electric dipole resonance (GDR). This resonance appears as a



well known feature in the nuclear excitation spectra produced by such photonuclear processes as photodisintegration, photon absorption, and inverse reactions such as the  $(p, \gamma)$  reaction. The presence of the GDR is characterized by a giant bump in photonuclear spectra at excitation energies from 10 to 25 MeV. It occurs throughout the periodic table and, according to the Goldhaber-Teller and to the Steinwedel-Jensen model, appears at excitation energies that are proportional to the inverse nuclear radius,  $E_x = 80A^{-1/3}$  [Ref. 9].

This conspicuous resonance appears also in inelastic electron scattering spectra. The advantage of electron scattering experiments over the photonuclear processes for exploring the nature of the giant resonances lies in the variable momentum transfer,  $q$ . In the case of the photonuclear processes, the momentum transfer is determined uniquely by the excitation energy. The giant resonance in the electron scattering cross section can be measured at different values of  $q$  depending on the scattering angle,  $\Theta$ , and the incident electron energy,  $E_i$ ; the dependence being given by

$$q^2 = 4E_i E_f \sin^2 \frac{\Theta}{2}, \quad \text{II-4}$$

where  $E_f$  is the final electron energy. In addition, at values of  $q$  higher than is available from photonuclear work, excitation modes other than  $E1$  become observable. Consequently, the term "giant resonance" now refers to all multipole excitation modes which exhibit strongly coherent properties.

Early attempts to explain the giant resonances were based mainly on collective models. The initial Goldhaber-Teller model assumed a rigid displacement of the protons as a whole against the neutrons as a whole, thus producing a large collective dipole moment. The



Generalized Goldhaber-Teller model assumes that nuclear matter consists of four interpenetrating fluids: neutrons with spin up ( $n\uparrow$ ), protons with spin up ( $p\uparrow$ ), neutrons with spin down ( $n\downarrow$ ), and protons with spin down ( $p\downarrow$ ). Any two of these fluids may oscillate  $180^\circ$  out of phase with the other two, giving rise to four possible collective modes. For example, nucleons with spin up could oscillate against nucleons with spin down producing a magnetic resonance [Ref. 5].

The Steinwedel-Jensen model or hydrodynamic model is based on the oscillation of neutron and proton fluids within a rigid nuclear surface. An extension of this model is the dynamic collective model [Ref. 9]. The dynamic collective model replaces the rigid nuclear surface of the Steinwedel-Jensen model by a surface that may vibrate. These surface vibrations, however, are much slower than the internal giant dipole oscillations so that the approximation can be made that the dipole oscillations take place in a deformed nucleus.

A different model of the dipole giant resonances, used mainly for light, closed-shell nuclei, is based on the shell model of the nucleus. This particle-hole (ph) or Brown model arises from the observation by Wilkinson [Ref. 10] that a  $1p$  nucleon, promoted into the  $2s-1d$  shell, will produce dipole strengths large enough to explain the observed giant resonance cross sections. The shell model description [Ref. 10] gave the strength of the absorption correctly, but the resonance energy was incorrect, and the effect of the particle-hole interaction was noted by Brown [Ref. 11].

It has been shown [Ref. 9] that such a shell model is separable in the relative co-ordinate of all the protons versus all the neutrons; and that, upon absorption of a dipole photon, a particular linear combination of single-particle excitations is produced which corresponds





to a motion of all the protons against all the neutrons; i.e., a Goldhaber-Teller collective vibration. Hence, the shell model and collective models are equivalent in their description of collective states.

Giants resonances, then, may be described as strongly collective modes of excitation in which a considerable number of nucleons take part. In the last several years, much progress has been made in the investigation of these states using models such as those described above. Several modes of excitation other than the GDR have been identified, including the giant electric quadrupole resonance (GQR) and the giant magnetic dipole resonance (GMDR). A short summary of some of these works follow.

Early electron scattering investigation of the giant resonance structure was carried out by Isabelle and Bishop on  $^{16}\text{O}$  [Ref. 7]. By studying the angular dependence of cross sections, the giant dipole was observed. These experiments demonstrated the feasibility of systematic studies of the level schemes and transitions in light nuclei.

Yamaguchi et. al. [Ref. 12] have carried out extensive investigations of the fine structure in the giant resonance region of  $^{12}\text{C}$ . By making measurements at both backward and forward angles while keeping the momentum transfer constant, they were able to extract the longitudinal and transverse form factors as functions of excitation energy. The resolution of the fine structure in this manner revealed seven peaks associated with the longitudinal form factors and nine with the transverse. The dipole strength is distributed among a number of states to produce this fine structure in the giant resonance region.

An investigation of the isotopic characteristics of the giant dipole resonance has been carried out by Carlos et. al. [Ref. 13]. The giant





dipole resonances for the Neodymium isotopes  $^{142}\text{Nd}$  to  $^{146}\text{Nd}$ ,  $^{148}\text{Nd}$ , and  $^{150}\text{Nd}$  were studied with a monochromatic photon beam by measuring the partial photonuclear cross sections  $\sigma(\gamma, n)$ ,  $\sigma(\gamma, pn)$  and  $\sigma(\gamma, 2n)$ . As the number of neutrons in the nucleus increased, the dipole resonance began to broaden. For  $^{150}\text{Nd}$ , the resonance actually split into two distinct peaks. An interpretation is given by the dynamic collective model in which there is a coupling between dipole oscillations and surface vibrations. The splitting of the resonance in  $^{150}\text{Nd}$  characterizes a deformed nucleus, and the experimental cross section was closely reproduced by superposition of two separate Lorentz lines whose centers are separated by four MeV. In the framework of the hydrodynamic model, one can assume for an axially deformed nucleus two modes of excitation at differing energies, equivalent to oscillations in the direction of the two axes.

Pitthan and Walcher [Ref. 14] did (e,e') investigations with Ce, La, and Pr. All three elements show distinct resonances at excitation energies of 9, 12, and 15 MeV. The resonance at 15 MeV was identified as the E1 giant resonance based on correlation of the excitation energy and strength with the GDR seen in photo-neutron experiments. The 12 MeV resonance exhibited longitudinal properties similar to the E1 but did not appear in  $(\gamma, n)$  studies. It was assigned an E2 or E0 excitation mode based on comparison to form factors developed from the hydrodynamic model. The 9 MeV resonance was found to be transverse and most likely an M1 transition.

Later, Lewis and Bertrand [Ref. 15] observed that the (p,p') spectra of  $^{27}\text{Al}$ ,  $^{54}\text{Fe}$ ,  $^{120}\text{Sn}$ , and  $^{209}\text{Bi}$  consistently produced enhancements at excitation energies slightly lower than those of the known giant



dipole resonance. Since the resonance energies of the giant dipole had been previously well measured, they concluded that it was unlikely that the giant dipole excitation was responsible for most of the observed enhancement. This apparent resonance shift was due to the presence of a giant isoscaler-quadrupole vibration located at  $E_x = 63 A^{-1/3} \text{ MeV}$ .

In forward angle (e,e') experiments, Fukuda and Torizuka observed [Ref. 16] similar E1 and E2 resonances in  $^{90}\text{Zr}$ . The usual dipole resonance was seen at 16.65 MeV and additional resonances were seen at 14 MeV and 28 MeV. Based on a prediction of the hydrodynamic model that a quadrupole giant resonance should appear at a position 1.6 times the peak energy of the usual giant dipole resonance [Ref. 17], the 28 MeV resonance was assigned an E2 transition. An assignment of either E0 or E2 was made for the 14 MeV peak. Based on similar results for  $^{59}\text{Fe}$ ,  $^{116}\text{Sn}$ , and  $^{208}\text{Pb}$  [Ref. 16] these new giant resonances seem to be universal features existing at excitation energies of  $65 A^{-1/3} \text{ MeV}$  and  $130 A^{-1/3} \text{ MeV}$  in all nuclei. Table I contains a listing of the  $A^{-1/3}$  dependence of the resonances identified in this work and a list of references where others have seen or predicted corresponding resonances in  $^{197}\text{Au}$  and other nuclei.

TABLE I  
 $A^{-1/3}$  Dependence of  $^{197}\text{Au}$  Resonances

$E_x$ (MeV)	$E_x$ ( $A^{-1/3}$ -MeV)	References
9.2	53	t [Refs. 21 & 22] Ce, La, Pr [Ref. 23]
10.8	63	Au [Refs. 3 & 18]
14.0	81	Au [Refs. 3, 18, and 25]
18.0	105	Pb [Refs. 24 & 26]
22.5	130	Au [Ref. 3] Pb [Ref. 24]
33.5	195	t [Ref. 22] Pb [Ref. 26]

t = Theoretical Prediction



### C. TRANSITION PROBABILITIES AND SUM RULES

The transverse-electric and magnetic form factors are given by [Ref. 27]

$$\left| F_{\vec{m}, L}^{\vec{E}, L}(\vec{q}_T) \right|^2 = \frac{4\pi}{Z^2} \left( \frac{L+1}{L} \right) \left[ \frac{q_T^L}{2(L+1)!!} \right]^2 B(\vec{E}_m, L, \vec{q}_T, J_i \rightarrow J_f) \quad \text{II-5}$$

The continuity equation provides similar relationships for the longitudinal form factor. The coefficients B are the reduced nuclear transition probabilities and, as such, contain all of the nuclear physics of the inelastic scattering problem. The theoretical form factors are calculated by the GBROW computer program as described by Ziegler [Ref. 27].

To determine whether or not the resonances observed are collective phenomena, comparison of experimental to single-particle reduced transition probabilities may be made. For a heavy nucleus, one would expect a strength greater than several times the single-particle value. Furthermore, a giant resonance should exhaust an appreciable fraction of the appropriate sum rule. If an observed resonance greatly exceeds the sum rule under the assumption of a particular transition multipolarity, it is unlikely that the multipolarity assignment is correct.

The single-particle reduced transition strengths are given by the Weisskopf unit [Ref. 28],

$$B(EL)_{SP} = \frac{e^2(2L+1)}{4\pi} \left( \frac{3}{L+3} R^L \right)^2 \quad \text{II-6}$$

and

$$B(ML)_{SP} = \frac{e^2(2L+1) \cdot 10}{\pi} \left( \frac{3}{L+3} \right)^2 R^{2L-2} (0.0111), \quad \text{II-7}$$



where  $R = 1.2 A^{-1/3} = 6.98$  fm for  $^{197}\text{Au}$ . Table II provides values of the single-particle transition strengths for  $^{197}\text{Au}$ .

Another, to some extent model-independent, evaluation of the transition strength of EL modes may be achieved by expressing the strength relative to the energy-weighted sum rule (EWSR). Sum rules for magnetic transitions are more dependent on the structure of the nucleus in question and, therefore, no comparisons to these have been made here. States at excitation energies less than 15 MeV appear to be  $\Delta T = 0$  except for the EL mode which is  $\Delta T = 1$ . The sum rule for isoscalar ( $\Delta T = 0$ ) excitation modes with  $L > 1$  is given by [Ref. 29]

$$S(EL, \Delta T=0) = \sum_f (E_f - E_i) B(EL, i \rightarrow f) = \frac{Z^2 e^2 L(2L+1)^2 \hbar^2 \langle R^{2L-2} \rangle}{8\pi A M} \quad \text{II-8}$$

The isovector ( $\Delta T = 1$ ) sum rule for  $L > 1$  is related to the isoscalar sum rule by

$$S(EL, \Delta T=1) = S(EL, \Delta T=0) \frac{N}{Z} \quad \text{II-9}$$

The corresponding sum rule for an isoscalar monopole (E0) excitation is [Ref. 30]

$$S(E0) = \sum_f (E_f - E_i) |M_{fi}|^2 = \frac{\hbar^2 Z}{M} \langle R^2 \rangle, \quad \text{II-10}$$

where  $M_{fi}$  is the monopole matrix element. The strength of E1 resonances can be expressed relative to the sum rule given by Warburton [Ref. 31]:

$$S(E1) = \frac{e^2 \hbar^2}{2 M_P} \left( \frac{9}{4\pi} \right) \frac{NZ}{A} \quad \text{II-11}$$

Table III gives values of these sums for  $^{197}\text{Au}$  using weighted radii of the ground state charge distribution.





TABLE II  
Weisskopf Units for  $^{197}\text{Au}$

L	$\frac{B(EL)}{e^2}$ *	$\frac{B(ML)}{e^2}$ *
1	6.6	.06
2	$3.35 \times 10^2$	3.1
3	$1.63 \times 10^4$	$1.44 \times 10^2$
* Units are $\text{fm}^{2L}$		

TABLE III  
Sum of EWSR For  $^{197}\text{Au}$

EL	$\Delta T$	Sum	Units
E0	0	$9.2 \times 10^4$	$\text{MeV-fm}^4$
E1	1	$7.04 \times 10^2$	$e^2\text{-MeV-fm}^2$
E2	0	$7.4 \times 10^4$	$e^2\text{-MeV-fm}^4$
	1	$1.1 \times 10^5$	$e^2\text{-MeV-fm}^4$
E3	0	$6.15 \times 10^6$	$e^2\text{-MeV-fm}^6$
	1	$9.16 \times 10^6$	$e^2\text{-MeV-fm}^6$



#### D. QUASI-ELASTIC SCATTERING

Electrons incident upon a nucleus may scatter elastically with individual nucleons within the nucleus. This process is called quasi-elastic scattering. If sufficient momentum is imparted to the nucleon, it will be knocked out of the nucleus with some kinetic energy. The simultaneous energy loss by the incident electron may result in a measurable quasi-elastic peak. It is of interest to see at which value of  $E_x$  this contribution appears on the spectrum of  $\frac{d^2\sigma}{d\Omega dE_x}$ , to see if it contributes to, and therefore must be subtracted from, the scattered electron spectrum in the region of the giant resonance.

The inelastic electron-nucleus cross section has been calculated in three energy-loss regions with the Fermi gas model for the nucleus [Ref. 32]: the quasi-elastic region; the threshold pion region with its pion electroproduction cross sections; and the region containing the 3-3 resonance of the proton. Pion production begins at 135 MeV and the 3-3 resonance is 298 MeV above the elastic peak, so these regions are not of interest in this work.

In the Fermi gas model, the electron scatters elastically from a single nucleon in the free Fermi sea, with the recoiling nucleon required to scatter out of the Fermi sphere. Because of the Pauli principle, the nucleon cannot scatter into an already occupied state. The final nucleon energy is given by

$$E_{kf} = [k^2 + M^2]^{\frac{1}{2}}, \quad \text{II-12}$$



where  $k$  is the nucleon momentum and  $M$  is the nucleon mass. The initial (bound) nucleon energy is

$$\epsilon_{ki} = \frac{k^2}{2M} + U(k^2) = \frac{k^2}{2M^*} + U(0), \quad \text{II-13}$$

where the effective nucleon mass,  $M^*$ , is given by [Ref. 33]

$$M^* = \frac{M}{1.4} \quad \text{II-14}$$

and  $U(k^2)$  is the effective single-particle potential in nuclear matter. This potential effectively shifts the electron energy loss,  $\omega$ , to take into account the nuclear binding. Moni treats  $U(k^2)$  as a constant,  $\bar{\epsilon}$ , equivalent to the average nucleon interaction energy. The quasi-elastic peak energy loss will increase with  $q$  and approach the effective free nucleon kinetic energy,  $q^2 / (2M^*)$ , at large values of  $q$ .

The magnitude of the above shift for the gold nucleus is of interest. The average nucleon interaction energy,  $\bar{\epsilon}$ , and the nuclear Fermi momentum,  $k_F$ , have been determined for nine elements from  ${}^6\text{Li}$  to  ${}^{208}\text{Pb}$  [Ref. 34] by a least squares fit of theory to quasi-elastic peak data. In the cross-section formula, an energy-conserving delta function appears, given by

$$\delta\left(\omega + \frac{k^2}{2M} - \bar{\epsilon} - \frac{(\vec{k} + \vec{q})^2}{2M}\right), \quad \text{II-15}$$



which determines  $\omega$ . The nucleon momentum is approximated by the Fermi momentum and is roughly constant at 260-265 MeV/c for heavy nuclei from nickel through lead.  $\bar{\epsilon}_{Au} = 43$  MeV [Ref. 34]. The energy-conserving delta function requires that

$$\omega = \bar{\epsilon} + \frac{kq}{m} + \frac{q^2}{2M}. \quad \text{II-16}$$

For the smallest momentum transfer of these experiments,  $\omega = 73$  MeV. The threshold value is  $\omega \cong \bar{\epsilon}_{Au} = 43$  MeV. Thus explicit corrections for the quasi-elastic contribution need not be made for the spectra of this work.





### III. DATA ANALYSIS

#### A. DATA REDUCTION

Five experiments were performed with an incident beam energy of 90 MeV and at scattering angles of  $60^\circ$ ,  $75^\circ$ ,  $90^\circ$ ,  $105^\circ$ , and  $120^\circ$ . The experimental information is given in Table IV. In general the data analysis procedure described by Warshawsky and Webber [Ref. 1] was followed throughout this work. The initial choice of resonance energies was determined from known resonances [Refs. 1, 2, 3, and 4] and by visual inspection. Previously determined resonances occur at 7.3, 9.2, 11.0, and 14.0 MeV. Visual inspection of the background subtracted data revealed wide structures at 18.0, 22.5, and 33.5 MeV.

The fitted inelastic spectra extended from 4 to 40 MeV, except the  $120^\circ$  spectrum which extends only to 31.5 MeV. In some cases the data extends to 50 MeV but no structures were observed at this high excitation energy. Since the  $120^\circ$  data has different boundary conditions for matching the background at high excitation energies, only cross sections for resonances below 20 MeV were extracted at this angle.

Breit-Wigner resonance forms were assumed for each resonance and computer code NAW (see appendix) simultaneously fits the entire inelastic spectrum including the radiation tail and background. The new resonance positions and widths were manually varied while the NAW code varied the peak strengths and background parameters until a best fit was obtained.

The criteria used to determine a reasonable fit were as follows:

1. The data and calculated spectrum should coincide visually.
2.  $\chi^2$  per degree of freedom should be less than one. The data is not strictly statistical because the detector momentum interval is larger



than the momentum increment of the spectrometer field and hence correlations exist between energy bins.

3. The difference between the total  $\chi^2$  for successive computer iterations was  $\Delta\chi^2 \approx 0.5$ . Typically there were 300 degrees of freedom and a change of 0.5 in the total chi squared resulted in  $1.7 \times 10^{-3}$  per degree of freedom.

4. All observed resonances should consistently fit each of the five spectra. Table V lists the positions and widths of the resonances required.

Figures 1 - 5 present the experimental data corrected for spectrometer dispersion effects with the fitted total background and the individual resonances drawn. Figures 6 - 10 present these spectra with the total background subtracted, revealing the resonance structures of Table V.

The areas under the inelastic resonances were calculated from the Breit-Wigner resonance shapes and the ratio of each inelastic resonance area to the area under the elastic peak was calculated ( $A_i/A_e$ ). The inelastic form factors squared are obtained by multiplying  $A_i/A_e$  by the square of the elastic scattering form factors ( $F_{e1}^2$ ). Table VII presents the experimental ratios of the inelastic to elastic areas and the square of the inelastic form factors. The elastic form factors were calculated using the program of Rawitscher and Fischer [Ref. 35]. The charge distribution parameters were  $c = 6.38$  fm and  $t = 2.32$  fm.



TABLE IV  
Experimental Conditions

Angle	Target Thickness	Incident Energy	$F_{el}^2$
60	.048 g/cm <sup>2</sup>	90.32 MeV	.1185
75	.096 g/cm <sup>2</sup>	90.24 MeV	.0444
90	.096 g/cm <sup>2</sup>	90.21 MeV	.0298
105	.144 g/cm <sup>2</sup>	89.98 MeV	.0217
120	.240 g/cm <sup>2</sup>	89.60 MeV	.0132

TABLE V  
Observed Resonances

Excitation Energy MeV	Widths MeV
7.3 $\pm$ 0.3	2.8 $\pm$ 0.4
9.1 $\pm$ 0.4	2.3 $\pm$ 1.0
10.9 $\pm$ 0.4	2.8 $\pm$ 0.4
14.0 $\pm$ 0.4	4.4 $\pm$ 0.4
18.0 $\pm$ 0.6	5.1 $\pm$ 0.6
22.5 $\pm$ 0.6	7.2 $\pm$ 0.6
33.5 $\pm$ 1.4	10.5 $\pm$ 2.5



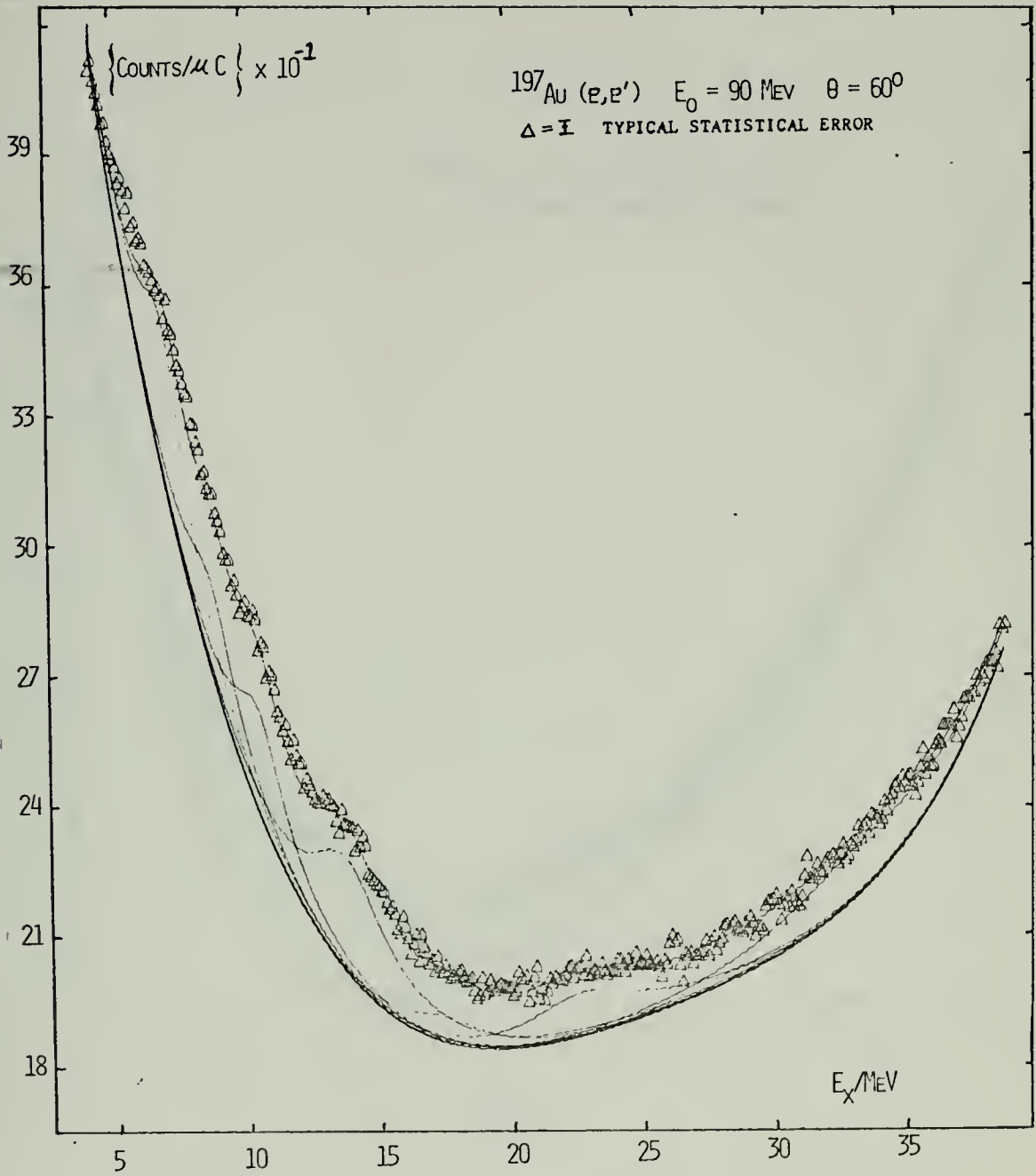


FIGURE 1





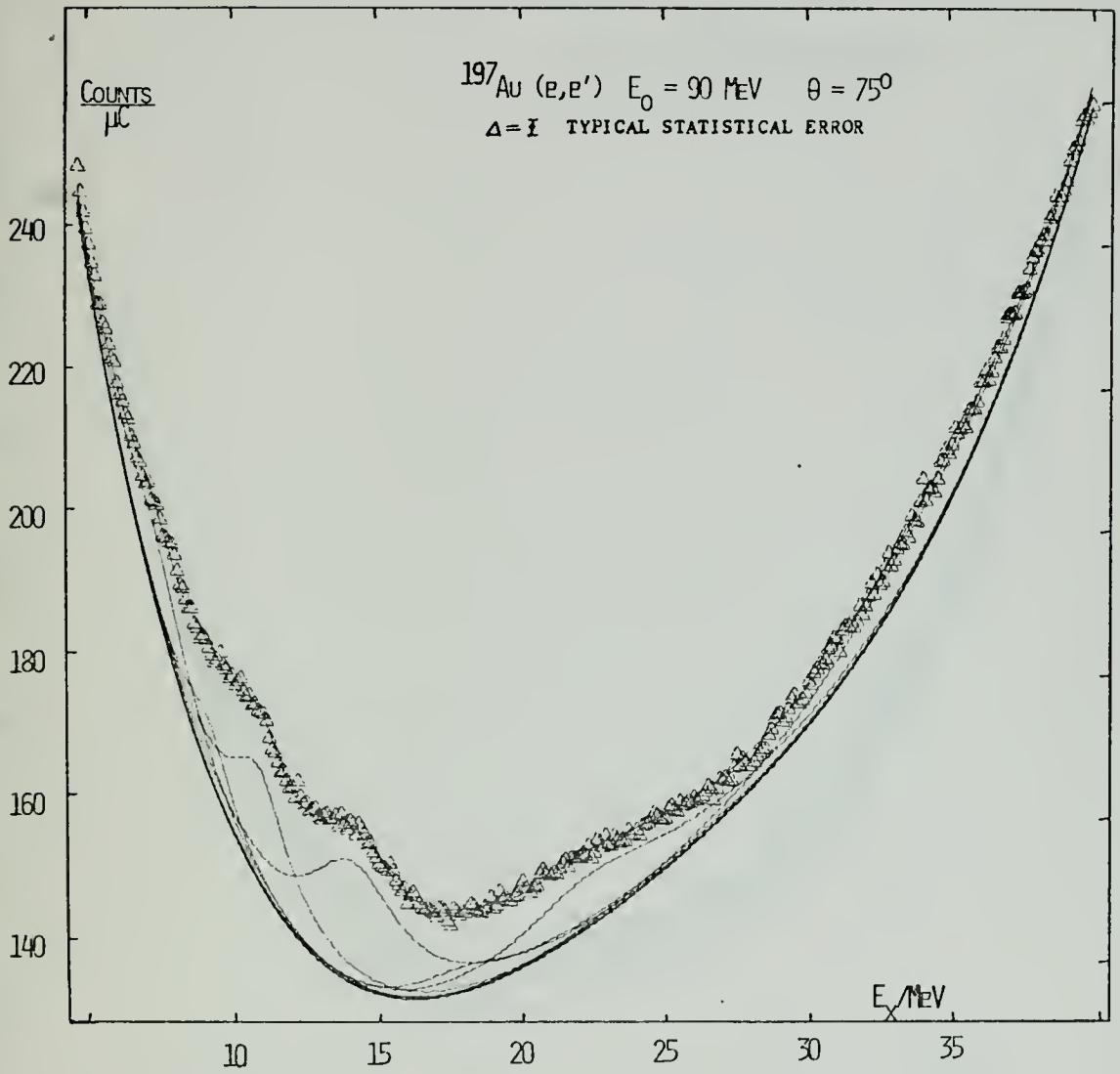


FIGURE 2



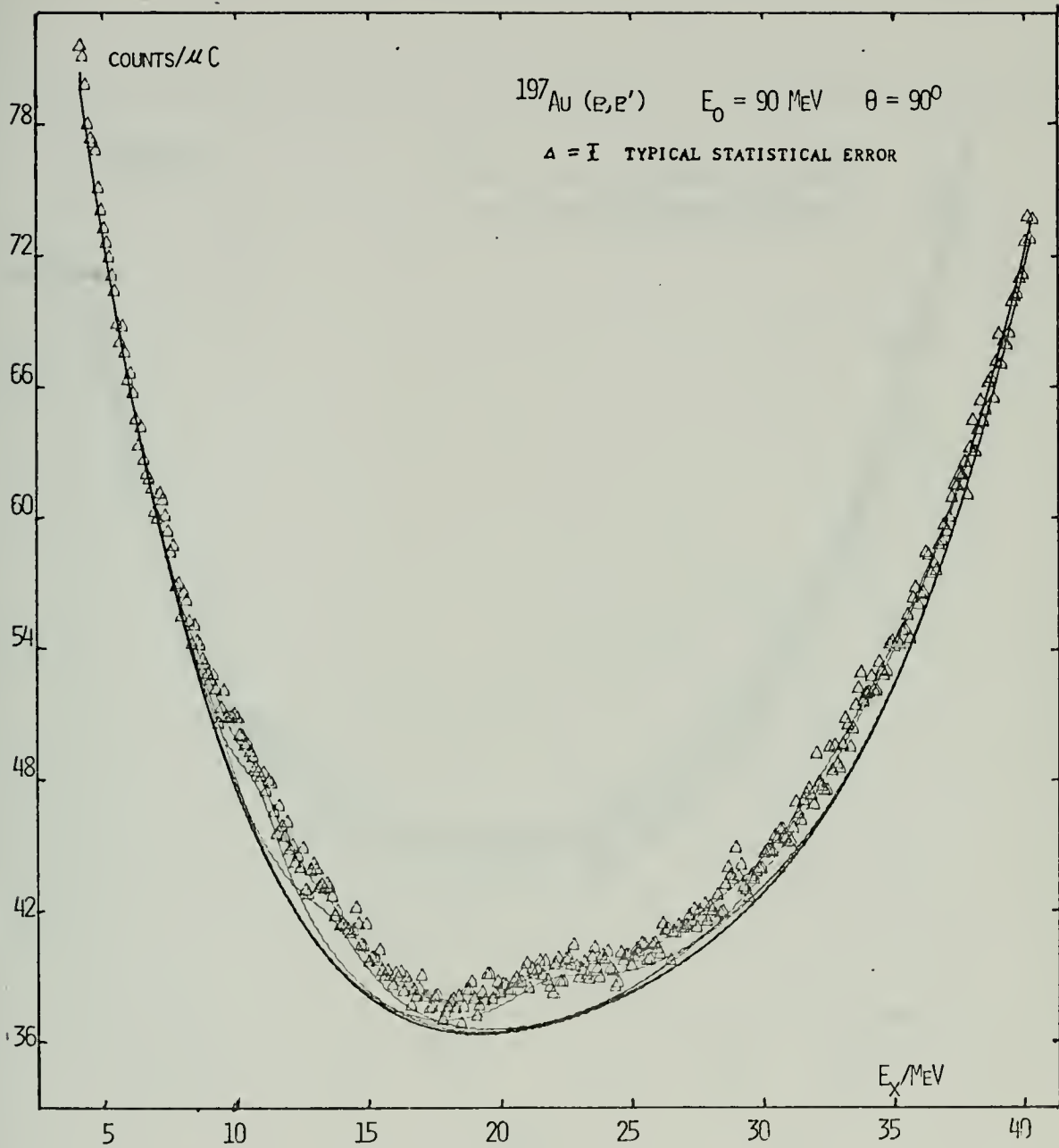


FIGURE 3



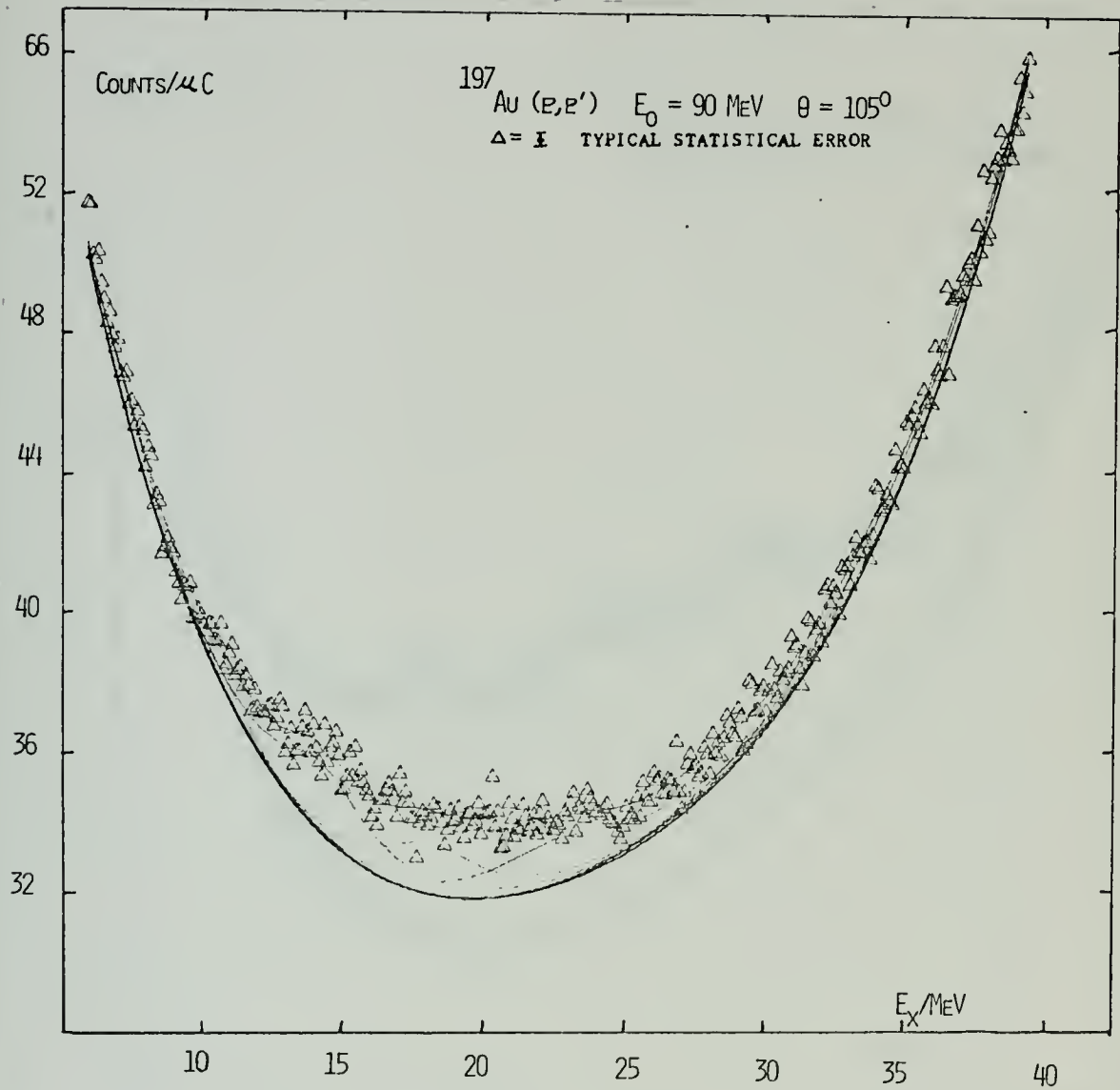


FIGURE 4



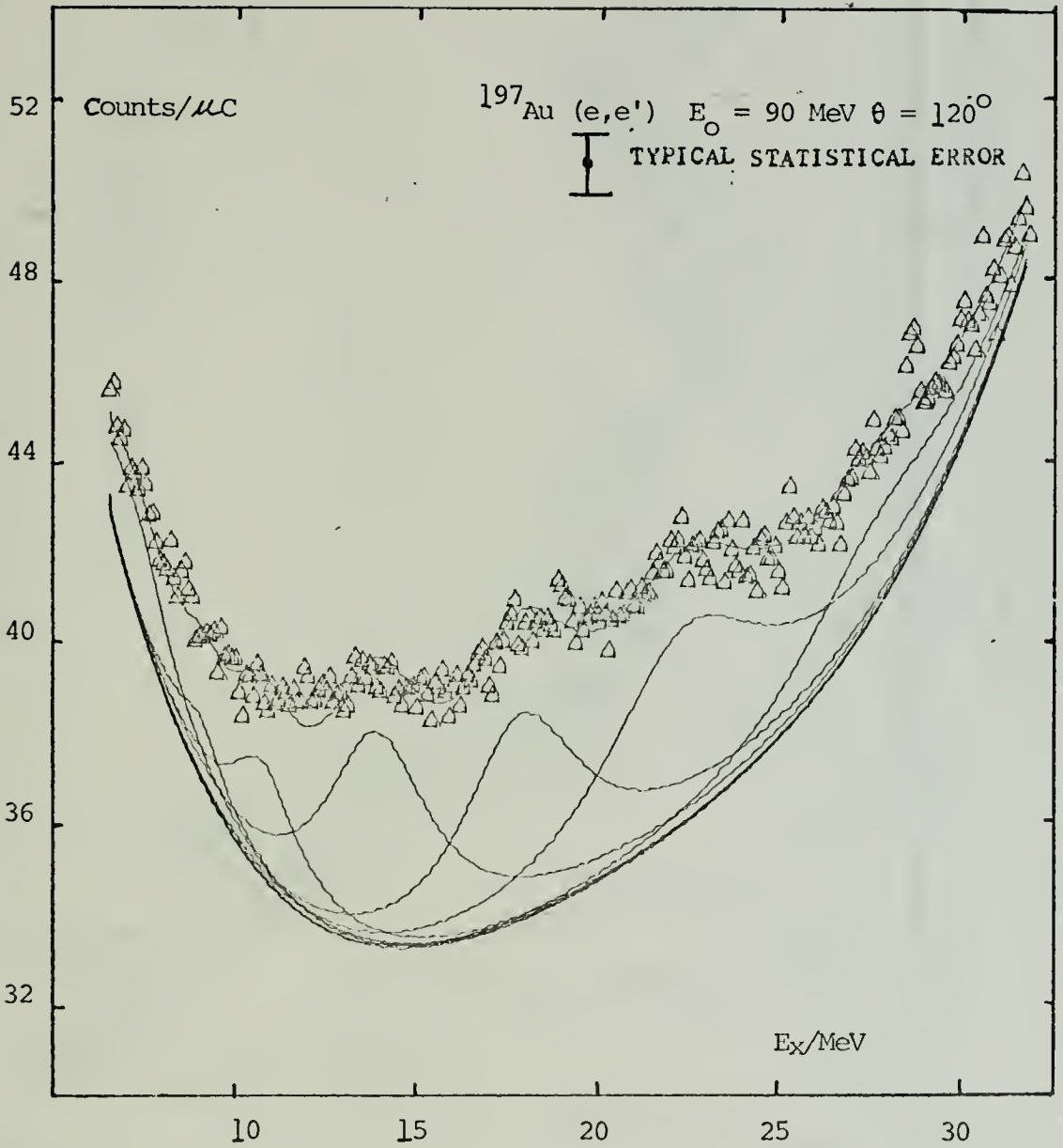


FIGURE 5





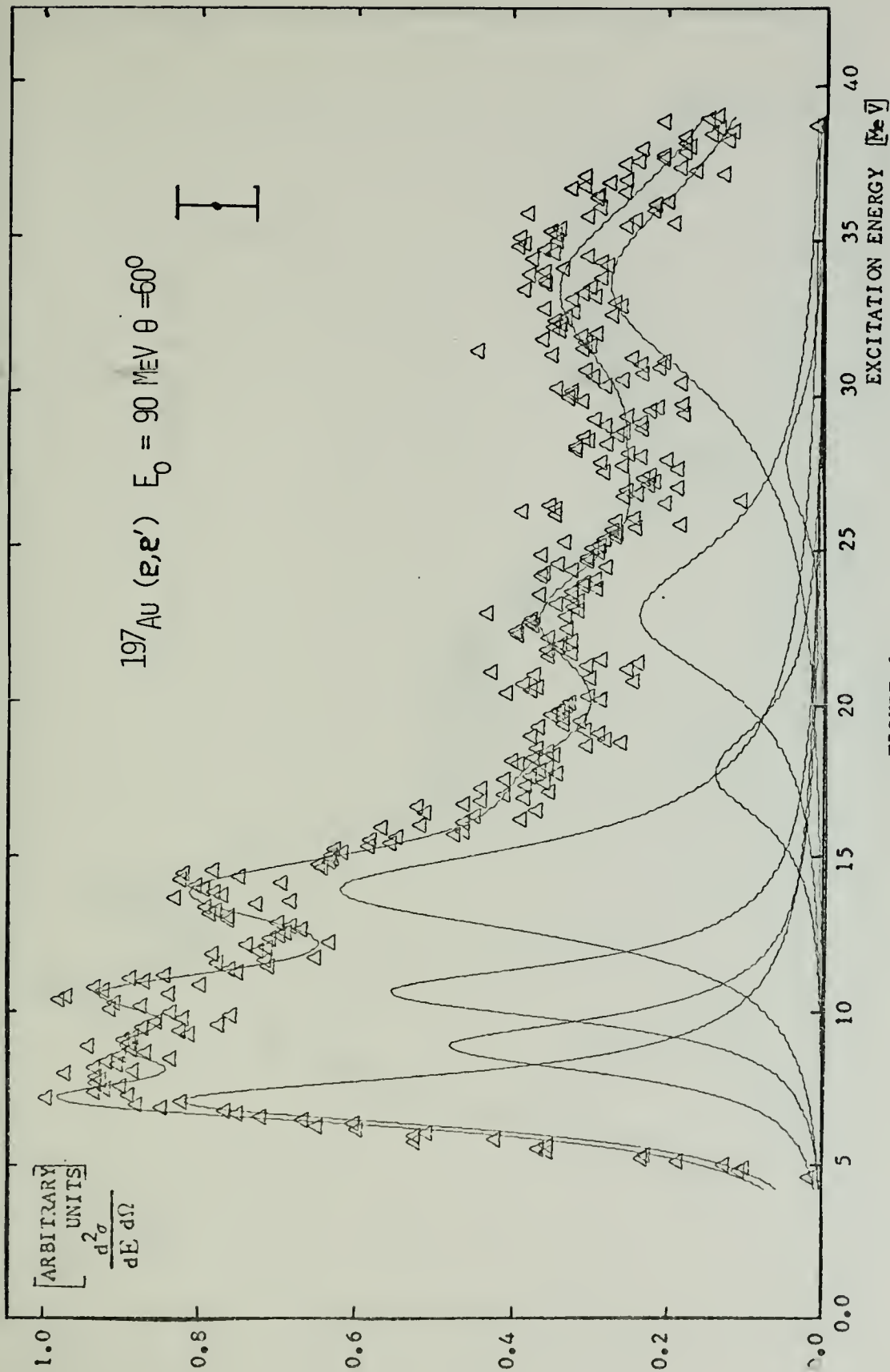
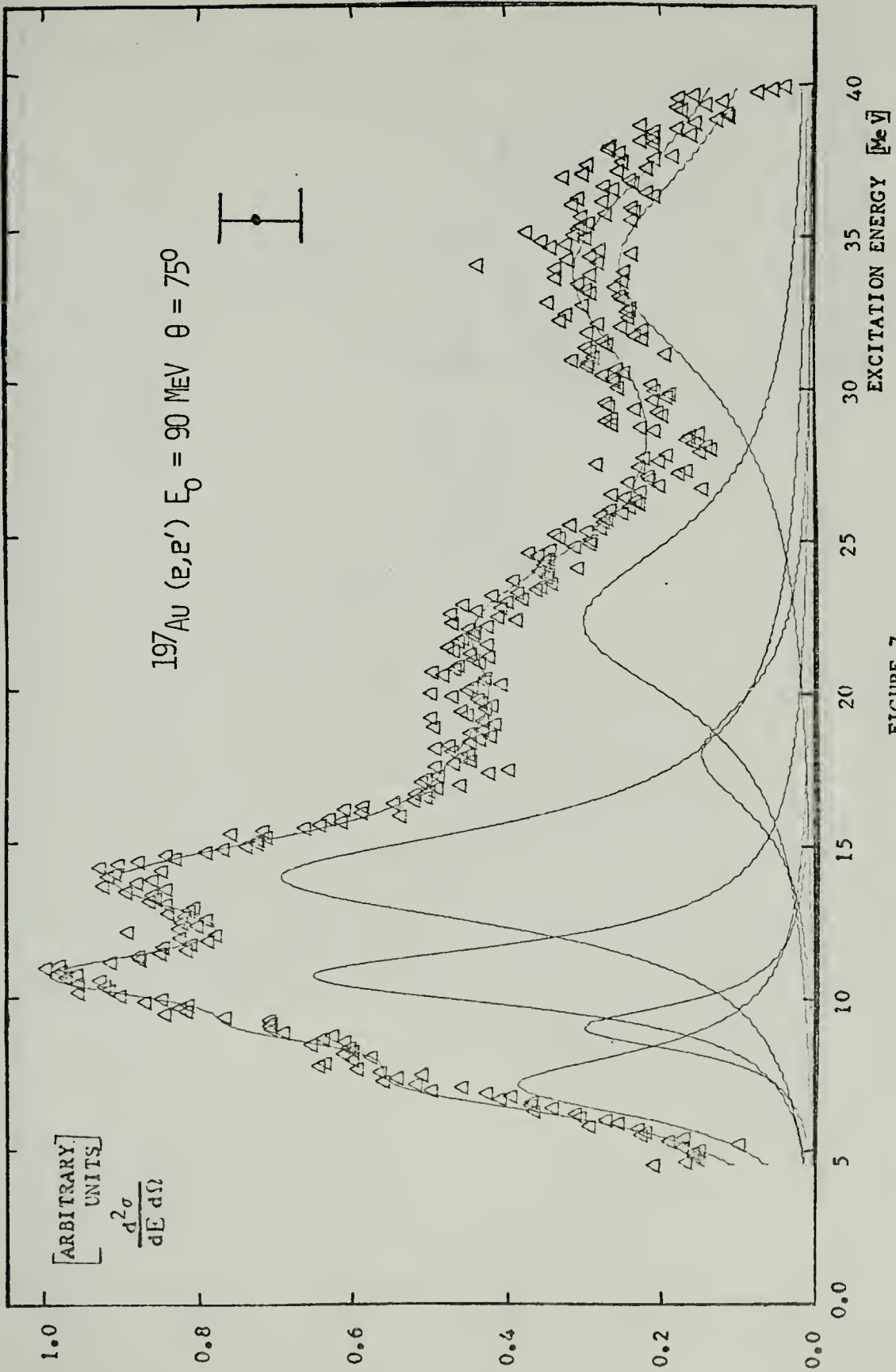


FIGURE 6







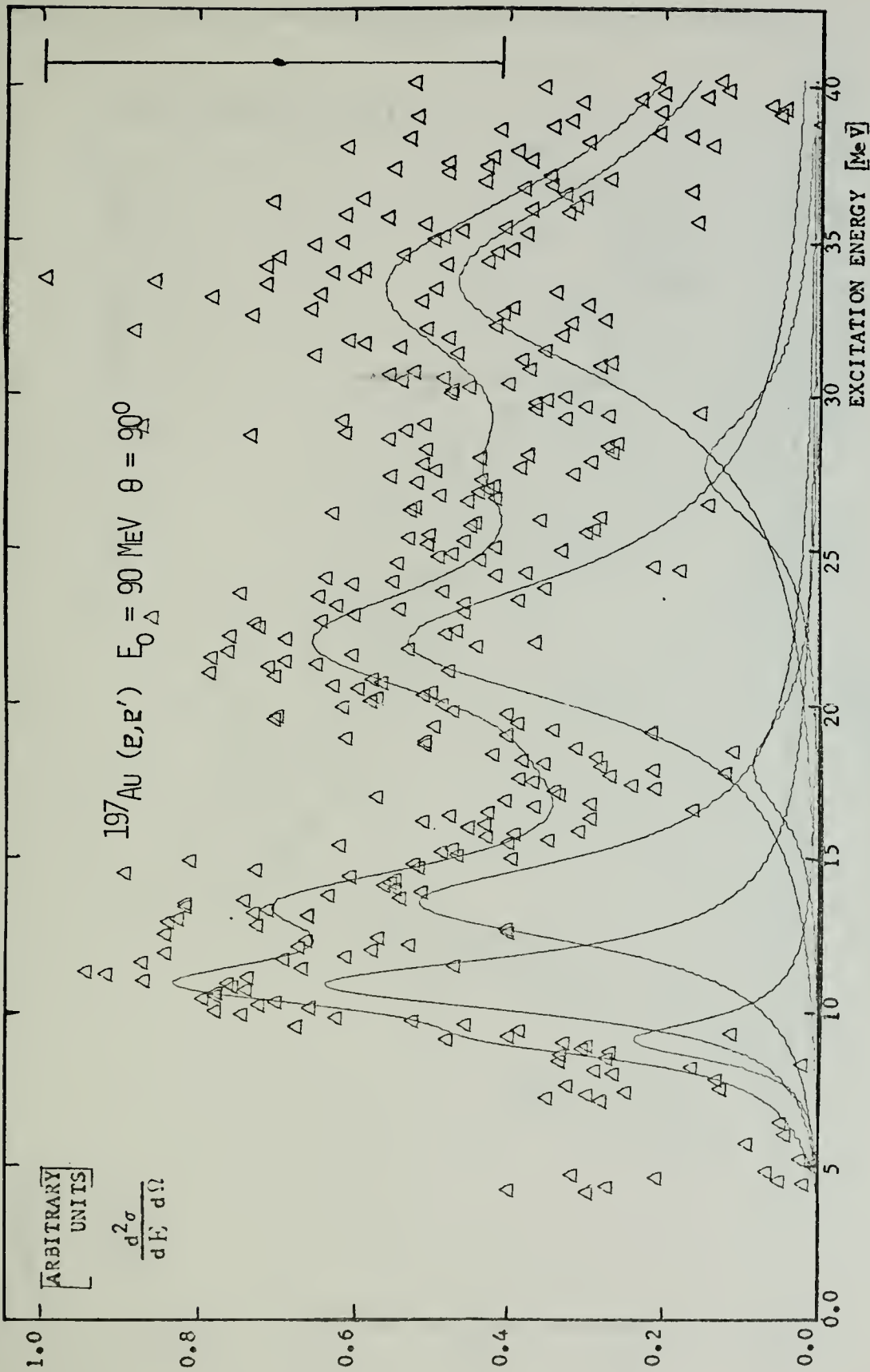


FIGURE 8



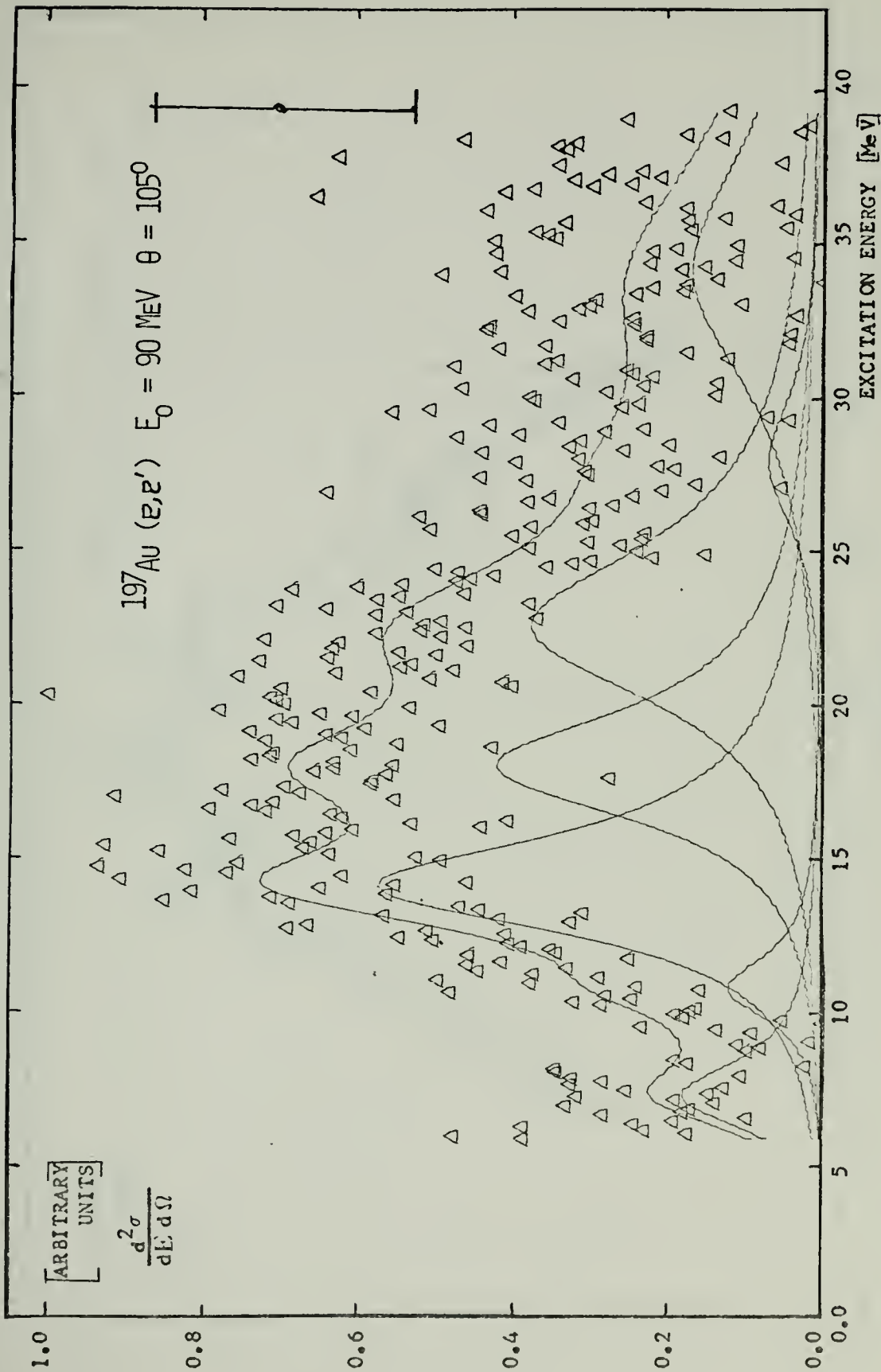


FIGURE 9





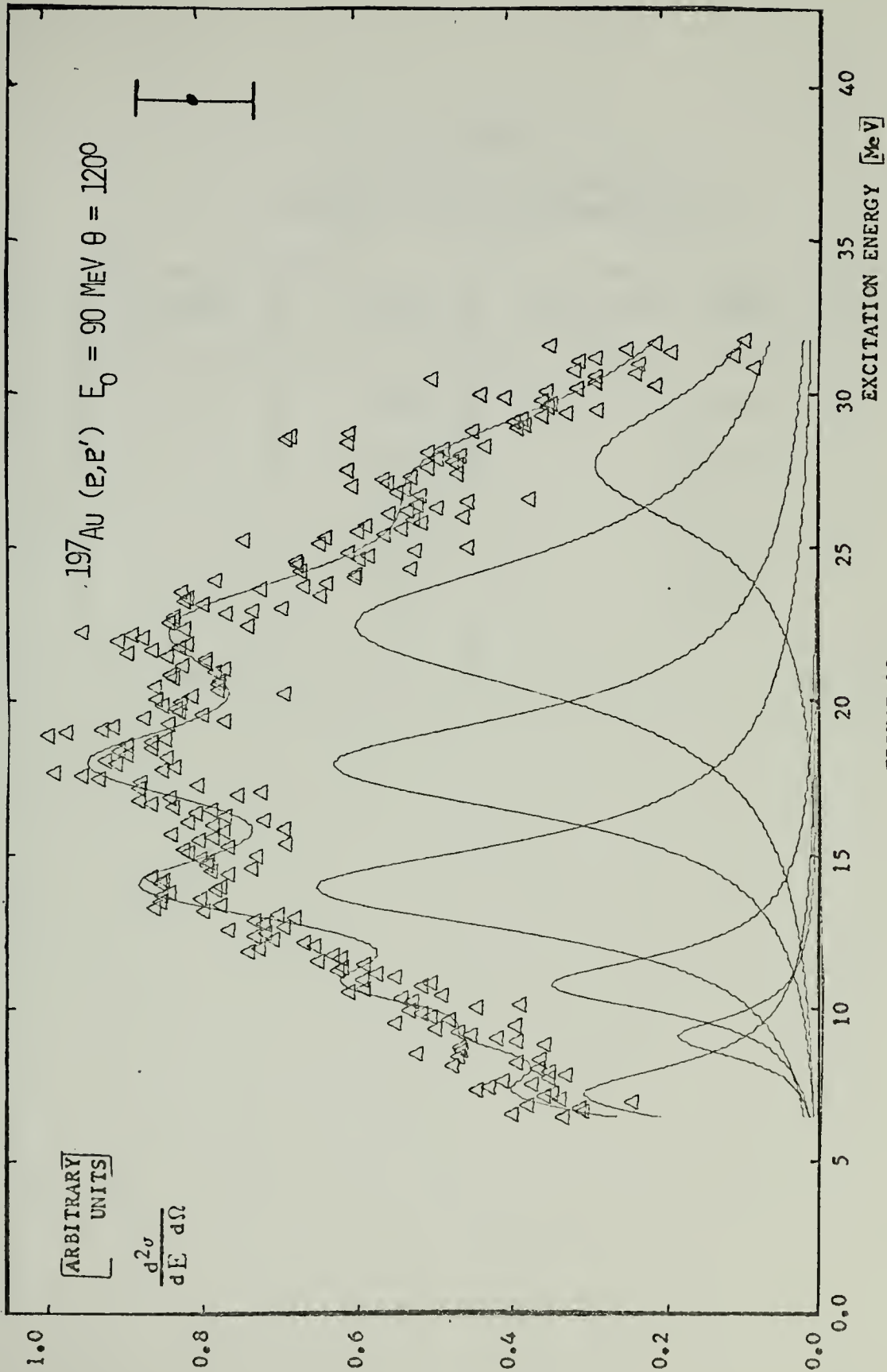


FIGURE 10



TABLE VI

Scale Factors for Converting  
Arbitrary Units to Experimental Units

Figure	Angle	Scale Factor $\frac{\text{fm}^2}{\text{MeV}\cdot\text{sr}}$ / Arb. units
6	$60^\circ$	$9.9 \times 10^{-4}$
7	$75^\circ$	$1.89 \times 10^{-4}$
8	$90^\circ$	$2.99 \times 10^{-5}$
9	$105^\circ$	$1.53 \times 10^{-5}$
10	$120^\circ$	$1.47 \times 10^{-5}$



TABLE VII

Area Ratios and Inelastic Form Factors Squared

$E_x$	Angle	$A_i/A_e$	Total % Error	$\sigma/\sigma_{\text{Mott}}$
7.3	60	$6.37 \times 10^{-3}$	21	$7.75 \times 10^{-4}$
7.3	75	$4.00 \times 10^{-3}$	25	$1.78 \times 10^{-4}$
7.3	90	$9.7 \times 10^{-4}$	--	$2.9 \times 10^{-5}$
7.3	105	$2.70 \times 10^{-3}$	85	$6.05 \times 10^{-5}$
7.3	120	$1.01 \times 10^{-2}$	50	$1.34 \times 10^{-4}$
9.2	60	$3.19 \times 10^{-3}$	60	$3.78 \times 10^{-4}$
9.2	75	$2.20 \times 10^{-3}$	60	$9.78 \times 10^{-5}$
9.2	90	$9.73 \times 10^{-4}$	85	$2.91 \times 10^{-5}$
9.2	105	$1.5 \times 10^{-4}$	--	$3.2 \times 10^{-6}$
9.2	120	$4.34 \times 10^{-3}$	75	$5.73 \times 10^{-5}$
10.8	60	$4.47 \times 10^{-3}$	25	$5.30 \times 10^{-4}$
10.8	75	$6.53 \times 10^{-3}$	20	$2.86 \times 10^{-4}$
10.8	90	$3.58 \times 10^{-3}$	30	$1.07 \times 10^{-4}$
10.8	105	$3.9 \times 10^{-3}$	--	$8.40 \times 10^{-5}$
10.8	120	$1.07 \times 10^{-2}$	35	$1.41 \times 10^{-4}$



TABLE VII Continued

$E_x$	Angle	$A_i/A_c$	Total % Error	$\sigma_{\text{Mott}}\%$
14.0	60	$7.60 \times 10^{-3}$	18	$9.02 \times 10^{-4}$
14.0	75	$1.12 \times 10^{-2}$	16	$4.98 \times 10^{-4}$
14.0	90	$4.42 \times 10^{-3}$	30	$1.32 \times 10^{-4}$
14.0	105	$8.46 \times 10^{-3}$	35	$1.84 \times 10^{-4}$
14.0	120	$3.18 \times 10^{-2}$	20	$4.2 \times 10^{-4}$
18.0	60	$1.82 \times 10^{-3}$	40	$2.16 \times 10^{-4}$
18.0	75	$2.70 \times 10^{-3}$	35	$1.20 \times 10^{-4}$
18.0	90	$1.7 \times 10^{-3}$	--	$5.10 \times 10^{-5}$
18.0	105	$6.91 \times 10^{-3}$	40	$1.50 \times 10^{-4}$
18.0	120	$3.41 \times 10^{-2}$	22	$4.50 \times 10^{-4}$
22.5	60	$4.21 \times 10^{-3}$	30	$5.00 \times 10^{-4}$
22.5	75	$7.83 \times 10^{-3}$	20	$3.47 \times 10^{-4}$
22.5	90	$7.17 \times 10^{-3}$	25	$2.14 \times 10^{-4}$
22.5	105	$8.60 \times 10^{-3}$	45	$6.02 \times 10^{-4}$
33.5	60	$7.66 \times 10^{-3}$	50	$9.09 \times 10^{-4}$
33.5	75	$9.80 \times 10^{-3}$	50	$4.35 \times 10^{-4}$
33.5	90	$9.90 \times 10^{-3}$	60	$2.48 \times 10^{-4}$
33.5	105	$1.5 \times 10^{-2}$	--	$3.2 \times 10^{-4}$





## B. ERROR ANALYSIS

The procedure outlined in the previous section provides values of the ratios of  $A_i/A_e$  but the uncertainties attached to these ratios requires additional considerations. The program NAW computes only statistical errors. There are other errors, both systematic and random which are not strictly statistical. Some of these are the errors resulting from the instrumental effects and are outlined by Bernard and Traverso [Ref. 36].

To obtain a more realistic estimate of the errors in the ratios  $A_i/A_e$ , the non-statistical uncertainties are assumed to be those which arise from the uncertainties in locating the resonance positions and determining their respective widths. These arise from two sources: the overall instrumental uncertainty which is estimated to be 0.15 MeV, and the freedom allowed in the resonance peak positions and widths in fitting these spectra. The total uncertainty of the energy in this work is the sum of the fitting possibilities and instrumental errors.

The Breit-Wigner shape is then analysed to see how these variations in energy effect the area under the resonances. The Breit-Wigner shape is given by

$$\frac{d\sigma(E, E_R, \Gamma)}{dE} = \frac{A}{(E - E_R)^2 + (\Gamma/2)^2} ; \quad \text{III-1}$$



where  $E_R$  is the resonance energy,  $\Gamma$  is the resonance width,  $A$  is the strength parameter, and  $E$  is the excitation energy. The resonance cross section is given by the integral of the Breit-Wigner form

$$\sigma(E_R, \Gamma) = \frac{A\pi}{\Gamma} + \frac{2A \text{TAN}^{-1}(E_R/\Gamma)}{\Gamma} \quad \text{III-2}$$

The variations in position and width introduce an uncertainty in the resonance cross section,

$$\Delta \sigma(E_R, \Gamma) = \left| \frac{\partial \sigma}{\partial E_R} \right| \Delta E_R + \left| \frac{\partial \sigma}{\partial \Gamma} \right| \Delta \Gamma \quad \text{III-3}$$

The absolute values of the partial derivatives are assumed since the peak position and width uncertainties may be correlated.

The fractional error in the resonance cross section is

$$\frac{\Delta \sigma}{\sigma} = \frac{\left[ \frac{2}{\Gamma^2 + E_R^2} \right] \Delta E_R + \left[ \frac{\pi}{\Gamma^2} + \frac{2 \text{TAN}^{-1}(E_R/\Gamma)}{\Gamma^2} + \frac{2E_R}{\Gamma^3 + \Gamma E_R^2} \right] \Delta \Gamma}{\frac{\pi}{\Gamma} + \frac{2 \text{TAN}^{-1}(E_R/\Gamma)}{\Gamma}} \quad \text{III-4}$$

This uncertainty is then added to the statistical error to provide the total error. Table VIII lists the uncertainties from the two sources.



TABLE VIII

Fractional Error Contributions

$E_r$	$\frac{\Delta \sigma}{\sigma}$	Statistical Errors				
		60°	75°	90°	105°	120°
7.2	.16	.05	.08	Upper Limit	.67	.32
9.2	.51	.08	.09	.35	Upper Limit	.27
10.8	.16	.07	.05	.34	Upper Limit	.17
14.0	.11	.07	.06	.20	.25	.10
18.0	.13	.27	.21	.10	.25	.08
22.5	.09	.20	.12	.16	.34	
33.5	.27	.25	.21	.30	Upper Limit	



## IV. DISCUSSION

### A. COLLECTIVITY

Data analysis techniques explained in Chapter III were used to subtract the background and extract the elastic and inelastic spectra of the scattered electrons. A least squares fit was used to evaluate these spectra and resulted in the identification of seven giant resonances at excitation energies of 7.3, 9.2, 10.8, 14.0, 18.0, 22.5, and 33.5 MeV. These resonances have strengths corresponding to an appreciable fraction of the appropriate sum rule as exhibited in Table IX. Additional support to their collective nature is given in Table X where it is seen that the transition strengths are greater than several Weisskopf units. True giant resonances are therefore being observed.

### B. MULTIPOLARITY ASSIGNMENTS

#### 7.3 MeV Resonance

The multipolarity assignment of the resonance at 7.3 MeV is uncertain. Buskirk et. al. [Ref. 18] and Lone et. al. [Ref. 19] have observed this structure in backward angle (e,e') experiments and indicate a transverse excitation. A comparison of the data with calculations using GBROW [Ref. 27] and a code by Drechsel [Ref. 37] is incapable of determining a definite multipolarity assignment. Among possible reasons for this situation are a multipolarity admixture, existence of unknown longitudinal resonances influencing the forward angle data, and experimental difficulties encountered by the high counting rate at forward scattering angles.





## 9.2 MeV Resonance

The analysis of the 9.2 MeV resonance ( $53 A^{-1/3}$ ) favors an E0 multipolarity assignment. This resonance exhausts 38% and 18% respectively of the E0 and E2 energy-weighted sum rules and its strength corresponds to five Weisskopf units when an E2 assignment is made. However, the angular dependence of the experimental form factors favors E0 (Fig. 12). The result would be in agreement with the isoscalar giant monopole resonance predicted at  $56 A^{-1/3}$  MeV [Ref. 21] and at  $58 A^{-1/3}$  MeV [Ref. 22]. Resonances exhibiting the characteristics of E0 or E2 transitions were also seen at  $E_x = 53 A^{-1/3}$  MeV in  $N = 82$  nuclei [Ref. 23] and in  $^{208}\text{Pb}$  [Ref. 24]. An E1 multipolarity assignment is ruled out by comparison with DWBA form factor calculations and furthermore by the fact that no resonance has been seen in  $\gamma$  - absorption measurements [Ref. 25] at this excitation energy.

## 10.8 MeV Resonance

The resonance at 10.8 MeV ( $63 A^{-1/3}$ ) was seen previously in  $^{197}\text{Au}$  [Refs. 3 and 18] and in many other nuclei (see e.g. [Ref. 5]). Generally an assignment of E2 was favored in these cases. In this experiment a transition strength of 14 Weisskopf units is obtained assuming an electric quadrupole excitation. Comparison of the experimental form factors to the angular distribution calculated from the DWBA form factors is inconclusive in distinguishing between an E0 and E2 modes (Fig. 13). However, the experimental strength corresponds to 113% and 68% of the E0 and E2 sum rules, respectively. Together with our results on the 9.2 MeV resonance, the E2 assignment is therefore favored. This resonance is thought to be the isoscalar quadrupole state predicted by the collective vibration model of Bohr and Mottelson [Ref. 38].



#### 14.0 MeV Resonance

The 14 MeV ( $81 A^{-1/3}$ ) giant dipole resonance was measured with monochromatic  $\gamma$ - experiments [Ref. 25] and has also been seen in other (e,e') experiments [Refs. 3 and 18]. For this experiment, form factor comparisons to DWBA calculations show reasonable agreement at forward angles, but exhibit deviations at backward angles. An attempt was made to further explore the nature of the giant dipole transition by assuming both volume and surface oscillations separately in the DWBA calculations (Fig. 14) [Ref. 23]. Transition strengths of six and eight Weisskopf units and 84% and 104% of the sum rule are obtained for the volume and surface oscillation modes respectively. However, neither model adequately describes the observed data. As suggested by Pitthan [Ref. 39], the discrepancy at backward angles might be explained by an electric spin flip contribution to the E1 resonance [Ref. 40]. Using only the forward angle data to extract the reduced matrix elements, the volume oscillation strength exhausts the sum rule (98%), whereas the surface oscillation exceeds the sum rule. Thus a volume oscillation is favored.

#### 18.0 MeV Resonance

A resonance of undetermined multipolarity is observed at 18.0 MeV ( $105 A^{-1/3}$ ) excitation energy. Nagao and Torizuka [Ref. 24] have observed E3 strength at the corresponding energy in  $^{208}\text{Pb}$ , but the state in  $^{197}\text{Au}$  does not conform to this multipolarity. E(0-3) and M (1-2) calculations have been compared to the data (Fig. 15), but no reasonable agreement is found. It appears that the cross-section at this excitation energy contains a mixture of various transitions.



### 22.5 MeV Resonance

Comparison of the experimental results for the 22.5 MeV ( $130A^{-1/3}$ ) resonance with DWBA calculations (Fig. 16) gives support to an E2 assignment. At forward angles the data is indistinguishable between the E0 and E2 transitions, but the  $105^\circ$  measurement is incompatible with the E0 assignment. The Sendai group has reported [Ref. 3] an electric monopole or quadrupole at this excitation energy in  $^{197}\text{Au}$  and also at the corresponding energy in  $^{208}\text{Pb}$  [Ref. 24]. Other multipolarity assignments are unsuccessful in predicting an angular distribution in agreement with the data. The 22.5 MeV state is thought to be the isovector quadrupole state predicted by the collective vibration model of Bohr and Mottelson [Ref. 38].

### 33.5 MeV Resonance

A new resonance, previously unreported in  $^{197}\text{Au}$ , is observed at 33.5 MeV ( $194 A^{-1/3}$ ). This corresponds roughly to the isovector giant monopole resonance predicted by Suzuki [Ref. 22] through calculations based on the collective vibration model of Bohr and Mottelson. Both the E0 and E2 assignments exceed their respective energy-weighted sum rule prediction. However, our line shape fit, while very sensitive to any structure in the spectrum, may tend to over-estimate cross-sections for very wide resonances. For this reason, the B-values may be too large. Comparison of the experimental form factors to the DWBA calculations allow both E0 and E2 multiplicities to be possible. A similar new resonance was observed in  $^{208}\text{Pb}$  at a corresponding energy of 32.5 MeV [Ref. 26] and tentatively given an E0 or E2 assignment. It is to be emphasized that consistent fits for all the spectra ( $^{197}\text{Au}$ ) over the entire range from 5 MeV to 40 MeV excitation energy were only possible with the inclusion of this resonance at 33.5 MeV.



TABLE IX

Ratio of Resonance Strengths to Energy Weighted Sum Rules

$E_r$	Assumed Transition	$\Delta T$	$B(\lambda L) e^{2-fm^{2L}}$ *	$E_x B(\lambda L)$	Sum	$\frac{E_x B(\lambda L)}{\text{Sum}}$
7.3	$M_1$	--	.35±.26	--	--	--
7.3	$M_2$	--	(2.4±.8)x10 <sup>3</sup>	--	--	--
9.2	$E_0$	0	(3.6±1.8)x10 <sup>3</sup>	3.4x10 <sup>4</sup>	9.3x10 <sup>4</sup>	.36
9.2	$E_1$	1	9.2±10.0	85	7.0x10 <sup>2</sup>	.10
9.2	$E_2$	0	(1.5±.4)x10 <sup>3</sup>	1.4x10 <sup>4</sup>	7.4x10 <sup>4</sup>	.19
9.2	$E_2$	1	(1.5±.9)x10 <sup>3</sup>	1.4x10 <sup>4</sup>	1.10x10 <sup>5</sup>	.13
10.8	$E_0$	0	(9.2±1.9)x10 <sup>3</sup>	1x10 <sup>5</sup>	9.3x10 <sup>4</sup>	1.0
10.8	$E_2$	0	(4.7±.8)x10 <sup>3</sup>	5.1x10 <sup>4</sup>	7.4x10 <sup>4</sup>	.68
10.8	$E_2$	1	(4.7±.8)x10 <sup>3</sup>	5.1x10 <sup>4</sup>	1.1x10 <sup>5</sup>	.46

\* For Monopole Matrix Elements are Given





TABLE IX Continued

$E_r$	Assumed Transition	$\Delta T$	$B(\lambda L) \quad e^2\text{-fm}^2L$	$E_x B(\lambda L)$	Sum	$\frac{E_x B(\lambda L)}{\text{Sum}}$
14.0	$E_1 [ \rho_{tr} = \rho_0 ]$	--	$(4.9 \pm .7) \times 10^1$	$6.9 \times 10^2$	$7.0 \times 10^2$	.98
14.0	$E_1 [ \rho_{tr} = \frac{d\rho_0}{dr} ]$	--	$(8.6 \pm 1.1) \times 10^1$	$1.2 \times 10^3$	$7.0 \times 10^2$	.17
18.0	$M_2$	--	$(1.5 \pm .8) \times 10^3$	--	--	--
18.0	$E_0$	0	$(5.0 \pm 1.5) \times 10^3$	$8.9 \times 10^4$	$9.3 \times 10^4$	.46
18.0	$E_2$	0	$(2.7 \pm 1.5) \times 10^5$	$4.8 \times 10^4$	$7.4 \times 10^4$	.65
18.0	$E_2$	1	$(2.7 \pm 1.5) \times 10^3$	$4.8 \times 10^4$	$1.1 \times 10^5$	.44
18.0	$E_3$	0	$(2.0 \pm 1.0) \times 10^5$	$3.6 \times 10^6$	$6.2 \times 10^6$	.58
18.0	$E_3$	1	$(2.0 \pm 1.0) \times 10^5$	$3.6 \times 10^6$	$9.2 \times 10^6$	.39



TABLE IX Continued

$E_r$	Assumed Transition	$\Delta T$	$B(\lambda L) e^2 \text{-fm}^2 L$	$E_x^B(\lambda L)$	Sum	$\frac{E_x^B(\lambda L)}{\text{Sum}}$
22.5	$E_0$	0	$(2 \pm .13) \times 10^3$	$4.3 \times 10^4$	$9.3 \times 10^4$	.46
22.5	$E_2$	0	$(7.1 \pm .5) \times 10^3$	$1.6 \times 10^5$	$7.4 \times 10^4$	2.2
22.5	$E_2$	1	$(7.1 \pm .5) \times 10^3$	$1.6 \times 10^5$	$1.1 \times 10^5$	1.5
22.5	$E_3$	0	$(3.4 \pm .8) \times 10^5$	$7.6 \times 10^6$	$6.2 \times 10^6$	1.2
22.5	$E_3$	1	$(3.4 \pm .8) \times 10^5$	$7.6 \times 10^6$	$9.2 \times 10^6$	.83
55.5	$E_0$	0	$(1.4 \pm .7) \times 10^4$	$8.2 \times 10^5$	$9.3 \times 10^4$	5.0
55.5	$E_2$	0	$(6.5 \pm 2.2) \times 10^3$	$2.1 \times 10^5$	$7.4 \times 10^4$	2.8
55.5	$E_2$	1	$(6.5 \pm 2.2) \times 10^4$	$2.1 \times 10^5$	$1.1 \times 10^5$	1.9



TABLE X

Transitions Strengths in Weisskopf Units

$E_r$	Transition	$B(\lambda L) [e^2 \cdot \text{fm}^{2L}]$	$B_{\text{sp}} (EL/ML)/e^2$	$B(\lambda L)/B_{\text{sp}}$
7.3	$M_1$	$.35 \pm .26$	.06	6
7.3	$M_2$	$(2.4 \pm .8) \times 10^3$	3.1	770
9.2	$E_0$	$(3.64 \pm 1.80) \times 10^3$	--	--
9.2	$E_1$	$9.2 \pm 10.0$	6.6	1
9.2	$E_2$	$(1.49 \pm .92) \times 10^5$	$3.4 \times 10^2$	4
10.8	$E_0$	$(9.22 \pm 1.87) \times 10^3$	--	--
10.8	$E_2$	$(4.68 \pm .80) \times 10^5$	$3.4 \times 10^2$	14
14.0	$E_1 [e_{\text{tr}} = \rho_0]$	$(4.9 \pm 7) \times 10^1$	6.6	8
14.0	$E_1 [e_{\text{tr}} = \frac{d\rho_0}{dR}]$	$(8.6 \pm 1.1) \times 10^1$	6.6	13
18.0	$M_2$	$(1.48 \pm .84) \times 10^3$	3.1	480
18.0	$E_0$	$(4.96 \pm 24.4) \times 10^3$	--	--
18.0	$E_2$	$(2.66 \pm 1.51) \times 10^5$	$3.4 \times 10^2$	8
18.0	$E_3$	$(1.97 \pm .09) \times 10^6$	$1.6 \times 10^4$	12



TABLE X Continued

$E_r$	Transition	$B(\lambda L)$ [ $e^2 \cdot \text{fm}^{2L}$ ]	$B_{sp} (EL/ML)/c^2$	$B(\lambda L)/B_{sp}$
22.5	$E_0$	$(1.90_{-}^{+}.13) \times 10^3$	--	--
22.5	$E_2$	$(7.10_{-}^{+}.48) \times 10^3$	$3.4 \times 10^2$	22
22.5	$E_3$	$(3.39_{-}^{+}.78) \times 10^5$	$1.63 \times 10^4$	21
33.5	$E_0$	$(1.40_{-}^{+}.72) \times 10^4$	--	--
33.5	$E_2$	$(6.25_{-}^{+}2.15) \times 10^3$	$3.4 \times 10^2$	19





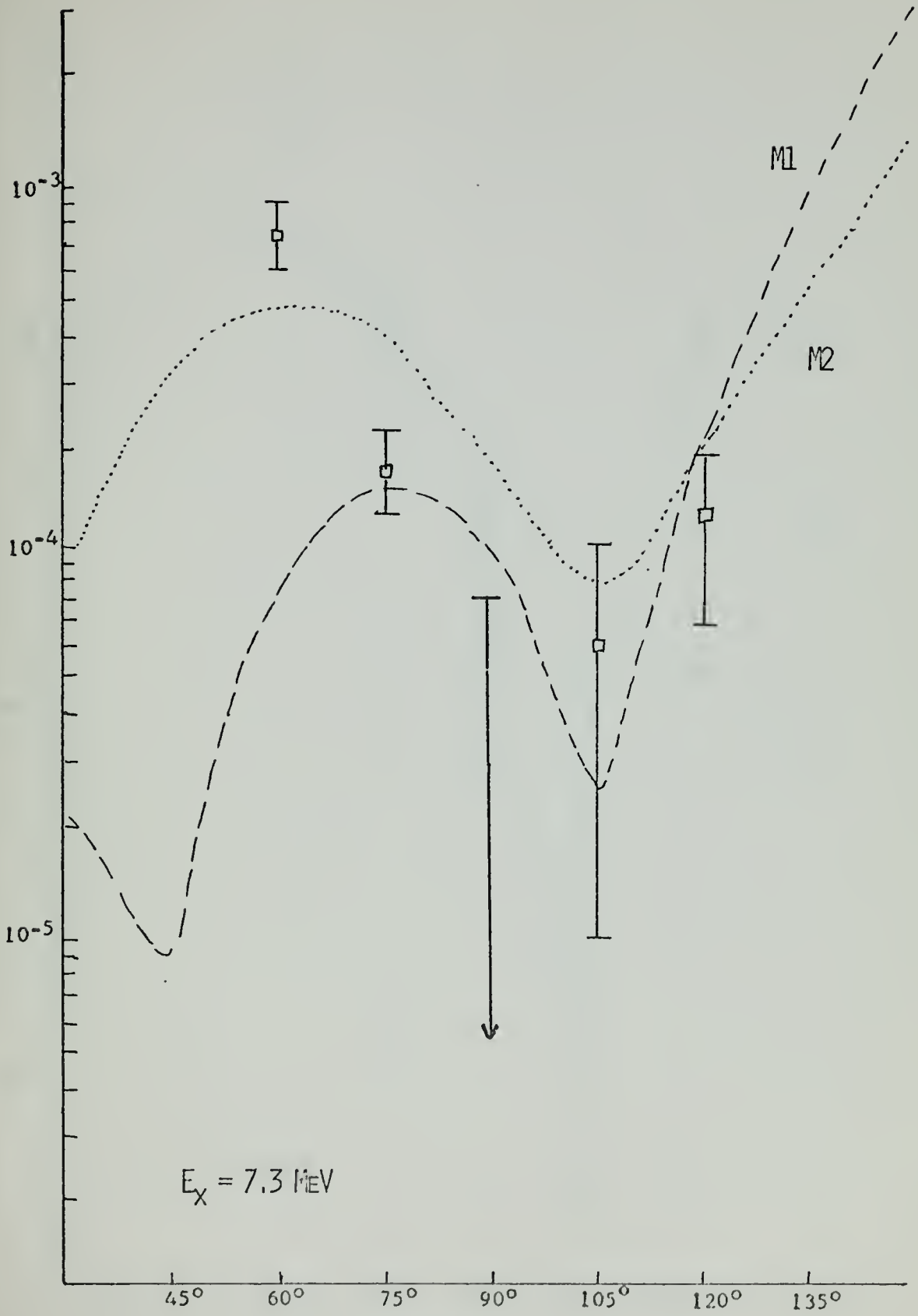


FIGURE 11



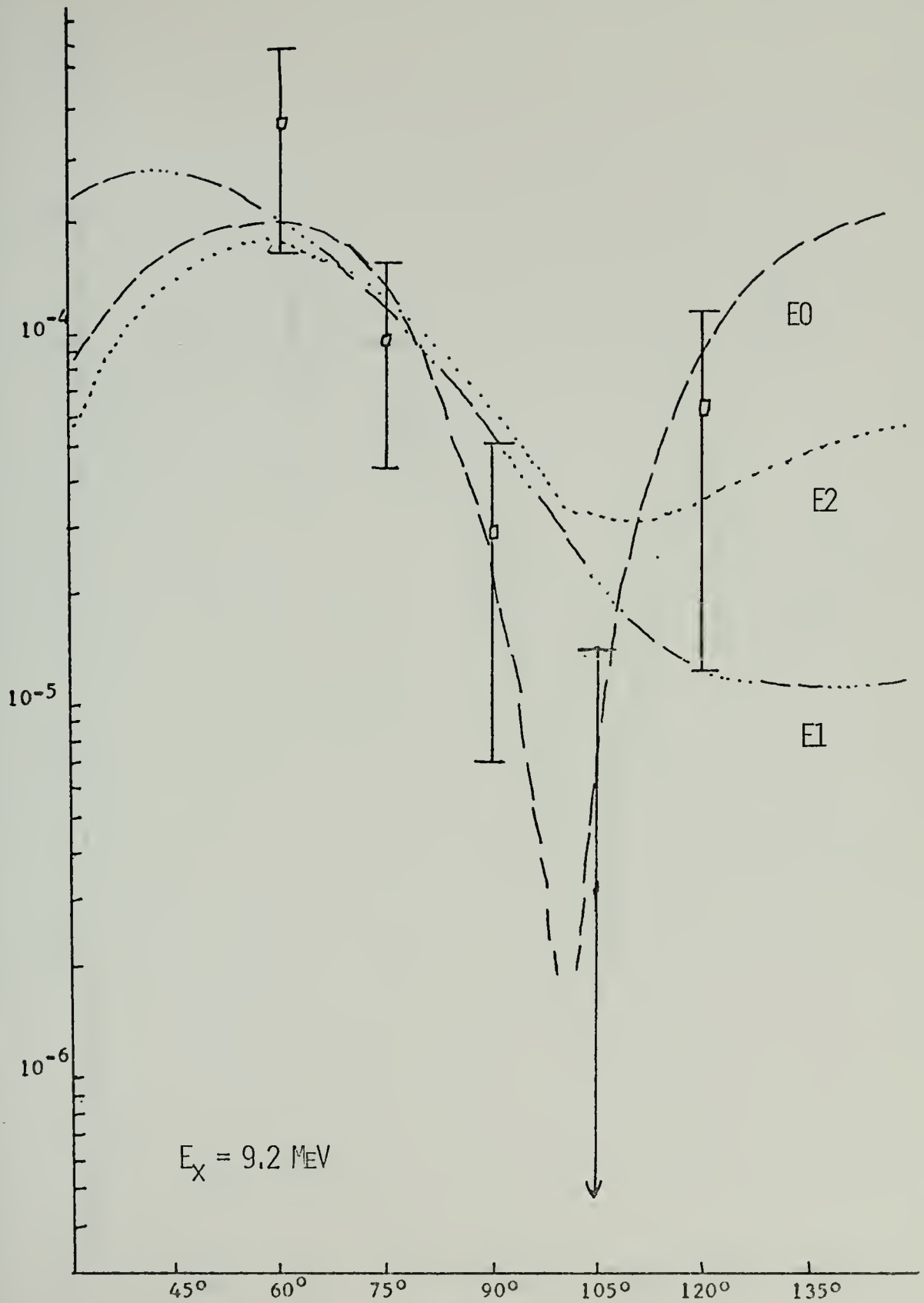


FIGURE 12



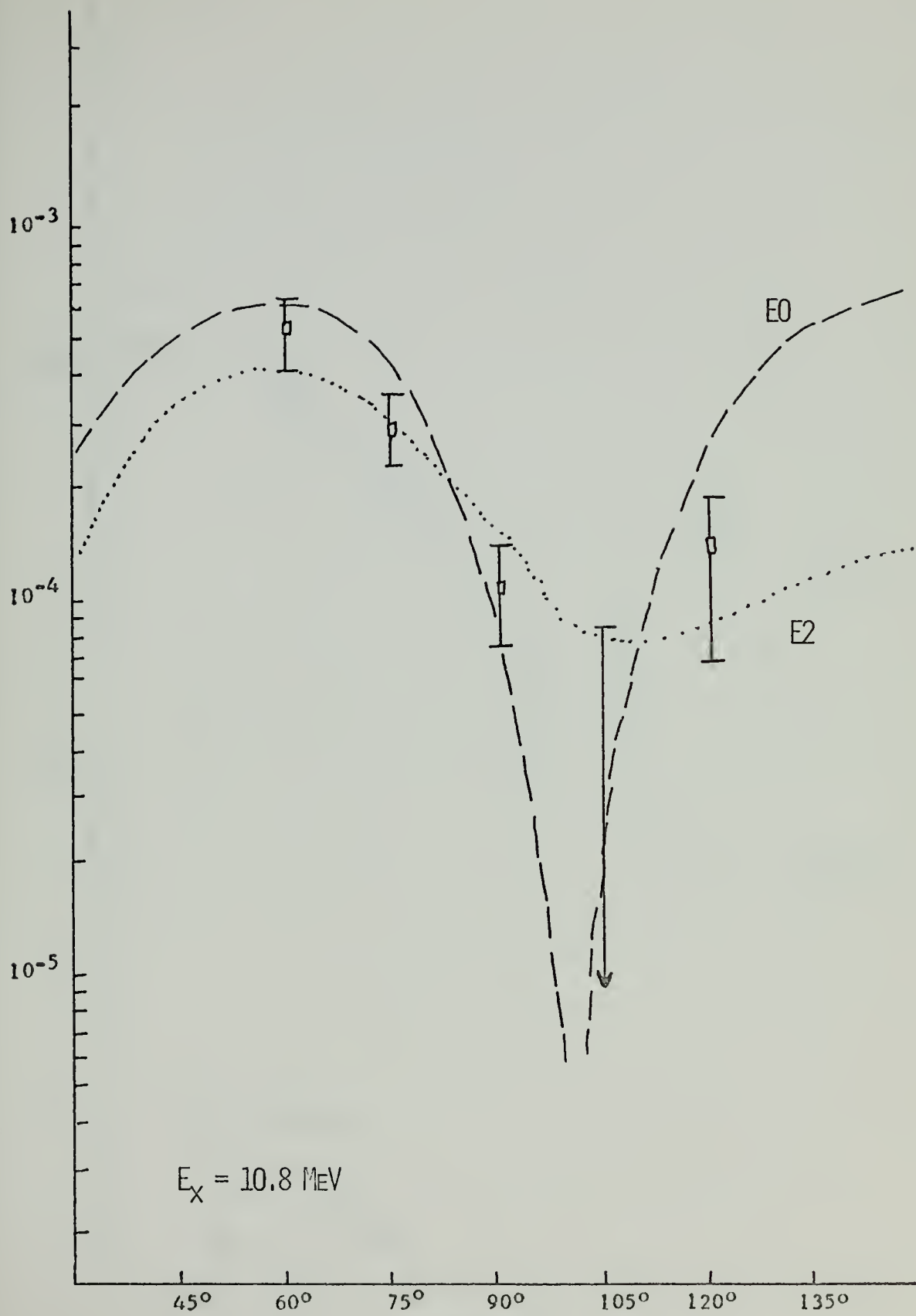


FIGURE 13



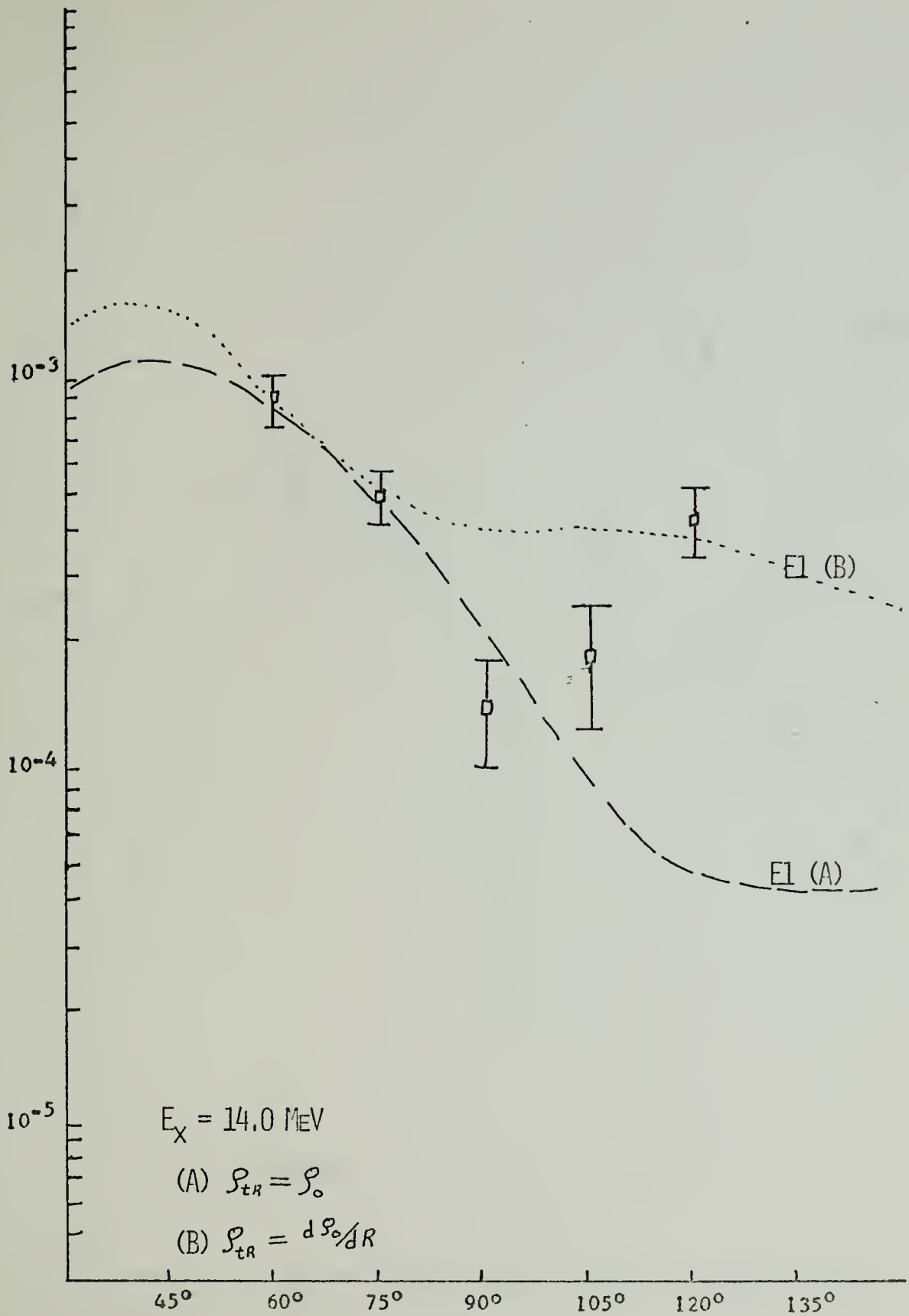


FIGURE 14





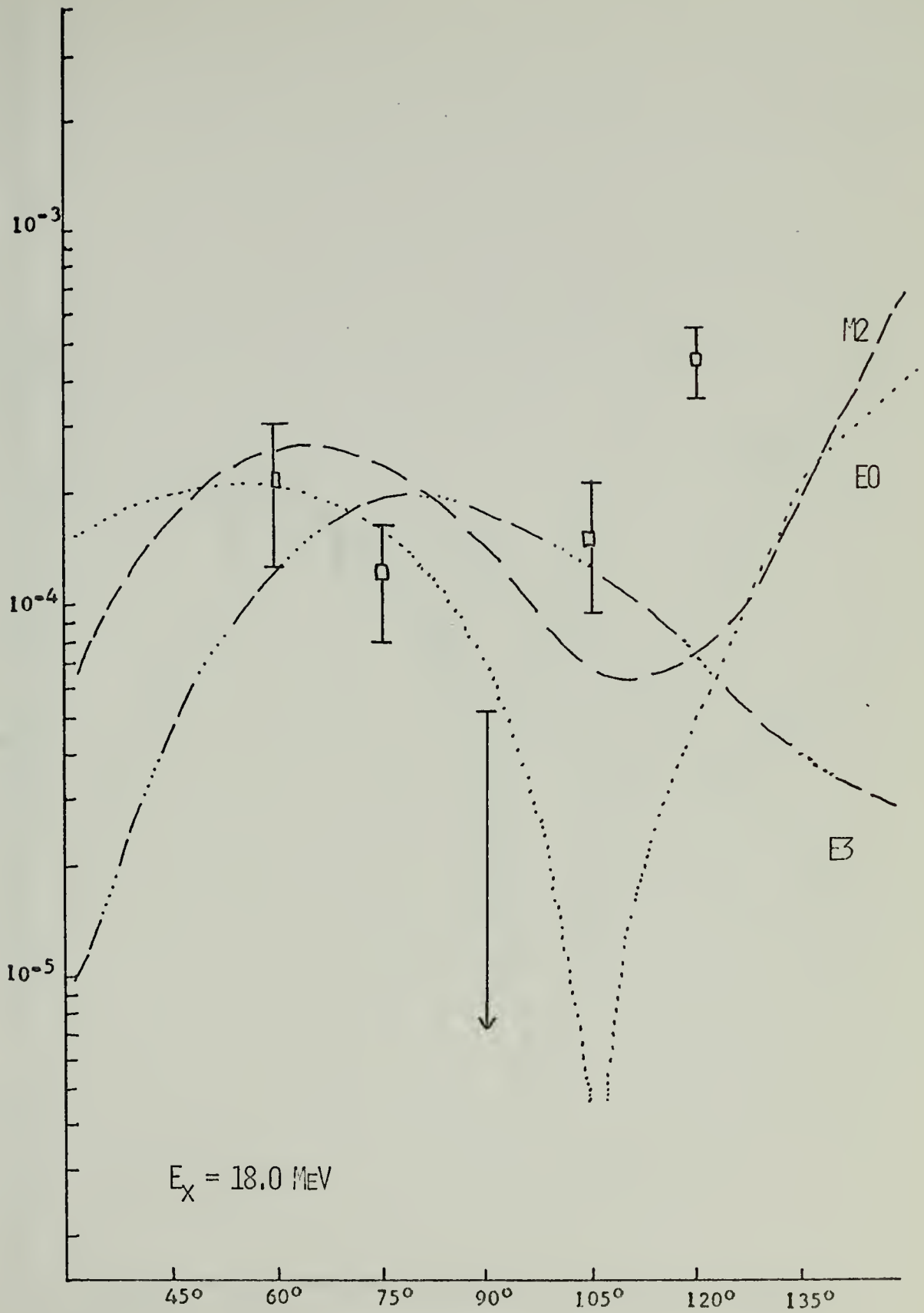


FIGURE 15



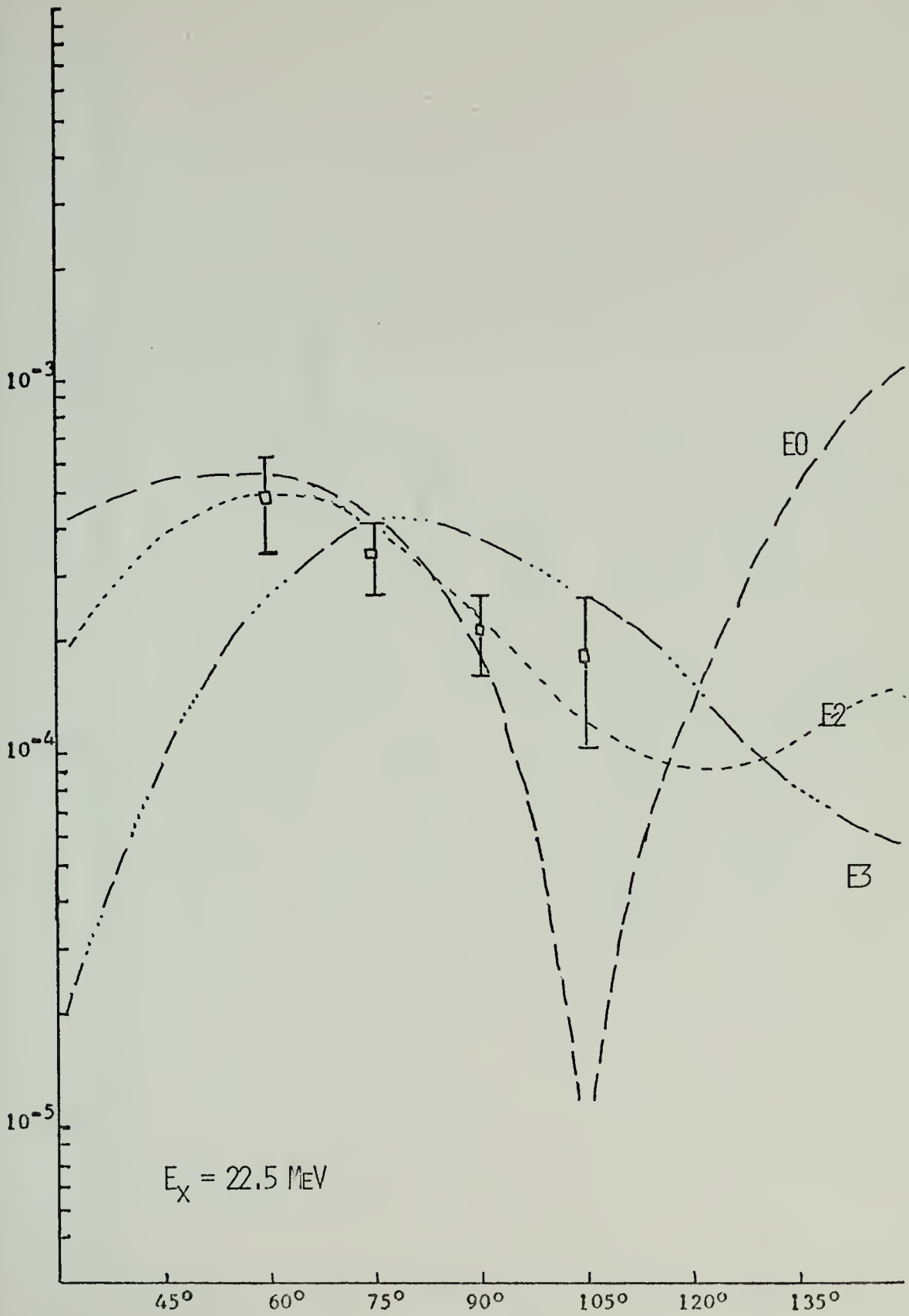


FIGURE 16



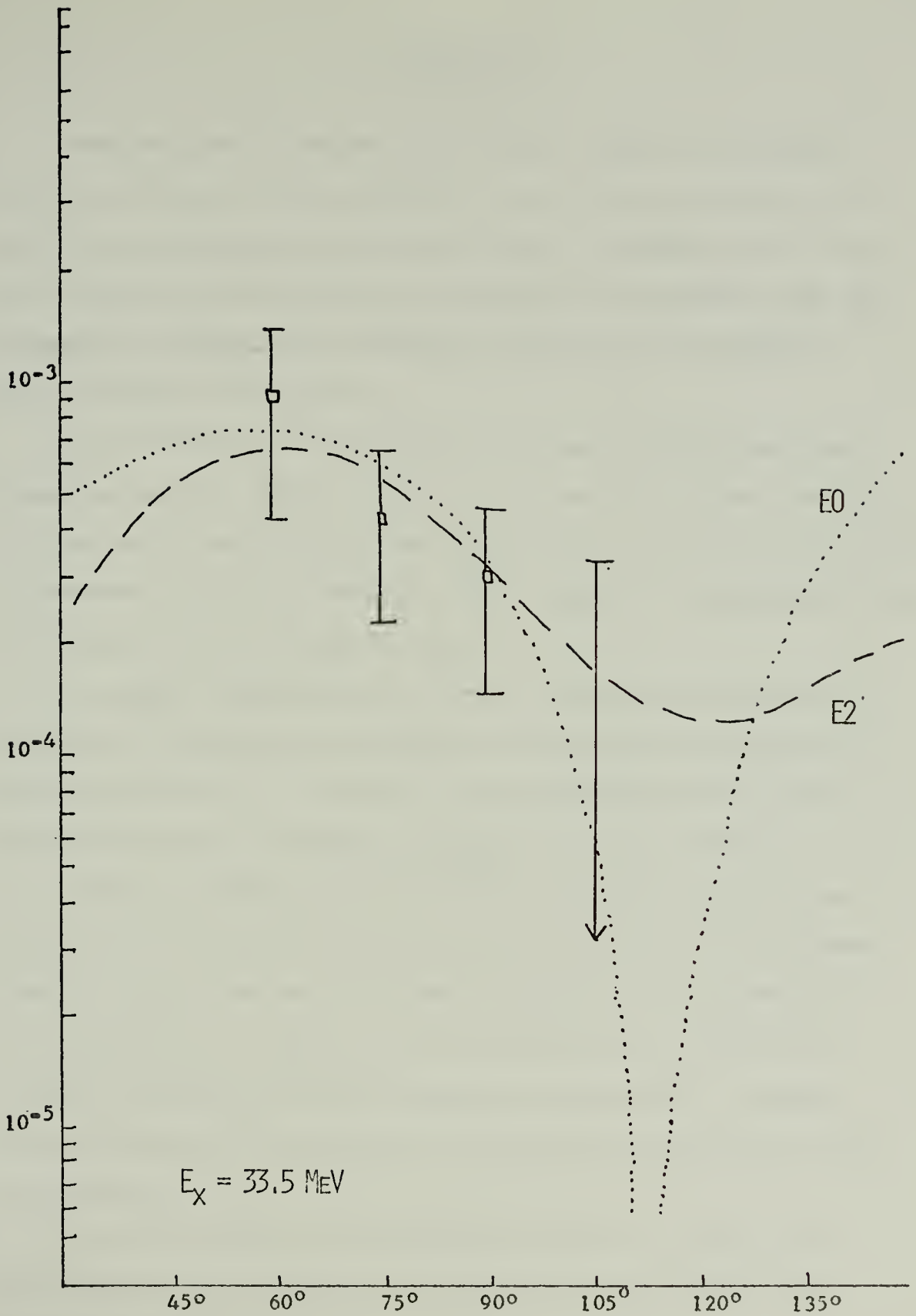


FIGURE 17



## V. CONCLUSIONS

Giant multipole resonances in  $^{197}\text{Au}$  were studied with inelastic electron scattering. Resonances at 7.3, 9.2, 10.8, 14.0, 18.0, 22.5, and 33.5 MeV excitation energy were observed. Of these, the 9.2, 18.0, 22.5, and 33.5 MeV states have not previously been reported in the open literature. Multipolarity assignments of these states were made as was discussed in section IV-B.

The 14.0 MeV state is identified as the much studied giant dipole resonance, see e.g. [Refs. 3, 5, & 25]. Our analysis suggests that this E1 state may be described by a volume oscillation of the charge distribution. Furthermore, the data is compatible with a suspected transverse spin flip contribution to the E1 cross section.

The angular distributions of the cross-sections for the 10.8 MeV and 22.5 MeV resonances are compatible with previous E2 multipolarity assignments [Ref. 3]. In addition, this investigation adds evidence for the exclusion of a monopole assignment for the 10.8 MeV state.

A state was found at 18 MeV in  $^{197}\text{Au}$ . In  $^{208}\text{Pb}$  a resonance was seen [Ref. 26] and cross-section analysis of another (e,e') experiment found E3 strength between 16 MeV and 22 MeV [Ref. 24]. The analysis of the experimental results for this energy region in  $^{197}\text{Au}$  is not capable of assigning a definite transition multipolarity. The data suggests admixtures of longitudinal and transverse contributions to the cross-section.

The 7.3 MeV transition is known to be magnetic. However, this investigation, like others [Ref. 19], was not able to determine the transition mode.





Theoretical investigation concerning the existence and location of the giant monopole resonance is found in the literature [Refs. 21 and 22]. New resonances, discovered in this experiment at 9.2 MeV and 33.5 MeV, have angular distributions in agreement with monopole assignments. Their excitation energies are comparable to theoretical predictions [Ref. 22] of the location of the isoscalar and isovector monopole states. For the lower state, an E0 transition is favored over an E2 assignment. However, for the 33.5 MeV state, the transition multipolarity is inconclusive. The excitation energies and transition multiplicities are presented in Table XI.



TABLE XI  
Multipolarity Assignments

Excitation Energy (MeV)	Transition Multipolarity	Comments
7.3	undetermined	probable magnetic, mixture
9.2	E0	Isoscalar
10.8	E2	Isoscalar
14.0	E1	Isovector, volume oscillation
18.0	undetermined	
22.5	E2	Isovector
33.5	E0 or E2	— Isovector

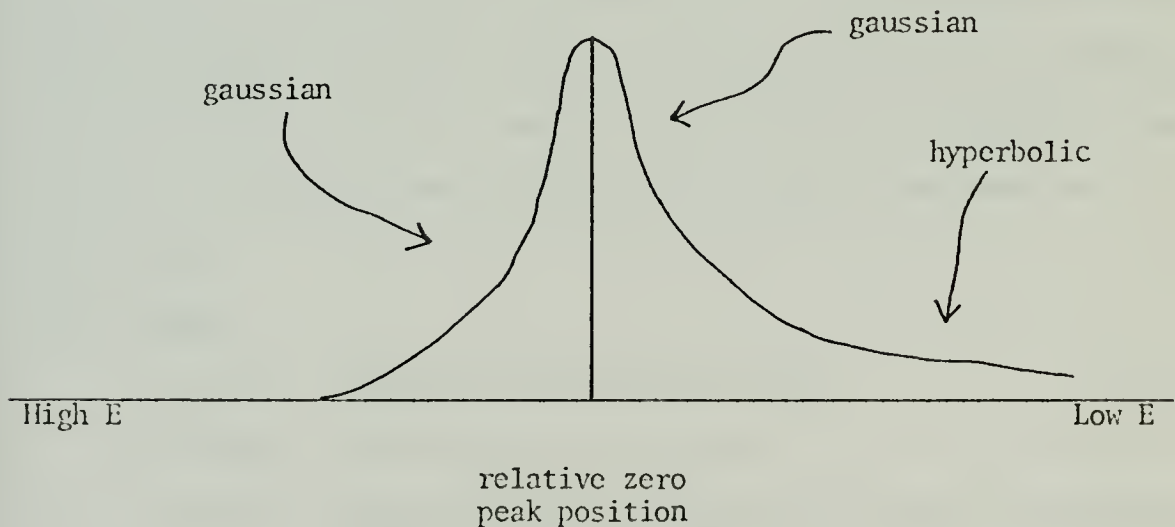


## APPENDIX A

### PROGRAM NAW, SPECTRUM FITTING CODE

#### A. GENERAL FEATURES

Naw is a Fortran program developed at Darmstadt [Refs. 41 & 43] which uses analytical expressions to approximate the inelastic and elastic line shapes and radiation tail to best fit the experimental spectra. The elastic line shape is assumed to be two half gaussians of differing half widths matched to a hyperbolic shape approximating the radiation tail. Work done by the experimenters at the Institut Für Technische Kernphysik at Darmstadt have found this parameterization to be an adequate estimate of the shape of the elastic peak (see Figure 18) [Ref. 41].



ASSUMED ELASTIC SPECTRUM SHAPE  
FIGURE 18

The elastic radiation tail results from processes occurring during elastic scattering that produce lower energy electrons. The radiation tail is calculated from the theory [Ref. 43] at several designated energies



and these theoretical values are fitted with a polynomial of the fourth degree. The total background, which includes the radiation tail, is then simultaneously fit with a series of Breit-Wigner line shapes representing the nuclear excitation structure to reproduce the experimental spectrum. The positions and widths of the resonances are free parameters to be adjusted by best fit or to be assigned from known values. The areas under these resonances are then calculated from which the ratios  $A_i/A_e$  (inelastic peak area divided by the elastic peak area) are determined. The values  $A_i/A_e$  are then multiplied by the calculated scattering cross section resulting in the experimental inelastic cross section for each of the resonances.

## B. CONTROL CARDS

NAW uses the output from program KIK as the input data. Since NAW's first use at the NPS IBM 360/67 [Ref.1] it has undergone various modifications. The present form is reproduced on the following pages, along with the control information. NAW accepts 200 elastic peak data points, 500 inelastic data points and 25 inelastic resonances per spectrum.

### NAW CONTROL

- |    |                 |   |          |
|----|-----------------|---|----------|
| 1. | TEXT            | Alphanumeric identifying remarks  | (20A4)   |
| 2. | NORM,           | PN <sub>2</sub> , PN <sub>3</sub> , PN <sub>4</sub> , PN <sub>5</sub> , PN <sub>6</sub> , PN <sub>7</sub>               | (7F10.5) |
|    | NORM            | Spectrometer energy normalization factor to change the original energy units (OEU) to MeV: [MeV] = NORM*[OEU]           |          |
|    | PN <sub>i</sub> | i=2-7, Starting values for the various parameters being used to fit the elastic spectrum with an analytical expression. |          |
|    | PN <sub>2</sub> | Half-width of a gaussian to fit the high energy side of the elastic peak.   |          |





- PN<sub>3</sub> Half-width of a gaussian to fit the low energy side of the elastic peak.
- PN<sub>4</sub> Matching point for the low energy gaussian and the hyperbolic segment. PN<sub>4</sub> is specified in MeV from the position of the zero point and is a negative number.
- PN<sub>5</sub> Dummy.
- PN<sub>6</sub> PN<sub>6</sub> ≠ 0, Height of a second elastic resonance peak if necessary.
- PN<sub>7</sub> Position of a second elastic resonance peak if necessary, only with PN<sub>6</sub> ≠ 0.
3. EE, THETA, D, UNT, (4F10.5)
- EE Energy of the elastic peak in MeV. EE is used as a starting value in the elastic fit.
- THETA Electron scattering angle in degrees.
- D Target thickness in g/cm<sup>2</sup>.
- UNT Constant background at the high energy end of the elastic resonance spectrum in counts/μ.c.
4. EEL<sub>u</sub>, EEL<sub>1</sub>, GSSW<sub>u</sub>, GSSW<sub>1</sub>, DSSW, E<sub>ref</sub> (6F10.5)
- EEL<sub>u</sub> High energy limit for fitting the elastic peak in OEU.
- EEL<sub>1</sub> Low energy limit for fitting the elastic peak in OEU.
- GSSW<sub>u</sub> Upper limit for calculating the radiation tail in OEU.
- GSSW<sub>1</sub> Lower limit for calculating the radiation tail in OEU.
- DSSW Energy step to be used in calculating the radiation tail in OEU.
- E<sub>ref</sub> See card 7, L<sub>2</sub>. If E<sub>ref</sub> = 0, E<sub>ref</sub> is calculated internally.
5. E<sub>u</sub>, E<sub>1</sub> (2F10.5)
- E<sub>u</sub> High electron energy limit for fitting the inelastic spectrum in excitation energy (MeV).
- E<sub>1</sub> Low electron energy limit for fitting the inelastic spectrum in excitation energy in (MeV).



6.  $GR_u^1, GR_1^1, GR_u^2, GR_1^2, GR_u^3, GR_1^3$  (6F10.5)

Regions where it is expected that only background contributes to the spectrum; i.e., between resonances in the inelastic spectrum. There must be at least one set of  $GR_u^1, GR_1^1$  input, with a maximum of three. It is recommended to set  $GR_u^1 = GSSW_e$ ,  $GR_1^1 = GSSW_u$  for giant resonances, energy in MeV.

7.  $L_i$   $i=1$  to 22 (22I3)

- $L_i$
- 1 A continuous curve consisting of two gaussians and a segment of a hyperbolic is used for both elastic and inelastic fits.
  - 2 The elastic resonance will be fit as if  $L_1=1$  and the inelastic resonance will be fit using the Breit-Wigner formula. Each curve will be multiplied by a straight line with an equation  $Y=1 + P_i(E_x - E_x^{(i)})$ .  $P_i$  will be specified on card #10,  $E_x$  is the excitation energy, and  $E_x^{(i)}$  is the position of the maximum as specified on card #8.
  - 3 The elastic resonance will be fit as if  $L_1=2$  except that NAW will determine the values  $P_i$  and card #10 will be left blank.

$L_2$  Parameters which choose the analytical expression to fit the background due to the radiative and experimental processes,  $L_2$  is a two digit number, the units give the order of a polynomial e.g. unit=3,  $P_1 + P_2A + P_3A^2$ , the  $P_i$ 's are free fitted parameters. The decades choose the function to be multiplied and the argument to be used.

$$\begin{aligned}
 1 \leq L_2 < 10 & \quad P_1 + P_2A + P_3A^2 + \dots + P_{L_2+1} \cdot \text{FSSW} \\
 10 < L_2 < 20 & \quad P_1 + P_2A + P_3A^2 + \dots + \text{FSSW} \\
 20 < L_2 < 30 & \quad \left\{ P_1 + P_2A + P_3A^2 + \dots + P_{L_2+1} \cdot \text{FSSW} \right\} \exp \left[ P_{L_2+2} \left( \frac{E - E_c}{E} \right) \right] \\
 30 < L_2 < 40 & \quad \left\{ P_1 + P_2A + P_3A^2 + \dots + \text{FSSW} \right\} \exp \left[ P_{L_2+1} \left( \frac{E - E_c}{E} \right) \right]
 \end{aligned}$$

FSSW is the calculated radiative tail,  $E$  is the energy of the outgoing electron,  $E_c$  is the energy of the elastic peak.  $A = (E - E_{\text{ref}})$  for  $L_2 < 50$ ,  $E_{\text{ref}}$  is read in on card #7. If  $50 < L_2 < 90$  the analytical form corresponding to  $L_2 = L_2 - 50$  is chosen, but  $A$  is replaced by  $1/A$ .

$L_3$  The number of resonances with fixed energies. If  $L_3 < 0$  then the resonance will be kept fixed relative to  $E_x^{(1)}$ , but  $E_x^{(1)}$  will be fitted. See card #8.

$L_4$  The number of resonances with energies to be fitted. Note that  $|L_3| + L_4$  must not exceed 25.



- $L_5$  The number of lines widths to be fitted.  $>0$  E.g.  $=3$ , means that three line widths are free and in fact are  $E(1)$ ,  $E(2)$ , and  $E(3)$ . If there are more than three resonances,  $E_x(3)$  through  $E_x(N)$  will all be fitted with the same line widths.
- $L_6$  For  $L_1 \geq 2$  the first  $L_6$  peaks will be fitted with Breit-Wigner forms, the remainders with gaussian shapes. For  $L_6=0$ ,  $L_6 = |L_3| + L_4$
- $L_7$  Target geometry:  
 $=0$  Transmission -normal thickness entered on card #3  
 $=1$  Reflection -normal thickness entered on card #3  
 $=2$  Effective thickness entered on card #3
- $L_8$  Output control:  
 $=0$  Maximum output  
 $=1$  Minimal output-no tables and no intermediate calculations  
 $=-1$  Limited output-no tables
- $L_9$  Plotting control:  
 $=0$  Results are not plotted.  
 $=2$  Results are plotted.
- $L_{10}$  Repetitions of NAW procedure.  
 $=0$  No repetitions are performed.  
 $>0$  In this case it is possible to evaluate the data with new parameters. The new parameters will be entered on additional cards placed after the control cards.
- $L_{11}$ - $L_{14}$  Inelastic spectrum fit iteration control. Only override to the default options need be entered.
- $L_{11}$  Number of iterations where all energies are fixed, default  $L_{11}=2$ .
- $L_{12}$  Maximum number of iterations; default is  $L_{12}=30$ .
- $L_{13}$  Minimum number of iterations; default is  $L_{13}=5$ .
- $L_{14}$  Number of iterations where width are held constant in case  $\chi^2_{\text{actual}} > 3\chi^2_{\text{theory}}$ .
- $L_{15}$   $=0$  The radiation tail will be calculated using the Ginsber's method [Ref. 43].  
 $<0$  The radiation tail will be calculated using Schiff's method [Ref. 44].



$L_{16}$ - $L_{18}$  Radiation tail fit control. Only override options need be entered.

$L_{16}$  Analytical representation of the theoretically calculated radiation tail; default is  $L_{16}=1$ .

$$=1 \quad SSW=P_1 + P_2A + P_3A^2 + P_4A^3 + P_5A^4$$

$$=2 \quad SSW=P_1 + P_2A + P_3A^2 + P_4E + P_5E^2$$

$$=3 \quad SSW=P_1 + P_2A + P_3E + P_4E^2 + P_5E^3$$

$$=4 \quad SSW=P_1 + P_2E + P_3E^2 + P_4E^3 + P_5E^4$$

$$=5 \quad SSW=P_1 + P_2U + P_3U^2 + P_4U^3 + P_5U^4$$

$P_i$  will be fitted internally by NAW

$$A \equiv 1/(E_{c1} - E)$$

$$U = E - E_b$$

$$E_b = GSSW_1 + (GSSW_u - GSSW_1)/L_{18}$$

$\langle 0$  the radiation tail parameters  $P_i$  and  $E_b$  will be specified on card #15.

$L_{17}$  The number of parameters for the radiation tail fit ( $3 \leq L_{17} \leq 5$ ); default is  $L_{17} = 4$ .

$L_{18}$  See  $L_{16}$ ; default is  $L_{18} = 2$ .

$L_{19}$   $>0$  Error ellipsoid control values will not be calculated.

$L_{20}$  = 0

$L_{21}$   $>0$  The calculation of the radiation tail will be modified by the factors entered on card #16.

$L_{22}$  =3 Inelastic spectrum results. Card images will be written on file 7 and may be stored on the data cell for further handling.

8.  $E_x(1), E_x(2), \dots, E_x(i)$  (8F10.5)

Resonance excitation energies in MeV; enter those with variable positions first.  $i = |L_3| + L_4 < 25$ .

9.  $V(1), V(2), \dots, V(i)$  (8F10.5)

Resonance line widths,  $i = |L_3| + L_4 < 25$   
If any  $V(i) = 0$ , the width of the elastic line will be assumed.





10. P(1), P(2), ..., P(i) (8F10.5)  
 The  $P_i$  factors when  $L_1 = 2$  on card #7,  $i = |L_3| + L_4 < 25$ .
11. Z, A, XRAD, FIRA, C T (4F10.5, 20X, 2F10.5)  
 Z Atomic number.  
 A Mass number.  
 XRAD Radiation length in  $\text{g}/\text{cm}^2$ .  
 FIRA Elastic form factor for  $E_E$  and THETA of the experiment.  
 C,T The fermi charge distribution parameters. Since these values are used as the parameters for the trapezoidal charge distribution potential used in the radiation tail calculations in Born Approximation, these values must be entered.
12. XMAS, H (2F10.5)  
 XMAS Nuclear mas equivalent in MeV.  
 H Isotope percentage (100% = 1.00).
13.  $A_i$  The  $A_i$  are the coefficients for the calculations of the elastic form factor using a power series expansion on  $q^2$ :  

$$F_{e1}^2 = \sum_{i=1}^8 A_i q^{2(i-1)}$$
 When  $A_1 = 0$  the elastic form factor will be calculated using the Born Approximation.
14.  $U_i$  (8F10.5)  
 $U_i$  may be used to get better starting values for the fit by taking the background parameters from a previous fit. If the  $U_i = 0$ , the starting values are chosen internally. See card #7.
15. P1, P2, P3, P4, P5,  $E_b$  (6F10.5)  
 This card is read only if  $L_{16} < 0$ .  
 P1-P5 Radiation tail parameters, see  $L_{16}$  for the expression.  
 $E_b$  This value is read only when  $L_{16} = -5$ .



16. XSSW, XM, XB,

(3F10.5)

This card is read only if  $L_{21} > 0$ .

XSSW Factor Multiplying the area under the elastic peak,  $A_e$

XM Factor multiplying the Møller term.

XB factor multiplying the B  
The equation for the radiation tail is now given by  
$$SSW = \frac{MAXI}{F^2 \sigma_{Mott}} + XM(\text{Møller term}) + XB(\text{Brems. term}) \cdot (XSSW)A_e$$
  
MAXI are form factors as described by Maximon and Isabelle.

The data are assigned to file #4. The elastic and the inelastic data must be followed by one card image having the data energy step width entered beginning in column 72. Repeated evaluations of the same set of data (i.e.  $L_{10} = 0$ ) must have the following set of cards placed after the 14 to 16 control cards as described above; one set for each repetition to be performed.

- 1)  $L_1 \dots L_{10}$  see card #7
- 2)  $E_u, E_1$  see card #5
- 3)  $E_x(1) \dots, E_x(i)$  see card #8
- 4)  $V(1), \dots, V(i)$  see card #9
- 5)  $P(1), \dots, P(i)$  see card #10



NAW LISTING

```

DIMENSION DUM1(600), LPAR(100), PAR(300)
DIMENSION WORD(20), LP(25), PR(50), FFF(200), BR(25)
DIMENSION BVAR(25), PVAR(25)
COMMON PAR, LPAR, DUM1, WORD, FFF
COMMON /BVAR1/ BVAR, BVAR2/ PVAR
COMMON /SSWD/ACO(25)
COMMON /U/START/UPAR(10)
CALL ERRSET(208,300,-1,1,0,207)
C EINLESEN ZUSATZDATEN LINIE, LORENTZKURVE FUER UNELSTISCHEN FIT
C ZWEITE ELASTISCHE JAN 72 PITTHAN
C EINGEFUEGT TR-440 IM FEBRUAR 1971
C UMGESCHRIEBEN FUER ERHL/4HERHL/
C DATA WIED/4HWIED/
CALL SETIME
PAR(39)=0.0
PAR(40)=0.0
1 READ(5,1000) (WORD(I),I=1,20)
DO 99 I=1,600
DUM1(I)=0.0
IF(I.LE.100) LPAR(I)=0
IF(I.LE.300) PAR(I)=0.
IF(I.LE.200) FFF(I)=0.
IF(I.LE.50) PR(I)=0.
IF(I.GT.25) GO TO 99
ACQ(I)=0.
PVAR(I)=0.
LP(I)=0
BVAR(I)=0.
IF(I.LE.10) UPAR(I)=0.
C CONTINUE
99 READ(5,1001) (PAR(I),I=1,5), (PAR(I),I=62,63), (PAR(I),I=6,12),
I=15,20)
1 READ(5,1002) (LPAR(I),I=1,20), LPAR(40), LPAR(38)
IF(LPAR(1).EQ.2) PAR(64)=1.
IF(LPAR(1).EQ.3) PAR(64)=2.
IF(LPAR(1).GE.2) LPAR(1)=1
IM=IABS(LPAR(3))+LPAR(4)
IF(LPAR(6).EQ.0) LPAR(6)=IM
IM100=IM+100
5 READ(5,1005) (PAR(I),I=101,IM100)
READ(5,1005) (BR(I),I=1,IM)
READ(5,1005) (BVAR(I),I=1,IM)
IF(NWDH.EQ.1) GO TO 980
READ(5,1005) (PAR(I),I=51,58)
READ(5,1005) (PAR(59),PAR(61))
C EINLESEN DER KOEFFIZIENTEN FUER FORMFAKTOR(Q**N)

```

EING0040  
EING0070  
EING0090  
EING0110  
EING0120  
EING0130  
EING0140  
EING0170  
EING0190  
EING0210  
EING0220  
EING0240  
EING0250  
EING0260  
EING0270  
EING0290  
EING0300  
EING0310  
EING0320  
EING0330  
EING0340  
EING0350  
EING0360  
EING0370  
EING0380  
EING0390  
EING0400  
EING0410  
EING0600  
EING0610  
EING0620  
EING0630



```

98 READ(5,1100)(ACO(I),I=1,8)
   READ(5,1002)(LPAR(I),I=41,50)
   READ(5,1100)(UPAR(I),I=1,8)
   IF(LPAR(16).LT.0) READ(5,1007) (PAR(2*N+89),N=1,5),PAR(81)
   IF(LPAR(40).GT.0) READ(5,1008) PAR(90),PAR(89),PAR(88),LPAR(39)
1000 FORMAT(20A4)
1001 FORMAT(7F10.5/5F10.5/6F10.5/7F10.5/6F10.5)
1002 FORMAT(22I3)
1003 FORMAT(7F10.5,I10)
1004 FORMAT(7F10.5,I10)
1005 FORMAT(8F10.5)
1006 DOUBLE PRECISION ABCDE
   IM=IABS(LPAR(3))+LPAR(4)
   LPAR(21)=IM
   LPAR(25)=0
   LWRITE(6,1006) (WORD(I),I=1,18)
1007 FORMAT(1H0,18A4)
1008 FORMAT(6F10.5)
1009 FORMAT(3F10.5,I3)
1100 FORMAT(8F10.5)
C   BERECHNUNG HILFSGROSSEN
   LA=0
   EGO=PAR(6)*PAR(1)
   THETA=PAR(7)*0.0174533
   CU2=COS(THETA/2.)
   SU2=SIN(THETA/2.)
   DEFF=PAR(8)
   IF(LPAR(7).EQ.0) DEFF=DEFF/CU2
   IF(LPAR(7).GT.0) DEFF=DEFF/SU2
   PAR(41)=DEFF
   PAR(46)=DEFF*1.4
   EGO=EGO+PAR(46)/2.
   PAR(47)=1.+2.*EGO*SU2**2/PAR(59)
   EGO=EGO*PAR(47)
   PAR(47)=1.+2.*EGO*SU2**2/PAR(59)
C   AUSGABE DER EINGABEDATEN
9001 FORMAT(3X,20I6)
9000 FORMAT(3X,10F12.3)
C   KATEGORIE, AUSWERTUNG ELASTISCH
   IF(LPAR(20).GT.1) CALL ZETKO
   REWIND 1
   REWIND 3
   IF(LPAR(20).GT.0) CALL ZETKO
   CALL EGB
   IF(LPAR(20).GT.0) CALL ZETKO
   REWIND L
   IF(LPAR(20).GT.0) CALL ZETKO
   ELST

```

```

EING0740
EING0750
EING0760
EING0780
EING0790
EING0800
EING0810
EING0820
EING0830
EING0840
EING0850
EING0860
EING0870
EING0880
EING0890
EING0900
EING0910
EING0930
EING0950
EING0960
EING0970
EING0980
EING0990
EING1000
EING1010
EING1020
EING1030
EING1040
EING1050
EING1060
EING1070
EING1080
EING1110
EING1120
EING1150
EING1160
EING1180
EING1200
EING1230
EING1240
EING1250
EING1260
EING1270
EING1280

```





```

IF(LPAR(20).GT.0) CALL ZETKO
LWD=0
C BER CHNEN EPS,BREITEN
110 EEO=PAR(31)*PAR(1)+PAR(46)/2.
EEO=EEO
PAR(47)=1.+2.*EEO*S2**2/PAR(59)
EEO=EEO*PAR(47)
PAR(47)=1.+2.*EEO*S2**2/PAR(59)
EEO=EEO*PAR(47)
PAR(36)=EEO
112 XBR=(2.6*PAR(31)/(1000.*PAR(34)))**2
DC 120 I=1,IM
PAR(I+150)=PAR(I+100)/PAR(1)*((1.0+PAR(I+100))/2.-PAR(I+150))/PAR(47)
PAR(I+125)=SQRT(1.-XBR*PAR(I+150))/PAR(31)*((2.-PAR(I+150))/PAR(31))
IF(PAR(64).GE.1.) BR(I)=BR(I)/PAR(34)/PAR(1)
IF(BR(I).NE.0.) PAR(I+125)=BR(I)
C CONTINUE
UNELASTISCH, PLOTT, WIEDERHOLUNG
IF(LWD.GT.0) REWIND 3
DC 130 J=1,20
LP(J)=0
130 IF(LPAR(20).GT.0) CALL ZETKO
CALL UNST
IF(LPAR(20).GT.0) CALL ZETKO
WRITE(6,9000)(PAR(I),I=1,300)
WRITE(6,9001)(LPAR(I),I=1,50)
WRITE(6,9000)(BVAR(I),I=1,25)
WRITE(6,9000)(PVAR(I),I=1,25)
CALL AUSG
IF(LPAR(20).GT.0) CALL ZETKO
IF(LPAR(9).GT.0) CALL DRUCK
IF(LPAR(20).GT.0) CALL ZETKO
LWD=1
IF(LPAR(10).EQ.0) GO TO 1
C EINLESEN DER WIEDERHOLUNGSDATEN
READ(5,1002,END=1050)(LP(I),I=1,20)
IF(LP(1).NE.LPAR(1)) LWD=0
READ(5,1001) PAR(13),PAR(14),PR64
IF(PR64.GT.0.) PAR(64)=PR64
NWDH=1
GO TO 5
980 NWDH=0
DC 150 I=1,25
IF(LP(I).EQ.0.AND.I.GT.10) GO TO 140
LPAR(I)=LP(I)
140 IF(PR(I).NE.0.0) PAR(I+100)=PR(I)

```

```

EING1290
EING1300
EING1320
EING1340
EING1350
EING1370
EING1380
EING1390
EING1400
EING1410
EING1420
EING1430
EING1440
EING1470
EING1490
EING1510
EING1520
EING1530
EING1540
EING1550
EING1560
EING1570
EING1580
EING1590
EING1600
EING1610
EING1620
EING1630
EING1640
EING1650
EING1660
EING1670
EING1680
EING1690
EING1700
EING1710
EING1720
EING1730
EING1740
EING1750
EING1760
EING1770
EING1780
EING1790
EING1800

```

-



EING1810  
 EING1820  
 EING1830  
 EING1840  
 EING1850  
 EING1860  
 EING1870  
 EING1880  
 EING1890  
 EING1900  
 EING1910  
 EING1920  
 EING1930  
 EING1940  
 EING1950  
 EING1960  
 EING1970  
 EING1980  
 EING1990  
 EING2000  
 EING2010  
 EING2020  
 EING2030  
 EING2040

```

150 IF (PR(I+25).NE..0) BR(I)=PR(I+25)
CONTINUE
IM=IABS(LP PAR(3))+LP AR(4)
LP AR(21)=IM
IMP=I00+IM
IMB=I25+IM
WORD(I7)=WERHL
WRITE(6,2070) (WORD(I),I=1,18)
WRITE(6,2075) (LP AR(I),I=1,20)
WRITE(6,2080) (PAR(I),I=1,14)
WRITE(6,2085) (BR(I),I=1,IM)
WRITE(6,2090) (BR(I),I=1,IM)
FORMAT(1H1,13H EINGABEDATEN / )
FORMAT(1H0,13HNEUE GRENZEN / )
FORMAT(1H0,13HNEUE GRENZEN 2F10.4)
FORMAT(1H0,19HANREGUNGSENERGIEN /2(10F10.4/))
FORMAT(1H0,19HNEUE BREITEN: /2(10F10.4/))
LP AR(31)=0
REWIND 1
IF (LP AR(20).GT.1) CALL ZETKO
IF (LWD.EQ.0) GO TO 105
GO TO 112
STOP
END

```

EING2140

```

SUBROUTINE ZETKO
DIMENSION PAR(300)
COMMON PAR
CALL GETIME(IET)
Z0=FLOAT(IET)*.000026
Z1=Z0-P AR(40)
PAR(40)=Z0
Z2=PAR(40)/60.0
Z3=INT(Z2)
Z4=AMOD(P AR(40),60.0)
Z1=Z1/60.0
LZ1=INT(Z1)
WRITE(6,1000) LZ2,Z3,LZ1,Z4
FORMAT(1H,5HZEIT 14,5H MIN F 7.3,12H SEC, DIFF. 13,5H MIN
F7.3,5H SEC)
RETURN
END

```

EING2300  
 EING2310  
 EING2330  
 EING2340  
 EING2350

```

SUBROUTINE INPUT(LB)
DIMENSION CO(500),HZ(500),FEH(500),PAR(300)
COMMON PAR

```



```

COMMON/DATA1/HZ,CO,FEH
DO 10 I=1,250
I1=2*I-1
I2=2*I
READ(4,100) (HZ(J),CO(J),FEH(J),J=I1,I2)
CC(I1)=CO(I1)-PAR(9)
CO(I2)=CO(I2)-PAR(9)
IF(HZ(I1).EQ.0.) GO TO 20
CONTINUE
10 IMAX=2*I
20 IF(FEH(IMAX).GT.0.0) PAR(LB+26)=FEH(IMAX)
CALL SORT(HZ,CO,FEH,IMAX)
WRITE(3) I=1,IMAX
DC 30 I=1,IMAX
30 WRITE(3) HZ(I),CO(I),FEH(I)
IF(LB.EQ.0) GO TO 40
PAR(49)=FLOAT(IFIX(50.*HZ(I)))/50.
PAR(50)=FLOAT(IFIX(50.*HZ(IMAX-1)))/50.
40 RETURN
100 FCRMAT(2(F9.4,1X,F15.7,F13.7,2X))
END

```

```

SUBROUTINE EGB
DIMENSION PAR(300),LPAR(100)
COMMON PAR,LPAR
PAR(26)=0.0
PAR(27)=0.0
LB=0
CALL INPUT(LB)
LBI=1
CALL INPUT(LBI)
REWIND 3
RETURN
END

```

```

SUBROUTINE SORT(HZ,CO,FEH,IMAX)
DIMENSION HZ(500),CO(500),FEH(500)
LA=1
1 K=IMAX-1
DO 10 L=LA,K
HERZ=HZ(L)
ZN=CO(L)
ZF=FEH(L)
M=L+1
9 I=M,IMAX
IF(HZ(I).NE.0.) GO TO 4
IMAX=IMAX-1

```

```

EING2360
EING2370
EING2380
EING2390
EING2400
EING2410
EING2420
EING2430
EING2440
EING2450
EING2460
EING2470
EING2480
EING2490
EING2500
EING2510
EING2520
EING2530
EING2540
EING2550

```

```

EING2580
EING2590
EING2600
EING2610
EING2620
EING2630
EING2640
EING2650
EING2660
EING2670
EING2680
EING2690
EING2700

```

```

EING2720
EING2730
EING2740
EING2750
EING2760
EING2770
EING2780
EING2790
EING2800
EING2810
EING2820
EING2830

```



DD = LL+1  
 L2(LL)=HZ(L2)  
 CO(LL)=CO(L2)  
 FEH(LL)=FEH(L2)  
 LA=L  
 GO TO 1  
 IF(HZ(L)-HZ(I)) 5,9,9  
 HZ(L)=HZ(I)  
 HZ(I)=HERZ  
 CO(L)=CO(I)  
 CC(I)=ZW  
 FEH(I)=ZF  
 FEH(L)=ZF  
 HERZ=HZ(L)  
 ZW=CO(L)  
 ZF=FEH(L)  
 CC=CONTINUE  
 RETURN  
 END

3 LL=I, IMAX  
 3 FEH(LL)=FEH(L2)  
 4 GO TO 1  
 5 IF(HZ(L)-HZ(I)) 5,9,9  
 9 CONTINUE  
 10 RETURN  
 END

SUBROUTINE ELST  
 ELST PEAK UND STRAHLUNGSSCHWANZ  
 ELST ISCHER PAR(300), LPAR(100), FREQ(200), ZAEH(200), FEH(200), FF(200)  
 DIMENSION PAR(300), LPAR(100), FREQ(200), ZAEH(200), FEH(200), FF(200),  
 DIMMON PAR, LPAR, FREQ, ZAEH, FEH,  
 IF(LPAR(20).GT.1) CALL ZETKO  
 WORD, FF  
 CCMMCN/DATA1/XX, YY, ERR  
 ZMAX=0.0  
 LPAR(24)=0  
 GAMMA=1000.  
 LEO=0  
 READ(3) IEL, LB  
 LPAR(32)=IEL  
 LPAR(9) I=1, IEL  
 M=IEL-I+1  
 READ(3) HZ, CO, FEHEL  
 ZAEH(M)=CO  
 FREQ(M)=HZ  
 FEH(M)=FEHEL  
 ZMAX=AMAX1(ZMAX, CO) GO TO 8  
 IF(LPAR(1).EQ.0) GO TO 8  
 IF(HZ.GT.PAR(11).OR.HZ.LT.PAR(12)) GO TO 9  
 LEC=LEO+1  
 XX(LEO)=HZ  
 YY(LEO)=CO

3 LL=I, IMAX  
 3 FEH(LL)=FEH(L2)  
 4 GO TO 1  
 5 IF(HZ(L)-HZ(I)) 5,9,9  
 9 CONTINUE  
 10 RETURN  
 END

SUBROUTINE ELST  
 ELST PEAK UND STRAHLUNGSSCHWANZ  
 ELST ISCHER PAR(300), LPAR(100), FREQ(200), ZAEH(200), FEH(200), FF(200),  
 DIMENSION PAR(300), LPAR(100), FREQ(200), ZAEH(200), FEH(200), FF(200),  
 DIMMON PAR, LPAR, FREQ, ZAEH, FEH,  
 IF(LPAR(20).GT.1) CALL ZETKO  
 WORD, FF  
 CCMMCN/DATA1/XX, YY, ERR  
 ZMAX=0.0  
 LPAR(24)=0  
 GAMMA=1000.  
 LEO=0  
 READ(3) IEL, LB  
 LPAR(32)=IEL  
 LPAR(9) I=1, IEL  
 M=IEL-I+1  
 READ(3) HZ, CO, FEHEL  
 ZAEH(M)=CO  
 FREQ(M)=HZ  
 FEH(M)=FEHEL  
 ZMAX=AMAX1(ZMAX, CO) GO TO 8  
 IF(LPAR(1).EQ.0) GO TO 8  
 IF(HZ.GT.PAR(11).OR.HZ.LT.PAR(12)) GO TO 9  
 LEC=LEO+1  
 XX(LEO)=HZ  
 YY(LEO)=CO

EING2840  
 EING2850  
 EING2860  
 EING2870  
 EING2880  
 EING2890  
 EING2900  
 EING2910  
 EING2920  
 EING2930  
 EING2940  
 EING2950  
 EING2960  
 EING2970  
 EING2980  
 EING2990  
 EING3000  
 EING3010  
 EING3020  
 EING3030  
 EING3040  
 EING3050  
 ELST0010  
 ELST0020  
 ELST0060  
 ELST0070  
 ELST0080  
 ELST0090  
 ELST0100  
 ELST0110  
 ELST0120  
 ELST0130  
 ELST0140  
 ELST0150  
 ELST0160  
 ELST0170  
 ELST0180  
 ELST0190  
 ELST0200  
 ELST0210  
 ELST0220  
 ELST0240





ELST0250  
 ELST0260  
 ELST0270  
 ELST0280  
 ELST0650  
 ELST0660  
 ELST0670  
 ELST0680  
 ELST0690  
 ELST0700  
 ELST0710  
 ELST0720  
 ELST0730  
 ELST0740  
 ELST0750  
 ELST0760  
 ELST0770  
 ELST0780  
 ELST0790  
 ELST0820  
 ELST0850  
 ELST0870  
 ELST0880  
 ELST0890  
 ELST0910  
 ELST0920  
 ELST0930  
 ELST0960  
 ELST0970  
 ELST0980  
 ELST1000  
 ELST1020  
 ELST1040  
 ELST1050  
 ELST1060  
 ELST1070  
 ELST1080  
 ELST1090  
 ELST1100  
 ELST1110  
 ELST1120  
 ELST1130  
 ELST1140  
 ELST1150  
 ELST1160  
 ELST1170  
 ELST1180  
 ELST1190

```

ERR(LEO)=FEHEL
CONTINUE
IF(LPAR(20).GT.2) CALL ZETKO
IF(LPAR(1).GT.0) GO TO 91
FIT ELASTISCH
91 P(1)=ZMAX
P(2)=PAR(2)
P(3)=PAR(3)
P(4)=PAR(4)
P(5)=PAR(6)
P(6)=PAR(62)
P(7)=PAR(63)
ME=5
LPPAR(28)=ME
NZEL=5
IF(PAR(62).NE.0) NZEL=7
CALL FIT(NZEL,P,LEO,XX,YY,ERR,O,GAMMA,1)
PAR(31)=PAR(79)

C HALBWERTSBREITEN,FLAECHEEN BEI ARC4
CALL HALBW(PAR(73),PAR(75),PAR(77),PAR(74), PAR(76),HW,DHW,FHW,MW)
PAR(34)=HW
PAR(33)=WW*PAR(71)
PAR(35)=PAR(71)
PAR(31)=PAR(79)-DHW
PAR(48)=DHW
DESCHW=2.0*PAR(34)*PAR(1)
97 CONTINUE

C AUSDRUCK HALBWERTSBREITEN
IF(LPAR(20).GT.1) CALL ZETKO
IF(LUNG HALBWERTSBREITEN
STR AHLUNGSSCHWANZ
  GO TO 40
  NAL=0
  NSW=40
  IF(LPAR(15).LT.0) NAL=1
  DCUBLE PRECISION U
  U=PAR(7)
  EI=PAR(31)*PAR(1)
  T=PAR(58)
  IF(LPAR(49).EQ.3) R=PAR(55)
  IF(LPAR(49).EQ.3) T=PAR(56)
  ATW=PAR(52)
  XRAD=PAR(53)
  Z=PAR(51)
  FLAE=PAR(33)*PAR(1)
  DES=PAR(34)*2.
  XGO=PAR(49)+0.3
  
```



```

XGCB=(PAR(31)-DES)*50.
IF(XGO.GT.(PAR(31)-DES)) XGO=FLOAT(IFIX(XGCB))/50.
XGU=PAR(50)-0.3
DIFF=XGO-XGU
DES=PAR(34)*2.*PAR(1)
ABST=0.02
IF(DIFF.GT.1.2) ABST=0.03
IF(DIFF.GT.1.6) ABST=0.04
IF(LP(20).GT.1) CALL ZETKO
C DIR EKT EINGABE STUETZSTELLEN
IF(PAR(21).GT.0.) XGO=PAR(21)
IF(PAR(22).GT.0.) XGU=PAR(22)
IF(PAR(23).GT.0.) ABST=PAR(23)
C
DO 20 I=1,500
XX(I)=XGO-ABST*FLOAT(I-1)
IF(XX(I).LT.XGU) GO TO 25
E=XX(I)*PAR(1)
CALL RADTL(U,YY(I),EI,E,R,T,ATW,Z,XRAD,DESCHW,FLAE,PAR(41),NAL,
1 .02)
ADISP=E/EI
IF(PAR(64).EQ.0.) YY(I)=YY(I)*ADISP
IF(LP(40).EQ.0) PAR(90)=1.
YY(I)=YY(I)*PAR(90)
20 ERR(I)=0.01*ABS(YY(I))
25 NSWI=I-1
IF(LP(39).NE.1) GO TO 31
XX(NSWI+1)=0.
YY(NSWI+1)=0.
ERR(NSWI+1)=0.
FCR MAT (IH1,20A4)
WRITE(7,501) (WORD(I),I=1,20)
501 DC 30
FORMAT (2(F9.4,1X,F15.7,F13.7,2X))
500 WRITE(7,500) XX(I),YY(I),XX(I+1),YY(I+1),ERR(I+1)
30 CONTINUE
C
31 IF(LP(20).GT.1) CALL ZETKO
DO 26 J=91,100
PAR(J)=0.0
26 LP(24)=1
28 NPSW=LP(17)
IF(LP(17).EQ.0) NPSW=4
IF(LP(16).EQ.0) LP(16)=1
J1=IABS(LP(16))
J=J1+1
P(I)=0.
P(4)=0.

```

```

ELST11200
ELST11210
ELST11220
ELST11230
ELST11240
ELST11250
ELST11260
ELST11270
ELST11280
ELST11300
ELST11320
ELST11330
ELST11340
ELST11350
ELST11360
ELST11370
ELST11380
ELST11390
ELST11400
ELST11410
ELST11420
ELST11430
ELST11460
ELST11470
ELST11490
ELST11500
ELST11510
ELST11520
ELST11530
ELST11560
ELST11570
ELST11580
ELST11590
ELST11600
ELST11620
ELST11630
ELST11640
ELST11670
ELST11680
ELST11690
ELST11700
ELST11710
ELST11750
ELST11760

```



```

P(5)=0.
X1=1./((PAR(31)-XX(1)),
X2=1./((PAR(31)-XX(NSW1))
Y2=YY(NSW1)
GC TO (50,50,53,53,55),J1
C
50 P(3)=(YY(1)/X1-Y2/X2)/(X1-X2)
P(2)=Y2/X2-P(3)*X2
GO TO 32
53 P(2)=(YY(1)*XX(NSW1)-Y2*XX(1))/((XX(NSW1)*X1-XX(1))*X2)
IF(J1.EQ.4) GO TO 54
GO TO 32
54 P(1)=P(2)
P(2)=P(3)
P(3)=0.
GO TO 32
55 IF(LPAR(18).EQ.0) LPAR(18)=2
PAR(81)=(XG0-XG0)/FLOAT(LPAR(18))+XG0
XK1=XX(1)-PAR(81)
XK2=XX(NSW1)-PAR(81)
P(3)=(YY(1)/XK1-Y2/XK2)/(XK1-XK2)
P(2)=YY(1)/XK1-P(3)*XK1
C
32 CALL FIT(NPSW,P,NSW1,XX,YY,ERR,0,GAMMA,J)
IF(LPAR(20).GT.1) CALL ZETKO
40 RETURN
END
C
SSWF STRAHLENSCHWANZFIT
SUBROUTINE SSWF(Y,X)
DIMENSION PAR(300),LPAR(100)
COMMON PAR,LPAR
J=IABS(LPAR(16))
IF(J.EQ.5) GO TO 50
A=1./((PAR(31)-X)
Y=PAR(91)+PAR(93)*A
GO TO (10,20,30,40),J
10 Y=Y+PAR(95)*A**2+PAR(97)*A**3+PAR(99)*A**4
RETURN
20 Y=Y+PAR(95)*A**2+PAR(97)*X+PAR(99)*X**2
RETURN
30 Y=Y+PAR(95)*X+PAR(97)*X**2+PAR(99)*X**3
RETURN
40 Y=PAR(91)+PAR(93)*X+PAR(95)*X**2+PAR(97)*A**2+PAR(99)*X**3
RETURN
50 A=X-PAR(81)
Y=PAR(91)+PAR(93)*A+PAR(95)*A**2+PAR(97)*A**3+PAR(99)*A**4

```

```

ELST1770
ELST1790
ELST1800
ELST1810
ELST1830
ELST1840
ELST1850
ELST1860
ELST1870
ELST1890
ELST1900
ELST1920
ELST1930
ELST1960
ELST1970
ELST1980
ELST1990
ELST2020
ELST2030
ELST2040
ELST2050
ELST2060
ELST2070
ELST2080
ELST2090
ELST2100
ELST2130
ELST2140
ELST2160
ELST2170
ELST2180
ELST2190
ELST2200
ELST2210
ELST2220
ELST2230
ELST2240
ELST2250
ELST2260
ELST2270
ELST2280
ELST2290
ELST2300
ELST2310
ELST2320
ELST2330
ELST2340

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ELST2350  
ELST2360

ELST5430  
ELST5460  
ELST5480

ELST5520  
ELST5540  
ELST5550  
ELST5580

ELST5600  
ELST5610  
ELST5620  
ELST5630  
ELST5640  
ELST5650  
ELST5660  
ELST5670  
ELST5680  
ELST5690

ELST5750  
ELST5770  
ELST5780  
ELST5790  
ELST5800  
ELST5810  
ELST5820  
ELST5830  
ELST5840  
ELST5850  
ELST5860  
ELST5870  
ELST5880  
ELST5890  
ELST5900  
ELST5910  
ELST5930

```

RETURN
END
C ANPASSUNG NACH DER METHODE DER KLEINSTEN QUADRATE
SLBROUTINE FIT(MZ,P,N,X,Y,DY,NRM,DELTA,NUM)
C NRM = 0 FALLS DIE MESSFEHLER ABSOLUT SIND
DIMENSION PAR(300),LPAR(100),DUM(600), WORD(20)
DIMENSION X(500),Y(500),DY(500), V1(25), V2(25),
P(25), DF(25), DZ(25), DZ2(50,50), U(50,50)
1 DIMENSION JOTA(10)
DIMENSION ZWP(25),PNB(25)
DIMENSION PAR,LPAR,DUM, WORD
COMMON TEXT,NAME1,NAME2,NAME3
COMMON/SSWD/ACO(25)
COMMON/LJINGL/ DZ2,U
DATA V1,V2,DP,DF,DZ,ZWP,PNB /175*.0/
DATA NAME3 /: GRD:/
GAM1=.0
NORM=.0
GAMMA=DELTA
M=MZ
MY=5
JCO6=6
MCO2=20
MCO1=5
DO 11 K=1,50
DO 11 L=1,50
DZ2(K,L)=.0
U(K,L)=.0
11 DIAGNOSTIK
IF(M) 16,16,15
13 IF(M-25) 17,17,16
15 IF(M-25) 17,17,16
16 WRITE(6,997)M
997 FORMAT (40H1 NUR 1 BIS 25 PARAMETER ERLAUBT STATT ,I7)
17 RETURN
18 IF(N-500) 20,20,18
19 WRITE(6,998)N
998 FORMAT (42H1 NUR 1 BIS 500 MESSPUNKTE ERLAUBT STATT ,I7 )
20 ZSCULL=N - M
24 IF(ZSOLL) 24,25,26
999 WRITE(6,999)M,N
999 FORMAT (19H1 MEHR PARAMETER (, I2, 20H) ALS MESSPUNKTE (, I2,
1 RETURN
25 NCRM = 0

```





```

26 DO 27 I = 1,N
27 IF(DY(I)) 28,28,27
CONTINUE
GC TO 14
28 WRITE (6,1000)I
1000 FORMAT (49H1 MESSFEHLER DY(I) IST NULL ODER NEGATIV FUER I = , I4)
RETURN
14 FEH=1.0
IF(NORM) 29,30,30
29 FEH=DY(1)**2
30 J = 0
C BEGINN DER ITERATION ZUR LOESUNG DER NORMAL - GLEICHUNGEN
IF(LPAR(8)) 31,31,32
31 WRITE (6,1002)
1002 FORMAT (1H1,17X,18HZWISCHENERGEBNISSE, 1.PARAMETER 2.PARAMETER
/85H NAEHERUNG ZWANG //)
3.PARAMETER USW.....
32 G = 1.0
DATA NAME1/'
TEXT=NAME1
C
50 Z = 0.0 K = 1,M
DO 60 K = 0.0
DC 60 L = 1,M
DZZ(K,L) = 0.0
DO 70 I = 1,NUM,X(I),P,F,DF
H1 = DY(I)**(-2)*FEH
H2 = Y(I) - F
Z = Z + H1*H2*H2
DC 70 K = 1,M
DZ(K) = DZ(K) - H1*H2*DF(K)
DO 70 L = 1,M
DZZ(K,L) = DZZ(K,L) + H1*DF(K)*DF(L)
70 ZY=Z/FEH 72,73,72
71 IF(NORM) 74,74,78
72 G=Z/ZSOLL 74,74,78
73 IF(LPAR(8)) 77,77,75
74 IF(NORM) 77,77,75 TEXT, Z, (P(K), K = 1,M)
75 WRITE (6,1003)J, IPE14.6,3X,7E15.6/(27X,7E15.6)
1003 FORMAT (1X,13,A6,1PE14.6,3X,7E15.6/(27X,7E15.6))
GO TO 78
77 WRITE (6,1004)J, TEXT,ZY, (P(K), K = 1,M)
1004 FORMAT (1X,13,A6,1PE15.6/(27X,7E15.6))
C PRUEFUNG AUF ITERATIONSSTOP
78 IF(J-MY) 100,80,80
80 IF (ABS(Z-ZALT)-1.0E-02) 140,140,79

```



ELST6470  
 ELST6480  
 ELST6490  
 ELST6500  
 ELST6510  
 ELST6520  
 ELST6530  
 ELST6550  
 ELST6570  
 ELST6580  
 ELST6590  
 ELST6600  
 ELST6610  
 ELST6620  
 ELST6630  
 ELST6640  
 ELST6650  
 ELST6660  
 ELST6670  
 ELST6680  
 ELST6700  
 ELST6720  
 ELST6730  
 ELST6740  
 ELST6750  
 ELST6770  
 ELST6780  
 ELST6790  
 ELST6800  
 ELST6810  
 ELST6820  
 ELST6830  
 ELST6840  
 ELST6860  
 ELST6870  
 ELST6880  
 ELST6890  
 ELST6900  
 ELST6910  
 ELST6920  
 ELST6930  
 ELST6940  
 ELST6950  
 ELST6960  
 ELST6970

```

79 IF (Z-ZALT) 90,102,102
102 DO 103 K=1,M
103 P(K)=ZWP(K)
   GAMMA=0.8*GAM1
1040 WRITE (6,1040)GAMMA
   WFORMAT(13X,10H GAMMA = F13.6)
90 IF (J-M00Z) 100,130,130
   DURCHFUEHRUNG DER ITERATION
C 100 DO 104 K = 1,M
   DP(K) = - DZ(K)
   DC 104 L = 1,M
104 U(K,L) = DZZ(K,L)
105 CALL LINGL (U,DP,M,1)
106 DC 106 K = 1,M
   H1 = H1 + DP(K)*DZZ(K,K)
   DATA NAME2/' NWT'/
   TEXT=NAME2(H1)
   GAM1=SQRT(GAMMA) 113,113,107
C IF (GAM1-GAMMA) 113,113,107
   GRADIENTEN METHODE STATT NEWTON
C 107 TEXT=NAME3
   H1 = 0.0
   H2 = 0.0
   DC 109 K = 1,M
   H3 = DZ(K)
   DZ(K) = H3/DZZ(K,K)
109 H1 = H1 + H3*DZ(K)
   DO 110 L = 1,M
   DC 110 L = 1,M
110 H2 = H2 + DZ(K)*DZZ(K,L)*DZ(L)
   H3 = AMIN1(H1/H2,GAMMA/SQRT(H1))
112 DO 112 K = 1,M
   DP(K) = - H3*DZ(K)
113 ZPALT = Z
   IF (J.LE.M001) GO TO 114
114 DO 120 K = 1,M
   I1=I1+2
120 ZWP(K)=P(K)
124 P(K) = P(K) + DP(K)
   J=J+1
   GC TO 50
C
C ALSDRUCK FALLS NOCH KEINE KONVERGENZ NACH 20 ITERATIONEN
120 WRITE (6,1005) M00Z
1005 FCRMAT (24H (KEINE KONVERGENZ NACH 13,15H ITERATIONEN) )
114C WRITE (6,1030)(WORD(I),I=1,20)
1030 FCRMAT(10H,20A4//)

```



```

131 IF(LPAR(24)) 131,131,132
1031 WRITE(6,1031)
1031 FORMAT(27H FITERGEBNISSE ELAST. PEAK )
1032 GO TO 133
1032 WRITE(6,1032)
1032 FCRMAT(20H FITERGEBNISSE SSW )
502 IF(ACO(1).NE.0.)WRITE(6,502)
1033 IF(NORM) 35,35,40
1033 IF(NORM) 35,35,40
C ERWARTUNGSWERT VON Z IST ZSOLL = N * M FALLS DIE MESSFEHLER DY(I)
C ABSOLUT SIND. STREUUNG IST SQRT (2*ZSOLL)
35 H1 = SQRT(2.0*ZSOLL)
A012=H1
B012=Z
WRITE(6,1007)ZSOLL, H1, ZY
1007 FCRMAT(9H ZWANG, 9X, 8HSOLLWERT, F8.1, 3H +- , F5.1
1  /18X,7HISTWERT ,F9.1)
40 WRITE(6,1008)
1008 FCRMAT(11H,32X,7HOPTIMAL,13X,6HFEHLER,
1 28X,27HKONTROLLE (SOLLWERT IST 1) ,//)
DC 160 K = 1,M
DC 150 L = 1,M
150 U(K,L) = 0.0
160 U(K,K) = G
170 CALL LINGL (D2Z, U, M, M)
180 DP(K) = SQRT(ABS(U(K,K)))
C
C KONTROLLE DURCH BERECHNUNG DES ZWANGES AN DEN EXTREMPUNKTEN
C DES FEHLERELLIPTISOIDS
DC 205 K = 1,M
DC 190 L = 1,M
U(K,L) = U(K,L)/DP(K)
V1(L) = P(L) - U(K,L)
V2(L) = P(L) + U(K,L)
190 U(K,L) = U(K,L)/QP(L)
J = J+1
H1 = 0.0
DC 195 I = 1,N
CALL THEORY (NUM,X(I),V1,F,DF)
H1 = H1 + ((Y(I) - F)/DY(I))**2*FEH
J = J+1
H2 = 0.0
DC 200 I = 1,N
CALL THEORY( NUM,X(I),V2,F,DF)

```

```

ELST77020
ELST77030
ELST77040
ELST77050
ELST77060
ELST77070
ELST77080
ELST77090
ELST77100
ELST77110
ELST77120
ELST77130
ELST77140
ELST77150
ELST77160
ELST77170
ELST77180
ELST77190
ELST77200
ELST77210
ELST77220
ELST77230
ELST77240
ELST77250
ELST77260
ELST77270
ELST77280
ELST77290
ELST77300
ELST77310
ELST77320
ELST77330
ELST77350
ELST77360
ELST77370
ELST77380
ELST77390
ELST77400
ELST77410
ELST77420
ELST77430
ELST77440
ELST77450
ELST77460
ELST77470
ELST77480
ELST77490
ELST77500

```



```

200 H2 = H2 + ((Y(I) - F)/DY(I))*2*FEH
H2 = (H2 - Z)/G
205 WRITE (6,1009)K, P(K), DP(K), H1, H2
1009 FORMAT (1X,I2,10H.PARAMETER,15X,1PE15.6,5H +- ,E15.6,11X,0P2F20.5
)
1010 WRITE (6,1010)
FORMAT (40H KORRELATIONEN ZWISCHEN DEN PARAMETERN
//12X,31H1.PA 2.PA 3.PA USW....
//)
185 DC 185 K = 1,M
1011 WRITE (6,1011)K,(U(K,L), L = 1,M)
1011 FORMAT (1X,I2,6H.PA ,15F8.3/9X,10F8.3)
C
C
TABELLIERUNG DER ANGEPAASSTEN FUNKTION
IF(LPAR(8)) 181,186,181
IF(LPAR(9)) 216,216,182
181 WRITE (6,1015)
186 WRITE (6,1015)
1015 FORMAT(1H0,54X,13HT A B E L L E ///56X,8HORDINATE/
114H NR. ABSZISSE
8BERECHNET DIFFERENZ MESSFEHLER GEMESSEN
//)
182 IF(LPAR(9)).LE.0.OR.LPAR(24).GT.0) GO TO 188
187 WRITE (1) LPAR(24),
LPAR(24), M,N,ZSOLL,A012,B012
188 DC 213 I = 1,N
CALL THEORY (NUM ,X(I),P,F,DF)
H1 = Y(I) - F
H14=H1/DY(I)
IF(LPAR(9)).EQ.0.OR.LPAR(24).GT.0) GO TO 215
UNT=0.0
WRITE (1) X(I),Y(I),F,H14,UNT,UNT
IF(LPAR(8)).NE.0) GO TO 213
214 WRITE (6,1016)I, X(I),Y(I),F,H1,DY(I)
1016 FORMAT (15,F19.4,3X,1P2(E27.6,E17.6))
213 CONTINUE
C
216 NO=70
IF(LPAR(24).GT.0) NO=90
1021 DC 1022 I=1,M
NCN=NO+2*I
PAR(NON-1)=P(I)
PAR(NON)=DP(I)
1022 RETURN
END
C
SUBROUTINE THEORY(J,X,P,Y,DA)
THEORY
ELAST
DIMENSION P(25),DA(25),PAR(300),LPAR(100)
COMMON PAR,LPAR
GO TO (10,20,30,40,50,60) , J

```





```

10 CALL ARC4(X,P,Y,DA)
RETURN
20 A=1./ (PAR(31)-X)
Y=P(1)+P(2)*A+P(3)*A**2+P(4)*A**3
IF(LPAR(33).EQ.5) Y=Y+P(5)*A**4
DA(1)=1
DA(2)=A
DA(3)=A**2
DA(4)=A**3
DA(5)=A**4
RETURN
30 A=1./ (PAR(31)-X)
Y=P(1)+P(2)*A+P(3)*A**2+P(4)*X+P(5)*X**2
DA(1)=1.
DA(2)=A
DA(3)=A**2
DA(4)=X
DA(5)=X**2
RETURN
40 A=1./ (PAR(31)-X)
Y=P(1)+P(2)*A+P(3)*X+P(4)*X**2+P(5)*X**3
DA(1)=1.
DA(2)=A
DA(3)=X
DA(4)=X**2
DA(5)=X**3
RETURN
50 A=1./ (PAR(31)-X)
Y=P(1)+P(2)*X+P(3)*X**2+P(4)*A**2+P(5)*X**3
DA(1)=A
DA(2)=X
DA(3)=X**2
DA(4)=A**2
DA(5)=X**3
RETURN
60 A=X-PAR(81)
Y=P(1)+P(2)*A+P(3)*A**2+P(4)*A**3+P(5)*A**4
DA(1)=1.
DA(2)=A
DA(3)=A**2
DA(4)=A**3
DA(5)=A**4
RETURN
END

```

SUBROUTINE ARC4(X,P,Y,DY)

ELST 9400



ELST9410  
 ELST9420  
 ELST9430  
 ELST9440  
 ELST9450  
 ELST9460  
 ELST9470  
 ELST9480  
 ELST9490  
 ELST9500  
 ELST9510  
 ELST9520  
 ELST9530  
 ELST9540  
 ELST9550  
 ELST9560  
 ELST9570  
 ELST9580  
 ELST9590  
 ELST9600  
 ELST9610  
 ELST9620  
 ELST9630  
 ELST9640  
 ELST9650  
 ELST9660  
 ELST9670  
 ELST9680  
 ELST9690  
 ELST9700  
 ELST9710  
 ELST9720  
 ELST9730  
 ELST9740  
 ELST9750  
 ELST9760  
 ELST9770  
 ELST9780  
 ELST9790  
 ELST9800  
 ELST9810  
 ELST9820  
 ELST9830  
 ELST9840  
 ELST9850  
 ELST9860  
 ELST9870  
 ELST9880

```

DIMENSION P(25),DY(25)
BT=1.44269505
DEZ,71: ZWEITE EL LINIE MIT HOEHE=P(6), LAGE=P(7)
U=(X-P(5))
IF(U) 20,10,10
CONTINUE
ALPHA=(U/P(2))**2/BT
FORMAT(5E15.4)
Y=P(1)*EXP(-ALPHA)
DY(1)=Y/P(1)
DY(2)=2.0*Y*U**2/(BT*P(2)**3)
DY(3)=0.0
DY(4)=0.0
DY(5)=2.0*U*Y/(BT*P(2)**2)
IF(P(6).EQ.0) RETURN
Y6=P(6)*EXP(-(V/P(2))**2/BT)
DY(6)=Y6/P(6)
DY(7)=2.0*V*Y6/(BT*P(2)**2)
Y=Y+Y6
DY(2)=DY(2)+ 2.*Y6*V**2/(BT*P(2)**3)
RETURN
YF=P(1)*EXP(-(P(4)/P(3))**2/BT)
IF(U-P(4)) 42,30,30
Y=P(1)*EXP(-(U/P(3))**2/BT)
DY(1)=Y/P(1)
DY(2)=0.0
DY(5)=2.0*Y*U/(BT*P(3)**2)
DY(4)=0.0
DY(3)=DY(5)*U/P(3)
IF(X.GE.P(7)) GOTO 60
Y6=P(6)*EXP(-(V/P(3))**2/BT)
DY(6)=Y6/P(6)
DY(7)=2.0*V*Y6/(BT*P(3)**2)
Y=Y+Y6
DY(3)= DY(3)+2.*Y6*V**2/(BT*P(3)**3)
RETURN
PP=2.0*(P(4)/P(3))**2/BT
K6=0
IF(K6.EQ.1) YP=YP6
C=Y*PP(4)**2*(2.0-PP)
D=Y*PP(4)**2*(PP-1.0)
PPP=YP*PP
CP3=PPP*PP(4)*(4.0-PP)/P(3)
DP3=PPP*PP*(P(4)**2/P(3)-3.0*BT*P(3)/2.0)
CP4=PPP*(2.0/PP-5.0+PP)
DP4=PPP*PP(4)*(5.0-PP-2.0/PP)

```



ELST19890  
 ELST19900  
 ELST19910  
 ELST19920  
 ELST19930  
 ELST19940  
 ELST19950  
 ELST19960  
 ELST19970  
 ELST19980  
 ELST19990  
 ELST20000  
 ELST20010  
 ELST20020  
 ELST20030  
 ELST20040  
 ELST20050  
 ELST20060  
 ELST20070  
 ELST20080  
 ELST20090  
 ELST20100

ELST01130  
 ELST01140  
 ELST01150  
 ELST01160  
 ELST01170  
 ELST01180  
 ELST01190  
 ELST02000  
 ELST02010  
 ELST02020  
 ELST02030  
 ELST02040  
 ELST02050  
 ELST02060  
 ELST02070  
 ELST02080  
 ELST02090  
 ELST03000  
 ELST03010  
 ELST03020  
 ELST03030  
 ELST03040  
 ELST03060  
 ELST03080

```

IF(K6.EQ.1) GOTO 63
Y=C/U+D/U**2
DY(1)=Y/P(1)
DY(2)=0.0
DY(3)=CP3/U+DP3/U**2
DY(4)=CP4/U+DP4/U**2
DY(5)=C/U**2+2.0*D/U**3
DY(P(6)).EQ.0) RETURN
IF(X.GE.P(7)) GOTO60
IF(X.LT.P(7)) .AND. X.GE.P(4)) GOTO 61
K6=1
YP6=P(6)*EXP(-P(4)/P(3))**2/BT
GOTO 62
Y6=C/V+D/V**2
DY(6)=Y6/P(6)
Y=Y+Y6
DY(7)=C/V**2+2.0*D/V**3
DY(4)=DY(4)+CP4/V + DP4/V**2
DY(3)=DY(3)+CP3/V + DP3/V**2
RETURN
END
  
```

63

```

SUBROUTINE HALBW(P2,P3,P4,DP2,DP3,HWB,DHWB,FHW,W)
  
```

1

```

IF(P2) 1,1,2
HNB=0.0
DHWB=0.0
FHW=0.0
W=0.0
  
```

2

```

RETURN
PVP=P4/P3
BT=1.44269505
YP=EXP(-PVP**2/BT)
HNB=P2+P3
DHWB=(P3-P2)/2.0
PP=2.0*PVP**2/BT
C=Yp**P4**2*(PP-1.0)
X=-PVP/SQRT(BT)
W=ERF(X)
W0=P2*SQRT(BT**3.14159265)*.5*(1.0+W**P3/P2)
ZA=2.
GRENZ=ZA*HNB+.5*DHWB
W2=C*ALOG(-P4/GRENZ)-D/P4-D/GRENZ
W=W0+W2
FHW=1.0
RETURN
END
  
```



```

SUBROUTINE RADIL(U1,SW,EIMEV,EMMV,R,T,ATW,Z,XRAD,DES,DS,D,NA1,ERR)
DIMENSION PAR(300),LPAR(100)
COMMON PAR,PRECISION
DCUBBLE PRECISION
DCUBBLE PRECISION
DCUBBLE PRECISION
U=U1
ERROR=ERR
EIMEV=.5109
DESCHW=DES
DSMOT=DS
EMIN=EMMV/.5109
SU2=COS(.0174533*SNGL(U)/2.0)
U=DCGS(.0174533*U)
DEFF=DS/(2.0*XRAD)
REF=EXP(-.3333*ALOG(Z))*.047829*.5
Z13=ALOG(.183*Z13)
XC=2.0*(SNGL(EI)*SU2)**2
CALL FORM(FO,XO,R,T)
CALL DMCTT(DM,CU2,SU2,Z,EI)
SIGMA=DM*
ES=EIMEV
W=EI-EF
PI=DSQRT(EI**2-1.00)
PF=DSQRT(EF**2-1.00)
A=DSQRT(PI**2-2.0*PI*PF*U+PF**2)
B=EI*EF-PI*PF*U-1.00
DSIGMA=0.0
IF(NA1) 303,303,304
XB=B+W*(W-PI+PF*U)+B
XMIN=.500*(A-W)**2
XMAX=.500*(A+W)**2
IF(SNGL(XA)-SNGL(XB)-15.0*SNGL(XB)-15.0*SNGL(XB)-15.0*SNGL(XB))/SNGL(EI)/SNGL(EF)) 5,6,6
1(EF)) 5,6,6
5 X2=XB
X3=XA
GC TO 7
6 X2=XB+10.00*W*EF/EI
X3=XA-10.00*W*EI/EF
7 IF(SNGL(XB)-15.0*SNGL(XB)-15.0*SNGL(XB)-15.0*SNGL(XB))/SNGL(EI)-SNGL(XMIN)) 8,9,9
8 X1=XB
GC TO 10

```





```

9 X1=XB-10.00*W*EF/EI
10 IF(SNGL(XMAX)-SNGL(XA)-15.0*SNGL(W)*SNGL(EI)/SNGL(EF)) 11,12,12
11 GC TO 2
12 X4=X4+10.00*W*EI/EF
2 CONTINUE
CALL WEDDLE(H1,XMIN,X1,ERROR,EI,EF,PI,PF,U,A,B,W,R,T)
CALL WEDDLE(H2,X1,X2,ERROR,EI,EF,PI,PF,U,A,B,W,R,T)
CALL WEDDLE(H3,X2,X3,ERROR,EI,EF,PI,PF,U,A,B,W,R,T)
CALL WEDDLE(H4,X3,X4,ERROR,EI,EF,PI,PF,U,A,B,W,R,T)
CALL WEDDLE(H5,X4,XMAX,ERROR,EI,EF,PI,PF,U,A,B,W,R,T)
DSIGMA =H1+H2+H3+H4+H5
DSIGMA =DSIGMA *90.2376*(SNGL(PF)/SNGL(PI))*Z**2
304 EFMEV =.5109*ES
XEF=2.0*(SNGL(EF)*SU2)**2
CALL FORM(FE,XE,R,T)
BRA L=(1.0+((SNGL(EI)*FE)/(SNGL(EF)*FO))**2)
SCHIFR =0.0
SCH=0
IF(NAI.GT.0) CALL SCHIFC(SCH,SU2,EI,EF,W)
SCHIFR=BRA *SCH*DSMOT
DSIGMA=DSIGMA*DSMOT/SIGMA
306 CALL HEIMOL(H,REFF,DEFF,GZ13,EI,EF)
HEM=DSIGMA+HEMO
SSW=SSW*(1.+DELT)
FCR MAT(5E15.6)
1001 RETURN
END

SUBROUTINE CAPR(Z,Y,EI,EF,PI,PF,U,A,B,W,R,T)
SUBROUTINE FOR CALCULATION OF R(X) FOR ELECTRON SPECTRUM
DOUBLE PRECISION A,D,DSQRT,X,EI,EF
DOUBLE PRECISION D1,DSQRT,X,B,D2,PI,PF,G1
DOUBLE PRECISION G2,G3,W,U,PI,PF,F
X=PI**2-PI*PF*U
B=91-A**2
D=2.00*X*((EI*EF-1.00)*U-PI*PF)*PI*PF
D1=DSQRT((PI*X)**2+D+(PI*B)**2)
D2=DSQRT((PI**2-X)*(B*B1+X*B2)/D2**3)
G1=(2.00*EF**2-X)*(B*B2+X*B1)/D1**3
G2=2.00*(EI**2+EF**2)+B-X-2.00*((1.00+B)*((2.00*EI*EF-X)+B*W**2)/(B
1-X)
F=((1.00/D2)-(1.00/D1))*G3-G2-G1-(2.00/A)
Z=$NGL(F)/SNGL(X)**2

```

ELST11020  
ELST11030  
ELST11040  
ELST11050  
ELST11060  
ELST11070  
ELST11080  
ELST11090  
ELST11100  
ELST11110  
ELST11120  
ELST11130  
ELST11140  
ELST11150  
ELST11160  
ELST11170  
ELST11180  
ELST11190  
ELST11200  
ELST11210  
ELST11220  
ELST11230  
ELST11240  
ELST11250  
ELST11260  
ELST11270  
ELST11280  
ELST11290  
ELST11300  
ELST11310  
ELST11320  
ELST11330  
ELST11340  
ELST11350  
ELST11360

ELST11390  
ELST11380  
ELST11400  
ELST11410  
ELST11420  
ELST11430  
ELST11440  
ELST11450  
ELST11460  
ELST11470  
ELST11480  
ELST11490  
ELST11500  
ELST11510  
ELST11520  
ELST11530  
ELST11540



```

XX=SNGL(X)
CALL FORM(FOR,XX,R,T)
Z=Z**2
RETURN
END
ELST11550
ELST11560
ELST11570
ELST11580
ELST11590

```

```

SUBROUTINE FORM(F,X,R,T)
WAHLWEISE TRAPEZT ODER EXAKTER FCRMFAKTOR IN POTENZEN VON Q**2
EINGEFUEGT AUGUST 72 PITTHAN,WALCHER
COMMON/SSWD/ACO(25)
IF(ACO(1).EQ.0.) GOTO 5
QC=X**2./386.1**2

```

```

DO 4 I=1,12
F=F+ACO(I)**QQ**((I-1)
RETURN
4 TRAPEZFORMFAKTOR:

```

```

5 Q=(R+5.0*T/8.0)*SQRT(2.0*X)/386.154
ETA=(R-5.0*T/8.0)/(R+5.0*T/8.0)
QN=ETA**Q
P=12.0*(QN*SIN(QN)-Q*SIN(Q)+2.0*COS(QN))-2.0*COS(Q)/(Q**4-QN**4)
WRITE(6,10)
FORMAT(2F20.10)
RETURN
END

```

```

SUBROUTINE DMOTT(X,CU2,SU2,Z,EI)
DOUBLE PRECISION U, W, PF, EI, A, EF, B
DCUBLE PRECISION U, W, PF, EI, A, EF, B
ANGL=CU2**2/(SU2**4)
X=ANGL*(Z/(2.0*EI))**2*79411.0
RETURN
END

```

```

SUBROUTINE SCHIFC(X,SU2,EI,EF,W)
SUBROUTINE RADIATION TAIL DURING SCATTERING
DOUBLE PRECISION U, W, PF, EI, A, EF, B
DCUBLE PRECISION U, W, PF, EI, A, EF, B
X1=(1.0+(EF/EI)**2)*(DLOG(2.0*EI*SU2)-0.5)
X=.002323*X1/(W**5109)
RETURN
END

```

```

SUBROUTINE SCHWINGER CORRECTION
SUBROUTINE DELTA(X,SU2,DESCHW,EI)
DOUBLE PRECISION U, W, PF, EI, A, EF, B
DCUBLE PRECISION U, W, PF, EI, A, EF, B
R1=DLOG(2.0*EI

```

```

ELST11630
ELST11610
ELST11620
ELST11650
ELST11660
ELST11670
ELST11680
ELST11690
ELST11700
ELST11710
ELST11720
ELST11730
ELST11740
ELST11750
ELST11760
ELST11770
ELST11780
ELST11790
ELST11820
ELST11830
ELST11840
ELST11850
ELST11860
ELST11870
ELST11880
ELST11910
ELST11900
ELST11920
ELST11930
ELST11940
ELST11950
ELST11960
ELST11970
ELST11990
ELST12000
ELST12010
ELST12020
ELST12030

```



```

A1= DESCHW/.5109
R2= DLOG(EI /A1)-1.0833
X=.009236*(R1*R2+.2361)
RETURN
END

```

```

SUBROUTINE WEDDLE(V,A,B,ERROR,EI,EF,PI,PF,U01,A01,B01,W01,RO1,TO1)
CWEEDLE EVALUATION OF QUADRATURES USING WEDDLE-S RULE
DIMENSION C(7),F(200)
C DIMENSION C(7),F(200)
DOUBLE PRECISION U01
DOUBLE PRECISION PI
FCRMAI(29H **INTEGRAL NOT CONVERGENT**, /)
FCRMAI(8E14.6)
FCRMAI(//17H EI4.6,24H TO,EI4.6,4H TO,EI4.6,10H GIVES V=,
1 E14.6,24H (PREVIOUS VALUE - W=,EI4.6,1H),//)
3 FCRMAI(21H **A=B, INTEGRAL=0**,//)
IF(B-A) 2,3,2
WRITE (6,103)
V=0.0
GO TO 1
DC 10 I=1,7,2
10 C(I)=1.0
C(2)=5.0
C(4)=6.0
C(6)=5.0
DY=(B-A)/192.0
SUM=0.0
L=0
DC 20 I=1,193,32
L=L+1
Y=A+DY*FLOAT(I-1)
CALL CAPR(Z01,Y,EI,EF,PI,PF,U01,A01,B01,W01,RO1,TO1)
20 F(I)=Z01
SUM=SUM+F(I)*C(L)
V=.3*SUM*DY*32.0
DC 50 N=1,5
h=V
K=2**(6-N)
M=K/2
K1=M+1
K2=193-M
DC 30 J=K1,K2,K
Y=A+DY*FLOAT(J-1)
CALL CAPR(Z01,Y,EI,EF,PI,PF,U01,A01,B01,W01,RO1,TO1)
30 F(J)=Z01
K3=3*K
J1=193-K3

```

```

ELST 2040
ELST 2050
ELST 2060
ELST 2070
ELST 2080

```

```

ELST 2110
ELST 2120
ELST 2130
ELST 2140
ELST 2150
ELST 2160
ELST 2170
ELST 2180
ELST 2190
ELST 2200
ELST 2210
ELST 2220
ELST 2230
ELST 2240
ELST 2250
ELST 2260
ELST 2270
ELST 2280
ELST 2290
ELST 2300
ELST 2310
ELST 2320
ELST 2330
ELST 2340
ELST 2350
ELST 2360
ELST 2370
ELST 2380
ELST 2390
ELST 2400
ELST 2410
ELST 2420
ELST 2430
ELST 2440
ELST 2450
ELST 2460
ELST 2470
ELST 2480
ELST 2490
ELST 2500
ELST 2510

```



```

SUM=0.0
DC 40 J=1,J1,K3
L2=J+K3
L=0
DC 40 I=J,J2,M
L=L+1
SUM=SUM+F(I)*C(L)
V=3*SUM*DY*FLOAT(M)
IF(ABS((V-W)/V)-ERROR) 1,1,50
CONTINUE
WRITE (6,100)
1 RETURN
END

```

```

CHEM
SUBROUTINE HEIMDL(X,REFF,DEFF,GZ13,EI,EF)
SUBROUTINE HEITLER - MOELLER
DIMENSION PAR(300),LPAR(100)
COMMON PAR,LPAR
DOUBLE PRECISION U , W , PF , EI , A , EF , B
K=EI-EF
BDEE=EF /EI *((EI **2+EF **2)/(EI *EF )-.66667)+1.0/(
19.0*GZ13))
XI=8DEE*DEFF/W
IF(LPAR(40).EQ.0) PAR(88)=1.
X1=X1*PAR(88)
X21=6.28319*REFF **2-(2.0*EI +1.0)/(EF *W *(EI +1.0)**2)
X22=1.0/EF **2+1.0/(EI **2)
X2=X21*X22
IF(LPAR(40).EQ.0) PAR(89)=1.
X2=X2*PAR(89)
X=(X1+X2)/.5109
RETURN
END

```

```

C UNELASTIC UNST
SUBROUTINE PAR(300),LPAR(100),FREQ(200),ZAEH(200),FEH(200)
DIMENSION PAR(300),LPAR(100),ERR(500),YY(500),P(50)
DIMENSION IHO(25)
DIMENSION DUMMY(50)
DIMENSION PVAR(25); BVAR(25)
COMMON PAR,LPAR,FREQ,ZAEH,FEH,
COMMON/BVAR1/BVAR
COMMON/BVAR2/BVAR
COMMON/DATA1/XX,YY,ERR
COMMON /USTART/UPAR(10)

```

ELST2520  
ELST2530  
ELST2540  
ELST2550  
ELST2560  
ELST2570  
ELST2580  
ELST2590  
ELST2600  
ELST2610  
ELST2620  
ELST2630  
ELST2640

ELST2670  
ELST2660  
ELST2680  
ELST2690  
ELST2700  
ELST2710  
ELST2720  
ELST2730  
ELST2740  
ELST2750  
ELST2760  
ELST2770  
ELST2780  
ELST2790  
ELST2800  
ELST2810  
ELST2820  
ELST2830  
ELST2840  
ELST2850  
ELST2860

ELST2880  
ELST2890  
ELST2900  
ELST2920  
ELST2930  
ELST2950  
ELST2970









ELST3460  
ELST3480  
ELST3500  
ELST3510  
ELST3520

ELST3630  
ELST3640  
ELST3660

ELST3720  
ELST3730

ELST3790  
ELST3800  
ELST3880

ELST3890  
ELST3900  
ELST3910  
ELST3920  
ELST3930  
ELST3940  
ELST3950  
ELST3960  
ELST3970  
ELST3980  
ELST3990  
ELST4000

```
C   NAEHERUNGSUNTERGRUND, HOEHENVERHAELTNISSE
      IF(LPAR(20).GT.1) CALL ZETKO
      DO 40 I=1,IM
      IZ=IHO(I)
      CALL SSWF(FSSW,XX(IZ))
      PAR(I+175)=(YY(IZ)-FNAH*FSSW/FLOAT(ISS))/PAR(35)
      IF(LPAR(21).GT.20)PAR(I+175)=(YY(IZ)-UNAH/FLOAT(ISS)-FSSW)/PAR(35)
      IF(PAR(I+175).LE.0.0) PAR(I+175)=0.01*PAR(176)
      IF(PAR(64).GT.1) PAR(I+175)=(YY(IZ)-FSSW)/PAR(35)
40  CCNTINUE
C   NAEHERUNGSPARAMETER FUER FIT
      NPJ=INT(FLOAT(LPAR(2))/10.)
      LPAR2=LPAR(2)
      IF(LPAR2.GT.50) LPAR2=LPAR2-50
      NST=NPA+1
      IF(LPAR2.GT.10) NST=NPA
      IF(LPAR2.GT.20) NST=NPA+2
      IF(LPAR2.GT.30) NST=NPA+1
      DC 41 I=1,NST
41  P(I)=UPAR(I)
      LBEZ=5.*(PAR(13)+PAR(14))
      PAR(45)=PAR(0.1*FLOAT(LBEZ))
      IF(LPAR2.GT.10 .AND. LPAR2.LT.20) GO TO 45
      IF(LPAR2.GT.30 .AND. LPAR2.LT.40) GO TO 45
      IF(LPAR(NPA+1).EQ.0.) P(NPA+1)=FNAH/FLOAT(ISS)
      GO TO 46
45  CCNTINUE
      P(1)=UNAH/FLOAT(ISS)
46  LPAR(22)=NST
      IF(UPAR(1).EQ.0.) GOTO 44
      DC 43 I=1,IM
      IZ=IHO(I)
      CALL UNTF(XX(IZ),P,DUMMY,UNT,FSSW)
43  PAR(I+175)=(YY(IZ)-UNT)/PAR(35)
44  CCNTINUE
      ENERGIENFEST,VARIABLEL
      IF(LFE.EQ.0) LFE=1
      DO 50 I=1,LFE
49  NP=NST+1
      NE=175+LVE+I
50  P(NP)=PAR(NE)
54  LFE=IABS(LPAR(3))
      IF(LVE.EQ.0) LVE=1
      DC 55 I=1,LVE
48  NP=NST+LFE+2*I-1
      P(NP)=PAR(I+175)
55  P(NP+1)=PAR(I+150)
```



ELST4010  
 ELST4020  
 ELST4030  
 ELST4040  
 ELST4050  
 ELST4060  
 ELST4070  
 ELST4080  
 ELST4090  
 ELST4100  
 ELST4110  
 ELST4120  
 ELST4130  
 ELST4140  
 ELST4150  
 ELST4160  
 ELST4170  
 ELST4180  
 ELST4190  
 ELST4200  
 ELST4210  
 ELST4220  
 ELST4230  
 ELST4240  
 ELST4250  
 ELST4260  
 ELST4270  
 ELST4290  
 ELST4300  
 ELST4310

```

56 LVE=LPAR(4)
   BREITEN EQ,0) GO TO 61
   IF(LBV EQ,1, LBV
   DC 60 I=1, LBV
   NP=NST+LFE+2*LVE+I
   P(NP)=PAR(I+125)
60 M=NST+LFE+2*LVE+LBV
61 LPAR(29)=M
   LPAR(24)=1
C
   NUM=LPAR(1)+1
   IF(LPAR(20).GT.1) CALL ZETKO
   CALL FITUN(M,P,LEO,XX,YY,ERR,-1,GAMMA,NUM,&80)
   IF(LPAR(20).GT.1) CALL ZETKO
C
   IF(LPAR(31).LE.0) GO TO 75
   DC 70 J=1,IM
   FEH(J)=PAR(31)-PAR(J+150)
   NP1=(NST+LFE)*2+200+1
   NP=NP1+4*J-2
   IF(J.LE.LVE) ZAEH(J)=PAR(31)-PAR(NP)
   IF(J.GT.LVE.AND.LPAR(3).GT.0) ZAEH(J)=FEH(J)
70 IF(J.GT.LVE.AND.LPAR(3).LT.0) ZAEH(J)=FEH(J)-PAR(NP1+2)+PAR(151)
   WRITE (1) (FEH(J),J=1,IM)
   WRITE (1) (ZAEH(J),J=1,IM)
75 IF(LPAR(20).GT.1) CALL ZETKO
   CCONTINUE
80 RETURN
   END
C
   SLBROUTINE UNTF(X,P,DY,UNT,FSSW)
   DIMENSION P(50),DY(50)
   COMMON PAR(300),LPAR(100)
   IF(PAR(24).NE.0) PAR(45)=PAR(24)
   A=X-PAR(45)
   LPAR2=LPAR(2)
   IF(LPAR2.GT.50) LPAR2=LPAR2-50
   IF(LPAR(2).GT.50) A=1./(-A)
   CALL SSWF(FSSW,X)
   NPJ=INT(FLOAT(LPAR(2))/10.)
   NPA=LPAR(2)-NPJ*10
C
   UNT=.0
   DC 5 I=1,NPA
   UNT=UNT+P(I)*A**(I-1)
5 DY(I)=A**(I-1)
   GC TO (10,20,10,20), NPJ
  
```



```

10 UNT=UNT+FSSW*P(NPA+1)
DY(NPA+1)=FSSW
IF(LPAR2.LE.10) RETURN
GOTO 30
20 UNT=UNT+FSSW
IF(LPAR2.LE.20) RETURN
30 UNT=UNT*EXP(-P(NPA+2)*(PAR(31)-X)/X)
DY(NPA+2)=-UNT*(PAR(31)-X)/X
RETURN
END

SUBROUTINE FITUN(MZ,P,N,X,Y,DY,NRM,DELTA,NUM,*)
NRM = 0 FALLS DIE MESSFEHLER ABSOLUT SIND
DIMENSION PAR(300),LPAR(100),DUM(600), WORD(20)
DIMENSION X(500),Y(500),DY(500),V1(50),V2(50),
1 P(50),DPTA(25), DZ(50), DZ2(50,50), U(50,50)
DIMENSION JOTA(25)
DIMENSION ZWP(50),PNB(50) WORD
COMMON NZP/LEO,LX,LV2
COMMON /XKURV/YKURV(10)
COMMON /LING1/DZ2,U
REAL*8 TEXT; NAME1,NAME2,NAME3
DATA TEXT / ' GRD' /
DATA NAME3 / ' GRD' /
GAM1=.0
NORM = NRM
GAMMA=DELTA
LEG=N
LV2=0
M=MZ
MY=5
J006=6
M00Z=30
LFIT=IABS(LPAR(3))
LFIT1=LPAR(4)
LNKURV=LFIT+LFIT1
K001=LPAR(22)+1
LPAR(31)=0
M001=2
DC 5 K=1,50
DC 5 L=1,50
DZ(K,L)=.0
5 U(K,L)=.0
DIAGNOSTIK
I1=1
IF(NUM-4) 11,13,11
11 LZ=M-LPAR(6)

```

```

UNST0030
UNST0050
UNST0090
UNST0100
UNST0110
UNST0120
UNST0150
UNST0160
UNST0170
UNST0180
UNST0190
UNST0200
UNST0210
UNST0220
UNST0230
UNST0240
UNST0250
UNST0260
UNST0270
UNST0280
UNST0290
UNST0300
UNST0310
UNST0320
UNST0340
UNST0360
UNST0370
UNST0380

```





```

UNST 0390
UNST 0400
UNST 0410
UNST 0420
UNST 0430
UNST 0440
UNST 0450
UNST 0460
UNST 0470
UNST 0480
UNST 0490
UNST 0500
UNST 0510
UNST 0520
UNST 0530
UNST 0540
UNST 0550
UNST 0560
UNST 0570
UNST 0580
UNST 0590
UNST 0600
UNST 0610
UNST 0620
UNST 0630
UNST 0640
UNST 0650
UNST 0660
UNST 0670
UNST 0690
UNST 0700
UNST 0710
UNST 0720
UNST 0740
UNST 0760
UNST 0770
UNST 0780
UNST 0790
UNST 0800
UNST 0810
UNST 0820
UNST 0830
UNST 0840
UNST 0850
UNST 0860
UNST 0870
UNST 0880
UNST 0890

IF(LPAR(11).GT.0) M001=LPAR(11)
IF(LPAR(12).GT.0) M002=LPAR(12)
IF(LPAR(13).GT.0) MY=LPAR(13)
IF(LPAR(14).GT.0) J006=LPAR(14)
DC 8 I=1,LFITI
JCTA(I)=LFIT+2*I+LPAR(22)
13 IF(M) 16,16,17,17,16
15 IF(M-50) 17,17,16
16 WRITE(6,997)M
997 FORMAT(40H1 NUR 1 BIS 50 PARAMETER ERLAUBT STATT ,I7)
RETURN
17 IF(N-500) 20,20,18
18 WRITE(6,998)N
998 FORMAT(42H1 NUR 1 BIS 500 MESSPUNKTE ERLAUBT STATT ,I7)
RETURN
20 ZSOLL = N - M
IF(ZSOLL) 24,25,26
24 WRITE(6,999)M,N
999 FORMAT(19H1 MEHR PARAMETER (, 12, 20H) ALS MESSPUNKTE (, 12,
1 IH)
RETURN
25 NORM = 0
DC 27 I = 1,N
26 IF(DY(I)) 28,28,27
27 CONTINUE
GC TO 14
28 WRITE(6,1000)I
1000 FORMAT(49H1 MESSFEHLER DY(I) IST NULL ODER NEGATIV FUER I = , I4)
RETURN
14 FEH=1.0
IF(NORM) 29,30,30
29 FEH=DY(1)**2
30 J BEGINN DER ITERATION ZUR LOESUNG DER NORMAL - GLEICHUNGEN
IF(LPAR(8)) 31,31,32
31 WRITE(6,1002)
1002 FORMAT(17H1,17X,18HZWISCHENERGEBNISSE, 1.PARAMETER 2.PARAMETER
1 /85H NAEHERUNG ZWANG //)
2 3.PARAMETER USW.....
32 G = 1.0
DATA NAME1/./
TEXT=NAME1
C 50 Z = 0.0
DC 60 K = 1,M
DZ(K) = 0.0
DC 60 L = 1,M

```



```

60 DZZ(K,L) = 0.0
DC 70 THEORU I = 1, NUM, X(I), P, F, DF)
CALL DY(I)**(-2)*FEH
H1 = Y(I) - F
H2 = Z + H1*H2*H2
DC 70 K = 1, M
DZ(K) = H1*H2*DF(K)
DO 70 L = 1, M
DZZ(K,L) = DZZ(K,L) + H1*DF(K)*DF(L)
71 ZY=Z/FEH 72, 73, 72
72 G=Z/ZSOLL 74, 74, 78
73 IF(LPAR(8)) 77, 77, 75
74 IF(NORM) 77, 1003)J, TEXT, Z, (P(K), K = 1, M)
75 WRITE(6,1004)J, TEXT,ZY, (P(K), K = 1, M)
1003 WCRMAT(IX,I3,A6,1PE14.6,3X,7E15.6/(27X,7E15.6))
GC TO 78
77 WRITE(6,1004)J, TEXT,ZY, (P(K), K = 1, M)
1004 FORMAT(IX,I3,A6,1PE14.6,3X,7E15.6/(27X,7E15.6))
C
78 IF(J-MY) 100, 80, 80
80 IF(Z-ZALT) 90, 101, 101
101 IF(Z-ZALT-1.0E-02) 90, 101, 91, 102, 102
91 IF(NDUR) 140, 140, 102
102 DC 103 K=1, M
103 P(K)=ZWP(K)GAM1
GAMMA=0.8*GAM1
WRITE(6,1040)GAMMA
WCRMAT(13X,10H GAMMA = F13.6)
90 IF(J-M00Z) 100, 130, 130
C 100 DURCHFUEHRUNG DER ITERATION
DC 104 K = 1, M
DP(K) = - DZ(K)
DO 104 L = 1, M
U(K,L) = DZZ(K,L)
CALL LINGL(U,DP, M, 1, DETERM)
IF(LPAR(20).GT.2) CALL ZETKO
IF(DETERM) 105, 107, 105
105 H1 = 0.0 K = 1, M
DO 106 K = 1, M
H1 = H1 + DP(K)*DP(K)*DZZ(K,K)
DATA NAME2/1. NWT, /
TEXT=NAME2
GAMI=SQRT(H1)
IF(GAMI-GAMMA) 113, 113, 107
C 107 GRADIENTEN METHODE STATT NEWTON
TEXT=NAME3

```

```

UNST0900
UNST0910
UNST0920
UNST0930
UNST0940
UNST0950
UNST0960
UNST0970
UNST0980
UNST0990
UNST1000
UNST1010
UNST1020
UNST1030
UNST1040
UNST1050
UNST1060
UNST1070
UNST1080
UNST1090
UNST1100
UNST1110
UNST1130
UNST1140
UNST1160
UNST1170
UNST1180
UNST1190
UNST1200
UNST1210
UNST1220
UNST1240
UNST1260
UNST1270
UNST1280
UNST1290
UNST1310
UNST1320
UNST1330
UNST1340
UNST1350
UNST1360
UNST1370
UNST1380
UNST1390
UNST1410
UNST1430

```



```

H1 = 0.0
H2 = 0.0
DC 109 K = 1, M
H3 = DZ(K)
DZ(K) = H3/D2Z(K, K)
H1 = H1 + H3*DZ(K)
DC 110 K = 1, M
DC 110 L = 1, M
H2 = H2 + DZ(K)*D2Z(K, L)*DZ(L)
H3 = AMIN1(H1/H2, GAMMA/SQRT(H1))
DC 112 K = 1, M
DP(K) = - H3*DZ(K)
ZALT = Z
IF(J.LE.M001) GO TO 114
I1=I1+2
DC 120 K = 1, M
ZWP(K)=P(K)
P(K) = P(K) + DP(K)
DC 123 K=K001, M
IF(I1-25) I16, I16, I15
NDUR=0
DC 118 KY=I1, 25
IF(K-JOTA(KY)) I18, I17, I18
NDUR=1
GO TO 119
CONTINUE
IF(K-(L*PAR(22)+LFIT +2*L*PAR(4))) I25, I25, I26
IF(Z-ZSOLL*FEH#3.0) I28, I28, I27
IF(J-J006) I19, I19, I28
P(K)=ZWP(K)
NDUR=1
GO TO 123
IF(ABS(DP(K))-0.15*ZWP(K)) I25, I25, I29
P(K)=ZWP(K)+.10*DP(K)*ZWP(K)/ABS(DP(K))
WRITE(6, I042) K
FORMAT(14X, I7HBREITE(PARAMETER 12, 22H) UM 10 PROZ. GEAEEND.
)
IF(P(K)) I22, I22, I23
WRITE(6, I041) K, P(K)
FCRMMAT(10X, I5, 20H. PARAMETER NEGATIV( E15.6, 3H ))
P(K)=ZWP(K)
CONTINUE
J=J+1
GO TO 50
WRITE(6, I005) M00Z
FCRMMAT(24H, I6) (KEINE KONVERGENZ NACH I3, I5H ITERATIONEN)
WRITE(6, I030) (WORD(I), I=1, I8)
FCRMMAT(14H, I18A4//)
WRITE(6, I032)

```

```

UNST11440
UNST11450
UNST11460
UNST11470
UNST11480
UNST11490
UNST11500
UNST11510
UNST11520
UNST11530
UNST11540
UNST11550
UNST11570
UNST11580
UNST11590
UNST11600
UNST11610
UNST11620
UNST11630
UNST11640
UNST11650
UNST11660
UNST11670
UNST11680
UNST11690
UNST11700
UNST11710
UNST11720
UNST11730
UNST11740
UNST11750
UNST11760
UNST11770
UNST11780
UNST11790
UNST11800
UNST11810
UNST11820
UNST11830
UNST11840
UNST11850
UNST11860
UNST11870
UNST11900
UNST11910
UNST11940
UNST11950
UNST11960

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```

1032 FORMAT(23H FITERGEBNISSE SPEKTRUM )
133 IF(NORM) 35,35,40
C
C ERWARTUNGSWERT VON Z IST ZSOLL = N - M FALLS DIE MESSFEHLER DY(I)
C ABSOLUT SIND. STREUUNG IST SQRT (2*ZSOLL)
C
35 H1 = SQRT(2.0*ZSOLL)
A012=H1
B012=Z
PAR(42)=ZSOLL
PAR(43)=A012
PAR(44)=B012
IF(LPAR(20).GT.2) CALL ZETKO
WRITE(6,1007)ZSOLL, H1, ZY
FCRMAI(9H HISTWERT, F8.1, 3H +- , F5.1
/18X,7HISTWERT, F9.1)
40 WRITE(6,1008)
1008 FCRMAT(1H,32X,7HOPTIMAL,13X,6HFEHLER,
28X,27HKONTROLLE (SOLLWERT IST 1) //)
DO 160 K = 1,M
DO 150 L = 1,M
U(K,L) = 0.0
U(K,K) = G
CALL LINGL (D2Z, U, M, M, DETERM)
IF(LPAR(20).GT.2) CALL ZETKO
IF(DETERM) 170, 210, 170
170 DC 180 K = 1,M
180 DP(K) = SQRT(ABS(U(K,K)))
DETERM = 1.0/DETERM
KONTROLLE DURCH BERECHNUNG DES ZWANGES AN DEN EXTREMPUNKTEN
DES FEHLERELLIPSOIDS
DC 205 K = 1,M
H1 = 0.0
H2 = 0.0
IF(LPAR(19).GT.0) GO TO 205
DC 190 L = 1,M DP(K)=1. E-15
IF(DP(K).EQ.0.0) DP(K)
U(K,L) = U(K,L)/DP(K)
V1(L) = P(L) - U(K,L)
V2(L) = P(L) + U(K,L)
IF(DP(L).EQ.0.0) DP(L)=1. E-15
U(K,L) = U(K,L)/DP(L)
DC 195 I = 1,N
DO 195 THEORU (NUM,X(I),V1,F,DF)
CALL H1 + ((Y(I) - F)/DY(I))**2*FEH
H1 = H1 + (H1 - Z)/G
DC 200 I = 1,N
CALL THEORU( NUM,X(I),V2,F,DF)
UNST1970
UNST1980
UNST1990
UNST2000
UNST2010
UNST2020
UNST2030
UNST2040
UNST2050
UNST2060
UNST2070
UNST2080
UNST2090
UNST2100
UNST2110
UNST2120
UNST2130
UNST2140
UNST2150
UNST2160
UNST2170
UNST2180
UNST2190
UNST2210
UNST2220
UNST2230
UNST2240
UNST2250
UNST2270
UNST2280
UNST2300
UNST2310
UNST2320
UNST2330
UNST2340
UNST2360
UNST2370
UNST2380
UNST2400
UNST2410
UNST2420
UNST2430
UNST2440
UNST2450
UNST2460

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```

200 H2 = H2 + ((Y(I) - F)/DY(I))**2*FEH
H2 = (H2 - Z)/G
205 WRITE (6,1009)K, P(K), DP(K), H1, H2
1009 FCORMAT (1X,I2,10H.PARAMETER,15X,IPE15.6,5H +--,E15.6,11X,0P2F20.5
)
1 IF(LPAR(19).GT.0) GO TO 333
WRITE (6,1010)
FCORMAT (40H KORRELATIONEN ZWISCHEN DEN PARAMETERN
//12X,31H1.PA 2.PA 3.PA USH....
//)
1 DC 185
WRITE (6,1011)K,(U(K,L), L = 1,M)
FCORMAT (1X,I2,6H.PA ,15F8.3,3(/9X,15F8.3) )
1011 FCORMAT (6,1014)DETERM
1014 FCORMAT (1H0,11X,8HDETERM =,1PE11.3 //)
C TABELLIERUNG DER ANGEPASSTEN FUNKTION
1 IF(LPAR(20).GT.2) CALL ZETKO
333 LPAR(19)=0 181,186,181
IF(LPAR(8)) 216,216,182
181 IF(LPAR(9))
186 IF(LPAR(6,1015))
1015 FCORMAT (1H0,54X,13HT A B E L L E //29X,14HZAEHLRATE
17H UNR. FREQUENZ,7X,8HGEMESSEN,8X,9HBERECHNET,8X,
7HY-UNT,8X,7HUNT WRITE(6,1019) ,8X,9HDIFFERENZ,7X,6HFEHLER//)
2 IF(PAR(64).GE.1.0) WRITE(6,1019)
1019 FCORMAT(1X,/, ZAEHLRATE AUF DISPERSION KORRIGIERT *)
182 IF(LPAR(19)) 2006,2006,187
187 WRITE (1) LPAR(24),M,N,ZSOLL,A012,B012
LPAR(25)=LPAR(25)+1
IF(LPAR(31)=1) NE.1) GO TO 188
WRITE(7,2001) (WORD(I),I=1,20)
2006 WRITE(7,2002) PAR(7),PAR(59),PAR(1)
WRITE(7,2003) (PAR(2*I+89),I=1,5),PAR(81),LPAR(16)
WRITE(7,2004) (P(I),I=1,4);PAR(45),PAR(24),LPAR(2)
FCORMAT(20A4)
2001 FCORMAT(4E12.4)
2002 FCORMAT(6E12.4,I3)
2003 FCORMAT(6E12.4,I3)
188 LV2=1
DO 213 I = 1,N
LX=1
CALL THEORU (NUM ,X(I),P,F,DF)
H1=Y(I)
H14=H1/DY(I)
CALL UNIF(X(I),P,DY,UNT,FSSW)
H7=Y(I)-UNT
H2=Y(I)-FSSW
H22=UNT-FSSW
197

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UNST2470
UNST2480
UNST2490
UNST2500
UNST2510
UNST2520
UNST2530
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UNST2580
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UNST2600
UNST2620
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UNST2650
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UNST2680
UNST2690
UNST2700
UNST2710

UNST2740
UNST2750
UNST2760
UNST2770

UNST2800
UNST2810
UNST2820

UNST2840
UNST2850
UNST2860
UNST2870
UNST2880
UNST2890
UNST2900
UNST2910
UNST2920

UNST3150
UNST3160
UNST3170

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UNST3180  
UNST3190  
UNST3210  
UNST3220  
UNST3230  
UNST3240  
UNST3250

UNST3280  
UNST3290  
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UNST3490  
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UNST3550  
UNST3560  
UNST3570  
UNST3580  
UNST3590  
UNST3600  
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UNST3620  
UNST3630  
UNST3640  
UNST3650  
UNST3660  
UNST3670  
UNST3680  
UNST3690

```
H13=F-UNT  
IF(LPAR(9)) 215,215,189  
WRITE(1)X(I),H7,H13,H14,UNT,(YKURV(K),K=1,NKURV),FSSW  
IF(LPAR(8)) 213,213,214,213  
WRITE(6,1016)I,X(I),Y(I),F,H7,UNT,FSSW,H1,DY(I)  
FORMAT(15,F12.4,2X,1P(2(E17.6,E15.6),E17.6,E15.6))  
IF(LPAR(38)).NE.2) GO TO 213  
H31=F-FSSW  
FORMAT(15,6(1PE12.4))  
WRITE(7,2005) I,X(I),Y(I),DY(I),F,UNT,FSSW  
CONTINUE  
C 213  
DO 1022 I=1,M  
PAR(2*I+199)=P(I)  
PAR(2*I+200)=DP(I)  
IF(LPAR(20)).GT.2) CALL ZETKO  
RETURN  
FALLS SICH DAS LINEARE GLEICHUNGSSYSTEM NICHT AUFLÖSEN LAESST  
1012 WRITE(6,1012) NORMAL-GLEICHUNGEN SIND NICHT AUFLÖSBAR //  
FCRMAT(42HO,NORMAL-GLEICHUNGEN SIND NICHT AUFLÖSBAR //)  
DO 220 K=1,M  
220 WRITE(6,1013)(U(K,L), L=1,M)  
1013 FCRMAT(1X,1P7E15.7)  
245 RETURN  
END
```

```
SUBROUTINE LINGL(A,B,N,M,DETERM)  
AUFLÖSUNG LINEARER GLEICHUNGSSYSTEME BIS ZUR GRDNUNG 25  
DIMENSION A(50,N),B(50,M)  
60 IF(N) 15,15,16  
16 IF(M) 15,15,17  
17 DETERM=1.0  
AMAX=ABS(A(1,1))  
LP=1  
DO 2 I=2,N  
IF(ABS(A(I,1))-AMAX) 2,2,1  
1 AMAX=ABS(A(I,1))  
LP=I  
CONTINUE  
C 2  
DO 11 J=1,N  
IP=LP  
T=A(IP,J)  
IF(T) 3,15,3  
3 IF(J-N) 4,8,8  
4 A(IP,J)=A(J,J)  
JPI=J+1
```



```

LP = JPI
AMAX = 0.0
DC 7 K=JPI,N
H = A(IP,K)/T
A(IP,K) = A(J,K)
A(J,K) = H
DO 7 I=JPI,N
  A(I,K) = A(I,K) - A(I,J)*H
  IF(K-JPI) 5,5,7
  IF(ABS(A(I,K))-AMAX) 7,7,6
5 AMAX = ABS(A(I,K))
6 LP = I
7 CCNTINUE
C
8 DC 11 K=1,M
  H = B(IP,K)/T
  B(IP,K) = B(J,K)
  B(J,K) = H
  IF(J-N) 9,11,11
9 DO 10 I=JPI,N
  B(I,K) = B(I,K) - A(I,J)*H
10 B(I,K) = B(I,K) - A(I,J)*H
11 CCNTINUE
C
12 IF(N-1) 15,15,12
DC 13 I = 2,N
  IP=N-I+1
  LP=IP+1
DC 13 J=1,M
  DC 13 K=LP,N
  B(IP,J) = B(IP,K)*B(K,J)
13 B(IP,J) = B(IP,K)*B(K,J)
15 RETURN
END
C
SUBROUTINE THEORU(J,X,P,Y,DA)
DIMENSION P(50),DA(50)
CALL ARCS(X,P,Y,DA)
RETURN
END
C
SUBROUTINE ARC4U(X,P,Y,DY,I)
DIMENSION P(25),DY(25),PVAR(25)
COMMON PAR(300),LPAR(100)
BT=1.44269505
IF(P(25)-1.0) 5,6,5
5 P(6)=1.0
6 U=(X-P(5))/P(6)
UNST3700
UNST3710
UNST3720
UNST3730
UNST3740
UNST3750
UNST3760
UNST3770
UNST3780
UNST3790
UNST3800
UNST3810
UNST3820
UNST3830
UNST3840
UNST3850
UNST3860
UNST3870
UNST3880
UNST3890
UNST3900
UNST3910
UNST3920
UNST3930
UNST3940
UNST3950
UNST3960
UNST3970
UNST3980
UNST3990
UNST4000
UNST4010
UNST4020
UNST4040
UNST4060
UNST4100
UNST4110
UNST4120
UNST5620
UNST5630
UNST5640
UNST5650
UNST5660
UNST5670
UNST5680
UNST5690

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```

IF(PAR(64).GE.1. .AND. I.LE.LPAR(6)) GOTO 110
UNST5700
UNST5710
UNST5720
UNST5730
UNST5740
UNST5750
UNST5760
UNST5770
UNST5780
UNST5790
UNST5800
UNST5810
UNST5820
UNST5830
UNST5840
UNST5850
UNST5860
UNST5870
UNST5880
UNST5890
UNST5900
UNST5910
UNST5920
UNST5930
UNST5940
UNST5950
UNST5960
UNST5970
UNST5980
UNST5990
UNST6000
UNST6010
UNST6020
UNST6030
UNST6040
UNST6050
UNST6060
UNST6070
UNST6080
UNST6090
UNST6100
UNST6110
UNST6120
UNST6130
UNST6140
UNST6150
UNST6160
UNST6170

10  IF(U) 20,10,10
    CCNTINUE
    ALPHA=(U/P(2))**2/BT
100  FORMAT(5E15.4)
11   Y=P(1)*EXP(-ALPHA)
    DY(1)=Y/P(1)
    DY(2)=2.0*Y*U**2/(BT*P(2)**3)
    DY(3)=0.0
    DY(4)=0.0
    DY(5)=2.0*U*Y/(BT*P(2)**2)/P(6)
    DY(6)=U*DY(5)
    RETURN
20   YP=P(1)*EXP(-(P(4)/P(3))**2/BT)
    IF(U-P(4)) 42,30,30
30   Y=P(1)*EXP(-(U/P(3))**2/BT)
    DY(1)=Y/P(1)
    DY(2)=0.0
    DY(5)=2.0*Y*U/(BT*P(3)**2)/P(6)
    DY(4)=0.0
    DY(3)=DY(5)*U/P(3)*P(6)
    DY(6)=U*DY(5)
    RETURN
42   PP=2.0*(P(4)/P(3))**2/BT
    C=Y*P(4)**(2.0-PP)
    D=Y*P(4)**2*(PP-1.0)
    PPP=Y*P*PP
    CFB=PPP*P(4)*(4.0-PP)/P(3)
    DFB=PPP*P(4)**2/P(3)-3.0*BT*P(3)/2.0)
    CP4=PPP*(2.0/PP-5.0*PP)
    DF4=PPP*P(4)*(5.0-PP-2.0/PP)
    Y=C/U+D/U**2
    DY(1)=Y/P(1)
    DY(2)=0.0
    DY(3)=CP3/U+DP3/U**2
    DY(4)=CP4/U+DP4/U**2
    DY(5)=C/U**2+2.0*D/U**3/P(6)
    DY(6)=U*DY(5)
    RETURN
C    BREIT-WIGNER-KURVE MIT FESTEM P-FAKTOR (DISS PIT 72) AB 110
C    MIT GEFITTETEM P-FAKTOR AB 120
C1=U/P(2)
C2=1.+C1**2
IF(PAR(64).GT.1.) GOTO 120
Y=P(1)/C2
Y=P(1)+PVAR(I)*X-P(5)
DY(1)=Y/P(1)
DY(2)=2.*P(1)*C1**2/P(2)/C2**2

```





UNST 6180  
 UNST 6190  
 UNST 6200  
 UNST 6210  
 UNST 6220  
 UNST 6230  
 UNST 6240  
 UNST 6250  
 UNST 6260  
 UNST 6270  
 UNST 6280  
 UNST 6290  
 UNST 6300  
 UNST 6310  
 UNST 6320  
 UNST 6330  
 UNST 6340  
 UNST 6350  
 UNST 6360  
 UNST 6380  
 UNST 6390  
 UNST 6420  
 UNST 6430  
 UNST 6460  
 UNST 6900  
 UNST 6910  
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 UNST 6970  
 UNST 6980  
 UNST 6990  
 UNST 7000  
 UNST 7010  
 UNST 7020  
 UNST 7030  
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 UNST 7070  
 UNST 7080  
 UNST 7090

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DY(3)=2.*P(1)*C1**2/P(3)/C2**2
DY(4)=0
DY(5)=2.*DY(2)*C1/P(6)
DY(6)=DY(5)*U
RETURN
IF(PAR(64).GT. 2.) GOTO 130
FLOR=P(1)/C2
Y=P(1)*(1+P(4)*U)/C2
DY(1)=Y/P(1)
DY(2)=2.*FLOR*(1+P(4)*U)*C1**2/P(2)/C2
DY(3)=DY(2)*P(2)/P(3)
DY(4)=FLOR*U
DY(5)=(-FLOR*P(4)+2.*DY(2)*C1)/P(6)
DY(6)=DY(5)*U
RETURN
CONTINUE
130 END

SUBROUTINE ARC5(X,P,Y,DY)
DIMENSION P(50),DY(50),PAR(300),LPAR(100),DUM(600)
DIMENSION PW(25),DYW(25)
DIMENSION BVAR(25)
COMMON PAR,LPAR,DUM
COMMON /BVAR1/ BVAR
COMMON /NZP/ LEO,LX,LV2
COMMON /XKURV/ YKURV(10)
INTERGRUND
CALL UNTF(X,P,DY,UNT,FSSW)
Y=UNT
MZH=LPAR(29)
LFIT=IABS(LP(3))
LFIT1=LPAR(4)
L=LPAR(5)
ME=LPAR(28)+LPAR(4)
KZ=M+LPAR(4)+L+LPAR(22)
PW(2)=PAR(73)
PW(3)=PAR(75)
PW(4)=PAR(77)
PW(25)=1.0
DPY(KZ)=0.0
DPY=0.0
NP=LFIT+LPAR(22)+2
DEPS=P(NP)-PAR(151)
DO 120 I=1,M
K=150+LPAR(4)+I
INS=I+LPAR(22)
KB=M+LPAR(4)+LFIT1+INS

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UNST71100  
UNST71110  
UNST71120  
UNST71130  
UNST71140  
UNST71150  
UNST71160  
UNST71170  
UNST71180  
UNST71190  
UNST72200  
UNST72210  
UNST72220  
UNST72230  
UNST72240  
UNST72250  
UNST72260  
UNST72270  
UNST72280  
UNST72290  
UNST73310  
UNST73320  
UNST73330  
UNST7370  
UNST7380  
UNST7390  
UNST7400  
UNST7410  
UNST7420  
UNST7430  
UNST7440  
UNST7450  
UNST7460  
UNST7470  
UNST7480  
UNST7490  
UNST7500  
UNST7510  
UNST7520  
UNST7530  
UNST7540  
UNST7550  
UNST7560  
UNST7570  
UNST7580  
UNST7590

```

35 KC=2*I-LFIT +LPAR(22)+1
36 KD=M+LPAR(4) -LFIT+INS
    EPS=PAR(K)
    IF(LPAR(3)) DEPS
    EPS=PAR(K)+DEPS
37 EPW(1)=P(INS)*PAR(71)
38 PW(5)=PAR(79)-EPS
39 PW(6)=PAR(K-25)
40 IF(LPAR(5)) 55,55,40
    IF(LFIT+I-L) 45,45,55
41 IF(LFIT)=P(KB) 80,80,60
42 IF(I-LFIT) 80,80,60
43 IF(1)=P(KC)*PAR(71)
44 PW(5)=PAR(79)-P(KC+1)
45 PW(6)=PAR(5) 75,75,65
65 IF(LPAR(KZ)) 70,70,80
    IF(I-LFIT-L)
    PW(6)=P(KD)
70 GC TO 80
71 KW=I-LFIT+125
72 PW(6)=PAR(KW)
C
80 CALL ARC4U(X,PW,YK,DYM,I)
    IF(I.LE.10 .AND. LV2.EQ.1) YKURV(I)=YK
    Y=Y+YK
81 IF(I-LFIT) 84,84,85
82 DY(INS)=DYM(1)*PAR(71)
83 CPY=DPY-DYW(5)
84 GC TO 90
85 DY(KC)=DYW(1)*PAR(71)
    DPY=DPY+DYW(2)
    DY(KC+1)=-DYW(5)
90 IF(LPAR(5)) 120,120,95
91 IF(I-LFIT) 100,100,130
100 IF(I-LFIT) 105,105,110
105 DY(KZ)=DY(KZ)+DYW(6)
    GC TO 120
110 IF(I-L+LFIT) 115,105,105
115 GC(KB)=DYW(6)
130 IF(I-LFIT-L) 135,140,140
135 DY(KD)=DYW(6)
    GC TO 120
140 DY(KZ)=DY(KZ)+DYW(6)
120 CONTINUE
150 IF(LPAR(3)) 150,160,160
    LJT=LFIT+LPAR(22)+2

```



UNST7600  
UNST7610  
UNST7620  
UNST7630

AUSG0020  
AUSG0050  
AUSG0060  
AUSG0080  
AUSG0090  
AUSG0100

AUSG0110  
AUSG0120  
AUSG0130  
AUSG0140  
AUSG0150  
AUSG0160  
AUSG0170  
AUSG0180  
AUSG0190  
AUSG0200  
AUSG0210  
AUSG0230

AUSG0250  
AUSG0260  
AUSG0390

AUSG0430  
AUSG0440  
AUSG0450  
AUSG0460  
AUSG0470  
AUSG0480  
AUSG0490  
AUSG0500  
AUSG0510  
AUSG0520  
AUSG0530  
AUSG0540  
AUSG0550  
AUSG0560  
AUSG0570  
AUSG0580  
AUSG0590  
AUSG0600

DY(LJT)=DY(LJT)+DPY  
CCNTINUE  
RETURN  
END

160

SLBROUTINE AUSG  
DIMENSION EPS(25), DEPS(25), BR(25), DBR(25), PIP(25), DPIP(25)  
DIMENSION FIF(25), DFIF(25)  
DIMENSION PAR(300), LPAR(100), DUM( 600), WORD(20)  
COMMON PAR, LPAR, DUM, WORD  
DIMENSION BVAR(25)  
COMMON/BVARI/BVAR  
COMMON/SSWD/ACO(25)  
DATA A001/4H

C

1000 FCRMAT(1H, 10HERGEBNISSE /1H, 18A4, F8.0, 2H /, F8.1, 2H /, F9.1 )  
1001 FCRMAT(1H, 34HELASTISCH (BERECHNET MIT A4, 4H )  
1002 FCRMAT(1H, 20HENERGIE ELAST. PEAK 14X, F9.4, 7H MHZ = F13.4, 4H MEV )  
1003 FCRMAT(1H, 15HPRIMAERENENERGIE, 19X, F9.4, 7H MHZ = F13.4, 4H MEV )  
1004 FCRMAT(1H, 23HHALBWEITENBREITE ELAST., 11X, 3PF9.4, 7H KHZ = F13.4,  
1 8H KEV = F8.2, 12H PROMIL EE )  
1005 FCRMAT(1H, 7HFLAECHE, 24X, F12.4, 7H MHCO = F13.4, 6H MEVCO )  
1006 FCRMAT(1H, 9HPEAKHOEHE, 22X, F12.4, 5H Z/AS )  
1007 FCRMAT(1H, 43HUNTERINTEGRATIONSGRENZE BEI EE - 2\*HNB )  
1008 FCRMAT(1H, 12HUNELASTISCH //35H STRAHLUNGSSCHWANZ  
1009 FCRMAT(1H, 17H SSW = 1PE12.4, 3H +E11.4, 8H\*U\*\*2 +E11.4, 8H\*U\*\*3  
2 +E11.4, 6H\*U\*\*4 )  
1010 FCRMAT(1H, 10HUNTERGRUNDSPARAMETER=, 7(1PE12.2) )  
1011 FCRMAT(1H, 22HUNELASTISCHE LINIEN / VARIABEL, F FEST, FB ENER  
1013 FCRMAT(1H, 11H3H (EPSO VORGEGEBEN, EPST --- BRV VORGEBENES /  
2 GIEBSTAND ZU EPST(1) BREITENVERHAELTNIS, U BREITE UNABHAENGIG, G  
3 110H BRB VAREIERT, F FEST --- Z PEAKHOEHE, HEN / IN KLAMMER  
4 GEMEINSAM VARIERT, AELTNIS, F/F FLAECHE NVERHAELTNIS --  
5 84H P/P HOEHENVERH, AELTNIS, F/5H  
6 FEHLER IN PROZENT  
7 6H EPST, 7X, 4HEPST, 17X, 3HBRV, 6X, 3HBRB, 16X, 1HZ, 10X, 3HP/P, 22X, 3HF/F )  
1014 FCRMAT(17.3, F9.4, 3H +-F7.4, A3, A2, F7.3, F8.3, 3H +-F7.4, 1PE11.2,  
1 E13.4, 3H +-E9.2, E13.4, 3H +-E9.2, 2H ( 2PF7.2, 2H )

C

NST=LPAR(22)  
LFI=LPAR(3)  
LFI=LPAR(4)  
LBI=LPAR(5)  
LBI=LPAR(21)  
DFWB=PAR(48)



```

WRITE (6,1000) (WORD(I),I=1,18), (PAR(I),I=42,44)
DATA ARC3/4HARC3 /
DATA ARC4/4HARC4 /
DATA EFT/4H F /
DATA EVA/4H V /
DATA EFB/4H FB /
DATA BUN/4H U /
DATA BGE/4H G /
DATA BFT/4H F /
A001=ARC3
IZT=PAR(51)
WRITE (6,1015) PAR(1),PAR(9)
WRITE (6,1016) IZT,(PAR(1),I=52,59)
WRITE (6,1017) (PAR(1),I=11,14)
WRITE (6,1018) (PAR(1),I=15,20)
FCRMAT(22H FITGRENZEN 2H - F8.4)
9HUNELAST F8.4,2H - F8.4,5X,
1017 1 FCRMAT(1H,13HEINGABEBERTE ELAST. PEAK F10.5, 5X
1015 1 FCRMAT(2H Z,I4,5H,ATW F7.2,6H,XRAD F7.3,4H, RM F6.3,3H,R F7.3,
1016 1 FCRMAT(3H,G F7.3,3H,C F8.4,3H,I F8.4,12H,MASSE(MEV) F10.1)
1018 1 FCRMAT(37H NAEHERUNGSUNTERGRUND ZWISCHEN (MHZ) ,3(F11.4,2H -F8.4))
IF(LPAR(1).GT.0) A001=ARC4
WRITE (6,1001) A001
EE=PAR(31)
EEMEVEV=E*PAR(1)
EOMEV=PAR(36)
EO=PAR(36)/PAR(1)
HM=PAR(34)
HMEV=HM*PAR(1)
HWPRO=HW/EE
FL=PAR(33)
FLMEV=FL*PAR(1)
WRITE (6,1002) EE,EEMEVEV
WRITE (6,1003) EO,EOMEV,HWPRO
WRITE (6,1004) HW,HMEV
WRITE (6,1005) FL,FLMEV
WRITE (6,1006) PAR(35)
IF(LPAR(1).GT.0) WRITE (6,1008)
JGLEITABS(LPAR(16))
WRITE (6,1009) (PAR(2*I+89),I=1,5)
WRITE (6,1010) (PAR(2),PAR(45)
WRITE (6,1011) (PAR(2*I+199),I=1,NST)
WRITE (6,1012)
WRITE (6,1013)
WRITE (PAR(64).GE.1) WRITE(6,1035)
IFCRMAT(1H, FUEER BREIT-WIGNER-FORM BERECHNETAUS PI*BRB*HWMEV*Z
/FLMEV/2. BREITEN ABSOLUT IN MEV , /)
AUSG0610
AUSG0620
AUSG0630
AUSG0640
AUSG0650
AUSG0660
AUSG0670
AUSG0680
AUSG0690
AUSG0700
AUSG0710
AUSG0720
AUSG0730
AUSG0740
AUSG0750
AUSG0760
AUSG0770
AUSG0780
AUSG0790
AUSG0800
AUSG0810
AUSG0820
AUSG0830
AUSG0840
AUSG0850
AUSG0860
AUSG0870
AUSG0880
AUSG0890
AUSG0900
AUSG0910
AUSG0920
AUSG0930
AUSG0940
AUSG0950
AUSG0960
AUSG0970
AUSG0980
AUSG1000
AUSG1010
AUSG1110
AUSG1120
AUSG1130
AUSG1140
AUSG1150

```





AUSG11160  
 AUSG11170  
 AUSG11180  
 AUSG11190  
 AUSG11200  
 AUSG11210  
 AUSG11220  
 AUSG11230  
 AUSG11240  
 AUSG11250  
 AUSG11260  
 AUSG11270  
 AUSG11280  
 AUSG11290  
 AUSG11300  
 AUSG11310  
 AUSG11320  
 AUSG11330  
 AUSG11340  
 AUSG11350  
 AUSG11360  
 AUSG11370  
 AUSG11380  
 AUSG11390  
 AUSG11400  
 AUSG11410  
 AUSG11420  
 AUSG11430  
 AUSG11440  
 AUSG11450  
 AUSG11460  
 AUSG11470  
 AUSG11480  
 AUSG11490  
 AUSG11500  
 AUSG11510  
 AUSG11520  
 AUSG11530  
 AUSG11540  
 AUSG11550  
 AUSG11560  
 AUSG11570  
 AUSG11580  
 AUSG11620  
 AUSG11630  
 AUSG11640  
 AUSG11650  
 AUSG11660

```

C
140 I=1,IM
DBR(I)=PAR(I+125)
IF(LBV.EQ.0) GO TO 100
MB=2*(NST+LF)+4*LF1+2*I-1
IF(I.GT.LBV) MB=2*(NST+LF)+4*LF1+2*LBV-1
DBR(I)=PAR(MB+200)
IF(I.GT.LF1) GO TO 110
MP=2*(NST+LF)+4*I-3+200
ME=MP+2
EPS(I)=PAR(ME)
IF(PAR(ME).EQ.0) GO TO 119
DEPS(I)=PAR(ME+1)/PAR(ME)
GC TO 120
MP=2*NST+2*(I-LF1)+199
EPS(I)=PAR(I+150)
DEPS(I)=0.0
IF(LPAR(3).LT.0) EPS(I)=EPS(I)+DHWB*(BR(I))-1.0)
NPI=(NST+LF)*2+203
IF(LT.0) EPS(I)=EPS(I)+PAR(NPI)--PAR(151)
GC TO 120
DEPS(I)=0.0
DIP(I)=PAR(MP)
IF(PAR(MP).EQ.0) GO TO 125
DPIP(I)=PAR(MP+1)/PAR(MP)
GC TO 126
DPIPI=0
CONTINUE
EPS(I)=EPS(I)*PAR(I)*PAR(47)
ILAH=0
EPSNH=EPS(I)
EPS(I)=EPSNH/(1.+EPS(I)/(2.*PAR(59)))
ILA=ILA+1
IF(ILA.LE.3) GO TO 131
FIF(I)=PIP(I)*BR(I)
DFF=EPS(I)*SQRT(DPIP(I)**2+DDBR(I)**2)
DPP=CFIP(I)*FIF(I)
DDB=DPPI(I)*PIP(I)
ZM=PIPI(I)*BR(35)
TXE=EFT
IF(I.LE.LF1) TXE=EVA
IF(I.GT.LF1.AND.LPAR(3).LT.0) TXE=EFB
IF(LBV.GT.0) TXB=BUN
  
```



```

C
IF(LBV.GT.0.AND.I.GE.LBV) TXB=BGE
PAR(64)=BREIT-WIGNER-FORM
IF(PAR(64).GE.1. .AND. I.LE.LPAR(6)) GOTO 136
136 GOTO 138
CONTINUE
FIF(I)=ZM*3.1416*BR(I)*HMEV/FLMEV/2.
PAR(I+125)=PAR(I)*PAR(34)
IF(PIP(I).NE.0.) DFF=DFF*FIF(I)/(PIP(I)*BR(I))
DEB=DBB*PAR(I)*PAR(34)
BR(I)=BR(I)*PAR(1)*PAR(34)
IF(I.EQ.LPAR(6)+1) WRITE(6,1036)
1036 IF(FCRMAT(1X,/,/)) FOLGENDE LINIE SIND GAUSSKURVEN:
1037 IF(BVAR(I).NE.0.) WRITE(6,1037) BVAR(I) DER GERADEN 1+X*,F10.3,
1 MULTIPLIZIERT:
140 WRITE (6,1014) PAR(I+100),EPS(I),DEP,TXE,TXB,PAR(I+125),BR(I),
1088B,ZM,PIP(I),DPP,FIF(I),DFF,DFIF(I)
CALL RESULT(EPS,DEPS,FIF,DFIF,BR)
RETURN
END
SUBROUTINE RESULT(EPS,DEPS,FIF,DFIF,BR)
DIMENSION DI(25)
DIMENSION PAR(300),LPAR(100),DUMMY(600),WORD(20)
DIMENSION FKOR(25),FORF(25)
COMMON PAR,LPAR,DUMMY,WORD
COMMON/SSWD/ACQ(25)
IF(PAR(61).EQ.0) PAR(61)=1.
ZET=PAR(51)
ATW=PAR(52)
XO =PAR(53)
FIRA=PAR(54) FIRA=1.
IF(FIRA.EQ.0)
DEFF=PAR(41)
DEE=PAR(34)*PAR(1)
IM=LPAR(21)
DC 2 I=1,IM
DI(I)=BR(I)
CC 6 I=1,IM
IF(PAR(64).EQ.0.) DI(I)=DI(I)*DEE
CONTINUE
FCRMAT(F10.5)
1102 FCRMAT(F10.4)
1103 FCRMAT(10I3)
1104 FCRMAT(10F6.3)
1105 FCRMAT(10F5.3)
AUSG1670
AUSG1590
AUSG1780
AUSG1790
AUSG1830
AUSG1840
AUSG1880
AUSG1940
AUSG1950
AUSG1960
AUSG2020
AUSG2030
AUSG2050
AUSG2090
AUSG2100
AUSG2110
AUSG2120
AUSG2130
AUSG2140
AUSG2150
AUSG2160
AUSG2170
AUSG2180

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