1	PASSIVITY BREAKDOWN OF TITANIUM IN LIBr SOLUTIONS
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13	The passive behavior of titanium and its susceptibility to undergo localized attacks in
14	different LiBr solutions at 25° C have been investigated using different electrochemical
15	techniques: potentiodynamic polarization curves, potentiostatic passivation tests, EIS
16	measurements and Mott-Schottky analysis. In low and moderately concentrated LiBr
17	solutions, the breakdown potential $E_b$ decreased with increasing bromide
18	concentrations, while in highly concentrated LiBr solutions, $E_b$ increased with
19	increasing LiBr concentration. In the most concentrated LiBr solution (11.42M) Ti did
20	not undergo passivity breakdown even at 9 $V_{\mbox{Ag/AgCl}}.$ These results have been related to a
21	decrease in the activity of water in highly concentrated LiBr solutions.
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24	Keywords: titanium; passivity breakdown; bromide solutions; activity of water; point
25	defect model.

### 26 **1. INTRODUCTION**

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28 Titanium is a very reactive metal that shows remarkable corrosion resistance in many 29 environments due to the formation of a very stable, continuous, highly adherent, and 30 protective passive oxide film on the metal surface. A freshly abraded titanium surface 31 immediately passivates when exposed to an oxygen source (air and/or moisture). The 32 nature, composition, and thickness of the passive film that forms on titanium depend on 33 the environmental conditions. In most aqueous environments, the oxide is typically TiO<sub>2</sub>, but may consist of mixtures of other titanium oxides, including TiO<sub>2</sub>, Ti<sub>2</sub>O<sub>3</sub>, and 34 35 TiO [1-8]. The oxygen content of the titanium oxide gradually decreases from  $TiO_2$  at 36 the surface, to  $Ti_2O_3$  and TiO as it approaches the metal/oxide interface [1]. Depending 37 on the environment, this oxide may be covered with an amorphous or hydrated surface 38 oxide, giving a two-layer oxide structure [1, 4, 5, 9]. The TiO<sub>2</sub> passive film is an n-type 39 semiconductor and its semiconductive properties determine the current/potential 40 behavior of the titanium/passive film/electrolyte system [1, 6, 10-13]. The metal/passive 41 film interface has been assumed to be an ohmic contact with minimal resistance. On the 42 other hand, a space charge region has been hypothesized to develop at the passive 43 film/electrolyte interface [6]. Oxygen vacancies have been reported to be the main defects within the passive film on titanium [13-15].  $Ti^{3+}$  interstitials have also been cited 44 45 as electron donor species [10, 12].

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Titanium exhibits excellent resistance to most salt solutions over a wide range of pH and temperatures [16]. Although titanium is resistant to these media, it is not immune and can be susceptible to pitting and crevice attack, depending on the environmental conditions [15-18]. Breakdown potentials ( $E_b$ ) of titanium in sulfate and phosphate 51 media are typically in excess of 80 V<sub>SCE</sub> [1, 19]. Values of 5 to 12 V<sub>SCE</sub> can be expected 52 in room-temperature chloride solutions [1, 15, 18, 19]. However, it has been observed in 53 numerous works that bromide anions are far more effective than chlorides or iodides in 54 facilitating breakdown of the passive film formed on Ti [15, 17-22]. In fact, many 55 authors have observed passivity breakdown of titanium in bromide containing solutions 56 taking place at relatively low potentials compared to those in chloride or iodide 57 containing solutions ( $E_b < 2-5 V_{SCE}$ ) [15, 17-24]. Sazou et al. [15] interpreted this high 58 susceptibility of TiO<sub>2</sub> to local breakdown in the presence of bromides by considering 59 electrochemical reactions taking place at the passive film/electrolyte interface, such as 60 the formation of surface titanium peroxo-species, leading to changes in the surface and 61 bulk properties (ionic and electronic) of the titanium oxide. The breakdown potential of 62 titanium in relatively low concentrated bromide solutions (up to 1M) has been found to 63 decrease with increasing bromide concentrations [15, 18, 20], with values in the range of 2-5 VAg/AgCl. Nevertheless, it has been observed in recent works [2, 25] that in very 64 65 concentrated LiBr solutions (700, 850 and 992 g/l LiBr, 8.06M, 9.78M and 11.42, 66 respectively), titanium did not undergo localized attacks at potentials as high as 67  $3V_{Ag/AgCl}$ . These results contradict the tendencies found in the literature for  $E_b$  with 68 bromide concentration, since in heavy brine LiBr solutions the values of  $E_b$  should be 69 lower than 2-3 V<sub>Ag/AgCl</sub>.

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The aim of the present work is to study the influence of bromide concentration on the passivity breakdown of titanium at 25° C and to clarify the anomalous behavior of passive films formed on titanium in highly concentrated LiBr solutions against localized attacks. These heavy brine LiBr solutions are widely used in absorption cooling systems [26-30]. The steady-state properties of the passive film formed on Ti electrodes in different LiBr solutions have been studied in this work, as well as the role of electronic
properties and the activity of water on the stability of Ti passive film.

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# 79 2. MATERIALS AND METHODS

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81 Materials, solution and electrochemical cell

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The material tested was commercially pure Grade 2 titanium (0.03 wt.% N, 0.1 wt.% C, 83 84 0.25 wt.% O, 0.3 wt.% Fe, 0.0125 wt.% H, Bal. Ti). Titanium electrodes were 85 cylindrically shaped and covered with a polytetrafluoroethylene (PTFE) coating, to expose an area of  $0.5 \text{ cm}^2$  to the test solution. All specimens were wet abraded from 200 86 87 to 4000 SiC grit, and finally rinsed with distilled water and air-dried. Titanium 88 electrodes were tested in different LiBr solutions, from 0.05M to 11.42M, at a constant temperature of 25° C (Table 1). Values of molalities  $(m_{LiBr})$ , mean ionic activity 89 90 coefficients  $(\gamma_{\pm})$ , activities of LiBr solutions  $(a_{LiBr})$ , osmotic coefficients  $(\phi)$  and 91 activities of water  $(a_W)$  are also shown in **Table 1**. Activities for LiBr solutions,  $a_{LiBr}$ , 92 have been calculated from the following equation:

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94 
$$a_{LiBr} = \frac{m_{LiBr}}{m^0} \gamma_{\pm}$$
(1)

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96 where the LiBr molality,  $m_{LiBr}$ , is defined as mol LiBr/kg H<sub>2</sub>O and  $m^0$  is the standard 97 LiBr molality ( $m^0 = 1$  mol LiBr/kg H<sub>2</sub>O). Activity coefficients and osmotic coefficients 98 have been obtained from the literature [31]. Activities of water have been calculated 99 from osmotic coefficients according to the following expression [32, 33]:

$$\ln(a_w) = -\phi M_w v m \tag{2}$$

102

103 where  $M_W$  is the molar mass of water and v is the stoichiometric parameter, i.e., the 104 number of ion moles produced by complete dissociation of one mole of salt (for a 1:1 105 salt such as LiBr, v equals 2); notice that molality m is expressed in kilomole of solute 106 per kilogram of solvent (or mole of solute per gram of solvent) so the units in eq. (2) are 107 consistent.

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To study the effect of solvent nature and the influence of the bromide salt on the passivity breakdown of Ti, additional experiments were carried out in different electrolytes: 0.5M LiBr solutions with different water-methanol proportions (70 mol%, 50 mol% and 0 mol% water), and a 3M KBr solution.

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The electrochemical tests were performed in a sheathed vertical electrochemical cell. A silver-silver chloride (Ag/AgCl 3M KCl) was used as reference electrode and a platinum wire as auxiliary electrode. Electrochemical measurements were performed using an Autolab PGSTAT302N potentiostat. In all cases the tests were repeated at least three times in order to verify reproducibility.

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### 120 Potentiodynamic polarization curves

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Polarization tests began at a cathodic potential value of -0.9  $V_{Ag/AgCl}$  and the potential was subsequently scanned anodically at a scan rate of 0.5 mV s<sup>-1</sup>. The passive current density,  $i_p$ , and the breakdown potential,  $E_b$  (defined as the potential at which the current density increases abruptly after the peak corresponding to the Br<sup>-</sup>/Br<sub>2</sub> oxidation), were

126	obtained from potentiodynamic polarization tests. After potentiodynamic polarization,
127	the morphological inspection of the Ti electrodes surface was done using a Leica DM
128	LA optical microscope.

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130 Potentiostatic passivation tests

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Before the potentiostatic passivation experiments, the surface of the Ti electrodes was pretreated cathodically at -0.9  $V_{Ag/AgCl}$  for 30 min to create reproducible initial conditions. Afterwards, the working electrode was polarized at a potential within the passive region of titanium (0.4  $V_{Ag/AgCl}$ ) for 1 hour in several LiBr solutions (0.1M, 0.5M, 3M and 11.42M) at 25° C, to form a stable passive film. During passivation experiments, transients in the current density were recorded against time.

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139 EIS and capacitance measurements

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EIS and capacitance measurements were performed under potentiostatic conditions after the passivation tests, once a stable passive film was formed on the surface of the Ti samples. EIS measurements were conducted using a sinusoidal signal of 10 mV in amplitude in the frequency range of 100 kHz-10 mHz. Capacitance measurements were performed at a frequency of 10 kHz applying a sinusoidal potential perturbation of 10 mV and scanning the potential from the formation value in the negative direction at a rate of 25 mV s<sup>-1</sup>.

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# 149 **3. RESULTS AND DISCUSSION**

#### 151 **Potentiodynamic polarization curves**

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153 The potentiodynamic polarization curves obtained for Ti in different LiBr solutions at 25° C are shown in Figure 1. A wide passive range can be observed in all LiBr 154 155 solutions where passive current density values  $(i_n)$  remain very low (of the order of 1-3  $\mu$ A cm<sup>-2</sup>), indicating the excellent passive behavior of Ti (**Table 2**). The values of  $i_p$ 156 157 have been estimated at a potential of 0.4 V, within the stable passive region. At first,  $i_p$ 158 slightly increases with increasing LiBr concentrations, but at concentrations higher than 0.7 M it starts decreasing until reaching its lowest value (1 µA cm<sup>-2</sup>) in the 11.42 LiBr 159 160 solution.

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162 Upon increasing the anodic polarization of Ti, the occurrence of a broad anodic wave 163 can be observed at potentials between 1-1.5 V, corresponding to the oxidation of bromides present in the solution  $(E^0(Br_2/Br) = 0.89 V_{Ag/AgCl})$ . This anodic wave is more 164 165 evident at high LiBr concentrations. In fact, in the potential range between 1-1.5 V a 166 change of the solution color around the titanium electrode surface (from colorless to 167 yellow-orange) was noticeable at high LiBr concentrations. This change in color is 168 related to the localized oxidation of bromide ions to bromine (Br<sub>2</sub>) at electroactive sites 169 at potentials around 1.5 V<sub>Ag/AgCl</sub> [15, 17-20, 22] and to the subsequent formation of 170 HBrO, according to the following equations [34]:

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$$2Br^{-} \Leftrightarrow Br_{2} + 2e^{-}$$
 (3a)

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174 
$$Br_2 + H_2O \rightarrow HBrO + Br^- + H^+$$
 (3b)

Apart from the oxidation of Br<sup>-</sup>, several processes may take place simultaneously during anodic polarization of titanium at high potentials, such as oxygen evolution reaction (OER) [2, 19, 21, 23, 25, 35], oxidation of TiO and Ti<sub>2</sub>O<sub>3</sub> to TiO<sub>2</sub> [2, 3, 5, 25, 35] or additional oxidation of TiO<sub>2</sub> to form soluble complex species, such as TiO<sup>2+</sup> [36], or/and peroxo-species, such as TiO<sub>3</sub> or TiOOH [7, 8, 10, 15].

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182 The abrupt increase of current density at potentials above the Br oxidation peak 183 indicates the onset of passivity breakdown of Ti. Once exposed to the Br solution, the Ti metal rapidly oxidizes to a soluble oxi-halide Ti<sup>4+</sup> species [18, 23]. In the low and 184 185 moderately concentrated LiBr solutions (Figure 1a) a decrease of  $E_b$  with increasing 186 bromide concentration can be observed, as reported in the literature [15, 18, 20] (Table 187 **2**). In the more concentrated LiBr solutions (Figure 1b), on the other hand,  $E_b$  does not 188 follow the previous tendency, since its value increases with increasing LiBr 189 concentration (Table 2). In fact, in the 11.42M LiBr solution Ti did not undergo 190 localized attacks under the experimental conditions (up to 9  $V_{Ag/AgCl}$ ), indicating that  $E_b$ 191 is apparently higher than 9  $V_{Ag/AgCl}$ , which is a behavior more typical of chloride- or 192 iodide- than bromide-containing solutions [18]. A similar behavior was observed by 193 Davydov for niobium in LiBr solutions [37]. By way of illustration, Figure 2 shows the 194 images of the Ti electrode surface before and after the polarization tests in 0.5M 195 (Figure 2a) and 11.42M (Figure 2b) LiBr solution. Visual inspection of the Ti 196 electrode polarized in the 0.5M LiBr solution (Figure 2a) reveals the presence of 197 several large corrosion pits. On the other hand, no damage is observed on the Ti 198 electrode polarized in the 11.42M LiBr solution up to 9 V (Figure 2b), in accordance 199 with the description of the polarization curves made above.

201 Figure 3 shows the plot of  $E_b$  vs. the logarithm of bromide activity (log  $a_{Br}$ ) for all the 202 LiBr solutions under study. It can be observed that  $E_b$  decreases linearly with log  $a_{Br}$ . 203 for the low and moderately concentrated LiBr solutions, while it increases linearly with 204 log  $a_{Br}$ . for the highly concentrated LiBr solutions (except in the 11.42M LiBr solution, 205 where passivity breakdown was not observed at all under the experimental conditions, as commented above). The linear relationship with negative slope between  $E_b$  and log 206 207  $a_{Br}$  (in low and moderately concentrated LiBr solutions) has been predicted analytically 208 by the Point Defect Model (PDM) [38, 39], according to the following equation:

209

210 
$$E_{b} = \left(\frac{4.606RT}{\chi\alpha F}\right)\log\left(\frac{b}{D}\right) - \frac{2.303RT}{\alpha F}\log(a_{Br^{-}})$$
(4)

211

where  $\gamma$  is the charge of the metal ions (Ti<sup>4+</sup>),  $\alpha$  is the polarizability of the passive 212 213 film/solution interface, F is Faraday's constant (96485.34 C mol<sup>-1</sup>), b is a constant, D is 214 the diffusivity of cation vacancies in the passive film, R is the gas constant and T is 215 absolute temperature. From the straight line with negative slope shown in **Figure 3**, the 216 value of  $\alpha$  can be determined, yielding a value of 0.045. This low value is comparable to 217 those reported in the literature for Ti [13, 15]. A low value of  $\alpha$  suggests that the 218 Helmholtz layer is very weakly polarized since the potential drop at the passive 219 film/solution interface is negligible. Thus, almost all the potential drop in the interfacial 220 regions takes place in the metal/passive film interface, where the reaction of passive film formation and growth occurs according to the PDM [40-42]. 221

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# 223 **Potentiostatic passivation tests**

225 The current density transients obtained after applying a potential step from a value of 226 -0.9 V (cathodic pre-treatment) to a value of 0.4 V within the passive region are shown 227 in Figure 4 for different LiBr solutions (0.1M, 0.5M, 3M and 11.42M). It can be 228 observed that current densities follow the same tendency as  $i_p$ , that is, the order from 229 highest to lowest values during the whole potentiostatic test is: 0.5M, 3M, 0.1M and 230 11.43 LiBr solutions. In all cases, current density values decrease exponentially with 231 passivation, and eventually reach a steady-state value  $(i_{SS})$ . These values are the following: 0.12±0.05 µA cm<sup>-2</sup> (0.1M), 0.22±0.07 µA cm<sup>-2</sup> (0.5M), 0.17±0.05 µA cm<sup>-2</sup> 232 (3M) and 0.059±0.03 µA cm<sup>-2</sup> (11.42M). Virtanen and Curty [20] observed a clear 233 234 increase in passive current density with increasing bromide concentrations for Ti in 235 NaBr. However, in the present study,  $i_{SS}$  increases with Br<sup>-</sup> concentration at relatively 236 low concentrations, but it decreases with concentration at high LiBr concentrations.

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# According to the PDM [42-44], the steady-state current density can be written as:

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240 
$$i_{SS} = \delta F \left[ k_{int} + k_{cat} + k_{dis} \left( \frac{C_{H^+}}{C_{H^+}^0} \right)^n \right]$$
(5)

241

where  $\delta$  is the charge of the cation ejected from the passive film (3 in the case of Cr<sup>3+</sup>),  $C_{H^+}$  is the proton concentration in the solution,  $C_{H^+}^0$  is a standard state  $H^+$  concentration and *n* is the kinetic order of the reaction of passive film dissolution with respect to  $(C_{H^+}/C_{H^+}^0)$ . The parameters  $k_{int}$ ,  $k_{cat}$  and  $k_{dis}$  are the rate coefficients for the reactions of formation of cation interstitials, cation vacancies and of passive film dissolution, respectively. Eq. (5) suggests that the steady-state current density  $i_{SS}$  is associated with the fluxes of point defects and also with the kinetic parameters for the interfacial reactions. Hence,  $i_{SS}$  consists of three components: the first ( $\delta F k_{int}$ ) resulting from the generation and transport of cation interstitials, the second ( $\delta F k_{cat}$ ) resulting from the generation and transport of cation vacancies and the third  $\delta F k_{dis} (C_{H^+}/C_{H^+}^0)^n$  resulting from the movement of oxygen vacancies. The latter term is expressed in terms of the dissolution rate of the passive film [42-45].

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The passive film on titanium is well known to behave as an *n*-type semiconductor [1, 6, 10-13] (see Mott-Schottky analysis below), whose donor species are oxygen vacancies and/or Ti<sup>3+</sup> interstitials [10, 12-15], so the term related to the generation and transport of cation vacancies,  $k_{cat}$ , can be neglected [43]. Therefore, the differences between current densities with LiBr concentration are either due to the process of generation and transport of cation interstitials or the dissolution of the passive film closely related to the generation and movement of oxygen vacancies in the steady-state [41].

- 262
- 263264 Mott-Schottky analysis

265 266 267

268 In order to check if the differences observed in  $i_{SS}$  at different LiBr concentrations is 269 due to the generation of donor species at the metal/passive film interface (oxygen 270 vacancies and/or cation interstitials), the electrochemical capacitance of the passive 271 system was measured as a function of the applied potential. Mott-Schottky (M-S) 272 analysis was performed to obtain the donor density,  $N_D$ , of the passive films formed on 273 Ti in different LiBr solutions. The total capacitance has been calculated from the 274 imaginary component of the impedance (Z') using the relationship  $C = -1/\omega Z''$ , where  $\omega$ 275 is the angular frequency.

277 As an illustration, Figure 5 shows the capacitance-potential curves obtained for Ti in 278 the 0.5M LiBr solution at 0.4 V and at different frequencies. It is clearly observed that 279 capacitance values depend strongly on frequency. Capacitance-potential curves shift to 280 lower values with increasing frequencies. Such behavior has been frequently observed 281 for semiconducting passive films formed on metals and alloys [12, 46-54]. Several 282 physical explanations have been suggested in the literature as possible sources of this 283 dependence [46-54], such as the presence of deep donor states, contribution of surface 284 states, surface roughness, ionic charges movement, frequency dependant dielectric 285 constant, etc. It can be seen from Figure 5 that capacitance becomes almost 286 independent of frequency at approximately 10 kHz. Therefore, a value of 10 kHz has 287 been used in this work to eliminate capacitance dependence on frequency.

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Figure 6 shows the Mott-Schottky plots at a frequency of 10 kHz of the passive film formed on Ti in 0.1M, 0.5M, 3M and 11.42M LiBr solutions at 0.4 V. The positive slopes of the M-S plots are characteristic of *n*-type semiconductors, with the dominant defects in the passive film being oxygen vacancies and/or  $Ti^{3+}$  interstitials, as it has been mentioned above.

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The donor density,  $N_D$ , can be determined from the slopes in the M-S plots, according to the following equation [10, 12]:

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$$N_D = \frac{2}{\varepsilon \varepsilon_0 e \sigma}$$
(6)

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300 where  $\varepsilon$  is the dielectric constant of the passive film (a value of 60 has been assumed for 301 the passive film formed on Ti [15]),  $\varepsilon_0$  is the vacuum permittivity (8.85·10<sup>-14</sup> F/cm), *e* is 302 the electron charge  $(1.60 \cdot 10^{-19} \text{ C})$  and  $\sigma$  is the positive slope of each straight line in the 303 M-S plots.

304

305 The values of  $N_D$  are presented in Figure 7 for different LiBr solutions (0.1M, 0.5M,

306 3M and 11.42M) at 0.4 V. The results show that  $N_D$  values are of the order of  $10^{19}$ - $10^{20}$ 

307 cm<sup>-3</sup>, regardless the LiBr concentration, which is in agreement with those reported in the

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literature for TiO<sub>2</sub> [10, 12].

An increase in LiBr concentration from 0.1M to 0.5M leads to a clear increase of  $N_D$ (Figure 7). The PDM [38, 40, 55] assumes that the initial event that takes place at the passive film/electrolyte interface in passivity breakdown is the occupation of surface oxygen vacancies by aggressive anions, such as bromides:

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315 
$$V_o^{\bullet\bullet} + \mathrm{Br}^- \to \mathrm{Br}_o^{\bullet} \tag{7}$$

316

317 where  $Br_0^{\bullet}$  represents a bromide anion adsorbed into an oxygen vacancy,  $V_0^{\bullet\bullet}$ , in the 318 film/electrolyte interface. This process is followed by a Schottky-pair type reaction:

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320 Null 
$$\rightarrow V_{T_i}^{4'} + \left(\frac{\chi}{2}\right) V_o^{\bullet \bullet}$$
 (8)

321

where  $V_{Ti}^4$  is a cation vacancy in the Ti passive film. Kroger-Vink notation is used for the effective positive (\*) and negative (<sup>4</sup>) charges in oxygen and cation vacancies, respectively.

The adsorbed Br<sup>-</sup> anions could desorb along with a surface cation to form a cation vacancy/oxygen vacancy pair  $(V_{T_i}^{4'} V_{O}^{\bullet\bullet})^{2'}$ . This pair is subsequently annihilated, leading to the formation of a subsurface cation vacancy and the regeneration of the surface  $V_{O}^{\bullet\bullet}$ : 329

330 
$$\left(V_{T_{i}}^{4'}V_{O}^{\bullet}\right)^{"} \rightarrow \left(V_{T_{i}}^{4'}\right)_{S} + V_{O}^{\bullet}$$
 (9)

331

where the subscript *S* indicates a subsurface position in the passive film. The oxygen vacancies continue reacting with additional  $Br^-$  at the film/electrolyte interface to generate additional cation and oxygen vacancies. Hence, the generation of cation and oxygen vacancies through eqs. (7)-(9) is autocatalytic.

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337 Anion adsorption leads to the generation of cation vacancies at the film/solution 338 interface and hence to an enhancement in the flux of cation vacancies across the passive 339 film to the film/metal interface due to the outward movement (towards the passive 340 film/electrolyte interface) of cations in the film. If the flux of cation vacancies exceeds 341 the rate at which they are annihilated at the metal/film interface, the excess vacancies 342 will condense at the interface and the passive film will be locally detached from the 343 metal. As passive film dissolution continues, the film thins locally, eventually leading to 344 film breakdown [40].

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According to the previous explanation, since oxygen vacancies are donor species, the higher the donor density in the passive film,  $N_D$ , the lower the resistance to passivity breakdown. This affirmation is in agreement with the values of the breakdown potential,  $E_b$ , shown in **Table 2**, where  $E_{b0.5M} < E_{b0.1M}$ . However, when increasing the LiBr concentration from 0.5M to 3M and 11.42M,  $N_D$  decreases, indicating that the higher bromide concentration in the 3M and 11.42M LiBr solutions does not lead to a higher degree of oxygen vacancies occupation by bromide ions. This result could explain the lower susceptibility of Ti to undergo passivity breakdown in the 3M and 11.42M LiBr solutions if compared with the 0.5M LiBr solution. Nevertheless, since the  $N_D$  values obtained in concentrated LiBr solutions are still higher than that in the 0.1M LiBr solution, this fact cannot satisfactorily explain the increase of  $E_b$  at high LiBr concentrations or even the absence of passivity breakdown in the 11.42M LiBr solution.

The tendency of  $N_D$  with LiBr concentration also confirms that the differences mentioned above in the steady-state current density ( $i_{SS}$ ) with LiBr concentration are not related to the density of donor species (oxygen vacancies and/or cation interstitials), since  $i_{SS0.IM} > i_{SS11.42M}$  but  $N_{D0.IM} < N_{D11.42M}$ .

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#### 364 EIS measurements

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Mott-Schottky analysis has proved that the behavior of the passive film on Ti against localized attacks in high LiBr concentrations is not directly related with its electronic properties. **Figure 8** shows the experimental and simulated Nyquist diagrams for Ti in 0.1M, 0.5M, 3M and 11.42M LiBr solutions at 0.4 V. The experimental data have been fitted to the electrical equivalent circuit shown in **Figure 9**, commonly used to account for the bilayer structure of the passive film on Ti (porous outer layer and compact inner layer) [4, 5, 56]

In all cases, Nyquist plots exhibit a typical passive state behavior characterized by
semicircular shape and high impedance values, suggesting that a highly stable film is
formed on all the electrodes [31-34].

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An increase in LiBr concentration from 0.1M to 0.5M reduces the amplitude of the 378 379 semicircle in Nyquist plots, indicating a decrease of the total impedance of the system 380 and a loss in the protective properties of the passive film. This result is consistent with 381 the higher  $i_{SS}$  and  $N_D$  values obtained in the 0.5M LiBr solution. However, upon 382 increasing the bromide concentration to 3M and 11.42M, the total impedance increases 383 again, which agrees with the tendency of  $i_{SS}$  observed before (the lower the  $i_{SS}$ , the 384 higher the total impedance of the system), but differs in the case of  $N_D$  values (Figure 385 7), since the lowest  $N_D$  value (0.1M LiBr solution) does not imply the highest 386 impedance values.

387

The values of the equivalent circuit parameters are shown in **Table 3**. In this equivalent circuit,  $R_S$  represents the solution resistance, the time constant at high frequencies  $(R_1CPE_1)$  is related to the outer layer of the passive film (hydroxides), whereas the lowfrequency behavior  $(R_2CPE_2)$  is attributed to the inner layer (mainly TiO<sub>2</sub>), more compact than the outer one. Constant Phase Elements (CPEs) have been used instead of pure capacitors, to describe the non-ideality of the system studied. CPEs have been converted into pure capacitance, *C*, by using the following equation [57, 58]:

395

$$C = \frac{(Q \cdot R)^{1/\alpha}}{R}$$
(10)

398 where Q is the impedance of the CPE and R corresponds to  $R_2$  when determining 399 capacitance values of the second time constant in **Figure 9**. To determine  $C_1$  from 400 *CPE*<sub>1</sub>, R has been calculated as follows:

401

402 
$$\frac{1}{R} = \frac{1}{R_s} + \frac{1}{R_1}$$
 (11)

403

404 It can be observed from **Table 3** that the resistance values of the compact inner layer 405  $(R_2)$  are significantly larger than the values associated with the outer porous layer  $(R_1)$ , 406 which is consistent with the chosen physical model. These results indicate that the 407 protection provided by the passive film is predominantly due to the barrier layer. In 408 general, both the outer and inner layers exhibit similar capacitance values, of the order of some µF cm<sup>-2</sup>. Regarding the influence of LiBr concentration on the parameters, it 409 410 can be observed that  $R_1$  decreases at high LiBr concentrations, while  $R_2$  decreases from 411 0.1M to 0.5M but then it increases again, reaching the highest value in the 11.42M LiBr 412 solution. These results are in agreement with those obtained for  $i_{SS}$  (the lower the value 413 of  $i_{SS}$ , the higher the resistance of the barrier inner layer of the passive film,  $R_2$ ).

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415 According to the M-S results, the increase of  $R_2$  in the 3M and 11.42M LiBr solutions 416 with respect to the value in the 0.1M LiBr solution cannot be related to the degree of 417 defectiveness of the passive film on Ti (Figure 7 and Table 3). This increase of  $R_2$ 418 could be explained in terms of the passive film dissolution and the influence of the 419 activity of water,  $a_W$ , on this process. Water is one of the reacting species in 420 electrochemical and corrosion mechanisms, but most corrosion works focus on the 421 properties of solutes whereas changes in the activity of water are rarely considered [33, 422 59]. Smart and Bockris [59] observed a decrease in the corrosion current density of iron 423 in concentrated chloride media with decreasing free water in the solution, that is, the 424 rate of iron dissolution (Fe/Fe<sup>2+</sup>) was shown to strongly decrease with decreasing 425 activity of water. Hence, in the case of Ti in highly concentrated LiBr solution, the 426 decrease of  $a_W$  (**Table 1**) might lead to similar results. The reaction of dissolution of the 427 passive film formed on Ti by proton attack can be expressed as follows:

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429 
$$\operatorname{TiO}_{2} + 4\mathrm{H}^{+} \rightarrow \operatorname{Ti}_{aq}^{4+} + 2\mathrm{H}_{2}\mathrm{O}$$
 (12)

430

Although the activity of protons,  $a_{H^+}$ , is constant, the concentration of H<sup>+</sup> will decrease 431 432 with decreasing free water (decreasing  $a_W$ ). Therefore, the dissolution reaction given by eq. (12) not only depends on pH  $(a_{H^+})$  but also on the number on H<sup>+</sup> at the passive 433 434 film/electrolyte interface, which in turn depends on  $a_W$ . Passive film dissolution is a 435 necessary step for passivity breakdown to occur, causing the thinning and subsequent 436 rupture of the passive film [40, 60]. Therefore, if passive film dissolution is somehow 437 hampered, for example by reducing the activity of water, passivity breakdown will 438 appear at higher potentials. This fact would explain the increase of  $E_b$  in highly 439 concentrated LiBr solutions (Table 2).

440

441 On the other hand, the formation of the precipitated outer layer (TiO<sub>2</sub>·*n*H<sub>2</sub>O) via the 442 hydrolysis of Ti<sup>4+</sup> cations ejected from the inner layer (eq. (12)) also depends on  $a_W$ :

443

444 
$$\operatorname{Ti}_{\mathrm{aq}}^{4+} + (n+2)\operatorname{H}_2\operatorname{O} \to \operatorname{TiO}_2 \cdot n\operatorname{H}_2\operatorname{O} + 4\operatorname{H}^+$$
 (13)

445

In fact, the values of the outer layer resistance,  $R_I$ , are higher in the 0.1M and 0.5M LiBr solutions (**Figure 7**), indicating that the formation of precipitated hydroxides is favored when  $a_W$  values are close to unity (**Table 1**). 449 However, the explanations given above concerning the dissolution of the passive film 450 do not completely explain why this behavior has not been observed, for example, in 451 stainless steels in highly concentrated LiBr solutions [35, 61, 62]. Apart from the 452 process of passive film dissolution (eq. (12)), there must be other Ti-specific surface 453 processes which are influenced by the decrease of  $a_W$  in highly concentrated LiBr 454 solutions. Davydov [37] suggested that the formation of an intermediate complex 455 between a valve metal (such as Ti) and the aggressive anions (Br) and its interaction 456 with the metal oxide could restrict the process of metal dissolution. Following the 457 descriptions of Davydov [37], this complex is formed when a bromide anion adsorbs on 458 a surface oxygen vacancy (eq. (7)) and then interacts with a Ti cation in the passive film 459  $(Ti_{Ti})$ , according to the following reaction:

460

461 
$$\operatorname{Ti}_{\mathrm{Ti}} + a \mathrm{Br}_{\mathrm{O}}^{\bullet} \rightarrow [\mathrm{Ti}\mathrm{Br}_{\mathrm{a}}^{z}]_{\mathrm{s}} + n\mathrm{e}^{-}$$
 (14)

462

463 where z = n + a and  $[TiBr_a^z]_s$  is the surface complex. Along with the complex formation, 464 a cation vacancy/anion vacancy pair is also generated via a Schottky-pair type reaction, 465 eq. (8). The hydrolysis of the complex will break the bond with the passive film, thus 466 favoring the process of dissolution with the formation of the product of activated 467 dissolution, TiO<sub>2</sub>·*n*H<sub>2</sub>O:

468

469 
$$2\left[\mathrm{TiBr}_{a}^{z}\right]_{s} + (5+n)\mathrm{H}_{2}\mathrm{O} \rightarrow \mathrm{TiO}_{2}\cdot n\mathrm{H}_{2}\mathrm{O} + 2a\mathrm{Br}^{-} + 10\mathrm{H}^{+}$$
(15)

470

471 As in the case of eq. (12), a decrease in the content of free water (i.e. decrease in  $a_W$ ) 472 will slow down the process of hydrolysis and, as a consequence, the full process of 473 passivity breakdown.

# 475 Influence of solvent and solute on the activity of water and passivity breakdown

476

477 In order to verify that the decrease of the activity of water is responsible for the decrease 478 of the susceptibility of Ti to undergo passivity breakdown in highly concentrated LiBr 479 solutions, additional electrolytes have been used to study the electrochemical behavior 480 of Ti. In the first case, the activity of water has been reduced by adding different 481 amounts of methanol to a 0.5M LiBr aqueous solution.  $a_W$  values in these solutions have 482 been estimated from literature [63] taking the values of  $a_W$  shown in **Table 1** as starting 483 point: water-free methanol solution ( $a_W \approx 0$ ), 50 mol% water – 50 mol% methanol ( $a_W \approx$ 0.60) and 70 mol% water – 30 mol% methanol ( $a_W \approx 0.76$ ). In the second case, the 484 485 activity of water has been increased by changing the bromide salt (LiBr was replaced by 486 KBr), since Li is the smallest alkali cation and it is most strongly hydrated, thus reducing the amount of free water in the solution. A 3M KBr solution has been used, 487 whose molality  $m_{KBr} = 3.40 \pmod{\text{KBr}} (\text{kg H}_2\text{O})^{-1}$  is similar to that of a 3M LiBr 488 solution ( $m_{LiBr} = 3.77$  (mol LiBr) (kg H<sub>2</sub>O)<sup>-1</sup>). The comparison of mean ionic activity 489 490 coefficients  $(\gamma_{\pm})$  and activities of water  $(a_W)$  between LiBr and KBr solutions is shown 491 in Figures 10a and 10b, respectively. It can be observed that  $\gamma_{\pm}$  values are higher in 492 LiBr solutions than in KBr solutions in almost the whole range of molalities, while  $a_W$ 493 values are visibly lower for LiBr solutions than for KBr solutions at molalities higher than 2 mol kg<sup>-1</sup> (for 3M KBr and LiBr solutions,  $a_{W \ KBr} = 0.89$  and  $a_{W \ LiBr} = 0.81$ ). 494

495

496 Figure 11 shows the potentiodynamic polarization curves of Ti in a 0.5M LiBr solution 497 with different water-methanol mol%. It can be observed that the four curves present a 498 clear passive range, regardless the water and methanol content. However, in the waterfree methanol solution, this passive range is significantly narrower than in the other solutions. Burstein and Whillock [64], as well as Banas et al [65], studied the passive behavior of Ti and other metals in methanolic solutions, and they concluded that methanol was responsible for the process of Ti dissolution, according to:

503

504 
$$Ti + nMeOH \rightarrow Ti(OMe)_n^{(4-n)+} + nH^+ + 4e^-$$
 (16)

505

506 In the presence of water, (50 mol% water, 70 mol% water and aqueous solution), Ti 507 presents considerably wider passive regions than in the free-water methanol solution 508 (**Figure 11**). Apart from the formation of  $TiO_2$  due to direct reaction of Ti with water 509 molecules, the dissolved complex of Ti(IV) formed according to eq. (16) could 510 hydrolyze to form additional  $TiO_2$  [64]:

511

512 
$$\operatorname{Ti}(\operatorname{OMe})_{n}^{(4-n)+} + 2H_{2}O \rightarrow \operatorname{TiO}_{2} + n\operatorname{MeOH} + (4-n)\mathrm{H}^{+}$$
 (17)

513

514 It is clearly observed in **Figure 11** that the breakdown potential in the water-methanol 515 mixed solutions is significantly higher than in the aqueous solution ( $E_{b70,30} = 4.9 \pm 0.2$  V;  $E_{b50 50} = 3.7 \pm 0.3$  V;  $E_{b_aqueous} = 2.6 \pm 0.2$  V). This increase in  $E_b$  can be explained taking 516 517 into account the decrease of free water in the electrolyte (i.e. decrease in  $a_W$ ), by 518 partially replacing the aqueous solvent by the methanolic one. Thus, although methanol 519 causes Ti dissolution (eq. (16)), the decrease of  $a_W$  slows down the dissolution of the 520 electrode surface via eqs. (12) and/or (15), in a similar way as in highly concentrated 521 LiBr solutions.

523 The potentiodynamic polarization curves of Ti in 3M KBr and 3M LiBr solutions are 524 presented in Figure 12. It can be observed that Ti is more susceptible to undergo 525 passivity breakdown in the 3M KBr solution ( $E_{b_{-KBr}} = 2.6 \pm 0.2$  V) than in the 3M LiBr 526 solution ( $E_{b \ KBr} = 3.3 \pm 0.3$  V). As it has been explained before, the mean ionic activity 527 coefficient reaches higher values in LiBr than in KBr solutions due to the high charge 528 density of the small lithium ion (Figure 10a). The higher activity coefficients in LiBr 529 solutions result in higher activity of bromide anions ( $a_{Br}$  in the 3M KBr solution is 2.0 530 and  $a_{Br}$  in the 3M LiBr solution is 6.6). As stated in eq. (4), the value of  $E_b$  should be 531 higher in the 3M KBr solution than in the 3M LiBr solution, since  $a_{Br}$  is significantly 532 lower in the former electrolyte. However, the higher values of  $\gamma_{\pm}$  in the LiBr solutions 533 indicate a very strong interaction between water and ions, which lead to a decrease of 534  $a_W$  (Figure 10b). Again, the reduction of  $a_W$  in the 3M LiBr solution restricts the 535 process of dissolution of the Ti surface (eqs. (12) and/or (15)), thus delaying the onset 536 of stable passivity breakdown.

537

## 538 Relationship between $E_b$ and $a_W$ in concentrated LiBr solutions

539

The linear relationship with positive slope found between  $E_b$  and log  $a_{Br}$  in highly concentrated LiBr solutions (**Figure 3**) provides an empirical equation to evaluate  $E_b$  in function of the activity of bromide anions. However, according to the previous results, it is more coherent to relate  $E_b$  to the activity of water,  $a_W$ , since it is this parameter that directly affects the value of  $E_b$  in these concentrated media. Starting from eqs. (1) and (2), the activity of water can be expressed in terms of  $a_{Br}$ .:

547 
$$\log a_{W} = \frac{-\phi M_{W} v \frac{a_{Br^{-}}}{\gamma_{\pm}}}{2.303 \cdot 1000} = \frac{-\phi \cdot 18 \cdot 2 \frac{a_{Br^{-}}}{\gamma_{\pm}}}{2.303 \cdot 1000} = \frac{-A\phi a_{Br^{-}}}{\gamma_{\pm}}$$
(18)

548

549 On the other hand,  $a_{Br}$  can be expressed in terms of  $E_b$  by using the linear relationship 550  $E_b$  vs. log  $a_{Br}$  mentioned above (**Figure 3**):

551

552 
$$a_{Br^-} = 10^{\frac{E_b - 2.50}{0.93}}$$
 (19)

553

554 Substituting eq. (19) into eq. (18):

555

556 
$$-\log a_W = \frac{A\phi}{\gamma_{\pm}} 10^{\frac{E_b - 2.50}{0.93}} \rightarrow \log(-\log(a_W)) = \log\left(\frac{A\phi}{\gamma_{\pm}}\right) + \left(\frac{E_b - 2.50}{0.93}\right)$$

557

558 Finally, solving for 
$$E_b$$
 yields:

559

560 
$$E_b = 0.93 \cdot \left[ \left( \log(-\log(a_W)) \right) - \log\left(\frac{A\phi}{\gamma_{\pm}}\right) \right] + 2.50$$
 (20)

561

Eq. (20) relates the value of  $E_b$  to a logarithmic function of  $a_W$ . Figure 13 shows the experimental and calculated values of  $E_b$  as a function of  $a_W$  in concentrated LiBr solutions. It can be observed that both experimental and calculated values of  $E_b$  are similar, indicating that eq. (20) can be used to satisfactorily calculate the breakdown potential of Ti in concentrated LiBr solutions ( $m_{LiBr} > \sim 0.7$  mol kg<sup>-1</sup>) as a function of the activity of water.

569 Finally, it is necessary to explain why Ti did not undergo passivity breakdwon in the 570 11.42M LiBr solution. According to eq. (20), the breakdown potential of Ti in the 571 11.42M LiBr solution ( $a_W = 0.12$ ) should be 5.6 V. However, passivity breakdown was 572 not observed even at potentials as high as 9 V. The absence of localized attacks in this 573 solution can be explained by changes in the  $TiO_2$  surface. The irreversible formation of 574 titanium peroxo species at sufficiently high potentials (normally higher than  $3V_{Ag/AgCl}$ ) 575 has been reported in the literature [7, 8, 15]. The formation of Ti peroxo species may 576 result in a very thin insulating or dielectric layer [8, 15] which would form in the 577 hydrated (outer) part of the passive film above 3V [8]:

578

579  $\operatorname{TiO}_2 + \operatorname{OH}^- \rightleftharpoons \operatorname{TiO}_3 + \operatorname{H}^+ + 2e^-$  (21)

580

581 The presence of these species, such as  $TiO_3$ , can be noted by a yellowish color on the 582 electrode surface [8]. This yellow color was observed on the undamaged areas of all the 583 electrodes after the potentiodynamic polarization tests. However, the surface of the 584 electrode polarized in the 11.42M LiBr solution had a dark golden color. The color-585 dependence of anodic Ti oxide layers on the polarization potential, and hence oxide 586 thickness, is a very well known phenomenon [66-73]. For example, it has been observed 587 [70, 71, 73] that at 5V the Ti surface has a light vellow color, while at 10-15V the 588 surface color changes to dark golden/brown. Hence, the absence of stable localized 589 attacks of Ti in the 11.42M LiBr solution can be attributed either to the formation of a 590 highly insulating or dielectric layer of Ti peroxo species at potentials higher than ~5V (the value of  $E_p$  in the 8.06M LiBr solution is 4.7V) or/and to the increase of the Ti 591 592 oxide film thickness. Both phenomena could block the interaction between aggressive anions and the surface of Ti, preventing passivity breakdown. This mechanism may be 593

similar to that reported by Sazou et al. [15] to explain passivity breakdown of Ti at high potentials in the presence of high concentrations of chlorides. In any case, it is evident that at sufficiently high potentials (higher than  $\sim$ 5V) the surface properties of the titanium oxide change to such an extent that the risk of passivity breakdown is completely eliminated, at least up to potentials as high as 9V.

599

# 600 **5. CONCLUSIONS**

601

A wide passive range where passive current density values  $(i_p)$  remained very low (of the order of 1-3  $\mu$ A cm<sup>-2</sup>) was observed in all LiBr solutions, indicating the excellent passive behavior of Ti.  $i_p$  values slightly increased with increasing LiBr concentrations up to 0.7M LiBr, where  $i_p$  started to decrease until reaching its lowest value (1  $\mu$ A cm<sup>-2</sup>) in the 11.42 LiBr solution.

607

In the low and moderately concentrated LiBr solutions, a decrease of  $E_b$  with increasing bromide concentrations was observed, while in the more concentrated LiBr solutions,  $E_b$ increased with increasing LiBr concentration. In the most concentrated LiBr solution (11.42M) Ti did not undergo passivity breakdown under the experimental conditions (up to 9 V<sub>Ag/AgCl</sub>).

613

614 The steady-state passive current density,  $i_{SS}$ , was found to increase with Br<sup>-</sup> 615 concentration at relatively low concentrations, but it decreased at high LiBr 616 concentrations, following the same tendency as  $i_p$ .

617

618 The passive film formed on Ti behaved as an *n*-type semiconductor, with the dominant defects in the passive film being oxygen vacancies and/or Ti<sup>3+</sup> interstitials. An increase 619 in LiBr concentration from 0.1M to 0.5M led to a clear increase of the donor density 620 621  $(N_D)$ , but when increasing the LiBr concentration from 0.5M to 3M and 11.42M,  $N_D$ 622 decreased. However, N<sub>D</sub> values obtained in concentrated LiBr solutions were still higher 623 than in the 0.1M LiBr solution, so the electronic properties of TiO<sub>2</sub> passive film could 624 not be used to satisfactorily explain the increase of  $E_h$  at high LiBr concentrations or the 625 absence of localized attacks in the 11.42M LiBr solution.

626

The resistance of the hydroxide outer layer of the passive film on Ti was higher in low concentrated LiBr solutions, whereas the resistance of the barrier layer of the film was found to be higher in concentrated LiBr solutions. These results could not be related to the degree of defectiveness of the passive film on Ti, but seemed to be associated with the passive film dissolution and the influence of the activity of water,  $a_W$ , on this process.

633

The activity of water,  $a_W$ , was found to have a strong influence on the susceptibility of Ti to undergo passivity breakdown. In view of the results, it may be concluded that a decrease in this parameter was responsible for the increase of the breakdown potential in highly concentrated LiBr solutions and water-methanol-LiBr mixtures. An empirical correlation between  $E_b$  and  $a_W$  was obtained for Ti in concentrated LiBr solutions.

639

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644

645

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# 763 <u>Tables</u>

$C_{LiBr}$ /mol l <sup>-1</sup>	$C_{LiBr}/\mathrm{g}\mathrm{\Gamma}^1$	$m_{LiBr}/(\text{mol LiBr}) (\text{kg H}_2\text{O})^{-1}$	$\gamma_{\pm}$	a <sub>LiBr</sub>	ø	a <sub>W</sub>
0.05	4.34	0.05	0.93	0.05	0.95	1.00
0.10	8.68	0.10	0.80	0.08	0.94	1.00
0.30	26.05	0.31	0.76	0.23	0.95	0.99
0.50	43.42	0.51	0.75	0.39	0.97	0.98
0.70	60.79	0.72	0.77	0.56	1.00	0.97
1.15	100	1.20	0.84	1.01	1.07	0.95
3.00	260.52	3.77	1.75	6.60	1.53	0.81
4.61	400	5.69	3.56	20.26	1.93	0.67
8.06	700	11.18	21.89	222.82	2.95	0.34
11.42	992	15.15	154.00	2333.67	3.93	0.12

# 765 Table 1. Physical properties of the different LiBr solutions used in this work.

768 Table 2. Values of passive current density  $(i_p)$  at 0.4 V and breakdown potential

769 ( $E_b$ ) of Ti in different LiBr solutions, obtained from potentiodynamic polarization

**curves.** 

$C_{LiBr}$ /mol l <sup>-1</sup>	$i_p/\mu A \text{ cm}^{-2}$	E <sub>b</sub> vs (Ag/AgCl)/V
0.05	$1.4 \pm 0.1$	$3.8 \pm 0.2$
0.10	$2.0 \pm 0.2$	$3.2 \pm 0.3$
0.30	$2.3 \pm 0.2$	$2.7 \pm 0.2$
0.50	$2.5 \pm 0.3$	$2.6 \pm 0.3$
0.70	$2.6 \pm 0.1$	$2.2 \pm 0.2$
1.15	$1.7 \pm 0.1$	$2.7 \pm 0.1$
3.00	$1.4 \pm 0.1$	$3.3 \pm 0.3$
4.61	$1.4 \pm 0.1$	$3.8 \pm 0.2$
8.06	$1.2 \pm 0.1$	$4.7 \pm 0.3$
11 42	$10 \pm 02$	

785 Table 3. Values of the equivalent circuit parameters for Ti in the 0.1M, 0.5M, 3M

#### 786 and 11.42M LiBr solutions.

787

$C_{LiBr}$ /mol l <sup>-1</sup>	$\frac{R_S/\Omega}{\mathrm{cm}^2}$	$R_1/k\Omega$ cm <sup>2</sup>	$C_1/\mu F$ cm <sup>-2</sup>	$\alpha_{I}$	$R_2/k\Omega$ cm <sup>2</sup>	$C_2/\mu F$ cm <sup>-2</sup>	$\alpha_2$	$\chi^2 (x10^{-3})$
0.1	$7 \pm 1$	$23 \pm 3$	$3.8 \pm 0.8$	0.91 ± 0.01	2391 ± 157	$4.6 \pm 0.5$	0.78 ± 0.24	0.9
0.5	$4\pm0$	$28 \pm 2$	$4.4 \pm 0.7$	0.90 ± 0.01	1804 ± 116	10.5 ± 0.4	0.89 ± 0.14	1.8
3	$1 \pm 0$	19 ± 2	$3.4 \pm 0.5$	0.91 ± 0.01	3144 ± 108	$1.3 \pm 0.3$	0.94 ± 0.09	3.5
11.42	$2 \pm 0$	$16 \pm 2$	3.2 ± 0.6	0.94 ± 0.01	$3280 \pm 101$	3.9 ± 0.4	$0.87 \pm 0.07$	4.1

 $\begin{array}{c} 788\\ 789\\ 790\\ 791\\ 792\\ 793\\ 794\\ 795\\ 796\\ 797\\ 798\\ 799\\ 800\\ 801\\ 802\\ 803\\ 804\\ 805\\ 806\\ 807\\ 808\\ 809\\ 810\\ 811\\ 812\\ 813\\ 814\\ 815\\ 816\\ 817\\ 818\\ 820\\ 821\\ 822\\ 823\\ 824\\ 825\\ 826\\ 827\\ 828\\ 830\end{array}$ 

831 832	Tables captions
833	Table 1. Physical properties of the different LiBr solutions used in this work.
834	
835	Table 2. Values of passive current density $(i_p)$ at 0.4 V and breakdown potential $(E_b)$ of Ti in different
836	LiBr solutions, obtained from potentiodynamic polarization curves.
837	
838	Table 3. Values of the equivalent circuit parameters for Ti in the 0.1M, 0.5M, 3M and 11.42M LiBr
839	solutions.
840	
841	Figures captions
842	
843	Fig. 1. Potentiodynamic polarization curves of Ti in different LiBr solutions at 25° C; (a) low and
844	moderately concentrated LiBr solutions; (b) concentrated LiBr solutions.
845	
846	Fig. 2. Images of the Ti electrode surface before and after the polarization tests in (a) 0.5M and (b)
847	11.42M LiBr solutions.
848	
849	Fig. 3. Plots of $E_b$ vs. log $a_{Br}$ for Ti in the different LiBr solutions.
850	
851	Fig. 4. Current density transients obtained after applying a potential step from a cathode value of $-0.9$ V
852	to a passive value of 0.4 V for Ti in different LiBr solutions (0.1M, 0.5M, 3M and 11.42M).
853	
854	Fig. 5. Capacitance-potential curves obtained for Ti in 0.5M LiBr solution at 0.4 V and at different
855	frequencies.
856	
857	Fig. 6. Mott-Schottky plots of the passive film formed on Ti in 0.1M, 0.5M, 3M and 11.42M LiBr
858	solutions at 0.4 V, obtained at a frequency of 10 kHz.
859	
860	Fig. 7. Values of donor density $(N_D)$ of the passive film on Ti in different LiBr solutions (0.1M, 0.5M, 3M
861	and 11.42M).

862	Fig. 8. Experimental and simulated Nyquist diagrams for Ti in 0.1M, 0.5M, 3M and 11.42M LiBr
863	solutions at 0.4 V.
864	
865	Fig. 9. Electrical equivalent circuit used to fit the experimental EIS data.
866	
867	Fig. 10. Comparison of (a) mean ionic activity coefficients ( $\gamma_{\pm}$ ) and (b) activities of water ( $a_W$ ) between
868	LiBr and KBr solutions.
869	
870	Fig. 11. Potentiodynamic polarization curves of Ti in a 0.5M LiBr solution with different water-methanol
871	mol%.
872	
873	Fig. 12. Potentiodynamic polarization curves of Ti in 3M KBr and 3M LiBr solutions.
874	
875	Fig. 13. Experimental and calculated (using eq. (20)) values of $E_b$ as a function of $a_W$ in the highly
876	concentrated LiBr solutions.



**(a)** 



**(b)** 





Initial mosaic (sample diameter = 8 mm)

Final mosaic (sample diameter = 8 mm)





Initial mosaic (sample diameter = 8 mm)

Final mosaic (sample diameter = 8 mm)

**(b)** 























