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# ORTHODIAGONAL ANTI-INVOLUTIVE KOKOTSAKIS POLYHEDRA 

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#### Abstract

We study the properties of Kokotsakis polyhedra of orthodiagonal anti-involutive type. Stachel conjectured that a certain resultant connected to a polynomial system describing flexion of a Kokotsakis polyhedron must be reducible. Izmestiev [3] showed that a polyhedron of the orthodiagonal anti-involutive type is the only possible candidate to disprove Stachel's conjecture. We show that the corresponding resultant is reducible, thereby confirming the conjecture. We do it in two ways: by factorization of the corresponding resultant and providing a simple geometric proof. We describe the space of parameters for which such a polyhedron exists and show that this space is non-empty. We show that a Kokotsakis polyhedron of orthodiagonal anti-involutive type is flexible and give explicit parametrizations in elementary functions and in elliptic functions of its flexion.


## 1. Introduction

A Kokotsakis polyhedron is a polyhedral surface in $\mathbb{R}^{3}$ which consists of one $n$-gon (the base), $n$ quadrilaterals attached to every side of the $n$-gon, and $n$ triangles placed between each pair of two consecutive quadrilaterals. Such surfaces are rigid in general. However, there exist special classes of flexible surfaces, that is, ones that allow isometric deformations. Flexible Kokotsakis polyhedra were studied by several authors: in the PhD thesis of Antonios Kokotsakis and in [4]; Nawratil [5] studied flexible Kokotsakis polyhedra with triangular base; and in [8] some classes of flexible Kokotsakis polyhedra with quadrilateral base were found. The question of complete classification for the case $n=4$ was a long standing open problem. Using Bricard's approach [1], Stachel and Nawratil [6, 7, 9] with the use of the following trigonometric substitution


Figure 1. Notation of dihedral angles. Zero angles correspond to a flat polyhedron with the lateral quadrilaterals facing outwards. Positive dihedral angles rotate the side quadrilaterals into the same (upper) hemisphere.

$$
\begin{equation*}
z=\tan \frac{\varphi}{2}, \quad w_{1}=\tan \frac{\psi_{1}}{2}, \quad u=\tan \frac{\theta}{2}, \quad w_{2}=\tan \frac{\psi_{2}}{2} \tag{1}
\end{equation*}
$$

expressed the dependencies between the dihedral angles $\varphi, \psi_{1}, \theta, \psi_{2}$ (see Figure 1) as a system of polynomial equations in $z, w_{1}, u, w_{2}$ :

$$
\begin{array}{lr}
P_{1}\left(z, w_{1}\right)=0, & P_{2}\left(z, w_{2}\right)=0 \\
P_{3}\left(u, w_{1}\right)=0, & P_{4}\left(u, w_{2}\right)=0 \tag{2}
\end{array}
$$

[^0]They studied this system with respect to different factorizations of the resultant $R_{12}\left(w_{1}, w_{2}\right)$ of $P_{1}$ and $P_{2}$ as polynomials in $z$, and they classified all classes of flexible Kokotsakis polyhedra with quadrilateral base for reducible $R_{12}$. Stachel conjectured that there is no flexible Kokotsakis polyhedron whose system has the irreducible resultant.

Considering a complexified version of system (2), Izmestiev [3] classified all possible classes of flexible polyhedra disregarding reality and embeddability. In particular, from the results of [3] it follows that there is only one possible candidate for a flexible Kokotsakis polyhedron with the irreducible resultant $R_{12}$. He called elements of this class Kokotsakis polyhedra of orthodiagonal anti-involutive type (OAI).

In this paper we study the properties of OAI Kokotsakis polyhedra. We find an explicit parametrization in elementary functions of all solutions of the system (2) for this type and show that there are solutions which may be embedded in the real three-dimensional space. We prove that the resultant $R_{12}$ is always reducible, which confirms Stachel's conjecture. In particular, it follows that our reducible compositions fulfill Dixon's angle condition (see [7, Theorem 2]) and their flexion is described by symmetric-symmetric case of type II in the classification given in [7]. In addition, we provide a simple geometric proof of this fact in Theorem 4 that directly follows from the results of [3].

The rest of the article is organized as follows. We give the definition of an OAI Kokotsakis polyhedron in Section 2 together with algebraic formulas for coefficients of the polynomial systems required for further analysis. In Section 3, we list our main results. In Theorem 2, we give an explicit parametrization in elementary functions of all branches of the solution of the system (2). In other words, here we come up with a description of all possible flexions of an OAI polyhedron when all planar angles are given. Then we show that the corresponding resultant is reducible by its explicit factorization. Finally, in Theorem 3, we give a full parametrization of the solution set of the system which describes planar angles of the lateral quadrilaterals in terms of the angles of the base quadrilateral and one additional parameter $\tau$. Since the proofs are quite technical, we prove these theorems in Appendix 6. In Section 4, we discuss the geometry behind the definitions and equations given in Section 2. We show that there is a noteworthy flattening effect in the case of the orthodiagonal anti-involutive surface, and give a simple proof of Stachel's conjecture using it. In fact, we briefly explain Izmestiev's approach and show that the above mentioned flattening effect is a nice geometric property that connects an algebraic method of Stachel and Nawratil and an algebraic geometry method of Izmestiev. Summarizing the results, we introduce in Section 5 an algorithm for constructing a flexible OAI Kokotsakis polyderon for a given base quadrilateral without right angles. In addition, we present results from the numerical screening of the space of parameters and give an example of a flexible OAI Kokotsakis polyhedron. In Appendix 6, we give technical results including the proofs of the main theorems.

Remark 1.1. We formulate our main results in such a manner that the reader can understand them right after reading the definitions of Section 2. This allows us to use formulas from these theorems and construct a flexible polyhedron without going into the details of the proofs.

## 2. Definitions

In this section we give an algebraic description of a Kokotsakis polyhedron of the orthodiagonal anti-involutive type. We do not discuss the nature of all equations given in this section. A comprehensive explanations can be found in [3]. However, we give a brief explanation in Section 4, where we show that Stachel's conjecture is a simple consequence of the results from [3].

Vertices of a Kokotsakis polyhedron and the values of its planar angles at the interior vertices are denoted as in Figure 2. Clearly, it is only a neighborhood of the base face $A_{1} A_{2} A_{3} A_{4}$ that matters: replacing, say, the vertex $B_{1}$ by any other point on the half-line $A_{1} B_{1}$ does not affect the flexibility or rigidity of the polyhedron.


Figure 2. Planar angles in a Kokotsakis polyhedron.

For each of the four interior vertices $A_{1}, A_{2}, A_{3}, A_{4}$, consider the intersection of the cone of adjacent faces with a unit sphere centered at the vertex. This yields four spherical quadrilaterals $Q_{i}$ with side lengths $\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}$ in this cyclic order. Let $\varphi, \psi_{1}, \theta$, and $\psi_{2}$ be the exterior dihedral angles at edges $A_{1} A_{2}, A_{2} A_{3}, A_{3} A_{4}$ and $A_{4} A_{1}$, respectively. Clearly, it suffices to parameterize these dihedral angles in order to describe a flexion of a Kokotsakis polyhedron.


Figure 3. Spherical quadrilateral $Q$
2.1. Orthodiagonal quadrilaterals. We call a spherical quadrilateral orthodiagonal if its diagonals are orthogonal. A spherical quadrilateral is called elliptic if its planar angles satisfy $\alpha \pm \beta \pm \gamma \pm \delta \neq 0(\bmod 2 \pi)$.

Let $Q$ be a spherical elliptic orthodiagonal quadrilateral with side lengths and angles denoted as in Figure 3. Then the orthodiagonality of $Q$ is equivalent to the following identity:

$$
\begin{equation*}
\cos \alpha \cos \gamma=\cos \beta \cos \delta \tag{3}
\end{equation*}
$$

(see Lemma 6.3 below).
By definition put $z=\tan \frac{\varphi}{2}, w=\tan \frac{\psi}{2}$. Then $(z, w)$ satisfy the following equation (see Lemma 4.13 in [3]):

$$
\begin{equation*}
\left(z^{2}+\lambda\right)\left(w^{2}+\mu\right)=\nu z w, \quad z, w \in \mathbb{R} P^{1}=\mathbb{R} \cup \infty, \tag{4}
\end{equation*}
$$

where $\lambda, \mu$ are the so-called involution factors defined as follows

$$
\begin{align*}
& \lambda:= \begin{cases}\frac{\tan \delta+\tan \alpha}{\tan \delta-\tan \alpha}, & \text { if } \alpha \neq \frac{\pi}{2} \text { or } \delta \neq \frac{\pi}{2}, \\
\cos \beta+\cos \gamma \\
\cos \beta-\cos \gamma & \text { if } \alpha=\delta=\frac{\pi}{2} ;\end{cases}  \tag{5}\\
& \mu:= \begin{cases}\frac{\tan \delta+\tan \gamma}{\tan \delta-\tan \gamma}, & \text { if } \gamma \neq \frac{\pi}{2} \text { or } \delta \neq \frac{\pi}{2}, \\
\frac{\cos \beta+\cos \alpha}{\cos \beta-\cos \alpha}, & \text { if } \gamma=\delta=\frac{\pi}{2} ;\end{cases} \tag{6}
\end{align*}
$$

$$
\nu:= \begin{cases}\frac{(\lambda-1)(\mu-1)}{\cos \delta}, & \text { if } \delta \neq \frac{\pi}{2},  \tag{7}\\ 2(\mu-1) \tan \alpha, & \text { if } \delta=\gamma=\frac{\pi}{2}, \\ 2(\lambda-1) \tan \gamma, & \text { if } \delta=\alpha=\frac{\pi}{2} .\end{cases}
$$

The involution factors $\lambda, \mu$ and $\nu$ are well-defined real numbers different from 0 .
2.2. Orthodiagonal anti-involutive type. A Kokotsakis polyhedron belongs to the orthodiagonal anti-involutive type if its planar angles satisfy the following conditions.

1) All quadrilaterals $Q_{i}$ are orthodiagonal and elliptic:

$$
\cos \alpha_{i} \cos \gamma_{i}=\cos \beta_{i} \cos \delta_{i}, \quad \alpha_{i} \pm \beta_{i} \pm \gamma_{i} \pm \delta_{i} \neq 0(\bmod 2 \pi)
$$

2) The involution factors at common vertices are opposite:

$$
\begin{equation*}
\lambda_{1}=-\lambda_{2}, \quad \mu_{1}=-\mu_{4}, \quad \mu_{2}=-\mu_{3}, \quad \lambda_{3}=-\lambda_{4} . \tag{8}
\end{equation*}
$$

3) The involution factors $\lambda_{i}, \mu_{i}$ and $\nu_{i}$ satisfy the following system:

$$
\begin{equation*}
\frac{\nu_{1}^{2}}{\lambda_{1} \mu_{1}}=\frac{\nu_{3}^{2}}{\lambda_{3} \mu_{3}}, \quad \frac{\nu_{2}^{2}}{\lambda_{2} \mu_{2}}=\frac{\nu_{4}^{2}}{\lambda_{4} \mu_{4}}, \quad \frac{\nu_{1}^{2}}{\lambda_{1} \mu_{1}}+\frac{\nu_{2}^{2}}{\lambda_{2} \mu_{2}}=16 . \tag{9}
\end{equation*}
$$

By (4) and (8), variables $\left(z, w_{1}, u, w_{2}\right)$ of a Kokotsakis polyhedron of the orthodiagonal antiinvolutive type satisfy the system

$$
\begin{align*}
& P_{1}=\left(z^{2}+\lambda_{1}\right)\left(w_{1}^{2}+\mu_{1}\right)-\nu_{1} z w_{1}=0, \\
& P_{2}=\left(z^{2}-\lambda_{1}\right)\left(w_{2}^{2}-\mu_{3}\right)-\nu_{2} z w_{2}=0, \\
& P_{3}=\left(u^{2}+\lambda_{3}\right)\left(w_{2}^{2}+\mu_{3}\right)-\nu_{3} u w_{2}=0,  \tag{10}\\
& P_{4}=\left(u^{2}-\lambda_{3}\right)\left(w_{1}^{2}-\mu_{1}\right)-\nu_{4} u w_{1}=0 .
\end{align*}
$$

We are interested only in the non-trivial solutions of this system, that is, the one-parametric branches of the solution set chosen in a way that each of $z, w_{1}, u, w_{2}$ is not a constant.

Remark 2.1. System (9) expresses the proportionality of the resultants $R_{12}\left(w_{1}, w_{2}\right)=\operatorname{res}_{z}\left(P_{1}, P_{2}\right)$ and $R_{34}\left(w_{1}, w_{2}\right)=\operatorname{res}_{u}\left(P_{3}, P_{4}\right)$.

Remark 2.2. There is an error in the last equation of (9) in [3], Subsection 3.1.2.
2.3. Five-parametric system. The problem of the description of all flexible Kokotsakis polyhedra of the orthodiagonal anti-involutive type is separated into two sub-problems in a natural way:

- to describe the planar angles $\left(\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}\right)$ such that the corresponding polyhedron is of the orthodiagonal anti-involutive type;
- to describe its flexion, i.e. the trajectory in the dihedral angles space.

The latter is equivalent to solving system (10), which describes the configuration space of $\left(z, w_{1}, u, w_{2}\right)$, with given coefficients $\lambda_{i}, \mu_{i}, \nu_{i}$ that satisfy (8) and (9). In Theorem 2 we give a complete parametrization with one variable of a solution set of this system.

It is more complicated to describe all admissible planar angles $\left(\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}\right)$. It is natural to consider the angles $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ as parameters. Since $\delta_{1}+\delta_{2}+\delta_{3}+\delta_{4}=2 \pi$, it is a three parametric family. From the algebraic point of view, system (9) together with identities (8) and definition (7) is a system of three equations in four variables $\left(\lambda_{1}, \mu_{1}, \lambda_{3}, \mu_{3}\right)$. It is enough to introduce one more parameter to describe a solution, which we do in Section 3.3. On the other hand, given $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ and ( $\lambda_{1}, \mu_{1}, \lambda_{3}, \mu_{3}$ ) together with (8) one can recover angles $\alpha_{i}, \gamma_{i}$ from

$$
\begin{equation*}
\tan \alpha_{i}=\frac{\lambda_{i}-1}{\lambda_{i}+1} \tan \delta_{i}, \quad \tan \gamma_{i}=\frac{\mu_{i}-1}{\mu_{i}+1} \tan \delta_{i} \tag{11}
\end{equation*}
$$

$$
\begin{equation*}
\cos \beta_{i}=\frac{\cos \alpha_{i} \cos \gamma_{i}}{\cos \delta_{i}} \tag{12}
\end{equation*}
$$

if the right-hand side is in $(-1,1)$. That is, the angles of the base quadrilateral and $\left(\lambda_{1}, \mu_{1}, \lambda_{3}, \mu_{3}\right)$ determine all the needed data to construct a Kokotsakis polyhedron.

It is not obvious that such a geometric construction exists even when one finds all the planar angles of an OAI Kokotsakis polyhedron. However, the following result gives the answer to this question.

Definition 2.3. We consider the following geometric assumptions on $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ and $\left(\lambda_{i}, \mu_{i}\right)$.

1. $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ are the angles of a quadrilateral without right angles, that is $\delta_{i}>0$ and $\delta_{i} \neq \frac{\pi}{2}$, where $i=1,2,3,4$ and $\delta_{1}+\delta_{2}+\delta_{3}+\delta_{4}=2 \pi$.
2. $\left(\lambda_{i}, \mu_{i}\right)_{1}^{4}$ satisfy systems (8), (9) with $\nu_{i}$ given by (7).
3. The angles $\left\{\alpha_{i}, \gamma_{i}, \beta_{i}\right\}_{1}^{4}$ given by (11) and (12) are well-defined, that is, the right-hand side of $(12)$ is in interval $(-1,1)$.
4. Spherical quadrilaterals $Q_{i}$ with sides $\left(\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}\right)$, where $1 \leq i \leq 4$, are elliptic.

Theorem 1. Let $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ and $\left(\lambda_{i}, \mu_{i}\right)_{1}^{4}$ meet the geometric assumptions of Definition 2.3. Then a Kokotsakis polyhedron with spherical quadrilaterals $Q_{i}$ with sides $\left(\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}\right)$ exists, it is flexible and all its possible flexion are parameterized as in Theorem 2.

In Theorem 3 of Section 3.3, we parameterize all $\left(\lambda_{i}, \mu_{i}\right)_{1}^{4}$ that satisfy systems (8), (9) with $\nu_{i}$ given by (7).

Remark 2.4. Interestingly enough, if $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ and $\left(\lambda_{i}, \mu_{i}\right)_{1}^{4}$ meet only first two of the geometric assumptions, a non-trivial solution of (10) still exists and is given by Theorem 2. However, the angles $\beta_{i}$ might fail to be determined by (12) if the absolute value of the righthand side of (12) is bigger than 1 . This follows from Lemma 6.2 and the condition of the existence of non-trivial solution in Theorem 2.

## 3. Results

3.1. Parametrization of motion. The form of the solution in real numbers of the system (10) depends on the signs of $\left(\lambda_{1}, \mu_{1}, \lambda_{3}, \mu_{3}\right)$. That is, it depends on the signs of the entries in the matrix

$$
\left(\begin{array}{ll}
\lambda_{1} & \mu_{1} \\
\lambda_{2} & \mu_{2} \\
\lambda_{3} & \mu_{3} \\
\lambda_{4} & \mu_{4}
\end{array}\right)=\left(\begin{array}{cc}
\lambda_{1} & \mu_{1} \\
-\lambda_{1} & -\mu_{3} \\
\lambda_{3} & \mu_{3} \\
-\lambda_{3} & -\mu_{1}
\end{array}\right)
$$

We call the arrangement of signs in this matrix a sign pattern. It will be shown below in Lemma 6.1 that there is a unique possible arrangement of signs up to re-enumeration of the vertices. Using this observation, we assume that

$$
\begin{equation*}
\lambda_{1}, \mu_{1}>0 \quad \text { and } \quad \lambda_{3}, \mu_{3}<0 . \tag{13}
\end{equation*}
$$

Theorem 2. Let $\lambda_{i}, \mu_{i}, \nu_{i}$ satisfy systems (9), (8) and assumption (13). Then system (10) has a real non-trivial one-parametric set of solutions if and only if $\zeta_{1}:=\frac{\left|\nu_{1}\right|}{4 \sqrt{\lambda_{1} \mu_{1}}}>1$. The solution

$$
\left\{\begin{array}{l}
z=\operatorname{sgn} \nu_{1} \cdot \sqrt{\lambda_{1}} \cdot F(t) F\left(t+\frac{\pi}{2}\right)  \tag{14}\\
w_{1}=\sqrt{\mu_{1}} \cdot F(t) F\left(t-\frac{\pi}{2}\right), \\
u=\operatorname{sgn} \nu_{4} \cdot \sqrt{-\lambda_{3}} \cdot \frac{\sqrt{\zeta_{1}+1} \pm\left(G(t)+G\left(t+\frac{\pi}{2}\right)\right)}{V(t)-V\left(t+\frac{\pi}{2}\right)} \\
w_{2}=\sqrt{-\mu_{3}} \cdot \frac{\operatorname{sgn}\left(\nu_{1} \nu_{2}\right) \sqrt{\zeta_{1}+1} \pm \operatorname{sgn}\left(\nu_{3} \nu_{4}\right)\left(G(t)+G\left(t-\frac{\pi}{2}\right)\right)}{V(t)-V\left(t-\frac{\pi}{2}\right)}
\end{array}\right.
$$

where

$$
\begin{aligned}
& F(t)=\sin t \sqrt{\zeta_{1}-1}+\sqrt{1+\left(\zeta_{1}-1\right) \sin ^{2} t} \\
& G(t)=\sin t \sqrt{1+\left(\zeta_{1}-1\right) \sin ^{2} t} \\
& V(t)=\sin t \sqrt{1+\left(\zeta_{1}-1\right) \cos ^{2} t}
\end{aligned}
$$

set has four branches, which, up to the symmetry $\left(z, w_{1}, u, w_{2}\right) \rightarrow-\left(z, w_{1}, u, w_{2}\right)$, are given by
$t \in[0,2 \pi)$ is a parameter and the choice of $\pm$ in $u$ and $w_{2}$ is simultaneous.
According to the assumption (13) and by condition $\zeta_{1}>1$, we see that all values under square roots in the previous formulas are positive.
3.2. Reducible resultant. To simplify the equations of system (10) and (9), we scale the variables and consider a reduced system (after the re-enumeration of vertices so that $\lambda_{1}, \mu_{1}>0$ and $\lambda_{3}, \mu_{3}<0$ ). Making the following substitution in system (10):

$$
\begin{array}{ll}
f_{1}=\frac{z}{\sqrt{\lambda_{1}}} \operatorname{sgn} \nu_{1}, & f_{3}=\frac{u}{\sqrt{-\lambda_{3}}} \operatorname{sgn} \nu_{4}, \\
g_{1}=\frac{w_{1}}{\sqrt{\mu_{1}}}, & g_{3}=\frac{w_{2}}{\sqrt{-\mu_{3}}} \operatorname{sgn}\left(\nu_{1} \nu_{2}\right), \\
\zeta_{i}=\frac{\left|\nu_{i}\right|}{4 \sqrt{\left|\lambda_{i} \mu_{i}\right|}} \text { for } i=1,2,4, & \zeta_{3}=\frac{\nu_{3}}{4 \sqrt{\lambda_{3} \mu_{3}}} \operatorname{sgn}\left(\nu_{1} \nu_{2} \nu_{4}\right) ; \tag{15}
\end{array}
$$

we get 4 equations:

$$
\left\{\begin{array}{l}
\left(f_{1}^{2}+1\right)\left(g_{1}^{2}+1\right)=4 \zeta_{1} f_{1} g_{1}  \tag{16}\\
\left(f_{1}^{2}-1\right)\left(g_{3}^{2}+1\right)=4 \zeta_{2} f_{1} g_{3} \\
\left(f_{3}^{2}-1\right)\left(g_{3}^{2}-1\right)=4 \zeta_{3} f_{3} g_{3} \\
\left(f_{3}^{2}+1\right)\left(g_{1}^{2}-1\right)=4 \zeta_{4} f_{3} g_{1},
\end{array}\right.
$$

where $\zeta_{1}, \zeta_{2}$ and $\zeta_{4}$ are positive.
System (9) which describes conditions for the proportionality of the resultants can be rewritten in the form

$$
\left\{\begin{array}{l}
\zeta_{1}^{2}=\zeta_{3}^{2}  \tag{17}\\
\zeta_{2}^{2}=\zeta_{4}^{2} \\
\zeta_{1}^{2}-\zeta_{2}^{2}=1
\end{array}\right.
$$

Since a linear substitution does not change the property of irreducibility of the resultant, we consider the resultant $R_{12}^{r}$ of the first two equations of reduced system (16) as polynomials in $f_{1} . R_{12}^{r}$ is given by:

$$
\begin{aligned}
\frac{1}{4} R_{12}^{r}=1+2\left(1-2 \zeta_{1}^{2}\right) g_{1}^{2}+2\left(1+2 \zeta_{2}^{2}\right) g_{3}^{2}+g_{1}^{4}+ & 4\left(1-2 \zeta_{1}^{2}+2 \zeta_{2}^{2}\right) g_{1}^{2} g_{3}^{2}+g_{3}^{4} \\
& +2\left(1+2 \zeta_{2}^{2}\right) g_{1}^{4} g_{3}^{2}+2\left(1-2 \zeta_{1}^{2}\right) g_{1}^{2} g_{3}^{4}+g_{1}^{4} g_{3}^{4}
\end{aligned}
$$

$$
\begin{equation*}
s=\frac{\delta_{1}-\delta_{2}+\delta_{3}-\delta_{4}}{4}, \quad x=\frac{\delta_{1}-\delta_{3}}{2}, \quad y=\frac{\delta_{2}-\delta_{4}}{2} \tag{18}
\end{equation*}
$$

Then there exists $\tau \in[0,2 \pi)$ and a proper choice of signs such that

$$
\begin{equation*}
\lambda_{i}=\frac{1 \pm \sqrt{1-r_{i}^{2}}}{r_{i}}, \quad \mu_{i}=\frac{1 \pm \sqrt{1-c_{i}^{2}}}{c_{i}} \quad \text { for } \quad i=1,3 \tag{19}
\end{equation*}
$$

where $r_{i}, c_{i}$ are functions of $(\tau, x, y, s)$ with the symmetries

$$
\begin{align*}
& c_{1}(\tau, x, y, s)=r_{1}(\tau+\pi, x,-y, s) \\
& r_{3}(\tau, x, y, s)=r_{1}(\tau,-x,-y, s)  \tag{20}\\
& c_{3}(\tau, x, y, s)=r_{1}(\tau+\pi,-x, y, s)
\end{align*}
$$

$$
\begin{equation*}
r_{1}(\tau, x, y, s)=\frac{N+S \sqrt{L}}{2 D} \tag{21}
\end{equation*}
$$

where

$$
\begin{aligned}
S & =s_{10} \cos \tau+s_{01} \sin \tau \\
L & =l_{20} \cos ^{2} \tau+l_{11} \sin \tau \cos \tau+l_{02} \sin ^{2} \tau \\
N & =n_{20} \cos ^{2} \tau+n_{11} \sin \tau \cos \tau+n_{02} \sin ^{2} \tau \\
D & =d_{20} \cos ^{2} \tau+d_{11} \sin \tau \cos \tau+d_{02} \sin ^{2} \tau
\end{aligned}
$$

with

$$
\begin{aligned}
s_{10} & =\cos 2 x+\cos 2 y+2 \cos 2 s \\
s_{01} & =\cos 2 x-\cos 2 y \\
l_{20} & =(\cos 2 x-\cos 2 y)^{2}+8 \cos 2 s(\cos 2 x+\cos 2 y) \\
l_{11} & =2(\cos 2 x-\cos 2 y)(\cos 2 x+\cos 2 y+2 \cos 2 s) \\
l_{02} & =(\cos 2 x+\cos 2 y-2 \cos 2 s)^{2} \\
n_{20} & =\sin (x+y)(\sin (x-3 y)+\sin (3 x-y)+6 \sin (x-y+2 s)-2 \sin (x-y-2 s)) \\
n_{11} & =8\left(\sin ^{2}(x+s) \cos ^{2}(x-s)+\sin ^{2}(y-s) \cos ^{2}(y+s)\right) \\
n_{02} & =(\cos 2 y-\cos 2 x)(\cos 2 x+\cos 2 y-2 \cos 2 s) \\
d_{20} & =\cos (x-y)(\cos (3 x+y)+\cos (x+3 y)+2 \cos (x+y-2 s)+4 \cos (x+3 s) \cos (y-s)) \\
d_{11} & =\cos 4 x-\cos 4 y-4 \sin (x+y) \sin (x-y+2 s) \\
d_{02} & =\sin (x-y)(\sin (x+3 y)-\sin (3 x+y)+2 \sin (x+y-2 s)-4 \cos (x+3 s) \sin (y-s))
\end{aligned}
$$

It is very intriguing that for each value of $\tau$ in this theorem such that all values under square roots are positive, system (10) has real non-trivial solutions (see Lemma 6.2 below).

## 4. Geometric properties

4.1. The configuration space of a spherical quadrilateral. Side lengths and angles of a spherical quadrilateral $Q$ are denoted as in Figure 4, and $z=\tan \frac{\varphi}{2}, w=\tan \frac{\psi}{2}$.


Figure 4. Spherical quadrilateral $Q$. The angle between $\beta$ and $\gamma$ is denoted as $\chi$.
As was discovered by Bricard [1], the equation of the configuration space of $Q$ has the form

$$
\begin{equation*}
P(z, w)=a_{22} z^{2} w^{2}+a_{20} z^{2}+a_{02} w^{2}+2 a_{11} z w+a_{00}=0, \tag{22}
\end{equation*}
$$

where the coefficients are trigonometric functions of the angles of $Q$.
By allowing $z$ to take complex values, we arrive at a complex algebraic curve

$$
\begin{equation*}
C=\left\{(z, w) \in \mathbb{C} P^{1} \times \mathbb{C} P^{1} \mid P(z, w)=0\right\} \tag{23}
\end{equation*}
$$

with two coordinate projections

$$
\begin{equation*}
P_{z}: C \rightarrow \mathbb{C P}^{1} \ni z \quad \text { and } \quad P_{w}: C \rightarrow \mathbb{C P}^{1} \ni w . \tag{24}
\end{equation*}
$$

If $Q$ is of elliptic type then $C$ is an elliptic curve. The projections in (24) are two-fold branched covers with exactly 4 points in the branch set. Since equation (22) is quadratic in one variable, one can easily find explicit formulas for points of the branch set by solving a quadratic equation. The complete classification of different classes of spherical quadrilaterals and their branch sets can be found in [3] Subsection 2.4 and Lemma 4.10.

Denote by

$$
\begin{array}{rlrl}
i: C & \rightarrow C & j: C & \rightarrow C \\
(z, w) & \mapsto\left(z, w^{\prime}\right) & (z, w) & \mapsto\left(z^{\prime}, w\right)
\end{array}
$$

the deck transformations of $P_{w}$ and $P_{z}$. If $z$ and $w$ are real, we can give a geometric interpretation of these involutions: $i$ and $j$ act by folding the quadrilateral along one of its diagonals, see Figure 5. This means that realizable branch points correspond to the degenerate case when quadrilateral becomes a triangle. Hence, we conclude

Lemma 4.1. Let $Q$ be a spherical quadrilateral of elliptic type. We use the notation as in Figure 4 and put $z=\tan \frac{\varphi}{2}$. Then $z$ belongs to the branch set of $P_{z}$ if and only if $\chi=0(\bmod \pi)$.


Figure 5. Involutions $i$ and $j$ on the configuration space of a quadrilateral.


Figure 6. Two coupled spherical quadrilaterals associated with the edge $A_{1} A_{2}$. The two marked angles are required to stay equal during the deformation.
4.2. Scissors-like linkage. Two adjacent vertices of the base quadrilateral share common dihedral angles, which means the corresponding spherical quadrilaterals are coupled by means of the angle (see Figure 6). There are four such scissors-like linkages corresponding to edges $A_{1} A_{2}, A_{2} A_{3}, A_{3} A_{4}$ and $A_{4} A_{1}$ of the base quadrilateral in a Kokotsakis polyhedron. Isometric deformations of the polyhedron correspond to motions of these linkages. Algebraically speaking, these linkages are described by the rows and columns of system (2), and the system itself describes possible dihedral angles of the polyhedron. Moreover, a Kokotsakis polyhedron is flexible if and only if the system of polynomial equations (2) has a one-parameter family of solutions over the reals ([3], Lemma 2.2). One possible approach to find a solution is as follows:
(1) consider a pair of opposite edges $A_{1} A_{2}$ and $A_{3} A_{4}$ together with the corresponding linkages, which are described by systems $\left\{P_{1}=0, P_{2}=0\right\}$ and $\left\{P_{3}=0, P_{4}=0\right\}$, respectively;
(2) exclude common variables $z$ and $u$ in the corresponding systems by computing the resultants $R_{12}\left(w_{1}, w_{2}\right)=\operatorname{res}_{z}\left(P_{1}, P_{2}\right)$ and $R_{34}\left(w_{1}, w_{2}\right)=\operatorname{res}_{u}\left(P_{3}, P_{4}\right)$.
A polyhedron is flexible if and only if the algebraic sets $R_{12}=0$ and $R_{34}=0$ have a common irreducible component. This means that they are reducible or irreducible simultaneously. Stachel and Nawratil described all flexible classes for reducible $R_{12}$ and $R_{34}$.

Remark 4.1. Geometrically speaking, the zero set of the resultant $R_{12}\left(w_{1}, w_{2}\right)$ gives an implicit dependence of the non-common angles of spherical quadrilaterals $Q_{1}$ and $Q_{2}$, and $R_{34}$ does the same for $Q_{3}$ and $Q_{4}$. A polyhedron is flexible if and only if these "dependencies" have a common branch, that is, on can join the scissors-like linkages $\left(Q_{1}, Q_{2}\right)$ and $\left(Q_{3}, Q_{4}\right)$ during the corresponding flexion.

Lemma 4.3. Let spherical orthodiagonal quadrilaterals $Q_{1}$ and $Q_{2}$ form a scissor-like linkage as in Figure 7 with the involution factors $\lambda_{1}$ and $\lambda_{2}$ in the common vertex. Then

$$
\begin{array}{r}
\lambda_{1}=\lambda_{2} \Leftrightarrow \xi_{1}=\xi_{2}, \xi_{1}^{\prime}=\xi_{2}^{\prime} \text { during flexion } \\
\lambda_{1}=-\lambda_{2} \Leftrightarrow \xi_{1}=\xi_{2}^{\prime}, \xi_{1}^{\prime}=\xi_{2} \text { during flexion }
\end{array}
$$



Figure 7. To the proofs of Lemma 4.3 and Lemma 4.4.
Izmestiev [3] considers the complexified configuration space of a spherical linkage (that is, the set of all complex solutions of the corresponding polynomial system). The following lemma is a direct consequence of the results of [3]. It follows from Lemma 5.1 and assertion (4) of Lemma 4.10 of that paper.

Lemma 4.2. Let polynomial system $\left\{P_{1}\left(z, w_{1}\right)=0=P_{2}\left(z, w_{2}\right)\right\}$ describe a scissors-like linkage of two elliptic spherical quadrilaterals $Q_{1}$ and $Q_{2}$. The resultant $R_{12}\left(w_{1}, w_{2}\right)$ is reducible if and only if the branch sets of $w_{1}$ and $w_{2}$ coincide.

Remark 4.2. We note that Izmestiev understands the property of irreducibility of a coupling of two spherical quadrilaterals as the irreducibility of an algebraic set that describes the corresponding complexified configuration spaces. One should not confuse the properties of irreducibility of the resultant with the irreducibility of the coupling.

As was mentioned above, Izmestiev showed that there is only one possible candidate for a flexible Kokotsakis polyhedron with irreducible resultant $R_{12}$. In particular, all scissors-like linkages of such a polyhedron must be of special type, which the author called anti-involutive coupling.
4.3. Anti-involutive coupling and its properties. A coupling of two orthodiagonal quadrilaterals is called anti-involutive if their involution factors at the common vertex are opposite, e.g. $\lambda_{1}=-\lambda_{2}$ for two coupled spherical quadrilaterals associated with the edge $A_{1} A_{2}$.

Proof. Equation (4) implies that the involutions of an orthodiagonal quadrilateral are given by

$$
i(z, w)=\left(\lambda z^{-1}, w\right), \quad j(z, w)=\left(z, \mu w^{-1}\right) .
$$

Therefore, the condition $\lambda_{1}=\lambda_{2}$ is equivalent to the compatibility of the involutions of the coupled quadrilaterals at their common vertex:

$$
i_{1}\left(z, w_{1}\right)=\left(z^{\prime}, w_{1}\right), \quad i_{2}\left(z, w_{2}\right)=\left(z^{\prime \prime}, w_{2}\right) \Rightarrow z^{\prime}=z^{\prime \prime}
$$

On the other hand, the involution $i_{1}$ changes the angle of the first quadrilateral at the common vertex from $\xi_{1}+\xi_{1}^{\prime}$ to $\xi_{1}-\xi_{1}^{\prime}$, and the angle of the second quadrilateral from $\xi_{2}+\xi_{2}^{\prime}$ to $\xi_{2}-\xi_{2}^{\prime}$. The compatibility means that the angles at the common vertex remain equal after the involution: $\xi_{1}-\xi_{1}^{\prime}=\xi_{2}-\xi_{2}^{\prime}$. Since $\xi_{1}+\xi_{1}^{\prime}=\xi_{2}+\xi_{2}^{\prime}$, we conclude that $\xi_{1}=\xi_{2}, \xi_{1}^{\prime}=\xi_{2}^{\prime}$.

Observe that exchanging $\alpha$ and $\delta$ changes the sign of the involution factor. Thus, if we rotate the left half of the scissors linkage, then the condition $\lambda_{1}=-\lambda_{2}$ transforms to $\lambda_{1}=\lambda_{2}$. On the other hand, this exchanges the angles $\xi_{2}$ and $\xi_{2}^{\prime}$. This proves the second part of the lemma.

Lemma 4.4. In an anti-involutive ( $\lambda_{1}=-\lambda_{2}$ ) linkage of two orthodiagonal quadrilaterals, if one of the angles $\chi_{1}$ and $\chi_{2}^{\prime}$ in Figure 7 becomes 0 or $\pi$, then the other one does as well.

Proof. Indeed, $\chi_{1}=0(\bmod \pi)$ if and only if $\xi_{1}=0(\bmod \pi)$.
4.4. Geometric properties of orthodiagonal anti-involutive type polyhedra. From the definition of an OAI polyhedron, we see that all four linkages of it are anti-involutive. Applying Lemma 4.4, we observe the following flattening effect.

Corollary 4.1. In a flexible OAI Kokotsakis polyhedron, the dihedral angles at the three bold edges in Figure 8 vanish or become $\pi$ at the same time.


Figure 8. A partial flattening of a flexible OAI polyhedron.

Using Theorem 2, it is quite easy to understand what kind of flattening has a flexible OAI Kokotsakis polyhedron.

Lemma 4.5. Two dihedral angles of a flexible OAI Kokotsakis polyhedron never become 0 or $\pi$. The other two dihedral angles become 0 and $\pi$ on each branch of the solution given by Theorem 2

Proof. We use the notation of Section 3.1. If one of the angles $\varphi, \psi_{1}, \theta, \psi_{2}$ is 0 or $\pi$ then the corresponding variable is 0 or $\infty$. Since substitution (15) just scales the variables, the same happens to the corresponding variable of system (16). Let us consider it for simplicity. By Subsection 6.1.2, the variables of the first equation of (16) (with positive signature) never become 0 or $\infty$. Also, each branch of the solution of the whole system contains all solutions of the first equation up to symmetry $\left(f_{1}, g_{1}\right) \rightarrow-\left(f_{1}, g_{1}\right)$. Moreover, since $\zeta_{1}>1$, then $f_{1}=1$ at two different points of the solution of the first equation. Therefore, by the second equation, $g_{3}$ becomes 0 or $\infty$ at these points on each branch of the solution of the whole system. It is easy to see that $g_{3}$ is different at these points. So it must be 0 and $\infty$ at these points, and it can have these values only at these points. A similar argument works for $g_{1}$ and $f_{3}$.

Summarizing all results, we prove Stachel's conjecture.
Theorem 4. Stachel's conjecture is true. The resultant $R_{12}$ of a real flexible OAI Kokotsakis polyhedron is reducible.

Proof. We use the notation of the first three sections. By Lemma 4.2, it is enough to show that the branch sets for $w_{1}$ and $w_{2}$ coincide. By Lemma 4.1, it suffices to prove that the dihedral angles at edges $A_{1} B_{1}$ and $A_{2} B_{2}$ (see Figure 2) become 0 or $\pi$ simultaneously at four different points. This follows from Corollary 4.1 and Lemma 4.5.

Remark 4.3. It is easy to show that the resultant $R_{12}$ is reducible by a direct calculation of the branch sets for $w_{1}$ and $w_{2}$ from the first two equations of system (10). The first equation of (8) and the third equation of (9) implies that the branch sets coincide. This observation works even for non-real polyhedra.

Lemma 4.6. The resultant $R_{12}^{r}\left(g_{1}, g_{3}\right)$ of the first two equations of system (16) is reducible if and only if $\zeta_{1}^{2}-\zeta_{2}^{2}=1$.

Proof. It is easy to see that $f_{1}$ belongs to the branch set of $g_{1}$ if and only if

$$
\frac{4 \zeta_{1} f_{1}}{f_{1}^{2}+1}= \pm 2
$$

From here, the branch set is

$$
\left\{ \pm \zeta_{1} \pm \sqrt{\zeta_{1}^{2}-1}\right\}
$$

Similar, the branch set of $g_{3}$ for the second equation is

$$
\left\{ \pm \zeta_{2} \pm \sqrt{\zeta_{2}^{2}+1}\right\}
$$

These sets coincide exactly when $\zeta_{1}^{2}-\zeta_{2}^{2}=1$.
4.5. Dixon's angle condition. As it was mentioned in the introduction, Nawratil gave a complete classification of all reducible scissors-like linkages. Checking his classification, we understood that any OAI Kokotsakis polyhedron satisfies so-called Dixon's angle condition. It follows from the algebraic description of this condition given in item b of Corollary 1 in [7], which is a simple consequence of (8) and (9). In our notations Dixon's angle condition can be expressed in the following form:
The dihedral angles at the edges $e_{1}$ and $e_{2}$ in Figure 8 have always the same absolute value.
By the above mentioned results, we see that Dixon's angle condition implies that the correspondent branch sets coincide. That is, this condition brings together an algebraic method of Stachel and Nawratil and an algebraic geometry method of Izmestiev.

## 5. Discussion

5.1. Symmetries. Here we discuss some symmetries hidden in our equations.

First of all, the symmetry $\left(z, w_{1}, u, w_{2}\right) \rightarrow\left(-z,-w_{1},-u,-w_{2}\right)$ or, equivalently, in terms of the dihedral angles $\left(\varphi, \psi_{1}, \theta, \psi_{2}\right) \rightarrow\left(-\varphi,-\psi_{1},-\theta,-\psi_{2}\right)$, is just the symmetry with respect to the plane of the base quadrilateral. Clearly, it does not change a configuration.

Secondly, the equations allow us to make the following substitution for planar angles at vertex $A_{i}$ of the base quadrilateral:
(1) $\left(\alpha_{i}, \gamma_{i}\right) \rightarrow\left(\alpha_{i}-\pi, \gamma_{i}-\pi\right)$;
(2) $\left(\alpha_{i}, \beta_{i}\right) \rightarrow\left(\alpha_{i}-\pi, \pi-\beta_{i}\right)$;
(3) $\left(\gamma_{i}, \beta_{i}\right) \rightarrow\left(\gamma_{i}-\pi, \pi-\beta_{i}\right)$.

Clearly, the corresponding spherical quadrilateral remains elliptic, and the involution factors do not change. These symmetries may help with avoiding self-intersections of a polyhedron.
5.2. Elliptic parametrization. In [3], the author parameterized a solution of (4) in terms of the Jacobi elliptic sines. Using the results of that paper, we provide a sketch of a proof of the following result in which we give an elliptic parametrization of the solution set of system (10). A comprehensive description of the OAI flexible Kokotsakis polyhedra with the use of elliptic functions can be found in [2].

Theorem 5. Let $\lambda_{i}, \mu_{i}, \nu_{i}$ satisfy systems (9), (8) and assumption (13). Then system (10) has a real non-trivial one-parametric set of solutions if and only if $\zeta_{1}:=\frac{\left|\nu_{1}\right|}{4 \sqrt{\lambda_{1} \mu_{1}}}>1$. The solution set has four branches, which, up to the symmetry $\left(z, w_{1}, u, w_{2}\right) \rightarrow-\left(z, w_{1}, u, w_{2}\right)$, are given by

$$
\left\{\begin{array}{l}
z=\operatorname{sgn} \nu_{1} \cdot \sqrt{\lambda_{1} k} \operatorname{sn}(K+i t),  \tag{25}\\
w_{1}=\sqrt{\mu_{1} k} \operatorname{sn}\left(K+\frac{i K^{\prime}}{2}+i t\right), \\
u=i \operatorname{sgn} \nu_{4} \cdot \sqrt{-\lambda_{3} k} \operatorname{sn}\left(K+\frac{i K^{\prime}}{2} \pm \operatorname{sgn}\left(\nu_{1} \nu_{2}\right)\left(K-\frac{i K^{\prime}}{2}\right)+i t\right), \\
w_{2}=i \operatorname{sgn}\left(\nu_{1} \nu_{2}\right) \cdot \sqrt{-\mu_{3} k}, \operatorname{sn}\left(K \mp \operatorname{sgn}\left(\nu_{3} \nu_{4}\right)\left(K-\frac{i K^{\prime}}{2}\right)+i t\right)
\end{array}\right.
$$

where $k$ is the elliptic modulus given by $k=\left(\zeta_{1}-\sqrt{\zeta_{1}^{2}-1}\right)^{2}, K$ and $\frac{i K^{\prime}}{2}$ are the quarter periods of elliptic sine with modulus $k ; t \in\left[0,2 K^{\prime}\right)$ is a parameter and the choice of $\pm$ and $\mp$ in $u$ and $w_{2}$ is simultaneous.
Sketch of the proof. It was shown in [3] that the solution set of (4) has a parametrization of the form

$$
z=p \operatorname{sn}(t, k), \quad w=q \operatorname{sn}(t+\tau, k),
$$

where $\tau$ is a quarter-period of sn , that is, $\tau=n K+\frac{i K^{\prime}}{2}$, and the amplitudes $p, q$ belong to $\mathbb{R}_{+} \cup i \mathbb{R}_{+}$. Moreover, there is a table with explicit formulas for $(\lambda, \mu, \nu)$ as functions of $(p, q, k)$ in Lemma 4.17 of [3]:

$$
(\lambda, \mu, \nu)= \begin{cases}\left(\frac{p^{2}}{k}, \frac{q^{2}}{k}, \frac{2(1+k)}{k \sqrt{k}} p q\right), & \text { if } \tau=\frac{i K^{\prime}}{2} \\ \left(\frac{p^{2}}{k}, \frac{q^{2}}{k},-\frac{2(1+k)}{k \sqrt{k}} p q\right), & \text { if } \tau=2 K+\frac{i K^{\prime}}{2} \\ \left(-\frac{p^{2}}{k},-\frac{q^{2}}{k}, \frac{2 i(1-k)}{k \sqrt{k}} p q\right), & \text { if } \tau=K+\frac{i K^{\prime}}{2} \\ \left(-\frac{p^{2}}{k},-\frac{q^{2}}{k},-\frac{2 i(1-k)}{k \sqrt{k}} p q\right), & \text { if } \tau=3 K+\frac{i K^{\prime}}{2}\end{cases}
$$

From this table, we see that the elliptic modulus $k$ is a function of $\frac{\nu^{2}}{\lambda \mu}$. Since $k \in(0,1)$ and by the choice of sign pattern (26), we get that the elliptic moduli of all four spherical quadrilaterals are the same in case of an OAI polyhedron. Then it is easy to recover what a quarter-period shift one can use in each equation of (16).
Remark 5.1. It follows that the OAI type of Kokotsakis polyhedra is a special class of the elliptic equimodular type in the notation of [3].

We note that the parametrizations from Theorem 5 cannot be transformed to the parametrizations given by Theorem 2 using linear substitution of the parameter $t$ as the first one uses elliptic functions that are not elementary. However, under the substitution $t \rightarrow-\frac{K^{\prime}}{\pi} t$, the parametrization given by Theorem 5 is quite close to the parametrization given by Theorem 2 for a sufficiently small $k$.
5.3. Algorithm. In the following Theorem, we summarize all needed steps to construct a flexible Kokotsakis polyhedron of the orthodiagonal anti-involutive type and to describe its flexion using given values of the angles of the base quadrilateral. We are always looking for angles in interval $(0, \pi)$. However, one can allow to the angles $\alpha_{i}$ and $\gamma_{i}$ to be in $(0, \pi) \cup(\pi, 2 \pi)$.

Theorem 6. Given parameters $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$, the algorithm of constructing a flexible Kokotsakis polyhedron of the orthodiagonal anti-involutive type is the following:
(1) Check whether $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ meet the first of the geometric assumptions, that is, they form a quadrilateral without right angles.
(2) Calculate ( $x, y, s$ ) using substitution (18).
(3) Check whether there exists $\tau$ such that $r_{1}, r_{3}, c_{1}, c_{3}$ given by (21) and (20) satisfy inequality (35) (see Subsection 5.4 for how this can be done). Set $r_{2}=-r_{1}, r_{4}=-r_{3}$, $c_{2}=-c_{3}, c_{4}=-c_{1}$.


Figure 9. Screening result. Due to the symmetry of the system, the set is centrally symmetric with respect to $\left(\frac{\pi}{4}, \frac{\pi}{4}, \frac{\pi}{4}\right)$ and is symmetric with respect to the plane $x=y$.
(4) Calculate the angles $\alpha_{i}$ and $\gamma_{i}$ from

$$
\tan \alpha_{i}=\sigma_{i}^{\alpha} \sqrt{\frac{1-r_{i}}{1+r_{i}}} \tan \delta_{i} \quad \text { and } \quad \tan \gamma_{i}=\sigma_{i}^{\gamma} \sqrt{\frac{1-c_{i}}{1+c_{i}}} \tan \delta_{i},
$$

where $\sigma_{i}^{\boldsymbol{\bullet}}= \pm 1$ are chosen in a way to satisfy: $\sigma_{1}^{\alpha} \sigma_{2}^{\alpha}=\sigma_{3}^{\alpha} \sigma_{4}^{\alpha}=\sigma_{1}^{\gamma} \sigma_{4}^{\gamma}=\sigma_{2}^{\gamma} \sigma_{3}^{\gamma}=1$. Select angles to be from 0 to $\pi$.
(5) Check that $\frac{\cos \alpha_{i} \cos \gamma_{i}}{\cos \delta_{i}} \in(-1,1)$. Find $\beta_{i}$ using (12).
(6) Check that obtained spherical quadrilaterals $Q_{i}$ with sides $\left(\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}\right)$ are elliptic.
(7) Find $\lambda_{i}, \mu_{i}, \nu_{i}$ using formulas (5), (6) and (7).
(8) Re-enumerate vertices to get (26) as a sign pattern of the system.
(9) Find ( $z, w_{1}, u, w_{2}$ ) using formulas of Theorem 2. Calculate dihedral angles using

$$
\left(\varphi, \psi_{1}, \theta, \psi_{2}\right)=\left(2 \arctan z, 2 \arctan w_{1}, 2 \arctan u, 2 \arctan w_{2}\right) .
$$

Proof. If all requirements are fulfilled, an existence and flexibility of a Kokotsakis polyhedron constructing using this algorithm follows from Theorem 1 and Theorem 2.
Remark 5.2. We do not check whether the obtained surface is not self-intersecting. There remains a possibility that at every moment during flexion, there are two intersecting facets.

It's possible to derive an implicit formula for a set of $(x, y, s)$ (or, equivalently, $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ ) for which there exists $\tau$ such that $-1<r_{i}, c_{i}<1$ and the angles $\beta_{i}$ are well-defined. However, the resulting formula is likely to be enormous and incomprehensive. We instead explain how to check the conditions and show the result of numerical screening in the next subsection.
5.4. Screening of the space of parameters $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$. In Theorem 6 , one need to find a proper parameter $\tau$ for a given angles $(x, y, s)$ (or, equivalently, $\left(\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right)$ ). We run a computer screening and found numerically the set of $(x, y, s)$ for which there exists $\tau$ such that $\left(\lambda_{i}, \mu_{i}\right)$ are given by Theorem 3 and $\left\{\delta_{1}, \delta_{2}, \delta_{3}, \delta_{4}\right\}$ meet the geometric assumptions. The result of our numerical computation is in Figure 9, here we do not show surfaces on which at least one $Q_{i}$ is not elliptic. Figure 10 shows the result of screening in ( $\delta_{1}, \delta_{3}, \delta_{2}$ ) coordinates for convex and non-convex base quadrilaterals.

We claim that it can be done with an arbitrary precision or implicit formulas for the boundary can be found, as all the restrictions on $\tan \tau$ are polynomial.


Figure 10. Sets of ( $\delta_{1}, \delta_{2}, \delta_{3}$ ) with an existing flexible OAI polyhedron (left) with convex base quadrilateral and (right) with non-convex base quadrilateral $\left(\delta_{4}>\pi\right)$.

More precisely,
(1) $\frac{L(\tau, x, y, z)}{\cos ^{2} \tau} \geq 0$ is clearly a polynomial in $\tan \tau$.
(2) $r_{i}<C=$ const (and others comparisons for $r_{i}, c_{i}$ with constant $C$ ) is equivalent to a proper system of the following polynomial inequalities

$$
\left(\frac{2 C D-N}{S \cos \tau}\right)^{2} \geq \frac{L}{\cos ^{2} \tau}, \quad\left(\frac{2 C D-N}{S \cos \tau}\right)^{2} \leq \frac{L}{\cos ^{2} \tau}
$$

and comparison of $\frac{2 C D-N}{S \cos \tau}, \frac{D}{\cos ^{2} \tau}$ and $\frac{S}{\cos \tau}$ with zero.
(3) $\beta_{i}$ is well-defined if and only if

$$
\frac{\cos ^{2} \delta_{i}}{\cos ^{2} \alpha_{i} \cos ^{2} \gamma_{i}} \leq 1
$$

By $\frac{1}{\cos ^{2} \alpha}=1+\tan ^{2} \alpha$, the latter is

$$
\cos ^{2} \delta_{i}\left(\frac{1-r_{i}}{1+r_{i}} \tan ^{2} \delta_{i}+1\right)\left(\frac{1-c_{i}}{1+c_{i}} \tan ^{2} \delta_{i}+1\right) \geq 1
$$

which is equivalent to a system of polynomial inequalities.
5.5. Example. In this section we consider the construction of a flexible OAI Kokotsakis polyhedron with angles of the base quadrilateral:

$$
\delta_{1}=1.36292, \quad \delta_{2}=1.41009, \quad \delta_{3}=1.80327, \quad \delta_{4}=2 \pi-\delta_{1}-\delta_{2}-\delta_{3}=1.70691
$$

We selected $\tau=-\arctan 60=-1.55413$. Following the procedure in Theorem 6, we get

$$
\begin{array}{llll}
\alpha_{1}=1.34086, & \alpha_{2}=1.42575, & \alpha_{3}=1.69859, & \alpha_{4}=1.81798, \\
\gamma_{1}=1.15746, & \gamma_{2}=1.13993, & \gamma_{3}=1.65410, & \gamma_{4}=1.63656 ; \\
\beta_{1}=1.11122, & \beta_{2}=1.18397, & \beta_{3}=1.61684, & \beta_{4}=1.68958 .
\end{array}
$$

In Figure 11 we show the state of the given polyhedron when we observe the flattening effect.


Figure 11. Partial flattening. Red arcs show flat dihedral angles.

For this example, we have elliptic modulus $k=0.0083$.

## 6. Appendix

### 6.1. Proof of Theorem 2.

6.1.1. Uniqueness of the sign pattern. Due to system (9), we have some restrictions to possible signs of $\left(\lambda_{i}, \mu_{i}\right)$.

Lemma 6.1. System (8-10) allows unique sign pattern of $\left(\lambda_{i}, \mu_{i}\right)$ up to cyclic re-enumeration of vertices. The only possible sign pattern is

$$
\operatorname{Sign}\left(\begin{array}{ll}
\lambda_{1} & \mu_{1}  \tag{26}\\
\lambda_{2} & \mu_{2} \\
\lambda_{3} & \mu_{3} \\
\lambda_{4} & \mu_{4}
\end{array}\right)=\left(\begin{array}{ll}
+ & + \\
- & + \\
- & - \\
+ & -
\end{array}\right) .
$$

Proof. We first show that for some vertex both $\lambda_{i}, \mu_{i}>0$. By the third equation in (9), we see that either $\lambda_{1} \mu_{1}>0$ or $\lambda_{2} \mu_{2}>0$. If in this positive product both factors are also positive, we have found the sought vertex. If however both coefficients are negative, from the fact that in a sign pattern there are always exactly four pluses and four minuses (follows from (8)), we conclude that if in sign pattern there is a row of minuses, there should be a row of pluses.

Since cyclic permutation just changes the order of vertices of the base quadrilateral, we assume that $\lambda_{1}>0$ and $\mu_{1}>0$. In other words, known signs are:

$$
\operatorname{Sign}\left(\begin{array}{ll}
\lambda_{1} & \mu_{1} \\
\lambda_{2} & \mu_{2} \\
\lambda_{3} & \mu_{3} \\
\lambda_{4} & \mu_{4}
\end{array}\right)=\left(\begin{array}{cc}
+ & + \\
- & * \\
* & * \\
* & -
\end{array}\right)
$$

We rewrite equation $P_{1}\left(z, w_{1}\right)=0$ of (10) in the following way:

$$
\left(\frac{z}{\sqrt{\lambda_{1}}}+\frac{\sqrt{\lambda_{1}}}{z}\right)\left(\frac{w_{1}}{\sqrt{\mu_{1}}}+\frac{\sqrt{\mu_{1}}}{w_{1}}\right)=\frac{\nu_{1}}{\sqrt{\lambda_{1} \mu_{1}}}
$$

Since absolute values of both parentheses are greater or equal to 2 , for this equation to have a one-parametric solution (and not discrete set of points), we obtain

$$
\frac{\nu_{1}^{2}}{\lambda_{1} \mu_{1}}>16
$$

By the third equation in (9)

$$
\frac{\nu_{2}^{2}}{\lambda_{2} \mu_{2}}=16-\frac{\nu_{1}^{2}}{\lambda_{1} \mu_{1}}<0 .
$$

Hence, $0>\lambda_{2} \mu_{2}=\lambda_{1} \mu_{3}$ or, equivalently, $\mu_{3}<0$. By the first equation in (9), we get that $\lambda_{3}<0$. Finally, the only possible sign pattern is (26).

$$
\begin{equation*}
\left(f_{1}^{2}+1\right)\left(g_{1}^{2}+1\right)=4 \zeta_{1} f_{1} g_{1} \tag{27}
\end{equation*}
$$

Recall $\zeta_{1}>0$.
One can see that $f_{1}$ and $g_{1}$ are of the same sign, and there is a symmetry of the solution $\left(f_{1}, g_{1}\right) \rightarrow\left(-f_{1},-g_{1}\right)$. We assume that $f_{1}, g_{1}>0$. Now we consider new coordinates $X=\log \left(f_{1} / g_{1}\right), Y=\log f_{1} g_{1}$. In these coordinates, the equation takes the form:

$$
\begin{equation*}
\cosh X+\cosh Y=2 \zeta_{1} \tag{28}
\end{equation*}
$$

Thus, there is a one-parametric real solution set if and only if

$$
\begin{equation*}
\zeta_{1}>1 \tag{29}
\end{equation*}
$$

which we have already shown in Lemma 6.1.
Solution to Equation (28) can be parameterized in an explicit way. Let $\omega=\operatorname{arccosh} \zeta_{1}$, $\varrho=\sqrt{2} \sinh \frac{\omega}{2}$. Then we can parameterize solution of (28) as

$$
X=\operatorname{arcsinh}(\varrho \cos t) \quad \text { and } \quad Y=\operatorname{arcsinh}(\varrho \sin t)
$$

where $t \in[0,2 \pi)$, which, by direct calculations, leads to the following representation of $g_{1}$ and $f_{1}$ :

$$
f_{1}=F(t) F\left(t+\frac{\pi}{2}\right) \quad \text { and } \quad g_{1}=F(t) F\left(t-\frac{\pi}{2}\right)
$$

where

$$
F(t)=\sqrt{\left(\zeta_{1}-1\right) \sin ^{2} t+1}+\sqrt{\zeta_{1}-1} \sin t
$$

Moreover, the solution of $(28)$ is bounded in the plane $(X, Y)$. This means that $f_{1}$ and $g_{1}$ cannot be equal 0 or $\pm \infty$. Or, speaking geometrically, if equation (27) describes the configuration space of the spherical quadrilateral of a Kokotsakis polyhedron, then the corresponding side quadrilaterals never cross the base plane (see Lemma 4.5).

### 6.1.3. Proof of Theorem 2.

Proof. By the results of Subsection 6.1 .2 and by (29) in particularly, we see that the real solution of the system is nonempty and non-trivial if and only if $\zeta_{1}>1$ (recall that $\zeta_{1}, \zeta_{2}$ and $\zeta_{4}$ are positive). Also in Subsection 6.1.2, the equation with positive sign signature was resolved, which leads to the explicit formulas for $z$ and $w_{1}$. One can prove that the formulas for $u$ and $w_{2}$ from (14) are the solution of the system by the direct substitution. However, we want to show that we have found all branches of the solution. We can consider system (16) and variables $\left(f_{1}, g_{1}, f_{3}, g_{3}\right)$ instead of $\left(z, w_{1}, \theta, w_{2}\right)$. One can see that the solution of the first equation is given in Subsection 6.1.2, the second equation is quadratic in $g_{3}$, the fourth equation is quadratic in $f_{3}$. Hence, we have two possible branches of solution for both $f_{3}$ and $g_{3}$, and there are at most four branches of the solution up to the symmetry $\left(f_{1}, g_{1}, f_{3}, g_{3}\right) \rightarrow\left(-f_{1},-g_{1},-f_{3},-g_{3}\right)$. So the only difficulty is to check that we have chosen the correct branches to meet the third equation of the system. We claim that the latter obstacle implies that there are exactly two branches of the solution with positive $f_{1}$ and $g_{1}$. Hence, we described all possible solution with positive $f_{1}$ and $g_{1}$. Indeed, fixing $f_{1}$, we fix $\left(g_{3}^{2}+1\right) / g_{3}$. Therefore, fixing $f_{1}, f_{3}, g_{1}$, we fix $\left(g_{3}^{2}-1\right) / g_{3}$. But $\left(g_{3}^{2}-1\right) / g_{3}=\left(g_{3}^{2}+1\right) / g_{3}-2 / g_{3}$ changes its value if we change the branch for $g_{3}$. Therefore, there is a unique branch of $g_{3}$ for each branch of $f_{3}$. By the symmetry $\left(f_{1}, f_{3}, g_{1}, g_{3}\right) \rightarrow\left(-f_{1},-g_{1},-f_{3},-g_{3}\right)$, we describe all branches of the solution.

### 6.2. Proof of Theorem 3.

$$
\begin{equation*}
Z_{i}=a_{i}\left(r_{i}^{-1}-1\right)\left(c_{i}^{-1}-1\right)=4 \frac{\nu_{i}^{2}}{\lambda_{i} \mu_{i}}=\frac{\left(\lambda_{i}-1\right)^{2}\left(\mu_{i}-1\right)^{2}}{4 \lambda_{i} \mu_{i} \cos ^{2} \delta_{i}} . \tag{32}
\end{equation*}
$$

Then system (30) is equivalent to the system of the following three equations $Z_{1}+Z_{2}=4$, $Z_{1}=Z_{3}$ and $Z_{2}=Z_{4}$. Adding a linear dependent equation, system (30) can be rewritten in the following equivalent way

$$
\begin{equation*}
\frac{1}{2} Z_{i}=1-(-1)^{i} \tan \tau \tag{33}
\end{equation*}
$$

where $\tau$ is parameter.
System (8) takes the form:

$$
\begin{equation*}
r_{1}=-r_{2}, \quad c_{1}=-c_{4}, \quad c_{2}=-c_{3}, \quad r_{3}=-r_{4} \tag{34}
\end{equation*}
$$

The equations of system (33) are linear in both $r_{i}$ or $c_{i}$. Since, by the definitions, $\lambda_{i}, \mu_{i}$ are real and $\nu_{i} \neq 0$, we obtain a restriction on the values of $r_{i}, c_{i}$ :

$$
\begin{equation*}
-1 \leq r_{i}, c_{i}<1 \tag{35}
\end{equation*}
$$

It is clear that $\left(r_{i}, c_{i}\right)$ has the same sign pattern as $\left(\lambda_{i}, \mu_{i}\right)$ and that substitution (31) is inverse to (19).

### 6.2.2. Proof of Theorem 3.

Proof. Restriction (35) yields positivity of radical expressions in (19). So, it suffices to prove identities (20) and (21). We start with the latter one.

Using exclusion technique to find $r_{1}(\tan \tau)$ from system (33) is equivalent to finding roots of an equation of the second degree with coefficients determined by the coefficients of the system (33). Indeed, from the equation $Z_{i}=$ const, we see that functions $c_{i}\left(r_{i}\right)$ or $r_{i}\left(c_{i}\right)$ are linear fractional transformations. Since composition of two linear fractional transformations is a linear fractional transformation, finding $c_{1}\left(r_{1}\right), r_{3}\left(c_{1}\right)$ and $c_{2}\left(r_{1}\right), r_{3}\left(c_{2}\right)$, we obtain two different linear fractional transformation functions for $r_{3}\left(r_{1}\right)$. Therefore, general solution has a form:

$$
\begin{equation*}
r_{1}(\tan \tau)=\frac{P(\tan \tau) \pm \sqrt{Q(\tan \tau)}}{R(\tan \tau)} \tag{36}
\end{equation*}
$$

where $P$ and $R$ are polynomials of 2 nd degree and $Q$ is a polynomial of 4th degree. However, we noticed that $Q$ always have a full square as a multiplier, when $\sum \delta_{i}=2 \pi$. Using our substitution (18), we managed to factorize $Q$, merge two branches in (36) and obtain (21). The explicit formulas for the coefficients were found and can be verified with the use of a system of computer algebra. We used Wolfram Mathematica 11 for this purpose.

Since equation (36) is obtained by the exclusion technique, it describes all possible $r_{1}$. Note that it has two branches, however we have only one in (21). So we need explain how we managed to merge the branches. The radical expression in (36) equals to $\frac{S^{2} L}{\cos ^{4} \tau}$. Since $S$ is the
only function that have linear terms in $\cos \tau$ and $\sin \tau$, substitution $\tau \rightarrow \tau+\pi$ changes the sign of $S$ and preserves the values of $N, L, D$ in (21), that is, it merges two branches of (36).

Let us explain identities (20). There are no more solutions as each equation of (33) is linear in either $r_{i}$ or $c_{i}$. Apart from checking (20) with direct calculations, we provide the following arguments. The dependencies on $x, y, s$ arise from re-enumeration of vertices, Explanation of a "phase shift" of $\tau$ is trickier. By the symmetry of the system (33), it is enough to explain a "phase shift" for $c_{1}$. Fixing $r_{1}$, we have exactly one possible value for $c_{1}$ by the equation $Z_{1}=2-2 \tan \tau$ and two possible branches in (36) rewritten for $c_{1}$. The "phase shift" $\tau \rightarrow \tau+\pi$ corresponds to the right choice of the branch.
6.2.3. Proof of Theorem 1. The theorem is a direct consequence of the following lemma.

Lemma 6.2. Let $\delta_{i}$ and $\left(\lambda_{i}, \mu_{i}\right)$ meet the first two of the geometric assumptions. Then (26) is the sign pattern of $\left(\lambda_{i}, \mu_{i}\right)$ up to re-enumeration of vertices (15) and $\frac{\nu_{j}^{2}}{\lambda_{j} \mu_{j}}>16$ for $j$ such that $\lambda_{j}, \mu_{j}>0$.

Indeed, to prove the theorem we need to check the following:

- spherical quadrilaterals with sides $\left(\alpha_{i}, \beta_{i}, \gamma_{i}, \delta_{i}\right)$ exist;
- system (10) has non-trivial solution.

The first assumption follows from Lemma 6.4, the second is guaranteed by Lemma 6.2.
Let us prove the lemma.
Proof. By the first equation in (30), either $r_{1}, c_{1}$ or $r_{2}, c_{2}$ are both non-zero and have the same sign. Again, by Pigeonhole principle, it implies that there is a pair with positive numbers. Without loss of generality, let it be pair $r_{1}, c_{1}$. By the second equation, the latter implies that $r_{3}, c_{3}$ has the same sign. We claim that they are negative, and, therefore, we have the case of the right sign pattern.

Assume the contrary that $r_{3}, c_{3}>0$. By the definition of $Z_{1}$ and system (33), we have

$$
2+2 \tan \tau=Z_{1}=\frac{1}{\cos ^{2} \delta_{1}}\left(\frac{1}{r_{1}}-1\right)\left(\frac{1}{c_{1}}-1\right)>0
$$

or, equivalently, $2>-2 \tan \tau$. Then, by the definition of $Z_{2}$, we have

$$
Z_{2}=\frac{1}{\cos ^{2} \delta_{2}}\left(1 / r_{1}+1\right)\left(1 / c_{3}+1\right)>\left(1 / r_{1}+1\right)\left(1 / c_{3}+1\right)>4>2-2 \tan \tau
$$

which contradicts system (33). Therefore, $r_{3}, c_{3}<0$. Then, by system (33) and (32), we obtain

$$
\frac{\nu_{1}^{3}}{\lambda_{3} \mu_{3}}=\frac{\nu_{1}^{1}}{\lambda_{1} \mu_{1}}=\frac{1}{4} Z_{3}=\frac{1}{4 \cos ^{2} \delta_{3}}\left(\frac{1}{\left|r_{3}\right|}+1\right)\left(\frac{1}{\left|c_{3}\right|}+1\right)>1 .
$$

### 6.3. Properties of orthodiagonal quadrilaterals.

Lemma 6.3. The diagonals of a spherical quadrilateral with side lenths $\alpha, \beta, \gamma, \delta$ (in this cyclic order) are orthogonal if and only if its side lengths satisfy the relation

$$
\cos \alpha \cos \gamma=\cos \beta \cos \delta
$$

Proof. Take the intersection point of two diagonals of the quadrilateral. (A pair of opposite vertices determines a big circle. We take one of the two intersection points of these two big circles.) Denote the lengths of the segments between the vertices and the intersection point of the diagonals as shown in Figure 12. By the spherical Pythagorean theorem we have $\cos a \cos b=\cos \alpha$, $\cos b \cos c=\cos \beta, \cos c \cos d=\cos \gamma, \cos d \cos a=\cos \delta$. It follows that

$$
\cos \alpha \cos \gamma=\cos a \cos b \cos c \cos d=\cos \beta \cos \delta
$$



Figure 12. To the proof of Lemma 6.3.

Lemma 6.4. Let $\alpha, \beta, \gamma, \delta \in(0, \pi)$ satisfy $\cos \alpha \cos \gamma=\cos \beta \cos \delta$. Then there exists a spherical orthodiagonal quadrilateral with side lengths $\alpha, \beta, \gamma, \delta$.
Proof. Assume that $\alpha+\beta$ is the minimum of the sum of two consecutive angles in $(\alpha, \beta, \gamma, \delta, \alpha)$. If $\delta=\alpha$ and $\beta=\gamma$, the construction is obvious. Otherwise, assume that $\delta>\alpha$. Put sides $V_{1} V_{2}$ and $V_{2} V_{3}$ with lengths $\alpha$ and $\beta$ respectively on the same great circle. Then there is a vertex $V_{4}$ such that $V_{1} V_{2} V_{4}$ is a right angle and the length of $V_{4} V_{1}$ is $\delta$. By the identity, the length of $V_{3} V_{4}$ is $\gamma$.

## References

[1] R. Bricard. Mémoire sur la théorie de l'octaèdre articulé. Journal de Mathématiques pures et appliquées, 3:113-148, 1897.
[2] I. Izmestiev. Flexible Kokotsakis polyhedra based on orthodiagonal spherical quadrilaterals. In preparation.
[3] I. Izmestiev. Classification of flexible Kokotsakis polyhedra with quadrangular base. International Mathematics Research Notices, 2017(3):715-808, 2016.
[4] A. Kokotsakis. Über bewegliche polyeder. Mathematische Annalen, 107(1):627-647, 1933.
[5] G. Nawratil. Flexible octahedra in the projective extension of the Euclidean 3-space. J. Geom. Graph., 14(2):147-169, 2010.
[6] G. Nawratil. Reducible compositions of spherical four-bar linkages with a spherical coupler component. Mechanism and Machine Theory, 46:725-742, 2011.
[7] G. Nawratil. Reducible compositions of spherical four-bar linkages without a spherical coupler component. Mechanism and Machine Theory, 49:87-103, 2012.
[8] R. Sauer and H. Graf. Über Flächenverbiegung in Analogie zur Verknickung offener Facettenflache. Mathematische Annalen, 105(1):499-535, 1931.
[9] H. Stachel. A kinematic approach to Kokotsakis meshes. Comput. Aided Geom. Design, 27(6):428-437, 2010.
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