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eMIL: Advanced emission Mössbauer spectrometer for measurements in versatile conditions

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ABSTRACT

The current work presents a contemporary design of an advanced emission Mössbauer Spectrometer: eMIL equipped with a parallel-plate avalanche detector, which has been devised and built for the Mössbauer collaboration at ISOLDE/CERN. The setup is based on emission geometry, combined with on-line/off-line isotope implantation and provides numerous advantages over conversion electron, common emission (where isotope is deposited chemically on a sample) or transmission Mössbauer spectroscopy. eMIL is designed to measure hyperfine interactions in solids under various exposures. The implemented design overcomes limitations and improves performance and handling. In the current revision, the chamber is supplied with an UV extension — allowing to perform studies of photo-catalytic materials under external light exposure. A specifically designed motorized lid-samples-holder is fully automatized, and makes it possible to study up to 4 samples loaded in a magazine within a temperature range from RT up to 1100 K and to perform angular dependent measurements in high vacuum. This work additionally briefly describes data acquisition with additional electronic blocks, vacuum and data-acquisition system construction.

1. Introduction

Mössbauer spectroscopy is among the powerful methods to locally probe structural and electronic properties in condensed matter. The method allows to analyze and quantify various atomic surroundings, magnetic states, infield magnetic arrangements of magnetic moments and coordination symmetry. Moreover, the method allows us to obtain unique information, which is out of reach for other experimental techniques [1,2]. The Mössbauer effect is based on the recoilless nuclear emission and resonant absorption of γ -radiation in studied samples and is commonly applied for absorption and emission experiments. Most of the absorption experiments are done in transmission geometry. In this case, one has the detector placed behind a thin absorber (sample) and a spectrum is measured as a transmission from a Doppler shifted radioactive source. The resulting spectrum is recorded as a function of the relative source velocity. In emission experiments (emission Mössbauer spectroscopy or eMS) the radioactive parent isotope is usually introduced in the sample of interest by doping or implantation.

In the case of the on-line eMS experiments, isotopes are produced and implanted at facilities such as ISOLDE/CERN [3,4] and measurements may take place simultaneously alongside with implantation or

time delayed (implant and measure few minutes later) [5]. The main advantage of such studies is the amount of required parent isotope. For instance, with a $^{57}\text{Mn}^+$ beam flux of $2 \times 10^8/\text{s}$ one obtains peak implanted ion concentration varying from 2.5×10^{-4} at.% for the first measurement to total amount of 2×10^{-3} at.% when the final implantation of a sample is performed (6–7 measurements) [6]. The remarkably low concentrations assured single ion implantation and no overlap of implantation damage cascades. In transmission experiments, a common requirement is averaging near 0.05–0.1 at.% (for enriched ^{57}Fe) [1].

The low quantity and quick measuring time allow one to measure without the risk of precipitation, for instance isolated defects in diluted magnetic systems [7–10]. When the implantation is used to introduce the radioactive probe, the damage imparted on the host material invokes additional concern on the determination of the exact location of implanted atoms and their properties. Limit of the concentration for eMS studies is below the amorphization threshold, which is generally in the order of $10^{13} - 10^{14}$ ions/cm² depending on the material. When a sample is implanted with fluences below 10^{13} ions/cm², the implanted atoms cease to move in amorphous zones (clusters, of few nm). These

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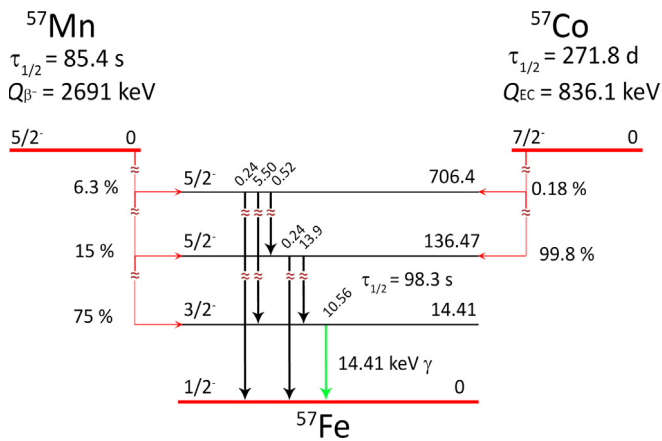


Fig. 1. Main decay branches of isotopes which are commonly employed in on-line (^{57}Mn) and off-line (^{57}Co) studies. Both decay into ^{57}Fe .

zones can be annealed (recrystallized) at relatively low temperatures, while amorphous layers (overlapping clusters form a continuous zone) remain stable up to elevated temperatures [11,12]. Unless done on purpose, a common fluence applied in on-line eMS seldom goes above the amorphization threshold.

The situation becomes even more tangled if during a decay to the ground state the recoil energy (e.g., $\langle E_R \rangle = 40$ eV for ^{57}Mn β^- decay) is sufficient enough to displace an atom to a random location in the lattice. However, it is possible to perform annealing in order to moderate the implantation damage.

The majority of the Mössbauer studies are based on the 14.4 keV γ -ray of ^{57}Fe , which is generally produced when ^{57}Mn or ^{57}Co isotopes decay into the first excited state, as presented in Fig. 1.

With the first attempts to perform aforementioned on-line experiments dating back to ISOLDE in 1996. Equipment used for the measurements has been constantly upgraded in the last 20 years providing new possibilities for measurements (such as quenching experiments), however all these modifications have been bearing a rather temporary character and the measurements with new settings could turn out fairly time consuming. In this paper, we present a newly developed emission Mössbauer spectrometer eMIL (emission Mössbauer system from ILMenau), which exploits contemporary electronic solutions and provides great flexibility towards its application providing a plug and play experience.

2. The spectrometer's design

Fig. 2 shows the electrical part of the setup, which follows the classical principle common for Mössbauer spectroscopy [1,13]. The main differences originate from the fact that the source (sample) is kept inside the implantation chamber (the eMIL's main chamber) and that imposes limitations on the chamber design. It is noteworthy, that before a sample becomes the source, it gets implanted with parent isotopes of interest. At ISOLDE/CERN facility [14,15] the procedure commences with the protons production, which are later accelerated up to 50 MeV in a linear accelerator. Consecutively, the protons are introduced into the Proton Synchrotron Booster, which has four superimposed rings with radius of 25 m where the protons undergo acceleration up to 1.4 GeV. Afterwards, radioactive ions are produced utilizing 1.4 GeV proton-induced nuclear spallation fragmentation and fission reactions in a UC_2 target. During proton irradiation the target is kept approximately at 2300 K, for the facilitation of the new radioactive nuclei diffusion towards the target's surface. These radioactive elements are later ionized by one of the numerous ion sources (lasers), then are accelerated by the extraction electrode to 30–60 kV and mass separated by the separator magnets. At ISOLDE there are two separator magnets:

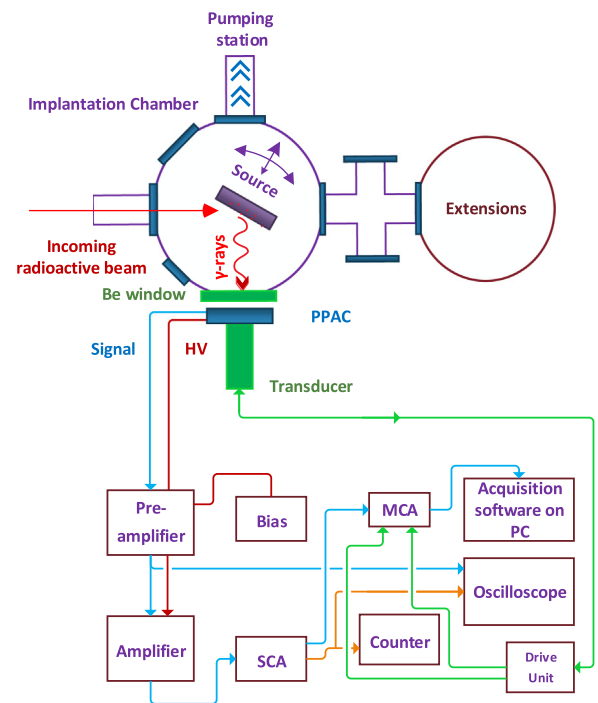


Fig. 2. Schematic diagram of the implantation chamber and electronics used in the current implementation. Arrows above the sample indicate the sample holder axis of freedom (Z and Y).

General Purpose Separator (GPS) and High Resolution Separator. The GPS separator allows the simultaneous extraction of three mass separated beams in two low mass (GLM) and high mass (GHM) beamlines. Hence, once mass separated, the ions are implanted into the material of interest [16,17] at the GLM beamline where the Mössbauer setup is plugged in. The lasers are used because of their element-specific selectivity and together with the isobar-specific selectivity of the separation magnet it becomes possible to suppress significant quantities of contaminants [18,19].

As it is shown in Fig. 2, the ion implantation is performed at the incident angle of $\theta_I = 30^\circ$, whilst measured at $\theta_M = 60^\circ$ relative to the sample surface normal. Due to its special design the detector can only be installed on the transducer outside the eMIL's chamber.

Gas filled detectors (which are also in use all over experimental nuclear physics) have been used since the 60s. The most known variation of such detector is gas-flow proportional counter, commonly used in conversion electron Mössbauer spectroscopy (CEMS). Another variation of a proportional counter is Parallel-Plate Avalanche Counter (PPAC). In contrast to the CEMS proportional counters, PPAC provides a higher signal-background ratio, higher yields and swifter time response (10^{-9} s) [20]. The main principle is based on recording of scattered radiation emitted by the absorber nuclei after resonant absorption and following de-excitation. In this case the nuclear excited state decays via internal conversion and a following emission of a conversion electron (mainly, $\approx 90\%$). This is followed by another emission of an Auger electron or X-ray, when a vacancy (primarily in K-shell) is replenished. As an alternative, a re-emission of a 14.4 keV γ -photon might take place [1]. From the construction point of view the PPAC detector follows the same idea, although is simpler built. [20,21]. PPAC detector generally consists of two electrodes (few mm between), namely a cathode made of a single line Mössbauer absorber (in the current case it is a steel plate doped with enriched ^{57}Fe) and an anode made of graphite. Just as in gas-flow proportional counters one has to rely on the appropriate gas (acetone or 90%He %10 CH_4) inside the detector. The events detection is attained by gas multiplication. Potential difference is applied to the



Fig. 3. A photo of the current electronic rack. The flight case includes a NIM crate with data acquisition modules (SCA, MCA, velocity unit), computer, oscilloscope, HV unit.

plates, which generates a uniform electric field repelling conversion electrons away from the surface. These electrons ionize the gas between the plates thus causing a formation of avalanches moving towards the electrodes (Townsend discharge). Since there is broad distribution of the ionization and the avalanches growth — there is no collectable energy resolution.

2.1. Hardware

In the current work, the same PPAC detector is used [20]. It is placed relative to the incoming beam (at 90°). In order to remove the high energetic electrons emitted during the β^- decay a Beryllium window is placed on CF40 flange of the implantation chamber. This Beryllium filter effectively blocks electrons and X-rays, allowing γ -radiation to pass through with approximate transmission of 90%. The detector is mounted on a transducer with constant acceleration. It is directly connected to an Ortec 142PC pre-amplifier input. The incoming signal is being transferred to both an amplifier and an oscilloscope. The latter one is used for the adjusting of pulse height and amplitude. The signal is amplified and forwarded to the single channel analyzer (590A amplifier combined with timing single channel analyzer or SCA, Ortec). The output of the single channel analyzer is connected to a counter unit. It is of big importance to control the pulses' shape as well as the countrate. Any peculiar behavior may hint on possible vacuum leaks, sample absence etc. Furthermore, the signal from SCA is transmitted to a multichannel analyzer (MCA), which allows multichannel scaling data acquisition (MCA4, by FASTComTech). MCA additionally collects information from the transducer (the start and channel advance signals). After the evaluation process, the MCA transfers the signal to a PC. All data exchange with the PC is carried out via both RS-485 and USB connections (see Section 2.2). Currently, the PC is built on a quad-core i7 6700 3.40 Ghz CPU with 32 Gb of RAM. All modules (omitting MCA) are powered by NIM8301 by CAEN. The only exception is the pre-amplifier, which is powered by 4-channel N1471 (CAEN) power supply. The complete electronic rack is shown in Fig. 3.

The pumping unit is a separate part of eMIL. Due to safety reasons (i.e., one cannot open the target line when the pressure is insufficiently low) it has been automatized entirely. At the moment it consists of a

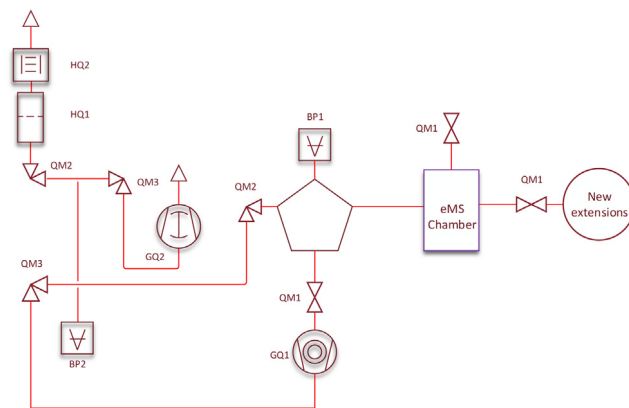


Fig. 4. Pumping unit scheme. Here QM stands for valve gates (QM1), electro-angular valves (QM2) and electro-pneumatic (QM3). GQ1 and GQ2 are turbomolecular and roots pumps, BP1 and BP2 are Pirani and Piezo gauges and HQ1-2 are particle and air filters, respectively.

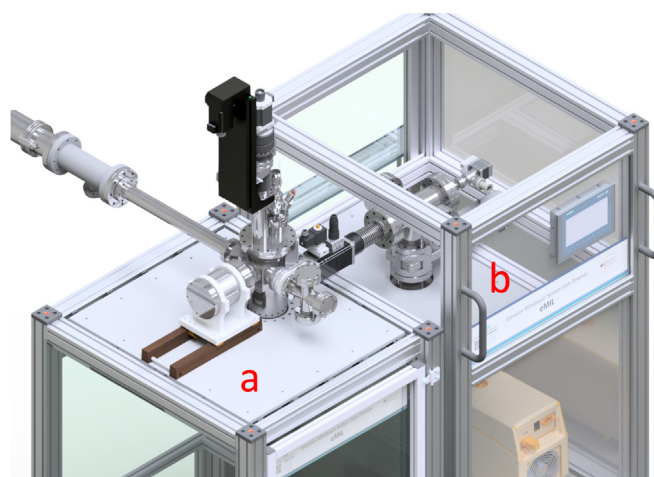


Fig. 5. The chamber unit (a), with the PPAC detector (on the rails) mounted on the transducer. The chamber itself is plugged to the pumping unit (b) via a soft bellow to minimize harmful vibrations originating from the pumps.

forepump (roots pump by Pfeiffer ACP15), turbo pump (Pfeiffer HiPace 300) and 5 valve gates and angle valves (by VAT), Pfeiffer RPT 200 and CPT 200 gauges all are controlled by RS-485 serial communication with Siemens Simatic 1212C (with SB1241RS485). Vacuum scheme is presented in Fig. 4.

This combination provides a robust base for establishing point-to-point communication via RS-485 among nodes. Besides a PC control, a human-machine-interface (Siemens Simatic KTP700) is placed on the side of the pumping unit. The units overview is shown in Fig. 5.

The remarkable flexibility of the setup is achieved by virtue of so-called lids (or sample-holders). Currently, there is solely one lid available. The lid is built on a CF100 flange with 2 CF40 ports. One CF40 flange is for the manipulation system (2 stepping motors), which allows samples to be moved along Z and Y axes (rotate and move up and down). While the second port is for the electrical feedthroughs and K-type thermocouples. The current design permits 4 samples to be measured without the need of breaking the vacuum in a temperature range from 300 to 1100 K. The temperature control is achieved with a special design and utilization of button-heaters. The heaters are made from pyrolytic boron nitride and pyrolytic graphite (by HeatWave Labs). An additional sample position is provided to a 5th sample, which can be a α -Fe foil. The optimal sample size is 1×1 cm, but additional installation of a collimator at the beam entering flange

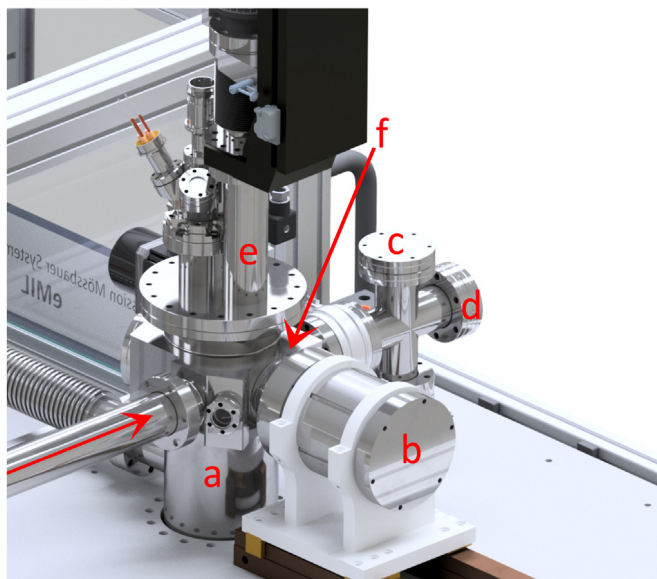


Fig. 6. The implantation chamber (a) with a hot lid mounted on top (e). On the left there is an incoming beam (marked by the arrow), next to it the quartz port, the Beryllium window (f) is placed behind the transducer (b). The transducer is equipped with a PPAC detector. On the far right is a distribution cross for chemical and magnetic extension chambers (d) and a Faraday cup mounted in a flange(c).

allows one to reduce the beam spot, and thus to implant in a small area. The rotation (Z-axis) can be employed for the determination of the angular dependence of the collected spectrum. In general, an on-line measurement (implanting and measuring) takes up to 4–5 min at $\theta_I = 30^\circ$ ($\theta_M = 60^\circ$). After this, the implantation is stopped and the sample is rotated in a way that it is placed parallel to the detector (see Fig. 6).

Last but not least, measurements as function of temperature allow the determination of the Debye–Waller factors and Debye temperatures.

At this point, the radiation during the decay is being collected for the same amount of time. Depending on the isotope the whole procedure can be repeated several times as a measure to improve the statistics. Further expansions (cryogenic, magnetic and chemical measurements) can be easily mounted on another lid, fed through the bottom flange or through the cross placed on the axis of the beam (see Figs. 2 and 5). The remote control is realized via an Arduino Uno Rev3 card, which is coupled with two of CNC single axis routers Toshiba TB6560AHQ (or 3A TB6560).

The chamber unit is based on an austenitic stainless steel (316LN ESR) tubular chamber ($\varnothing = 100$ mm). Use of this steel provides very low magnetic permeability and can withstand the harmful impact of corrosion. Everything else (CF and KF flanges, crosses) is made from 316L steel. Furthermore, when the implantation chamber was devised we had to take into account numerous factors. On the one hand, most of them were attempts to overcome the limitations one came across before, while flexibility and expanding of applications were prioritized. This resulted in a bigger chamber, with a Beryllium window which permits the usage of 2 PPAC detectors, pumping coming from the pumping station and an additional port for light experiments. The light port is inspired by the fact that one may see changes in the quadrupole splitting and isomer shift, whilst a semiconductor sample is being irradiated by a light source [22]. Additional information, including electric and vacuum schemes, drawings of the sample holder as well as a Faraday cup design, are presented in Ref. [23].

2.2. Software

In order to make the measuring experience reliable and productive two main software applications were devised and programmed. The

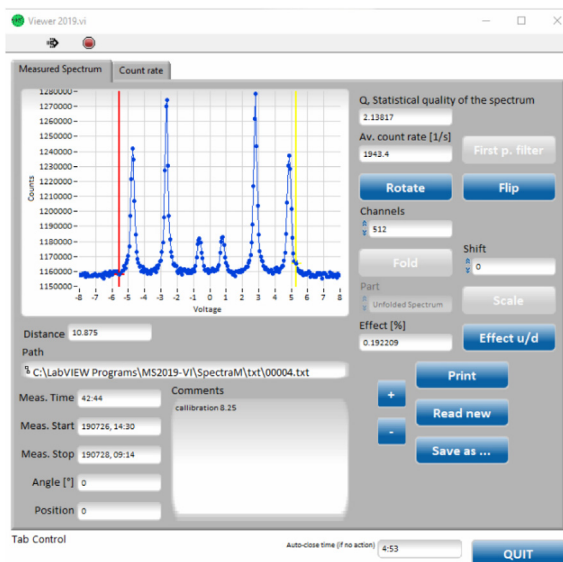


Fig. 7. eMIL's spectra viewer window, with statistical information.

first is a LabView-based data acquisition part, which is the core of eMIL and is responsible for data acquisition. In the current work we drew the main inspiration from software of a Mössbauer spectrometer built on National Instruments' solutions — Nuclear DSP System [24]. The second is a Siemens Simatic application written for a PLC-controller. Simatic can be operated remotely from a PC or directly via the human-machine-interface mounted on the pumping unit. A comprehensive description of the developed applications can be found in Ref. [23].

When a new sample is about to be measured a user can define basic parameters on the sample (velocity range, XY-relative position, temperature of the button-heaters etc.). Once acquisition has started, it can be paused and then continued later on. Additionally, software allows to directly print spectra, read previously saved measurements and convert the output data into various formats. The second tab in the Viewer Pop-Up shows a graph with the count rate over time. The preferable output format is either ASCII (.dat, txt) or binary (.ms0) which are commonly used at ISOLDE/CERN alongside with VINDA program [25] for further spectra evaluation. The viewer window is shown in Fig. 7.

A user can easily control parameters such as voltage, current and observe changes in the real time of the HV-supply and NIM crate. Several auxiliary functions were additionally implemented. With this, it becomes feasible to estimate spectral quality (Q-factor). The statistical quality of a measured Mössbauer spectrum as follows:

$$Q = \frac{\epsilon^2}{\epsilon + 2} N_\infty, \quad (1)$$

here N_∞ is the total number of the counts and ϵ stands for the resonant effect:

$$\epsilon = \frac{|N_\infty - N_0|}{N_\infty}, \quad (2)$$

where N_0 represents the number of counts at the maximum (minimum for transmission) of a resonance peak. In principle, the developed program has the flexibility to be used for Mössbauer spectroscopy in all transmission, emission or back-scattering configurations. The pumping unit is completely automatized by means of TIA Portal environment (Siemens).

3. Experimental procedure and results

In order to confirm the correct functionality in the emission geometry we have performed a couple of test measurements on both

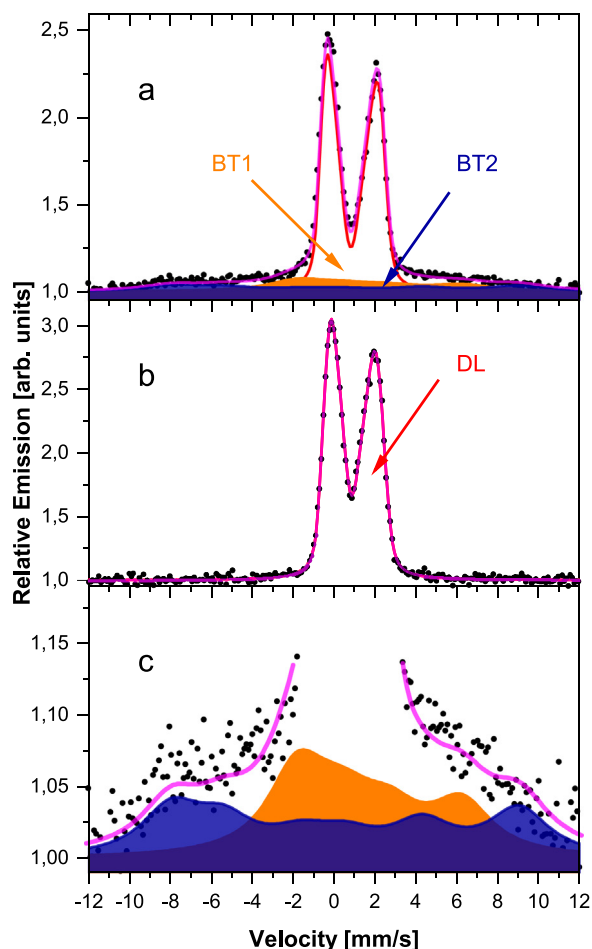


Fig. 8. ^{57}Fe emission Mössbauer spectra captured at room temperature: (a) pristine rutile; (b) rutile hydrogenated at 473 K; (c) magnified view of the pristine rutile spectrum showing the effect of the magnetic splitting.

hydrogenated and pristine TiO_2 (rutile) thin films (500 nm). The thin film were deposited onto Si substrates by radio frequency sputtering (LA 440S, VON ARDENNE), employing a TiO_2 target (99.9% FHR Anlagenbau, Germany). Prior to hydrogenation, the samples were annealed at 1173 K for 5 h to ensure their crystallization into rutile. Consequently, one sample was hydrogenated in inductively coupled hydrogen plasma (Plasma Lab 100 ICP-CVD, Oxford Instruments). The film was treated for 30 min at 2000 W at 473 K (sample temperature during the procedure).

The samples were implanted at ISOLDE/CERN after acceleration to 40 keV with ^{57}Mn ions. Emission spectrum was recorded during the implantation of a fluence 4×10^{11} ^{57}Mn ions/cm² for each sample. The energy of the incoming beam results in an implantation depth up to ~ 45 nm (average 29 nm, SRIM calculations). Velocity and isomer-shift (IS) values are given relative to the center of $\alpha\text{-Fe}$ spectrum at room temperature.

Fig. 8 shows spectra of hydrogenated TiO_2 (rutile) taken at room temperature (5.7 mm/s). Pristine sample resembles results captured earlier [26], where one may notice the presence of magnetic splitting and thus were evaluated similarly. The spectrum was analyzed in terms of three components by virtue of *VINDA* program [25].

The dominating double peak (DL) in the central part was analyzed with a quadrupole split component, which was simulated as a distribution of quadrupole splittings $P(\Delta E_Q)$ [27]. A coupling between the quadrupole splitting and the isomer-shift was assumed in the evaluation of the lineshape, ($\delta = \delta_0 + \delta_1 \cdot \Delta E_Q$) where δ_1 and δ_0 were fitting variables. Here, δ_0 was used as a free parameter, δ_1 as a global parameter

(same in all spectra of the temperature series) [27]. DL was assumed to be high-spin Fe^{2+} .

Derived average isomer-shift and quadrupole splitting are $\delta_{RT} = 0.88(2)$ and $\Delta E_Q = 2.16(3)$ mm/s. The rest two components (BT1 and BT2) were analyzed with the semi-empirical Blume–Tjøn model. On the basis of current results one may assume that the behavior matches with the earlier defects annealing studies [26], where the magnetic contribution was evaluated as a paramagnetic Fe^{3+} slowly relaxing with a temperature increase.

After hydrogenation at 473 K, the spectrum changes and demonstrates no Fe^{3+} contribution. It was sufficient to employ only one Fe^{2+} component as a distribution of quadrupole splittings to reasonably fit the spectrum. Obtained isomer-shift and quadrupole splitting are $\delta_{RT} = 0.92(2)$ and $\Delta E_Q = 1.92(3)$ mm/s, respectively. The reason responsible for the absence of Fe^{3+} is likely related to a hydrogen lattice incorporation, as have been seen in the case of variously hydrogenated anatase [28]. Probably, hydrogen diffuses easily into the bulk, since the atomic probes implanted up to 45 nm are affected, thus showing no Fe^{3+} component. It is noteworthy, that annealing behavior taking place is different in anatase and rutile [23].

4. Conclusions and outlook

This paper describes the emission Mössbauer Spectrometer setup, which provides a user with great versatility and reliability. eMIL has been designed for ISOLDE/CERN and is directly connected to the GLM implantation branch. The possibility to easily perform experiments has been demonstrated. At the moment the setup supports measurements from room up to elevated (1100 K) temperatures, with the possibility to perform angular-dependent measurements and experiments under applied light irradiation. Furthermore, eMIL provides a solid ground for additional modifications. One of the ways to improve data acquisition, for instance, is the installation of the second PPAC detector. Two detectors could significantly boost the data collection time. The bottom flange can serve as the base for a DC sputtering source, if samples are going to be treated by plasma.

Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

CRediT authorship contribution statement

Dmitry V. Zybkin: Data curation, Investigation, Formal analysis, Visualization, Writing - original draft. **Ulrich Vetter:** Conceptualization, Supervision. **Fredericus M.A. Linderhof:** Software, Writing - review & editing. **Haraldur P. Gunnlaugsson:** Validation, Writing - review & editing. **Peter Schaaf:** Funding acquisition, Resources, Supervision, Writing - review & editing.

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