

# Simultaneous analysis of elastic scattering and transfer/breakup channels for the ${}^6\text{He}+{}^{208}\text{Pb}$ reaction at energies near the Coulomb barrier

J. P. Fernández-García<sup>a,b</sup>, M. A. G. Alvarez<sup>a,b</sup>, A. M. Moro<sup>a</sup>, M. Rodríguez-Gallardo<sup>a,c</sup>

<sup>a</sup>Departamento de FAMN, Universidad de Sevilla, Apartado 1065, 41080 Sevilla, Spain.

<sup>b</sup>Centro Nacional de Aceleradores, Universidad de Sevilla, 41092 Sevilla, Spain.

<sup>c</sup>Instituto de Estructura de la Materia, CSIC, Serrano 123, 28006 Madrid, Spain.

---

## Abstract

The elastic and  $\alpha$ -production channels for the  ${}^6\text{He}+{}^{208}\text{Pb}$  reaction are investigated at energies around the Coulomb barrier ( $E_{\text{lab}}=14, 16, 18, 22,$  and  $27$  MeV). The effect of the two-neutron transfer channels on the elastic scattering has been studied within the Coupled-Reaction-Channels (CRC) method. We find that the explicit inclusion of these channels allows a simultaneous description of the elastic data and the inclusive  $\alpha$  cross sections at backward angles. Three-body Continuum-Discretized Coupled-Channels (CDCC) calculations are found to reproduce the elastic data, but not the transfer/breakup data. The trivially-equivalent local polarization potential (TELP) derived from the CRC and CDCC calculations are found to explain the features found in previous phenomenological optical model calculations for this system.

**Keywords:** Nuclear reaction  ${}^{208}\text{Pb}({}^6\text{He}, {}^6\text{He}), {}^{208}\text{Pb}({}^6\text{He}, {}^4\text{He})$ , halo nuclei, Continuum-Discretized Coupled-Channels and Coupled-Reaction-Channels calculations.

**PACS:** 25.60.Dz, 25.60.Gc, 25.60.Bx, 21.10.Gv, 27.20.+n.

---

## 1. Introduction

Since the early nineties, a considerable amount of experimental data of reactions induced by the Borromean nucleus  ${}^6\text{He}$  on a variety of targets has been accumulated. In the energy regime around the Coulomb barrier, and for medium-mass and heavy targets, these data show some common remarkable features. First, the elastic scattering angular distribution does not follow the expected Fresnel pattern, that characterizes the scattering of heavy ions at these energies. Instead, the characteristic Fresnel peak is partially or completely suppressed and the angular distribution shows a smooth decrease as a function of the scattering angle. Second, these reactions exhibit a large yield of  $\alpha$  particles [1–5]. This is clearly a consequence of the weak binding of the  ${}^6\text{He}$  nucleus, which can be easily broken up by the strong couplings induced by the target.

Heavy ion reactions between stable nuclei can be understood in terms of the *strong absorption* picture [6]. This implies the existence of an interaction distance ( $R_{sa}$ ) such that impact parameters below  $R_{sa}$  are dominated by strong absorbing processes (i.e. deep inelastic collisions, compound nucleus, etc), whereas for separations larger than  $R_{sa}$ , the nuclei barely interact, and elastic scattering dominates. Direct reactions, such as nucleon transfer or diffractive breakup, take place only for peripheral or grazing collisions. However, the analysis of recent experiments suggest that this picture may be inadequate

to interpret collisions involving weakly bound nuclei. In the  ${}^6\text{He}+{}^{208}\text{Pb}$  case, optical model (OM) calculations performed at Coulomb-barrier energies [7–9] show that, in order to reproduce the elastic data, one requires optical potentials with a very diffuse imaginary tail. If the interacting potential is parametrized using a standard complex Woods-Saxon form, the diffuseness parameter of the imaginary component required to reproduce the elastic data turns out to be around  $a_i \approx 2$  fm, a value that is considerably larger than the diffuseness derived from the matter distribution ( $\approx 0.6$  fm). This phenomenon, recently referred to as *long range absorption* effect, suggests the presence of reaction channels that remove flux from the elastic channel at distances well beyond the strong absorption radius, in contrast to the picture suggested by the strong absorption model. Continuum-Discretized Coupled-Channels (CDCC) calculations using either a simplified *di-neutron* model [10, 11] or a more realistic three-body model for the  ${}^6\text{He}$  nucleus [12–18] indicate that this long range absorption phenomenon can be explained in terms of the strong couplings to the breakup channels due mainly to the dipole Coulomb interaction.

Despite their success to describe the elastic data angular distributions, it was shown in [19] that these CDCC calculations failed to reproduce the energy and angular distributions of the  $\alpha$  particles observed at large angles in the same experiment. On the other hand, it was shown in the same work that these alpha particles distributions could be well accounted for by assuming a two-neutron transfer mechanism leading to very excited states of the target. These calculations were performed using the distorted-wave Born approximation (DWBA), assuming a

---

Email addresses: jpfernandez@us.es (J. P. Fernández-García), malvarez@us.es (M. A. G. Alvarez), moro@us.es (A. M. Moro), mrodri@us.es (M. Rodríguez-Gallardo)

two-neutron mechanism leading to both bound and unbound final states of the target. The relevance of the one- and two-neutron transfer channels has been evidenced in other reactions involving  ${}^6\text{He}$  [2–5, 20].

The main difference between the CDCC and DWBA calculations is due to the couplings included in each method. This is schematically illustrated in Fig. 1. The CDCC method assumes a *direct breakup* mechanism in which the breakup process is treated as an inelastic excitation of the projectile to its continuum spectrum (top panel). On the other hand, the DWBA method (or more elaborated forms of the transfer amplitude [21]) is based on a *transfer to the continuum* picture, which treats the  ${}^6\text{He}$  breakup assuming a neutron transfer mechanism populating  ${}^{210}\text{Pb}^*$  states, as shown in the bottom panel of Fig. 1. Since both the projectile and target states form a complete set, we expect both methods to give similar results provided that the underlying interactions are the same and the model space is sufficiently large to achieve convergence of the studied observables [21]. In this respect, the direct breakup and transfer to the continuum methods can be regarded as alternative methods to describe the projectile dissociation. Therefore, the failure of the CDCC approach to describe the  $\alpha$  yield could be attributed to the truncation of the model space required in practical calculations, and/or to the neglect of the transfer to the bound states of the target which, at these energies, may contribute significantly to the  $\alpha$  cross section.

The transfer to the continuum calculations presented in [19] are based on the DWBA approximation and describe the elastic channel by means of a phenomenological OM, with parameters adjusted to the elastic data. The features of this OM potential must reflect the effect of channel couplings on the elastic channel. Since the two-neutron transfer channels were found to be very important to describe the  $\alpha$  yields, it would be desirable to investigate to what extent the explicit inclusion of these transfer couplings can explain simultaneously the observed features of the elastic scattering and the  $\alpha$  channel, without the need of a phenomenological OMP adjusted to the data. To study the effect of the transfer channels on the elastic channel one needs to go beyond the Born approximation, that is, to use the Coupled-Reaction-Channels (CRC) method.

With this purpose in this work we present CRC and CDCC calculations for the  ${}^6\text{He}+{}^{208}\text{Pb}$  reaction at near-barrier energies, comparing the calculated elastic and  $\alpha$  cross sections with the available data. In the CRC calculations we rely on the scheme illustrated in the bottom panel of Fig. 1, but we allow back couplings from the transfer to the elastic channels. We show that, unlike the CDCC, the CRC formalism is able to explain simultaneously the long-range absorption effect on the elastic scattering and the large yield of  $\alpha$  particles for this reaction.

Although the calculations are restricted to the  ${}^6\text{He}+{}^{208}\text{Pb}$  case, we believe that the conclusions can be extrapolated to other reactions induced by weakly bound projectiles.

## 2. CDCC calculations

Within the Continuum-Discretized Coupled-Channels (CDCC) method, the dissociation of the projectile is treated

assuming a direct breakup picture, in which the projectile is excited to its unbound states (see Fig. 1). For  ${}^6\text{He}$  reactions, this has been done using a three-body reaction model (based on a simple di-neutron model of the projectile,  ${}^4\text{He}+2n$  [10, 11]) or a four-body reaction model (using a realistic three-body model of  ${}^6\text{He}={}^4\text{He}+n+n$  [12, 13, 16–18]). In this work, we rely on the first model because this will permit a more meaningful comparison with the CRC calculations presented in Section 3, which are also based on a two-body model of the  ${}^6\text{He}$  nucleus. In particular, we will make use of the improved di-neutron model proposed in [11, 19], in which the  ${}^4\text{He}+2n$  relative wavefunctions are calculated in a Woods-Saxon potential with radius  $R = 1.9$  fm and diffuseness parameter  $a_0 = 0.39$  fm. The ground state wavefunction is calculated assuming a  $2S$  configuration and an effective separation energy of the two-neutron cluster  $S_{2n} = 1.6$  MeV. For the  $\alpha+{}^{208}\text{Pb}$  interaction we took the potential of Barnett and Lilley [22], whereas the  $2n+{}^{208}\text{Pb}$  interaction was approximated by the deuteron- ${}^{208}\text{Pb}$  global potential of Ref. [23].  ${}^6\text{He}$  continuum states with relative angular momentum  $l_i = 0, 1, 2$  for the  ${}^4\text{He}+2n$  relative motion were considered. For the  $l_i = 2$  continuum, the potential depth was adjusted to reproduce the known  $2^+$  resonance at  $E_x = 1.8$  MeV above the g.s. For each value of  $l_i$ , the continuum was discretized using the standard binning method. The maximum excitation energy depended somewhat on the incident energy, ranging from 5 MeV (for  $E_{\text{lab}}=14$  MeV) to 8 MeV (for  $E_{\text{lab}}=27$  MeV). These calculations were performed with the code FRESKO [24].

The calculated elastic differential cross sections are compared in Fig. 2 (dashed lines) with the experimental data from Refs. [7–9]. Despite the simplified structure model used for the  ${}^6\text{He}$  nucleus and the absence of any free adjustable parameter, the overall agreement with the data is very good at the five considered energies. Below the barrier ( $E_{\text{lab}}=14, 16,$  and  $18$  MeV), the inclusion of the breakup couplings produces a decrease of the cross section with respect to Rutherford at c.m. angles beyond  $60^\circ$ . Above the barrier ( $E_{\text{lab}}=22$  and  $27$  MeV) the main effect is the absence of the Fresnel peak.

In Figs. 3 and 4, the calculated energy and angular distributions of the  $\alpha$  particles (dotted lines) are compared with the data from Ref. [19] (solid circles). Clearly, the calculations fail to describe both observables and, therefore, we conclude that the direct breakup model, at least within this restricted model space, cannot explain the large yield of  $\alpha$  particles produced at large angles. On the other hand, it was shown in Ref. [19] that these distributions could be well reproduced by means of DWBA calculations, assuming a two-neutron transfer mechanism.

## 3. CRC calculations

In the distorted-wave Born approximation (DWBA) calculations of Ref. [19], the entrance channel was described using a phenomenological OM potential, with parameters adjusted to reproduce the experimental elastic angular distribution. In the Coupled-Reaction-Channels (CRC) calculations, we aim at explaining the features of the elastic scattering as a consequence

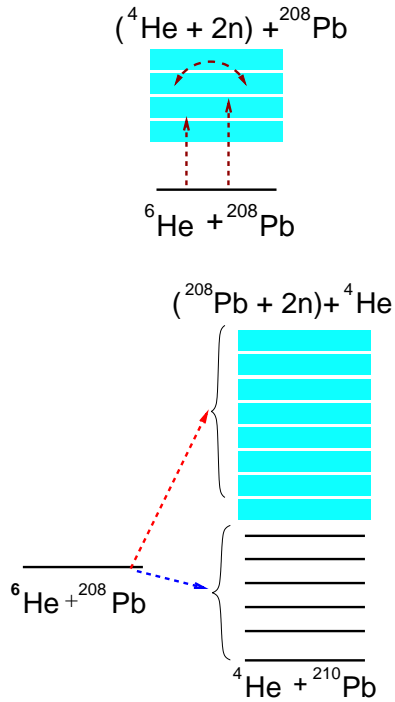


Figure 1: Schematic representation of the couplings included in the direct breakup (CDCC) (top) and transfer to the continuum (DWBA, CRC) (bottom) calculations.

of the same couplings which are responsible for the production of the  $\alpha$  particles with the objective of reproducing both channels within the same framework. Therefore, in order to avoid double counting, in the CRC calculations one has to use instead a *bare* interaction, that is, a potential that represents the interaction between the colliding nuclei in absence of couplings that will be explicitly included.

In particular, we use for the bare interaction the São Paulo potential (SPP) [25, 26]. This is a microscopic potential obtained by means of a double-folding procedure, using the matter densities of the colliding nuclei and an effective nucleon-nucleon interaction, multiplied by an energy-dependent term, which accounts for part of the non-locality of the optical potential. This provides a parameter-free prescription, which has been successfully applied to a large variety of systems [25–32]. Indeed, the choice of the bare interaction is not unique. The present choice is convenient because, being based on the matter densities of the colliding nuclei, reduces many ambiguities associated with more phenomenological potentials. The bare interaction includes also a short-range imaginary potential of Woods-Saxon shape and parameters  $W = 50$  MeV,  $R_0 = 1$  fm,  $a_0 = 0.1$  fm, to simulate the in-going boundary condition for fusion. This bare interaction was recently used in the OM analysis performed in [33] for the same elastic data. As it was shown in that work, the bare interaction alone is unable to reproduce the elastic data. This is illustrated in Fig. 2, where the one-channel calculation with the bare interaction (dotted line) is compared

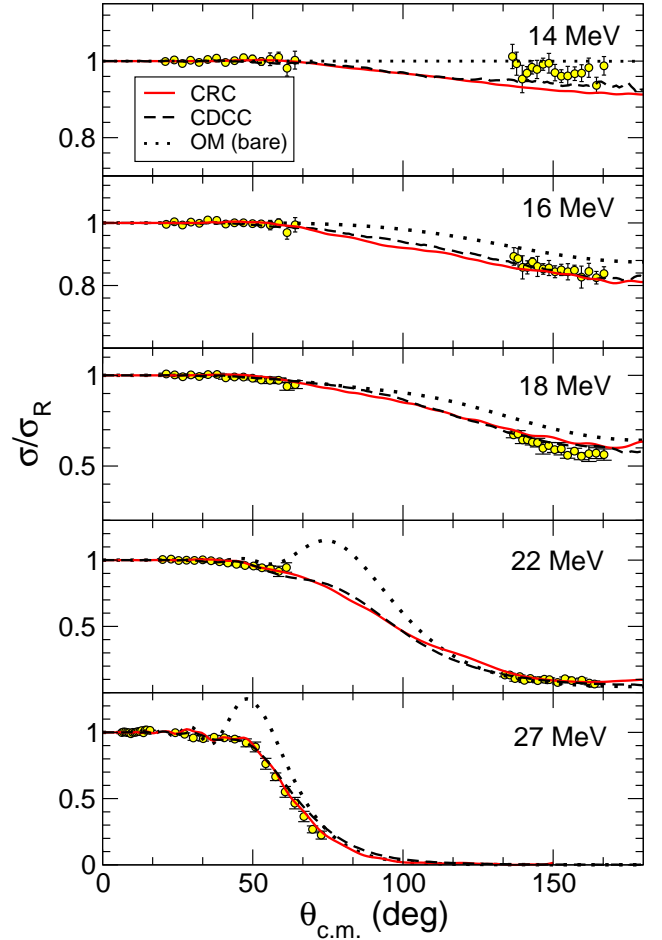


Figure 2: Elastic scattering angular distribution as a function of the c.m. scattering angle for the reaction  ${}^6\text{He}+{}^{208}\text{Pb}$  at  $E_{\text{lab}}=14, 16, 18, 22,$  and  $27$  MeV. The dotted line is the one-channel calculation with the bare interaction,  $U_{\text{SPP}}(r)$ , the dashed line is the three-body CDCC calculation and the solid line is the full CRC calculation. The circles are the experimental data from Ref. [7, 9].

with the experimental data. At  $E_{\text{lab}} = 22$  and  $27$  MeV (above the nominal Coulomb barrier) this calculation predicts a pronounced rainbow, which is absent from the data. Below the barrier ( $E_{\text{lab}} = 14, 16,$  and  $18$  MeV), the calculation overestimates the data for angles beyond  $50^\circ$ . Thus, this prescription fails to describe the elastic data both below and above the Coulomb barrier. The failure is clearly due to the influence of channel couplings, mainly, the  $2n$  removal channels, which are expected to be very important for a loosely bound nucleus like  ${}^6\text{He}$ . In Ref. [33], the effect of these channels on the elastic channel was taken into account adding a phenomenological component to the bare interaction of Woods-Saxon form. It was found that the real and imaginary parts of the Woods-Saxon potentials required a large diffuseness parameter in order to reproduce the data.

In the present work, we start from the same bare interaction but, instead of adding any phenomenological component, we introduce the two-neutron transfer couplings explicitly, using the CRC formalism. Using the prior representation, the transfer couplings involve a matrix element of the operator

$V_{[2n+Pb]} + U_{[\alpha+Pb]} - U_{\text{bare}}$ , where  $V_{[2n+Pb]}$  is a real interaction describing two neutron states of the  $2n+^{208}\text{Pb}$  system,  $U_{[\alpha+Pb]}$  is the optical potential describing the  $\alpha+^{208}\text{Pb}$  elastic scattering and  $U_{\text{bare}}$  is the bare potential defined above, given by the São Paulo prescription. In the CRC calculations we took the same  $\alpha+^{208}\text{Pb}$  and  $2n+^{208}\text{Pb}$  interactions used in the CDCC calculations of the preceding subsection, in order to have a meaningful comparison between both formalisms. Nevertheless, for the latter only the real part of the interaction was considered, in order to allow the inclusion of two-neutron bound states in  $^{210}\text{Pb}$ . For the  $^6\text{He}$  nucleus, we adopt the same di-neutron model used in the CDCC calculations. For the  $^{210}\text{Pb}^*$  final states, we considered several values for the angular momentum of the relative angular motion for  $^{208}\text{Pb}+2n$  ( $l_f$ ). For each value of  $l_f$ , the energy spectrum was described by a set of representative states, following the procedure described in [19]. These representative states include both bound and unbound states, with respect to the two-neutron breakup threshold. Above this threshold, the  $2n+^{208}\text{Pb}$  continuum was discretized using 1 MeV bins, up to a maximum relative energy of 5 MeV. For energies below the threshold, we considered 5 bound states, evenly spaced in steps of 2 MeV. The number of partial waves  $l_f$  was progressively increased until convergence of the observables was found, in the kinematic region where comparison with data is possible (see discussion below). The coupled equations were solved iteratively using the code FRESKO [24], until the desired degree of convergence of the elastic and transfer cross sections was achieved. We include also the non-orthogonality correction [6].

The elastic differential cross sections obtained with these CRC calculations are displayed in Fig. 2 with solid lines. For all the considered energies, the agreement with the CDCC calculation and hence with the data is significantly improved with respect to the single-channel calculation performed with the bare interaction alone. Therefore, the inclusion of the transfer channels produces essentially the same effect on the elastic cross section as the inclusion of the  $^6\text{He}$  continuum states in the CDCC calculations. This is an important result which supports the conclusion that both, CDCC and CRC methods, populate to a large extent the same final states, although these states are expressed in different basis representations.

Despite the good agreement between both methods, the rate of convergence is very different. The CRC calculations were performed with  $l_f = 0 - 8$  partial waves for the  $2n+^{208}\text{Pb}$  relative motion, whereas the three-body CDCC calculations required only  $l_i = 0, 1, 2$  partial waves for  $2n+^4\text{He}$  in order to achieve convergence of the elastic scattering [11]. This indicates that the elastic scattering is mostly affected by the coupling to continuum states with small energy and angular momentum between  $2n$  and  $\alpha$ , and hence a representation based on the continuum of the projectile is more efficient to describe this observable.

The calculated energy and angular distributions of the  $\alpha$  particles are compared with the data in Figs. 3 and 4, respectively. For comparison, the DWBA calculations of Ref. [19] are also included (dashed lines). We see that the CRC calculations, represented by solid lines, reproduce the data with a similar quality of the DWBA calculations, although the magnitude is some-

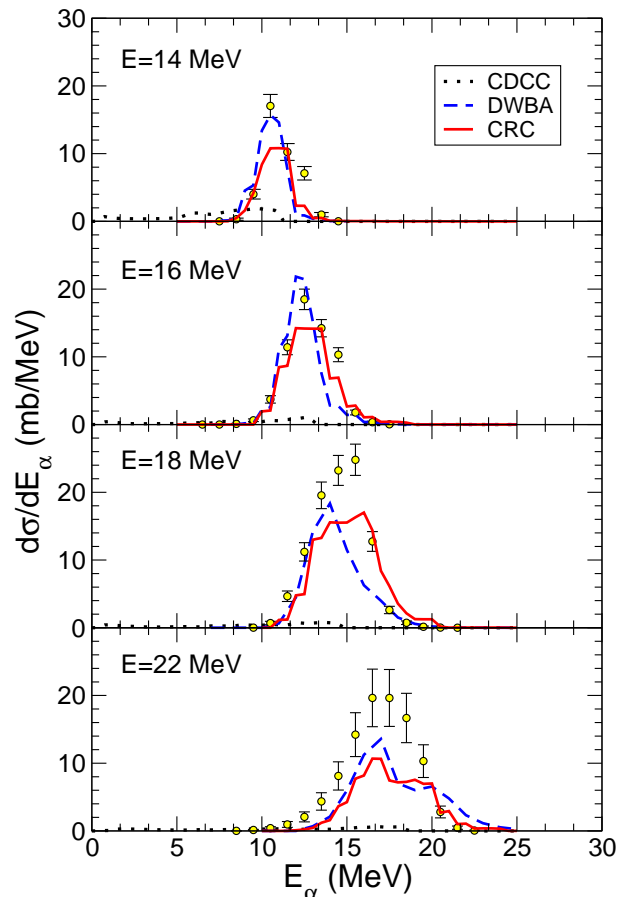


Figure 3: Energy distribution of the  $\alpha$  particles produced in the reaction  $^6\text{He}+^{208}\text{Pb}$  at  $E_{\text{lab}}=14, 16, 18,$  and  $22$  MeV, integrated in the angular interval  $\theta_{\text{lab}} = 132^\circ - 164^\circ$ . The dotted and solid lines are the CDCC and CRC calculations. The dashed lines and circles are, respectively, the DWBA calculations and the experimental data from Ref. [19].

what underestimated. This discrepancy might be due to the limitations of our di-neutron model (for both the projectile and target states), to the choice of the underlying interactions, or to the contribution of other channels. Considering the simplicity of our model, we can say that the overall agreement is fairly good. Therefore, unlike DWBA calculations, the CRC calculations are able of reproducing simultaneously the elastic and  $\alpha$  production data without requiring any phenomenological fit of the elastic cross section.

For a meaningful comparison with the data, we have verified that the number of partial waves used for the  $2n+^{208}\text{Pb}$  motion ( $l_f = 0 - 8$ ) provides also convergence of the  $\alpha$  cross section in the angular range where data exist. This is shown in Fig. 5, where we plot the contribution of each  $l_f$  to the angle-integrated transfer cross section for  $E_{\text{lab}}=22$  MeV. The dark histograms correspond to the angular range  $\theta_{\text{lab}} = 132^\circ - 164^\circ$ , whereas the dashed bars are for the full angular range. It is clear from this figure that the model space used in the CRC calculations yields convergence within the angular range of interest and hence we conclude that this model space is suitable to describe the  $\alpha$  particles at backward angles. At forward angles, the histograms

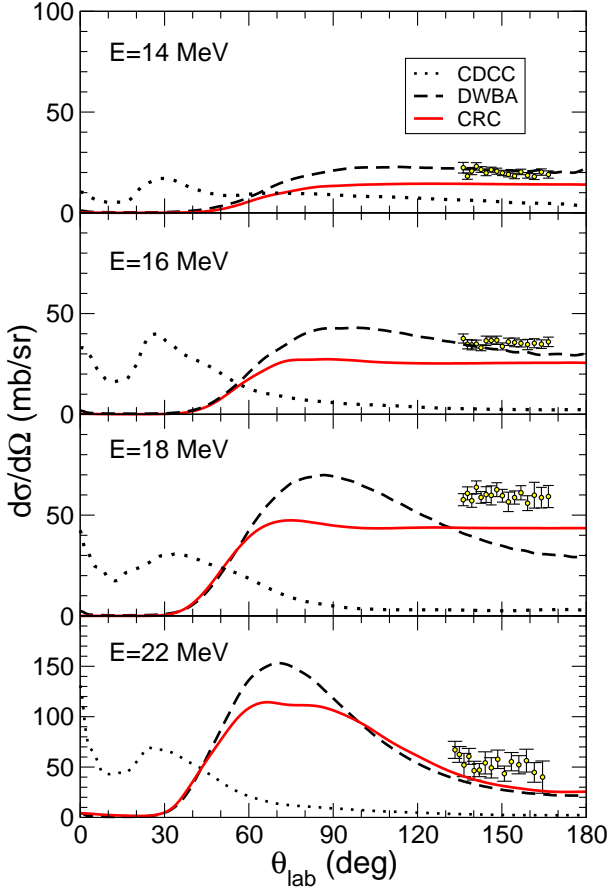


Figure 4: Angular distribution of the  $\alpha$  particles produced in the reaction  ${}^6\text{He}+{}^{208}\text{Pb}$  at  $E_{\text{lab}}=14, 16, 18,$  and  $22$  MeV. The dotted and solid lines are the CDCC and CRC calculations. The dashed lines and circles are, respectively, the DWBA calculations and the experimental data from Ref. [19].

of Fig. 5 suggest a non-negligible contributions from  $l_f > 8$  waves. Although there are no data for the  $\alpha$  particles emitted at these forward angles, it is expected that these are more easily described by the CDCC calculations.

#### 4. Trivially-equivalent local polarization potentials

The analysis of the preceding section demonstrates that the two-neutron stripping channel can explain the long-range effect found in previous phenomenological analyses of the elastic data [7–9]. To corroborate this conclusion in a more quantitative way, we have evaluated the so called *trivially-equivalent local polarization potential* (TELP). This is a local and  $L$ -independent potential constructed from the solution of the coupled equations in such a way that it simulates the net effect of the couplings on the elastic scattering. Here, the TELP is calculated following the prescription proposed by Thompson *et al.* [34], which involves two steps. First, for each total angular momentum, a local polarization potential is calculated from the source term of the elastic channel equation. Then, a  $L$ -independent potential is constructed by averaging these  $L$ -dependent polarization potentials, weighted with the trans-

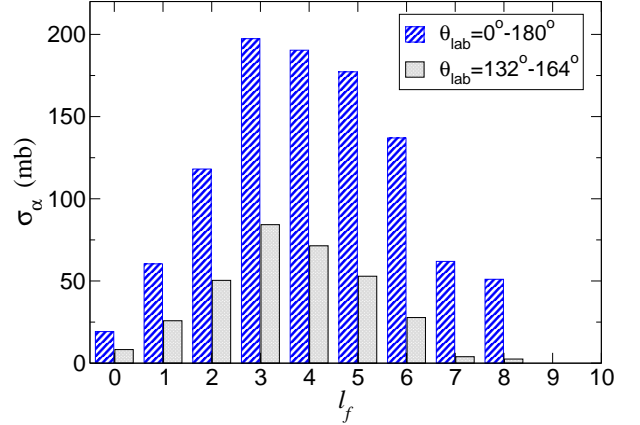


Figure 5: Two-neutron transfer cross section as a function of the  $2n-{}^{208}\text{Pb}$  orbital angular momentum  $l_f$ , for the CRC calculations performed at an incident energy of  $22$  MeV. The dashed bars and solid bars correspond, respectively, to the full angular range and the integrated angular distribution for  $\theta_{\text{lab}} = 132^\circ - 164^\circ$ .

fer/breakup cross section for each angular momentum. The TELP obtained by this procedure should be regarded as a simplified local approximation of a complicated coupled-channels system. If the TELP extracted from the coupled-channels calculation is a good representation of the overall effect of the couplings, the solution of the single-channel Schrödinger equation with the effective potential  $U_{\text{eff}} = U_{\text{bare}} + U_{\text{TELP}}$  should reproduce the elastic scattering similar to the one obtained with the full coupled-channels calculation.

In Fig. 6 we show the effective potentials ( $U_{\text{eff}}$ ) extracted from the CDCC and CRC calculations at  $E_{\text{lab}} = 22$  MeV. The top and bottom panels correspond to the real and imaginary parts, respectively. The arrows indicate the radius of sensitivity of the real and imaginary parts of the optical potential, according to the OM analysis performed in [9].

The dot-dashed and solid lines correspond, respectively, to the effective potential  $U_{\text{eff}}$  extracted from the CDCC and CRC calculations. For comparison, we have included also the double-folding SPP potential (dotted line in top panel), which corresponds to the bare interaction in the CRC calculations, and the phenomenological optical potential extracted in Ref. [9] from the fit of the elastic data (dashed line).

For distances  $r \gtrsim 15$  fm, the real part of the effective potential (in both the CRC and CDCC calculations) is dominated by a long-range attractive tail. This attractive component is known to arise from the dipole Coulomb interaction [33, 35, 36]. Although the direct breakup model provides a more natural representation for this effect, it is noticeable that the TELP extracted from the CRC shows up also this behaviour.

At distances around the strong absorption radius the TELP becomes repulsive, making the effective potential shallower than the bare interaction. This repulsive component is mainly due to nuclear couplings [33]. The OM potential is also less intense than the bare interaction, but it does not exhibit the long-range tail observed in the TELPs. This might be due to the fact that the real part of the optical potential is mostly sensitive to

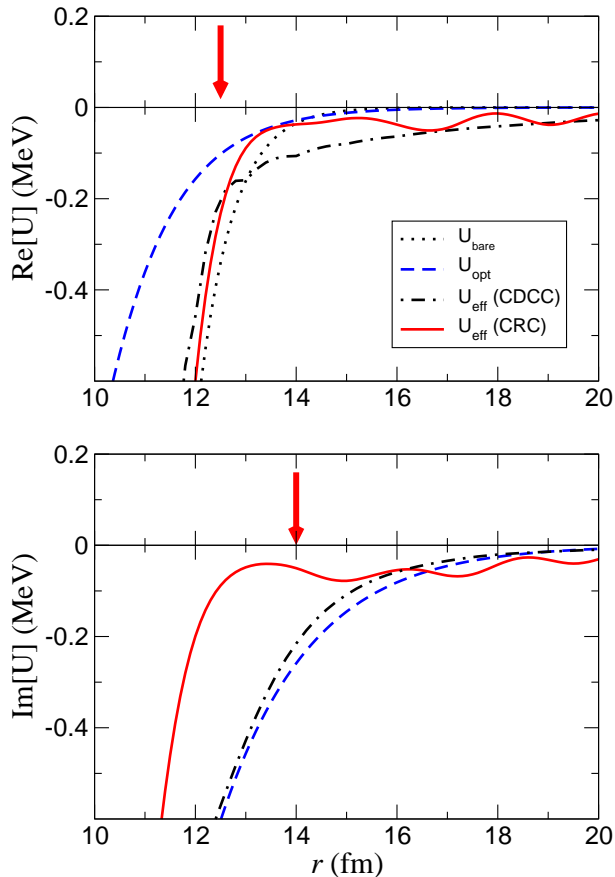


Figure 6: Effective potential,  $U_{\text{eff}}=U_{\text{bare}} + U_{\text{TELP}}$ , for  ${}^6\text{He}+{}^{208}\text{Pb}$  at 22 MeV, extracted from the CDCC (dot-dashed lines) and CRC calculations (solid lines). The dashed line is the phenomenological optical model potential extracted in Ref. [9] from the fit of the elastic data. The dotted line is the bare interaction used in the CRC calculations. The top and bottom panels correspond to the real and imaginary parts, respectively. The inset in the top figure compares the polarization potential  $U_{\text{TELP}}$  extracted from the CDCC and CRC calculations with the bare interaction.

distances around the strong absorption radius.

In the bottom panel of Fig. 6 we see that the imaginary part of the effective potential extracted from both, the CDCC and CRC calculations, is absorptive and exhibits also a diffuse tail. This behaviour is also observed in the optical potential (dashed line) and is related to the long-range absorption effect discussed in previous optical model analyses of this reaction [7–9]. Therefore, another important conclusion of this work is that the CDCC and CRC calculations provide a microscopic interpretation of this long-range absorption effect in terms of channel couplings. Quantitatively, the effective potential extracted from the CDCC calculations is closer to the phenomenological optical potential. We emphasize, however, that due to the  $L$  average, the TELP is not strictly equivalent to the full set of equations, and hence only the qualitative behavior is meaningful. In addition, the transfer couplings are intrinsically non-local and thus their approximation by a local object has to be interpreted with caution.

## 5. Summary and conclusions

We have studied the elastic and  $\alpha$ -production channel for the  ${}^6\text{He}+{}^{208}\text{Pb}$  reaction at energies around the Coulomb barrier. The available experimental data have been compared with three-body Continuum-Discretized Coupled-Channels (CDCC) and Coupled-Reaction-Channels (CRC) calculations using in both cases a di-neutron model of the  ${}^6\text{He}$  nucleus.

In the CDCC calculations, the projectile dissociation is taken into account by including the coupling with the  ${}^4\text{He}+2\text{n}$  continuum states. The inclusion of these couplings have a significant effect on the calculated elastic cross section, and provides a good description of the experimental elastic data. However, the calculated  $\alpha$  cross sections largely underestimate the experimental data, in agreement with the findings of Ref. [19].

In the CRC calculations, we consider the coupling to  $2\text{n}$  transfer channels and its influence on the elastic scattering cross section. For the bare interaction, we have used the microscopic double-folding São Paulo potential, supplemented with a short-range Woods-Saxon imaginary potential to account for complete fusion. The inclusion of two-neutron transfer channels produces a strong effect on the elastic cross section, providing a very good agreement with the experimental data. Moreover, these calculations explain also reasonably well the energy and angular distribution of the  $\alpha$  particles measured at backward angles. In particular, the shape of the energy distribution is very well accounted for, although the magnitude of the cross section is somewhat underestimated. This discrepancy might be due to the limitations of our simple di-neutron model or the choice of the interactions. In addition, besides the two-neutron transfer channels, other channels could contribute to this reaction. For example, in Ref. [37] the authors studied the effect of the one-neutron stripping channel, ( ${}^6\text{He}, {}^5\text{He}$ ), on the elastic and fusion cross section for several targets. In the  ${}^6\text{He}+{}^{208}\text{Pb}$  case, they found that the inclusion of these channels produce a reduction of the elastic cross section in the region of the Coulomb rainbow, and an increase at backward angles. The simultaneous inclusion of both mechanisms would be of interest.

The trivially-equivalent local polarization potentials (TELP) derived from the CRC and CDCC calculations exhibit the characteristic long-range real and absorptive parts, in agreement with the behavior observed in the phenomenological optical potentials extracted in previous OM analyses of the same data.

## Acknowledgements

This work has been supported by the Spanish Ministerio de Ciencia e Innovación under projects FPA2009-08848, FPA2009-07653, by the local government of Junta de Andalucía under project P07-FQM-02894 and the Spanish Consolider-Ingenio 2010 Programme CPAN (CSD2007-00042).

## References

## References

- [1] E. F. Aguilera, et al., Phys. Rev. Lett. 84 (2000) 5058.

- [2] P. A. DeYoung, et al., Phys. Rev. C 71 (2005) 051601.
- [3] A. Di Pietro, et al., Phys. Rev. C 69 (2004) 044613.
- [4] A. Navin, et al., Phys. Rev. C 70 (2004) 044601.
- [5] R. Raabe, et al., Nature 431 (2004) 823.
- [6] G. R. Satchler, Direct Nuclear Reactions, Oxford University Press, New York, 1983.
- [7] O. R. Kakuee, et al., Nucl. Phys. A728 (2003) 339.
- [8] O. R. Kakuee, et al., Nucl. Phys. A765 (2006) 294.
- [9] A. M. Sánchez-Benítez and others, Nucl. Phys. A803 (2008) 30.
- [10] K. Rusek, I. Martel, J. Gómez-Camacho, A. M. Moro, R. Raabe, Phys. Rev. C 72 (2005) 037603.
- [11] A. M. Moro, K. Rusek, J. M. Arias, J. Gómez-Camacho, M. Rodríguez-Gallardo, Phys. Rev. C 75 (2007) 064607.
- [12] T. Matsumoto, E. Hiyama, K. Ogata, Y. Iseri, M. Kamimura, S. Chiba, M. Yahiro, Phys. Rev. C 70 (2004) 061601(R).
- [13] T. Matsumoto, T. Egami, K. Ogata, Y. Iseri, M. Kamimura, M. Yahiro, Phys. Rev. C 73 (2006) 051602(R).
- [14] A. M. Moro, M. Rodríguez-Gallardo, R. Crespo, I. J. Thompson, Phys. Rev. C 75 (2007) 017603.
- [15] R. Crespo, A. M. Moro, I. J. Thompson, M. Rodríguez-Gallardo, J. Gómez-Camacho, J. M. Arias, Eur. Phys. J. S.T. 150 (2007) 13.
- [16] M. Rodríguez-Gallardo, J. M. Arias, J. Gómez-Camacho, R. C. Johnson, A. M. Moro, I. J. Thompson, J. A. Tostevin, Eur. Phys. J. S.T. 150 (2007) 51.
- [17] M. Rodríguez-Gallardo, J. M. Arias, J. Gómez-Camacho, R. C. Johnson, A. M. Moro, I. J. Thompson, J. A. Tostevin, Phys. Rev. C 77 (2008) 064609.
- [18] M. Rodríguez-Gallardo, J. M. Arias, J. Gómez-Camacho, A. M. Moro, I. J. Thompson, J. A. Tostevin, Phys. Rev. C 80 (2009) 051601(R).
- [19] D. Escrib, et al., Nucl. Phys. A792 (2007) 2.
- [20] E. F. Aguilera, et al., Phys. Rev. C 63 (2001) 061603.
- [21] A. M. Moro, F. M. Nunes, Nucl. Phys. A767 (2006) 138.
- [22] A. R. Barnett, J. S. Lilley, Phys. Rev. C 9 (1974) 2010.
- [23] W. W. Daehnick, J. D. Childs, Z. Vrcelj, Phys. Rev. C 21 (1980) 2253.
- [24] I. J. Thompson, Comp. Phys. Rep. 7 (1988) 167.
- [25] L. C. Chamon, et al., Phys. Rev. C 66 (2002) 014610.
- [26] L. C. Chamon, D. Pereira, M. S. Hussein, M. Cândido-Ribeiro, D. Galetti, Phys. Rev. Lett. 79 (1997) 5218.
- [27] L. C. Chamon, D. Pereira, M. S. Hussein, Phys. Rev. C 58 (1998) 576.
- [28] D. Galetti, S. S. Mizrahi, L. C. Chamon, D. Pereira, M. S. Hussein, M. A. Cândido Ribeiro, Phys. Rev. C 58 (1998) 1627.
- [29] M. A. G. Alvarez, L. C. Chamon, M. S. Hussein, D. Pereira, L. R. Gasques, E. S. Rossi, E. S. Rossi, C. P. Silva, Nucl. Phys. A723 (2003) 93.
- [30] M. A. G. Alvarez, L. C. Chamon, D. Pereira, E. S. Rossi, E. S. Rossi, C. P. Silva, L. R. Gasques, H. Dias, M. O. Roos, Nucl. Phys. A656 (1999) 187.
- [31] C. P. Silva, et al., Nucl. Phys. A679 (2001) 287.
- [32] E. S. Rossi, et al., Nucl. Phys. A707 (2002) 325.
- [33] J. P. Fernández-García, M. Rodríguez-Gallardo, M. A. G. Alvarez, A. M. Moro, Nucl. Phys. A840 (2010) 19–38.
- [34] I. J. Thompson, M. A. Nagarajan, J. S. Lilley, M. J. Smithson, Nucl. Phys. A505 (1989) 84.
- [35] K. Rusek, Eur. Phys. J. A 41 (2009) 399.
- [36] R. S. Mackintosh, N. Keeley, Phys. Rev. C 79 (2009) 014611.
- [37] N. Keeley, N. Alamanos, Phys. Rev. C 77 (2008) 054602.