


# Photoelectric fields in doped lithium niobate crystals

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## ABSTRACT

Photoinduced light scattering (PILS) in nominally pure stoichiometric and congruent lithium niobate single crystals ( $\text{LiNbO}_3$ ), and ones doped with  $\text{B}^{3+}$ ,  $\text{Cu}^{2+}$ ,  $\text{Zn}^{2+}$ ,  $\text{Mg}^{2+}$ ,  $\text{Gd}^{3+}$ ,  $\text{Y}^{3+}$ ,  $\text{Er}^{3+}$  cations was studied. All crystals have a relatively low effect of photorefraction and are promising materials for frequency conversion, electro-optical modulators and shutters. It was found that the photovoltaic and diffusion fields for some crystals have a maximum at a wavelength of 514.5 nm. All the crystals studied are characterized by a maximum of the integral intensity of the speckle structure of the PILS at a wavelength of 514.5 nm.

## ARTICLE HISTORY

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## KEYWORDS

Single crystal of lithium niobate; photorefractive effect; Rayleigh photoinduced light scattering; photovoltaic and diffusion fields

## 1. Introduction

The ferroelectric photorefractive lithium niobate crystal ( $\text{LiNbO}_3$ ) is characterized by a high value of spontaneous polarization, large values of the photoinduced electric fields intensity [1, 2]. The control of the photorefraction (optical damage) magnitude is usually carried out by doping of the crystal with cations of various metals [1, 2]. Photorefractive effect causes Rayleigh photoinduced light scattering, which occurs on static and dynamic (fluctuating) defects with an altered refractive index induced by laser radiation [3, 4]. In this case, the value of the electro-optical effect determines the value of the angle  $\theta$  of the PILS indicatrix opening in the crystal [4].

In this paper, the angular distribution of the intensity of the speckle structure of the PILS as a function of the wavelength of the exciting laser radiation was studied. Experiments were performed in nominally pure stoichiometric ( $\text{LiNbO}_{3\text{stoich}}$ ) and congruent ( $\text{LiNbO}_{3\text{cong}}$ ) lithium niobate single crystals, and ones doped with  $\text{B}^{3+}$ ,  $\text{Cu}^{2+}$ ,  $\text{Zn}^{2+}$ ,  $\text{Mg}^{2+}$ ,  $\text{Gd}^{3+}$ ,  $\text{Y}^{3+}$ ,  $\text{Er}^{3+}$  cations. The following laser lines were used in PILS experiment: 476.5, 488.0, 514.5 and 532.0 nm. Quantitative estimation of the photovoltaic ( $E_{pv}$ ) and diffusion ( $E_D$ ) fields values were made using the approach described in [3].

## 2. Experiment setup

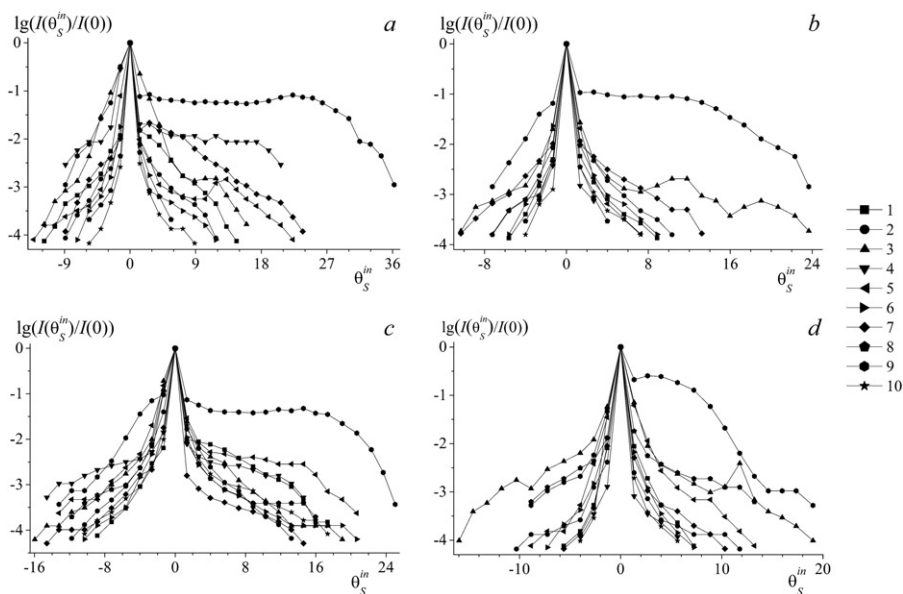
$\text{LiNbO}_3$  crystals were grown from the congruent melt at the “Crystal-2” installation by the Czochralski technique in air [5].  $\text{LiNbO}_{3\text{stoich}}$  crystal was grown from the melt with

58.6 mol. % of  $\text{Li}_2\text{O}$ . PILS registration was carried out using an installation, described in details in [4]. For the PILS registration the following lines of Spectra Physics (2018-RM) argon-krypton laser were used: 476.5 nm ( $P=216$  mW), 488.0 nm ( $P=98$  mW), 514.5 nm ( $P=282$  mW) and 532.0 nm ( $P=160$  mW). The value of the intensity of the photovoltaic and diffusion fields in crystals was calculated in the Mathcad 15.0 program using the approach proposed in [3]. The refractive indices of the extraordinary and ordinary rays were determined from empirical equations [6].

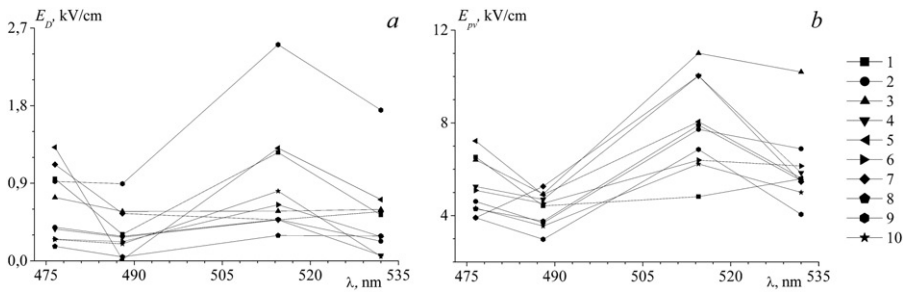
### 3. Results and discussion

The speckle structure of the PILS pattern of  $\text{LiNbO}_3$  crystals is determined by the features of the secondary structure of the crystal, which depends significantly on the composition and growing technology [1, 2, 4]. The birefringence of the crystal, both intrinsic and induced by laser radiation, is also important. At a power of excitation laser radiation of 160 mW, the indicatrix of the speckle structure of the PILS in  $\text{LiNbO}_3$ ,  $\text{LiNbO}_3$ : Zn (2.93),  $\text{LiNbO}_3$ : Gd (0.002): Mg (0.04),  $\text{LiNbO}_3$ : Er (3.1 wt. %) crystals is not revealed, and only circular scattering on crystal lattice defects is observed. For all other crystals investigated the indicatrix of the PILS is asymmetric with the form of a figure eight or an ellipse stretched along the polar axis.

It should be noted that for the  $\text{LiNbO}_3$ : Y (0.46 wt. %) crystal at  $P=160$  mW, the indicatrix of the speckle structure of the PILS is revealed very rapidly, in a time of about 1 s. For all other crystals, the opening time of the speckle structure of the PILS is about 60 s.



**Figure 1.** The angular distribution of the scattered light intensity at  $\lambda = 476.5$  (a), 488.0 (b), 514.5 (c), 532 (d) nm for the following crystals:  $\text{LiNbO}_3$ :Zn (0.018) (1);  $\text{LiNbO}_3$ :Zn (2.93) (2);  $\text{LiNbO}_3$ :Y (0.46) (3);  $\text{LiNbO}_3$ :Cu (0.007):Gd (0.02) (4);  $\text{LiNbO}_3$ :Gd (0.05) (5);  $\text{LiNbO}_3$ :Gd (0.002):Mg (0.4) (6);  $\text{LiNbO}_3$ :B (0.08 in the reacted mixture) (7);  $\text{LiNbO}_3$ :Er (3.1 wt. %) (8);  $\text{LiNbO}_{3\text{stoich}}$  (9);  $\text{LiNbO}_{3\text{cong}}$  (10).



**Figure 2.** Dependence of  $E_D$  (a) and  $E_{pv}$  (b) on laser line wavelength for  $\text{LiNbO}_3$  crystals with various composition:  $\text{LiNbO}_3:\text{Zn}$  (0.018) (1);  $\text{LiNbO}_3:\text{Zn}$  (2.93) (2);  $\text{LiNbO}_3:\text{Y}$  (0.46) (3);  $\text{LiNbO}_3:\text{Cu}$  (0.007):Gd (0.02) (4);  $\text{LiNbO}_3:\text{Gd}$  (0.05) (5);  $\text{LiNbO}_3:\text{Gd}$  (0.002):Mg (0.4) (6);  $\text{LiNbO}_3:\text{B}$  (0.08 in the reacted mixture) (7);  $\text{LiNbO}_3:\text{Er}$  (3.1 wt. %) (8);  $\text{LiNbO}_{3\text{stoich}}$  (9);  $\text{LiNbO}_{3\text{cong}}$  (10).

It can be seen from Figure 1 that crystals  $\text{LiNbO}_{3\text{stoich}}$ ,  $\text{LiNbO}_3:\text{Gd}$  (0.05) and  $\text{LiNbO}_3:\text{Y}$  (0.46 wt. %) possess the greatest asymmetry and the scattered radiation angle  $\theta$  in the series of crystals studied, regardless of the wavelength of the exciting line. At the same time, for a  $\text{LiNbO}_{3\text{stoich}}$  crystal, the shape of the scattering curve when excited by laser lines 476.5, 488.0, and 532.0 nm is approximately the same, but differs significantly from them when excited by a 514.5 nm laser line, Figure 1.

Figure 2 shows the dependences of the  $E_{pv}$  and  $E_D$  intensities in the investigated crystals on the wavelength of the exciting radiation. For  $\text{LiNbO}_3:\text{Zn}$  (0.018),  $\text{LiNbO}_3:\text{Zn}$  (2.93),  $\text{LiNbO}_3:\text{Gd}$  (0.05 wt. %) crystals, a maximum in the  $E_D$  dependence at the length of the exciting laser line of 514.5 nm is observed. However, the maximum in the  $E_D$  dependence are not observed for  $\text{LiNbO}_3:\text{Er}$  (3.1),  $\text{LiNbO}_3:\text{B}$  (0.08),  $\text{LiNbO}_3:\text{Y}$  (0.46 wt. %) crystals.

For  $\text{LiNbO}_3:\text{Y}$  (0.46),  $\text{LiNbO}_3:\text{Cu}$  (0.007):Gd (0.02),  $\text{LiNbO}_3:\text{B}$  (0.08 wt. %),  $\text{LiNbO}_3:\text{Gd}$  (0.05),  $\text{LiNbO}_3:\text{Zn}$  (2.93 wt. %),  $\text{LiNbO}_{3\text{stoich}}$  crystals the maximum in the  $E_{pv}$  dependence is also observed at the length of the exciting laser line of 514.5 nm. But at the same time, the maximum is absent for  $\text{LiNbO}_3:\text{Zn}$  (0.018),  $\text{LiNbO}_3:\text{Gd}$  (0.02):Mg (0.4) crystals. It is also seen from Figure 2 that the  $\text{LiNbO}_{3\text{stoich}}$  crystal at wavelengths of the exciting radiation of 476.5, 488.0 and 532.0 nm and  $\text{LiNbO}_3:\text{Zn}$  (0.018 wt. %) at 514.5 nm possess the smallest value of the  $E_{pv}$ .

#### 4. Summary

Photoinduced light scattering in nominally pure  $\text{LiNbO}_{3\text{stoich}}$  and  $\text{LiNbO}_{3\text{cong}}$  single crystals ( $\text{LiNbO}_3$ ), and ones doped with  $\text{B}^{3+}$ ,  $\text{Cu}^{2+}$ ,  $\text{Zn}^{2+}$ ,  $\text{Mg}^{2+}$ ,  $\text{Gd}^{3+}$ ,  $\text{Y}^{3+}$ ,  $\text{Er}^{3+}$  cations was studied. According to the characteristics of the PILS, a quantitative estimation of the intensity of the photovoltaic and diffusion fields was made. It was found that the  $E_{pv}$  and  $E_D$  for some crystals have a maximum at a wavelength of 514.5 nm. However, for the  $\text{LiNbO}_3:\text{Y}$  crystal (0.46 wt. %) there is no maximum in the  $E_{pv}$  dependence, but one at a wavelength of 514.5 nm in  $E_D$  dependence.

#### References

1. N. V. Sidorov, T. R. Volk, B. N. Mavrin, and V. T. Kalinnikov, *Lithium Niobate: Defects, Photorefraction, Vibrational Spectrum and Polaritons* (Nauka, Moscow, Russia 2003).

2. T. Volk, and M. Wohlecke, *Lithium Niobate: Defects, Photorefractive and Ferroelectric Switching* (Springer, Berlin, Germany 2008).
3. M. Gouklov, M. Imlau, and T. Woike, Photorefractive parameters of lithium niobate crystals from photoinduced light scattering, *Phys. Rev. B.* 77, 235110-1 (2008).
4. V. A. Maksimenko, A. V. Syuy, and Y. M. Karpets, *Photoinduced Processes in Lithium Niobate Crystals* (Physmatlit, Moscow, Russia 2008).
5. M. N. Palatnikov *et al.*, Growth and concentration dependencies of rare-earth doped lithium niobate single crystals, *J. Cryst. Growth.* 291(2), 390 (2006). DOI: [10.1016/j.jcrysgro.2006.03.022](https://doi.org/10.1016/j.jcrysgro.2006.03.022).
6. G. G. Gurzadyan, V. G. Dmitriev, and D. N. Nikogosyan, *Nonlinear Optical Crystals: Properties and Applications in Quantum Electronics* (Radio and Communication, Moscow, Russia 1991).