

## Comment on “Phase Contribution of Image Potential on Empty Quantum Well States in Pb Islands on the Cu(111) Surface”

The Letter by Yang *et al.* [1] presents an experimental scanning tunneling spectroscopy (STS) study of unoccupied quantum well states (QWSs) in Pb islands grown on the Cu(111) surface. The departure from particle-in-a-box energy dispersion has been observed with decreasing energy spacing between QWSs for the energies above 3 eV with respect to the Fermi level ( $E_F$ ). This is attributed to the image potential felt by an electron at the Pb-vacuum interface. It is proposed that the experiment probes the quasi-image potential states (ISs) of the Rydberg-like series converging to the vacuum level ( $E_V$ ) at 4.6 eV above  $E_F$ . We show here that this interpretation is incorrect and offer an alternative explanation.

Two well-documented facts invalidate the discussion presented in the Letter: (i) The work function of 4.6 eV used by the authors substantially differs from the values of up to 4.2 eV obtained in photoemission experiments and *ab initio* calculations [2,3]. (ii) For the bias of a few eV the tip-induced electric field in the junction overrides the image potential. The ISs at surfaces experience a Stark energy shift, and evolve into field emission resonances (FERS) [4–6]. Thus, the description of the metal-vacuum interface with image potential only is incorrect. We further illustrate [Fig. 1(a)] point (ii) with calculation of the energies of QWSs in free-electron Pb/Cu(111) as a function of a uniform electric field within a 1D model [7]. The states ( $E \lesssim 3$  eV) localized inside the Pb film are only mildly sensitive to the applied field. As to the QWSs with essential IS character close to  $E_V$ : the field as low as 0.05 eV/ $a_0$  (corresponding to the tip surface distance as large as 42 Å for the bias of 4 eV) destroys the Rydberg-like series in full accord with *ab initio* results [4].

Here we conjecture that the Pb band structure along the  $\Gamma$ - $L$  direction perpendicular to the surface of the film is at the origin of the results reported in [1]. Conclusive evidence supporting our explanation is achieved with data analysis as developed in Ref. [8]. Within the phase accumulation model, QWSs induced by the Pb overlayer of thickness  $D$  at  $\bar{\Gamma}$  are characterized by the phase relation  $\phi(E_n) + 2Dk(E_n) = 2\pi n$ .  $\phi(E_n)$  is the scattering phase shift accumulated at the interfaces of the overlayer. If, for the overlayers  $D$  and  $D'$  there is a corresponding pair of quantum numbers  $n$  and  $n'$  such that  $E_n \simeq E_{n'} = E$ , the  $\phi$  can be approximately canceled out, and the energy-dependent wave vector is  $k(E) = \pi(n' - n)/(D' - D)$ . Under the assumption that  $E_n = E_{n'}$  for the states within 40 meV energy window ( $\Delta$ ), we obtain from the data of Ref. [1] the Pb band along  $\Gamma$ - $L$ . Results are shown in Fig. 1(b) together with data from [8] for Pb/Ag(111) ( $\Delta = 20$  meV) and *ab initio* band structure calculations [9]. The calculated band dispersion saturates at 5.4 eV for

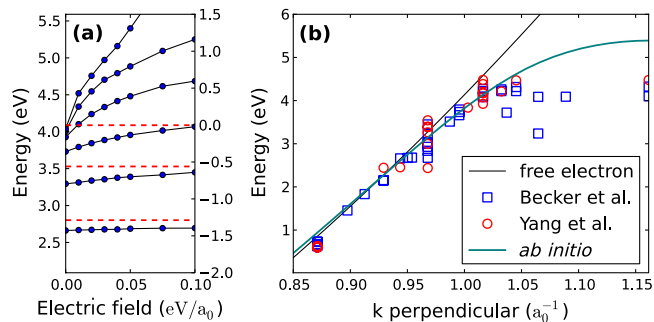


FIG. 1 (color online). (a) Calculated QWS energies (lines with dots) for 14ML-Pb/Cu(111) as a function of the electric field in the STM junction. The left (right) axis: the energy with respect to the  $E_F$  ( $E_V$ ). Dashed horizontal lines: experimental data of [1]. (b) Bulk Pb band dispersion along  $\Gamma$ - $L$  derived from experimental QWS energies (with respect to  $E_F$ ). Results are compared with free-electron dispersion from the present 1D model and with *ab initio* calculations [9] (shifted by +0.4 eV to coincide with photoemission data in [10]).

wave vector  $k$  approaching the reciprocal lattice vector  $G = 1.161a_0^{-1}$ , i.e., at  $\Gamma$  point. The agreement between the *ab initio* results, photoemission data [10], and these extracted from experimental STS data confirms the validity of our interpretation.

This work was partially funded by MCINN(FIS2010-19609-C02-01) and G.V-UPV/EHU(IT-366-07).

A. Zugarramurdi,<sup>1,2</sup> N. Zabala,<sup>1,2</sup>

A. G. Borisov,<sup>3</sup> and E. V. Chulkov<sup>2,4</sup>

<sup>1</sup>Elektrizitatea eta Elektronika, FCT-ZTF, UPV/EHU, 48080, Bilbao, Spain

<sup>2</sup>CFM, Centro Mixto CSIC-UPV/EHU and DIPC, 20018, San Sebastián, Spain

<sup>3</sup>ISMO, UMR 8214 CNRS-Université Paris-Sud, 91405 Orsay CEDEX, France

<sup>4</sup>Física de Materiales, UPV/EHU, 20080, San Sebastián, Spain

Received 25 March 2011; published 14 June 2011

DOI: [10.1103/PhysRevLett.106.249601](https://doi.org/10.1103/PhysRevLett.106.249601)

PACS numbers: 68.37.Ef, 68.65.Fg, 73.21.Fg

- [1] M. C. Yang *et al.*, *Phys. Rev. Lett.* **102**, 196102 (2009).
- [2] D. Yu and M. Scheffler, *Phys. Rev. B* **70**, 155417 (2004).
- [3] P. S. Kirchmann and U. Bovensiepen, *Phys. Rev. B* **78**, 035437 (2008).
- [4] P. Wahl *et al.*, *Phys. Rev. Lett.* **91**, 106802 (2003).
- [5] A. Hanuschkin D. Wortmann, and S. Blügel, *Phys. Rev. B* **76**, 165417 (2007).
- [6] D. B. Dougherty *et al.*, *Phys. Rev. B* **76**, 125428 (2007).
- [7] A. Zugarramurdi *et al.*, *Phys. Rev. B* **80**, 115425 (2009).
- [8] M. Becker and R. Berndt, *Phys. Rev. B* **81**, 205438 (2010).
- [9] I-Po Hong *et al.*, *Phys. Rev. B* **80**, 081409(R) (2009).
- [10] K. Horn *et al.*, *Phys. Rev. B* **30**, 1711 (1984).