# ON POSITIVE SOLUTION FOR A CLASS OF QUASILINEAR ELLIPTIC SYSTEMS WITH SIGN-CHANGING WEIGHTS 

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Abstract. In this paper, we consider the problem for the existence of positive solutions of quasilinear elliptic system

$$
\left\{\begin{array}{lrl}
-\Delta_{p} u=\lambda a(x) u^{\alpha} v^{\gamma}, & x \in \Omega \\
-\Delta_{q} v=\lambda b(x) u^{\eta} v^{\beta}, & x \in \Omega \\
u=v=0, & x \in \partial \Omega
\end{array}\right.
$$

where the $\lambda>0$ is a parameter, $\Omega$ is a bounded domain in $\mathbb{R}^{N}(N>1)$ with smooth boundary $\partial \Omega$, and the $\Delta_{p} z=\operatorname{div}\left(|\nabla z|^{p-2} \nabla z\right)$ is the $p$-Laplacian operator. Here $a(x)$ and $b(x)$ are $C^{1}$ sign-changing functions that maybe are negative near the boundary. Using the method of sub-super solutions and comparison principle, which studied the existence of positive solutions for quasilinear elliptic system. The main results of the present paper are new and extend the previously known results.

## 1. Introduction

In this note we consider the existence of positive solutions for the system

$$
\begin{cases}-\Delta_{p} u=\lambda a(x) u^{\alpha} v^{\gamma}, & x \in \Omega  \tag{1.1}\\ -\Delta_{q} v=\lambda b(x) u^{\eta} v^{\beta}, & x \in \Omega \\ u=v=0, & x \in \partial \Omega\end{cases}
$$

where $\lambda>0$ is a parameter, $1<p, q<N$, and $\Omega$ is a bounded domain in $\mathbb{R}^{N}(N>1)$ with smooth boundary $\partial \Omega$, and the $\Delta_{p} z=\operatorname{div}\left(\left|\nabla_{z}\right|^{p-2} \nabla_{z}\right)$ is the $p$-Laplacian operator. Here $a(x)$ and $b(x)$ are $C^{1}$ sign-changing functions that maybe are negative near the boundary.

Problem (1.1) arises in the theory of quasi-regular and quasi-conformal mappings as well as in the study of non- Newtonian fluids. In the latter case, the pair $(p, q)$ is a characteristic of the medium. Media with $(p, q)>(2,2)$ are called dilatant fluids and those with $(p, q)<(2,2)$ are called pseudo-plastics. If $(p, q)=(2,2)$, they are Newtonian fluids.

[^0]When $p=q=2$, the following system

$$
\begin{cases}\Delta u=a(|x|) v^{\alpha}, & x \in \mathbb{R}^{N}, \\ \Delta v=b(|x|) u^{\beta}, & x \in \mathbb{R}^{N}\end{cases}
$$

for which existence results for boundary blow-up positive solutions can be found in a recent paper by Lair and Wood [12]. The authors established that all positive entire radial solutions of systems above are boundary blow-up provided that

$$
\int_{0}^{\infty} t a(t) d t=\infty, \quad \int_{0}^{\infty} t b(t) d t=\infty .
$$

On the other hand, if

$$
\int_{0}^{\infty} t a(t) d t<\infty, \quad \int_{0}^{\infty} t b(t) d t<\infty
$$

then all positive entire radial solutions of this system are bounded.
F. Cìrstea and V.Rǎdulescu [5] extended the above results to a larger class of systems

$$
\begin{cases}\Delta u=a(|x|) g(v), & x \in \mathbb{R}^{N} \\ \Delta v=b(|x|) f(u), & x \in \mathbb{R}^{N}\end{cases}
$$

Z.D. Yang [15] extended the above results to a class of systems

$$
\begin{cases}\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right)=a(|x|) g(v), & x \in \mathbb{R}^{N}, \\ \operatorname{div}\left(|\nabla v|^{q-2} \nabla v\right)=b(|x|) f(u), & x \in \mathbb{R}^{N}\end{cases}
$$

Caisheng Chen [3] discussed the existence and non-existence of positive weak solution to the following system

$$
\begin{cases}-\Delta u=\lambda u^{\alpha} v^{\gamma}, & x \in \Omega \\ -\Delta v=\lambda u^{\delta} v^{\beta}, & x \in \Omega \\ u(x)=v(x)=0, & x \in \partial \Omega\end{cases}
$$

D.D. Hai [11] studied the existence and nonexistence of positive solutions for the quasilinear system

$$
\begin{cases}-\Delta_{p} u=\lambda a(x) f(u, v), & x \in \Omega,  \tag{E}\\ -\Delta_{q} v=\mu b(x) g(u, v), & x \in \Omega \\ u=v=0, & x \in \partial \Omega\end{cases}
$$

where $\Omega$ is a bounded domain in $\mathbb{R}^{N}$ with smooth boundary $\partial \Omega, p, q>1, \lambda, \mu$ are positive parameters, $a(x), b(x)$ are bounded functions that can change sign, which obtained existence results for the quasilinear system $(E)$ when $f(t, t)$ is p-sublinear in 0 and $g(t, t)$ is $q$-sublinear at 0 , and $\lambda, \mu$ are small. Nonexistence results are also obtained.

Motivated by the above results, we focus on further extending the study in [3] to the system (1.1) and supplementary the results in [11]. In fact, we study the existence of positive solution to the system (1.1) with sign-changing weight functions $a(x)$ and $b(x)$. Due to this weight functions, the existence are challenging and nontrivial. Our approach is based on the method of sub-super solutions, see [3, 14].

To precisely state our existence result we need the eigenvalue problem

$$
\begin{align*}
& -\Delta_{p} \phi_{1}=\lambda_{1}\left|\phi_{1}\right|^{p-2} \phi_{1}, \quad x \in \Omega, \quad \phi_{1}=0, \quad x \in \partial \Omega  \tag{1.2}\\
& -\Delta_{q} \phi_{1}^{*}=\lambda_{1}^{*}\left|\phi_{1}^{*}\right|^{q-2} \phi_{1}^{*}, \quad x \in \Omega, \quad \phi_{1}^{*}=0, x \in \partial \Omega \tag{1.3}
\end{align*}
$$

Let $\lambda_{1}>0$ be the principal eigenvalue and $\phi_{1}>0$ with $\left\|\phi_{1}\right\|_{\infty}=1$ the corresponding eigenfunction of $-\Delta_{p}$ and $\lambda_{1}^{*}>0$ be the principal eigenvalue and $\phi_{1}^{*}>0$ with $\left\|\phi_{1}^{*}\right\|_{\infty}=1$ the corresponding eigenfunction of $-\Delta_{q}$, with the Dirichlet boundary condition. It is well known that

$$
\frac{\partial \phi_{1}}{\partial v}<0 \quad \text { and } \quad \frac{\partial \phi_{1}^{*}}{\partial v}<0 \quad \text { on } \partial \Omega
$$

where $v$ is the unit outward normal, while $\phi_{1}, \phi_{1}^{*}=0$ on $\partial \Omega$. This result is well known and hence, depending on $\Omega$, there exist $\sigma, \sigma^{*} \in(0,1], \delta>0$ and $m>0$ such that (see [14])

$$
\begin{gather*}
\lambda_{1} \phi_{1}^{p}-\left|\nabla \phi_{1}\right|^{p} \leqslant-m, \text { on } \bar{\Omega}_{\delta},  \tag{1.4}\\
\phi_{1} \geqslant \sigma, \text { on } \Omega_{0}=\Omega \backslash \bar{\Omega}_{\delta}, \tag{1.5}
\end{gather*}
$$

and

$$
\begin{gather*}
\lambda_{1} \phi_{1}^{* q}-\left|\nabla \phi_{1}^{*}\right|^{q} \leqslant-m, \text { on } \bar{\Omega}_{\delta}  \tag{1.6}\\
\phi_{1}^{*} \geqslant \sigma^{*}, \text { on } \Omega_{0}=\Omega \backslash \bar{\Omega}_{\delta} \tag{1.7}
\end{gather*}
$$

where $\Omega_{\delta}=\{x \in \Omega \mid d(x, \partial \Omega)<\delta\}$. We will also consider the unique solution, $e_{1}(x)$, $e_{2}(x) \in C^{1}(\bar{\Omega})$, of the boundary value problem

$$
\begin{array}{ll}
-\Delta_{p} e_{1}=1, & x \in \Omega, \quad e_{1}=0, \quad x \in \partial \Omega \\
-\Delta_{q} e_{2}=1, & x \in \Omega, \quad e_{2}=0, \quad x \in \partial \Omega \tag{1.9}
\end{array}
$$

to discuss our existence result. It is known that $e_{i}(x)>0(i=1,2)$ in $\Omega$ and

$$
\frac{\partial e_{i}(x)}{\partial v}<0 \quad \text { on } \partial \Omega(i=1,2) \quad(\text { see }[8,9,10])
$$

## 2. Existence results

In this section, we shall establish our existence result via the method of sub and supersolutions. A pair of nonnegative functions $\left(\psi_{1}, \psi_{2}\right),\left(z_{1}, z_{2}\right)$, are called a subsolution and supersolution of (1.1) if they satisfy $\left(\psi_{1}, \psi_{2}\right)=(0,0)=\left(z_{1}, z_{2}\right)$, on $\partial \Omega$

$$
\begin{aligned}
& \int_{\Omega}\left|\nabla \psi_{1}\right|^{p-2}\left|\nabla \psi_{1}\right| \cdot \nabla f_{1} d x \leqslant \lambda \int_{\Omega} a(x) \psi_{1}^{\alpha} \psi_{2}^{\gamma} f_{1} d x \\
& \int_{\Omega}\left|\nabla \psi_{2}\right|^{q-2}\left|\nabla \psi_{2}\right| \cdot \nabla f_{2} d x \leqslant \lambda \int_{\Omega} b(x) \psi_{1}^{\eta} \psi_{2}^{\beta} f_{2} d x,
\end{aligned}
$$

and

$$
\begin{aligned}
& \int_{\Omega}\left|\nabla z_{1}\right|^{p-2}\left|\nabla z_{1}\right| \cdot \nabla f_{1} d x \geqslant \lambda \int_{\Omega} a(x) z_{1}^{\alpha} z_{2}^{\gamma} f_{1} d x \\
& \int_{\Omega}\left|\nabla z_{2}\right|^{p-2}\left|\nabla z_{2}\right| \cdot \nabla f_{2} d x \geqslant \lambda \int_{\Omega} a(x) z_{1}^{\eta} z_{2}^{\beta} f_{2} d x
\end{aligned}
$$

for all test functions $f_{1}(x) \in W_{0}^{1, p}(\Omega)$ and $f_{2}(x) \in W_{0}^{1, q}(\Omega)$ with $f_{1}, f_{2} \geqslant 0$. Then the following result holds:

Lemma 1. (See [17]) Suppose there exist sub and super-solutions $\left(\psi_{1}, \psi_{2}\right)$ and $\left(z_{1}, z_{2}\right)$ respectively of $(1.1)$ such that $\left(\psi_{1}, \psi_{2}\right) \leqslant\left(z_{1}, z_{2}\right)$. Then (1.1) has a solution $(u, v)$ such that $(u, v) \in\left[\left(\psi_{1}, \psi_{2}\right),\left(z_{1}, z_{2}\right)\right]$.

We make the following assumptions:
(i) $\alpha, \beta \geqslant 0, \gamma, \eta>0$ and $(p-1-\alpha)(q-1-\beta)>\gamma \eta$;
(ii) Assume that there exist positive constants $a_{0}, a_{1}, b_{0}$ and $b_{1}$, such that

$$
a(x) \geqslant-a_{0}, b(x) \geqslant-b_{0} \text { on } \bar{\Omega}_{\delta}
$$

and

$$
a(x) \geqslant a_{1}, b(x) \geqslant b_{1} \text { on } \Omega \backslash \bar{\Omega}_{\delta}
$$

(iii) Suppose that there exists $\varepsilon>0$ such that:

$$
\frac{\lambda_{1}}{m} a_{0}<\min \left\{c_{1}, c_{2} \varepsilon^{d_{2}-d_{1}}\right\}, \quad \frac{\lambda_{1}}{m} a_{0}<\min \left\{c_{1} \varepsilon^{-\left(d_{2}-d_{1}\right)}, c_{2}\right\}
$$

and

$$
\max \left\{\frac{\lambda_{1}}{c_{1}} \varepsilon^{1-d_{1}}, \frac{\lambda_{1}^{*}}{c_{2}} \varepsilon^{1-d_{2}}\right\} \leqslant \min \left\{\frac{1}{\|a\|_{\infty}}, \frac{1}{\|b\|_{\infty}}\right\}
$$

where

$$
\begin{gathered}
c_{1}=a_{1}\left(\frac{p-1}{p} \sigma^{\frac{p}{p-1}}\right)^{\alpha}\left(\frac{q-1}{q} \sigma^{* \frac{q}{q-1}}\right)^{\gamma}, c_{2}=b_{1}\left(\frac{p-1}{p} \sigma^{\frac{p}{p-1}}\right)^{\eta}\left(\frac{q-1}{q} \sigma^{*} \frac{q}{q-1}\right)^{\beta}, \\
d_{1}=\frac{\alpha}{p-1}+\frac{\gamma}{q-1}, \quad d_{2}=\frac{\eta}{p-1}+\frac{\beta}{q-1} .
\end{gathered}
$$

Now we are ready to state our existence results.

THEOREM 1. Let (i) - (iii) hold. Then there exists a positive solution of (1.1) for every $\lambda \in[\underline{\lambda}(\varepsilon), \bar{\lambda}(\varepsilon)]$, where

$$
\begin{gather*}
\bar{\lambda}=\min \left\{\frac{m}{a_{0} \varepsilon^{d_{1}-1}}, \frac{m}{b_{0} \varepsilon^{d_{1}-1}}, \frac{1}{\|a\|_{\infty}}, \frac{1}{\|b\|_{\infty}}\right\},  \tag{2.1}\\
\underline{\lambda}=\max \left\{\frac{\lambda_{1}}{c_{1} \varepsilon^{d_{1}-1}}, \frac{\lambda_{1}^{*}}{c_{2} \varepsilon^{d_{2}-1}}\right\} . \tag{2.2}
\end{gather*}
$$

REMARK 1. Note that (iii) implies $\underline{\lambda}<\bar{\lambda}$.

Proof. Let

$$
\left(\psi_{1}, \psi_{2}\right)=\left(\frac{p-1}{p} \varepsilon^{\frac{1}{p-1}} \phi_{1}^{\frac{p}{p-1}}, \frac{q-1}{q} \varepsilon^{\frac{1}{q-1}} \phi_{1}^{* \frac{q}{q-1}}\right)
$$

we shall verify that $\left(\psi_{1}, \psi_{2}\right)$ is a sub-solution of (1.1). Let $f_{1} \in W_{0}^{1, p}(\Omega)$, then a calculation shows that

$$
\begin{align*}
\int_{\Omega}\left|\nabla \psi_{1}\right|^{p-2} \nabla \psi_{1} \cdot \nabla f_{1} d x & =\varepsilon \int_{\Omega} \phi_{1}\left|\nabla \phi_{1}\right|^{p-2} \nabla \phi_{1} \cdot \nabla f_{1} d x \\
& =\varepsilon\left\{\int_{\Omega}\left|\nabla \phi_{1}\right|^{p-2} \nabla\left(\phi_{1} f_{1}\right) d x-\int_{\Omega}\left|\nabla \phi_{1}\right|^{p} f_{1} d x\right\}  \tag{2.3}\\
& =\varepsilon\left\{\int_{\Omega}\left[\lambda_{1} \phi_{1}^{p}-\left|\nabla \phi_{1}\right|^{p}\right] f_{1} d x\right\}
\end{align*}
$$

A similarly calculation shows that

$$
\begin{equation*}
\int_{\Omega_{2}}\left|\nabla \psi_{2}\right|^{q-2} \nabla \psi_{2} \cdot \nabla f_{2} d x=\varepsilon\left\{\int_{\Omega^{2}}\left[\lambda_{1}^{*} \phi_{1}^{* q}-\left|\nabla \phi_{1}^{*}\right|^{q}\right] f_{2} d x\right\} \tag{2.4}
\end{equation*}
$$

First, we consider the case when $x \in \bar{\Omega}_{\delta}$. We have $\lambda_{1} \phi_{1}^{p}-\left|\nabla \phi_{1}\right|^{p} \leqslant-m$ on $\bar{\Omega}_{\delta}$ and since $\lambda \leqslant \bar{\lambda}$, we have $\lambda \leqslant \frac{m}{a_{0} \varepsilon^{d_{1}-1}}$. Then

$$
\begin{equation*}
-\varepsilon m \leqslant-\lambda a_{0} \varepsilon^{\frac{\alpha}{p-1}+\frac{\gamma}{q-1}} . \tag{2.5}
\end{equation*}
$$

Hence

$$
\begin{align*}
\varepsilon\left(\lambda_{1} \phi_{1}^{p}-\left|\nabla \phi_{1}\right|^{p}\right) & \leqslant-m \varepsilon \\
& \leqslant-\lambda a_{0} \varepsilon^{\frac{\alpha}{p-1}+\frac{\gamma}{q-1}} \\
& \leqslant-\lambda a_{0}\left(\frac{p-1}{p} \varepsilon^{\frac{1}{p-1}}\left\|\phi_{1}\right\|_{\infty}^{\frac{p}{p-1}}\right)^{\alpha}\left(\frac{q-1}{q} \varepsilon^{\frac{1}{q-1}}\left\|\phi_{1}^{*}\right\|_{\infty}^{\frac{q}{q-1}}\right)^{\gamma}  \tag{2.6}\\
& \leqslant \lambda a(x) \psi_{1}^{\alpha} \psi_{2}^{\gamma} .
\end{align*}
$$

A similar argument shows that:

$$
\begin{equation*}
\varepsilon\left(\lambda_{1}^{*} \phi_{1}^{* q}-\left|\nabla \phi_{1}^{*}\right|^{q}\right) \leqslant \lambda b(x) \psi_{1}^{\eta} \psi_{2}^{\beta} . \tag{2.7}
\end{equation*}
$$

Then we obtain from (2.3), (2.4) and (2.6), (2.7) that

$$
\begin{align*}
& \int_{\bar{\Omega}_{\delta}}\left|\nabla \psi_{1}\right|^{p-2}\left|\nabla \psi_{1}\right| \cdot \nabla f_{1} d x \leqslant \lambda \int_{\bar{\Omega}_{\delta}} a(x) \psi_{1}^{\alpha} \psi_{2}^{\gamma} f_{1} d x  \tag{2.8}\\
& \int_{\bar{\Omega}_{\delta}}\left|\nabla \psi_{2}\right|^{q-2}\left|\nabla \psi_{2}\right| \cdot \nabla f_{2} d x \leqslant \lambda \int_{\bar{\Omega}_{\delta}} b(x) \psi_{1}^{\eta} \psi_{2}^{\beta} f_{2} d x \tag{2.9}
\end{align*}
$$

On the other hand, on $\Omega \backslash \bar{\Omega}_{\delta}$, we note that

$$
\phi_{1} \geqslant \sigma>0, \phi_{1}^{*} \geqslant \sigma^{*}>0, a(x) \geqslant a_{1}, b(x) \geqslant b_{1}
$$

and since $\lambda \geqslant \underline{\lambda}$, we have $\lambda \geqslant \frac{\lambda_{1}}{c_{1} \varepsilon^{d_{1}-1}}$. Then

$$
\begin{equation*}
\varepsilon \lambda_{1} \leqslant \lambda a_{1}\left(\frac{p-1}{p} \varepsilon^{\frac{1}{p-1}} \sigma^{\frac{p}{p-1}}\right)^{\alpha}\left(\frac{q-1}{q} \varepsilon^{\frac{1}{q-1}} \sigma^{* \frac{q}{q-1}}\right)^{\gamma} . \tag{2.10}
\end{equation*}
$$

Hence

$$
\begin{align*}
\varepsilon\left(\lambda_{1} \phi_{1}^{p}-\left|\nabla \phi_{1}\right|^{p}\right) & \leqslant \varepsilon \lambda_{1} \phi_{1}^{p} \\
& \leqslant \varepsilon \lambda_{1}\left\|\phi_{1}\right\|_{\infty}^{p} \\
& \leqslant \varepsilon \lambda_{1}  \tag{2.11}\\
& \leqslant a_{1}\left(\frac{p-1}{p} \varepsilon^{\frac{1}{p-1}} \sigma^{\frac{p}{p-1}}\right)^{\alpha}\left(\frac{q-1}{q} \varepsilon^{\frac{1}{q-1}} \sigma^{*} \frac{q}{q-1}\right)^{\gamma} \\
& \leqslant \lambda a(x) \psi_{1}^{\alpha} \psi_{2}^{\gamma} .
\end{align*}
$$

A similar argument shows that:

$$
\begin{equation*}
\varepsilon\left(\lambda_{1}^{*} \phi_{1}^{* q}-\left|\nabla \phi_{1}^{*}\right|^{q}\right) \leqslant \lambda b(x) \psi_{1}^{\eta} \psi_{2}^{\beta} \tag{2.12}
\end{equation*}
$$

Then we obtain from (2.3), (2.4) and (2.11), (2.12) that

$$
\begin{align*}
& \int_{\Omega-\bar{\Omega}_{\delta}}\left|\nabla \psi_{1}\right|^{p-2}\left|\nabla \psi_{1}\right| \cdot \nabla f_{1} d x \leqslant \lambda \int_{\Omega-\bar{\Omega}_{\delta}} a(x) \psi_{1}^{\alpha} \psi_{2}^{\gamma} f_{1} d x  \tag{2.13}\\
& \int_{\Omega-\bar{\Omega}_{\delta}}\left|\nabla \psi_{2}\right|^{q-2}\left|\nabla \psi_{2}\right| \cdot \nabla f_{2} d x \leqslant \lambda \int_{\Omega-\bar{\Omega}_{\delta}} b(x) \psi_{1}{ }^{\eta} \psi_{2}^{\beta} f_{2} d x \tag{2.14}
\end{align*}
$$

Since $\Omega=\bar{\Omega}_{\delta} \cup\left(\Omega \backslash \bar{\Omega}_{\delta}\right)$, We obtain from (2.8), (2.9) and (2.13), (2.14) that

$$
\begin{align*}
& \int_{\Omega}\left|\nabla \psi_{1}\right|^{p-2}\left|\nabla \psi_{1}\right| \cdot \nabla f_{1} d x \leqslant \lambda \int_{\Omega} a(x) \psi_{1}^{\alpha} \psi_{2}^{\gamma} f_{1} d x  \tag{2.15}\\
& \int_{\Omega}\left|\nabla \psi_{2}\right|^{q-2}\left|\nabla \psi_{2}\right| \cdot \nabla f_{2} d x \leqslant \lambda \int_{\Omega} b(x) \psi_{1}^{\eta} \psi_{2}^{\beta} f_{2} d x \tag{2.16}
\end{align*}
$$

we have shown that $\left(\psi_{1}, \psi_{2}\right)$ is sub-solution.
Now, we will construct a super-solution $\left(z_{1}, z_{2}\right)$ of (1.1). It is clear that:

$$
\begin{equation*}
-\operatorname{div}\left(\left|\nabla z_{1}\right|^{p-2} \nabla z_{1}\right)=A, \quad x \in \Omega \tag{2.17}
\end{equation*}
$$

$$
\begin{equation*}
-\operatorname{div}\left(\left|\nabla z_{2}\right|^{q-2} \nabla z_{2}\right)=B, \quad x \in \Omega \tag{2.18}
\end{equation*}
$$

We denote

$$
\begin{equation*}
z_{1}(x)=A e_{1}(x), z_{2}(x)=B e_{2}(x) \tag{2.19}
\end{equation*}
$$

where the constants $A, B>0$ are large and to be chosen later. We shall verify that is a super-solution of (1.1).

Next, since $\lambda \leqslant \bar{\lambda}$ we have $\lambda \leqslant \frac{1}{\|b\|_{\infty}}, \lambda \leqslant \frac{1}{\|a\|_{\infty}}$. Let $l_{1}=\left\|e_{1}\right\|_{\infty}, l_{2}=\left\|e_{2}\right\|_{\infty}$, since (i) hold, it is easy to prove that there exist positive large constants $A, B$ such that([3]):

$$
\begin{align*}
A^{p-1-\alpha} & \geqslant B^{\gamma} l_{1}^{\alpha} l_{2}^{\gamma} \\
& \geqslant \lambda\|a\|_{\infty} B^{\gamma} l_{1}^{\alpha} l_{2}^{\gamma} \tag{2.20}
\end{align*}
$$

and

$$
\begin{align*}
B^{q-1-\beta} & \geqslant A^{\eta} l_{1}^{\eta} l_{2}^{\beta}  \tag{2.21}\\
& \geqslant \lambda\|b\|_{\infty} A^{\eta} l_{1}^{\eta} l_{2}^{\beta}
\end{align*}
$$

These imply that:

$$
\begin{equation*}
A^{p-1} \geqslant \lambda a(x) z_{1}^{\alpha} z_{2}^{\gamma}, \quad B^{q-1} \geqslant \lambda b(x) z_{1}^{\delta} z_{2}^{\beta} \tag{2.22}
\end{equation*}
$$

Let $f_{1} \in W_{0}^{1, p}(\Omega), f_{2} \in W_{0}^{1, q}(\Omega)$ with $f_{1}, f_{2}>0$. Then we obtain from (2.12), (2.13) and (2.17) that

$$
\begin{align*}
& \int_{\Omega}\left|\nabla z_{1}\right|^{p-2}\left|\nabla z_{1}\right| \cdot \nabla f_{1} d x=A^{p-1} \int_{\Omega} f_{1}(x) d x \geqslant \lambda \int_{\Omega} a(x) z_{1}{ }^{\alpha} z_{2}^{\gamma} f_{1} d x,  \tag{2.23}\\
& \int_{\Omega}\left|\nabla z_{2}\right|^{p-2}\left|\nabla z_{2}\right| \cdot \nabla f_{2} d x=B^{q-1} \int_{\Omega} f_{2}(x) d x \geqslant \lambda \int_{\Omega} a(x) z_{1}^{\eta} z_{2}^{\beta} f_{2} d x . \tag{2.24}
\end{align*}
$$

a.e. in $\Omega$. Thus, $\left(z_{1}, z_{2}\right)$ is a super-solution of (1.1). Obviously, we have $z_{i}(x) \geqslant$ $\psi_{i}(x)$ in $\Omega$ with large $A, B$ for $i=1,2$. Thus, by Lemma 1 , there exists a positive solution $(u, v)$ of (1.1) such that $\left(\psi_{1}, \psi_{2}\right) \leqslant(u, v) \leqslant\left(z_{1}, z_{2}\right)$. This completes the proof of Theorem 1 .

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