

Research Article

Near-Threshold Production of W^\pm , Z^0 , and H^0 at a Fixed-Target Experiment at the Future Ultrahigh-Energy Proton Colliders

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We outline the opportunities to study the production of the Standard Model bosons, W^\pm , Z^0 , and H^0 , at “low” energies at fixed-target experiments based on possible future ultrahigh-energy proton colliders, that is, the High-Energy LHC, the Super proton-proton Collider, and the Future Circular Collider hadron-hadron. These can be indeed made in conjunction with the proposed future colliders designed to reach up to $\sqrt{s} = 100$ TeV by using bent crystals to extract part of the halo of the beam which would then impinge on a fixed target. Without disturbing the collider operation, this technique allows for the extraction of a substantial amount of particles in addition to serving for a beam-cleaning purpose. With this method, high-luminosity fixed-target studies at centre-of-mass energies above the W^\pm , Z^0 , and H^0 masses, $\sqrt{s} \approx 170$ –300 GeV, are possible. We also discuss the possibility offered by an internal gas target, which can also be used as luminosity monitor by studying the beam transverse shape.

1. Introduction

In this paper, we consider the possibility of performing fixed-target experiments with the beams of the proposed high-energy LHC (HE-LHC) (see, e.g., [1]), the Future Circular Collider hadron-hadron (FCC-hh) (see, e.g., [2]), and the Super proton-proton Collider (SppC) (see, e.g., [3]), which, as their names indicate, are primarily intended for collider physics. The beam energy of these possible future facilities ranges from 16.5 up to 50 TeV, allowing for fixed-target collisions at centre-of-mass system (c.m.s.) energies ranging from 175 to 300 GeV.

Just as lower energy beams, these can in principle be extracted over the course of about a meter by using the channelling of particles in a bent crystal. This phenomenon is well documented (see, e.g., [4–8]) and has been experimentally studied for protons and heavy ions at energies per nucleon up to 900 GeV. Recently, studies performed at SLAC have shown that the beam bending by means of bent crystals is also possible for high-energy positrons and electrons [9]. In [10], it was discussed specifically for the LHC beams.

Channelling experiments at the LHC have been proposed [11], are installed [12, 13], and will be performed during Run-2 for beam collimation studies. The bent crystal extraction technique allows for the extraction of particles from the beam halo only, so that the collider experiments can be kept running simultaneously. These particles would anyway be lost to collimation and would not be used in the collider mode.

In [14], a comprehensive list of physics opportunities offered by the use of the multi-TeV proton and lead LHC beams on a fixed target was presented. Let us recall the critical assets of the fixed-target mode as compared to the collider mode, that is,

- (i) a quasi unlimited target-species versatility,
- (ii) a full access to the target rapidity, which corresponds to the far backward region in the centre-of-mass frame,
- (iii) the possibility to polarise the target,
- (iv) very large luminosities with modest beam intensity thanks to the high target density.

These advantages already translate, with the proton and lead LHC beams, into an impressive list of possible physics studies [14] beyond their respective state of the art, in particular as regards precision studies of spin-related observables with a polarised target and of quark-gluon plasma formation in lead-nucleus collisions in the target rapidity region at c.m.s. energies between those of SPS and RHIC and of QCD at large momentum fractions x in proton-proton, proton-deuteron, and proton-nucleus collisions, and so forth. We refer to [15] for the more specific case of quarkonium studies, to [16–19] for spin physics, and to [15, 20, 21] for heavy-ion physics. First simulation studies at the generator level have been presented in [22] and have demonstrated the great potential for both charmonium and bottomonium studies at $\sqrt{s} = 72$ and 115 GeV at a fixed-target experiment with the LHC beams (thereafter referred to as AFTER@LHC).

With beams of higher energies at future facilities, the available c.m.s. energies can nearly be three times as large as at AFTER@LHC and allow for even more systematic studies of systems whose masses are well above that of the bottomonia, that is, 10 GeV. At fixed-target LHC energies, W and Z production, sometimes generically referred to as the Drell-Yan-like processes, is just reachable with very low expected rates but with the advantage of potentially providing unique information about the nucleon structure at momentum fractions x close to unity ($x \approx M/\sqrt{s}e^{\pm y_{\text{c.m.s.}}}$) and about QCD corrections near the threshold (let us note here that the production of heavy Beyond-the-Standard-Model particles produced at the LHC in the collider mode might also be subject to similar QCD threshold corrections) and the advantage of offering interesting information on hadronic W decays. With beams of higher energies, rates would significantly be larger allowing, among other things, for rapidity dependent measurements. In general, the combination of high-luminosity hadron-hadron collisions at \sqrt{s} well above 100 GeV and a backward c.m.s coverage provide the opportunity to study the interplay between the—genuinely nonperturbative—confinement of partons at large momentum fractions x and the perturbative behaviour of the short-distance parton scatterings. Such extremely hard reactions are indeed believed to be well understood within the perturbative regime of QCD. Using a polarised target allows one to advance further the precision and the refinements of such studies of the hadron inner structure with information on the helicity of the partons and on their angular momentum when they carry most of the hadron momentum. The confinement properties of Quantum Chromodynamics, the theory of strong interaction, is still an open problem which deserves novel and innovative studies, even at high-energy facilities.

Although, exactly as for AFTER@LHC, the potential for physics studies go well beyond that of Drell-Yan-like studies in the mass region of W and Z bosons, we wish to focus on it here as an illustrative example of the gain offered by even higher energy beams—in what we believe to be the very first paper on the use of ultrahigh-energy beam in the fixed-target mode. The case for such W and Z studies is clear. So far, the production of W^\pm and Z bosons at RHIC could only be performed at $\sqrt{s} = 500$ GeV (see, e.g., [23, 24]) with a couple of thousand W^\pm candidates and less than one

hundred Z counts. Studies at lower energies, in order to reach x higher than 0.2 and measure the \bar{d}/\bar{u} flavour asymmetry at larger x and high Q^2 , require luminosities on the order of an inverse femtobarn which is out of reach at RHIC in particular if operated at $\sqrt{s} = 200$ GeV where its luminosity is significantly lower. In addition, the most accessible region in the fixed-target mode at high energies is that of backward c.m.s. rapidities, where x in the target can even be larger. Of course, H^0 production has so far only been observed at the LHC at 7 and 8 TeV [25, 26] and it is of interest to have a look at the conditions in which it could be studied not far from the threshold.

The structure of this paper is thus as follows. In Section 2, we discuss the beam extraction in a parasitic mode using the bent crystal technique at ultrahigh energies. In Section 3, we recall the main characteristics of possible future colliders as they are currently discussed and derive reasonable expectations for luminosities in the fixed-target mode. In Section 4, we discuss the case of the weak boson production as a benchmark of what can be achieved with luminosities up to 100 times larger than at RHIC. In addition to the expected rates, we briefly discuss the potential backgrounds whose precise size can however only be assessed with a proposed detector setup. Section 5 gathers our conclusions.

2. Beam Extraction by Means of a Bent Crystal with Ultrahigh-Energy Protons

As previously mentioned, these possible future collider facilities would use proton beams from 16.5 up to 50 TeV. The bending of GeV beams of protons and ions has been studied extensively during the past three decades. As a first approximation, one may calculate the approximate deflection efficiency as a function of crystal length as, for example, done in [6]. For example, at a deflection angle of 0.5 mrad, as approximately required for the passage of a septum blade downstream required for further extraction, the efficiency (excluding surface transmission) in Si (110) is 84% for a 50 TeV beam. This efficiency is obtained at the optimum crystal length of $L/L_D = 0.085$ (see Figure 1), corresponding to a length of 1.6 m.

However, the optimal distance calculated in this simulation may significantly be overestimated because it does not take into account the probability for dechanneled particles to get extracted on a later encounter with the crystal. Several studies have shown that this multipass mechanism may result in a significant shortening of the optimal distance. The optimal crystal thickness also depends on the beam optics used. Results of the order 20–30 cm are certainly not unrealistic.

3. Main Parameters of Future Colliders and Their Corresponding Characteristics in the Fixed-Target Mode

In order to derive the luminosities which can reasonably be expected in the fixed-target mode with the beams of future colliders, we start by recalling their main parameters as currently discussed. Indeed, efforts are now being made in

TABLE 1: Beam parameters for the proposed next-generation colliders and the corresponding fixed-target energies along with the boost ($\gamma_{\text{c.m.s.}}^{\text{lab}}$) and the rapidity shift ($\Delta y_{\text{c.m.s.}}^{\text{lab}}$) between the centre-of-mass frame of the fixed-target collision and the laboratory frame, which identifies to the centre-of-mass frame in the collider mode.

	SppC-1	SppC-2	HE LHC	FCC-hh
Beam energy (E_p) [TeV]	20	45	16.5	50
Fixed-target centre-of-mass energy ($\sqrt{2E_p m_N}$) [GeV]	194	291	176	307
Number of bunches stored	3000	6000	1404 (50 ns spacing)	10600/53000 (25 and 5 ns spacing)
Number of protons (N_p) per bunch [10^{11}]	$1.7 \cdot 10^{-3}$	$0.98 \cdot 10^{-3}$	1.3	1/0.2
$\gamma_{\text{c.m.s.}}^{\text{lab}} = \frac{\sqrt{s}}{2m_p}$	103	155	94	163
$\Delta y_{\text{c.m.s.}}^{\text{lab}} = \ln\left(\gamma_{\text{c.m.s.}}^{\text{lab}} + \sqrt{(\gamma_{\text{c.m.s.}}^{\text{lab}})^2 - 1}\right)$	5.3	5.7	5.2	5.8

investigating the best location and technology for future collider projects. These proposed circular colliders have circumferences of 50 to 100 km and the ability to circulate protons with 15 to 50 TeV energies. We list their most relevant characteristics in Table 1. In particular, we consider phases one and two of the SppC, denoted by SppC-1 and SppC-2, as discussed in [3]. We also consider the High-Energy Large Hadron Collider (HE LHC) [1] and the Future Circular Collider hadron-hadron (FCC-hh) [2].

With colliding beams of equal energies, the c.m.s. frame obviously corresponds to the laboratory frame. In the fixed-target mode, with the LHC 7 TeV protons, for instance, the boost ($\gamma_{\text{c.m.s.}}^{\text{lab}}$) and the rapidity shift ($\Delta y_{\text{c.m.s.}}^{\text{lab}}$) between the c.m.s. frame of the fixed-target collision and the laboratory frame are, respectively, $\gamma_{\text{c.m.s.}}^{\text{lab}} = \sqrt{s}/(2m_p) \approx 60$ and $\Delta y_{\text{c.m.s.}}^{\text{lab}} = \ln(\gamma_{\text{c.m.s.}}^{\text{lab}} + \sqrt{(\gamma_{\text{c.m.s.}}^{\text{lab}})^2 - 1}) \approx 4.8$. The region of central c.m.s. rapidities, $y_{\text{c.m.s.}} \approx 0$, is thus highly boosted at an angle with respect to the beam axis of about one degree in the laboratory frame. The entire backward hemisphere, $y_{\text{c.m.s.}} < 0$, is thus easily accessible with standard experimental techniques. With the future facilities, the rapidity shift is on the order of 5-6; see Table 1. A detector covering $\eta_{\text{lab}} \in [2, 6]$ would thus cover nearly half of the physical phase space of the fixed-target mode.

As we discussed in the previous section, the extraction of such high-energy beams by a bent crystal should not pose more challenges than at the LHC where it will be tested in the coming year. In this case, the main accelerator parameter fixing the luminosities achievable is the flux of the extracted beam. In the following discussion, we will assume (see [14] for a discussion on the LHC conditions where it corresponds to half of the beam loss) that it amounts to 5% of the protons stored in the beam over a fill lasting 10 hours. In the case of the LHC, such a parasitic mode corresponds to a proton flux of 5×10^8 per second and, on the average, to the extraction of minibunches of about 15 protons per bunch per pass with a 25 ns bunch spacing. In such a case, with a target thickness of 5–10% of interaction length, which is the case we consider here, the pileup is not an issue. The corresponding numbers for future facilities are given in Table 2.

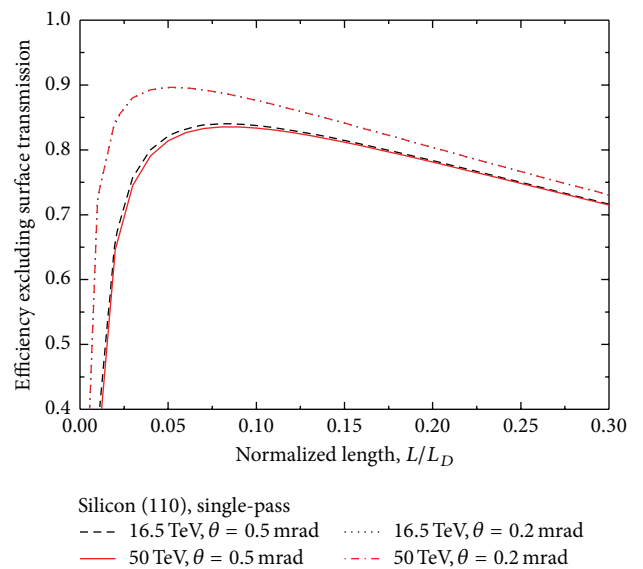


FIGURE 1: Single-pass efficiency, excluding surface transmission, calculated as in [6]. The angles and energies are given in the legends.

Yet, it may not be necessary to extract the beam from the collider to perform fixed-target experiments. By injecting a small amount of gas into the detector region of a running collider, one can sufficiently increase the probability of beam-gas interactions such as to reach decent luminosities yet without affecting at all the beam lifetime. At the LHC, the LHCb experiment has implemented such a system initially to monitor the beam luminosity [34–36] referred to as SMOG for System for Measuring Overlap with Gas (SMOG). SMOG has so far proved to be functioning well while not disturbing the primary beam. LHCb is currently analysing data in proton-neon and lead-neon collisions taken during beam tests in 2012 and 2013. In [19], the corresponding luminosities are given for the LHC. One may think that switching from a dense solid or liquid target to a dilute (10^{-9} Bar) gas necessarily decreases the luminosity. In fact, it may not be always so. Indeed, this decrease is compensated since the entire collider beam, amounting to a current close to an ampere for the LHC, traverses the gas cell, as opposed to the extraction beam which is similar to the beam loss.

TABLE 2: Luminosities reachable in the fixed target mode. The proton flux is calculated by assuming that 5% of the beam is used per fill of 10 hours. The luminosities are calculated for the case of targets that are 1 cm thick. The three values displayed represent luminosities for three different targets: liquid hydrogen, beryllium, and tungsten. The gas-target values are calculated using the same parameters as in [19] for a perfect gas at a pressure of 10^{-9} Bar in a zone of 100 cm.

	SppC-1	SppC-2	HE LHC	FCC-hh
Extracted beam				
Proton flux	$7.1 \cdot 10^5$	$8.1 \cdot 10^5$	$2.5 \cdot 10^8$	$1.5 \cdot 10^9$
$\mathcal{L}(\mu\text{b}^{-1}\text{s}^{-1})$	0.028/0.088/0.044	0.032/0.10/0.05	10/31/15	30/93/45
$\int dt \mathcal{L}(\text{pb}^{-1}\text{yr}^{-1})$	0.28/0.88/0.44	0.32/1.0/0.5	100/310/155	300/930/450
Gas target				
$\mathcal{L}(\mu\text{b}^{-1}\text{s}^{-1})$	0.014	0.016	5	30
$\int dt \mathcal{L}(\text{pb}^{-1}\text{yr}^{-1})$	0.14	0.16	50	300

Table 2 summarises all these numbers for 1 cm thick targets in the case of an extracted beam with a bent crystal and an internal gas target of 1 meter. In the former case and for light target materials, the luminosity can be increased by using a target much thicker/longer than 1 cm: NA51 at SPS used, for instance, a 1.2 m long hydrogen and deuterium target [37] with 450 GeV protons, E866 at Fermilab used 3 target cells of 50 cm [38] with 900 GeV protons, and COMPASS at SPS uses [39, 40] a 110 cm polarised NH_3 target with 160 GeV pions. In such a case, one can obtain luminosities *per annum* well above the inverse femtobarn. Even in this case, the thickness of the target does not reach more than 10% of interaction length.

Much higher luminosities could be achieved by using the full amount of the remaining protons stored in the beam at the end of each fill. Such a “dumping” mode, which could last an hour without unreasonably impacting the schedule of the machine, would of course provide luminosities orders-of-magnitude higher, probably up to 3, than the ones quoted in Table 2. However, this would also be done at the cost of carrying out the experiment in a highly activated environment, which may not be feasible in practice, and at the cost of a significant pileup. For some specific studies to look for rare events with very large momentum tracks, the latter may however not be an unsolvable issue.

4. Fixed-Target Mode and Boson Production

4.1. Expected Signal Rates. As announced we have decided to focus on the production of SM bosons production as an illustrative example of what the high luminosities reachable with the fixed-target mode can allow for. Physicswise, by measuring the production of Standard Model bosons W^\pm and Z^0 , the distribution of quarks and antiquarks at large x can be probed both in the proton with a hydrogen target and in the nuclei with nuclear targets. In particular, it allows one to determine the \bar{d}/\bar{u} flavour asymmetry at large x and large scales. The study of these reactions not far from the threshold also allows one to validate the theoretical methods to account for the so-called threshold resummation effects (see [41] for W production at RHIC).

To evaluate the cross section at NLO accuracy, we used the library MCFM [28] and set the $\mu_F = \mu_R$ equal to the boson mass. Since we are mainly interested in illustrating how

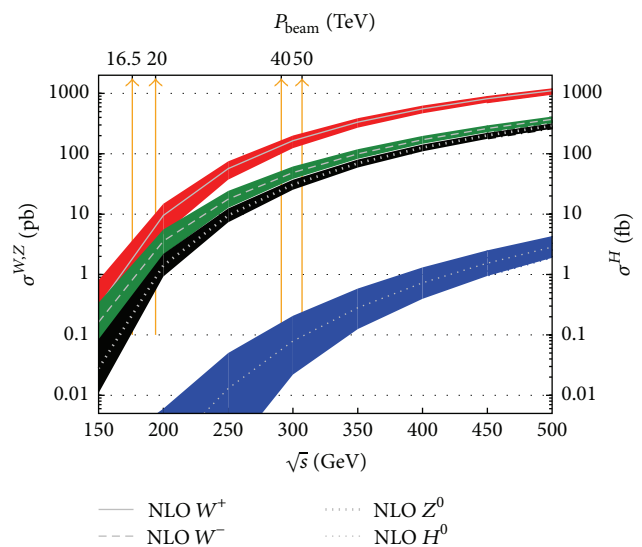


FIGURE 2: Cross sections for Standard Model boson production in proton-proton collisions at various centre-of-mass (lower x -axis) and corresponding beam energies in the fixed-target mode (upper x -axis). The color bands indicate the $1\text{-}\sigma$ values coming from the PDF uncertainty. The NLO calculations were performed with CT10NLO [27] by using MCFM [28] with $\mu_F = \mu_R$ set to the mass of the produced particle. The 4 orange arrows point at the beam energy of the 4 setups which we have considered.

such a measurement would help to better constrain parton distribution (PDF), we only show the theoretical uncertainty from these as they are currently determined. To this end, we use the NLO PDF set CT10 [27] and its associated eigenvector sets. MCFM takes all these into account and provide a $1\text{-}\sigma$ uncertainty which we have depicted on Figure 2. One observes an increasing PDF uncertainty for decreasing energies. With a 20 TeV beam, the uncertainty of W^+ production is as large as a factor of 3 for a total cross section about 10 pb.

If one sticks to the conventional leptonic decay channels, the branching is on the order of 10%. At 200 GeV, with a detector covering pseudorapidity from 2 to 6—a detector similar to LHCb [42] with a slightly more forward coverage, for instance—and imposing the lepton transverse momentum

(as well as the missing transverse momentum carried by the neutrino) to be larger than 20 GeV as usually done to cut the background (see later), a quick evaluation shows that the acceptance is on the order of 45% for the W as well as for the Z , although with a smaller branching ratio it is on the order of 3%. The central values for the cross sections times the corresponding leptonic branching at $\sqrt{s} = 200$ GeV in this fiducial volume ($P_{T,\ell,\nu_\ell} > 20$ GeV, $2 < \eta_{\ell,\nu_\ell} < 6$) are therefore 400, 150, and 20 fb, respectively, for the W^+ , W^- , and Z .

Relying on the performance of LHCb for similar studies [43], the efficiency including that of the triggering, the tracking, and additional selections (for the W : the lepton isolation, a cut on the energy deposit to limit the punch-through, the absence of a second lepton with a minimum P_T , and a vertex cut to remove the heavy-flavour decays) is around 40% for the W and 67% for the Z .

With a yearly luminosity of 15 fb^{-1} using, for instance, a 50 cm long liquid hydrogen target at a facility similar to a FCC-hh (see Table 2) one would expect a couple of thousands of W^+ events to be measured at 200 GeV. It would be one order of magnitude more at 300 GeV.

Z boson production is also at reach at 200 GeV with 200 dimuon events, especially if the rather clean environment associated with the rather low centre-of-mass energy allows for the use—and the study—of hadron decay channels. With a detector such as LHCb, both electron and muon decay channels can be used (see, e.g., [44]).

For the H^0 , the situation is very different with much smaller rates (on Figure 2, the cross sections for H^0 are shown in fb for readability). Since the process is dominated by gluon fusion—and we have checked that the vector-boson-fusion contribution is negligible down to low energies—the PDF uncertainties are very large. Fixed-target luminosities of 100 fb^{-1} seem to be needed to be able to have a hope to see a H^0 signal even with the 50 TeV FCC-hh proton beams.

For the W and Z bosons, with a very successful experiment, one may however be able to access regions more backward than $y_{\text{c.m.s.},V} = \ln(M_V/\sqrt{s})$ and hence to probe the quark distribution in the target for x larger than unity in the nuclear case (the yields in proton-nucleus collisions with a 1 cm thick gold or lead target should be similar to that in proton-proton collisions with a 1 m long liquid hydrogen target).

4.2. Expected Background. At such low energies, very few processes can mimic a dilepton pair of a mass around 90 GeV or an isolated lepton of a transverse momentum around 40 GeV accompanied with a missing transverse energy of a similar size. In fact, one expects the main background to be of electroweak origin such as the $Z \rightarrow \ell\ell$ for the $W \rightarrow \ell\nu_\ell$ channel where one lepton from the Z is lost. $W \rightarrow \tau\nu_\tau$ is also known to sometimes mimic a $\mu\nu_\mu$ final state, although usually at lower P_T . These should be tractable with data. As compared to studies at the LHC or at RHIC, QCD backgrounds are in general expected to be smaller.

For the W^\pm , Figure 3 illustrates that, even without any specific cuts usually used to reduce the background (see below), the P_T spectrum from heavy-flavour-decay electrons

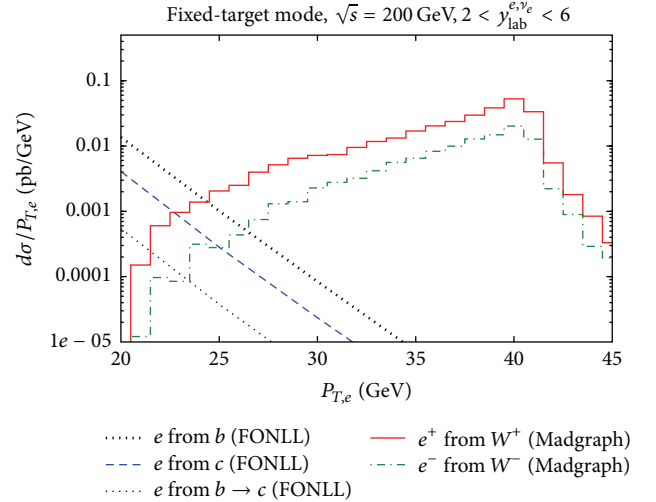


FIGURE 3: P_T spectra for heavy-flavour-decay electrons (and positrons) as predicted by FONLL compared to that from the predicted signal from W^\pm decays obtained with Madgraph 4 [29] at $\sqrt{s} = 200$ GeV. The FONLL cross sections have been evaluated with the default setup of <http://www.lpthe.jussieu.fr/~cacciari/fonll/fonll-form.html> along the lines of [30–32] with the PDFs CTEQ6.6M [33].

is extremely suppressed at large P_T and the electron yield close to $m_W/2$ is essentially purely from W decays. The same holds for muons. Let us however stress the fact that the main QCD background in the LHCb study of [43] is from hadron decay-in-flight which contributes less than 10%. Overall the purity of their W^\pm yield is 78%.

At RHIC, electron channels were used both by STAR [45, 46] and PHENIX [23, 47] and they required its isolation. Since PHENIX has an incomplete azimuthal coverage, the requirement for a missing transverse energy could not be imposed. Yet, the signal could be extracted. Once these requirements are applied, the W peak in the electron P_T spectrum around 40 GeV is evident. Let us note here that the background electron spectra reported in Figure 3 would severely be reduced once cuts for a missing transverse energy, for the isolation of the electron and for a maximum distance from the primary vertex, are imposed. Conversely, these could be measured—and then subtracted—by selecting displaced leptons.

With muons, it is even possible not to impose any lepton isolation to extract the W signal as demonstrated by the CMS study [48] in the very busy lead-lead environment at 2.76 TeV. At lower P_T , QCD backgrounds contribute essentially converted photons from π^0 and η decays, whereas muons from converted photon are usually negligible. The P_T spectrum of this background can be evaluated from dedicated simulations at the detector level but its size is however very difficult to predict without a precise knowledge of the hadron detector response. Usually the normalisation of such a background is simply adjusted on the data. In any case, it always has been found to be smaller than the signal at electron P_T above 30 GeV. We also note that the STAR detector is rather slow as compared to the LHCb detector, for instance, and tracks of

particles produced in earlier collisions can pile up in the data acquisition system, which further increases the background.

In the Z case [43], the main background at LHC energies is that of heavy-flavour decays with on average 3 background events per thousand! At lower energies, it should also be even more suppressed [23]. At 200 GeV, for instance, $\sigma_{b\bar{b}}$ ($2 < \eta_b < 6$; $P_{T,b} > 20$ GeV, $m_{b\bar{b}} > 80$ GeV) ~ 220 fb. Even if one neglects the momentum difference between the b quark, the beauty hadron, and the lepton, which is however a far too conservative approximation, one should multiply it by $f(B \rightarrow \ell) \simeq 0.1$ squared. The dilepton background from b decays is thus in any case much smaller than 2 fb whereas the signal size, accounting for the acceptance, is 20 fb. As regards the uncorrelated background from hadrons, no same-sign dimuon was found by LHCb with 40 pb^{-1} of data [43] and, with 1 fb^{-1} of data, they determined [49] it to be 0.2% of the signal with an overall purity of more than 0.99. At 500 GeV, STAR did not report any same-sign events [45] although with a, necessary less clean, dielectron sample.

Given the likely smaller background at energies below 500 GeV, the fact that the muon channel, with smaller background than that for RHIC studies, would be preferred with a much reduced background and the likely strong dependence of any background simulation on specific detector performances, we chose not to perform any generator level simulation and tend to advocate in view of past experiments at higher energies that, using conventional cuts resulting in tractable acceptances and efficiencies (see above), such signals should easily be extractable.

5. Conclusions

The current planning of future proton colliders necessitates a discussion of whether these facilities could also be used in a fixed-target mode. There is a long list of physics arguments that supports this case at the LHC. We have shown calculations specific to Standard Model bosons at the HE-LHC, SppC, and FCC-hh. This next generation of fixed-target experiments would provide access to high-luminosity measurements at unique laboratory energies and momentum transfers. Using a bent crystal is a viable option to extract high-energy beams of protons and perform—in a parasitic mode—fixed-target experiment at $\sqrt{s} \simeq 170\text{--}300$ GeV with annual luminosities on the order of tens of inverse femtobarn, that is, with high enough rates to produce a significant amount of Standard Model bosons. Although it offers a priori smaller luminosities, an internal gas target is an option which probably requires less civil engineering.

Conflict of Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

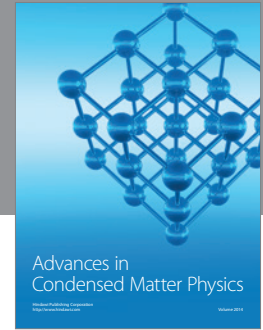
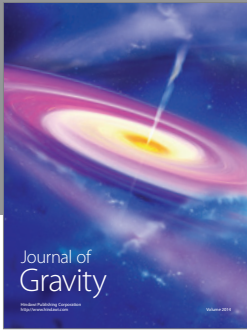
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