

Research Article

Numerical Solutions of Stochastic Differential Equations Driven by Poisson Random Measure with Non-Lipschitz Coefficients

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The numerical methods in the current known literature require the stochastic differential equations (SDEs) driven by Poisson random measure satisfying the global Lipschitz condition and the linear growth condition. In this paper, Euler's method is introduced for SDEs driven by Poisson random measure with non-Lipschitz coefficients which cover more classes of such equations than before. The main aim is to investigate the convergence of the Euler method in probability to such equations with non-Lipschitz coefficients. Numerical example is given to demonstrate our results.

1. Introduction

In finance market and other areas, it is meaningful and significant to model the impact of event-driven uncertainty. Events such as corporate defaults, operational failures, market crashes, or central bank announcements require the introduction of stochastic differential equations (SDEs) driven by Poisson random measure (see [1, 2]), since such equations were initiated in [3].

Actually, we can only obtain the explicit solutions of a small class of SDEs driven by Poisson random measure and so numerical methods are necessary. In general, numerical methods can be divided into strong approximations and weak approximations. Strong approximations focus on pathwise approximations while weak approximations (see [4, 5]) are fit for problems such as derivative pricing.

We give an overview of the results on the strong approximations of stochastic differential equations (SDEs) driven by Poisson random measure in the existing literature. In [6], a convergence result for strong approximations of any given order $\gamma \in \{0.5, 1, 1.5, \dots\}$ was presented. Moreover, N. Bruti-Liberati and E. Platen (see [7]) obtain the jump-adapted order 1.5 scheme, and they also give the derivative-free or implicit jump-adapted schemes with

desired order of strong convergence. And for the specific case of pure jump SDEs, they (see [8]) establish the strong convergence of Taylor's methods under weaker conditions than the currently known. In [5, 7], the drift-implicit schemes which achieve strong order $\gamma \in \{0.5, 1\}$ are given. Recently, Mordecki et al. [9] improved adaptive time stepping algorithms based on a jump augmented Monte Carlo Euler-Maruyama method, which achieve a prescribed precision. M. Wei [10] demonstrates the convergence of numerical solutions for variable delay differential equations driven by Poisson random measure. In [11], the developed Runge-Kutta methods are presented to improve the accuracy behaviour of problems with small noise to SDEs with Poisson random measure.

Clearly, the results above require that the SDEs driven by Poisson random measure satisfy the global Lipschitz condition and the linear growth condition. However, there are many equations which do not satisfy above conditions, and we can see such equations in Section 5 in our paper. Our main contribution is to present Euler's method for these equations with non-Lipschitz coefficients. Here non-Lipschitz coefficients are interpreted in [12], that is to say, the drift coefficients and the diffusion coefficients satisfy the local Lipschitz conditions, the jump coefficients satisfy the global Lipschitz conditions, and the one-sided linear growth condition is considered. Our work is motivated by [12] in which the existence of global solutions for these equations with non-Lipschitz coefficients is proved, while there is no numerical method is presented in our known literature. And our aim in this paper is to close this gap.

Our work is organized as follows. In Section 2, the property of SDEs driven by Poisson random measure with non-Lipschitz coefficients is given. In Section 3, Euler method is analyzed for such equations. In Section 4, we present the convergence in probability of the Euler method. In Section 5, an example is presented.

2. The SDEs Driven by Poisson Random Measure with Non-Lipschitz Coefficients

Throughout this paper, unless specified, we use the following notations. Let $u_1 \vee u_2 = \max\{u_1, u_2\}$ and $u_1 \wedge u_2 = \min\{u_1, u_2\}$. Let $|\cdot|$ and $\langle \cdot, \cdot \rangle$ be the Euclidean norm and the inner product of vectors in $\mathbf{R}^d, d \in \mathbf{N}$. If A is a vector or matrix, its transpose is denoted by A^T . If A is a matrix, its trace norm is denoted by $|A| = \sqrt{\text{trace}(A^T A)}$. Let $L^2_{\mathcal{F}_0}(\Omega; \mathbf{R}^d)$ denote the family of \mathbf{R}^d -valued \mathcal{F}_0 -measurable random variables ξ with $\mathbf{E}|\xi|^2 < \infty$. $[z]$ denotes the largest integer which is less than or equal to z in \mathbf{R} . $I_{\mathcal{A}}$ denotes the indicator function of a set \mathcal{A} .

The following d -dimensional SDE driven by Poisson random measure is considered in our paper:

$$dx(t) = a(x(t-))dt + b(x(t-))dW(t) + \int_{\varepsilon} c(x(t-), v)\tilde{p}_{\phi}(dv \times dt), \quad (2.1)$$

for $t > 0$ with initial condition $x(0-) = x(0) = x_0 \in L^2_{\mathcal{F}_0}(\Omega; \mathbf{R}^d)$, where $x(t-)$ denotes $\lim_{s \rightarrow t-} x(s)$ and $\tilde{p}_{\phi}(dv \times dt) := p_{\phi}(dv \times dt) - \phi(dv)dt$.

The drift coefficient $a : \mathbf{R}^d \rightarrow \mathbf{R}^d$, the diffusion coefficient $b : \mathbf{R}^d \rightarrow \mathbf{R}^{d \times m}$, and the jump coefficient $c : \mathbf{R}^d \times \varepsilon \rightarrow \mathbf{R}^d$ are assumed to be Borel measurable functions.

The randomness of (2.1) is generated by the following (see [9]). An m -dimensional Wiener process $W = \{W(t) = (W^1(t), \dots, W^m(t))^T\}$ with independent scalar components is defined on a filtered probability space $(\Omega^W, \mathcal{F}^W, (\mathcal{F}_t^W)_{t \geq 0}, \mathbf{P}^W)$. A Poisson random measure

$p_\phi(\omega, dv \times dt)$ is on $\Omega^J \times \varepsilon \times [0, \infty)$, where $\varepsilon \subseteq \mathbf{R}^r \setminus \{0\}$ with $r \in \mathbf{N}$, and its deterministic compensated measure $\phi(dv)dt = \lambda f(v)dvdt$, that is, $\mathbf{E}(p_\phi(dv \times dt)) = \phi(dv)dt$. $f(v)$ is a probability density, and we require finite intensity $\lambda = \phi(\varepsilon) < \infty$. The Poisson random measure is defined on a filtered probability space $(\Omega^J, \mathcal{F}^J, (\mathcal{F}_t^J)_{t \geq 0}, \mathbf{P}^J)$. The Wiener process and the Poisson random measure are mutually independent. The process $x(t)$ is thus defined on a product space $(\Omega, \mathcal{F}, (\mathcal{F}_t)_{t \geq 0}, \mathbf{P})$, where $\Omega = \Omega^W \times \Omega^J$, $\mathcal{F} = \mathcal{F}^W \times \mathcal{F}^J$, $(\mathcal{F}_t)_{t \geq 0} = (\mathcal{F}_t^W)_{t \geq 0} \times (\mathcal{F}_t^J)_{t \geq 0}$, $\mathbf{P} = \mathbf{P}^W \times \mathbf{P}^J$ and \mathcal{F}_0 contains all \mathbf{P} -null sets.

Now, the condition of non-Lipschitz coefficients is given by the following assumptions.

Assumption 2.1. For each integer $k \geq 1$, there exists a positive constant C_k , dependent on k , such that

$$|a(x) - a(y)|^2 \vee |b(x) - b(y)|^2 \leq C_k |x - y|^2, \quad (2.2)$$

for $x, y \in \mathbf{R}^d$ with $|x| \vee |y| \leq k$, $k \in \mathbf{N}$. And there exists a positive constant C such that

$$\int_\varepsilon |c(x, v) - c(y, v)|^2 \phi(dv) \leq C |x - y|^2, \quad (2.3)$$

for $x, y \in \mathbf{R}^d$.

Assumption 2.2. There exists a positive constant L such that

$$2\langle x, a(x) \rangle + |b(x)|^2 + \int_\varepsilon |c(x, v)|^2 \phi(dv) \leq L(1 + |x|^2), \quad (2.4)$$

for $x \in \mathbf{R}^d$.

A unique global solution of (2.1) exists under Assumptions 2.1 and 2.2, see [12].

Assumption 2.3. Consider

$$|a(0)|^2 + |b(0)|^2 + \int_\varepsilon |c(0, v)|^2 \phi(dv) \leq \tilde{L}, \quad \tilde{L} > 0. \quad (2.5)$$

Actually, Assumptions 2.1 and 2.3 imply the linear growth conditions

$$|a(x)|^2 \vee |b(x)|^2 \leq \tilde{C}_k (1 + |x|^2), \quad (2.6)$$

for $x \in \mathbf{R}^d$ with $|x| \leq k$ and $\tilde{C}_k > 0$, and

$$\int_\varepsilon |c(x, v)|^2 \phi(dv) \leq \tilde{C} (1 + |x|^2), \quad (2.7)$$

for $x \in \mathbf{R}^d$ and $\tilde{C} > 0$.

The following result shows that the solution of (2.1) keeps in a compact set with a large probability.

Lemma 2.4. *Under Assumptions 2.1 and 2.2, for any pair of $\epsilon \in (0, 1)$ and $T > 0$, there exists a sufficiently large integer k^* , dependent on ϵ and T , such that*

$$\mathbf{P}(\tau_k \leq T) \leq \epsilon, \quad \forall k \geq k^*, \quad (2.8)$$

where $\tau_k = \inf\{t \geq 0 : |x(t)| \geq k\}$ for $k \geq 1$.

Proof. Using Itô's formula (see [1]) to $|x(t)|^2$, for $t \geq 0$, we have

$$\begin{aligned} |x(t)|^2 &= |x_0|^2 + \int_0^t \left(\langle 2x(s-), a(x(s-)) \rangle + |b(x(s-))|^2 \right) ds \\ &\quad + \int_0^t \int_{\epsilon} \left(|x(s-) + c(x(s-), v)|^2 - |x(s-)|^2 - \langle 2x(s-), c(x(s-), v) \rangle \right) \phi(dv) ds \\ &\quad + \int_0^t \langle 2x(s-), b(x(s-)) \rangle dW(s) + \int_0^t \int_{\epsilon} \left(|x(s-) + c(x(s-), v)|^2 - |x(s-)|^2 \right) \tilde{p}_{\phi}(dv \times ds), \end{aligned} \quad (2.9)$$

which gives

$$\begin{aligned} \mathbf{E}|x(t \wedge \tau_k)|^2 &= \mathbf{E}|x_0|^2 + \mathbf{E} \int_0^{t \wedge \tau_k} \left(\langle 2x(s-), a(x(s-)) \rangle + |b(x(s-))|^2 \right) ds \\ &\quad + \mathbf{E} \int_0^{t \wedge \tau_k} \int_{\epsilon} |c(x(s-), v)|^2 \phi(dv) ds \\ &= \mathbf{E}|x_0|^2 + \mathbf{E} \int_0^{t \wedge \tau_k} \left(\langle 2x(s-), a(x(s-)) \rangle + |b(x(s-))|^2 + \int_{\epsilon} |c(x(s-), v)|^2 \phi(dv) \right) ds, \end{aligned} \quad (2.10)$$

for $t \in [0, T]$. By Assumption 2.2, we thus have

$$\begin{aligned} \mathbf{E}|x(t \wedge \tau_k)|^2 &\leq \mathbf{E}|x_0|^2 + \mathbf{E} \int_0^{t \wedge \tau_k} L(1 + |x(s-)|^2) ds \\ &\leq \mathbf{E}|x_0|^2 + LT + L \int_0^t \mathbf{E}|x(s \wedge \tau_k-)|^2 ds, \end{aligned} \quad (2.11)$$

for $t \in [0, T]$. Consequently by using the Gronwall inequality (see [13]), we obtain

$$\mathbf{E}|x(t \wedge \tau_k)|^2 \leq (\mathbf{E}|x_0|^2 + LT)e^{LT}, \quad (2.12)$$

for $t \in [0, T]$. We therefore get

$$(\mathbf{E}|x_0|^2 + LT)e^{LT} \geq \mathbf{E}|x(T \wedge \tau_k)|^2 \geq \mathbf{E}(|x(\tau_k)|^2 I_{\{\tau_k \leq T\}}) \geq k^2 \mathbf{P}(\tau_k \leq T), \quad (2.13)$$

which means

$$\mathbf{P}(\tau_k \leq T) \leq \frac{e^{LT}}{k^2} (\mathbf{E}|x_0|^2 + LT). \quad (2.14)$$

So for any $\epsilon \in (0, 1)$, we can choose

$$k^* = \left[\sqrt{\frac{e^{LT} \mathbf{E}|x_0|^2 + LT e^{LT}}{\epsilon}} \right] + 1, \quad (2.15)$$

such that

$$\mathbf{P}(\tau_k \leq T) \leq \epsilon, \quad \forall k \geq k^*. \quad (2.16)$$

Hence, we have the result (2.8). \square

3. The Euler Method

In this section, we introduce the Euler method to (2.1) under Assumptions 2.1, 2.2, and 2.3. Subsequently, we give two lemmas to analyze the Euler method over a finite time interval $[0, T]$, where T is a positive number.

Given a step size $\Delta t \in (0, 1)$, the Euler method applied to (2.1) computes approximation $X_n \approx x(t_n)$, where $t_n = n\Delta t$, $n = 0, 1, \dots$, by setting $X_0 = x_0$ and forming

$$X_{n+1} = X_n + a(X_n)\Delta t + b(X_n)\Delta W_n + \int_{t_n}^{t_{n+1}} \int_{\epsilon} c(X_n, v) \tilde{p}_{\phi}(dv \times dt), \quad (3.1)$$

where $\Delta W_n = W(t_{n+1}) - W(t_n)$.

The continuous-time Euler method is then defined by

$$\bar{X}(t) := X_0 + \int_0^t a(Z(s))ds + \int_0^t b(Z(s))dW(s) + \int_0^t \int_{\epsilon} c(Z(s), v) \tilde{p}_{\phi}(dv \times ds), \quad (3.2)$$

where $Z(t) = X_n$ for $t \in [t_n, t_{n+1})$, $n = 0, 1, \dots$

Actually, we can see in [8], $p_\phi = \{p_\phi(t) := p_\phi(\varepsilon \times [0, t])\}$ is a stochastic process that counts the number of jumps until some given time. The Poisson random measure $p_\phi(dv \times dt)$ generates a sequence of pairs $\{(t_i, \xi_i), i \in \{1, 2, \dots, p_\phi(T)\}\}$ for a given finite positive constant T if $\lambda < \infty$. Here $\{t_i : \Omega \rightarrow \mathbf{R}_+, i \in \{1, 2, \dots, p_\phi(T)\}\}$ is a sequence of increasing nonnegative random variables representing the jump times of a standard Poisson process with intensity λ , and $\{\xi_i : \Omega \rightarrow \varepsilon, i \in \{1, 2, \dots, p_\phi(T)\}\}$ is a sequence of independent identically distributed random variables, where ξ_i is distributed according to $\phi(dv)/\phi(\varepsilon)$. Then (3.1) can equivalently be the following form:

$$X_{n+1} = X_n + \left(a(X_n) - \int_\varepsilon c(X_n, v) \phi(dv) \right) \Delta t + b(X_n) \Delta W_n + \sum_{i=p_\phi(t_n)+1}^{p_\phi(t_{n+1})} c(X_n, \xi_i). \quad (3.3)$$

The following lemma shows the close relation between the continuous-time Euler method (3.2) and its step function $Z(t)$.

Lemma 3.1. *Under Assumptions 2.1 and 2.3, for any $T > 0$, there exists a positive constant $K_1(k)$, dependent on k and independent of Δt , such that for all $\Delta t \in (0, 1)$ the continuous-time Euler method (3.2) satisfies*

$$\mathbf{E} \left| \bar{X}(t) - Z(t) \right|^2 \leq K_1(k) \Delta t, \quad (3.4)$$

for $0 \leq t \leq T \wedge \tau_k \wedge \rho_k$, where $\rho_k = \inf\{t \geq 0 : |\bar{X}(t)| \geq k\}$ for $k \geq 1$ and τ_k is defined in Lemma 2.4.

Proof. For $0 \leq t \leq T \wedge \tau_k \wedge \rho_k$, there is an integer n such that $t \in [t_n, t_{n+1})$. So it follows from (3.2) that

$$\bar{X}(t) - Z(t) = X_n + \int_{t_n}^t a(Z(s)) ds + \int_{t_n}^t b(Z(s)) dW(s) + \int_{t_n}^t \int_\varepsilon c(Z(s), v) \tilde{p}_\phi(dv \times ds) - X_n. \quad (3.5)$$

Thus, by taking expectations and using the Cauchy-Schwarz inequality and the martingale properties of $dW(t)$ and $\tilde{p}_\phi(dv \times dt)$, we have

$$\begin{aligned} \mathbf{E} \left| \bar{X}(t) - Z(t) \right|^2 &\leq 3\mathbf{E} \left| \int_{t_n}^t a(Z(s)) ds \right|^2 + 3\mathbf{E} \left| \int_{t_n}^t b(Z(s)) dW(s) \right|^2 \\ &\quad + 3\mathbf{E} \left| \int_{t_n}^t \int_\varepsilon c(Z(s), v) \tilde{p}_\phi(dv \times ds) \right|^2 \end{aligned}$$

$$\begin{aligned}
&\leq 3\mathbf{E}\left(\int_{t_n}^t 1^2 ds \int_{t_n}^t |a(Z(s))|^2 ds\right) + 3\mathbf{E}\int_{t_n}^t |b(Z(s))|^2 ds \\
&\quad + 3\mathbf{E}\int_{t_n}^t \int_{\varepsilon} |c(Z(s), v)|^2 \phi(dv) ds \\
&\leq 3\Delta t \mathbf{E}\int_{t_n}^t |a(Z(s))|^2 ds + 3\mathbf{E}\int_{t_n}^t |b(Z(s))|^2 ds + 3\mathbf{E}\int_{t_n}^t \int_{\varepsilon} |c(Z(s), v)|^2 \phi(dv) ds,
\end{aligned} \tag{3.6}$$

where the inequality $|u_1 + u_2 + u_3|^2 \leq 3|u_1|^2 + 3|u_2|^2 + 3|u_3|^2$ for $u_1, u_2, u_3 \in \mathbf{R}^d$ is used. Therefore, by applying Assumptions 2.1 and 2.3, we get

$$\begin{aligned}
\mathbf{E}\int_{t_n}^t |a(Z(s))|^2 ds &\leq \widetilde{C}_k \mathbf{E}\int_{t_n}^t (1 + |Z(s)|^2) ds \leq \widetilde{C}_k \Delta t + \widetilde{C}_k k^2 \Delta t, \\
\mathbf{E}\int_{t_n}^t |b(Z(s))|^2 ds &\leq \widetilde{C}_k \Delta t + \widetilde{C}_k k^2 \Delta t, \\
\mathbf{E}\int_{t_n}^t \int_{\varepsilon} |c(Z(s), v)|^2 \phi(dv) ds &\leq \widetilde{C} \Delta t + \widetilde{C} k^2 \Delta t,
\end{aligned} \tag{3.7}$$

which lead to

$$\mathbf{E}\left|\overline{X}(t) - Z(t)\right|^2 \leq \Delta t \left(3\widetilde{C}_k \Delta t + 3k^2 \widetilde{C}_k \Delta t + 3\widetilde{C}_k + 3k^2 \widetilde{C}_k + 3\widetilde{C} + 3k^2 \widetilde{C}\right), \tag{3.8}$$

for $t \in [0, T \wedge \tau_k \wedge \rho_k]$. Therefore, we obtain the result (3.4) by choosing

$$K_1(k) = 6\widetilde{C}_k + 6k^2 \widetilde{C}_k + 3\widetilde{C} + 3k^2 \widetilde{C}. \tag{3.9}$$

□

In accord with Lemma 2.4, we give the following lemma which demonstrates that the solution of continuous-time Euler method (3.2) remains in a compact set with a large probability.

Lemma 3.2. *Under Assumptions 2.1, 2.2, and 2.3, for any pair of $\varepsilon \in (0, 1)$ and $T > 0$, there exist a sufficiently large k^* and a sufficiently small Δt_1^* such that*

$$\mathbf{P}(\rho_{k^*} \leq T) \leq \varepsilon, \quad \forall \Delta t \leq \Delta t_1^*, \tag{3.10}$$

where ρ_{k^*} is defined in Lemma 3.1.

Proof. Applying generalized Itô's formula (see [1]) to $|\bar{X}(t)|^2$, for $t \geq 0$, yields

$$\begin{aligned} |\bar{X}(t)|^2 &= |X_0|^2 + \int_0^t \left(\langle 2\bar{X}(s), a(Z(s)) \rangle + |b(Z(s))|^2 \right) ds \\ &\quad + \int_0^t \int_{\varepsilon} \left(|\bar{X}(s) + c(Z(s), v)|^2 - |\bar{X}(s)|^2 - \langle 2\bar{X}(s), c(Z(s), v) \rangle \right) \phi(dv) ds \\ &\quad + \int_0^t \langle 2\bar{X}(s), b(Z(s)) \rangle dW(s) + \int_0^t \int_{\varepsilon} \left(|\bar{X}(s) + c(Z(s), v)|^2 - |\bar{X}(s)|^2 \right) \tilde{p}_{\phi}(dv \times ds). \end{aligned} \quad (3.11)$$

By taking expectations, we thus have

$$\begin{aligned} \mathbf{E}|\bar{X}(t \wedge \rho_k)|^2 &= \mathbf{E}|X_0|^2 + \mathbf{E} \int_0^{t \wedge \rho_k} \left(\langle 2\bar{X}(s), a(Z(s)) \rangle + |b(Z(s))|^2 + \int_{\varepsilon} |c(Z(s), v)|^2 \phi(dv) \right) ds \\ &= \mathbf{E}|X_0|^2 + \mathbf{E} \int_0^{t \wedge \rho_k} \left(\langle 2\bar{X}(s), a(\bar{X}(s)) \rangle + |b(\bar{X}(s))|^2 \right. \\ &\quad \left. + \int_{\varepsilon} |c(\bar{X}(s), v)|^2 \phi(dv) \right) ds \\ &\quad + \mathbf{E} \int_0^{t \wedge \rho_k} \langle 2\bar{X}(s), a(Z(s)) - a(\bar{X}(s)) \rangle ds \\ &\quad + \mathbf{E} \int_0^{t \wedge \rho_k} \left(|b(Z(s))|^2 - |b(\bar{X}(s))|^2 \right) ds \\ &\quad + \mathbf{E} \int_0^{t \wedge \rho_k} \int_{\varepsilon} \left(|c(Z(s), v)|^2 - |c(\bar{X}(s), v)|^2 \right) \phi(dv) ds. \end{aligned} \quad (3.12)$$

For $t \in [0, T]$. Now, by using the inequalities $\langle u_1, u_2 \rangle \leq |u_1||u_2|$ for $u_1, u_2 \in \mathbf{R}^d$, (2.2) in Assumption 2.1, Fubini's theorem, Cauchy-Schwarz's inequality, and Lemma 3.1, we get

$$\begin{aligned} \mathbf{E} \int_0^{t \wedge \rho_k} \langle 2\bar{X}(s), a(Z(s)) - a(\bar{X}(s)) \rangle ds &\leq 2\mathbf{E} \int_0^{t \wedge \rho_k} |\bar{X}(s)| |a(Z(s)) - a(\bar{X}(s))| ds \\ &\leq 2k\sqrt{C_k} \int_0^t \mathbf{E} |Z(s \wedge \rho_k) - \bar{X}(s \wedge \rho_k)| ds \\ &\leq 2k\sqrt{C_k} \int_0^t \left(\mathbf{E} |Z(s \wedge \rho_k) - \bar{X}(s \wedge \rho_k)|^2 \right)^{1/2} ds \\ &\leq 2kT\sqrt{C_k K_1(k)\Delta t}. \end{aligned} \quad (3.13)$$

And, similarly as above, we have

$$\begin{aligned}
\mathbf{E} \int_0^{t \wedge \rho_k} \left(|b(Z(s))|^2 - |b(\bar{X}(s))|^2 \right) ds &\leq \mathbf{E} \int_0^{t \wedge \rho_k} \left(|b(Z(s))| + |b(\bar{X}(s))| \right) \\
&\quad \times \left(|b(Z(s))| - |b(\bar{X}(s))| \right) ds \\
&\leq 2\sqrt{\tilde{C}_k(1+k^2)} \mathbf{E} \int_0^{t \wedge \rho_k} |b(Z(s)) - b(\bar{X}(s))| ds \\
&\leq 2\sqrt{C_k \tilde{C}_k(1+k^2)} \int_0^t \mathbf{E} |Z(s \wedge \rho_k) - \bar{X}(s \wedge \rho_k)| ds \\
&\leq 2T\sqrt{C_k \tilde{C}_k K_1(k)(1+k^2)} \Delta t. \tag{3.14}
\end{aligned}$$

Moreover, in the same way, we obtain

$$\begin{aligned}
&\mathbf{E} \int_0^{t \wedge \rho_k} \int_{\varepsilon} \left(|c(Z(s), v)|^2 - |c(\bar{X}(s), v)|^2 \right) \phi(dv) ds \\
&= \mathbf{E} \int_0^{t \wedge \rho_k} \int_{\varepsilon} \left(|c(Z(s), v) - c(\bar{X}(s), v) + c(\bar{X}(s), v)|^2 - |c(\bar{X}(s), v)|^2 \right) \phi(dv) ds \\
&\leq \mathbf{E} \int_0^{t \wedge \rho_k} \int_{\varepsilon} \left(2|c(Z(s), v) - c(\bar{X}(s), v)|^2 + |c(\bar{X}(s), v)|^2 \right) \phi(dv) ds \\
&\leq 2C\mathbf{E} \int_0^{t \wedge \rho_k} |Z(s) - \bar{X}(s)|^2 ds + \tilde{C}\mathbf{E} \int_0^{t \wedge \rho_k} \left(1 + |\bar{X}(s)|^2 \right) ds \\
&\leq 2C \int_0^t \mathbf{E} |Z(s \wedge \rho_k) - \bar{X}(s \wedge \rho_k)|^2 ds + \tilde{C}\mathbf{E} \int_0^{t \wedge \rho_k} \left(1 + |\bar{X}(s)|^2 \right) ds \\
&\leq 2CTK_1(k)\Delta t + \tilde{C}T + \tilde{C}\mathbf{E} \int_0^{t \wedge \rho_k} |\bar{X}(s)|^2 ds, \tag{3.15}
\end{aligned}$$

where the inequality $|u_1 + u_2|^2 \leq 2|u_1|^2 + 2|u_2|^2$ for $u_1, u_2 \in \mathbf{R}^d$, (2.3) in Assumptions 2.1 and 2.3, Fubini's theorem, and Lemma 3.1 are used. Subsequently, substituting (3.13), (3.14), and (3.15) into (3.12) together with Assumption 2.2 leads to

$$\begin{aligned}
\mathbf{E} |\bar{X}(t \wedge \rho_k)|^2 &\leq \mathbf{E} |X_0|^2 + LE \int_0^{t \wedge \rho_k} \left(1 + |\bar{X}(s)|^2 \right) ds + \tilde{C}\mathbf{E} \int_0^{t \wedge \rho_k} |\bar{X}(s)|^2 ds \\
&\quad + 2kT\sqrt{C_k K_1(k)\Delta t} + 2T\sqrt{C_k \tilde{C}_k K_1(k)(1+k^2)\Delta t} + 2CTK_1(k)\Delta t + \tilde{C}T
\end{aligned}$$

$$\begin{aligned} &\leq (L + \tilde{C}) \int_0^t \mathbf{E} \left| \bar{X}(s \wedge \rho_k) \right|^2 ds + \mathbf{E}|X_0|^2 + LT + \tilde{C}T \\ &\quad + \left(2kT\sqrt{C_k K_1(k)} + 2T\sqrt{C_k \tilde{C}_k K_1(k)(1+k^2)} \right) \sqrt{\Delta t} + 2CTK_1(k)\Delta t, \end{aligned} \quad (3.16)$$

for $0 \leq t \leq T$. Therefore, by the Gronwall inequality (see [13]), for $0 \leq t \leq T$, we get

$$\mathbf{E} \left| \bar{X}(t \wedge \rho_k) \right|^2 \leq \alpha_1 \alpha_4 + \alpha_4 \alpha_2(k) \sqrt{\Delta t} + \alpha_4 \alpha_3(k) \Delta t, \quad (3.17)$$

where

$$\begin{aligned} \alpha_1 &= \mathbf{E}|X_0|^2 + LT + \tilde{C}T, \\ \alpha_2(k) &= 2kT\sqrt{C_k K_1(k)} + 2T\sqrt{C_k \tilde{C}_k K_1(k)(1+k^2)}, \\ \alpha_3(k) &= 2CTK_1(k), \\ \alpha_4 &= \exp(LT + \tilde{C}T). \end{aligned} \quad (3.18)$$

We thus obtain that

$$k^2 \mathbf{P}(\rho_k \leq T) \leq \mathbf{E} \left(\left| \bar{X}(\rho_k) \right|^2 I_{\{\rho_k \leq T\}} \right) \leq \mathbf{E} \left| \bar{X}(T \wedge \rho_k) \right|^2 \leq \alpha_1 \alpha_4 + \alpha_4 \alpha_2(k) \sqrt{\Delta t} + \alpha_4 \alpha_3(k) \Delta t. \quad (3.19)$$

So for any $\epsilon \in (0, 1)$, we can choose sufficiently large integer $k = k^*$ such that

$$\frac{\alpha_1 \alpha_4}{k^{*2}} \leq \frac{\epsilon}{2}, \quad (3.20)$$

and choose sufficiently small $\Delta t_1^* \in (0, 1)$ such that

$$\frac{\alpha_4 \alpha_2(k^*) \sqrt{\Delta t_1^*} + \alpha_4 \alpha_3(k^*) \Delta t_1^*}{k^{*2}} \leq \frac{\epsilon}{2}. \quad (3.21)$$

Hence, we have

$$\mathbf{P}(\rho_{k^*} \leq T) \leq \epsilon, \quad \forall \Delta t \leq \Delta t_1^*. \quad (3.22)$$

□

4. Convergence in Probability

In this section, we present two convergence theorems of the Euler method to the SDE with Poisson random measure (2.1) over a finite time interval $[0, T]$.

At the beginning, we give a lemma based on Lemma 3.1.

Lemma 4.1. *Under Assumptions 2.1 and 2.3, for any $T > 0$, there exists a positive constant $K_2(k)$, dependent on k and independent of Δt , such that for all $\Delta t \in (0, 1)$ the solution of (2.1) and the continuous-time Euler method (3.2) satisfy*

$$\mathbf{E} \left(\sup_{0 \leq t \leq T} \left| x(t \wedge \tau_k \wedge \rho_k) - \bar{X}(t \wedge \tau_k \wedge \rho_k) \right|^2 \right) \leq K_2(k) \Delta t, \quad (4.1)$$

where τ_k and ρ_k are defined in Lemmas 2.4 and 3.1, respectively.

Proof. From (2.1) and (3.2), for any $0 \leq t' \leq T$, we have

$$\begin{aligned} & \mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| x(t \wedge \tau_k \wedge \rho_k) - \bar{X}(t \wedge \tau_k \wedge \rho_k) \right|^2 \right) \\ & \leq 3 \mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| \int_0^{t \wedge \tau_k \wedge \rho_k} (a(x(s-)) - a(Z(s))) ds \right|^2 \right) \\ & \quad + 3 \mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| \int_0^{t \wedge \tau_k \wedge \rho_k} (b(x(s-)) - b(Z(s))) dW(s) \right|^2 \right) \\ & \quad + 3 \mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| \int_0^{t \wedge \tau_k \wedge \rho_k} \int_{\varepsilon} (c(x(s-), v) - c(Z(s), v)) \tilde{p}_{\phi}(dv \times ds) \right|^2 \right), \end{aligned} \quad (4.2)$$

where the inequality $|u_1 + u_2 + u_3|^2 \leq 3|u_1|^2 + 3|u_2|^2 + 3|u_3|^2$ for $u_1, u_2, u_3 \in \mathbf{R}^d$ is used. Therefore, by using the Cauchy-Schwarz inequality, (2.2) in Assumption 2.1, Lemma 3.1 and Fubini's theorem, we obtain

$$\begin{aligned} & \mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| \int_0^{t \wedge \tau_k \wedge \rho_k} (a(x(s-)) - a(Z(s))) ds \right|^2 \right) \\ & \leq \mathbf{E} \left(\sup_{0 \leq t \leq t'} \int_0^{t \wedge \tau_k \wedge \rho_k} 1^2 ds \int_0^{t \wedge \tau_k \wedge \rho_k} |a(x(s-)) - a(Z(s))|^2 ds \right) \\ & \leq T \mathbf{E} \left(\int_0^{t' \wedge \tau_k \wedge \rho_k} |a(x(s-)) - a(Z(s))|^2 ds \right) \end{aligned}$$

$$\begin{aligned}
&\leq 2TC_k \mathbf{E} \left(\int_0^{t' \wedge \tau_k \wedge \rho_k} |\bar{X}(s) - Z(s)|^2 ds \right) + 2TC_k \mathbf{E} \left(\int_0^{t' \wedge \tau_k \wedge \rho_k} |x(s-) - \bar{X}(s)|^2 ds \right) \\
&\leq 2TC_k \int_0^{t'} \mathbf{E} |\bar{X}(s \wedge \tau_k \wedge \rho_k) - Z(s \wedge \tau_k \wedge \rho_k)|^2 ds \\
&\quad + 2TC_k \int_0^{t'} \mathbf{E} |x(s \wedge \tau_k \wedge \rho_k-) - \bar{X}(s \wedge \tau_k \wedge \rho_k)|^2 ds \\
&\leq 2T^2 C_k K_1(k) \Delta t + 2TC_k \int_0^{t'} \mathbf{E} \left(\sup_{0 \leq u \leq s} |x(u \wedge \tau_k \wedge \rho_k-) - \bar{X}(u \wedge \tau_k \wedge \rho_k)|^2 \right) ds. \quad (4.3)
\end{aligned}$$

Moreover, by using the martingale properties of $dW(t)$ and $\tilde{p}_\phi(dv \times dt)$, Assumption 2.1, Lemma 3.1, and Fubini's theorem, we have

$$\begin{aligned}
&\mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| \int_0^{t \wedge \tau_k \wedge \rho_k} (b(x(s-)) - b(Z(s))) dW(s) \right|^2 \right) \\
&\leq 4\mathbf{E} \int_0^{t' \wedge \tau_k \wedge \rho_k} |b(x(s-)) - b(Z(s))|^2 ds \\
&\leq 8C_k \mathbf{E} \int_0^{t' \wedge \tau_k \wedge \rho_k} |\bar{X}(s) - Z(s)|^2 ds + 8C_k \mathbf{E} \int_0^{t' \wedge \tau_k \wedge \rho_k} |x(s-) - \bar{X}(s)|^2 ds \\
&\leq 8C_k \int_0^{t'} \mathbf{E} |\bar{X}(s \wedge \tau_k \wedge \rho_k) - Z(s \wedge \tau_k \wedge \rho_k)|^2 ds \\
&\quad + 8C_k \int_0^{t'} \mathbf{E} |x(s \wedge \tau_k \wedge \rho_k-) - \bar{X}(s \wedge \tau_k \wedge \rho_k)|^2 ds \\
&\leq 8TC_k K_1(k) \Delta t + 8C_k \int_0^{t'} \mathbf{E} \left(\sup_{0 \leq u \leq s} |x(u \wedge \tau_k \wedge \rho_k-) - \bar{X}(u \wedge \tau_k \wedge \rho_k)|^2 \right) ds, \\
&\mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| \int_0^{t \wedge \tau_k \wedge \rho_k} \int_\varepsilon (c(x(s-), v) - c(Z(s-), v)) \tilde{p}_\phi(dv \times ds) \right|^2 \right) \\
&\leq 4\mathbf{E} \left| \int_0^{t' \wedge \tau_k \wedge \rho_k} \int_\varepsilon (c(x(s-), v) - c(Z(s-), v)) \tilde{p}_\phi(dv \times ds) \right|^2 \\
&= 4\mathbf{E} \int_0^{t' \wedge \tau_k \wedge \rho_k} \int_\varepsilon |c(x(s-), v) - c(Z(s-), v)|^2 \phi(dv) ds \\
&\leq 8TC K_1(k) \Delta t + 8C \int_0^{t'} \mathbf{E} \left(\sup_{0 \leq u \leq s} |x(u \wedge \tau_k \wedge \rho_k-) - \bar{X}(u \wedge \tau_k \wedge \rho_k)|^2 \right) ds. \quad (4.4)
\end{aligned}$$

Hence, by substituting (4.3) and (4.4) into (4.2), we get

$$\begin{aligned} & \mathbf{E} \left(\sup_{0 \leq t \leq t'} \left| x(t \wedge \tau_k \wedge \rho_k) - \bar{X}(t \wedge \tau_k \wedge \rho_k) \right|^2 \right) \\ & \leq \Delta t \left(6T^2 C_k K_1(k) + 24TC_k K_1(k) + 24TC K_1(k) \right) + (6TC_k + 24C_k + 24C) \\ & \quad \times \int_0^{t'} \mathbf{E} \left(\sup_{0 \leq u \leq s} \left| x(u \wedge \tau_k \wedge \rho_k) - \bar{X}(u \wedge \tau_k \wedge \rho_k) \right|^2 \right) ds. \end{aligned} \quad (4.5)$$

So using the Gronwall inequality (see [13]), we have the result (4.1) by choosing

$$K_2(k) = \left(6T^2 C_k K_1(k) + 24TC_k K_1(k) + 24TC K_1(k) \right) \exp \left(6T^2 C_k + 24TC_k + 24TC \right). \quad (4.6)$$

□

Now, let's state our theorem which demonstrates the convergence in probability of the continuous-time Euler method (3.2).

Theorem 4.2. *Under Assumptions 2.1, 2.2, and 2.3, for sufficiently small $\epsilon, \varsigma \in (0, 1)$, there is a Δt^* such that for all $\Delta t < \Delta t^*$*

$$\mathbf{P} \left(\sup_{0 \leq t \leq T} \left| x(t) - \bar{X}(t) \right|^2 \geq \varsigma \right) \leq \epsilon, \quad (4.7)$$

for any $T > 0$.

Proof. For sufficiently small $\epsilon, \varsigma \in (0, 1)$, we define

$$\bar{\Omega} = \left\{ \omega : \sup_{0 \leq t \leq T} \left| x(t) - \bar{X}(t) \right|^2 \geq \varsigma \right\}. \quad (4.8)$$

According to Lemmas 2.4 and 3.2, there exists a pair of k^* and Δt_1^* such that

$$\begin{aligned} \mathbf{P}(\tau_{k^*} \leq T) & \leq \frac{\epsilon}{3}, \\ \mathbf{P}(\rho_{k^*} \leq T) & \leq \frac{\epsilon}{3}, \quad \forall \Delta t \leq \Delta t_1^*. \end{aligned} \quad (4.9)$$

We thus have

$$\begin{aligned}
\mathbf{P}(\overline{\Omega}) &\leq \mathbf{P}(\overline{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) + \mathbf{P}(\tau_{k^*} \wedge \rho_{k^*} \leq T) \\
&\leq \mathbf{P}(\overline{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) + \mathbf{P}(\tau_{k^*} \leq T) + \mathbf{P}(\rho_{k^*} \leq T) \\
&\leq \mathbf{P}(\overline{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) + \frac{2\epsilon}{3},
\end{aligned} \tag{4.10}$$

for $\Delta t \leq \Delta t_1^*$. Moreover, according to Lemma 4.1, we have

$$\begin{aligned}
\varsigma \mathbf{P}(\overline{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) &\leq \mathbf{E} \left(I_{\{\tau_{k^*} \wedge \rho_{k^*} > T\}} \sup_{0 \leq t \leq T} |x(t) - \overline{X}(t)|^2 \right) \\
&\leq \mathbf{E} \left(\sup_{0 \leq t \leq T} |x(t \wedge \tau_{k^*} \wedge \rho_{k^*}) - \overline{X}(t \wedge \tau_{k^*} \wedge \rho_{k^*})|^2 \right) \\
&\leq K_2(k^*) \Delta t,
\end{aligned} \tag{4.11}$$

which leads to

$$\mathbf{P}(\overline{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) \leq \frac{\epsilon}{3}, \tag{4.12}$$

for $\Delta t \leq \Delta t_2^*$. Therefore, from the inequalities above, we obtain

$$\mathbf{P}(\overline{\Omega}) \leq \epsilon, \tag{4.13}$$

for $\Delta t \leq \Delta t^*$, where $\Delta t^* = \min\{\Delta t_1^*, \Delta t_2^*\}$. \square

We remark that the continuous-time Euler solution $\overline{X}(t)$ (3.2) cannot be computed, since it requires knowledge of the entire Brownian motion and Poisson random measure paths, not just only their Δt -increments. Therefore, the last theorem shows the convergence in probability of the discrete Euler solution (3.1).

Theorem 4.3. *Under Assumptions 2.1, 2.2, and 2.3, for sufficiently small $\epsilon, \varsigma \in (0, 1)$, there is a Δt^* such that for all $\Delta t < \Delta t^*$*

$$\mathbf{P}(|x(t) - Z(t)|^2 \geq \varsigma, 0 \leq t \leq T) \leq \epsilon, \tag{4.14}$$

for any $T > 0$.

Proof. For sufficiently small $\epsilon, \varsigma \in (0, 1)$, we define

$$\tilde{\Omega} = \left\{ \omega : |x(t) - Z(t)|^2 \geq \varsigma, 0 \leq t \leq T \right\}. \tag{4.15}$$

A similar analysis as Theorem 4.2 gives

$$\mathbf{P}(\tilde{\Omega}) \leq \mathbf{P}(\tilde{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) + \frac{2\epsilon}{3}. \tag{4.16}$$

Recalling that

$$\begin{aligned} \zeta \mathbf{P}(\tilde{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) &\leq \mathbf{E}(|x(T) - Z(T)|^2 I_{\{\tau_{k^*} \wedge \rho_{k^*} > T\}}) \\ &\leq \mathbf{E}|x(T \wedge \tau_{k^*} \wedge \rho_{k^*}) - Z(T \wedge \tau_{k^*} \wedge \rho_{k^*})|^2 \\ &\leq 2\mathbf{E}\left(\sup_{0 \leq t \leq T} |x(t \wedge \tau_{k^*} \wedge \rho_{k^*}) - \bar{X}(t \wedge \tau_{k^*} \wedge \rho_{k^*})|^2\right) \\ &\quad + 2\mathbf{E}|\bar{X}(T \wedge \tau_{k^*} \wedge \rho_{k^*}) - Z(T \wedge \tau_{k^*} \wedge \rho_{k^*})|^2 \\ &\leq 2K_1(k^*)\Delta t + 2K_2(k^*)\Delta t, \end{aligned} \tag{4.17}$$

and using Lemmas 3.1 and 4.1, we get that

$$\mathbf{P}(\tilde{\Omega} \cap \{\tau_{k^*} \wedge \rho_{k^*} > T\}) \leq \frac{\epsilon}{3}, \tag{4.18}$$

for sufficiently small Δt . Consequently, the inequalities above show that

$$\mathbf{P}(\tilde{\Omega}) \leq \epsilon, \tag{4.19}$$

for all sufficiently small Δt .

So we complete the result (4.14). □

5. Numerical Example

In this section, a numerical example is analyzed under Assumptions 2.1, 2.2, and 2.3 which cover more classes of SDEs driven by Poisson random measure.

Now, we consider the following equation:

$$dx(t) = a(x(t-))dt + b(x(t-))dW(t) + \int_{\epsilon} c(x(t-), v)\tilde{p}_{\phi}(dv \times dt), \quad t > 0, \tag{5.1}$$

with $x(0) = x(0-) = 0$, where $d = m = r = 1$. The coefficients of this equation have the form

$$a(x) = \frac{1}{2}(x - x^3), \quad b(x) = x^2, \quad c(x, v) = xv. \tag{5.2}$$

The compensated measure of the Poisson random measure $p_\phi(dv \times dt)$ is given by $\phi(dv)dt = \lambda f(v)dvdt$, where $\lambda = 5$ and

$$f(v) = \frac{1}{\sqrt{2\pi v}} \exp\left(-\frac{(\ln v)^2}{2}\right), \quad 0 \leq v < \infty \quad (5.3)$$

is the density function of a lognormal random variable.

Clearly, the equation cannot satisfy the global Lipschitz conditions and the linear growth conditions. On the other hand, we have

$$\begin{aligned} 2\langle x, a(x) \rangle + |b(x)|^2 + \int_{\varepsilon} |c(x, v)|^2 \phi(dv) &= x(x - x^3) + x^4 + \int_{\varepsilon} x^2 v^2 \lambda \frac{1}{\sqrt{2\pi v}} \exp\left(-\frac{(\ln v)^2}{2}\right) dv \\ &\leq (1 + 5e^2)(1 + x^2), \end{aligned} \quad (5.4)$$

that is to say, Assumptions 2.1, 2.2, and 2.3 in Section 2 are satisfied. Therefore, Albeverio et al. [12] guarantee that (5.1) has a unique global solution on $[0, \infty)$.

Given the stepsize Δt , we can have the Euler method

$$X_{n+1} = X_n + \frac{1}{2}(X_n - X_n^3)\Delta t + X_n^2 \Delta W_n + X_n \int_{t_n}^{t_{n+1}} \int_{\varepsilon} v \tilde{p}_\phi(dv \times dt), \quad (5.5)$$

with $X_0 = 0$.

And in Matlab experiment, each discretized trajectory is actually given in detail by the following.

Algorithm

Simulate $X_{n+1}^- := X_n + (1/2)(X_n - X_n^3)\Delta t + X_n^2 \Delta W_n$;

Simulate variable $p_\phi(t_{n+1}) - p_\phi(t_n)$, where $p_\phi(t_n)$ is from Poisson distribution with parameter λt_n ;

Simulate $p_\phi(t_{n+1}) - p_\phi(t_n)$ independent random variables t_i uniformly distributed on the interval $[p_\phi(t_n), p_\phi(t_{n+1})]$;

Simulate $p_\phi(t_{n+1}) - p_\phi(t_n)$ independent random variables ξ_i with law $f(v)$;

obtain $X_{n+1} = X_{n+1}^- + X_n \sum_{i=p_\phi(t_n)+1}^{p_\phi(t_{n+1})} I_{t_n \leq t_i < t_{n+1}} \xi_i$.

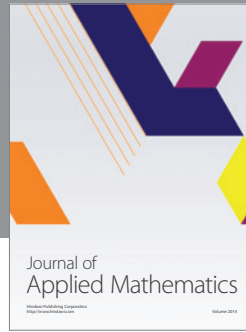
Subsequently, we can get the results in Theorems 4.2 and 4.3.

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