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# Exact Displacement Analysis of Four-Link Spatial Mechanisms by the Direction Cosine Matrix Method 


#### Abstract

A method of displacement analysis of the four-link spatial mechanism is developed. The results through this analysis will be exact solutions that can be obtained without resorting to numerical or iteration schemes. In the analysis, the position of a link in a mechanism can be fully defined if its direction and length are known. Therefore, this analysis involves the calculation of the unknown direction cosines and length of each link for a given configuration of the mechanism. In finding the direction cosines of the unknown unit vectors involved for each link and rotating axis, two types of coordinates, the global and the local, are generally used. Then, a direction cosine matrix between each local coordinate system and the global coordinates is established. Thus, the unknown direction cosines of the local coordinates, the links, and the rotating axes are obtained in global coordinates. In this development, direction cosine matrices are used throughout the analysis. As an illustration, the application of this method to the study of four-link spatial mechanisms, $R G G R$, $R G C R, R R G G$, and $R R G C$ will be presented.


## Introduction

A recent survey of space mechanism research [1], which covers analytical methods developed mainly since the 1950s with numerous pertinent references, serves as an extensive and informative source of background material. However, several selective references of well-known methods for the displacement analysis of spatial linkage may be mentioned. Among them are the $4 \times 4$ matrix iterative method [ 2,3 ], the dual number quarternion method $[4,5]$, the geometric transformation method [6, 7], the vector method [8-10], the screw method [11-15], the tensor method [16], the line geometric method [17], and the geometrical configuration method [18], etc. Most of these methods involve high level mathematics of complicated mathematical manipulation, and all require numerical or iterative schemes for solutions.

A method of displacement analysis using direction cosine matrices as transformation matrices for the four-link spatial mechanisms is developed and applied to various four-bar spatial linkages in this paper. The mathematics involved are elementary; the operations are simple without loss of geometric interpretation, and the solutions are exact. The

[^0]analysis starts by choosing an appropriate local coordinate system and assigns direction cosines to the related unit vectors. These direction cosines in the local coordinates are obtained by applying the dot product of unit vectors and using given angle data. Then, a direction cosine matrix between the global and the local coordinates is obtained, by using known unit vectors, the direction cosines of the local coordinates, and a special property of the direction cosine matrix. Using this special property of the direction cosine matrix, that is, that each element of the matrix equals its own cofactor, we obtain exact analytical solutions without resorting to numerical or iterative schemes. When the direction cosine matrix is known, the unknown unit vectors in the global coordinates can be fully calculated.

Analytical solutions in closed-form, input-output relations for a few spatial four-bar linkages are obtained in [4, 6, 7]. In [6,7], the rotation matrix is used, together with one or two constraints particular to the linkage concerned. However, the solutions for these closed-form, input-output relations will have to be obtained by numerically solving transcendental equations. In the present paper, the direction cosine matrices are used in successive steps of the analysis, from the input end to the output end of the linkage mechanism. In the process, the constraints of the mechanism, such as the constant length of a link or the constant angle between two links, are taken care of automatically. The solutions are obtained without resorting to numerical or iteration schemes. It should be mentioned that, for a simple case of the 2R-2G mechanism, the input-output relation, developed in [20], can be reduced to a closed-form solution.


Fig. 1 A RGGR four-link spatial mechanism


Fig. 2 The schematic diagram of a RGGR mechanism

In this paper, this direction cosine matrix method will be applied to obtain analytical solutions for the four-bar spatial linkages, an RGGR, an RGCR, an RRGG, and an RRGC, with numerical illustration.

## 1 Displacement Analysis of the RGGR Mechanism

The RGGR four-link spatial mechanism as shown in Fig. 1 is a generalization of the planar four-bar mechanism RRRR. It is one of the most versatile and practical configurations of three-dimensional mechanisms and will function as a single degree of freedom linkage with a passive degree of freedom in the connecting link. A schematic diagram of an RGGR mechanism is shown in Fig. 2.

The known quantities of the mechanism are the lengths, $l_{1}$, $l_{2}, l_{3}, l_{4}$, the vector $\bar{l}_{4}$, the directions of rotations, $\hat{p}_{1}, \hat{p}_{2}$, the angles $\zeta, \eta, \alpha, \beta$, from the construction of mechanism, and the input angle $\theta$. The unknown quantities are $\hat{l}_{1}, \hat{l}_{2}$, and $\hat{l}_{3}$.
(a) Input Angle $\theta$. The input angle $\theta$ for the rotation about the $\hat{p}_{1}$-axis can be measured with any arbitrary reference. As shown in Fig. 3, $\theta$ is chosen as the angle between the two planes formed by $\hat{p}_{1}$ and $\hat{l}_{4}$, and $\hat{p}_{1}$ and $\hat{l}_{1}$ in which the $\hat{p}_{1} \hat{l}_{4}$-plane chosen as a reference. Both the angle $\zeta$ between $\hat{p}_{1}$ and $\hat{l}_{1}$, and the angle $\eta$ between $\hat{p}_{1}$ and $\hat{l}_{4}$ are chosen to be less than $\pi$.

Let the local coordinates $x_{1}, y_{1}, z_{1}$ associated with $\hat{p}_{1}$ with the origin at 0 be chosen as follows:

The $x_{1}$-axis is set along the known rotating axis $\hat{p}_{1}$ and has the same positive direction as $\hat{p}_{1}$.


Fig. 3 A local coordinate system and input angle measurement

The $y_{1}$-axis is set in the plane of $\hat{p}_{1}$ and $\hat{l}_{4}$ and the angle between $\hat{l}_{4}$ and the $y$-axis is less than $\pi / 2$.

The $z_{1}$-axis follows the right-hand rule.
With this local coordinate system the input angle $\theta$ can now be measured between the $x_{1} y_{1}$-plane and the $x_{1} l_{1}$-plane with the $x_{1} y_{1}$-plane as references.
(b) Analysis of $\hat{l}_{1}$. The direction cosines of the unit vectors $\hat{p}_{1}, \hat{l}_{1}$, and $\hat{l}_{4}$ in the local coordinates system associated with $\hat{p}_{1}$ are expressed in the parenthesis for each unit vector as $\hat{p}_{1}(1,0,0), \hat{l}_{1}(\cos \zeta, \sin \zeta \cos \theta, \sin \zeta \sin \theta)$, and $\hat{\iota_{4}}\left(a_{1}^{\prime}, a_{2}^{\prime}, 0\right)$, respectively.

To find $\hat{l}_{1}$ in global coordinates, the direction cosine transformation matrix [ $T_{i j}$ ] should be defined. With $\hat{p}_{1}$ and $\hat{l}_{4}$ known in global coordinates we have

$$
\begin{align*}
&\left\{\hat{p}_{1 x}, \hat{p}_{1 y}, \hat{p}_{1 z}\right\}_{\text {global }}^{T}= {\left[T_{i j}\right]_{1}\{1,0,0\}_{\text {local }}^{T} }  \tag{1}\\
& L \rightarrow G
\end{align*}
$$

from which

$$
\begin{equation*}
\left(T_{11}\right)_{1}=\hat{p}_{1 x}, \quad\left(T_{21}\right)_{1}=\hat{p}_{1 y}, \quad\left(T_{31}\right)_{1}=\hat{p}_{1 z} \tag{2}
\end{equation*}
$$

$L \rightarrow G$ underneath $\left[T_{i j}\right]_{1}$ indicates transformation from local to global coordinates. From now on the subscripts global, local, and $L \rightarrow G$ will be omitted for simplicity. Similarly,

$$
\begin{equation*}
\left\{\hat{l}_{4 x}, \hat{l}_{4 y}, \hat{l}_{4 z}\right\}^{T}=\left[T_{i j}\right]_{1}\left\{a_{1}^{\prime}, a_{2}^{\prime}, 0\right\}^{T} \tag{3}
\end{equation*}
$$

in which

$$
\begin{align*}
& a_{1}^{\prime}=\hat{l}_{4} \bullet \hat{p}_{1}=\cos \eta  \tag{4}\\
& a_{2}^{\prime}=\sqrt{1-\left(a_{1}^{\prime}\right)^{2}}=\sin \eta \tag{5}
\end{align*}
$$

where positive sign is taken for the square root as a result of construction of the local coordinate system. Thus

$$
\begin{equation*}
\left\{\hat{l}_{4 x}, \hat{l}_{4 y}, \hat{l}_{4 z}\right\}^{T}=\left[T_{i j}\right]_{1}(\cos \eta, \sin \eta, 0\}^{T} \tag{6}
\end{equation*}
$$

or

$$
\begin{align*}
& \hat{l}_{4 x}=\left(T_{11}\right)_{1} \cos \eta+\left(T_{12}\right)_{1} \sin \eta \\
& \hat{l}_{4 y}=\left(T_{21}\right)_{1} \cos \eta+\left(T_{22}\right)_{1} \sin \eta \\
& \hat{l}_{4 z}=\left(T_{31}\right)_{1} \cos \eta+\left(T_{32}\right)_{1} \sin \eta \tag{7}
\end{align*}
$$

Solving $\left(T_{12}\right)_{1},\left(T_{22}\right)_{1}$, and $\left(T_{32}\right)_{1}$ gives

$$
\begin{align*}
& \left(T_{12}\right)_{1}=\left(\hat{l}_{4 x}-\hat{p}_{1 x} \cos \eta\right) / \sin \eta \\
& \left(T_{22}\right)_{1}=\left(\hat{l}_{4 y}-\hat{p}_{1 y} \cos \eta\right) / \sin \eta \\
& \left(T_{32}\right)_{1}=\left(\hat{l}_{4 z}-\hat{p}_{1 z} \cos \eta\right) / \sin \eta \tag{8}
\end{align*}
$$

Now $\left(T_{13}\right)_{1},\left(T_{23}\right)_{1}$, and $\left(T_{33}\right)_{1}$ can be found as their cofactors. Thus


Fig. 4 A local coordinate system and links $I_{2}$ and $I_{3}$

$$
\begin{align*}
& \left(T_{13}\right)_{1}=\left(\hat{p}_{1 y} \hat{l}_{4 z}-\hat{p}_{1 z} \hat{l}_{4 y}\right) / \sin \eta \\
& \left(T_{23}\right)_{1}=-\left(\hat{p}_{11} \hat{l}_{4 z}-\hat{p}_{1 z} \hat{l}_{4 x}\right) / \sin \eta \\
& \left(T_{33}\right)_{1}=\left(\hat{p}_{1 x} \hat{l}_{4 y}-\hat{p}_{1 y} \hat{l}_{4 x}\right) / \sin \eta \tag{9}
\end{align*}
$$

Therefore
$\left[T_{i j}\right]_{1}$

$$
=\frac{1}{\sin \eta}\left[\begin{array}{lll}
\hat{p}_{1 x} \sin \eta & \hat{l}_{4 x}-\hat{p}_{1 x} \cos \eta & \hat{p}_{1 y} \hat{l}_{4 z}-\hat{p}_{1 z} \hat{l}_{4 y}  \tag{10}\\
\hat{p}_{1 y} \sin \eta & \hat{l}_{4 y}-\hat{p}_{1 y} \cos \eta & \hat{p}_{1 z} \hat{l}_{4 x}-\hat{p}_{1 x} \hat{l}_{4 z} \\
\hat{p}_{1 z} \sin \eta & \hat{l}_{4 z}-\hat{p}_{1 z} \cos \eta & \hat{p}_{1 x} \hat{l}_{4 y}-\hat{p}_{1 y} \hat{l}_{4 x}
\end{array}\right]
$$

Now $\hat{l}_{1}$ in global coordinates can be found as

$$
\begin{equation*}
\left\{\hat{l}_{1 x}, \hat{l}_{1 y}, \hat{l}_{1 z}\right\}^{T}=\left[T_{i j}\right]_{1}\{\cos \zeta, \sin \zeta \cos \theta, \sin \zeta \sin \theta\}^{T} \tag{11}
\end{equation*}
$$

(c) Analysis of $\hat{l}_{3}$. Let vector $\bar{c}$ be defined such that it forms a closed loop with vectors $\bar{l}_{2}$ and $\bar{l}_{3}$ as shown in Fig. 4. It also forms a closed loop with vectors $\bar{l}_{1}$ and $\bar{l}_{4}$. From the figure

$$
\begin{equation*}
\bar{c}=\bar{l}_{1}-\bar{l}_{4} \quad \text { and } \quad \hat{c}=\bar{c} /|\bar{c}| \tag{12}
\end{equation*}
$$

A local coordinate system with the origin at joint $C$ is set such that the $y_{2}$-axis is along the known vector $\hat{p}_{2}$ as shown in Fig. 4. The $x_{2}$-axis is in the plane that consists of known vectors $\bar{c}$ and $\hat{p}_{2}$, perpendicular to the $y_{2}$-axis and the angle between $\hat{c}$ and $x_{2}$-axis is less than $\pi / 2$. The $z$-axis will follow the righthand rule. The direction cosines of the unit vectors $\hat{p}_{2}, \hat{c}$, and $\hat{l_{3}}$ in the local coordinate system are expressed in the parenthesis for each unit vector as $\hat{p}_{2}(0,1,0), \hat{c}\left(b_{1}, b_{2}, 0\right)$, and $\hat{l}_{3}\left(a_{1}, a_{2}, a_{3}\right)$, respectively.

The unknown direction cosines in the local coordinate system will be found by applying the dot product of unit vectors and using the known angles from its design. The direction cosine $b_{2}$ can be obtained from unit vectors $\hat{p}_{2}$ and $\hat{c}$ and the known angle $\gamma$ as follows.
$\hat{p}_{2} \cdot \hat{c}=\hat{p}_{2 x} \hat{c}_{x}+\hat{p}_{2 y} \hat{c}_{y}+\hat{p}_{2 z} \hat{c}_{z}$ in global coordinates
$=\cos \gamma$
$=b_{2}$ in local coordinates
from which

$$
\begin{equation*}
b_{2}=\cos \gamma \tag{14}
\end{equation*}
$$

where the angle $\gamma$ between $\hat{p}_{2}$ and $\hat{c}$ is taken to be less than $\pi$. Therefore

$$
\begin{equation*}
b_{1}=\sqrt{1-b_{2}^{2}}=\sin \gamma \tag{15}
\end{equation*}
$$

where the positive sign is taken for the square root as a result of construction of the local coordinate system.

The direction cosines $a_{1}, a_{2}$, and $a_{3}$ of unit vector $\hat{l}_{3}$ can be obtained as follows. From the dot product of $\hat{l}_{3}$ and $\hat{p}_{2}$, and $\hat{l}_{3}$ and $\hat{c}$ we have

$$
\begin{align*}
& -\hat{l}_{3} \cdot \hat{p}_{2}=-a_{2}=\cos \alpha  \tag{16}\\
& -\hat{l}_{3} \cdot \hat{c}=-a_{1} b_{1}-a_{2} b_{2}=\cos \delta \tag{17}
\end{align*}
$$

where the known angle $\alpha$ between $\hat{p}_{2}$ and $-\hat{l}_{3}$, and the unknown angle $\delta$ between $\hat{c}$ and $-\hat{l}_{3}$ are chosen to be less than $\pi$. By applying the cosine law to $\triangle A B C$

$$
\begin{equation*}
\cos \delta=\left(l_{3}^{2}+c^{2}-l_{2}^{2}\right) / 2 l_{3} c \tag{18}
\end{equation*}
$$

Therefore

$$
\begin{equation*}
a_{1}=-\left(\cos \delta+a_{2} b_{2}\right) / b_{1} \tag{19}
\end{equation*}
$$

in which $b_{1}, b_{2}, a_{2}$, and $\cos \delta$, have been just defined in equations (14)-(16), and (18). The direction cosine $a_{3}$ can now be calculated as

$$
\begin{equation*}
a_{3}= \pm \sqrt{1-a_{1}^{2}-a_{2}^{2}} \tag{20}
\end{equation*}
$$

The positive and negative signs of $a_{3}$ correspond to two possible positions of joint $B$ for the given problem. If $1-a_{1}^{2}$ $-a_{2}^{2}=0$, then the mechanism is not working.

Summarizing, we have

$$
\begin{array}{ll}
b_{1}=\sin \gamma & a_{1}=-(\cos \delta-\cos \alpha \cos \gamma) / \sin \gamma \\
b_{2}=\cos \gamma & a_{2}=-\cos \alpha \\
& a_{3}= \pm \sqrt{1-a_{1}^{2}-a_{2}^{2}} \tag{21}
\end{array}
$$

The transformation matrix from local to global coordinates can now be determined as follows. As

$$
\begin{equation*}
\left\{\hat{p}_{2 x}, \hat{p}_{2 y}, \hat{p}_{2 z}\right\}^{T}=\left[T_{i j}\right]_{2} \quad(0,1,0]^{T} \tag{22}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\{\hat{c}_{x}, \hat{c}_{y}, \hat{c}_{z}\right\}^{T}=\left[T_{i j}\right]_{2} \quad\left\{b_{1}, b_{2}, 0\right\}^{T} \tag{23}
\end{equation*}
$$

from which

$$
\begin{equation*}
\left(T_{12}\right)_{2}=\hat{p}_{2 x},\left(T_{22}\right)_{2}=\hat{p}_{2 y},\left(T_{32}\right)_{2}=\hat{p}_{2 z} \tag{24}
\end{equation*}
$$

and

$$
\begin{align*}
& \left(T_{11}\right)_{2}=\left[\hat{c}_{x}-\left(T_{12}\right)_{2} b_{2}\right] / b_{1} \\
& \left(T_{21}\right)_{2}=\left[\hat{c}_{y}-\left(T_{22}\right)_{2} b_{2}\right] / b_{1}  \tag{25}\\
& \left(T_{31}\right)_{2}=\left[\hat{c}_{z}-\left(T_{32}\right)_{2} b_{2}\right] / b_{1}
\end{align*}
$$

Since each element in a direction cosine matrix is equal to its own cofactor, therefore

$$
\begin{align*}
& \left(T_{13}\right)_{2}=\left(T_{21}\right)_{2}\left(T_{32}\right)_{2}-\left(T_{22}\right)_{2}\left(T_{31}\right)_{2} \\
& \left(T_{23}\right)_{2}=\left(T_{12}\right)_{2}\left(T_{31}\right)_{2}-\left(T_{11}\right)_{2}\left(T_{32}\right)_{2}  \tag{26}\\
& \left(T_{33}\right)_{2}=\left(T_{11}\right)_{2}\left(T_{22}\right)_{2}-\left(T_{12}\right)_{2}\left(T_{21}\right)_{2}
\end{align*}
$$

The resulting transformation matrix is
$\left[T_{i j}\right]_{2}=\frac{1}{\sin \gamma}\left[\begin{array}{lll}\hat{c}_{x}-\hat{p}_{2 x} \cos \gamma & \hat{p}_{2 x} \sin \gamma & \hat{p}_{2 z} \hat{c}_{y}-\hat{p}_{2 y} \hat{c}_{z z} \\ \hat{c}_{y}-\hat{p}_{2 y} \cos \gamma & \hat{p}_{2 y} \sin \gamma & \hat{p}_{2 x} c_{z}-\hat{p}_{2 z} \hat{c}_{x} \\ \hat{c}_{z}-\hat{p}_{2 z} \cos \gamma & \hat{p}_{2 z} \sin \gamma & \hat{p}_{2 y} \hat{c}_{x}-\hat{p}_{2 x} \hat{c}_{y}\end{array}\right]$

Now the unit vector $\hat{l}_{3}$ in global coordinates can be obtained from

$$
\begin{equation*}
\left\{\hat{l}_{3 x}, \hat{l}_{3 y}, \hat{l}_{3 z}\right\}^{T}=\left[T_{i j}\right]_{2}\left(a_{1}, a_{2}, a_{3}\right\}^{T} \tag{28}
\end{equation*}
$$

and the vector $\bar{l}_{3}$ is

$$
\begin{equation*}
\bar{l}_{3}=l_{3} \hat{l}_{3} \tag{29}
\end{equation*}
$$

Note that
if $\alpha+\beta>\gamma, \quad l_{3}$ has two positions
if $\alpha+\beta=\gamma, \quad l_{3}$ has one position
if $\alpha+\beta<\gamma$, it is an impossible case


Fig. 5 A local coordinate system and output angle measurement
(d) Analysis of $\hat{l}_{2}$. Let us consider the equation

$$
\begin{equation*}
\bar{l}_{2}+\bar{l}_{3}+\bar{c}=0 \tag{30}
\end{equation*}
$$

or its scalar form

$$
\begin{align*}
& l_{2} \hat{l}_{2 x}+l_{3} \hat{l}_{3 x}+c \hat{c}_{x}=0 \\
& l_{2} \hat{l}_{2 y}+l_{3} \hat{l}_{3 y}+c \hat{c}_{y}=0  \tag{31}\\
& l_{2} \hat{l}_{2 z}+l_{3} \hat{l}_{3 z}+c \hat{c}_{z}=0
\end{align*}
$$

From these equations the components of unit vector $\hat{l}_{2}$ can be solved as

$$
\begin{align*}
& \hat{l}_{2 x}=-\left(l_{3} \hat{l}_{3 x}+c \hat{c}_{x}\right) / l_{2} \\
& \hat{l}_{2 y}=-\left(l_{3} \hat{l}_{3 y}+c \hat{c}_{y}\right) / l_{2}  \tag{32}\\
& \hat{l}_{2 z}=-\left(l_{3} \hat{l}_{3 z}+c \hat{c}_{z}\right) / l_{2}
\end{align*}
$$

Therefore

$$
\begin{equation*}
\bar{l}_{2}=l_{2} \hat{l}_{2} \tag{33}
\end{equation*}
$$

(e) Output Angle $\phi$. The output angle $\phi$ for the rotation about the $\hat{p}_{2}$-axis can be measured with any arbitrary reference. As shown in Fig. 5, $\phi$ is chosen as the angle between the two planes formed by $\hat{p}_{2}$ and $\hat{l}_{4}$, and $\hat{p}_{2}$ and $\hat{l}_{3}$, and $\hat{p}_{2} \hat{l}_{4}$ plane is chosen as a reference. Both the angle $\alpha$ between $\hat{p}_{2}$ and $-\hat{l}_{3}$ and the angle $\beta$ between $\hat{p}_{2}$ and $-\hat{l}_{4}$ are chosen to be less than $\pi$ and they are known angles from the mechanism.

Let the local coordinates $x_{3}, y_{3}$, and $z_{3}$ with the origin at $C$ be chosen as follows:
$x_{3}$-axis is along the known rotating axis $\hat{p}_{2}$ and has the same positive direction as $\hat{p}_{2}$.
$y_{3}$-axis is in the plane of $\hat{p}_{2}$ and $\hat{l}_{4}$ and the angle between $-\hat{l}_{4}$ and the $y$-axis is less than $\pi / 2$.
$z_{3}$-axis follows the right-hand rule.
It is seen that the angle $\phi$ can now be measured between $x_{3} y_{3}$ plane and $x_{3} l_{3}$-plane with $x_{3} y_{3}$-plane as reference.

The direction cosines of the unit vectors $\hat{p}_{2}, \hat{l}_{3}$, and $\hat{l}_{4}$ in the local coordinate system are expressed in the parenthesis for each unit vector as $\hat{p}_{2}(1,0,0), \hat{l}_{3}(-\cos \alpha,-\sin \alpha \cos \phi,-\sin \alpha$ $\sin \phi$ ), and ( $a_{1}^{\prime \prime}, a_{2}^{\prime \prime}, 0$ ), respectively.

To find $\phi$, the same procedure used to find $\hat{l}_{1}$ will be followed. The transformation matrix $\left[T_{i j}\right]_{3}$ is obtained as

$$
\begin{align*}
& {\left[T_{i j}\right]_{3}} \\
& =\frac{1}{\sin \beta}\left[\begin{array}{ccc}
\hat{p}_{2 x} \sin \beta & -\hat{l}_{4 x}-\hat{p}_{2 x} \cos \beta & -\hat{p}_{2 y} \hat{l}_{4 z}+\hat{p}_{2 z} \hat{l}_{4 y} \\
\hat{p}_{2 y} \sin \beta & -\hat{l}_{4 y}-\hat{p}_{2 y} \cos \beta & -\hat{p}_{2 z} \hat{l}_{4 x}+\hat{p}_{2 x} \hat{l}_{4 z} \\
\hat{p}_{2 z} \sin \beta & -\hat{l}_{4 z}-\hat{p}_{2 z} \cos \beta & -\hat{p}_{2 x} \hat{l}_{4 y}+\hat{p}_{2 y} \hat{l}_{4 x}
\end{array}\right] \tag{34}
\end{align*}
$$

Therefore

$$
\begin{align*}
\begin{array}{l}
\left\{\hat{l}_{3 x}, \hat{l}_{3 y}, \hat{l}_{3 z}\right\}^{T}
\end{array} \\
\quad=-\left[T_{i j}\right]_{3}\{\cos \alpha, \sin \alpha \cos \phi, \sin \alpha \sin \phi\}^{T} \tag{35}
\end{align*}
$$

This equation in $\hat{l}_{3}$ can be expanded into

$$
\begin{array}{r}
\hat{l}_{3 x}=-\left[\left(T_{11}\right)_{3} \cos \alpha+\left(T_{12}\right)_{3} \sin \alpha \cos \phi\right. \\
\left.+\left(T_{13}\right)_{3} \sin \alpha \sin \phi\right] \\
\hat{l}_{3 y}=-\left[\left(T_{21}\right)_{3} \cos \alpha+\left(T_{22}\right)_{3} \sin \alpha \cos \phi\right. \\
\left.+\left(T_{23}\right)_{3} \sin \alpha \sin \phi\right]  \tag{36}\\
\hat{l}_{3 z}=-\left[\left(T_{31}\right)_{3} \cos \alpha+\left(T_{32}\right)_{3} \sin \alpha \cos \phi\right. \\
\left.+\left(T_{33}\right)_{3} \sin \alpha \sin \phi\right]
\end{array}
$$

From any two of these three equations $\sin \phi$ and $\cos \phi$ can be solved. If the first two equations are chosen, we obtain $\sin \phi=$
$\frac{\left[-\left(T_{22}\right)_{3} \hat{l}_{3 x}+\left(T_{12}\right)_{3} \hat{l}_{3 y}\right]-\left[\left(T_{11}\right)_{3}\left(T_{22}\right)_{3}-\left(T_{12}\right)_{3}\left(T_{21}\right)_{3}\right] \cos \alpha}{\left[\left(T_{13}\right)_{3}\left(T_{22}\right)_{3}-\left(T_{12}\right)_{3}\left(T_{23}\right)_{3}\right] \sin \alpha}$
$\cos \phi$
$=\frac{\left[-\left(T_{23}\right)_{3} \hat{l}_{3 x}-\left(T_{13}\right)_{3} \hat{l}_{3 y}\right]-\left[\left(T_{11}\right)_{3}\left(T_{23}\right)_{3}-\left(T_{13}\right)_{3}\left(T_{21}\right)_{3}\right] \cos \alpha}{\left[\left(T_{12}\right)_{3}\left(T_{23}\right)_{3}-\left(T_{13}\right)_{3}\left(T_{22}\right)_{3}\right] \sin \alpha}$

These two equations can be simplified as
$\sin \phi=\frac{\left(T_{22}\right)_{3} \hat{l}_{3 x}-\left(T_{12}\right)_{3} \hat{l}_{3 y}+\left(T_{33}\right)_{3} \cos \alpha}{\left(T_{31}\right)_{3} \sin \alpha}$
$\cos \phi=\frac{-\left(T_{23}\right)_{3} \hat{l}_{3 x}+\left(T_{13}\right)_{3} \hat{l}_{3 y}+\left(T_{32}\right)_{3} \cos \alpha}{\left(T_{31}\right)_{3} \sin \alpha}$
From these two equations the angle $\phi$ can be completely determined.

## 2 Displacement Analysis of the RGCR Mechanism

The RGCR four-link spatial mechanism whose schematic diagram is shown in Fig. 6 is similar to the RGGR mechanism except that one of the spherical joints is replaced by a cylindrical joint.

The known data of the mechanism for the analysis are the lengths $l_{1}, l_{2}$, the vector $\bar{l}_{4}$, the directions of rotation of $\hat{p}_{1}$, $\hat{p}_{2}$, the angles $\zeta, \eta, \alpha, \beta$, and the input angle $\theta$. The unknown quantities are $\hat{l}_{1}, \hat{l}_{2}$, and $\hat{l}_{3}$. The angles, $\zeta$ between $\hat{p}_{1}$ and $\hat{l}_{1}, \eta$ between $\hat{p}_{1}$ and $\hat{l}_{4}, \alpha$ between $\hat{p}_{2}$ and $\hat{l}_{3}$, and $\beta$ between $\hat{p}_{2}$ and $\hat{l_{4}}$ are chosen to be less than $\pi$ and they are known angles from the construction of the mechanism.
In the analysis, the only difference from the RGGR is the calculation of $\bar{l}_{3}$. The analysis of $\hat{l}_{3}$ is as follows. From Fig. 6,

$$
\begin{equation*}
\bar{c}=\bar{l}_{1}-\bar{l}_{4} \quad \text { and } \hat{c}=\bar{c} /|\bar{c}| \tag{41}
\end{equation*}
$$

By applying cosine law to $\triangle A B C$ we obtain

$$
\begin{equation*}
c^{2}=l_{2}^{2}+l_{3}^{2}-2 l_{2} l_{3} \cos \psi \tag{42}
\end{equation*}
$$



Fig. 6 The schematic diagram of a RGGR four-line spatial mechanism and its known data

To solve $l_{3}$ the preceding equation is rewritten as

$$
\begin{equation*}
l_{3}^{2}-\left(2 l_{2} \cos \psi\right) l_{3}+\left(l_{2}^{2}-c^{2}\right)=0 \tag{43}
\end{equation*}
$$

This gives

$$
\begin{equation*}
l_{3}=l_{2} \cos \psi \pm \sqrt{c^{2}-l_{2}^{2}} \sin ^{2} \psi \tag{44}
\end{equation*}
$$

in which the rules of choice of positive and negative signs before the square root are as follows:
(i) For $\psi \geq \pi / 2$ positive sign only and there is only one solution.
(ii) For $\psi<\pi / 2$ and $\sqrt{c^{2}-l^{2} \sin ^{2} \psi}>l_{2} \cos \psi$, positive sign only and there is only one solution.
(iii) For $\psi<\pi / 2$ and $l_{2} \cos \psi>\sqrt{c^{2}-l^{2} \sin \psi}$, positive and negative signs correspond to two positions for two solutions of $l_{3}$.

Figure 7 illustrates the choice of these rules. Also by applying the cosine law to $\triangle A B C$, we obtain

$$
\begin{equation*}
\cos \delta=\left(l_{3}^{2}+c^{2}-l_{2}^{2}\right) / 2 l_{3} c . \tag{45}
\end{equation*}
$$

Now, the unknown unit vectors $\hat{l}_{3}$ can be obtained in the same manner as finding $\hat{l}_{3}$ in the RGGR mechanism by using the local coordinate system and their direction cosines of corresponding unit vectors.

## 3 Displacement Analysis of the RRGG Mechanism

The RRGG four-link spatial mechanism is a variation of the popular RGGR four-link spatial linkage. Each mechanism has a single degree of freedom, with a passive degree of freedom in the GG link. A schematic diagram of a RRGG mechanism is shown in Fig. 8.

The known quantities of the mechanism are the lengths $l_{1}$, $l_{2}$, and $l_{3}$, the vector $\bar{l}_{4}$, the direction of rotation $\hat{p}_{1}$, the angles $\zeta, \eta, \alpha$, and $\beta$ from the linkage design, and the input angle $\theta$. The angle $\zeta$, between $\hat{p}_{1}$ and $\hat{l}_{1}$, and the angle $\eta$, between $\hat{p}_{1}$ and $\hat{l}_{4}$, are chosen to be less than $\pi$.

Both the angle $\alpha$, between $\hat{p}_{2}$ and $\hat{l}_{2}$, and the angle $\beta$, between $\hat{p}_{2}$ and $-\hat{l}_{1}$, are also chosen to be less than $\pi$. The unknown quantities are $\hat{l}_{1}, \hat{l}_{2}$, and $\hat{l}_{3}$.

After calculating $\hat{l}_{1}$ by equation (11), $\hat{p}_{2}$ is determined by the direction cosine matrix method as illustrated earlier. $\hat{p}_{2}$ is calculated as

(a) $\psi \geq \pi / 2$


(b) $\psi<\pi / 2, c^{\prime}>c^{\prime \prime}$

$$
c^{\prime}=\sqrt{c^{2}-l^{2} \sin ^{2} \psi}
$$

Fig. 7 Transmission angle and $I_{3}$


Fig. 8 The schematic diagram of the RRGG mechanism and its known data
where

$$
\begin{aligned}
& b_{1}=(\cos \rho-\cos \zeta \cos \beta) / \sin \zeta \\
& b_{2}=\cos \beta \\
& b_{3}= \pm \sqrt{1-b_{1}^{2}-b_{2}^{2}}
\end{aligned}
$$

The positive and negative signs of $b_{3}$ can be decided from the observation of the given mechanism construction in the local coordinate system.

Once $\hat{p}_{2}$ is determined, the analyses of $\hat{l}_{2}$ and $\hat{l}_{3}$ of the RRGG are similar procedure as the analysis of $\hat{l}_{2}$ and $\hat{l}_{3}$ of the RGGR by letting $\hat{l}_{2}$ and $\hat{l}_{3}$. Therefore, the transformation matrix equation (27) could be used with negative $\hat{c}$ components in calculation of $\hat{l}_{2}$.

The output angle $\phi$ for the rotation about the $\hat{p}_{2}$-axis can be measured with any arbitrary reference. The ouptut angle $\phi$ is chosen as the angle between the two planes formed by $\hat{p}_{2}$ and $\hat{l}_{1}$, and $\hat{p}_{2}$ and $\hat{l}_{2}$. The $p_{2} l_{2}$-plane is chosen as a reference. The output angle could be completely determined from the following two equations

$$
\begin{equation*}
\sin \phi=\frac{-\left(T_{22}\right)_{4} \hat{l}_{2 x}+\left(T_{12}\right)_{4} \hat{l}_{2 y}+\left(T_{33}\right)_{4} \cos \alpha}{\left(T_{31}\right)_{4} \sin \alpha} \tag{47}
\end{equation*}
$$

and
$\left\{\begin{array}{l}\hat{p}_{2 x} \\ \hat{p}_{2 y} \\ \hat{p}_{2 z}\end{array}\right\}=\frac{1}{\sin \zeta}\left[\begin{array}{lll}\hat{p}_{1 x}+\hat{l}_{1 x} \cos \zeta & -\hat{l}_{1 x} \sin \zeta & \hat{l}_{1 y} \hat{p}_{1 z}-\hat{l}_{1 z} \hat{p}_{1 y} \\ \hat{p}_{1 y}+\hat{l}_{x y} \cos \zeta & -\hat{l}_{1 y} \sin \zeta & \hat{l}_{1 \hat{p}_{1 z}} \hat{p}_{1 x}-\hat{l}_{1 x} \hat{p}_{1 z} \\ \hat{p}_{1 z}+\hat{l}_{1 z} \cos \zeta & -\hat{l}_{1 z} \sin \zeta & \hat{l}_{1 x} \hat{p}_{1 y}-\hat{l}_{1 y} \hat{p}_{1 x}\end{array}\right]\left\{\begin{array}{c}b_{1} \\ b_{2} \\ b_{3}\end{array}\right\}$

$$
\begin{equation*}
\cos \phi=\frac{\left(T_{23}\right)_{4} \hat{l}_{2 x}-\left(T_{13}\right)_{4} \hat{l}_{2 y}+\left(T_{32}\right)_{4} \cos \alpha}{\left(T_{31}\right)_{4} \sin \alpha} \tag{48}
\end{equation*}
$$

where the transformation matrix $\left[T_{i j}\right]_{4}$ is obtained as

$$
\left[T_{i j}\right]_{4}=
$$

$$
\frac{1}{\sin \beta}\left[\begin{array}{ccc}
\hat{p}_{2 x} \sin \beta & -\hat{l}_{1 x}-\hat{p}_{2 x} \cos \beta & -\hat{p}_{2 y} \hat{l}_{1 z}+\hat{p}_{2 z} \hat{l}_{1 y}  \tag{49}\\
\hat{p}_{2 x} \sin \beta & -\hat{l}_{x}-\hat{p}_{2 y} \cos \beta & -\hat{p}_{2 z} \hat{l}_{1 x}+\hat{p}_{2 x} \hat{l}_{1 z} \\
\hat{p}_{2 z} \sin \beta & -\hat{l}_{1 z}-\hat{p}_{2 z} \cos \beta & -\hat{p}_{2 x} \hat{l}_{1 y}+\hat{p}_{2 y} \hat{l}_{1 x}
\end{array}\right]
$$

## 4 Displacement Analysis of the RRGC Mechanism

The displacement analysis of the RRGC four-link spatial mechanism shown in Fig. 9 can be performed by using part of the analysis scheme of the RRGG.

The known data of the RRGC mechanism shown in Fig. 9 are the lengths $l_{1}$ and $l_{2}$, the vectors $\bar{l}_{4}$, and $\hat{l}_{3}$, the direction of rotation $\hat{p}_{1}$, the angles $\zeta, \eta, \alpha, \beta, \epsilon$, and the input angle $\theta$. The definition of the angles is the same as the RRGG, except angle $\epsilon$ between $\hat{l}_{3}$ and $\hat{l}_{4}$, which is chosen to be less than $\pi$. The unknown quantities of this mechanism are $\hat{l}_{1}, \hat{l}_{2}$, and $l_{3}$.

The displacement analysis of the RRGC can be performed easily by using part of the analysis scheme of the RRGG for $\hat{p}_{2}$ and part of the similar analysis scheme of the RGCR for $l_{3}$ and the output angle.

## 5 Numerical Examples

For the numerical illustrations, the dimensions and other known data of the spatial four-bar linkages are given as follows.

## Example 1

RGGR mechanism

| $l_{1}$ | $=101.6 \mathrm{~mm}$ |
| ---: | :--- |
| $l_{2}$ | $=381.0 \mathrm{~mm}$ |
| $l_{3}$ | $=254.0 \mathrm{~mm}$ |
| $l_{4}$ | $=314.2 \mathrm{~mm}$ |
| $l_{4}$ | $=(304.8,0,76.2)$ |
| $\hat{p}_{1}$ | $=(0,0,1)$ |
| $\hat{p}_{2}$ | $=(1,0,0)$ |
| $\rho$ | $=90 \mathrm{deg}$ |
| $\alpha$ | $=90 \mathrm{deg}$ |

## Example 3

RRGG mechanism

> Example 2
> RGCR mechanism
> $l_{1}=203.2 \mathrm{~mm}$
> $l_{2}=381.0 \mathrm{~mm}$
> $l_{4}=314.2 \mathrm{~mm}$
> $\bar{l}_{4}=(304.8,0,76.2)$
> $\hat{p}_{1}=(0,0,1)$
> $\hat{p}_{2}=(0,0,1)$
> $\rho=90 \mathrm{deg}$
> $\alpha=90 \mathrm{deg}$
> $\psi=74 \mathrm{deg}$

## Example 4

RRGC mechanism
$l_{1}=119.0 \mathrm{~mm}$

$$
\begin{aligned}
& l_{1}=119.0 \mathrm{~mm} \\
& l_{2}=248.6 \mathrm{~mm} \\
& l_{4}=228.6 \mathrm{~mm} \\
& \hat{l}_{3}=(0.4082,0.8165,04082) \\
& \hat{l}_{4}=(1,0,0) \\
& \hat{p}_{1}=(0,0,1) \\
& \zeta=90 \mathrm{deg} \\
& \alpha=90 \mathrm{deg} \\
& \epsilon=108 \mathrm{deg}
\end{aligned}
$$

With input angle $\theta$ as a parameter, the vectors $\bar{l}_{1}, \bar{l}_{2}$, and $\bar{l}_{3}$, the output angle $\phi$, and the transmission angle $\psi$ are determined. The transmission angle is defined as an angle between $-\hat{l}_{2}$ and $\hat{l}_{3}$. The results of Example 1 are tabulated in Table 1 for two possible configurations of the RGGR mechanism, i.e., there are two output angles $\phi_{1}$ and $\phi_{2}$ for a given input angle $\theta$. The transmission angles are the same in both configurations. Table 2 shows three components of each of the


Fig. 9 The schematic diagram of RRGC four-link spatial mechanism and its known data

Table 1 Output angle $\phi$ and transmission angle $\psi$ versus input angle $\theta$ of the RGGR mechanism, all angles in degree

| $\theta$ | $\phi_{1}$ | $\phi_{2}$ | $\psi$ |
| :---: | :---: | :---: | :---: |
| 110 | 91.6 | 165.6 | 65.8 |
| 100 | 75.1 | 179.5 | 62.3 |
| 90 | 63.0 | 190.8 | 58.7 |
| 80 | 53.3 | 201.3 | 54.9 |
| 70 | 45.3 | 211.9 | 51.0 |
| 60 | 38.4 | 223.4 | 47.2 |
| 50 | 32.5 | 236.3 | 43.5 |
| 40 | 27.4 | 251.4 | 40.1 |
| 30 | 23.4 | 269.2 | 37.1 |
| 20 | 21.0 | 290.0 | 34.8 |
| 10 | 21.8 | 312.1 | 33.4 |
| 0 | 29.9 | 330.1 | 32.9 |
| - 10 | 47.9 | 338.2 | 33.4 |
| -20 | 70.0 | 339.0 | 34.8 |
| -30 | 90.8 | 336.6 | 37.1 |
| -40 | 108.6 | 332.6 | 40.1 |
| -50 | 123.7 | 327.5 | 43.5 |
| -60 | 136.6 | 321.6 | 47.2 |
| -70 | 148.1 | 314.7 | 51.0 |
| -80 | 158.7 | 306.7 | 54.9 |
| -90 | 169.2 | 297.0 | 58.7 |
| -100 | 180.5 | 284.9 | 62.3 |
| -110 | 194.4 | 268.4 | 65.8 |

Table 2 Positions of the RGGR mechanism in system coordinates at $\theta=\mathbf{6 0 ~ d e g}$, all lengths in mm

| Configuration 1 |  |  | Configuration $2 \hat{j}$ |  |  | Resultant |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\hat{i}$ |  | $\hat{k}$ | $\hat{i}$ | $j$ | $\hat{k}$ |  |
| $\bar{l}_{1} 50.8$ | 88.0 | 0.0 | 50.8 | 88.0 | 0.0 | 101.6 |
| $\bar{I}_{2} 254.0$ | 69.8 | 275.3 | 254.0 | -262.5 | - 108.4 | 381.0 |
| $\begin{array}{lll}\bar{I}_{3} & 0.0\end{array}$ | -157.2 | -199.1 | 0.0 | 174.5 | 184.6 | 254.0 |
| $\bar{I}_{4} 304.8$ | 0.0 | 76.2 | 304.8 | 0.0 | 76.2 | 314.2 |

vector $\bar{l}_{1}, \bar{l}_{2}$, and $\bar{l}_{3}$ in the global coordinates corresponding to a particular input angle, in this case, $\theta=60 \mathrm{deg}$ for the RGGR mechanism. For Example 2, Table 3 shows output angle and the calculated length of link 3 corresponding to input angle $\theta$ in the RGGR mechanism.
For Examples 3 and 4, the results are tabulated in Table 4 for the RRGG mechanism and in Table 5 for the RRGC mechanism. Tables 4 and 5 show the results of the transmission angle and output angle for each input angle.

Table 3 Output angle $\phi$ and the length of link 3 versus input angle $\theta$ of the RGCR mechanism, angles in degree and lengths in mm

| $\theta$ | $\phi_{1}$ | $l_{3}$ |
| :---: | :---: | :---: |
| 90 | 248.4 | 181.6 |
| 100 | 264.3 | 270.5 |
| 110 | 274.5 | .324 .6 |
| 120 | 282.6 | 365.4 |
| 130 | 289.5 | 397.4 |
| 140 | 295.7 | 422.5 |
| 150 | 301.3 | 441.4 |
| 160 | 306.3 | 454.7 |
| 170 | 311.0 | 46.6 |
| 180 | 315.2 | 465.2 |
| 190 | 318.9 | 462.6 |
| 200 | 322.3 | 454.7 |
| 210 | 325.1 | 441.4 |
| 220 | 327.4 | 422.5 |
| 230 | 328.9 | 397.4 |
| 240 | 329.4 | 365.4 |
| 250 | 328.5 | 324.6 |
| 260 | 325.1 | 270.5 |
| 270 | 315.8 | 181.6 |

In the numerical example of configuration 2 of the RRGC mechanism, the input angle is limited to the range of 80 to -160 deg due to the construction of the mechanism. Another limitation should be observed. The motion that joint $B$ of the mechanism passes through the cylindrical joint $C$ during an increment of the input angle from -100 to -110 deg is impossible in actual cases, although it is theoretically possible in the analysis. Therefore, the range of motion of configuration 2 is divided into two intervals -80 to -100 deg and -110 to -160 deg .

## Conclusion

The direction cosine matrix method has been developed for a displacement analysis applicable to all types of four-link, spatial mechanisms. The analyses of the RGGR, RGCR, RRGG, and RRGC mechanisms have been illustrated to demonstrate this method. The advantage of this method is that the analysis yields exact solutions without loss of geometric interpretation and without the need for either numerical or iterative schemes.
The special property of the direction cosine matrix, that each element equals its own cofactor, is the focus of this analysis. Using this property, we avoid the inherent difficulties in the displacement analysis of four-link spatial mechanisms. For example, without using this property, equation (9) would be replaced by

$$
\begin{aligned}
& T_{13}= \pm \sqrt{1-T_{11}^{2}-T_{12}^{2}} \\
& T_{23}= \pm \sqrt{1-T_{21}^{2}-T_{22}^{2}} \\
& T_{33}= \pm \sqrt{1-T_{31}^{2}-T_{32}^{2}}
\end{aligned}
$$

and only one of the eight sets of possible combinations would be the solution. In another example, an algebraic equation of up to eighth-degree polynomial in reference [10] has to be solved numerically.

The extension of this method to a displacement analysis of mechanisms with more than four links and the continuation of kinematic analyses for determining velocities and accelerations of mechanisms with four or more links will be the topics of forthcoming papers.

Table 4 Output angle $\phi$ and transmission angle $\psi$ versus input angle $\theta$ of the RRGG mechanism, all angles in degree

| $\theta$ | $\phi_{1}$ | $\phi_{2}$ | $\psi$ |
| :---: | :---: | :---: | :---: |
| 0 | 137.5 | 225.8 | 24.2 |
| 10 | 156.0 | 245.4 | 25.4 |
| 20 | 172.9 | 263.6 | 28.8 |
| 30 | 188.3 | 279.5 | 33.6 |
| 40 | 202.3 | 293.1 | 39.2 |
| 50 | 215.1 | 304.6 | 45.4 |
| 60 | 227.1 | 314.3 | 51.8 |
| 70 | 238.3 | 322.8 | 58.3 |
| 80 | 248.7 | 330.3 | 64.7 |
| 90 | 258.6 | 337.1 | 71.0 |
| 100 | 268.0 | 343.4 | 77.1 |
| 110 | 276.8 | 349.3 | 82.8 |
| 120 | 285.1 | 355.0 | 33.2 |
| 130 | 293.0 | 0.6 | 93.0 |
| 140 | 300.5 | 6.1 | 97.1 |
| 150 | 307.6 | 11.6 | 100.5 |
| 160 | 374.3 | 17.3 | 103.0 |
| 170 | 320.8 | 23.1 | 104.6 |
| 180 | 326.9 | 29.1 | 105.1 |
| 190 | 332.9 | 35.4 | 104.6 |
| 200 | 338.6 | 41.9 | 103.0 |
| 210 | 344.2 | 48.8 | 100.5 |
| 220 | 349.8 | 56.0 | 97.1 |
| 230 | 355.3 | 63.6 | 93.0 |
| 240 | 0.9 | 71.7 | 88.2 |
| 250 | 6.8 | 80.1 | 82.8 |
| 260 | 12.9 | 89.1 | 77.1 |
| 270 | 19.4 | 98.6 | 71.0 |
| 280 | 26.6 | 108.6 | 64.7 |
| 290 | 34.7 | 119.4 | 58.3 |
| 300 | 43.9 | 130.9 | 51.8 |
| 310 | 54.6 | 143.3 | 45.4 |
| 320 | 67.2 | 156.8 | 39.2 |
| 330 | 82.1 | 171.7 | 33.6 |
| 340 | 99.3 | 188.1 | 28.8 |
| 350 | 118.1 | 206.4 | 25.4 |
| 360 | 137.5 | 225.8 | 24.2 |

Table 5 Output angle $\phi$ and transmission angle $\psi$ versus input angle $\theta$ of the RRGC mechanism, all angles in degree

| $\theta$ | $\phi_{1}$ | $\psi_{1}$ | $\phi_{2}$ | $\psi_{2}$ |
| :---: | :---: | :---: | :---: | :---: |
| Interval 1 |  |  |  |  |
| 80 | 293.5 | 79.1 | 276.2 | 100.9 |
| 70 | 297.1 | 65.3 | 258.2 | 114.7 |
| 60 | 297.0 | 55.8 | 242.3 | 124.2 |
| 50 | 29.3 | 47.6 | 225.8 | 13.4 |
| 40 | 292.5 | 40.2 | 208.0 | 139.8 |
| 30 | 288.7 | 33.6 | 189.2 | 146.4 |
| 20 | 283.8 | 22.9 | 170.5 | 152.1 |
| 10 | 277.8 | 23.2 | 153.3 | 156.8 |
| 0 | 270.8 | 19.8 | 138.3 | 160.2 |
| -10 | 262.5 | 17.8 | 125.4 | 162.2 |
| -20 | 252.8 | 17.0 | 114.4 | 163.0 |
| -30 | 241.3 | 17.3 | 104.9 | 162.7 |
| -40 | 22.7 | 18.3 | 96.5 | 16.7 |
| -50 | 212.0 | 20.0 | 89.0 | 160.0 |
| -60 | 194.9 | 22.3 | 82.3 | 157.7 |
| -70 | 177.3 | 25.1 | 76.1 | 155.0 |
| -80 | 160.8 | 28.3 | 70.3 | 151.7 |
| -90 | 145.9 | 32.0 | 64.8 | 148.0 |
| -100 | 132.7 | 36.1 | 59.6 | 143.9 |
| Interval 2 |  |  |  |  |
| -110 | 120.8 | 40.8 | 54.7 | 139.2 |
| -120 | 109.7 | 45.9 | 49.9 | 134.1 |
| -130 | 98.9 | 51.7 | 45.3 | 128.3 |
| -140 | 87.8 | 58.2 | 41.2 | 121.8 |
| -150 | 75.5 | 66.0 | 37.9 | 114.0 |
| -160 | 59.7 | 76.7 | 37.2 | 103.3 |

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