The 4D, 4th Rank, Antisymmetric Tensor and the 4D Equivalent to the Cross Product or
More Fun with Tensors!!!
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This document describes the four dimensional fourth rank antisymmetric tensor, $\varepsilon 4$, and some of its uses. The $\varepsilon 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{l})$ tensor is a vary close relative of Mathcad's $\varepsilon(\mathrm{i}, \mathrm{j}, \mathrm{k})$ tensor. First the tensor itself will be examined. After that its use to evaluate 4 by 4 determinants and to calculate the 4 dimensional equivalent of the cross product will be explored.

Below is the definition of $\varepsilon 4$, the antisymmetric tensor of rank 4, for a 4-dimensional vector space. The component's value is 0 if any of the indices are equal, and 1 if the indices are ( $0,1,2,3$ ). The tensor component is equal to -1 if the indices are an odd permutation of ( $0,1,2,3$ ), and 1 if the indices are an even permutation of ( $0,1,2,3$ ).

| perm4 (A) := |  |
| :---: | :---: |

If the indices are all different then $\varepsilon 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{l})$ uses perm4(A) to count the number of permutations needed to arrange the indices in ascending order, (0,1,2,3).

If any of the indices are the same, then that component is equal to zero.
If they are all different, examine the permutation count. If the count is odd, then the component is equal to -1 . If the count was even, then that component of the tensor is equal to 1 .

Below are a few of the 256 components of the fourth rank antisymmetric tensor.

$$
\begin{array}{lll}
\varepsilon 4(0,1,2,3)=1 & \varepsilon 4(0,0,2,3)=0 & \varepsilon 4(3,2,1,0)=1 \\
\varepsilon 4(1,2,3,0)=-1 & \varepsilon 4(0,1,1,3)=0 & \varepsilon 4(2,1,0,3)=-1 \\
\varepsilon 4(2,3,0,1)=1 & \varepsilon 4(0,1,2,2)=0 & \varepsilon 4(1,0,3,2)=1 \\
\varepsilon 4(3,0,1,2)=-1 & \varepsilon 4(3,1,2,3)=0 & \varepsilon 4(0,3,2,1)=-1
\end{array}
$$

The range variables; $\mathrm{ij}, \mathrm{j}, \mathrm{l}$ range over zero to three for the indexes of the $\varepsilon 4$ tensor.
$\mathrm{i}:=0 . .3 \quad \mathrm{j}:=0 . .3 \quad \mathrm{k}:=0 . .3 \quad 1:=0 . .3$

## Properties of a 4th rank, 4 dim. tensor

The number of components in a 4th rank, 4 dimensional tensor is:

$$
\sum_{1} \sum_{\mathrm{k}} \sum_{\mathrm{j}} \sum_{\mathrm{i}} 1=256
$$

from contraction of $\delta$ the dimension of our space is: $\quad \sum_{i} \delta(i, i)=4$
Coordinate Transformation Law is: $T_{i, j, k, 1}=\sum_{d=0}^{3} \sum_{c=0}^{3} \sum_{b=0}^{3} \sum_{a=0}^{3} T_{a, b, c, d} \cdot \lambda_{a, i} \cdot \lambda_{b, j} \cdot \lambda_{c, k} \cdot \lambda_{d, 1}$

## Some of the properties of the $\varepsilon 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{l})$ tensor

The number of non-zero components of $\varepsilon 4$ is: $\quad \sum_{1} \sum_{k} \sum_{j} \sum_{i}|\varepsilon 4(i, j, k, i)|=24$
The number of components of $\varepsilon 4$ equal to 1 is: $\quad \quad \sum_{1} \sum_{k} \sum_{j} \sum_{i}(\varepsilon 4(i, j, k, 1)=1)=12$
The number of components of $\varepsilon 4$ equal to -1 is: $\quad \sum_{1} \sum_{\mathrm{k}} \sum_{\mathrm{j}} \sum_{\mathrm{i}}(\varepsilon 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{l})=-1)=12$
The number of components of $\varepsilon 4$ equal to 0 is: $\quad \sum_{1} \sum_{\mathrm{k}} \sum_{\mathrm{j}} \sum_{\mathrm{i}}(\varepsilon 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{l})=0)=232$

Now that we have created the 4th rank totally anti-symmetric tensor, what can it be used for? The $\varepsilon 4$ tensor can be used to evaluate the determinant of a $4 \times 4$ matrix, as shown below. To evaluate the determinant one converts the columns of the matrix to column vectors Next apply the tensor to all four vectors yielding a single number, the determinant. Below are some examples, using the $\varepsilon 4$ tensor to calculate determinants and the built in Mathcad calculation for verification of the results.

$$
\begin{aligned}
& \mathrm{i}:=0 . .3 \quad \mathrm{j}:=0 . .3 \quad \mathrm{k}:=0 . .3 \quad 1:=0 . .3 \\
& \mathrm{D} 4(\mathrm{~m}):=\sum_{1} \sum_{\mathrm{k}} \sum_{\mathrm{j}} \sum_{\mathrm{i}} \mathrm{E} 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{l}) \cdot\left(\mathrm{m}^{\langle 0\rangle}\right)_{\mathrm{i}} \cdot\left(\mathrm{~m}^{\langle 1\rangle}\right)_{\mathrm{j}} \cdot\left(\mathrm{~m}^{\langle 2\rangle}\right)_{\mathrm{k}} \cdot\left(\mathrm{~m}^{\langle 3\rangle}\right)_{1} \\
& \mathrm{~m}:=\left[\begin{array}{cccc}
1 & 2 & 3 & 4 \\
0 & 1 & 2 & 3 \\
0 & 0 & 1 & 2 \\
0 & 0 & 0 & 1
\end{array}\right] \quad \mathrm{D} 4(\mathrm{~m})=1 \\
& \mathrm{~m}:=\left[\begin{array}{ccc}
0 & 1 & 0 \\
-1 & 0 & 1 \\
0 & -1 & 0 \\
-1 & 0 & -1 \\
-1
\end{array}\right] \quad \\
& \mathrm{m}:=\left[\begin{array}{cccc}
1 & 1 & 1 & 1 \\
0 & 1 & 1 & 1 \\
-1 & 0 & 1 & 1 \\
-1 & -1 & 0 & 1
\end{array}\right] \quad \mathrm{D} 4(\mathrm{~m})=4 \\
& \mathrm{~m}:=\left[\begin{array}{cccc}
1 & 0 & 3 & 1 \\
5 & 6 & 7 & 1 \\
9 & 10 & 11 & 2
\end{array}\right] \\
& \mathrm{D} 4(\mathrm{~m})=1
\end{aligned}
$$

Since the $\varepsilon 4$ tensor can be used to define the 4 by 4 determinant, one should be able to define the $\varepsilon 4$ tensor in terms of the 4 by 4 determinant and the $\delta$ function or $\delta$ tensor.

The $4 \times 4$ determinant of delta functions is used to define the eighth rank $\delta$ tensor with 4 contravariant indexes and 4 covariant indexes. The components of the $\varepsilon 4$ a tensor, defined from the 4 by 4 determinant, are a subset of the eighth rank $\delta$ tensor's components. If you are still using Mathcad-3.1 you would have to use the definition below for the 4D, 4th rank anti-symetric tensor.
$\left.\delta 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, 1, \mathrm{r}, \mathrm{s}, \mathrm{t}, \mathrm{u}):=\left\lvert\, \begin{array}{llll}\boldsymbol{\delta}(\mathrm{i}, \mathrm{r}) & \delta(\mathrm{i}, \mathrm{s}) & \delta(\mathrm{i}, \mathrm{t}) & \delta(\mathrm{i}, \mathrm{u}) \\ \delta(\mathrm{j}, \mathrm{r}) & \delta(\mathrm{j}, \mathrm{s}) & \delta(\mathrm{j}, \mathrm{t}) & \boldsymbol{\delta}(\mathrm{j}, \mathrm{u}) \\ \boldsymbol{\delta}(\mathrm{k}, \mathrm{r}) & \delta(\mathrm{k}, \mathrm{s}) & \delta(\mathrm{k}, \mathrm{t}) & \boldsymbol{\delta}(\mathrm{k}, \mathrm{u}) \\ \delta(1, \mathrm{r}) & \delta(1, \mathrm{~s}) & \delta(1, \mathrm{t}) & \delta(1, \mathrm{u})\end{array}\right.\right] \mid \quad \varepsilon 4 \mathrm{a}(\mathrm{r}, \mathrm{s}, \mathrm{t}, \mathrm{u}):=\delta 4(0,1,2,3, \mathrm{r}, \mathrm{s}, \mathrm{t}, \mathrm{u})$

Later in this paper both $\varepsilon 4$ and $\varepsilon 4$ a will be used to calculate the same 4 dimensional cross product to verify their results.

The second interesting thing you can do with the 4th rank totally antisymmetric tensor is to perform the 4 dimensional equivalent of the cross product in Euclidean 3-space, R3. First let's review some of the properties of the cross product in R3. The R3 cross product is geometrically defined as:
$\mathrm{C}=\mathrm{A} \times \mathrm{B}=|\mathrm{A}| \cdot|\mathrm{B}| \cdot \sin (\theta) \cdot \frac{\mathrm{E}}{|\mathrm{E}|} \quad \begin{aligned} & \mathrm{E} \text { is perpendicular to } \mathrm{A} \text { and } \mathrm{B} \text {. The direction is given by the } \\ & \text { right hand rule in cartesianl coordinate systems. }\end{aligned}$

$$
A \cdot E=0 \text { and } B \cdot E=0
$$

Vector C's magnitude is equal to the length of A times the length of $B$ times the sin of the angle between $A$ and $B$. The cross product of $A$ and $B$ gives a normal vector to both $A$ and $B$, who's length is equal to the area of of the parallelogram constructed from $A$ and $B$. A cross $B$ is anti-symetric so:

$$
\begin{aligned}
& \mathrm{A} \times \mathrm{B}=-\mathrm{B} \times \mathrm{A} \\
& \text { given } \mathrm{e} \text {, the matrix of basis vectors: } \quad \mathrm{e}:=\left[\begin{array}{lll}
1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right] \quad \begin{array}{l}
\text { The non-zero elements of the cross } \\
\text { product of the basis vectors in R3 are: }
\end{array} \\
& \mathrm{e}^{\langle 0\rangle} \times \mathrm{e}^{\langle 1\rangle}=\mathrm{e}^{\langle 2\rangle} \quad \mathrm{e}^{\langle 1\rangle} \times \mathrm{e}^{\langle 2\rangle}=\mathrm{e}^{\langle 0\rangle} \quad \mathrm{e}^{\langle 2\rangle} \times \mathrm{e}^{\langle 0\rangle}=\mathrm{e}^{\langle 1\rangle} \\
& \mathrm{e}^{\langle 1\rangle} \times \mathrm{e}^{\langle 0\rangle}=-\left(\mathrm{e}^{\langle 2\rangle}\right) \\
& \mathrm{e}^{\langle 2\rangle} \times \mathrm{e}^{\langle 1\rangle}=-\left(\mathrm{e}^{\langle 0\rangle}\right) \\
& \mathrm{e}^{\langle 0\rangle} \times \mathrm{e}^{\langle 2\rangle}=-\left(\mathrm{e}^{\langle 1\rangle}\right)
\end{aligned}
$$

$$
\begin{gathered}
\mathrm{e}_{0} \\
\mathrm{e}_{1}
\end{gathered} \mathrm{e}_{2}{ }_{\mathrm{e}_{0}}^{\mathrm{e}_{1}} \begin{array}{ccc}
0 & \mathrm{e}_{2} & -\mathrm{e}_{1} \\
\mathrm{e}_{2}
\end{array}\left[\begin{array}{ccc}
\mathrm{e}_{2} & \mathrm{e}_{0} \\
\mathrm{e}_{1} & -\mathrm{e}_{0} & 0
\end{array}\right]
$$

The multiplication table to the left summarizes the action of the cross product on the orthonormal basis vector of Euclidean 3-space. From this table one can see that the 3rd rank, 3 dimensional antisymmetric tensor, which is Mathcad's built in $\varepsilon(i, j, k$,$) tensor, applied to two$ independent 3 -vectors, can be used to evaluate the cross product.

The cross product implemented with the 3 dimensional anti-symetric tensor is:
$\mathrm{i} 3:=0 . .2 \quad \mathrm{j} 3:=0 . .2 \quad \mathrm{k} 3:=0 . .2$
$C P 3(A, B):=\sum_{k 3} \sum_{j 3} \sum_{i 3} e^{\langle k 3\rangle} \cdot\left(\varepsilon(i 3, j 3, k 3) \cdot A_{i 3} \cdot B_{j 3}\right)$
The k3 summation is used to collect the components, from the application of the $\varepsilon$ tensor to vectors $A$ and $B$, into the results vector. Note mathcad's order of operations is used to always force a vector result.

$$
\begin{array}{lll}
\mathrm{A}:=\left[\begin{array}{lll}
1 & 2 & 3
\end{array}\right]^{\mathrm{T}} & \mathrm{~B}:=\left[\begin{array}{lll}
4 & 5 & 6
\end{array}\right]^{\mathrm{T}} & \text { Two independent 3-vectors for the example below. } \\
\mathrm{Cp}:=\mathrm{CP} 3(\mathrm{~A}, \mathrm{~B}) & \mathrm{Cp}^{\mathrm{T}}=\left[\begin{array}{lll}
-3 & 6 & -3
\end{array}\right] & (\mathrm{A} \times \mathrm{B})^{\mathrm{T}}=\left[\begin{array}{lll}
-3 & 6 & -3
\end{array}\right] \\
\mathrm{Cn}:=\mathrm{CP} 3(\mathrm{~B}, \mathrm{~A}) & \mathrm{Cn}^{\mathrm{T}}=\left[\begin{array}{lll}
3 & -6 & 3
\end{array}\right] & (\mathrm{B} \times \mathrm{A})^{\mathrm{T}}=\left[\begin{array}{lll}
3 & -6 & 3
\end{array}\right] \quad \text { it is anti-communitive }
\end{array}
$$

$A \cdot C p=0 \quad B \cdot C p=0 \quad$ The cross product vector is orthogonal to both $A$ and $B$
$\mathrm{CP} 3(\mathrm{~A}, \mathrm{~A})^{\mathrm{T}}=\left[\begin{array}{lll}0 & 0 & 0\end{array}\right]$

$$
(A \times A)^{T}=\left[\begin{array}{lll}
0 & 0 & 0
\end{array}\right] \quad \begin{aligned}
& \text { The Cross product of } \\
& \text { dependent vectors is zero }
\end{aligned}
$$

The $\varepsilon 4$ tensor will now be used to generate the four-dimensional equivalent of the cross product. This function takes three linearly independent 4 dimensional vector that span a 3D hyper-plane and returns a 4-vector perpendicular to all the input vectors. Note: A fourth rank tensor must be applied to three vector in order to produce a single vector as its output.

$$
\mathrm{i} 4:=0 . .3 \quad \mathrm{j} 4:=0 . .3 \quad \mathrm{k} 4:=0 . .3 \quad 14:=0 . .3
$$

$\mathrm{id} 4:=\left[\begin{array}{cccc}1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1\end{array}\right]$
The id4 matrix has the four dimensional orthonormal basis vectors as its columns.

The four dimensional extension of the Cross product using the $\varepsilon 4$ tensor and three 4 -vectors is:

$$
\mathrm{CP} 4(\mathrm{~V} 1, \mathrm{~V} 2, \mathrm{~V} 3):=\sum_{14} \mathrm{id} 4^{\langle 14\rangle} \cdot \sum_{\mathrm{k} 4} \sum_{\mathrm{j} 4} \sum_{\mathrm{i} 4}\left(\varepsilon 4(\mathrm{i} 4, \mathrm{j} 4, \mathrm{k} 4,14) \cdot \mathrm{V} 1_{\mathrm{i} 4} \cdot \mathrm{~V} 2_{\mathrm{j} 4} \cdot \mathrm{~V} 3_{\mathrm{k} 4}\right)
$$

Again, Mathcad's order of operations is used to force CP4 to always return a vector. The zero 4 -vector is returned instead of just scalar zero.


CP4a is base on the 4th rank, 4 dimensional antisymmetric tensor generated from the determinant and delta function definition. This had better yield the same results as CP4.

$$
\mathrm{T} 1:=\mathrm{id} 4^{\langle 0\rangle} \mathrm{T} 2:=\mathrm{id} 4^{\langle 1\rangle} \mathrm{T} 3:=\mathrm{id} 4^{\langle 2\rangle} \mathrm{T} 4:=\mathrm{id} 4^{\langle 3\rangle} \quad \begin{aligned}
& \mathrm{T} 1, \mathrm{~T} 2, \mathrm{~T} 3, \mathrm{~T} 4 \text { are set to the four } \\
& \text { orthonormal bais vectors }
\end{aligned}
$$

Below are the results of the four dimensional cross product applied to sets of three basis vectors at a time.

| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 4, \mathrm{~T} 3, \mathrm{~T} 2)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}1 & 0 & 0 & 0\end{array}\right]$ | The 4D cross product of three of the four basis vectors yields the fourth basis vector that's perpendicular to the other three. |
| :---: | :---: | :---: |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 1, \mathrm{~T} 3, \mathrm{~T} 4)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}0 & 1 & 0 & 0\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 2, \mathrm{~T} 1, \mathrm{~T} 4)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}0 & 0 & 1 & 0\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 1, \mathrm{~T} 2, \mathrm{~T} 3)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}0 & 0 & 0 & 1\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 2, \mathrm{~T} 3, \mathrm{~T} 4)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}-1 & 0 & 0 & 0\end{array}\right]$ | The 4D cross product is anti-communitive just like the 3 dimensional cross product |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 2, \mathrm{~T} 4, \mathrm{~T} 3)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}1 & 0 & 0 & 0\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 3, \mathrm{~T} 2, \mathrm{~T} 4)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}1 & 0 & 0 & 0\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 2, \mathrm{~T} 3, \mathrm{~T} 4)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}-1 & 0 & 0 & 0\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 1, \mathrm{~T} 2, \mathrm{~T} 2)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}0 & 0 & 0 & 0\end{array}\right]$ | The 4D cross product returns the zero vector if any of the input vectors are linearly dependent. |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 2, \mathrm{~T} 2, \mathrm{~T} 3)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}0 & 0 & 0 & 0\end{array}\right]$ |  |
| $\mathrm{O} 4:=\mathrm{CP} 4(\mathrm{~T} 4, \mathrm{~T} 4, \mathrm{~T} 4)$ | $\mathrm{O} 4^{\mathrm{T}}=\left[\begin{array}{llll}0 & 0 & 0 & 0\end{array}\right]$ |  |

Below are three linearly independent test vectors and their four-dimensional cross product calculated by both $\varepsilon 4(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{I})$ and $\varepsilon 4 \mathrm{a}(\mathrm{i}, \mathrm{j}, \mathrm{k}, \mathrm{I})$. This was done to check the results of both the programmed version of the $\varepsilon$ tensor and the version generated from the 8th rank $\delta$ tensor. We get the same result from both calculations.

3 independent test 4-vectors

$$
\text { A1 }:=\left[\begin{array}{l}
1 \\
2 \\
3 \\
4
\end{array}\right] \quad \text { A2 }:=\left[\begin{array}{l}
5 \\
6 \\
7 \\
8
\end{array}\right] \quad \text { A3 }:=\left[\begin{array}{c}
-9 \\
10 \\
11 \\
12
\end{array}\right]
$$

Results of the 4 dimensional cross products are equal
$\mathrm{CP} 4(\mathrm{~A} 1, \mathrm{~A} 2, \mathrm{~A} 3)=\left[\begin{array}{c}0 \\ 72 \\ -144 \\ 72\end{array}\right] \quad \mathrm{CP} 4 \mathrm{a}(\mathrm{A} 1, \mathrm{~A} 2, \mathrm{~A} 3)=\left[\begin{array}{c}0 \\ 72 \\ -144 \\ 72\end{array}\right]$

The results of the 4D cross product, R4, is orthogonal to A1, A2, and A3 as expected.

$$
\begin{aligned}
& \mathrm{R} 4:=\mathrm{CP} 4(\mathrm{~A} 1, \mathrm{~A} 2, \mathrm{~A} 3) \\
& \mathrm{R} 4^{\mathrm{T}}=\left[\begin{array}{lll}
0 & 72 & -144 \\
72
\end{array}\right] \\
& \mathrm{R} 4 \cdot \mathrm{~A} 1=0 \quad \mathrm{R} 4 \cdot \mathrm{~A} 2=0 \quad \mathrm{R} 4 \cdot \mathrm{~A} 3=0
\end{aligned}
$$

## Closing Notes:

The Author would like to point out that the results of the four dimensional cross product are exactly correct only for Euclidean vector spaces with orthonormal basis vectors. In these spaces contravariant vectors, or just plain "vectors", and covariant vectors, or one-forms have identical components. This is because the metric is equal to the identity matrix.

The $\varepsilon 4$ tensor defined in this paper is of type ( 0,4 ), or it operates on four (contravariant) vectors and returns a real number. if this tensor operates on three vectors, as in the four dimensional cross product, the result is a (covariant vector) one-form. Only in a Euclidean space with orthonormal basis vectors will the one-form's components, returned by $\varepsilon 4$, be equal to its dual vector's components.

To turn an one-form into a vector, must apply the inverse (contravariant) metric tensor to it. The inverse metric maps one-forms into vectors and the (covariant) metric tensors maps vectors into one-forms. The metric allows the mapping of vectors into one-forms and visa-versa. This establishes the duality between vectors and one-forms in general metric spaces.

A more general definition of the vector cross product with non-orthogonal basis vectors would be:

$$
\overrightarrow{\mathrm{CP} 4}=\mathrm{G}^{-1} \cdot \varepsilon 4\left(\begin{array}{lll}
\overrightarrow{\mathrm{a}} 1, \overrightarrow{\mathrm{a}} 2, \mathrm{a} 3
\end{array}\right)
$$

$\mathrm{G}^{-1}$ is inverse of the metric
$\varepsilon 4$ is the 4 dimensional anti-symmetric tensor
a1, a2, a3, and CP4 are contravariant vectors
Although the one form result of $\varepsilon 4(\vec{a} 1, \vec{a} 2, \vec{a})$ is just as useful.

## References

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