

Research Article

New Types of Doubly Periodic Standing Wave Solutions for the Coupled Higgs Field Equation

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Based on the Hirota bilinear method and theta function identities, we obtain a new type of doubly periodic standing wave solutions for a coupled Higgs field equation. The Jacobi elliptic function expression and long wave limits of the periodic solutions are also presented. By selecting appropriate parameter values, we analyze the interaction properties of periodic-periodic waves and periodic-solitary waves by some figures.

1. Introduction

As is well known, the investigation of the explicit exact solutions to nonlinear evolution equations (NLEEs) plays an important role in the study of nonlinear physical phenomena [1, 2]. In the past decades, much effort has been spent on the construction of various forms of periodic wave solutions for NLEEs, and many powerful methods have been developed such as the algebrogeometrical approach [3, 4], nonlinearization approach of Lax pairs [5, 6], Weierstrass elliptic function expansion method [7, 8], Jacobi elliptic function expansion method [9–11], and subsidiary ordinary differential equation methods [12–18].

In this paper, we will focus on a coupled Higgs field equation with important physical interests [19],

$$\begin{aligned}u_{tt} - u_{xx} - \beta u + \delta |u|^2 u - 2uv &= 0, \\v_{tt} + v_{xx} - \delta (|u|^2)_{xx} &= 0,\end{aligned}\tag{1}$$

which describes a system of conserved scalar nucleons interacting with neutral scalar mesons in particle physics. Here β and δ are constants, and the function $v = v(x, t)$ represents a real scalar meson field and $u = u(x, t)$ a complex scalar nucleon field. Equation (1) is related to some nonlinear models with physical interests. Equation (1) is the coupled nonlinear Klein-Gordon equations for $\beta < 0$ and $\delta < 0$ and the Higgs equations for $\beta > 0$ and $\delta > 0$. Much attention

has been paid to investigate exact explicit solutions and integrable properties of (1). The symmetry reductions, the homoclinic orbits, N -soliton solutions, rogue wave solutions, Jacobi periodic solutions, and other types of travelling wave solutions have been presented [19–24].

The Hirota bilinear method is a powerful tool for constructing various exact solutions for NLEEs, which include soliton, negaton, rogue waves, rational solutions, and quasiperiodic solutions [25–35]. Recently, by means of Hirota bilinear method and theta function identities [36–38], Fan et al. obtained a class of doubly periodic standing wave solutions of (1) [39], which was expressed as rational functions of elliptic/theta functions of different moduli. A significant portion of these solutions represents travelling wave, that is, those which will remain steady in an appropriate frame of reference. Physically, the envelope of these oscillations is bounded by a pattern periodic in both time and space. The focus of this work is to investigate new types of doubly periodic standing wave solutions for (1).

This paper is organized as follows. In Section 2, we briefly illustrate some properties of theta functions and Jacobi elliptic functions. In Section 3, we construct a new kind of doubly periodic wave solutions for the coupled Higgs field equation. In Section 4, for the obtained periodic solution, we derive its Jacobi elliptic function representation and analyze interaction properties by some figures. Some conclusions are given in Section 5.

2. The Theta and Jacobi Elliptic Functions

The main tools used in this paper are Hirota operators and theta function formulas, which will be discussed here, to fix the notations and make our presentation self-contained. More formulas for the theta functions can be found in [36, 37].

The Riemann theta functions of genus one $\vartheta_n(x, \tau)$ ($n = 1-4$), the parameter q (the nome), and τ (pure imaginary) are defined by [40]

$$\begin{aligned} \vartheta_1(x, \tau) &= 2 \sum_{n=0}^{\infty} (-1)^n q^{(n+1/2)^2} \sin(2n+1)x, \\ \vartheta_2(x, \tau) &= 2 \sum_{n=0}^{\infty} q^{(n+1/2)^2} \cos(2n+1)x, \\ \vartheta_3(x, \tau) &= 1 + 2 \sum_{n=1}^{\infty} q^{n^2} \cos 2nx, \\ \vartheta_4(x, \tau) &= 1 + 2 \sum_{n=1}^{\infty} (-1)^n q^{n^2} \cos 2nx, \end{aligned} \tag{2}$$

$$0 < q < 1, \quad q = \exp(i\pi\tau), \quad q = \exp\left(-\frac{\pi K'}{K}\right).$$

Here, K, K' are the complete elliptic integrals of the first kind:

$$\begin{aligned} K &= \int_0^{\pi/2} (1 - k^2 \sin^2 \xi)^{-1/2} d\xi, \\ K_1 &= \int_0^{\pi/2} (1 - k_1^2 \sin^2 \xi)^{-1/2} d\xi. \end{aligned} \tag{3}$$

ϑ_1 is an odd function while the other three are even functions. The zeros of $\vartheta_1, \vartheta_2, \vartheta_3$, and ϑ_4 are at $M\pi + N\pi\tau, (M + (1/2))\pi + N\pi\tau, (M + (1/2))\pi + (N + (1/2))\pi\tau$, and $M\pi + (N + (1/2))\pi\tau$, respectively, and M and N are integers. Since ϑ_1, ϑ_2 , and $(\vartheta_3, \vartheta_4)$ are related by a phase shift of $\pi/2$, there are roughly two groups of theta functions.

There exists a large class of bilinear identities involving products of theta functions, some of which are listed here:

$$\begin{aligned} D_x^2 \vartheta_1(x, \tau) \cdot \vartheta_2(x, \tau) &= (b_2 + \vartheta_2^4(0, \tau)) \vartheta_1(x, \tau) \vartheta_2(x, \tau), \\ D_x^2 \vartheta_1(x, \tau) \cdot \vartheta_4(x, \tau) &= (b_2 - \vartheta_4^4(0, \tau)) \vartheta_1(x, \tau) \vartheta_4(x, \tau), \\ D_x^2 \vartheta_2(x, \tau) \cdot \vartheta_3(x, \tau) &= (b_2 - \vartheta_4^4(0, \tau)) \vartheta_2(x, \tau) \vartheta_3(x, \tau), \\ D_x^2 \vartheta_1(x, \tau) \cdot \vartheta_3(x, \tau) &= b_3 \vartheta_1(x, \tau) \vartheta_3(x, \tau), \\ D_x^2 \vartheta_2(x, \tau) \cdot \vartheta_4(x, \tau) &= b_3 \vartheta_2(x, \tau) \vartheta_4(x, \tau), \\ D_x^2 \vartheta_2(x, \tau) \cdot \vartheta_2(x, \tau) &= b_2 \vartheta_2^2(x, \tau) - b_1 \vartheta_4^2(x, \tau), \\ D_x^2 \vartheta_4(x, \tau) \cdot \vartheta_4(x, \tau) &= b_1 \vartheta_2^2(x, \tau) + b_2 \vartheta_4^2(x, \tau), \end{aligned} \tag{4}$$

where for simplicity we have used the notations

$$\begin{aligned} b_1 &= 2\vartheta_2^2(0, \tau) \vartheta_4^2(0, \tau), \quad b_2 = 2 \frac{\vartheta_3''(0, \tau)}{\vartheta_3(0, \tau)}, \\ b_3 &= \frac{\vartheta_2''(0, \tau)}{\vartheta_2(0, \tau)} + \frac{\vartheta_4''(0, \tau)}{\vartheta_4(0, \tau)}, \end{aligned} \tag{5}$$

and the formulas (4) can be derived from product identities of theta functions; the details can be found in [36, 37].

There are close connections between theta functions and elliptic functions as follows:

$$\begin{aligned} \frac{\vartheta_1(\gamma x, \tau)}{\vartheta_4(\gamma x, \tau)} &= k^{1/2} \operatorname{sn}(\lambda x, k) \equiv S, \\ \frac{\vartheta_2(\gamma x, \tau)}{\vartheta_4(\gamma x, \tau)} &= k^{1/2} (1 - k^2)^{-1/4} \operatorname{cn}(\lambda x, k) \equiv C, \\ \frac{\vartheta_3(\gamma x, \tau)}{\vartheta_4(\gamma x, \tau)} &= (1 - k^2)^{-1/4} \operatorname{dn}(\lambda x, k) \equiv D, \end{aligned} \tag{6}$$

where $\lambda = \gamma \vartheta_3^2(0, \tau)$ and $k = \vartheta_2^2(0, \tau) / \vartheta_3^2(0, \tau)$,

$$\begin{aligned} \frac{\vartheta_1(\omega t, \tau_1)}{\vartheta_4(\omega t, \tau_1)} &= k_1^{1/2} \operatorname{sn}(\rho t, k_1) \equiv S_1, \\ \frac{\vartheta_2(\omega t, \tau_1)}{\vartheta_4(\omega t, \tau_1)} &= k_1^{1/2} (1 - k_1^2)^{-1/4} \operatorname{cn}(\rho t, k_1) \equiv C_1, \\ \frac{\vartheta_3(\omega t, \tau_1)}{\vartheta_4(\omega t, \tau_1)} &= (1 - k_1^2)^{-1/4} \operatorname{dn}(\rho t, k_1) \equiv D_1, \end{aligned} \tag{7}$$

where $\rho = \omega \vartheta_3^2(0, \tau_1)$ and $k_1 = \vartheta_2^2(0, \tau_1) / \vartheta_3^2(0, \tau_1)$. It is clearly shown that arguments of the theta and elliptic functions are related by a scale factor.

3. A New Class of Doubly Periodic Wave Solutions

In this section, we construct a new class of doubly periodic wave solutions by Hirota bilinear method [2]. For (1), substituting the following transformation

$$u(x, t) = \frac{g}{f}, \quad v(x, t) = v_0 + 2(\ln f)_{xx} \tag{8}$$

into (1) and integrating with respect to x yield the bilinear forms

$$\begin{aligned} (D_t^2 + D_x^2 - \Lambda) f \cdot f &= \delta g \cdot g^*, \\ [D_t^2 - D_x^2 - (\beta + 2v_0 + \Lambda)] g \cdot f &= 0, \end{aligned} \tag{9}$$

where v_0 is a constant, $\Lambda = \Lambda(t)$ is an integration constant, and D is the well-known Hirota bilinear operator. Equation (10) is slightly different from the results given in [39] by adding one integration constant term.

The crucial step to derive doubly periodic wave solutions is to suppose f and g in (9) and (10) as suitable combination

of different theta functions. To obtain new doubly periodic wave solutions, we make a new ansatz,

$$\begin{aligned}
 g &= \eta [\vartheta_3(\omega t, \tau_1) \vartheta_2(\gamma x, \tau) \\
 &\quad + i\vartheta_1(\omega t, \tau_1) \vartheta_4(\gamma x, \tau)] \exp(ipt), \tag{11} \\
 f &= \vartheta_4(\omega t, \tau_1) \vartheta_4(\gamma x, \tau) + \vartheta_2(\omega t, \tau_1) \vartheta_2(\gamma x, \tau),
 \end{aligned}$$

where the parameters ω , γ , η , and p are constants to be determined and the period τ in the spatial direction and the period τ_1 in the temporal direction are purely imaginary constants.

Inserting the ansatz (11) into (9), together with the theta function identities given in Section 2, we set the coefficients of the terms $\vartheta_2^2(\gamma x, \tau)$, $\vartheta_4^2(\gamma x, \tau)$, and $\vartheta_2(\gamma x, \tau)\vartheta_4(\gamma x, \tau)$ to be zero and get

$$\begin{aligned}
 &(\gamma^2 b_2 + \omega^2 \tilde{b}_2 - \Lambda - \delta\eta^2 \tilde{b}_4) \vartheta_2^2(\omega t, \tau_1) \\
 &\quad + (\gamma^2 b_1 - \omega^2 \tilde{b}_1 - \delta\eta^2 \tilde{b}_5) \vartheta_4^2(\omega t, \tau_1) = 0, \\
 &(\gamma^2 b_1 - \omega^2 \tilde{b}_1 - \delta\eta^2 \tilde{b}_5) \vartheta_2^2(\omega t, \tau_1) \\
 &\quad - (\gamma^2 b_2 + \omega^2 \tilde{b}_2 - \Lambda - \delta\eta^2 \tilde{b}_4) \vartheta_4^2(\omega t, \tau_1) = 0, \\
 &(\gamma^2 b_3 + \omega^2 \tilde{b}_3 - \Lambda) \vartheta_2(\omega t, \tau_1) \vartheta_4(\omega t, \tau_1) = 0,
 \end{aligned} \tag{12}$$

where the theta constants \tilde{b}_j ($j = 1, \dots, 5$) are defined by

$$\begin{aligned}
 \tilde{b}_1 &= 2\vartheta_2^2(0, \tau_1) \vartheta_4^2(0, \tau_1), & \tilde{b}_2 &= \frac{2\vartheta_3''(0, \tau_1)}{\vartheta_3(0, \tau_1)}, \\
 \tilde{b}_3 &= \frac{\vartheta_2''(0, \tau_1)}{\vartheta_2(0, \tau_1)} + \frac{\vartheta_4''(0, \tau_1)}{\vartheta_4(0, \tau_1)}, \tag{13} \\
 \tilde{b}_4 &= \frac{\vartheta_2^2(0, \tau_1)}{\vartheta_3^2(0, \tau_1)}, & \tilde{b}_5 &= \frac{\vartheta_4^2(0, \tau_1)}{\vartheta_3^2(0, \tau_1)}.
 \end{aligned}$$

Due to the linear independence of theta functions, we set the coefficients of $\vartheta_2^2(\omega t, \tau_1)$, $\vartheta_4^2(\omega t, \tau_1)$, and $\vartheta_2(\omega t, \tau_1)\vartheta_4(\omega t, \tau_1)$ in (12) to be zero and obtain

$$\begin{aligned}
 &\gamma^2 b_1 - \omega^2 \tilde{b}_1 - \delta\eta^2 \tilde{b}_5 = 0, \\
 &\gamma^2 b_2 + \omega^2 \tilde{b}_2 - \Lambda - \delta\eta^2 \tilde{b}_4 = 0, \tag{14} \\
 &\gamma^2 b_3 + \omega^2 \tilde{b}_3 - \Lambda = 0.
 \end{aligned}$$

Similarly, substituting the ansatz (11) into (10) and applying identities for Hirota derivatives of theta functions yield

$$\begin{aligned}
 &(\omega^2 (\tilde{b}_2 - \vartheta_4^4(0, \tau_1)) - \gamma^2 b_2 - (\beta + 2\nu_0 + \Lambda + p^2)) \\
 &\quad \times \vartheta_2(\omega t, \tau_1) \vartheta_3(\omega t, \tau_1) \vartheta_2^2(\gamma x, \tau) \\
 &\quad + (\gamma^2 b_1 - 2p\omega\vartheta_4^2(0, \tau_1)) \vartheta_2(\omega t, \tau_1) \vartheta_3(\omega t, \tau_1) \vartheta_4^2(\gamma x, \tau) \\
 &\quad + (\omega^2 (\tilde{b}_2 + \vartheta_2^4(0, \tau_1)) - \gamma^2 b_3 - 2p\omega\vartheta_2^2(0, \tau_1) \\
 &\quad \quad - (\beta + 2\nu_0 + \Lambda + p^2)) \vartheta_3 \\
 &\quad \times (\omega t, \tau_1) \vartheta_4(\omega t, \tau_1) \vartheta_2(\gamma x, \tau) \vartheta_4(\gamma x, \tau) \\
 &\quad + i [- (\gamma^2 b_1 - 2p\omega\vartheta_4^2(0, \tau_1)) \vartheta_1(\omega t, \tau_1) \vartheta_4(\omega t, \tau_1) \vartheta_2^2(\gamma x, \tau) \\
 &\quad \quad + (\omega^2 (\tilde{b}_2 + \vartheta_2^4(0, \tau_1)) - \gamma^2 b_3 - 2p\omega\vartheta_2^2(0, \tau_1) \\
 &\quad \quad - (\beta + 2\nu_0 + \Lambda + p^2)) \\
 &\quad \quad \times \vartheta_1(\omega t, \tau_1) \vartheta_2(\omega t, \tau_1) \vartheta_2(\gamma x, \tau) \vartheta_4(\gamma x, \tau) \\
 &\quad \quad + (\omega^2 (\tilde{b}_2 - \vartheta_4^4(0, \tau_1)) - \gamma^2 b_2 - (\beta + 2\nu_0 + \Lambda + p^2)) \\
 &\quad \quad \times \vartheta_1(\omega t, \tau_1) \vartheta_4(\omega t, \tau_1) \vartheta_4^2(\gamma x, \tau)] = 0. \tag{15}
 \end{aligned}$$

For the real part and the imaginary part of (15), setting the coefficients of the terms $\vartheta_2^2(\gamma x, \tau)$, $\vartheta_4^2(\gamma x, \tau)$, and $\vartheta_2(\gamma x, \tau)\vartheta_4(\gamma x, \tau)$ to be zero yields an algebraic system as follows:

$$\begin{aligned}
 &\gamma^2 b_1 - 2p\omega\vartheta_4^2(0, \tau_1) = 0, \\
 &\omega^2 (\tilde{b}_2 - \vartheta_4^4(0, \tau_1)) - \gamma^2 b_2 - (\beta + 2\nu_0 + \Lambda + p^2) = 0, \tag{16} \\
 &\omega^2 (\tilde{b}_2 + \vartheta_2^4(0, \tau_1)) - \gamma^2 b_3 - 2p\omega\vartheta_2^2(0, \tau_1) \\
 &\quad - (\beta + 2\nu_0 + \Lambda + p^2) = 0.
 \end{aligned}$$

Solving the algebraic system given by (14) and (16) with respect to the variables ν_0 , η , p , γ , ω , τ , τ_1 , and Λ , one obtains a set of nontrivial solutions:

$$\begin{aligned}
 \omega^2 &= \left(\frac{b_1 \tilde{b}_4 - \tilde{b}_5 (b_2 - b_3)}{\tilde{b}_1 \tilde{b}_4 + (\tilde{b}_2 - \tilde{b}_3) \tilde{b}_5} \right) \gamma^2, & p &= \frac{b_1 \gamma^2}{2\omega\vartheta_4^2(0, \tau_1)}, \\
 \Lambda &= b_3 \gamma^2 + \tilde{b}_3 \omega^2, & \eta^2 &= \frac{\gamma^2 (b_2 - b_3) + \omega^2 (\tilde{b}_2 - \tilde{b}_3)}{\delta\tilde{b}_4}, \tag{17} \\
 \nu_0 &= \frac{\omega^2 (\tilde{b}_2 - \vartheta_4^4(0, \tau_1))}{2} - \frac{b_2 \gamma^2 + \beta + \Lambda + p^2}{2},
 \end{aligned}$$

where the parameters γ , τ , and τ_1 are arbitrary constants.

Therefore, we obtain a new doubly periodic wave solution of (1):

$$\begin{aligned}
 u(x, t) &= \frac{\eta [\vartheta_3(\omega t, \tau_1) \vartheta_2(\gamma x, \tau) + i \vartheta_1(\omega t, \tau_1) \vartheta_4(\gamma x, \tau)] \exp(ipt)}{\vartheta_4(\omega t, \tau_1) \vartheta_4(\gamma x, \tau) + \vartheta_2(\omega t, \tau_1) \vartheta_2(\gamma x, \tau)}, \\
 v(x, t) &= v_0 + 2 (\ln [\vartheta_4(\omega t, \tau_1) \vartheta_4(\gamma x, \tau) \\
 &\quad + \vartheta_2(\omega t, \tau_1) \vartheta_2(\gamma x, \tau)])_{xx}, \tag{18}
 \end{aligned}$$

where $\gamma, \tau,$ and τ_1 are arbitrary constants, and the parameters $\omega, p, \eta, v_0,$ and Λ are given by (17). To the author's knowledge, the solution (18) is firstly reported here.

In fact, the coupled Higgs field equation (1) admits abundant families of doubly periodic wave solutions. For example, we can suppose the solutions of the bilinear equations (9) and (10) as

$$\begin{aligned}
 g &= \eta [\vartheta_1(\omega t, \tau_1) \vartheta_3(\gamma x, \tau) \\
 &\quad + i \vartheta_2(\omega t, \tau_1) \vartheta_4(\gamma x, \tau)] \exp(ipt), \tag{19}
 \end{aligned}$$

$$f = \vartheta_4(\omega t, \tau_1) \vartheta_4(\gamma x, \tau) + \vartheta_3(\omega t, \tau_1) \vartheta_3(\gamma x, \tau),$$

or

$$\begin{aligned}
 g &= \eta [\vartheta_3(\omega t, \tau_1) \vartheta_4(\gamma x, \tau) \\
 &\quad + i \vartheta_1(\omega t, \tau_1) \vartheta_2(\gamma x, \tau)] \exp(ipt), \tag{20}
 \end{aligned}$$

$$\begin{aligned}
 f &= \vartheta_4(\omega t, \tau_1) \vartheta_2(\gamma x, \tau) \\
 &\quad + \vartheta_2(\omega t, \tau_1) \vartheta_4(\gamma x, \tau),
 \end{aligned}$$

and so on. For the sake of simplicity, the tedious computations are omitted here. It is noted that these two types of periodic solutions have also two independent periods in the spatial and temporal directions.

4. Jacobi Elliptic Function Expressions and Long Wave Limit

In order to analyze the periodic property by some figures, we may first convert solution (18) into Jacobi elliptic function expressions. Together with (6) and (7), solution (18) can be expressed as rational forms of Jacobi elliptic functions:

$$\begin{aligned}
 u(x, t) &= \pm \sqrt{\frac{\lambda^2(1-2k^2) + \rho^2(1-2k_1^2)}{\delta k_1} \frac{D_1 C + i S_1}{1 + C_1 C}} \exp(ipt), \\
 v(x, t) &= v_0 + 2\lambda^2 \left(1 - k^2 - \frac{E}{K}\right) \\
 &\quad + \frac{2\lambda^2 [k\sqrt{1-k^2}(C^2 - C_1^2) + (2k^2 - 1)CC_1]}{(1 + C_1 C)^2}, \tag{21}
 \end{aligned}$$

where λ is an arbitrary constant, and the parameters $\rho, p, \Lambda,$ and v_0 are given by

$$\begin{aligned}
 \rho^2 &= \lambda^2 [2k^2 - 1 + 2kk_1(1 - k^2)^{1/2} (1 - k_1^2)^{-1/2}], \\
 p &= \frac{\lambda^2}{\rho} k(1 - k^2)^{1/2} (1 - k_1^2)^{-1/2}, \\
 \Lambda &= \lambda^2 \left(1 - \frac{2E}{K}\right) + \rho^2 \left(1 - \frac{2E_1}{K_1}\right), \\
 v_0 &= \rho^2 \left(\frac{(1 - k_1^2)}{2} - \frac{E_1}{K_1}\right) - \lambda^2 \left(1 - k^2 - \frac{E}{K}\right) - \frac{\beta + p^2 + \Lambda}{2}, \tag{22}
 \end{aligned}$$

which indicates that the period k in the spatial direction and the period k_1 in the temporal direction are related by

$$2k^2 - 1 + 2kk_1(1 - k^2)^{1/2} (1 - k_1^2)^{-1/2} > 0. \tag{23}$$

In (21) and (22), the complete elliptic integrals E and E_1 are defined by

$$E = \int_0^{\pi/2} \sqrt{1 - k^2 \sin^2 \xi} d\xi, \quad E_1 = \int_0^{\pi/2} \sqrt{1 - k_1^2 \sin^2 \xi} d\xi. \tag{24}$$

From (21), it is easy to check that

$$\begin{aligned}
 u(x, t) &= u\left(x + \frac{4K}{\lambda}, t\right), & |u(x, t)| &= \left|u\left(x, t + \frac{4K_1}{\rho}\right)\right|, \\
 v(x, t) &= v\left(x + \frac{4K}{\lambda}, t\right), & v(x, t) &= v\left(x, t + \frac{4K_1}{\rho}\right), \tag{25}
 \end{aligned}$$

which implies that the solution $(|u|, v)$ is periodic in the x -direction with a period $4K/\lambda$ and the t -direction with a period $4K_1/\rho$.

By selecting appropriate parameter values in (21), the interactions of doubly periodic waves are shown in Figures 1 and 2. It is clearly seen that $(|u|, v)$ is periodic in the x -direction and the t -direction. For the solution in [39], the periodic waves are both bell shaped in the spatial and temporal directions. However, with regard to solution (21), the periodic waves are of different shapes in the spatial and temporal directions.

When the modulus $k \rightarrow 1,$ $(\operatorname{sn}x, \operatorname{cn}x, \operatorname{dn}x) \rightarrow (\tanh x, \operatorname{sech}x, \operatorname{sech}x).$ And if $k \rightarrow 0,$ $(\operatorname{sn}x, \operatorname{cn}x, \operatorname{dn}x) \rightarrow (\sin x, \cos x, 1).$ Therefore, the long wave limits of the periodic wave solutions can be readily obtained. A long wave limit of the solution (21) can be taken by assuming

$$\frac{\sqrt{k}}{(1 - k_1^2)^{1/4}} = m, \quad 0 < m < 1. \tag{26}$$

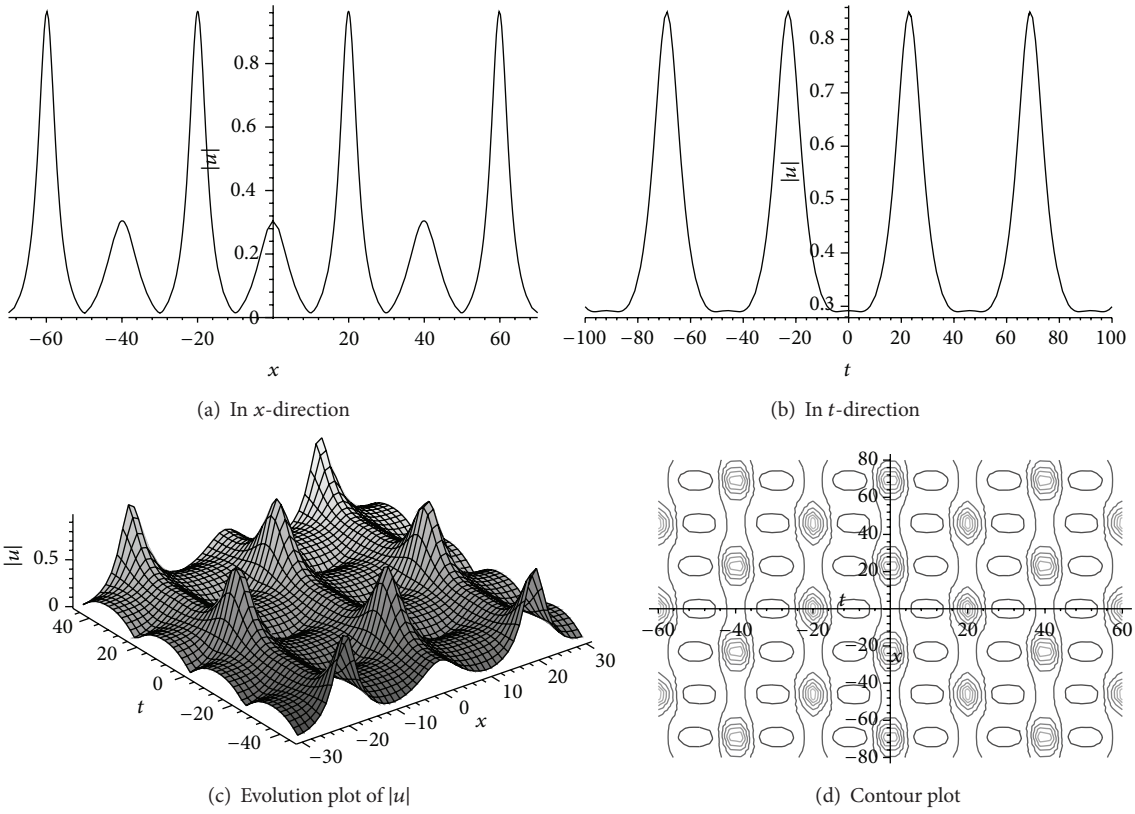


FIGURE 1: Doubly periodic wave solution $u(x, t)$ given by (21) with $\delta = 0.2$, $\beta = 0.4$, $\lambda = 0.2$, $k = 0.8$, and $k_1 = 0.2$.

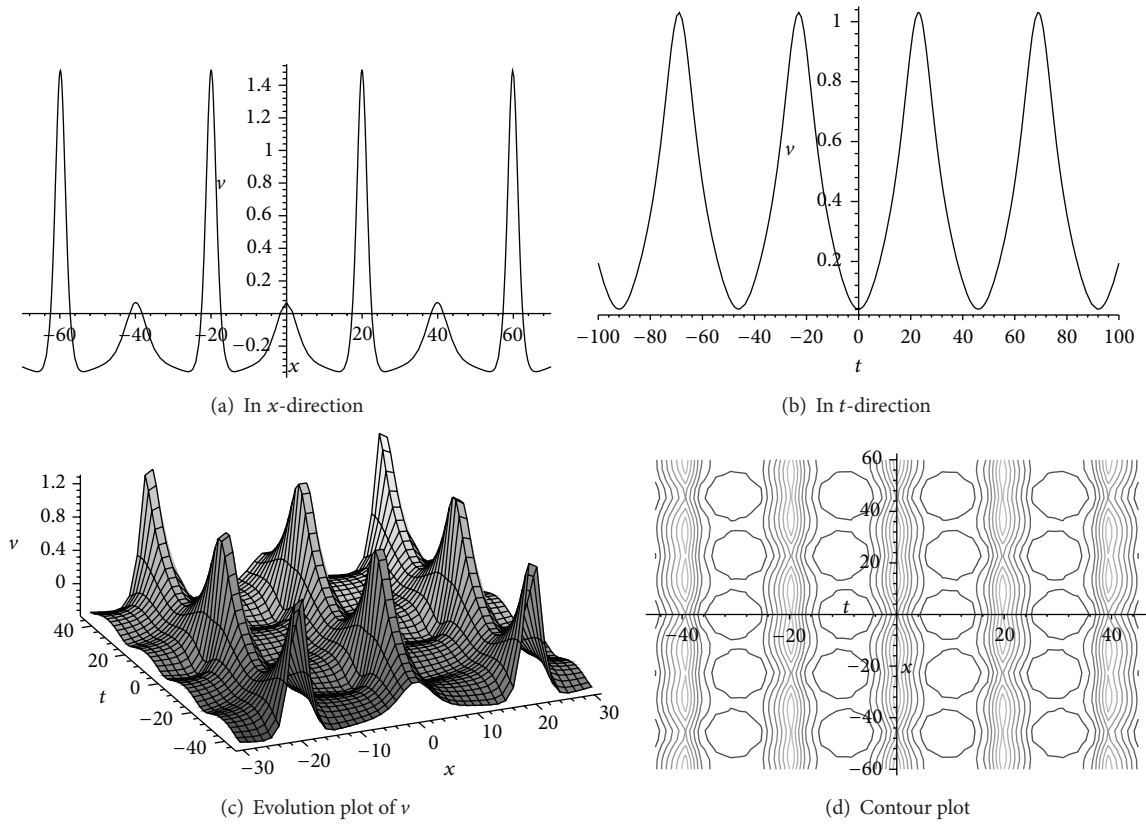


FIGURE 2: Doubly periodic wave solution $v(x, t)$ given by (21) with $\delta = 0.2$, $\beta = 0.4$, $\lambda = 0.2$, $k = 0.8$, and $k_1 = 0.2$.

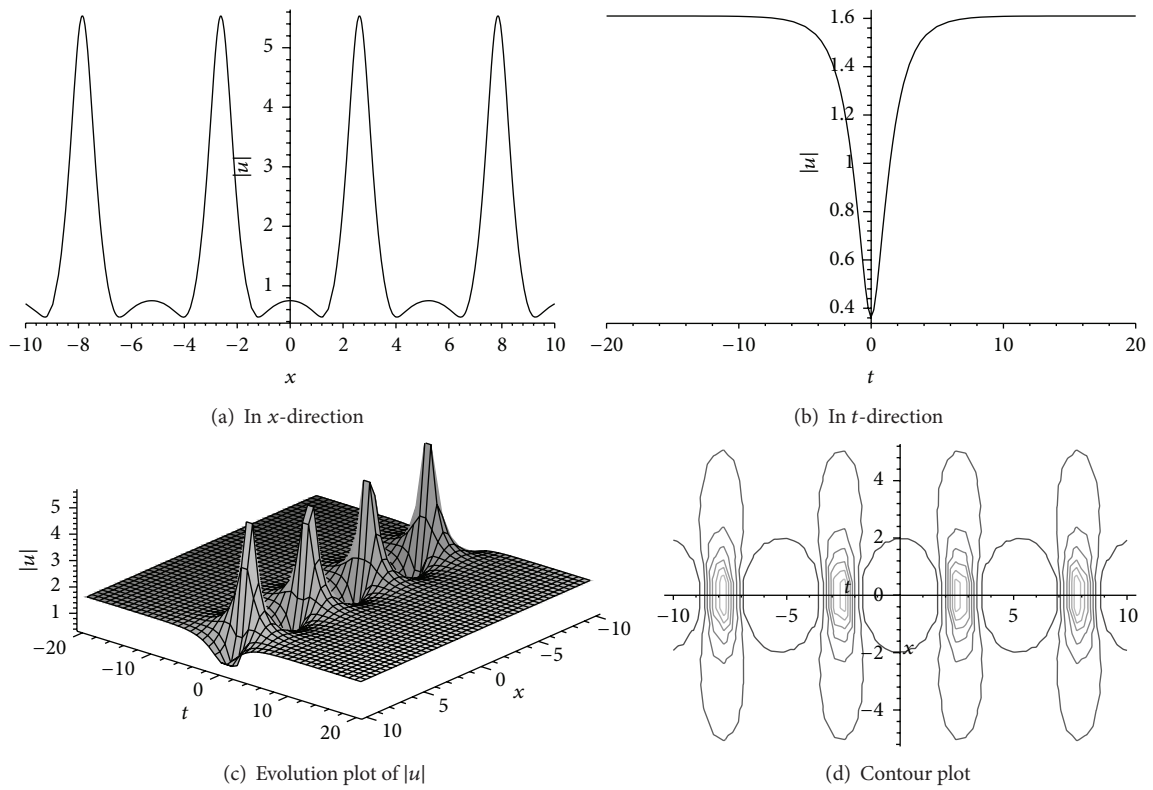


FIGURE 3: The periodic-solitary wave solution $u(x, t)$ given by (27) with $\delta = 0.4$, $\beta = 2.0$, $\lambda = 1.2$, and $m = 0.8$.

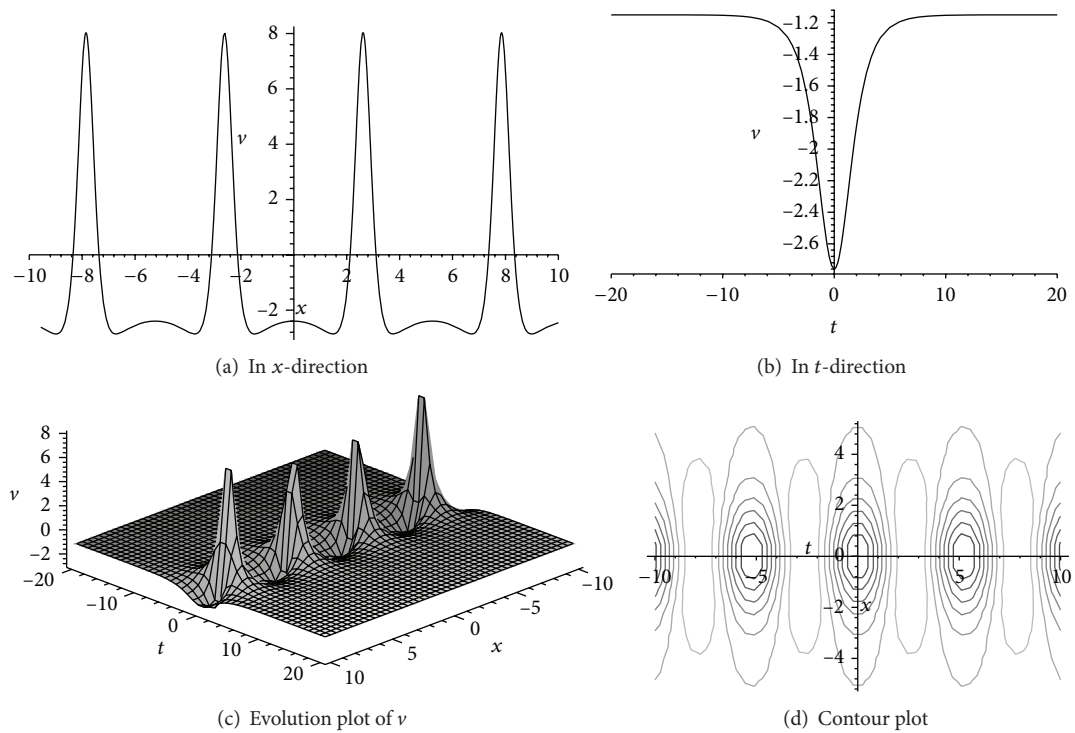


FIGURE 4: The periodic-solitary wave solution $v(x, t)$ given by (27) with $\delta = 0.4$, $\beta = 2.0$, $\lambda = 1.2$, and $m = 0.8$.

When the modulus $k \rightarrow 0$ and $k_1 \rightarrow 1$, one obtains a new periodic-solitary wave solution as follows:

$$u = \eta \left[\frac{m \cos(\lambda x) \operatorname{sech}(\rho t) + i \tanh(\rho t)}{1 + m \cos(\lambda x) \operatorname{sech}(\rho t)} \right] \exp(ipt),$$

$$v = -\frac{\beta}{2} - \frac{\lambda^2 (5m^4 - 6m^2 + 2)}{2(2m^2 - 1)} - \frac{2m\lambda^2 \operatorname{sech}(\rho t) (m \operatorname{sech}(\rho t) + \cos(\lambda x))}{(1 + m \cos(\lambda x) \operatorname{sech}(\rho t))^2}, \quad (27)$$

where the parameters ρ , p , and η are given by

$$\rho = \lambda \sqrt{2m^2 - 1}, \quad p = \frac{\lambda m^2}{\sqrt{2m^2 - 1}},$$

$$\eta = \lambda \sqrt{\frac{2(1 - m^2)}{\delta}}. \quad (28)$$

With proper selections of the values of β , δ , λ , and m , the interactions of periodic solitary waves (27) are shown in Figures 3 and 4. The solution $(|u|, v)$ displays the feature of a dark soliton in the t -direction; the cosine function causes periodic modulation and thus it is periodic in the x -direction.

With the aid of the computer algebra software **Maple**, the validity of the new solutions (18) and (27) are verified by putting them back into the original systems (1).

5. Conclusions

The combination of the Hirota bilinear method and theta function identities is demonstrated to be a powerful tool in finding periodic waves for the coupled Higgs field equation. As a result, we have derived a new kind of doubly periodic standing wave solutions for the coupled Higgs field equation, which is different from those of the known solutions reported in the literature. The interaction properties of periodic-periodic waves and periodic-solitary waves are analyzed by some figures.

The key of the combination method is that the solutions are supposed as rational expressions of elliptic functions of different moduli, which should be applicable to other nonlinear evolution equations or systems with bilinear forms in mathematical physics. The doubly periodic solutions will prove to be beneficial and instructive in modeling and understanding nonlinear phenomenon.

Conflict of Interests

The author declares that there is no conflict of interests regarding the publication of this paper.

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