Ferroelectric Control of the Conduction at the LaAlO₃/SrTiO₃ Hetero-interface

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Abstract

The LaAlO₃/SrTiO₃ (LAO/STO) interface serves as a model system in which a highly mobile quasi-twodimensional electron gas (2DEG) forms between two band insulators [1,2], exhibiting 2D superconductivity [3] and unusual magnetotransport properties [4]. In the study reported here, we added a ferroelectric Pb($Zr_{0.2}Ti_{0.8}$)O₃ (PZT) layer near the LAO/STO interface. The ferroelectric polarization of the PZT layer serves as a control parameter to modulate the 2DEG conducting behavior. The as-grown polarization (P_{up} state) leads to charge depletion and enhances the conduction. Switching the polarization direction (P_{down} state) results in charge accumulation and enhances the conduction at the LAO/STO interface. The origin of this modulation is attributed to a change in the electronic structure due to the ferroelectric polarization states, evidenced by X-ray photoelectron spectroscopy (XPS). Control of the conduction at this oxide interface suggests that the concept can be generalized for other oxide systems to design functional interfaces.

Keywords – LaAlO₃/SrTiO₃, 2DEG, ferroelectric Pb(Zr_{0.2}Ti_{0.8})O₃, transport, XPS.

Introduction

The LAO/STO interface is served as a model system in which a highly mobile quasi-two dimensional electron gas (2DEG) forms between two band insulators [1,2], exhibiting 2D superconductivity [3]. The ferroelectric polarization of PZT layer serves as a control parameter to modulate the 2DEG conducting behaviors. The as-grown polarization (Pup state) leads to charge depletion and consequently a low conduction. Switching the polarization direction (P_{down} state) results in a charge accumulation and enhances the conduction at the interface of LAO/STO. The origin of this modulation is attributed to a change in the electronic structure due to the ferroelectric polarization states, evidenced by x-ray photoelectron spectroscopy (XPS). Control of the conduction at this oxide interface suggests that the concept can be generalized for other oxide systems to design functional interfaces.

Experiments

In order of to understand the influence ferroelectricity LAO/STO on the interface experimentally, we firstly carried out electrical transport and XPS measurements on the heterostructure. In the ferroelectric PZT/LAO/STO heterojunction with a naturally upward PZT layer, the polarization was reversed by the scanning probe technique (probe voltage set to 8 V) to obtain a downward polarization area of about 1 mm² that is larger than x-ray beam size of about 400 µm². This ferroelectric-pattern assisted spectral technique makes us to study the polarization reversal

using the same sample, not having to compare the polarization state in two different samples, avoiding all the discussions regarding sample quality and other differences.

Results

carried out electrical transport We firstly measurements on the heterostructure (Figure 1a), in which the PZT layer with the spontaneous polarization P_{PZT} functions as the polarized dielectric slab to modulate the conduction of LAO/STO heterointerface. Figure 1b shows the sheet resistance versus temperature (R-T curves) of the PZT/LAO/STO samples with various PZT laver thicknesses (0~40 nm), while the LAO thickness was fixed (6 u.c.). The R-T curves for these samples show that the sheet resistance at room temperature is low (50 K Ω /sq) and decreased with temperature, showing the metallic behavior. For the samples with PZT on top, the ferroelectric effect sets in and the sheet resistance starts increasing with PZT thickness, showing the strong impact of the intrinsic polarization (Pup) of PZT on electron conduction at the interface. This is anticipated since the ferroelectric field effect provides one more degree of freedom to compensate the charge unbalance at interface¹⁵. We the also conducted transport measurements on the samples with different LAO thickness, while PZT thickness was kept constant (20 nm) (Figure 1c). Several studies suggested a critical LAO thickness (4 u.c.) for the formation of 2DEG at this interface⁷. Under such a circumstance, all the samples with LAO above a critical thickness behave like a metal.

The sheet resistance is decreased with lowering the temperature and increased when the thickness of LAO is reduced.



Fig. 1. Transport measurement of PZT/LAO/STO devices.

As depicted in Figure 2a, positive (Pdown state) and negative (Pup state) bound charge sheets at bottom PZT layer affected at the conducting LAO/STO heterointerface, leading to an accumulation or depletion of free electrons in STO layer. The schematics (Figure 2a) also shows the ferroelectric-pattern-assisted XPS technique on PZT (3.5 nm)/LAO (4 u.c)/STO heterjunction, the binding energy of the Sr core-level from buried layer thus decreases $(E_{B}^{Sr} - eV_{up})$ or increases $(E_{B}^{Sr} + eV_{up})$ depending on the downward and upward polarization directions, respectively, by comparing the binding energy of the Sr core-level ($E_{\rm B}^{\rm Sr}$) from LAO/STO heterojunction. Evidently, the concept of that interface capacitor is also confirmed from the lack of change of the core-level emissions from the top PZT (Pb 4f core-level) and LAO (La 4d and Al 2s core-levels) layers. Figure 2b shows the characteristic core-level photoelectron spectra in the ferroelectric PZT/LAO/STO and conducting LAO/STO heterojunction samples. The binding energy of Sr $3d_{5/2}$ core-level shifts associated with the potential drops of the PZT-modulated LAO/STO interface capacitors were obtained at eV values of +0.15 eV ($P_{up}\xspace$ state) and –0.1 eV $(P_{\text{down}} \mbox{ state}),$ and the core-level difference between Sr $3d_{5/2}$ and La $4d_{5/2}$ (ΔE_{CL}) were determined of 31.15 eV (Pup state) and 30.90 eV (Pdown state).



Figure 2 X-ray Photoemission spectroscopy study on the upward / downward ferroelectric polarization.

In conclusion, we have demonstrated that the ferroelectric polarization can tune and modulate the conduction at the LAO/STO hetero-interface. The XPS results reveal the electrostatic predictions of the conducting state modulation, and demonstrate the possibility of nonvolatile control and provide compelling evidence in favor of ferroelectric doping at the conducting polar-nonpolar oxide heterointerface.

Acknowledgments

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