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### Exploiting the Asymmetric Energy Barrier in Multi-stable Origami to Enable Mechanical Diode Behavior

A Thesis Presented to the Graduate School of Clemson University

In Partial Fulfillment of the Requirements for the Degree Master of Science Mechanical Engineering

> by Nasim Baharisangari August 2020

Accepted by: Dr. Suyi Li, Committee Chair Dr. Oliver Myers Dr. Lonny Thompson

#### Abstract

Recently, multi-stable origami have drawn many attentions for their potential applications in multi-functional structures and material systems. Especially, origami folding is essentially a three-dimensional mechanism, which induces unorthodox properties that distinguish this mechanism from its traditional counterparts. This study proposes a multi-stable origami cellular structure that can exhibit mechanical diode behavior in compression. Furthermore, with a small variation in the unit cell of the proposed structure, a extension diode can be achieved. Such structures consist of many stacked Miura-ori sheets, and can be divided into unit cells that pose two different stable configurations. To understand and elucidate the underlying mechanisms, two adjacent unit cells were considered as the most fundamental constituents of the cellular structures that display the desired diode behavior. This study examines how folding can impose a kinematic constraint onto the deformation of these two dual cell chains via estimating the elastic potential energy landscapes of two dual assemblies. For the compression diode, this folding-induced constraint increase the energy barrier for compressing from a certain stable state to another, however, the same constraint does not increase the energy barrier of the opposite extension. Thus, one should apply a large force to compress the chain, but a small force to extend it. As a result, a compression mechanical diode is achieved. This constraint acts the opposite way in extension diode. Then, four prototypes were fabricated to experimentally validate the analytical results. The results of this study can open new avenues towards multi-functional structure and materials systems capable of motion rectifying, wave propagation control, and even mechanical computation.

#### Keywords: Origami, Multi-stability, Mechanical Diode

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## Nomenclature

$a_I$	-	Muira-ori sheet I edge	
$a_{II}$	-	Muira-ori sheet II edge	
$b_I$	-	Muira-ori sheet I edge	
$b_{II}$	-	Muira-ori sheet II edge	
$\gamma_I$	-	Muira-ori sheet I sector angle	
$\gamma_{II}$	-	Muira-ori sheet II sector angle	
$\theta_I$	-	Muira-ori sheet I folding angle	
$\theta_{II}$	-	Muira-ori sheet II folding angle	
$\phi_i$	-	Dihedral angles	
$psi_I$	-	Muira-ori sheet I spine angle	
$\psi_{II}$	-	Muira-ori sheet II spine angle	
L	-	Individual unit cell length	
$l_I$	-	Sheet I length	
$l_{II}$	-	Sheet II length	
$l_c$	-	Connecting sheet length	
$L^A$	-	Cell A length	
$L^B$	-	Cell B length	
$L^{o}$	_	outer connecting sheet length	

- $L^t$  Dual cell-chain length
- $\prod^{A}$  Cell A elastic potential
- $\Pi^B$  Cell B elastic potential
- $\prod^{o}$  Outer connecting sheet elastic potential
- $\prod^{t}$  Dual cell chain total elastic potential
- $k_i$  crease stiffness
- $k^*$  kinematic constraint stiffness
- $A_1$  Sheet I surface area
- $A_2$  Sheet II surface area
- $A_3$  Connecting sheet surface area
- $F_c^A$  Cell A compression reaction force
- $F_e^A$  Cell A extension reaction force
- $F_c^B$  Cell B compression reaction force
- $F_e^B$  Cell B extension reaction force
- $F_c^{C} \quad$  Dual cell chain with crease-connection compression reaction force
- ${\cal F}^{\cal R}_c$  Dual cell chain with rod-connection compression reaction force
- $F_e^C$  Dual cell chain with crease-connection extension reaction force
- $F_e^R$  Dual cell chain with rod-connection extension reaction force

## Chapter 1

## Introduction and Literature Review

## 1.1 Introduction

A structure or material system is considered multi-stable if they exhibit more than one stable equilibrium (or stable state) within the deformation range so that each stable sate corresponds to a potential energy minimum [28]. Multi-stability can be used as an alternative mechanism in enabling a wide variety of functionalities such as stiffness adaptation [33], energy harvesting [32] [7]. Origami the ancient art of folding paper into aesthetic shapes has drawn the attention of the researches from various fields like aerospace [45] [27], architecture [43], robotics [29], and biomedical [18] [17] industries.

Recently, origami capability to create programmable and re-programmable systems

that can change shape, function and, mechanical properties has opened up innovation doors [10] from macro scale to nanoscale. For example, a sheet, with pre-defined fold lines, capable of reshaping autonomously into different 3D structures was created by Hawkes *etal*. [14], or Marras *etal*. [25] showed that with folding DNA nano-scale mechanisms with programmable mechanical functions can be built.

These structures exhibit unique mechanical properties such as negative Poisson's ratio [22] [39], discrete stiffness jumps [20] [3], elastic multistability [44] [16]. Recent studies have shown that origami-based cellular structures and materials are promising platforms to achieve bi-stability [24]. If the crease bending stiffness between two adjacent sheets in the cellular structure differ notably [27] [43], or its facets are deformed between different configurations [31] [21], the origami structure exhibit multiutility. Moreover, utilizing origami, a three-dimensional shape transformation mechanism, leads to obtaining multi-stability in higher dimensions [41]. This privilege of origami over currently employed bi-stable mechanisms such as the curved beams or their close relatives, prestressed bilayer shells and, axially constrained springs [28] [12], open avenues to create adaptive materials and functional materials [41]. The infinite possibilities of folding combinations [8] [15], and robust manufacturability [35] [30] of folded sheets make them a high potential candidate to construct multi-functional materials. One of the most used multi-stable origami structures is the stacked Miura-ori which is constructed by assembling geometrically compatible Miura-ori sheets along their creases [38]. In stacked Miura-ori, the multi-stability is induced by the Miura-ori sheets considerable stiffness difference [23].

Through the transition to obtaining multi-stability in higher dimensions with origami's 3D nature, the stacked Miura-ori has been shown to exhibit rapid deformation via pressure-induced snapping [43] and elastic modulus programming [26]. In one study,

it is observed that the multi-stable stacked Miura-ori exhibits unique asymmetric energy barriers and mechanical diode behavior [41]. Folding induces a kinematic constraint that causes a significant increase in the energy barrier when the structure is being stretched while the required energy for compression does not experience a notable change. Thus, a large amount of force must be applied to extend the stacked Miura-ori, but only a small force to compress it.

In this design, static diode behavior is observable only in the extension direction. This finding has brought up this question that how we can come up with a design to see the diode behavior in compression, and if it is possible to transform it into extension diode with a small change in the designed structure with the existing constituents.

The goal of this research to propose a cellular origami structure capable of exhibiting static diode behavior in compression. The proposed origami unit cell can be counted as a variation of the traditional stacked Miura-ori (Figure 1.1). This dual cell assembly is compared to an electronic diode or a mechanical ratchet. The three structures are designed to rectify the operating direction; The electricity current flow is one-way in the electronic diode, or the rotational movement of the ratchet is unidirectional, and finally the compression diode dual assembly facilitates the deformation in extension direction only.



Figure 1.1: The mechanical compression diode design. The dual cell chain is easy to extend but hard to compress. There can be an analogy between this structure and, electronic diode and, mechanical ratchet.

By using the rigid facet and spring hinge assumption the energy landscape of a newly designed origami cellular structure is calculated and it has shown the desired energy asymmetric barrier and static diode behavior. The calculated energy landscape shows that the kinematic constraint induced by folding causes a significant increase in the energy barrier in shifting between two consecutive stable configurations in compression direction, but no notable change in the energy barrier in the opposite switch was noticed. An experimental examination has been conducted to validate the theoretical results.

Followingly, the theoretical model of the extension diode was developed based on

the compression diode model (Figure 1.2). The energy landscape of the derivative stacked origami was calculated based on the same assumption of rigid foldability. The energy landscape of the extension diode showed the expected asymmetric energy barrier in the extension direction. Meaning that one should apply a large amount of force to stretch the cellular structure, but a small amount of force to compress it. The attained theoretical results were accompanied by experimental examination, and the extension diode behavior was observed.



Figure 1.2: The mechanical extension diode design. The dual cell chain is easy to compress but hard to extend. There can be an analogy between this structure and electronic diode and, mechanical ratchet.

One of the potential applications of this static diode is to be deployed in mechanical programming. The current transistor-based computing circuits use multiple interconnected transistors to create a single Boolean logic gate. These electronic computational components cannot function properly in harsh environments and because of excessive heat dissipation, they demand involved thermal management. Besides, transistor circuits are not capable of dynamically reconfigure their functionality in real-time [5].

The mechanical computing is being investigated by many research groups due to its advantages over its electronic counterparts. For example, in comparison with electronic parts, mechanical parts can resist much higher temperature and radiation exposure [4] [6]. Another advantage of mechanical logic devices they don't need power source because they use energy in mechanical form [36] [40]. Moreover, studies on reversible-computing have suggested that designing a mechanical logic system with small energy dissipation is theoretically possible [19] [11]. Currently, several mechanical computations systems have been introduced. For example, Yuanping Song *etal.* performed Boolean computations based on the mechanical forces and displacements of multi-stable micro-flexures [40]. Raney *etal.* and the coworkers have architected a medium composed of elastomeric bistable beams elements connected by elastomeric linear springs that propagate mechanical signals. This architected structure can be used to design mechanical diode and logic gates [34]. Origami structures have shown a rich potential to be adopted to soft actuation materials and mechanisms [13].

Another potential application of origami mechanical diode is to be integrated into soft robots and materials to serve different tasks. For example, The central unit processing units in soft robots that manage the decision step in the interaction process of the robot with the environment are composed of rigid electronics. Integrating these stiff parts in soft robots is not thoroughly compatible with the compliant body of soft robots. Treml *etal*. and his coworkers have developed a mechanical computation unit with an origami waterbomb as the experimental platform to be Incorporated in soft robots as an solution to the mentioned problem [42]. In what follows, Chapter 2 discusses the mechanics modeling and the theoretical analysis of the compression and tension diodes; Chapter three presents the experimental validation of the theoretical results proposed cellular designs. Chapter four investigates and optimization study on the compression diode unit cell, and eventually in Chapter five concludes this study with summary and future work.

## Chapter 2

# Mechanics Modeling and Theoretical Analysis of the Mechanical Diode

## 2.1 Design of The New Cellular Origami Structure

In this study, a new multistable cellular origami structure is introduced. This unit cell is fabricated by stacking geometrically compatible Miura-ori sheets and zigzag shaped "connect sheets" in an alternating arrangement (Figure 2.1.a). By connecting two unit cells via a connecting sheet, the most fundamental multi-stable structure that can exhibit diode behavior is obtained (Figure 2.1.b). The designed unit cell is essentially a variation of a classical stacked Miura-ori [38].

In the unit cell discussed here, the orientation of the Miura-ori with respect to each



Figure 2.1: Design of the new multistable stacked origami cellular structure. (a) An overview showing the alternating sequence of different Miura-ori sheets and zig-zag "connect sheets".

other is flipped. More clearly stated, the Miura-ori sheet with the bigger dimension in one edge, also referring to as sheet II, is reversed in the new design (Figure 2.2).

The new unit cell still follows the rigid-folding kinematics of traditional Miura-ori [38]. The crease design of a unit cell is determined by crease lengths  $(a_I, b_I, a_{II}, b_{II}, l_c)$  and the sector angles  $(\gamma_I, \gamma_{II})$  (Figure 2.3.a). Here, subscript I and II denote the two different Miura-ori sheets in a unit cell and  $l_c$  is the length of the connecting sheet. To satisfy the geometric compatibility the following restraints [38] should be imposed on these parameters values :

$$b_{\rm II} = b_{\rm I},\tag{2.1}$$

$$\frac{\cos\gamma_{\rm II}}{\cos\gamma_{\rm I}} = \frac{a_{\rm I}}{a_{\rm II}}.$$
(2.2)



Figure 2.2: (a) Miura-ori sheet II is flipped in the new compression diode unit cell. (b) The arrangement of sheet I and II with respect to each other in traditional stacked Miuor-ori

To describe the external geometry of a unit cell during rigid-folding, one can use dihedral folding angles  $\theta_I$  and  $\theta_{II}$  defined between the facets of the two Miura-ori sheets and the x-y reference plane, respectively (Figure 2.3.b).

In the geometric design, it assumed that the unit cell ideally satisfies the rigid-folding condition, which is essentially a one-degree-of-freedom motion [9]. This condition is stated by the following relationship between the two sector angles and the folding dihedral angles [38]:

$$\cos\theta_{\rm I}\tan\gamma_{\rm I} = \cos\theta_{\rm II}\tan\gamma_{\rm II}.\tag{2.3}$$

The summation of the different components of the unit cell gives the total length of the unit cell.



Figure 2.3: Detailed design of a unit cell in this study. (a)  $\Phi_i$  (I = 1...6) are unique dihedral angles between two adjacent facets along the difference creases.  $\psi_i$ is the spine angle, which is also the dihedral angle of in the connect sheet. The two drawings on the right show the design of Miura-ori sheets

$$L = l_{\rm I} + l_{\rm II} + l_c, \tag{2.4}$$

Where  $l_I$  and  $l_{II}$  are the length of the two constituent Miura-ori sheets respectively.

To describe dihedral folding angles between the facets in the connect sheet, a spine angle can be defined [9]:

$$\psi = 2\tan^{-1}(\cos\theta_{\rm I}\tan\gamma_{\rm I}),\tag{2.5}$$

At it was mentioned, it is assumed that the unit cell facets are ideal rigid the crease lines act as perfect hinges with prescribed torsional stiffness. This assumption satisfies the rigid-folding condition kinematics. Thus, the total elastic potential energy of the structure can be calculated using the following equation [1]:

$$\Pi = \frac{1}{2}k_i(\varphi_i - \varphi_i^o)^2 + \frac{1}{2}k_c(\psi - \psi^o)^2, \qquad (2.6)$$

Where  $_i$  is the dihedral crease opening angle denoted in Figure 2.3.a; These angles measure the angles between intersecting planes forming the compression diode unit cell's geometry, and  $\varphi_i^o$  is the initial value of the corresponding dihedral angle (it is worth to remind that all the angles defining the unit cell's geometry are functons of the folding angle  $\theta_i$ .).  $k_i$  is the corresponding torsional spring stiffness in the connect sheets. The initial stress-free configuration angles are denoted by subscripts o. The crease opening angles are the function of independent variable  $\theta_I$  and can be described using the following equations: (equations 2.8, 2.9 and 2.10 are adapted from previous publications [23]):

$$\varphi_1 = \pi - 2\theta_{\mathrm{I}},\tag{2.7}$$

$$\varphi_2 = 2\sin^{-1} \left( \frac{\cos \theta_{\rm I}}{\sqrt{1 - \sin^2 \theta_{\rm I} \sin^2 \gamma_{\rm I}}} \right),\tag{2.8}$$

$$\varphi_3 = \pi - 2\cos^{-1} \left( \tan \gamma_{\rm II} \tan^{-1} \gamma_{\rm I} \cos \theta_{\rm I} \right), \qquad (2.9)$$

$$\varphi_4 = 2\sin^{-1} \left( \frac{\sin \gamma_{\rm I}}{\sin \gamma_{\rm II}} \sin \frac{\varphi_2}{2} \right), \qquad (2.10)$$

$$\varphi_5 = \frac{\pi}{2} + \theta_{\rm I}.\tag{2.11}$$

$$\varphi_6 = \frac{\pi}{2} - \theta_{\rm II}.\tag{2.12}$$

Although the torsional springs added to the creases are linearly elastic as can be seen in equation 2.6, the correlations between folding and external deformation are geometric and strongly nonlinear. The desired diode behavior originates from this nonlinearity.  $k_i$  and  $k_I$  are the crease torsional stiffness per unit length of the Miuraori sheet I and II, respectively, and  $k_c$  is the crease torsional spring stiffness per unit length of the connecting sheet. The stiffness coefficients in equation 2.6 are  $k_1 = 2k_Ib_I, k_2 = 2k_Ia_I, k_3 = 2k_{II}b_I, k_4 = 2k_{II}a_{II}, k_5 = 4k_cb_I, k_6 = 4k_cb_I$ , and  $k_c = 2k_cl_c$ , where the numerical coefficients in these equations show determines the similar creases in one unit cell.

In order to achieve bi-stability in a stacked Miura-ori unit cell , the stiffness of the larger sheet II should be much higher than the crease stiffness of the sheet I and the connecting sheet (also known as  $k_{II} > k_I$  and  $k_c$ ). Moreover, the initial stress-free folding configuration should drift from 0 [23]. As it is shown in Figure 2.3-b different values of  $\theta_I$  can be chosen as the initial value of this angles to enable bi-stability except for  $\theta_I = 0$ . Figure 2.4 illustrates the energy landscape of two unit cells ( referred to as cell A and B hereafter ) of the geometric parameters value of  $a_I = b_I = 2 \text{cm}, a_{II} = 1.25 a_I, \gamma_I = 45^\circ, l_c = 2.5 \text{a}_I, k_I = k_c$ , and  $k_{II} = 20 k_I$  (equation 2.6). The initial dihedral angle of cell A is chosen to be 60 and cell B to be -60 degrees.

The two potential energy wells of each cell (Figure 2.4) exhibit the bi-stability of this group of geometric design parameters. In purpose of more clarity, the positive folding angle of sheet I is denoted as state (1) and the negative stable configuration as (0) so that the unit cell has the shortest length L at state (0). Throughout the entire



Figure 2.4: The energy landscape of the two unit cells used in this study.

thesis, these design values are kept the same consistently, unless noted otherwise. The initial dihedral angle of cell A is chosen to be 60 and cell B to be -60 degrees. These stress-free configurations dictate the force relation between the individual cells as follows:  $F_c^A < F_c^B$ , and  $F_e^B < F_e^A$ .

After formulating the unit cell external geometry and potential energy, the overall energy and dimension of the dual cell assembly can be calculated as:

$$\Pi^{t} = \Pi^{A} + \Pi^{B} + \Pi^{0}, \qquad (2.13)$$

$$L^{t} = L^{A} + L^{B} + L^{0}. (2.14)$$

 $\Pi^A$ , and  $\Pi^B$  are the strain energy of the unit cell A and B with the definition stated in equation 2.6.  $\Pi^0$  is the strain energy of the connecting sheet between the two unit cells and defined as:

$$\Pi^{0} = \frac{1}{2} k^{*} (\psi^{A} - \psi^{B})^{2}, \qquad (2.15)$$

Where  $k^*$  is the constraint stiffness of the "connecting sheet". This parameter is the

key element of this study since it quantifies the strength of kinematical constraint induced by folding. Ideally, if the rigid-folding assumption is observed (aka. all facets in the dual cell chain are fully rigid and all creases act perfectly as hinges), the spine angles of the two unit cells should be equal. ( $\psi_A = \psi_B$ ). In this way, the admissible deformations of the dual-cell chain are restricted to the "kinematic paths" shown in Figure 2.5.a.



Figure 2.5: Kinematic properties of the compression diode structure due to the folding induced constraint (or the lack of). (a) Admissible deformation of the dual cell assembly. The two kinematic paths based on ideal rigid-folding condition are shown by the solid and dashed curves. The gray area represents deformations that are not kinematically admissible. (b) The geometry of the dual cell assembly at different locations along these to kinematic paths.

In ideally rigid-folding condition, one possible path would be  $\theta_I^A = \theta_I^B$ , and the other path be  $\theta_I^A = -\theta_I^B$ . However, the facets are not ideally rigid, and the creases do not behave like perfect hinges. More specifically, the facets have small bending and creasing wrapping will take place. Thus, there would be some mismatch between the two spine angles of the two unit cells. In conclusion, the configuration of the dual-cell assembly can occur at any point within the parallelogram shown in Figure 2.5.a. This deviation from ideal rigidity can apply additional elastic potential energy that can be characterized by the constraint stiffness  $k^*$ .

In the next sections, first, the nonlinear elastic behavior of dual-cell assembly in the absence of the kinematic constraint stiffness  $(k^* = 0)$  is examined, and then the situation at which the kinematic strain energy is added to the system is studied.

#### 2.2 Diode effect in compression

Figure 2.6.a illustrates the total energy landscape of the dual cell chain according to equation 2.13 with  $k^* = 0$ . This scenario represents a hypothetical case in which the sheet that connects the two cells are soft so that it does not provide any resistance to the mismatch between the spine angles of the two cells ( $\psi^A$  and  $\psi^B$ ). The "equilibrium paths" corresponding to the potential energy minima at a given total length can be determined, and the dotted line shows the potential energy maxima at that length. During deformation (changing from the minimum length to maximum length) the dual cell assembly would pave these minima paths. Here, the continuous equilibrium path that connects the three stable states of "0-0", "0-1", and "1-1" is of interest. The energy landscape of the dual-cell assembly along this path is plotted in Figure2.6.b.

The extension energy barrier ( $\Delta E$ ) for shifting from "0-0" stable state to "0-1" stable state and the compression energy barrier for the opposite switch  $\Delta C$  can be seen in Figure 2.6. The corresponding reaction force can be calculated as the variation of total potential energy with respect to the change in total length:

$$F = \frac{\partial \Pi^{t}}{\partial L^{t}}.$$
(2.16)

The reaction force corresponding to the continuous equilibrium path is shown in Figure 2.6.c. Based on this plot, two important forces can be calculated. One is the critical reaction force  $(F_e)$  during switching from "0-0" stable configuration to "0-1" stable configuration. The other important force  $(F_c)$  is the critical reaction force to make the opposite switch (from "0-1" to "0-0") happen. Essentially,  $(F_e)$  is the force required to stretch the dual-cell chain from "0-0" to "0-1", and  $(F_c)$  is the needed amount of force to compress the structure back to "0-0" state.

The discussed scenario above showed a hypothetical case in which the connecting sheet between the two cells are soft enough that it does not impose any kinematical resistance. However, to exhibit the realistic structural behavior of the assembly under imposed kinematic constraint, it should be assumed that the connecting sheet in between is not soft (stiffer connection results in more resistance against the mismatch), and displays resistance to the mismatch between the spine angles during deformation ( $\psi^A$ , and  $\psi^B$ ). In terms of theoretically modeling this case, the magnitude of parameter  $k^*$  is crucial here. Figure 2.7 illustrates the potential energy landscape and reaction force of the dual-cell assembly along the continuous equilibrium path when the constraint stiffness  $k^*$  increases (the dotted lines in the first row of Figure 2.7 correspond to the potential energy maxima at a given length). As the constraint increases, the potential energy barrier for compression switch from "0-1" stable state to "0-0" increases significantly, but the energy barrier for the extension switch does not increase by the same degree. Moreover, when the kinematical stiffness reaches a threshold value  $\left(\frac{k^*}{k_I} = 140\right)$  in this case study), the initially continuous equilibrium path that connects three stable states splits into two separate ones (see the first two rows of Figure 2.7.b and .c). As a result, when the dual structure is extended from the "0-0" stable state, it will deform to point P at the end of one equilibrium path and then "leap" to the other path. In the compressing direction from "0-1" to "0-0", the dual structure deforms to Q first before leaping (see the insert figure in the first row in Figure2.7.c). The asymmetry in the energy barrier caused by kinematic constraint resulted from folding makes the required energy to reach mentioned leaps significantly different between the extension and compression direction.

By examining the changes in critical forces as the kinematical constraint  $k^*$  increases, the presence of the asymmetric energy barrier can be further emphasized (the third row of Figure 2.7). From Figure 2.7, one can see that with the increase of  $k^*$ , the required force to compress the dual structure from "0-1" stable state to "0-0" stable state is notably increasing while the required force to extend it back to "0-1" does not change much (Table 2.1).(Section 2.2 and Section 2.3 are published [1].)

Table 2.1: The normalized critical forces in the extension and compression switches between the (00) and (01) stable states based on the reaction force plots in Figure 6 and 7.

$\frac{k^*}{k_I}$	$\frac{F_e}{k_I}$	$\frac{F_c}{k_I}$
0	26.5	-91.7
50	32.5	-467.3
140	36.3	-1261.7
600	39.9	-2079.7

### 2.3 Diode effect in extension direction

In the unit cell of the extension diode, the orientation of sheet II is flipped. In other words, the orientation of sheet I and sheet II with respect to each other is the same as the traditional stacked Miura-ori. The dual cell chain of extension diode can be seen in Figure 2.8. The same as the compression diode unit cell, this unit cell also is consisted of geometrically compatible Miura-ori sheets and satisfies the rigid folding condition (equations 2.1 and 2.2). This change did not change the relations between the individual cells reactions forces ( $F_c^A < F_c^B$ , and  $F_e^A > F_e^B$ ). Moreover, the three achievable stable states by global extension or compression are the same as the compression diode.

In ideal rigid condition, where  $k^*$  is infinitely high, there is no mismatch between the spine angles. In this way, the admissible deformations of the extension dual cell chain are bounded to the "kinematic paths" shown in Figure 2.8.a. The total length of the dual stricture can be defined as follows (Figure 2.8.b):

$$L^{t} = L^{A} + L^{B} + L^{0}. (2.17)$$

One should note that the unit cell total length (cell A or B) is the summation of the sheet I and sheet II and the connecting sheet length while the unit cell length in the compression diode is calculated using equation 2.18:

$$L^{A} = L^{B} = L^{I} + L^{II} + L^{c}.$$
(2.18)

The total potential energy of the dual cell assembly is calculated with the same approach used for the compression diode. The changes in the energy landscape of this structure with the increase in  $k^*$  can be seen in the first row of Figure 2.10. The energy of the equilibrium path and the corresponding reaction force are shown in the second row and third row of Figure 2.10 respectively.

As  $k^*$  increases, the required force to switch from "0-1" to "1-1" stable configuration is increasing while, the reaction force for shifting from "1-1" back to "0-1" is not changed as much. Moreover the extension reaction force of changing from "0-0" to "1-1" does not experience a large change. Similarly with extension compression diode, when the reaches  $k^*$  the threshold  $\left(\frac{k^*}{k_I} = 220\right)$ , the continuous equilibrium path splits into two separate path (Figure 2.10.b). When the dual structure is extended from the "0-1" stable state to "1-1, it will deform to point Q at the end of one equilibrium path and then "leap" to the other path. In the compressing direction from "1-1" to "0-1", the dual structure deforms to P first before leaping (see the insert figure in the first row in Figure 2.10.c). With a closer look at the second row of Figure 2.10, it is evident that the extension energy barrier between state '0-1' and '1-1' is growing as  $k^*$  increases. However, this growing rate is not observed in the compression energy barrier between '1-1' and '0-1' stable configurations. Thus, to switch from '0-1' to '1-1' stable state, a larger force is required with the increase of  $k^*$ . The second row of Figure 2.10 also shows that the required force for compressing the structure from '1-1' to '0-1' is not increased much. In conclusion, this dual cell structure exhibits diode behavior in extension direction, a structure hard to extend but, easy to compress.



Figure 2.6: Mechanics of the dual-cell assembly assuming zero constraint stiffness  $k^*$ : (a) the total potential energy landscape, (b) the equilibrium path, and (c) the reaction force along the equilibrium path. The colormap in (a) represents the total potential energy, darker color means lower energy. It is worth nothing that in this figure and the following Figure 6, only the equilibrium path containing the (00), (01), and (11) stable states are shown in the energy landscape and reaction force plots. This is because the (10) state is not achievable by global extension or compression.


Figure 2.7: The energy contours (first row), energy landscapes (second row), and the reaction force (third row) corresponding to an increasingly stronger folding induced kinematic constraint: (a)  $\frac{k^*}{k_I} = 50$ , (b)  $\frac{k^*}{k_I} = 140$ , and (c)  $\frac{k^*}{k_I} = 600$ . The "leap" between the equilibrium paths are illustrated as dashed arrows in the insert figure in the first row of (c).



Figure 2.8: The dual cell chain of the extension mechanical diode. Miura-ori sheet II is flipped back to the configuration it poses in traditional stacked Miura-ori.



Figure 2.9: Kinematic properties of the extension diode structure due to the folding induced constraint (or the lack of). (a) Admissible deformation of the dual cell assembly. The two kinematic paths based on ideal rigid-folding condition are shown by the solid and dashed curves. The gray area represents deformations that are not kinematically admissible. (b) The geometry of the dual cell assembly at different locations along these to kinematic paths.



Figure 2.10: The energy contours (first row), energy landscapes (second row), and the reaction force (third row) corresponding to an increasingly stronger folding induced kinematic constraint: (a)  $\frac{k^*}{k_I} = 50$ , (b)  $\frac{k^*}{k_I} = 220$ , and (c)  $\frac{k^*}{k_I} = 600$ . The "leap" between the equilibrium paths are illustrated as dashed arrows in the insert figure in the first row of (c).

### Chapter 3

# Experimental Investigation of the Compression and Extension diode

# 3.1 Experimental observation of the diode behavior in compression

After numerous modifications, a carefully designed prototype was fabricated to experimentally validate the analytical results. The patterns of the smallest components of the geometry including the parallelograms of sheet I and II, and the connecting sheets were designed in SolidWorks<sup>TM</sup>. (The drawings can be found in the appendix) In the rigid folding condition, the planes are ideal rigid. In order to make the experimental setup as close to the rigid folding assumption as possible, fatigue-resistant 301 stainless steel spring temper sheet of 0.01" thickness was used to cut the parts from. This steel provided enough rigidity to the experimental setup to satisfy the

rigid folding condition up to a reasonable extent. The parts were water-jet cut, and two similar unit cells were fabricated using adhesive UHMW Polyethylene film of .005" thick. One of the main issues in developing the experimental setup was to find the proper method of fabricating the dual assembly to achieve an ideal multi-stable force-displacement curve. For example, initially the UHMW film was attached on both sides of the cut parts to assemble the dual chain setup. In order to archive a multi-stable measured F-D curve, one should minimize the amount usage of the UHMW adhesive film, specially on the crease regions. In another unsuccessful attempt, the dual assembly parts were broken down to smaller subsets to be fabricated separately and attached together to build the final structure. This method required using extra adhesive film resulting in failure in obtaining proper curves. Eventually, a fabrication process was designed that gave better multi-stable force-displacement curves (Appendix A). The weakness of this process shows itself in experimental investigation of the diode behavior  $(k^* > 0)$ . This method can not deliver strong enough connection between the two unit cells when they are connected to each other along their zig-zag creases to represent the added kinematical stiffness k\*.

For mounting the unit cells to the Universal Tensile Tester, an additional part was designed, and water-jet cut on the from the same materials and was attached to the assembly on carefully determined places on the unit cells. Beside this steel part, a customized connector was 3D printed. After many modifications, the best design for the 3D printed connectors was used to enable the structure to be mounted on the machine.

As it was mentioned above, for a stacked Miura-ori unit cell to be bi-stable, the torsional stiffness of certain creases of the unit cell should be notably higher than the other creases. The two adjacent creases on one side of the connecting sheets were chosen to add the stiffness. Many different metals with different thicknesses were used to choose the one that delivers the best experimental results. For example, the thickness of 0.007 18-8 stainless steel was not strong enough and deformed instantly after one loading cycle, and could not maintain the elastic behavior; Or the thickness of 0.009 of the same steel was too strong for the force scale of the designated experiment and the delicacy of the fabricated setup. Finally, the 18-8 stainless steel shim Stock (0.008" thickness) was employed to add the torsional stiffness on the creases of connecting sheets. This thickness provided enough strength, and at the same time maintains its elastic spring behavior during deformation. The rectangle parts from shim stock in  $2.5cm \times 3cm$ , were bent into equal angles to be attached on the creases. to avoid any loss of accuracy in the experiment, the angles of all the springs must be as equal as possible. In order to obtain the best consistency, a pair of fixtures with the mating surfaces angle of  $65^{\circ}$  was cut.

Two sets of experiments were conducted. In the first experimental setup, the two unit cells were simply connected to each other in series using M6 rigid rod with a balanced internal force (Figure 3.1.a). For the second test set, the two cells were connected to each other along their zig-zag crease lines by adhesive films. This setup is consistent with the stacked origami construction shown in Figure 3.1.b. In both sets, the arrangement of the cell is in a way that cell B is always on top.

Several single tension and compression load cycles, using the displacement control method, were conducted with the two setups. The increase in the compression force in crease connection in comparison with rod connection was noticed in all of them consistently. In what follows, a pair of numbers (i-j) is used to represent the stable configuration of Cell A and B respectively. All the tests were done in a way that the dual-cell chain is first compressed from the "convex- convex" (i.e. 1-1) stable config-



Figure 3.1: The photos of the rod-connected test (a) and crease-connected test of compression diode dual cell-chain (b) show the three stable states ('1-1','1-0', and '0-0')

uration to the "concave-concave" (i.e. 0-0) state and extended back to 1-1 state. The force-displacement graph of one test is shown in Figure 3.2. The noticeable hysteresis in the experimental graphs is due to using adhesive films. After the completion of half of the loading cycle (compression direction), this film goes through plastic deformation and does not provide the desired elastic behavior.

The crosshead speed of this test was 0.08 mm/s. Four different stable configuration combinations are possible for the dual-cell chain: '0-0', '0-1', '1-1', and '1-0'. These switches are evidenced by the negative slopes in both theoretical and experimental force-displacement curves (i.e. negative stiffness). The theoretical model depicts that the relation between the individual cells critical forces is as follows:  $F_e^A > F_e^B$ and  $F_c^A < F_c^B$ . This relation needed to beheld in the experimental setup as well. Experimentally, this relation could be applied by attaching proper numbers of steel stripe on both cells with a certain proportion. It was chosen to attach 4 stripes on



Figure 3.2: Compression and tension tests on compression diode dual-cell chain prototypes with rod connection (red curve) and crease connection (blue curve).

Cell B and 2 stripes on Cell A to achieve the desired force relationships.

The snapping sequence of the dual-cell chain is dictated by this relation both in compression and tension. Thus, only the first three of the possible combinations are achievable via displacement control. In other words, since the critical compression force of Cell B is higher than that of Cell A. '1-0' combination is not attainable. In conclusion, during compression, always Cell A nests in first because of its critical force for snap-through  $(F_c^A)$  is lower than that of the top Cell B  $(F_c^B)$ , and during tension, Cell A bulges out first and then Cell B. Regardless of the inter-cellular connection, the switching sequence in both compression and tension is dictated by the individual cells force relation. In compression, the switching sequence is '1-1','0-1','0-0', and in the tension is '0-0', '0-1', '1-1'. The experimental force-displacement curves of Cell A and B are shown in Figure 3.3.



Figure 3.3: Measured force-displacement curves of the two unit cell prototype of the compression diode.

With a comparison of the measured force-displacement curves of the two sets in Figure 3.2, it can be seen that the compression force from switching (01) to (00) is increased in the second setup (crease connection) while no notable increase was seen in extending from '0-0' to '0-1'. This experimentally validates the diode behavior

that was noticed in analytical results. However, not much increase was noticed from switching from '1-1' to '0-1'. (The critical forces of the dual chain are denoted with 'F' and 'R' superscript that are referring to crease (film) connection and rod connection, respectively). More specifically, in film connection, the required force to switch from "0-1" to "0-0" was 1.9 N and in rod connection was 1 N.

# 3.2 Experimental observation of the diode behavior in extension

In the previous sections, the extension diode behavior obtained from the analytical model. It was necessary to experimentally observe diode behavior in extension. The same assembly was used for this experiment with a slight variation in the Sheet II configuration. To obtain the proper setup, the Sheet II of each cell was flipped. That is the only change applied to the compression diode assembly (Figure 3.4).

According to the theoretical model, the same relation between the individual cells critical forces is held and that is,  $F_e^A > F_e^B$  and  $F_c^A < F_c^B$ . The Force-Displacement curve of the individual cells is shown in Figure 3.5. It is evident from the graph that the experimental setup is properly set to satisfy the desired force relation.

Similarly, with compression dual structure, the extension dual-chain poses four possible stable states regardless of the inter-cellular connection: '0-0', '0-1', '1-1', and '1-0'. Again, due to the assigned force relation, only the three stable arrangements are achievable via the displacement-control method (Figure 3.6).

The extension loading starts with the assembly at the 'convex-convex' configuration.



Figure 3.4: The extension diode individual cell prototype

That relation incurs a switching sequence of '0-0' to '0-1' to '1-1' during extension, and by the end of compression loading, the structure goes back to 'convex-convex' configuration. More specifically, Cell B bulges out first and then Cell A during extension, and Cell A nests in first and Cell B followingly (Figure 3.7).

Figure 3.6 shows that in crease-connected setup (blue curve), the required force to extend the dual cell assembly from '0-1' to '1-1' is higher than the needed force to apply this switch in the rod-connected. More specifically, in film connection, the required force of going from '0-1' to '0-0' is 2.1 N, and in the rod connection 1.5 N However, not much difference in magnitude of tension critical force of switch between '0-0' and '0-1' is seen (Figure 3.6).



Figure 3.5: Measured force-displacement curves of the two unit cells prototype of the extension diode.

### 3.3 Conclusion

The only difference between the two sets (whether in "compression diode" or "extension diode") is the inter-cellular connection, one connected with the rod, and the other with the adhesive film along the zig-zag creases. This illustrates that the connection between two bi-stable cells can impose a kinematic constraint on the static behavior of the structure. In other words, the increase of critical force for shifting from one stable state to the next in the crease connection reveals the significant increase in the energy barrier in extending from '0-1' to '1-1' in extension diode setup and in compressing from '0-1' to '0-0' in compression diode.



Figure 3.6: Tension and compression tests on extension diode dual-cell chain prototypes with rod connection (red curve) and crease connection (blue curve).

Observing the diode effect in these two dual cell chains (extension diode and compression diode) elucidates that this unique asymmetric energy barrier is the result of the coupling between unit cell length change in the z-axis and the connecting creases displacement along x and y axes at the boundary between two cells. This accentuates the importance of the three-dimensional nature of origami folding in obtaining mechanical diode behavior and multi-stability.



Figure 3.7: The photos of the rod-connected test (a) and crease-connected test of compression diode dual cell-chain (b) show the three stable states ('1-1','1-0', and '0-0')

### Chapter 4

# Optimization

### 4.1 Introduction

The primary objective of this study was to introduce a mechanical diode effect obtained by connecting two origami unit cells along their zigzag crease using adhesive film. Therefore, strengthening the diode effect is the prior concern in the design. On the other hand, minimizing the required material to fabricate a cellular structure is crucial to minimize the production cost. These two criteria were considered to find an optimum feasible solution region. The independent geometrical parameters were  $a_I, b_I, l_c$ , and  $\gamma_I$ ). As it was mentioned in previous sections, the ideal diode behavior is originated from the nonlinearity in the describing correlations of the geometry. From the four involved parameters,  $\gamma_I$  is responsible for the biggest portion of the nonlinearity in the correlations describing the unit cell geometry (equations 2-7,..., 2-12). The strength of the diode behavior can be measured by the ratio of the critical compression force to the extension force of the dual cell chain. The first objective is to maximize this ratio. The higher this ratio is, the stronger the mechanical diode can be achieved. In practice, the importance of this matter is revealed in the design of a mechanical logic gate to bear a high loading threshold.

Referring to previous sections, the derivative of the total stored elastic energy with respect to total length change gives the reaction force of the dual origami structure. Therefore, the design variables in the optimization process are the geometrical parameters used in potential energy evaluation including  $\gamma_I$ , a, b, and  $l_c$ .

To decrease the fabrication cost, the surface area of the unit cell should be minimized (note that the two cells are identical). The main constituents of the unit cell are: sheet I, sheet II and the connecting sheet.

Here, the inter-cellular connecting sheets are and the connectors that attach the two cells are assumed to have equal dimensions. The total surface area of the unit cell can be calculated as the summation of the areas of 4 parallelograms with sheet I dimensions and 4 parallelograms with sheet II dimensions and 2 rectangles with dimensions of the connecting sheet.

The total surface area is evaluated by the summation of the three following equations:

$$A_1 = 2a_I b_I \sin \gamma_I. \tag{4.1}$$

$$A_2 = 2a_{II}b_{II}\sin\gamma_{II}.\tag{4.2}$$

$$A_3 = 2b_I l_c. \tag{4.3}$$

### 4.2 Design problem

#### 4.2.1 Design objective

The goal of this bi-objective optimization problem is to strengthen the diode behavior and at the same time to minimize the used material required for fabrication. In other words, the first objective is to maximize the defined force ratio and the second objective is minimizing the unit cell surface area. The next important step is to define all the geometrical constraints.

#### 4.2.2 Design constraints

As it can be seen in figure 2-52.5, the unit cell poses the smallest length at  $\theta_I = -90$ . Therefore, it is important to take into consideration that the length of the intercellular connecting sheet is long enough to prevent any contact between sheet I and sheet II at this stable configuration. This geometrical constraint can be stated by the following inequality:

$$l_c > |l_I^{\max}| + |l_{II}^{\max}|.$$
(4.4)

Where  $l_I$  is the height of sheet one at  $\theta_I = -90$  and  $l_{II}$  is the height of sheet II at

the  $\theta_{II}$  (2.3). LI and LII are calculated using following equations:

$$L_I = a_I \sin \theta_I \sin \gamma_I. \tag{4.5}$$

$$L_{II} = a_{II} \sin \theta_{II} \sin \gamma_{II}.. \tag{4.6}$$

#### 4.2.3 Optimization problem setup

modeFrontier® 2017 R1 was used as the optimizer platform and MATLAB® 2017b was linked to the optimizer as the solver. DOE properties were adjusted at Basic mode with "Uniform Latin Hypercube" as the space filler. For this bi-criteria optimization problem "MOGA-II" method was used to carry out the optimization. "MOGA-II" which stands for "Multi-Objective Genetic Algorithm II" is an efficient evolutionary optimization algorithm for a constrained problem. Solving a multi-objective problem with the traditional form of the "Genetic Algorithm" can face deficiency in converging to the true "Pareto Front", and misidentify the true optima. "MOGA-II" algorithm tackles this issue with smart multi-search elitism. This new elitism operator has the advantage to preserve some desirable solutions without bringing the premature convergence into the local optimal fronts. In this method, the constraints are tackled by applying "Penalty Method", and it can handle both continuous and discrete design space (In the process of the optimization using this algorithm, the continuous design space is discretized internally [37]. In modeFrotier(R), the optimization algorithm configuration is set to the "automatic mode". In this configuration, The "Number of Generations" is 100, the "Probability of the Directional Cross-Over" is 0.5, the "Probability of Mutation" is 0.1, the "DNA String Mutation Ratio" is 0.05, and the "Number of Evaluations" is chosen to be 2000. The high values of the "Probability of Directional Cross-Over" decreases the robustness of the algorithm and this may cause the optimization process to get trapped at a local optima without touring the whole design space. This consideration matters in highly nonlinear problems, such as the "compression diode unit cell's geometry". Another important point is to enable the "Elitism" operator to enhance the convergence of the algorithm [37]. The workflow diagram can be seen in Figure 4.1.



Figure 4.1: The developed workflow in modFrontier®to obtain optimized designs admissible region.

For  $a_I$  and  $b_I$  the lower bound of 0.5 cm and upper bound of 5 cm was chosen.  $\gamma_I$  varies between 45° and 75°, and  $l_c$  is changing from 0.8 cm to 5 cm. One should note that the different combinations of these design parameters can lead to mono-stability. Thus, it is crucial to allow only eligible designs into the optimization process. This constraint is defined in the MATLAB script linked to the optimizer.

### 4.3 Optimization results

Among the four design variables,  $\gamma_I$  has a nonlinear relationship with angles defining the unit cell geometry; Therefore, the strength of the diode effect is highly sensitive to the variation of  $\gamma_I$ . More specifically, the spine angles, which quantify the kinematic constraint, are function of  $\gamma_I$ . This further demonstrates the impact of  $\gamma_I$  on the strength of the diode behavior of the structure. Sheet I and sheet II crease stiffnesses are function of a and b and the connecting sheet crease stiffness is a function of  $l_c$ . These stiffnesses have a linear relationship with the unit cell potential energy. These three design variables play a very important role in the surface area of the unit cell.

The obtained Pareto front from the optimizer is shown in Figure 4.2.

In the front, the force ratio varies from 18.5 to 36.9 and the surface area changes from 5.5  $cm^2$  to 17.3  $cm^2$ . The design variables at the beginning of this range are  $a_I = 0.5cm$ ,  $b_I = 0.5cm$ ,  $l_c = 1.1cm$ , and  $\gamma_I = 1.3rad$ , and the at the end of the range are  $a_I = 0.5cm$ ,  $b_I = 1.8cm$ ,  $l_c = 1.2cm$ , and  $\gamma_I = 1.3rad$  (Figure 4.3). It is worth to note that  $\gamma_I$  value on the "Pareto Front" is consistently 1.3rad. This shows that the optimum diode effect strength is achieved at this  $\gamma_I$  value and further illustrates the importance of this design parameter. However, parameter  $a_I$  value varies within the optimized designs range. Based on what type of application this structure is deployed in, and factors such as size and cost one objective can be the priority to the other one.



Figure 4.2: The"Pareto Front" for compression diode dual-cell chain.



Figure 4.3: a) Miura-ori sheet I's geometry at the beginning of the optimum range. b) Miura-ori sheet I's geometry at the end of the optimum range.

### Chapter 5

# **Conclusion and Future Work**

### 5.1 Conclusion and Summary

This study proposes and examines a mechanics model to theoretically examine the static behavior of a multi-stable cellular origami structure that exhibits diode behavior in compression, presents experimental results to validate the theoretical results and, investigates, theoretically and experimentally, a transformed version of the compression diode unit cell that results in mechanical extension diode. Each unit cell in this cellular structure is essentially a bistable unit that consists of two geometrically consistent Miura-ori sheets and zig-zag shaped connecting sheet. The bi-stability in each cell originates from the nonlinear correlations between folding and crease deformation. It was shown that with a small change in the designed unit cell of the compression cellular structure, a new unit cell is obtained that shows diode behavior in extension. Each unit cell in this cellular structure is also a bi-stable unit that follows the same nonlinear correlation between the crease rotation and the overall

#### folding.

The fundamental construction of this study was a dual-cell chain to investigate the static behavior of that. Using this dual-cell structure, the desired diode behavior was achieved in both compression and extension.

These origami unit cells (compression and extension diode designs) are essentially a three-dimensional transformation mechanism. This 3D nature of origami imposes a unique kinematic constraint onto the deformations of the two connected unit cells. More specifically, the magnitude mismatch between the spine angles of the two cells during deformation can determine the strength of the imposed kinematical constraint. This constraint can be quantified with an equivalent stiffness,  $k^*$ . The higher  $k^*$  results in a higher energy barrier either in compression direction or extension. In compression diode structure, due to this constraint, a higher force is required to deform the dual cell chain in compression direction while the required force to extend the structure is significantly smaller. In extension diode, the structure is easy to compress but hard to extend due to this added stiffness.

Then, two experimental setups were developed for both extension diode and compression diode to experimentally support the theoretical results. In the first set, the two unit cells are connected via a rigid rod with a balanced internal force. In the second setup, the two cells are connected via a zig-zag shaped connecting sheets along their zig-zag creases using adhesive film. The experimental observations revealed that the crease connection increases the energy barrier for switching between certain stable states. This increase for compression diode is in compressing the dual structure from one stable state to the next one and in the extension, it happens when the chain is being stretched from one stable state to the next. Obtaining diode-like behavior by employing nonlinear elastic properties to achieve diode-like mechanical behavior is a progressively important subject of research [9]. In recent studies, it has been shown that multi-stable structures and materials with carefully designed microstructures are capable of attaining unidirectional acoustic [2] and elastic wave propagation [34]. The diode effect reported in this study can open new avenues toward multifunctional structures and material systems that can be deployed in motion rectifying, wave propagation control, and mechanical computation.

### 5.2 Future Work

In chapter 3, a notable hysteresis was observed in the experimental force-displacement curves. This hysteresis was due the plastic deformation of "Adhesive UHMW Polyethylene Film" during the full loading cycle. If one runs the experiment for more than one cycle, the second or third cycle does not result in desirable  $F_D$  curves and the multi-stability behavior can not be concluded from the second or third loading cycles graphs. Thus, it is important to find an alternative to assemble the structure that can address the hysteresis issue and enhance the experiment repeatability. Moreover, the magnitude of increased forced between rod connection and crease connection (the increase in the tension reaction force in extension diode and increase in the compression reaction force in compression diode ) was expected to be higher. One improvement in the experimental setup should be strengthening the crease connection between two unit cells.

Developing a finite element model of the unit cell, and the stacked cellular structure is a very efficient tool to further investigate the mechanics of the developed theoretical model. This FEA model can be used to do a stress analysis on the static diodes and study the kinematics of the structure. Moreover, based on the results of these studies one can apply required modifications to further enhance the unit cells design and performance according to the anticipsted application. Additionally, developing a logic gate based on these two unit cells is the next phase of this project to demonstrate the capabilities of these designs in practice.

# Appendix A

# **Prototype Fabrication Process**





Figure A.1: SolidWork drawings of the fabricated prototype



Figure A.2: The used outlines of the subsets that are connected to each to assemble the dual structure using the method introduced in Appendix A



Figure A.3: Step 1. The cut parts should be fixed on the proper place on the outline as shown in Figure A.3



Figure A.4: Step 2.



Figure A.5: Step 3.



Figure A.6: Step 4: Fixing a piece of adhesive film with the sticky side facing upward.



Figure A.7: Step 5: The "slippery sheet" should be kept to be used in next step.



Figure A.8: Step 6: Attache the "slippery sheet" as it shown in the figure form the slippery side to the adhesive film.


Figure A.9: Step 7. Drawing the flaps of the connecting sheets using the sheet I and sheet II parallelograms.



Figure A.10: Step 8



Figure A.11: Step 9



Figure A.12: Step 10. Drawing the flaps of the sheet I using the side connecting sheet part.



Figure A.13: Step 11



Figure A.14: Step 12



Figure A.15: Step 13. Attaching the subsets to each other along their flaps.



Figure A.16: Step 14.

#### Appendix B

### MATLAB®Codes

## B.1 Theoretical Analysis Main MATLAB Script for $k^* = 50$ Compression Diode

1 clear
2 clc
3 clf
4 close all
5 %%CELL A parameters
6 aIa=2;
7 % bIa=aIa;
8 bIa=4;
9 aIIa=1.25\*aIa;
10 landaIa=pi/4;
11 landaIIa=acos((aIa\*cos(landaIa))/aIIa); % rigid-folding condition
12 KIa=1;
13 K\_A=zeros(1,5);

```
14 K_A(1, 1)=KIa;
   L_{-}c = 5;
15
   L_{-}o = 5;
16
   theta0_a=-pi/3;
17
   theta_a1_2_0 = acos((cos(theta0_a).*tan(landaIa))./(tan(landaIIa)));
18
19
  K_A(1,2) = 20 * KIa;
20
  K_A(1,3) = KIa;
21
  K_A(1, 4) = KIa;
22
  K_A(1,5) = 0;
23
   theta_a1=linspace(-pi/2, pi/2, 1555);
^{24}
   theta_a1_2 = acos((cos(theta_a1).*tan(landaIa))./(tan(landaIIa)));
25
   theta_a1_2_0 = acos((cos(theta0_a).*tan(landaIa))./(tan(landaIIa)));
26
   s1a=sin(landaIa);
27
   t2a=tan(landaIIa);
28
   t1a=tan(landaIa);
29
   L1_a = aIa . * sin(theta_a1) . * sin(landaIa);
30
   L2_a = aIIa . * sin(theta_a1_2) . * sin(landaIIa);
^{31}
   n = length(L1_a);
32
   L2_a_max = max(L2_a);
33
   L1_a_max = max(L1_a);
34
   L2_a_min=min(L2_a);
35
   L1_a_min=min(L1_a);
36
   LA_min=L2_a_min+L1_a_min+L_c;
37
   LA_max=L2_a_max+L1_a_max+L_c;
38
  % LA=linspace(LA_min,LA_max,n);
39
   LA_I=linspace(L1_a_min, L1_a_max, length(L1_a)); %length of sheet I in new
40
        design
```

```
41
```

```
42 %%
```

```
43 %CELL B parameters
```

```
aIb=2;
44
   bIb=aIb;
45
   aIIb = 1.25 * aIb;
46
   landaIb=pi/4;
47
   landaIIb=acos((aIb*cos(landaIb))/aIIb); % rigid-folding condition
48
   KIb = 1;
49
  K_B = zeros(1,5);
50
   K_B(1, 1) = KIb;
51
52
   theta0_b=pi/3;
53
   theta_b2_0 = acos((cos(theta0_b).*tan(landaIb))./(tan(landaIb)));
54
55
  K_B(1,2) = 20 * KIb;
56
  K_B(1,3) = KIb;
57
  K_B(1, 4) = KIb;
58
  K_B(1,5) = 0;
59
   theta_b1 = linspace(-pi/2, pi/2, n);
60
   theta_b2=acos((cos(theta_b1).*tan(landaIb))./(tan(landaIIb)));
61
   s1b=sin(landaIb);
62
   t2b=tan(landaIIb);
63
   t1b=tan(landaIb);
64
   L1_b=aIb.*sin(theta_b1).*sin(landaIb);
65
   L2_b=aIIb.*sin(theta_b2).*sin(landaIIb);
66
   L_b=L_c+L1_b-L2_b;
67
   L2_b_max = max(L2_b);
68
   L1_b_max = max(L1_b);
69
   L2_b_{\min} = \min(L2_b);
70
   L1_b_{\min}(L1_b);
71
  %%
72
  %Total length CEll A_I_II in original design
73
```

```
74 A=aIa * sin(land aIa);
```

```
B = \left( \left( \tan \left( \operatorname{landaIIa} \right) \right)^2 \right) / \left( \left( \tan \left( \operatorname{landaIa} \right) \right)^2 \right);
75
76
77
    L_a = aIa * s1a * ((sqrt(((t2a^2)/(t1a^2)) - ((cos(theta_a1).^2)))) - sin(theta_a1))))
 78
         ));
    L_A_{max} = max(L_a)
 79
    L_A_{\min} = \min(L_a)
 80
    %%
 81
    %Mismatch parameters
82
    K_{star} = 50;
 83
84
    say_A0=2*atan(cos(theta0_a)*tan(landaIa));
 85
    say_B0=2*atan(cos(theta0_b)*tan(landaIb));
86
87
    n = length (L1_a)
88
    %%
 89
90
 91
    %%
92
    %Calculating maximum and minium
 93
    i5 = 0;
94
    for i3=linspace(-pi/2, pi/2, n)
 95
96
          i5=i5+1;
97
          i6 = 0;
98
          for i4 = linspace(-pi/2, pi/2, n)
99
               i6 = i6 + 1;
100
               theta_a1_2(i5) = acos((cos(i3).*tan(landaIa))./(tan(landaIIa)));
101
               theta_b_2(i6) = acos((cos(i4).*tan(landaIb))./(tan(landaIIb)));
102
               total_length(i5, i6) = (L_c + (aIa. * sin(i3). * sin(landaIa)) - abs((aIIa)))
103
                    .* sin (theta_a1_2(i5)).* sin (landaIIa)))))+(L_c+(aIb.* sin (i4).*
```

```
\sin(\text{landaIb}))-abs((aIIb.*sin(theta_b_2(i6)).*sin(\text{landaIIb})))
                  )+L_{-0};
         end
104
105
    end
106
    L_total_min=min(min(total_length));
107
    L_total_max = max(max(total_length));
108
    L_T=linspace(L_total_min, L_total_max, n); %New design total length
109
   %%
110
111
    i1 = 0;
112
    for LT=linspace(L_total_min, L_total_max, n)
113
         i1 = i1 + 1;
114
      for i2=1:length(L1_a)
115
              [theta_A1(i1,i2), theta_A2(i1,i2), E_A(i1,i2), LA_II(i1,i2),
116
                  L_A_I_II(i1,i2)]=CellA(LA_I(i2),landaIa,aIa,bIa,K_A,theta0_a
                  , n, L_{-}o, L_{-}c);
              [E_B(i1,i2), LB_I_II(i1,i2), theta_B(i1,i2), E1_b2(i1,i2), E2_b2(i1,
117
                  (i2), E3_b2(i1, i2), E4_b2(i1, i2), E7_b2(i1, i2)] = CellB(aIb, aIIb, aIIb)
                  landaIb, LA_I(i2), LA_II(i1,i2), K_B, theta0_b, L2_a_min, L2_a_max
                  , L_A_I_{I_i}(i1, i2), n, LT, L_o, L_c);
              if E_B(i1, i2) = 0
118
                 say_A(i1, i2) = 2*atan((cos(theta_A1(i1, i2)))*tan(landaIa));
119
                 say_B(i1, i2) = 2*atan((cos(theta_B(i1, i2)))*tan(landaIb));
120
                 E_say(i1, i2) = K_star * bIa * (((say_A(i1, i2) - say_B(i1, i2)).^2)./2)
121
                      ;
                 E6_{a2}(i1, i2) = (K_A(3) \cdot (say_A(i1, i2) - say_A0) \cdot 2) \cdot /2;
122
                 E6_b2(i1, i2) = (K_B(3) \cdot (say_B(i1, i2) - say_B0) \cdot 2) \cdot /2;
123
                 E8_A_B(i1, i2) = (K_A(5) \cdot (say_B(i1, i2) - say_A(i1, i2)) \cdot (2) \cdot (2)
124
125
                 Et_A_B(i1, i2) = E_A(i1, i2) + E_B(i1, i2);
126
```

```
Et_say(i1, i2) = Et_A_B(i1, i2) + E_say(i1, i2) + E8_A_B(i1, i2) + E6_a2(i1) + E6_a2(i1)
127
                                                                                       i1, i2)+E6_b2(i1, i2);
                                                         else
128
                                                                       say_A(i1, i2) = 0;
129
                                                                       say_B(i1, i2) = 0;
130
                                                                       E_{-say}(i1, i2) = 0;
131
                                                                       E6_{-a2}(i1,i2)=0;
132
                                                                       E6_b2(i1,i2)=0;
133
                                                                       E8_A_B(i1, i2) = 0;
134
                                                                       Et_A_B(i1, i2) = 0;
135
                                                                       Et_{-say}(i1, i2) = 0;
136
137
                                                        end
138
                                    end
139
                end
140
141
              %%
142
              %
143
               % %Plotting results
144
                 for i=1:n
145
                                    L1_a_2(i) = aIa.*sin(theta_A1(1,i)).*sin(landaIa);
146
                                    L2_a_2(i) = aIIa . * sin(theta_A2(1,i)) . * sin(landaIIa);
147
                                   LA(i) = L1_a_2(i) - L2_a_2(i) + L_c;
148
                end
149
150
                [X1,Y1] = meshgrid(L_T,LA);
151
               XX1=X1.';
152
               YY1=Y1.';
153
              % %%
154
              % %Total energy
155
              % figure(1)
156
```

```
% V1=linspace(0,1000,100); % energy level
157
   %
158
   % contourf(XX1,YY1,Et_say)
159
   % xlabel('Total Length')
160
   % ylabel('Cell A_Sheet I length')
161
   % title ('Total Energy Contour')
162
   %
163
   % %%
164
   % %Thet_A1
165
   \% figure (2)
166
   % plot3(XX1,YY1,theta_A1)
167
   % title ('Theta_A_I')
168
   % %%
169
   % %Theta_A2
170
   \% figure (3)
171
   % plot3 (XX1, YY1, theta_A2)
172
   % title ('Theta_A_II')
173
   % %%
174
   % %Theta_B
175
   \% figure (4)
176
   % plot3(XX1,YY1,theta_B)
177
   % title ('Theta_B')
178
   % %%
179
   %
180
   %%plotting energy landscape of cell A
181
   % theta_A1_2=linspace(-pi/2, pi/2, n);
182
   %% L_A_I_III_2=abs(aIa*s1a*((sqrt(((t2a^2)/(t1a^2))-((cos(theta_A1_2)
183
       (^{2})))) - \sin( \text{theta}_A 1_2)));
   % L1_A_I_2=aIa.*sin(theta_A1_2).*sin(landaIa);
184
185 % phi10_a=pi-2.*theta0_a;
```

186 % phi20\_a=2.\*asin  $(\cos(\text{theta0_a})./\operatorname{sqrt}(1-((\sin(\text{theta0_a})).^2).*(\sin($ 

 $landaIa)).^2));$  % it should not be more than 1 or -1

- 187 % phi30\_a=pi (2.\*acos(tan(landaIa).\*(1/(tan(landaIIa))).\*cos(theta0\_a)))
  ;
- 188 % phi40\_a=2.\*asin(((sin(landaIa))./(sin(landaIIa))).\*sin(phi20\_a./2));
- 189 %  $phi50_a = (pi/2) + theta0_a$ ; % it should not be more than 1 or -1
- 190 % s1a=sin(landaIa);
- 191 % t2a=tan(landaIIa);
- 192 % t1a=tan(landaIa);
- 193 %  $phi1_a1=pi-2.*theta_A1_2;$
- 194 % phi2\_a1=2.\*asin(cos(theta\_A1\_2)./sqrt(1-((sin(theta\_A1\_2)).^2).\*(sin( landaIa)).^2));
- 195 % phi3\_a1=pi (2.\*acos(tan(landaIa).\*(1/(tan(landaIIa)))).\*cos(theta\_A1\_2)
  ));
- 196 % phi4\_a1=2.\*asin(((sin(landaIa))./(sin(landaIIa))).\*sin(phi2\_a1./2));

197 % 
$$phi5_a1 = (pi/2) + theta_A1_2;$$

- 198 % say\_ $A1_2=2*$ atan (cos (theta\_ $A1_2$ )\*tan (landaIa));
- 199 % say\_ $A1_2_0=2*atan(cos(theta0_a)*tan(landaIa));$

%Kc3\_a

- 200 % K1\_a= $2*K_A(1)*bIa$ ; %KIa
- 201 % K2\_a= $2*K_A(1)*aIa$ ; %KIa
- 202 % K3\_a= $2*K_A(2)*bIa$ ; %KIIa
- 203 % K4\_a= $2*K_A(2)*aIIa$ ; %KIIa
- 204 % K6\_a=2\*L\_c\*K\_A(3); %inter celllar crease stiffness . there are tow creases with their associated stiffness in each unit cell
- 205 %
- 206 % K7\_a=8.\*K\_A(4).\*bIa; %Kc2\_a
- 207 % K8\_a=K\_A(5).\*bIa; %K\_external
- 208 % E1\_a1=K1\_a.\*(phi1\_a1-phi10\_a).^2;
- 209 % E2\_a1=K2\_a.\*( $phi2_a1-phi20_a$ ).^2;
- 210 % E3\_a1=K3\_a.\*(phi3\_a1-phi30\_a).^2;
- 211 % E4\_a1=K4\_a.\*( $phi4_a1-phi40_a$ ).<sup>2</sup>;
- 212 % E6\_a1=K6\_a.\*(say\_A1\_2-say\_A1\_2\_0).^2;

- 213 % E7\_a1=K7\_a.\*( $phi5_a1-phi50_a$ ).<sup>2</sup>;
- $214 \ \% Et_a = (E1_a1 + E2_a1 + E3_a1 + E4_a1 + E7_a1 + E6_a1 + E7_a1) . / 2;$
- 215 % figure (5)
- 216 % plot ( $L1_A_I_2$ ,  $Et_a$ )
- 217 % %%
- 218 % % Eenrgy lanscape
- 219 % figure (6)
- 220 % plot3(XX1,YY1,Et\_say)
- 221 % title ('Total energy Landscape')
- 222 % %%
- $_{223}$  % % Plotting CELL B energy landscape
- 224 % figure (7)
- 225 % plot3(XX1,YY1,E\_B)
- 226 % title('Cll B energy Landscape')
- $_{227}$  % %%
- $_{228}$  % %Plotting CELL B energy landscape
- 229 % figure (8)
- 230 % plot3(XX1,YY1,E\_A)
- 231 % title('Cll A energy Landscape')
- 232 % %%
- 233 % % Kinematic Constraint Energy Contour
- 234 % V2=linspace(0,100,100); % energy level
- 235 % figure (9)
- 236 % contour (XX1,YY1, real( $E_say$ ),V2)
- 237 % xlabel('Total Length')
- 238 % ylabel('Cell A\_Sheet I length')
- 239 % title ('Kinematic Constraint Energy Contour')
- 240 % %%
- $_{241}$  % %Plotting total energy contour with total length of cell A on y axis  $_{242}$  %
- 243 % LA1\_min=L2\_a\_min+L1\_a\_min+L\_c;

```
244 % LA1_max=L2_a_max+L1_a_max+L_c;
```

```
245 % LA1_range=linspace (LA1_min, LA1_max, n);
```

246 %  $[X2, Y2] = meshgrid(L_T, LA1_range);$ 

```
247 % XX2=X2.';
```

```
248 % YY2=Y2.';
```

```
249 % figure (10)
```

```
250 % contour(XX2,YY2, real(E_say),V2)
```

```
251 % xlabel('Total Length')
```

```
252 % ylabel ('Cell A length')
```

```
253 % title ('Total Energy Contour')
```

```
254 %
```

```
_{255} % %%
```

```
256 % %Plotting say
```

```
257 % figure (18)
```

```
258 % plot3(XX1,YY1,say_A)
```

```
259 % title ('Say_A')
```

```
260 % %%
```

```
261 % figure (19)
```

```
262 % plot3(XX1,YY1,say_B)
```

```
263 % title('Say_B')
```

```
264 %
```

```
265 % %%
```

```
266 % % Plotting E_say
```

```
267 % figure (20)
```

```
268 % plot3(XX1,YY1,E_say)
```

```
269 % title ('ESay')
```

270 % %%

```
271 \% %Plotting E8_A_B
```

272 % figure (21)

```
273 % plot3 (XX1,YY1,E8_A_B)
```

```
274 % title ('E8_A_B')
```

```
275 % %%
   % %Plotting E6_a2
276
   % figure(22)
277
   % plot3(XX1,YY1,E6_a2)
278
   % title('E6_a2')
279
   %
280
   % %%
281
   %%Plotting E6_b2
282
   % figure (23)
283
   % plot3(XX1,YY1,E6_b2)
284
   % title('E6_b2')
285
286
287
    figure(24)
288
289
290
    Maxima = [];
291
    Index = [];
292
    \operatorname{Index} 2 = [];
293
    Index3 = [];
294
    Index4 = [];
295
    Minima1 = [];
296
    Minima2 = [];
297
    Minima3 = [];
298
    Minima4 = [];
299
    Index4 = [];
300
    Index5 = [];
301
    satreMinima1 = [];
302
    satreMinima2 = [];
303
    satreMinima3 = [];
304
    satreMinima4 = [];
305
```

```
sotooneMinima1 = [];
306
    sotooneMinima2 = [];
307
    sotooneMinima3 = [];
308
    sotooneMinima4 = [];
309
    satreMaxima = [];
310
    Maxima8 = [];
311
    satreMaxima8 = [];
312
    Index8 = [];
313
    first_false_maxima = [];
314
315
    Minima11 = [];
316
    Minima22 = [];
317
318
    Index22 = [];
319
    satreMaxima2 = [];
320
    Maxima2 = [];
321
    for z1=1:n
322
323
          if z1<=40
324
               [\max, \operatorname{index} 1] = \max(\operatorname{Et}_{\operatorname{say}}(z1, :));
325
               Maxima=[Maxima maxima];
326
               satreMaxima =[satreMaxima z1];
327
               Index=[Index index1];
328
          elseif z1<=1400
329
               aaaaa=\operatorname{sum}(\operatorname{Et}_{\operatorname{say}}(\operatorname{z1},:)>0);
330
               if Et_say(z1, n) == 0
331
                     [maxima, index1]=max(Et_say(z1, floor(aaaaa*0.2):aaaaa-floor(
332
                         aaaaa * 0.4)));
                     index1=index1+floor(aaaaa*0.2)-1;
333
                     Maxima=[Maxima maxima];
334
                     satreMaxima =[satreMaxima z1];
335
```

336	<pre>Index = [Index index1];</pre>		
337	$elseif$ Et_say(z1,1)==0		
338	$[maxima index1] = max(Et_say(z1, n-aaaaa+floor(aaaaa*0.2):n-aaaaaa+floor(aaaaa*0.2):n-aaaaa$		
	floor(aaaaa*0.2)));		
339	index1=index1+n-aaaaa+floor(aaaaa*0.2)-1;		
340	Maxima=[Maxima maxima];		
341	<pre>satreMaxima =[satreMaxima z1];</pre>		
342	<pre>Index = [Index index1];</pre>		
343	else		
344	$[\max \operatorname{index} 1] = \max(\operatorname{Et}_{\operatorname{say}}(z1, 30:995));$		
345	index1=index1+29		
346	Maxima=[Maxima maxima];		
347	<pre>satreMaxima =[satreMaxima z1];</pre>		
348	<pre>Index = [Index index1];</pre>		
349			
350	end		
351	else		
352	$[\max \inf ex1] = \max(Et_say(z1, :));$		
353	Maxima=[Maxima maxima];		
354	<pre>satreMaxima =[satreMaxima z1];</pre>		
355	Index = [Index index 1];		
356	end		
357			
358			
359	minima1=maxima;		
360	minima2=maxima;		
361	$\operatorname{index} 2 = [];$		
362	$\operatorname{index} 3 = [];$		
363	for ii=1:index1-1		
364	if Et_say(z1, ii)>0		
365	if minimal>Et_say(z1,ii)		

```
minima1=Et_say(z1, ii);
366
                      index2=ii;
367
                  end
368
369
             end
        end
370
        index4=ii;
371
372
        for ii=index1+1:n
373
             if Et_say(z1, ii) > 0
374
                  if minima2>Et_say(z1,ii)
375
                      minima2=Et_say(z1, ii);
376
                      index3=ii;
377
                  end
378
             end
379
        end
380
381
        if abs(minima1-maxima)>1e-6
382
             Minima1=[Minima1 minima1];
383
             Minima11=[Minima11 minima1];
384
             satreMinima1 = [satreMinima1 z1];
385
             sotooneMinima1 = [sotooneMinima1 index2];
386
        else
387
388
               \min a1=0;
389
               Minimal1=[Minimal1 minima1];
390
        end
391
392
        if abs(minima2-maxima)>1e-6
393
             Minima2=[Minima2 minima2];
394
             Minima22=[Minima22 minima2];
395
             satreMinima2 = [satreMinima2 z1];
396
```

```
sotooneMinima2 = [sotooneMinima2 index3];
397
        else
398
399
               minima2=0;
400
               Minima22=[Minima22 minima2];
401
        end
402
   end
403
   %%
404
   %after finding the minuimums and maximums , i need to plot the valuee vs
405
   %tital length and length of A
406
   L1X = [];
407
   L1Y = [];
408
   z8 = 0;
409
    for z7=1:length(sotooneMinima1)
410
        z8=z8+1;
411
        l1x=XX1(satreMinima1(z8),sotooneMinima1(z7));
412
        L1X = [L1X \ l1x];
413
        l1y=YY1(satreMinima1(z8),sotooneMinima1(z7));
414
        L1Y = [L1Y \ l1y];
415
416
   end
417
418
   L2X = [];
419
   L2Y = [];
420
   z10=0;
421
    for z9=1:length(sotooneMinima2)
422
        z10=z10+1;
423
        12x=XX1(satreMinima2(z10),sotooneMinima2(z9));
424
        L2X = [L2X \ 12x];
425
        l2y=YY1(satreMinima2(z10),sotooneMinima2(z9));
426
        L2Y = [L2Y \ 12y];
427
```

```
428
    end
429
430
431
432 L4X = [];
   L4Y = [];
433
    z13=0;
434
435
    for z14=1:length(Maxima)
436
         z13=z13+1;
437
         l4x=XX1(satreMaxima(z13),Index(z14));
438
        L4X = [L4X \ l4x];
439
         l4y=YY1(satreMaxima(z13),Index(z14));
440
        L4Y = [L4Y \ 14y];
441
442
   end
443
444
   L4X2 = [];
445
   L4Y2 = [];
446
    z132=0;
447
448
    for z14=1:length(Maxima2)
449
         l4x=XX1(satreMaxima2(z14),Index22(z14));
450
        L4X2 = [L4X2 \ l4x];
451
         l4y=YY1(satreMaxima2(z14),Index22(z14));
452
        L4Y2 = [L4Y2 \ 14y];
453
    end
454
455
456
457
458
```

```
459
    hold on
460
461
    for i = 1:30
462
       hold on
463
       LT=XX1(780+i,:);
464
       LA=YY1(780+i,:);
465
       ZZZ3 = Et_say(780 + i, :);
466
       plot3 (LT, LA, ZZZ3)
467
    end
468
469
   % plot3(XX1,YY1,Et_say)
470
   % colormap('pink')
471
    plot3(L1X,L1Y,Minima1,'or')
472
    plot3(L2X,L2Y,Minima2, 'og')
473
    plot3(L4X,L4Y,Maxima, 'oc')
474
475
    bbbb=length(L4X)-sum((L4X>10.5))
476
    dddd=sum((L4X<12.8))
477
478
    plot3(L4X(bbbb:dddd),L4Y(bbbb:dddd),Maxima(bbbb:dddd),'oc')
479
480
    axis([L_T(1) \ L_T(n) \ LA(1) \ LA(n)])
481
    set(gca, 'XTick', [], 'YTick', [])
482
   % plot3 (L4X,L4Y, Maxima, 'oc')
483
484
   % plot3(L4X2,L4Y2,Maxima2,'oy')
485
486
487
   1%
488
489
```

```
%Computing enregy derivative
490
491
    Maxima = [];
492
    Index = [];
493
    Index2 = [];
494
    Index3 = [];
495
    Index4 = [];
496
    Minima1 = [];
497
    Minima2 = [];
498
    Minima3 = [];
499
    Minima4 = [];
500
    Index4 = [];
501
    Index5 = [];
502
    satreMinima1 = [];
503
    satreMinima2 = [];
504
    satreMinima3 = [];
505
    satreMinima4 = [];
506
    sotooneMinima1 = [];
507
    sotooneMinima2 = [];
508
    sotooneMinima3 = [];
509
    sotooneMinima4 = [];
510
    satreMaxima = [];
511
    Maxima8 = [];
512
    satreMaxima8 = [];
513
    Index8 = [];
514
    first_false_maxima = [];
515
516
    Minima11 = [];
517
    Minima22 = [];
518
519
```

```
Index22 = [];
521
       satreMaxima2 = [];
522
       Maxima2 = [];
523
524
       for z1=1:n
525
526
527
                if z1<=400000000
528
                         [\max, \operatorname{index} 1] = \max(\operatorname{Et}_{\operatorname{say}}(z1, :));
529
                        Maxima = [Maxima maxima];
530
                        satreMaxima =[satreMaxima z1];
531
                        Index=[Index index1];
532
                elseif z1<=1545
533
                         aaaaa=\operatorname{sum}(\operatorname{Et}_{\operatorname{say}}(\operatorname{z1},:)>0);
534
                         if Et_say(z1, 999) == 0
535
                                  [maxima, index1]=max(Et_say(z1, floor(aaaaa*0.2):aaaaa-floor(
536
                                         aaaaa *0.4)));
                                 index1=index1+floor(aaaaa*0.2)-1;
537
                                 Maxima=[Maxima maxima];
538
                                 satreMaxima =[satreMaxima z1];
539
                                 Index=[Index index1];
540
                         elseif Et_say(z1,1)==0
541
                                  [\max \inf dex1] = \max (Et_say(z1, n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa
542
                                         floor(aaaaa*0.2)));
                                 index1=index1+n-aaaaa+floor(aaaaa*0.2)-1;
543
                                 Maxima=[Maxima maxima];
544
                                 satreMaxima =[satreMaxima z1];
545
                                 Index=[Index index1];
546
                         else
547
                                  [\max \operatorname{index1}] = \max(\operatorname{Et}_{\operatorname{say}}(z1, 30:995));
548
                                 index1=index1+29
549
```

```
Maxima=[Maxima maxima];
550
                    satreMaxima =[satreMaxima z1];
551
                    Index=[Index index1];
552
553
               end
554
         else
555
               [\max \operatorname{index} 1] = \max(\operatorname{Et}_{\operatorname{say}}(z1, :));
556
              Maxima=[Maxima maxima];
557
               satreMaxima =[satreMaxima z1];
558
               Index=[Index index1];
559
         end
560
561
562
         minima1=maxima;
563
         minima2=maxima;
564
         \operatorname{index} 2 = [];
565
         index3 = [];
566
         for ii=1:index1-1
567
               if Et_say(z1, ii) > 0
568
                    if minimal>Et_say(z1, ii)
569
                         minima1 = Et_say(z1, ii);
570
                         index2=ii;
571
                    end
572
               end
573
         end
574
         index4=ii;
575
576
         for ii=index1+1:n
577
               if Et_say(z1, ii) > 0
578
                    if minima2>Et_say(z1,ii)
579
                         minima2=Et_say(z1, ii);
580
```

```
index3=ii;
581
                 end
582
             end
583
584
        end
585
        if abs(minima1-maxima)>1e-6
586
             Minima1=[Minima1 minima1];
587
             Minima11=[Minima11 minima1];
588
             satreMinima1 = [satreMinima1 z1];
589
             sotooneMinima1 = [sotooneMinima1 index2];
590
        else
591
               \min a1=0;
592
               Minimal1=[Minimal1 minimal];
593
        end
594
595
        if abs(minima2-maxima)>1e-6
596
             Minima2=[Minima2 minima2];
597
             Minima22=[Minima22 minima2];
598
             satreMinima2 = [satreMinima2 z1];
599
             sotooneMinima2 = [sotooneMinima2 index3];
600
        else
601
602
               \min a2=0;
603
               Minima22=[Minima22 minima2];
604
        end
605
606
   end
607
608
   L1X = [];
609
   L1Y = [];
610
   z8 = 0;
611
```

```
for z7=1:length(sotooneMinima1)
612
        z8=z8+1;
613
        l1x=XX1(satreMinima1(z8),sotooneMinima1(z7));
614
        L1X = [L1X \ l1x];
615
        l1y=YY1(satreMinima1(z8),sotooneMinima1(z7));
616
        L1Y = [L1Y \ l1y];
617
618
    end
619
620
   L2X = [];
621
   L2Y = [];
622
    z10=0;
623
    for z9=1:length(sotooneMinima2)
624
        z10=z10+1;
625
        l2x=XX1(satreMinima2(z10),sotooneMinima2(z9));
626
        L2X = [L2X \ 12x];
627
        l2y=YY1(satreMinima2(z10),sotooneMinima2(z9));
628
        L2Y = [L2Y \ 12y];
629
630
    end
631
632
633
634
   L4X = [];
635
   L4Y = [];
636
    z13=0;
637
638
    for z14=1:length (Maxima)
639
        z13=z13+1;
640
        l4x=XX1(satreMaxima(z13),Index(z14));
641
        L4X = [L4X \ 14x];
642
```

```
l4y=YY1(satreMaxima(z13),Index(z14));
643
        L4Y = [L4Y \ 14y];
644
645
646
    end
647
   L4X2 = [];
648
   L4Y2 = [];
649
    z132=0;
650
651
    for z14=1:length (Maxima2)
652
        14x = XX1(satreMaxima2(z14), Index22(z14));
653
        L4X2 = [L4X2 \ l4x];
654
        l4y=YY1(satreMaxima2(z14),Index22(z14));
655
        L4Y2 = [L4Y2 \ 14y];
656
    end
657
658
    minima1_length=length(Minima11);
659
    minima2_length=length(Minima22);
660
    total_minima=Minima11+Minima22;
661
662
    figure(25)
663
    plot (L<sub>-</sub>T, total_minima)
664
665
   % figure (26)
666
   % plot(L_T(2:500),df)
667
   \% dx = L_T(2) - L_T(1);
668
   \% dE = [];
669
   % for i11=2:n-1
670
   %
           dU(i11) = (total_minima(i11+1) - total_minima(i11-1)) / (2*dx);
671
   %
           dE = [dE \ dU];
672
673 % end
```

```
dx = L_T(2) - L_T(1);
674
    dE = [];
675
    for i11=2:n-1
676
        dU(i11) = (total_minima(i11+1) - total_minima(i11-1)) / (2*dx);
677
        dE = [dE \ dU];
678
    end
679
    plot(L_T(4:end-1), dU(3:end-1))
680
    figure(26)
681
    mesh(XX1,YY1,Et_say)
682
   % MyPink=pink;
683
   % id= find (MyPink(:,1) <0.2 & MyPink(:,2) <0.2 & MyPink(:,3) <0.2);
684
   % for i=1:size(id,1)
685
   %
           MyPink(id(i), :) = [1 \ 1 \ 1];
686
   % end
687
688
   % colormap(MyPink)
689
    colormap('pink')
690
    axis([L_T(1) \ L_T(n) \ LA(1) \ LA(n)])
691
    set(gca, 'XTick', [], 'YTick', [])
692
693
694
695
   % figure (27)
696
   % plot3(L1X,L1Y,Minima1,'ok')
697
   \% axis ([L_T(1) L_T(n) LA(1) LA(n)])
698
   % figure (28)
699
   % plot3 (L2X, L2Y, Minima2, 'ok')
700
   \% axis ([L<sub>-</sub>T(1) L<sub>-</sub>T(n) LA(1) LA(n)])
701
   % figure (29)
702
   % plot3(L4X(bbbb:dddd),L4Y(bbbb:dddd),Maxima(bbbb:dddd),'r')
703
```

```
704 % axis ([L_T(1) L_T(n) LA(1) LA(n)])
```

```
figure(30)
705
   \% \text{ plot}(L_T(2:500), df)
706
    dx=L_T(2)-L_T(1);
707
   dE = [];
708
    for i11=2:n-1
709
         dU(i11) = (total_minima(i11+1) - total_minima(i11-1)) / (2*dx);
710
           dE = [dE \ dU];
   %
711
    end
712
    hold on
713
714
    plot (L<sub>-</sub>T, total_minima, 'k')
715
716
    legend('energy')
717
    \operatorname{xlim}([L_T(1) \ L_T(n)])
718
    ylim ([0 350])
719
    figure(31)
720
    plot(L_T(4:end-1), dU(3:end-1), 'k')
721
722
    legend('force')
723
    axis([L_T(1) \ L_T(n) \ -2500 \ 500])
724
725
726
727
    figure(100)
728
729
730
    hold on
731
    plot3(XX1,YY1,Et_say)
732
    plot3(L1X,L1Y,Minima1,'or')
733
    plot3(L2X,L2Y,Minima2, 'og')
734
   %
735
```

- 736 bbbb=length(L4X)-sum((L4X>10.5));
- 737 dddd=sum((L4X<13));

```
738 % %
```

- 739 % plot3(L4X(bbbb:dddd),L4Y(bbbb:dddd),Maxima(bbbb:dddd),'oc')
- 740 plot3 (L4X, L4Y, Maxima, 'oc')
- 741  $axis([L_T(1) \ L_T(n) \ LA(1) \ LA(n)])$

742 set (gca, 'XTick', [], 'YTick', [])

### B.2 Theoretical Analysis MATLAB Function for Cell A Compression Diode

```
1
  function [theta_A1, theta_A2, E_A, LA_II, L_A_I_II] = CellA (LA_I, landaIa, aIa,
\mathbf{2}
       bIa, K_A, theta0_a, n, L_o, L_c)
з %
  %Designe Parameters
4
5 %Lengths
   aIIa = 1.25 * aIa;
6
   landaIIa=acos((aIa*cos(landaIa))/aIIa); % rigid-folding condition
\overline{7}
8
  %Stiffness
9
  K1_a=2*K_A(1)*bIa;
                         %KIa
10
   K2_a=2*K_A(1)*aIa;
                         %KIa
11
   K_{3-a}=2*K_{A}(2)*bIa;
                         %KIIa
12
   K4_a=2*K_A(2)*aIIa;%KIIa
13
   K6_a=2*L_c*K_A(3); %inter cellar crease stiffness . there are tow
14
       creases with theire associated stiffness in each unit cell
                       %Kc3_a
15
  K7_a=8.*K_A(4).*bIa; %Kc2_a
16
```

- 17 K8\_a=K\_A(5).\*bIa; %K\_external
- 18 %Angles
- 19  $phi10_a = pi 2.*theta0_a;$

```
phi20_a=2.*asin(cos(theta0_a)./sqrt(1-((sin(theta0_a)).^2).*(sin(landaIa)).^2)); %it should not be more than 1 or -1
```

- <sup>22</sup> phi40\_a=2.\*asin(((sin(landaIa))./(sin(landaIIa))).\*sin(phi20\_a./2));

```
23 phi50_a = (pi/2) + theta0_a; %it should not be more than 1 or -1
```

```
24 s1a=sin(landaIa);
```

```
t_{25} t_{2a=tan}(landaIIa);
```

```
26 t1a=tan(landaIa);
```

```
27
```

```
28
```

```
29 %%Calculation of theta_a
```

- 30 theta\_A1=asin(LA\_I./(aIa.\*sin(landaIa))); %theta angle of sheet I in new design
- 31 theta\_A2=acos((cos(theta\_A1).\*tan(landaIa))./(tan(landaIIa))); %theta
  angle of sheet II in new design
- 32 LA\_II=aIIa.\*sin(theta\_A2).\*sin(landaIIa); %length of sheet II of cell A in new design
- 33 L\_A\_I\_II=abs(aIa\*s1a\*((sqrt(((t2a^2)/(t1a^2))-((cos(theta\_A1).^2))))-sin (theta\_A1))); %total length of Cell A in new design base of original designe equation

		./2));
40		$phi5_a2 = (pi/2) + theta_A1;$
41		
42		$E1_a2=K1_a.*(phi1_a2-phi10_a).^2;$
43		$E2_a2=K2_a.*(phi2_a2-phi20_a).^2;$
44		$E3_a2=K3_a.*(phi3_a2-phi30_a).^2;$
45		$E4_a2=K4_a.*(phi4_a2-phi40_a).^2;$
46		$E7_a2 = K7_a.*(phi5_a2-phi50_a).^2;$
47		
48		
49		$ E_{-}A = \left(\left( E1_{-}a2 + E2_{-}a2 + E3_{-}a2 + E4_{-}a2 + E7_{-}a2 \right) / 2 \right); $
50	else	
51		$phi1_a2=0;$
52		$phi2_a2=0;$
53		$phi3_a2=0;$
54		$phi4_a2=0;$
55		$phi5_a2=0;$
56		
57		$E1_{-}a2=0;$
58		$E2_{-}a2=0;$
59		$E3_{-}a2=0;$
60		$E4_{-}a2=0;$
61		$E7_{-}a2=0;$
62		$E_A = 0;$
63		
64	end	

# B.3 Theoretical Analysis MATLAB Function for Cell B Compression Diode
- 1 function [E\_B, LB\_I\_II, theta\_B, E1\_b2, E2\_b2, E3\_b2, E4\_b2, E7\_b2]=CellB(aIb, bIb, landaIb, LA\_I, LA\_II, K\_B, theta0\_b, L2\_a\_min, L2\_a\_max, L\_A\_I\_II, n, LT, L\_o, L\_c)

```
3
```

- 4 % Design Parameters
- 5 %Lengths
- aIIb = 1.25 \* aIb;
- 7 landaIIb=acos((aIb\*cos(landaIb))/aIIb); % rigid-folding condition
- 8

```
9 %Stiffness
```

- 10  $K1_b=2*K_B(1)*bIb;$  %KIa
- 11  $K2_b=2*K_B(1)*aIb;$  %KIa
- 12  $K3_b=2*K_B(2)*bIb;$  %KIIa
- 13 K4\_b= $2*K_B(2)*aIIb$ ; %KIIa
- 14  $K6_b=2*L_c*K_B(3)$ ; %inter cellar crease stiffness . there are tow creases with their associated stiffness in each unit cell
- 15 16

```
K7_b=8.*K_B(4).*bIb; %Kc2_a
```

- <sup>17</sup> K8\_b=K\_B(1,5).\*bIb; %K\_external
- 18 %Angles
- 19 s1b=sin(landaIb);
- 20 t2b=tan(landaIIb);

```
_{21} t1b=tan(landaIb);
```

```
A=aIb*sin(landaIb);
```

<sup>23</sup> B=((tan(landaIIb))^2)/((tan(landaIb))^2);

%Kc3\_a

- theta\_b=linspace(-pi/2, pi/2, n);
- <sup>25</sup> LB\_I\_II\_range=aIb\*s1b\*((sqrt(((t2b^2)/(t1b^2))-((cos(theta\_b).^2))))-sin (theta\_b));
- <sup>26</sup>  $L_B_I_II_max = max(LB_I_II_range);$

```
L_B_I_II_min=min(LB_I_II_range);
   phi10_b=pi-2.*theta0_b;
28
   phi20_b = 2.*asin(cos(theta0_b)./sqrt(1-((sin(theta0_b)).^2).*(sin(landaIb))))))))
29
       )).^{2}); % it should not be more than 1 or -1
   phi30_b=pi-(2.*acos(tan(landaIb).*(1/(tan(landaIIb)))).*cos(theta0_b)));
30
   phi40_b = 2.*asin(((sin(landaIb))./(sin(landaIIb))).*sin(phi20_b./2));
^{31}
   phi50_b = (pi/2) + theta0_b; %it should not be more than 1 or -1
32
33
34
  %
35
  %
               if (L_A_I_I_I \leq L_B_I_I_m ax) \& \& (L_A_I_I_I > L_B_I_I_m in)
36
37 % %
                     if (L2_a_min \leq LA_II) \& \& (LA_II \leq L2_a_max)
  %
                           LB_I_II = -(LT_L_c - LA_I - L_o - L_c + LA_II);
38
39 % %
                        if (L_B_I_II_max>=LB_I_II)&&(LB_I_II>=L_B_I_II_min)
40 %
                         theta_B=real (\operatorname{asin}((A./(2.*LB_I_II))).*(B-((LB_I_II)^2))
       (A^2) - 1));
41 %
                      else
                         theta_B = 0;
42 %
43 % %
                        end
44 %
                         L_B2=aIIb.*sin(theta_B).*sin(landaIIb);
  % %
                    end
45
46
47
            if (L2_a_min \leq LA_II) \& \& (LA_II \leq L2_a_max)
48
                   if (L_A_I_I = L_B_I_I_max) \& \& (L_A_I_I) = L_B_I_I_min)
49
50
                         LB_I_II = -(LT - L_c - LA_I - L_o - L_c + LA_II);
51
                   if (L_B_I_II_max \ge LB_I_II) \& \& (LB_I_II \ge L_B_I_II_min)
52
                       theta_B=real (asin((A./(2.*LB_I_I)).*(B-((LB_I_II.^2)./
53
                          A^2) - 1)));
                   else
54
```

```
95
```

 $theta_B = 0;$ 55end 56 $L_B2=aIIb.*sin(theta_B).*sin(landaIIb);$ 57end 5859end 60i f  $(\text{theta}_B = 0)$ 6162 $phi1_b2=pi-2.*theta_B;$ 63 $phi2_b2=2.*asin(cos(theta_B)./sqrt(1-((sin(theta_B)).^2).*(sin)))$ 64(landaIb)).<sup>2</sup>);  $phi3_b2=pi-(2.*acos(tan(landaIb).*(1/tan(landaIIb))).*cos($ 65 $theta_B)));$ phi4\_b2=2.\*asin(((sin(landaIb))./(sin(landaIIb))).\*sin(phi2\_b2 66./2));  $phi5_b2 = (pi/2) + theta_B;$ 6768 69 70 $E1_b2 = K1_b . * (phi1_b2 - phi10_b) . 2;$ 71 $E2_b2 = K2_b \cdot (phi2_b2 - phi20_b) \cdot 2;$ 72 $E3_b2 = K3_b \cdot (phi3_b2 - phi30_b) \cdot 2;$ 73 $E4_b2 = K4_b . * (phi4_b2 - phi40_b) . 2;$ 74 $E7_b2 = K7_b . * (phi5_b2 - phi50_b) . 2;$ 75 $E_B = ((E1_b2 + E2_b2 + E3_b2 + E4_b2 + E7_b2)/2);$ 76 else 77 $phi1_b2=0;$ 78 $phi2_{b}2=0;$ 79 $phi3_b2=0;$ 80  $phi4_b2=0;$  $^{81}$  $phi5_b2=0;$  $^{82}$ 

```
83
^{84}
85
                  E1_b2=0;
86
                  E2_{b}2=0;
87
                  E3_{b}2=0;
88
                  E4_{-}b2=0;
89
                  E7_{b}2=0;
90
^{91}
                  E_B=0;
92
93
94 end
```

## B.4 Optimization analysis MATLAB Script

```
1
2 %%CELL A parameters
3 % a=2;
4 %
5 % b=2.8;
6 aIIa=1.25*a;
```

```
7 % landa=1.33;
```

```
s landaIIa=acos((a*cos(landa))/aIIa); % rigid-folding condition
```

```
9 KIa=1;
```

```
10 K_A=zeros (1,5);
```

```
11 K_A(1,1)=KIa;
```

```
12 % Lc=
```

n = 5;

```
^{13} L_o=Lc;
```

```
theta0_a=-pi/3;
14
15
16
  K_A(1,2) = 20 * KIa;
17
  K_A(1,3) = KIa;
18
  K_A(1, 4) = KIa;
19
  K_A(1,5) = 0;
20
   theta_a1=linspace(-pi/2, pi/2, 777);
^{21}
   theta_a1_2 = acos((cos(theta_a1).*tan(landa))./(tan(landaIIa))));
22
   theta_a1_2_0=acos((cos(theta0_a).*tan(landa))./(tan(landaIIa))));
^{23}
   s1a=sin(landa);
24
   t2a=tan(landaIIa);
25
   t1a = tan(landa);
26
   L1_a = a.*sin(theta_a1).*sin(landa);
27
   L2_a = aIIa . * sin(theta_a1_2) . * sin(landaIIa);
^{28}
   n = length(L1_a);
29
   L2_a_max = max(L2_a);
30
   L1_a_max = max(L1_a);
^{31}
   L2_a_min=min(L2_a);
```

```
38
```

33

34

35

36

37

```
%%
39
```

%%CELL B parameters 40

design

 $L1_a_{\min}(L1_a);$ 

LA\_min=L2\_a\_min+L1\_a\_min+Lc;

LA\_max=L2\_a\_max+L1\_a\_max+Lc;

% LA=linspace(LA\_min,LA\_max,n);

- aIb=a;41
- bIb=aIb; 42
- aIIb = 1.25 \* aIb;43

LA\_I=linspace(L1\_a\_min, L1\_a\_max, length(L1\_a)); %length of sheet I in new

```
44 landaIb=landa;
45 landaIIb=acos((aIb*cos(landaIb))/aIIb); % rigid-folding condition
46 KIb=1;
47 K_B=zeros(1,5);
```

```
48 K_B(1,1)=KIb;
```

```
50 theta0_b=pi/3;
```

```
theta_b2_0=acos((cos(theta0_b).*tan(landaIb))./(tan(landaIIb)));
```

```
53 K_B(1,2)=20*KIb;
```

```
54 K_B(1,3)=KIb;
```

```
55 K_B(1, 4)=KIb;
```

```
56 K_B(1,5)=0;
```

```
57 \text{theta}_{b1} = \text{linspace}(-\text{pi}/2, \text{pi}/2, n);
```

```
58 theta_b2=acos((cos(theta_b1).*tan(landaIb))./(tan(landaIIb)));
```

```
59 s1b=sin(landaIb);
```

```
60 t2b=tan(landaIIb);
```

```
61 t1b=tan(landaIb);
```

```
62 L1_b=aIb.*sin(theta_b1).*sin(landaIb);
```

```
63 L2_b=aIIb .*sin(theta_b2).*sin(landaIIb);
```

```
{}_{64}\quad L_{-}b{=}Lc{+}L1_{-}b{-}L2_{-}b\ ;
```

```
65 L2_b_max = max(L2_b);
```

```
66 L1_b_max=max(L1_b);
```

```
_{67} L2_b_min=min(L2_b);
```

```
68 L1_b_min=min(L1_b);
```

```
69 %
```

```
70 %Total length CEll A_I_II in original design
```

```
71 A=a*sin(landa);
```

```
72 B=((\tan(\tan(\tan(\tan)))^2)/((\tan(\tan()))^2);
```

```
73
```

```
74
```

```
75 L_a = a * s1a * ((sqrt(((t2a^2)/(t1a^2)) - ((cos(theta_a1).^2)))) - sin(theta_a1)))
           L_A_{max} = \max(L_a);
  76
             L_A_{\min}(L_a);
  77
            %%
  78
            %Mismatch parameters
  79
             K_{-}star = 50;
  80
  81
             say_A0=2*atan(cos(theta0_a)*tan(landa));
  82
             say_B0=2*atan(cos(theta0_b)*tan(landaIb));
  83
  84
            n = length(L1_a);
  85
            %%
  86
  87
  88
            %%
  89
            %Calculating maximum and minium
  90
             i5 = 0;
  91
             for i3=linspace(-pi/2, pi/2, n)
  92
  93
                             i5=i5+1;
  94
                             i6 = 0;
  95
                             for i4 = linspace(-pi/2, pi/2, n)
  96
                                             i6 = i6 + 1;
  97
                                             theta_a1_2(i5) = acos((cos(i3).*tan(landa))./(tan(landaIIa)));
  98
                                             theta_b_2(i6) = acos((cos(i4).*tan(landaIb))./(tan(landaIIb)));
  99
                                             total_length(i5, i6) = (Lc+(a.*sin(i3).*sin(landa)) - abs((aIIa.*sin(basis))) - abs(basis)) - abs(basis)) - abs(basis) -
100
                                                          theta_a1_2(i5)).* sin(landaIIa)))))+(Lc+(aIb.* sin(i4).* sin(
                                                         landaIb))-abs((aIIb.*sin(theta_b_2(i6)).*sin(landaIIb))))+
                                                         L_0;
```

end

103	end
104	$L_total_min=min(min(total_length));$
105	$L_total_max = max(max(total_length));$
106	$L_T = linspace(L_total_min, L_total_max, n);$ %New design total length
107	78%
108	
109	i1 = 0;
110	for LT=linspace(L_total_min,L_total_max,n)
111	i1=i1+1;
112	for $i2=1:length(L1_a)$
113	$[$ theta_A1(i1,i2), theta_A2(i1,i2), E_A(i1,i2), LA_III(i1,i2),
	$L_A_I_I(i1,i2)] = CellA_opt(LA_I(i2), landa, a, b, K_A, theta0_a, n)$
	$, L_{-}o$ $, Lc$ $);$
114	$[E_B(i1,i2), LB_II_I(i1,i2), theta_B(i1,i2), E1_b2(i1,i2), E2_b2(i1,i2), E2_b2(i1,i2$
	$i2$ ), $E3_b2(i1,i2)$ , $E4_b2(i1,i2)$ , $E7_b2(i1,i2)$ ] = CellB_opt(aIb,
	$aIIb\ , landaIb\ , LA\_I(i2)\ , LA\_II(i1\ , i2)\ , K\_B\ , theta0\_b\ , L2\_a\_min\ ,$
	$L2_{a}max$ , $L_{A}_{I}II(i1,i2)$ , n, $LT$ , $L_{o}$ , $Lc$ );
115	$if E_B(i1, i2) = 0$
116	$say_A(i1,i2)=2*atan((cos(theta_A1(i1,i2)))*tan(landa));$
117	$say_B(i1,i2)=2*atan((cos(theta_B(i1,i2)))*tan(landaIb));$
118	$E_{say}(i1,i2) = K_{star*b*}(((say_A(i1,i2)-say_B(i1,i2)).^2)./2);$
119	$E6_a2(i1,i2) = (K_A(3).*(say_A(i1,i2)-say_A0).^2)./2;$
120	$E6_b2(i1,i2) = (K_B(3).*(say_B(i1,i2)-say_B0).^2)./2;$
121	$E8_{A}B(i1,i2) = (K_A(5) . * (say_B(i1,i2) - say_A(i1,i2)) . 2) . 2;$
122	
123	$Et_A_B(i1,i2) = E_A(i1,i2) + E_B(i1,i2);$
124	$Et_{say}(i1, i2) = Et_A_B(i1, i2) + E_{say}(i1, i2) + E_{a}A_B(i1, i2) + E6_{a}(i1, i2) +$
	$i1, i2$ )+E6_b2( $i1, i2$ );
125	else
126	$say_{A}(i1,i2)=0;$

```
say_B(i1, i2) = 0;
127
                  E_{-say}(i1, i2) = 0;
128
                  E6_a2(i1,i2)=0;
129
                  E6_b2(i1, i2) = 0;
130
                  E8_A_B(i1, i2) = 0;
131
                  Et_A_B(i1, i2) = 0;
132
                  Et_{-say}(i1, i2) = 0;
133
134
              end
135
         end
136
    end
137
    %%
138
    \% {\rm zero\_column} of CEll A
139
    kh = pi/4;
140
141
142
    if abs(landa-kh) > 0
143
    \operatorname{zero_column_n} = (n+1)/2;
144
145
    lt_0column=(linspace(L_total_min,L_total_max,n))';
146
147
148
149
150
151
152
153
    LA1_0column=0;
154
    la_0column=LA1_0column;
155
    theta_A1_0column = zeros(n,1);
156
    theta_A1_Ocolumn(:,:) = asin(LA1_Ocolumn./(a.*sin(landa)));
157
```

```
102
```

```
theta_A2_0column=zeros(n,1);
158
   theta_A2_0column(:,:) = acos((cos(theta_A1_0column).*tan(landa))./(tan(
159
       landaIIa)));
   LA2_Ocolumn=aIIa.*sin(theta_A2_Ocolumn).*sin(landaIIa);
160
161
162
163
164
   %
165
   %%Designe Parameters
166
   %Lengths
167
   aIIa = 1.25 * a;
168
   landaIIa=acos((a*cos(landa))/aIIa); % rigid-folding condition
169
170
   %Stiffness
171
   K1_a=2*K_A(1)*b;
                       %KIa
172
   K2_a=2*K_A(1)*a;
                       %KIa
173
   K3_a=2*K_A(2)*b;
                       %KIIa
174
   K4_a=2*K_A(2)*aIIa;%KIIa
175
   K6_a=2*Lc*K_A(3); %inter cellar crease stiffness . there are tow
176
       creases with their associated stiffness in each unit cell
                       %Kc3_a
177
   K7_a=8.*K_A(4).*b; \% Kc2_a
178
   K8_a=K_A(5).*b; \%K_external
179
   %Angles
180
   phi10_a = pi - 2.* theta0_a;
181
   phi20_a = 2.*asin(cos(theta0_a))/sqrt(1-((sin(theta0_a)))^2).*(sin(landa))
182
       (^2); % it should not be more than 1 or -1
   phi30_a = pi - (2 \cdot (a \cos(tan(landa) \cdot (1/(tan(landaIIa))) \cdot (cs(theta0_a))));
183
   phi40_a = 2.*asin(((sin(landa))./(sin(landaIIa))).*sin(phi20_a./2));
184
   phi50_a = (pi/2) + theta0_a; %it should not be more than 1 or -1
185
```

```
103
```

```
phi1_a2_0column=pi-2.*theta_A1_0column;
188
    phi2_a2_0column = 2.*asin(cos(theta_A1_0column))./sqrt(1-((sin(
189
        theta_A1_Ocolumn)).^2).*(sin(landa)).^2);
    phi3_a2_0column=pi-(2.*acos(tan(landa).*(1/(tan(landaIIa))).*cos(
190
        theta_A1_Ocolumn)));
    phi4_a2_0column=2.*asin(((sin(landa))./(sin(landaIIa))).*sin(
191
        phi2_a2_0column./2);
    phi5_a2_0column = (pi/2) + theta_A1_0column;
192
193
    E1_a2_0column=K1_a.*(phi1_a2_0column-phi10_a).^2;
194
    E2_a2_0column=K2_a.*(phi2_a2_0column-phi20_a).^2;
195
    E3_a2_0column=K3_a.*(phi3_a2_0column-phi30_a).^2;
196
    E4_a2_0column=K4_a.*(phi4_a2_0column-phi40_a).^2;
197
    E7_a2_0column=K7_a.*(phi5_a2_0column-phi50_a).^2;
198
199
200
   E_A_0column=((E1_a2_0column+E2_a2_0column+E3_a2_0column+E4_a2_0column+
201
       E7_a2_0column)/2);
   E_A_constant_column = zeros(n,1);
202
   E_A_constant_column (:,:)=E_A_0column;
203
204
   Et_say_0column = [];
205
   for i10=1:n
206
        if Et_say(i10, zero_column_n)~=0
207
          Edummy=Et_say(i10, zero_column_n)+E_A_Ocolumn(i10);
208
          Et_say_0column = [Et_say_0column; Edummy];
209
        else
210
          Edummy=0;
211
          Et_say_0column = [Et_say_0column; Edummy];
212
```

```
213
         end
    end
214
    Et_say(:,zero_column_n)=Et_say_0column;
215
    end
216
217
   %%
218
   %
219
   %%Plotting results
220
    for i=1:n
221
         L1_{a_2}(i) = a \cdot sin(theta_A1(1,i)) \cdot sin(landa);
222
         L2_{a_2}(i) = aIIa . * sin(theta_A2(1,i)) . * sin(landaIIa);
223
        LA(i) = L1_a_2(i) + L2_a_2(i) + Lc;
224
    end
225
226
    [X1,Y1] = meshgrid(L_T,LA);
227
   XX1=X1.';
228
    YY1=Y1.';
229
230
231
232
233
234
   %%
235
236
   %Computing enregy derivative
237
238
    Maxima = [];
239
    Index = [];
240
    Index2 = [];
241
    Index3 = [];
242
    Index4 = [];
243
```

```
_{244} Minima1 = [];
```

- $_{245}$  Minima2 = [];
- 246 Minima3 = [];
- 247 Minima4 = [];
- $_{248}$  Index4 = [];
- $_{249}$  Index5 = [];
- $_{250}$  satreMinima1 = [];
- 251 satreMinima2 = [];
- $_{252}$  satreMinima3 = [];
- $_{253}$  satreMinima4 = [];
- $_{254}$  sotooneMinima1 = [];
- $_{255}$  sotooneMinima2 = [];
- $_{256}$  sotooneMinima3 = [];
- $_{257}$  sotooneMinima4 = [];
- $_{258}$  satreMaxima = [];
- 259 Maxima8 = [];

```
_{260} satreMaxima8 = [];
```

 $_{261}$  Index8 = [];

```
_{262} \quad \text{first\_false\_maxima} = [];
```

263

```
264 Minima11 = [];
```

```
265 Minima22 = [];
```

```
266
```

```
_{268} Index22 = [];
```

```
_{269} satreMaxima2 = [];
```

```
270 Maxima2 = [];
```

```
271
```

```
272 n1=ceil(n*1545/1555);
```

```
273 n2=ceil(n*30/1555);
```

```
274 n3 = ceil(n*995/1555);
```

275	
276	for z1=1:n
277	
278	
279	if $z1 <= 4000000000$
280	$[\max, \operatorname{index} 1] = \max(\operatorname{Et}_{\operatorname{say}}(z1, :));$
281	Maxima=[Maxima maxima];
282	<pre>satreMaxima =[satreMaxima z1];</pre>
283	Index = [Index index 1];
284	elseif z1<=n1
285	$aaaaa = sum(Et_say(z1,:) > 0);$
286	$if$ Et_say(z1,n)==0
287	$[\max , \operatorname{index} 1] = \max (Et_say (z1, floor (aaaaa * 0.2): aaaaa - floor ($
	aaaaa*0.4))));
288	index1=index1+floor(aaaaa*0.2)-1;
289	Maxima=[Maxima maxima];
290	<pre>satreMaxima =[satreMaxima z1];</pre>
291	Index = [Index index1];
292	$elseif$ Et_say(z1,1)==0
293	$[\max \text{ index 1}] = \max(\text{Et}_{say}(z1, n-aaaaa+floor(aaaaa*0.2):n-aaaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa+floor(aaaaa*0.2):n-aaaaa$
	floor(aaaaa*0.2)));
294	index1=index1+n-aaaaa+floor(aaaaa*0.2)-1;
295	Maxima=[Maxima maxima];
296	<pre>satreMaxima =[satreMaxima z1];</pre>
297	<pre>Index=[Index index1];</pre>
298	else
299	$[\max \operatorname{index} 1] = \max(\operatorname{Et}_{-}\operatorname{say}(z1, n2: n3));$
300	index1=index1+n2-1
301	Maxima=[Maxima maxima];
302	<pre>satreMaxima =[satreMaxima z1];</pre>
303	<pre>Index=[Index index1];</pre>

```
304
             end
305
         else
306
             [\max index1] = \max(Et_say(z1,:));
307
             Maxima=[Maxima maxima];
308
             satreMaxima =[satreMaxima z1];
309
             Index=[Index index1];
310
        end
311
312
313
        minima1=maxima;
314
        minima2=maxima;
315
        index2 = [];
316
        index3 = [];
317
         for ii=1:index1-1
318
             if Et_say(z1, ii)>0
319
                  if minimal>Et_say(z1, ii)
320
                       minima1=Et_say(z1, ii);
321
                       index2=ii;
322
                  end
323
             end
324
        end
325
        index4=ii;
326
327
         for ii=index1+1:n
328
             if Et_say(z1, ii) > 0
329
                  if minima2>Et_say(z1, ii)
330
                       minima2=Et_say(z1, ii);
331
                       index3=ii;
332
                  end
333
             end
334
```

```
end
335
336
        if abs(minima1-maxima)>1e-6
337
             Minima1=[Minima1 minima1];
338
             Minimal1=[Minimal1 minima1];
339
             satreMinima1 = [satreMinima1 z1];
340
             sotooneMinima1 = [sotooneMinima1 index2];
341
        else
342
343
               \min a1=0;
344
               Minimal1=[Minimal1 minimal];
345
        end
346
347
        if abs(minima2-maxima)>1e-6
348
             Minima2=[Minima2 minima2];
349
             Minima22=[Minima22 minima2];
350
             satreMinima2 =[satreMinima2 z1];
351
             sotooneMinima2 = [sotooneMinima2 index3];
352
        else
353
354
               \min a2=0;
355
               Minima22=[Minima22 minima2];
356
        end
357
358
359
   end
360
361
362
363
    minima1_length=length(Minima11);
364
   minima2_length=length(Minima22);
365
```

```
109
```

```
total_minima=Minima11+Minima22;
366
367
368
369
370
371
    dx=L_{-}T(2)-L_{-}T(1);
372
    dE = [];
373
    for i11=2:n-1
374
         dU(i11) = (total_minima(i11+1) - total_minima(i11-1)) / (2*dx);
375
   %
           dE = [dE \ dU];
376
    \operatorname{end}
377
378
   % figure (31)
379
   \% \text{ plot}(L_T(4: \text{end} - 1), dU(3: \text{end} - 1), 'k')
380
   %
381
   % legend('force')
382
   \% axis ([L_T(1) L_T(n) -2500 500])
383
384
385
   %%surface area
386
   %cell a
387
    A1_A=4*a*b*sin(landa)/2;
388
    A2_A=4*aIIa*b*sin(landaIIa)/2;
389
    A3_A=2*b*Lc;
390
391
392
    %cell b
393
394
    A1_B=4*aIb*bIb*sin(landaIb)/2;
395
    A2_B=4*aIIb*bIb*sin(landaIIb)/2;
396
```

```
110
```

```
A3_B=2*b*Lc;
397
398
   % connecting sheets
399
400
   A4_C=2*b*Lc;
401
402
   %total surface area
403
404
   A_total=A1_A+A2_A+A3_A+A1_B+A2_B+A3_B+A4_C;
405
406
407
   %%Finding Fe and Fc
408
   nn=numel(dU);
409
410
   %total surface area
411
412
        A_total=A1_A+A2_A+A3_A+A1_B+A2_B+A3_B+A4_C;
413
414
415
        %%Finding Fe and Fc
416
        myArray=dU(100:nn-20);
417
   %
          dU(1:15) = [];
418
          dU(end:-1:nn-15) = [];
   %
419
   %
          n_reduced=numel(dU)
420
        [Fc_max, Fc_max_index] = min(myArray);
421
        fe_rnage=myArray(1:Fc_max_index-20);
422
        Fe_max=max(fe_rnage);
423
        \% length_portion=L_T(end)-L_T(1);
424
        %
425
        \% portion=L_T(1)+0.4*length_portion;
426
        % idx=length(L_T)-sum(L_T>portion);
427
```

428
429
430
431 % index\_fe=idx+1;
432 % index\_fc=idx+2;
433 % fe\_rnage=dU(3:index\_fe);
434 % fc\_range=dU(index\_fe+1:end-1);
435 % Fe\_max=max(fe\_rnage);
436 % Fc\_max=min(fc\_range);

437 for  $ce_ratio = abs(Fc_max)/Fe_max;$ 

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