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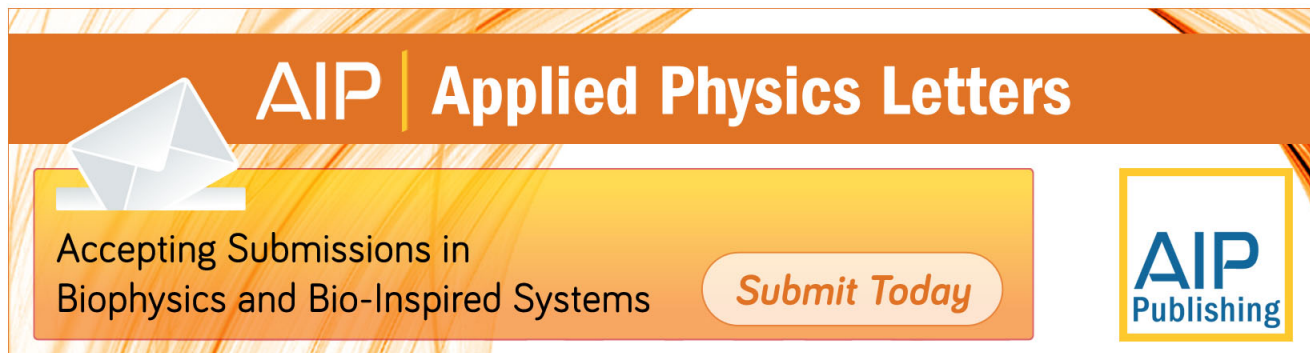
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# Disorder induced Coulomb gaps in graphene constrictions with different aspect ratios

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We present electron transport measurements on lithographically defined and etched graphene nanoconstrictions with different aspect ratios including different lengths ( $l$ ) and widths ( $w$ ). A roughly length-independent disorder induced effective energy gap can be observed around the charge neutrality point. This energy gap scales inversely with the width even in regimes where the length of the constriction is smaller than its width ( $l < w$ ). In very short constrictions, we observe both resonances due to localized states or charged islands and an elevated overall conductance level ( $0.1 - 1e^2/h$ ), which is strongly length-dependent in the gap region. This makes very short graphene constrictions interesting for highly transparent graphene tunneling barriers. © 2011 American Institute of Physics. [doi:10.1063/1.3544580]

Graphene nanoribbons and constrictions<sup>1-11</sup> are promising candidates to overcome the gapless nature of graphene.<sup>12</sup> Graphene, a truly two-dimensional (2D) semimetal with unique electronic properties, is becoming increasingly interesting for high mobility ultrafast nanoelectronics.<sup>13</sup> However, the missing band gap in graphene makes it difficult to realize graphene switches, transistors, and logic devices in general. By tailoring graphene into narrow ribbons or constrictions, the formation of an energy gap has been shown.

Graphene constrictions have been successfully demonstrated to act as tunneling barriers in single electron transistors,<sup>14</sup> quantum dots,<sup>15-17</sup> and more recently in double quantum dots.<sup>18-20</sup> The nature of the observed energy gaps has been discussed in great detail. From the theoretical point of view, different models have been put forward to describe the observed energy gap including quasi-one-dimensional (quasi-1D) Anderson localization,<sup>21-26</sup> percolation,<sup>27</sup> or Coulomb blockade on a series of localized states.<sup>28</sup> Experimentally, the nanoribbon width-dependence of the energy gap has been studied by a number of groups and it has been found that the gap size is mainly given by the charging energy of small isolated islands forming in the nanoribbon due to disorder.<sup>6,7</sup> However, only recently has the influence of the nanoribbon length,  $l$ , been investigated systematically for nanoribbons with width  $w < 40$  nm.<sup>10,11</sup> So far, the experimentally studied graphene nanostructures exhibited an aspect ratio of  $l > w$ , such that the width was the smallest and dominating length scale.

Here, we present a systematic study of electronic transport in graphene constrictions with different aspect ratios including both different widths and lengths. In particular, we focus on graphene constrictions in the regime of  $l \approx w$  and  $l < w$  [see Fig. 1(a)]. We show that in these regimes, the Coulomb charging energy  $E_g$  of the smallest disorder induced charged island in the constriction<sup>6,7</sup> is still predominantly determined by the width. However, the overall conductance level is strongly influenced by the length of the constriction. Most interestingly, we show that in very short

constrictions, Coulomb blockade diamonds with high transparency can be observed.

The samples have been fabricated from graphene that has been mechanically exfoliated from natural bulk graphite and deposited onto 295 nm SiO<sub>2</sub> on a highly p-doped Si substrate. Raman spectroscopy is used to verify the single-layer character of the investigated graphene nanostructures.<sup>29</sup> Electron-beam (e-beam) lithography is used to pattern the etch mask ( $\approx 110$  nm of e-beam resist) for structuring the graphene devices.

Reactive ion etching based on an Ar/O<sub>2</sub> plasma is employed to remove unprotected graphene. Scanning force microscope images of etched graphene nanostructures after removing the residual resist are shown in Fig. 1(b). Finally, the

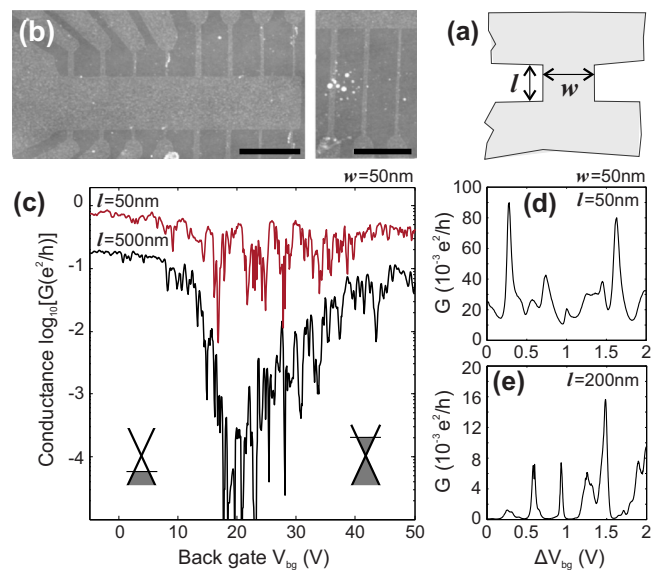


FIG. 1. (Color online) (a) Illustration of a constriction with length  $l$  and width  $w$ . (b) Scanning force microscope images of etched graphene constrictions with different aspect ratios. The length of the scale bars is 1  $\mu\text{m}$ . (c) Back gate characteristics at  $V_b = 500 \mu\text{V}$  for two graphene constrictions with width  $w = 50$  nm and different lengths  $l = 50$  nm (upper curve) and  $l = 500$  nm (lower curve). [(d) and (e)] Conductance peaks in the transport gap region for two different constrictions,  $w = l = 50$  nm (d) and  $w = l/4 = 50$  nm (e) at  $V_b = 100 \mu\text{V}$ .

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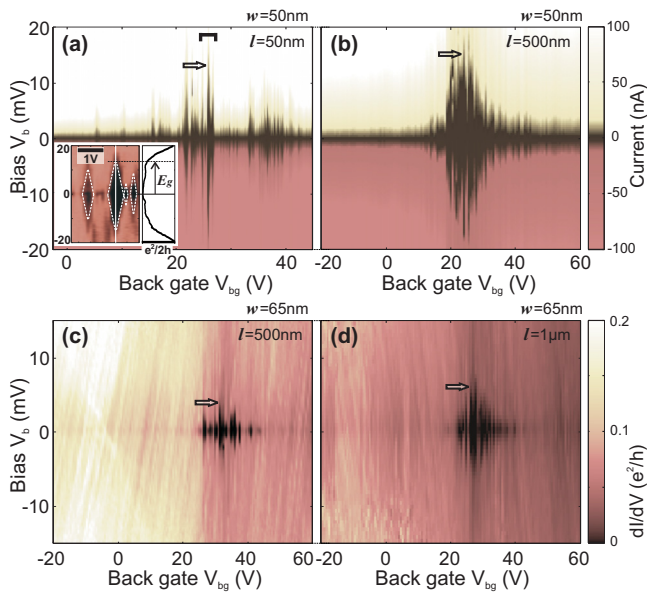


FIG. 2. (Color online) [(a) and (b)] Color plot of the source-drain current as a function of  $V_{bg}$  and  $V_b$  for two constrictions with  $w=50$  nm and  $l=50$  nm (a) and  $l=500$  nm (b). The inset shows differential conductance measurements of the region where the transport gap is most pronounced [see black bar in panel (a)]. A cross section through the center of the largest diamond is shown on the right side of the inset, highlighting the definition of the gap  $E_g$ . [(c) and (d)] Differential conductance maps for graphene constrictions with  $w=65$  nm and length  $l=500$  nm (c) and  $l=1$  μm (d).

graphene devices are contacted by e-beam patterned 5 nm Cr and 50 nm Au electrodes.

The measurements have been performed in a pumped <sup>4</sup>He system at  $T \approx 1.3$  K. We have measured the current through the two-terminal constrictions by applying a small symmetric dc bias voltage  $V_b$ . The differential conductance has been measured directly by low-frequency lock-in techniques by adding an ac bias of 100 μV.

In Fig. 1(c), we show the conductance  $G$  as function of the back gate (BG) voltage  $V_{bg}$  for two 50 nm wide constrictions with different lengths  $l=50$  nm and  $l=500$  nm ( $V_b=500$  μV). The transport gap, i.e., the  $V_{bg}$  region of suppressed conductance (roughly between 17 and 30 V), separates hole from electron transport, as indicated by the two lower inserts in Fig. 1(c).<sup>30</sup> While for the 500 nm constriction the conductance is strongly suppressed in this regime (down to  $10^{-5}e^2/h$ ), we observe a significantly increased conductance for the shorter constriction. Nevertheless, a number of pronounced resonances reaching lower conductance values mark the gap region, in rather good agreement with the measurement on the longer constriction [compare also Figs. 2(a) and 2(b)]. A close-up of such resonances is shown in Fig. 1(d). These reproducible resonances have been taken at  $V_b=100$  μV within the BG range of minimum conductance. We observe neither complete pinch-off nor strong Coulomb blockade behavior, as is found for longer constrictions with the same width. For example, in Fig. 1(e) we show data recorded on a 200 nm long ( $w=50$  nm) constriction. The overall conductance level differs significantly; however, the typical  $V_{bg}$  spacing between the resonances is comparable.

In Figs. 2(a) and 2(b), we show 2D plots of current as a function of bias and back gate voltage for the two 50 nm wide constrictions with  $l=50$  nm (a) and 500 nm (b). In

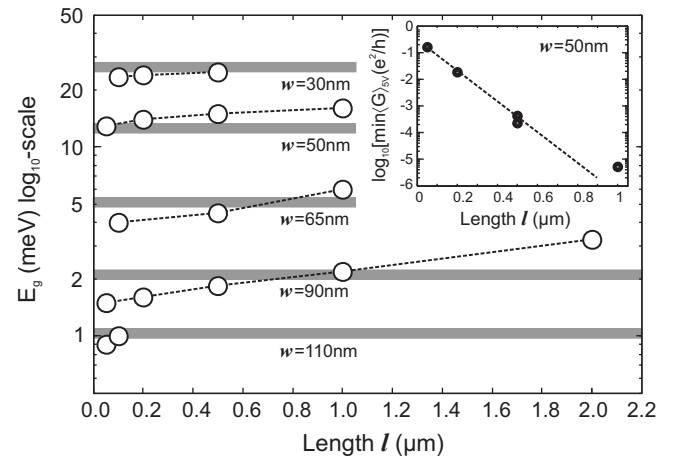


FIG. 3. Energy gap  $E_g$  as a function of length  $l$  for a number of different graphene constrictions with different widths  $w$  (see labels and data points connected by dashed lines). The horizontal gray lines represent  $E_g = a/we^{-bw}$ , where  $a=2$  eV and  $b=0.026$  nm<sup>-1</sup> are taken from Ref. 9. The inset shows the minimum value of the running averaged conductance  $G$  as function of length (running average interval 5 V in back gate voltage).

good agreement with earlier studies,<sup>2,5–11</sup> we find regions of suppressed current with an extension of  $E_g/e$  in the  $V_b$  direction (see arrows). By plotting higher-resolution measurements of the differential conductance, one can see that the region of suppressed current is composed of individual diamonds [see inset in Fig. 2(a)]. The energy gap  $E_g$  corresponds to the charging energy of the largest diamond,<sup>6,7</sup> as highlighted in the inset [see arrow in Fig. 2(a)]. Interestingly, this energy scale only weakly depends on the constriction length. By comparing Figs. 2(a) and 2(b), we observe only a small difference in  $E_g$  between the 50 nm ( $E_g=13$  meV) and 500 nm ( $E_g=15$  meV) long constriction, allowing the conclusion that the smallest island or localized state is predominantly a function of the width,  $w$ . The main difference can be found in the back gate coverage of the observed gaps. The shorter the constriction, the fewer islands [or localized (edge) states] are in the constriction, leading to fewer charging events. Consequently, the current is suppressed in smaller and fewer gate voltage regimes. However, the smallest island is found to be roughly length-independent (this is also true for  $l < w$ ). The definition of the transport gap in BG voltage,<sup>5,9</sup> used as a figure of merit in earlier work, is hard to define for very short constrictions since it is considered to strongly depend on the disorder potential,<sup>6,7,10</sup> which becomes very sample dependent, due to lack of averaging.

Similar behavior can also be seen for a 65 nm wide, 500 nm, and 1 μm long constriction, as shown in Figs. 2(c) and 2(d), respectively. Here, we plot the differential conductance as a function of  $V_b$  and  $V_{bg}$ . In both measurements, a maximum charging energy of around 4.5 and 6 meV can be observed (see arrows). In total, we studied roughly 20 graphene constrictions on three different samples. In Fig. 3, we summarize the extracted energy gaps  $E_g$  as a function of the length for five different widths.<sup>31</sup> It can be seen that  $E_g$  strongly depends on the width and we find good agreement with earlier experiments and theoretical models.<sup>11,28</sup> In particular, we compare our results with the model from Ref. 28, where the energy gap is approximated by  $E_g = a/we^{-bw}$  (see gray bars in Fig. 3) (see also Refs. 2, 9, and 28). In contrast to the weak  $E_g$ -length dependence, we observe a rather strong length dependence of the minimum conductivity, as

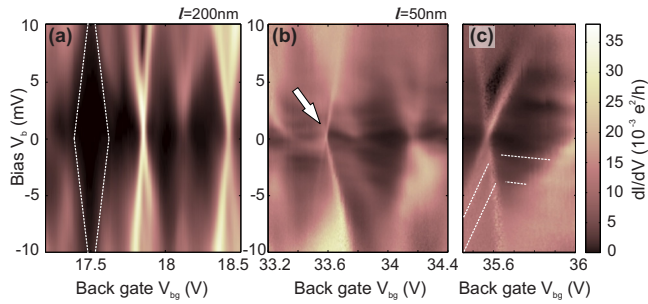


FIG. 4. (Color online) Differential conductance vs BG voltage and bias voltage of (a) 200 nm and [(b) and (c)] 50 nm long graphene constriction with a width of 50 nm. (a) Distinct diamonds of fully suppressed conductance can be observed (see white dashed line). [(b) and (c)] Similar data taken on a 50 nm long constriction. Diamonds at an elevated conductance level can be observed. Sharp resonances [white arrow in panel (b)] and features of increased conductance inside the diamonds [horizontal dashed lines in panel (c)] can be observed.

shown by the inset in Fig. 3. The minimum value of the running averaged conductance decreases exponentially with increasing constriction length. This fits well with the scenario where localizations or tunneling processes are dominating the transport through the constrictions. Moreover, it shows that by making graphene constrictions very short, we can obtain conductance level close to  $e^2/h$ .

More insights on the transparency of the shorter constrictions is gained by focusing on a small BG voltage range, as shown in Fig. 4. Here, we show high-resolution differential conductance  $dI/dV_b$  plots for different constriction lengths with constant width ( $w=50$  nm). Measurements on a 200 nm long constriction [Fig. 4(a)] show well-distinguishable diamonds of suppressed conductance, in good agreement with earlier studies. In shorter constrictions [ $w=50$  nm; Figs. 4(b) and 4(c)], we observe diamonds where the conductance is not fully suppressed, most likely because of strong coupling to localized states or isolated charged islands. Moreover, we observe faint lines of increased conductance inside the diamonds aligned with features outside the diamonds, which might be due to inelastic cotunneling [see, e.g., dashed lines in Fig. 4(c)]. According to recent experiments by Han *et al.*,<sup>11</sup> the average hopping length  $L_c$  is in the range of  $w \leq L_c < 2w$ . Consequently, it is likely that for very short constrictions,  $L_c$  exceeds  $l$  ( $l < w \leq L_c$ ) so that transport in short constrictions is no longer effectively one dimensional, resulting in several transport channels in the constrictions. This would account for (i) the strong increase in conductance and (ii) the appearance of very sharp features in the conductance [see, e.g., arrow in Fig. 4(b)], which could result from interference effects or Fano resonances.

In summary, we presented transport measurements on graphene nanoconstrictions with different aspect ratios. We showed that the strongly width-dependent disorder induced energy gap is roughly length-independent, whereas the overall conductance level depends strongly on the length. In the gap region of very short graphene constrictions, a conductance level close to  $0.1e^2/h$  can be reached, making these structures potential candidates for exploring Fano resonances

and Kondo physics in graphene nanostructures.

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- <sup>1</sup>Z. Chen, Y.-M. Lin, M. Rooks, and P. Avouris, *Physica E (Amsterdam)* **40**, 228 (2007).
- <sup>2</sup>M. Y. Han, B. Özyilmaz, Y. Zhang, and P. Kim, *Phys. Rev. Lett.* **98**, 206805 (2007).
- <sup>3</sup>X. Li, X. Wang, L. Zhang, S. Lee, and H. Dai, *Science* **319**, 1229 (2008).
- <sup>4</sup>X. Wang, Y. Ouyang, X. Li, H. Wang, J. Guo, and H. Dai, *Phys. Rev. Lett.* **100**, 206803 (2008).
- <sup>5</sup>F. Molitor, A. Jacobsen, C. Stampfer, J. Güttinger, T. Ihn, and K. Ensslin, *Phys. Rev. B* **79**, 075426 (2009).
- <sup>6</sup>C. Stampfer, J. Güttinger, S. Hellmüller, F. Molitor, K. Ensslin, and T. Ihn, *Phys. Rev. Lett.* **102**, 056403 (2009).
- <sup>7</sup>K. Todd, H.-T. Chou, S. Amasha, and D. Goldhaber-Gordon, *Nano Lett.* **9**, 416 (2009).
- <sup>8</sup>X. L. Liu, J. B. Oostinga, A. F. Morpurgo, and L. M. K. Vandersypen, *Phys. Rev. B* **80**, 121407(R) (2009).
- <sup>9</sup>F. Molitor, C. Stampfer, J. Güttinger, A. Jacobsen, T. Ihn, and K. Ensslin, *Semicond. Sci. Technol.* **25**, 034002 (2010).
- <sup>10</sup>P. Gallagher, K. Todd, and D. Goldhaber-Gordon, *Phys. Rev. B* **81**, 115409 (2010).
- <sup>11</sup>M. Y. Han, J. C. Brant, and P. Kim, *Phys. Rev. Lett.* **104**, 056801 (2010).
- <sup>12</sup>A. K. Geim and K. S. Novoselov, *Nature Mater.* **6**, 183 (2007).
- <sup>13</sup>M. I. Katsnelson, *Mater. Today* **10**, 20 (2007).
- <sup>14</sup>C. Stampfer, E. Schurtenberger, F. Molitor, J. Güttinger, T. Ihn, and K. Ensslin, *Nano Lett.* **8**, 2378 (2008).
- <sup>15</sup>L. A. Ponomarenko, F. Schedin, M. I. Katsnelson, R. Yang, E. W. Hill, K. S. Novoselov, and A. K. Geim, *Science* **320**, 356 (2008).
- <sup>16</sup>J. Güttinger, C. Stampfer, F. Libisch, T. Frey, J. Burgdörfer, T. Ihn, and K. Ensslin, *Phys. Rev. Lett.* **103**, 046810 (2009).
- <sup>17</sup>J. Moser and A. Bachtold, *Appl. Phys. Lett.* **95**, 173506 (2009).
- <sup>18</sup>F. Molitor, S. Dröscher, J. Güttinger, A. Jacobsen, C. Stampfer, T. Ihn, and K. Ensslin, *Appl. Phys. Lett.* **94**, 222107 (2009).
- <sup>19</sup>S. Moriyama, D. Tsuya, E. Watanabe, S. Uji, M. Shimizu, T. Mori, T. Yamaguchi, and K. Ishibashi, *Nano Lett.* **9**, 2891 (2009).
- <sup>20</sup>X. L. Liu, D. Hug, and L. M. K. Vandersypen, *Nano Lett.* **10**, 1623 (2010).
- <sup>21</sup>M. Evaldsson, I. V. Zozoulenko, H. Xu, and T. Heinzel, *Phys. Rev. B* **78**, 161407(R) (2008).
- <sup>22</sup>D. Querlioz, Y. Apertet, A. Valentin, K. Huet, A. Bournel, S. Galdin-Retailleau, and P. Dollfus, *Appl. Phys. Lett.* **92**, 042108 (2008).
- <sup>23</sup>D. Gunlycke, D. A. Areshkin, and C. T. White, *Appl. Phys. Lett.* **90**, 142104 (2007).
- <sup>24</sup>A. Lherbier, B. Biel, Y.-M. Niquet, and S. Roche, *Phys. Rev. Lett.* **100**, 036803 (2008).
- <sup>25</sup>E. R. Mucciolo, A. H. Castro Neto, and C. H. Lewenkopf, *Phys. Rev. B* **79**, 075407 (2009).
- <sup>26</sup>I. Martin and Y. M. Blanter, *Phys. Rev. B* **79**, 235132 (2009).
- <sup>27</sup>S. Adam, S. Cho, M. S. Fuhrer, and S. Das Sarma, *Phys. Rev. Lett.* **101**, 046404 (2008).
- <sup>28</sup>F. Sols, F. Guinea, and A. H. Castro Neto, *Phys. Rev. Lett.* **99**, 166803 (2007).
- <sup>29</sup>D. Graf, F. Molitor, K. Ensslin, C. Stampfer, A. Jungen, C. Hierold, and L. Wirtz, *Nano Lett.* **7**, 238 (2007).
- <sup>30</sup>All samples are slightly p-doped most likely due to residual resist on the graphene structures.
- <sup>31</sup>Since the extracted gap  $E_g$  is strongly disorder dependent the error bars are around 5%–10%, in agreement with earlier studies. For clarity, error bars are not shown explicitly in Fig. 3.