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## Segmental motions, not a two-state concerted switch, underlie allostery in CheY

Leanna R. McDonald<sup>a</sup>, Joshua A. Boyer<sup>a</sup>, and Andrew L. Lee<sup>a,b</sup>

<sup>a</sup>Department of Biochemistry and Biophysics, School of Medicine, University of North Carolina at Chapel Hill, Chapel Hill, North Carolina, 27599, USA

<sup>b</sup>Division of Chemical Biology and Medicinal Chemistry, Eshelman School of Pharmacy, University of North Carolina at Chapel Hill, Chapel Hill, North Carolina, 27599, USA

### Summary

The switch between an inactive and active conformation is an important transition for signaling proteins, yet the mechanisms underlying such switches are not clearly understood. *Escherichia coli* CheY, a response regulator protein from the two-component signal transduction system that regulates bacterial chemotaxis, is an ideal protein for the study of allosteric mechanisms. By utilizing <sup>15</sup>N CPMG relaxation dispersion experiments, we monitored the inherent dynamic switching of unphosphorylated CheY. We show that CheY does *not* undergo a two-state concerted switch between the inactive and active conformations. Interestingly, partial saturation of Mg<sup>2+</sup> enhances the intrinsic allosteric motions. Taken together with chemical shift perturbations, these data indicate that the  $\mu$ s-ms timescale motions underlying CheY allostery are segmental in nature. We propose an expanded allosteric network of residues, including W58, that undergo asynchronous, local switching between inactive and active-like conformations as the primary basis for the allosteric mechanism.

### Introduction

Allosteric conformational change is critical for the function of many proteins. At the present time, it is not generally understood how allosteric conformational changes are executed or how many different execution strategies exist. As a simpler alternative to defining precise conformational change trajectories, there has been intense focus directed at the “selected fit” versus “induced fit” paradigms (Okazaki and Takada, 2008; Hammes et al., 2009; Wlodarski and Zagrovic, 2009; Zhou, 2010; Csermely et al., 2010), as well as at concerted vs. propagated conformational changes (Cui and Karplus, 2008; Tsai et al., 2009). The pairing of selected fit with a concerted conformational change – that is, a simple “switch” – has been particularly popular. It is widely recognized that dynamics are central to these processes, yet there are few experiments that have been carried out to directly assess some of the basic assumptions of allostery. Hence there is a need for experimental dynamics data on the timescale of conformational equilibria in allosteric proteins.

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Correspondence should be addressed to Andrew Lee: University of North Carolina, Eshelman School of Pharmacy, Beard Hall, CB 7568, Chapel Hill, NC, 27599. drewlee@unc.edu. Phone: 919-966-7821. Fax: 919-966-0204.

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Receiver domains from response regulator (RR) proteins have been studied extensively and, because of their small size, have become favored as models for understanding conformational allostery. RRs, along with sensor kinases, comprise the two-component system ubiquitous in prokaryotes. RRs usually consist of an input receiver domain that is activated by phosphorylation, and an output domain that transmits the signal into various activities such as DNA binding (Bourret, 2010; Galperin, 2010). Accordingly, the ability of the receiver domain to undergo a well-defined conformational change is a vital component of RR function (Lee et al., 2001; Gao et al., 2007). The chemotaxis protein Y (CheY) from *Escherichia coli* is a RR that regulates chemotactic flagellar rotation (Clegg and Koshland, 1984; Matsumura et al., 1984). Because it lacks an output domain, the CheY receiver domain must both accept a phosphoryl group and directly activate its downstream effector; in response to phosphorylation at D57 (Sanders et al., 1989), which requires the presence of  $Mg^{2+}$  (Lukat et al., 1990), CheY undergoes a conformational change that enables it to directly bind the flagellar motor switch protein FliM at a surface distal to D57. CheY binding to FliM promotes a change in the flagellar rotation from counterclockwise to clockwise (Baker et al., 2006). Unphosphorylated CheY also interacts with FliM, though with considerably reduced affinity (Barak and Eisenbach, 1992; Welch et al., 1993).

As is common in allosteric signaling molecules, CheY samples an active conformation (“A”, FliM binding competent) and an inactive conformation (“I”, FliM binding incompetent). In the I conformation, FliM binding is sterically hindered by the location of Y106 in a solvent exposed or “out” position (Figure 1A). Upon activation by phosphorylation, Y106 (~11 Å away from the phosphoaspartate) rotates to the “in”, buried position, relieving the hindrance. T87 hydrogen bonds with the phosphoryl group and has been shown to be an important link for Y106 rotation (Appleby and Bourret, 1998; Lee et al., 2001a). Motion of the  $\beta 4$ - $\alpha 4$  loop, consisting of residues 88–91, facilitates burial of Y106 (Ma and Cui, 2007; Simonovic and Volz, 2001), leading to the possibility that residues in the loop (in addition to T87) may be involved in the signaling. Together with small changes in the  $\beta 5$ - $\alpha 5$  loop and side-chain motions of K109 and F14, the I and A conformations differ mainly in the Y106 rotation and the location of the  $\beta 4$ - $\alpha 4$  loop (Lee et al., 2001a).

The prevailing conceptual framework for receiver domain activation/allostery has been that the protein exists in a dynamic equilibrium that accesses both the I and A conformations, in line with the Monod-Wyman-Changeux (MWC) model of allostery. The following findings provide evidence of an I-to-A state conformational equilibrium: (1) in the absence of phosphorylation, CheY still has the ability to stimulate clockwise flagellar rotation (Barak and Eisenbach, 1992); (2) in the crystal structure of unphosphorylated,  $Mg^{2+}$ -free CheY, both “out” and “in” conformations of the Y106 side chain were observed (Volz and Matsumura, 1991); (3) binding of FliM, CheA, and CheZ peptides to CheY affect its ability to phosphorylate and can be explained by a ligand-induced shift of the I-to-A equilibrium (Schuster et al., 2001); (4) NMR studies of the receiver domains of Spo0F and NtrC showed that in regions where structural changes upon phosphorylation were observed, there was enhanced transverse relaxation due to conformational exchange on the  $\mu s$ -ms time scale (Volkman et al., 2001; Feher and Cavanagh, 1999); and (5) mutant NtrC proteins revealed a correlation between activity/inactivity and a two-state equilibrium between active and inactive conformations (Volkman et al., 2001; Gardino et al., 2009). While these data are consistent with a two-state switch of RRs, direct detection of a two-state dynamic process between active and inactive conformations has been elusive.

Complicating the simple idea of a pre-existing equilibrium between I and A conformations are several crystal structures of CheY that show intermediate conformations (Simonovic and Volz, 2001; Dyer and Dahlquist, 2006; Guhaniyogi et al., 2006). Additionally, molecular dynamics simulations indicated that Y106 rotation and the formation of a hydrogen bond

between T87 and the phosphoryl group are independent of one another (Ma and Cui, 2007). These studies showed that CheY is not restricted to the two end states and taken together suggest that it can be trapped in metastable states that presumably are sampled along the allosteric conformational change trajectory. It is unknown, however, whether such trapped states are functionally relevant or are merely artifacts from crystallization; it also remains unknown what the relevant timescales are for conformational switching in CheY.

In order to further understand the allosteric switch CheY undergoes upon phosphorylation, we investigated the conformational equilibrium that occurs in the unphosphorylated protein. NMR relaxation dispersion was used to measure the dynamics of CheY switching to test for consistency with a two-state model. We found that a physiological level of  $Mg^{2+}$  likely plays a critical role in promoting allosteric conformational changes. Nevertheless, whether  $Mg^{2+}$  is present or absent, unphosphorylated CheY appears *not* to undergo two-state concerted switching between I and A conformations. Rather, the data are more suggestive of a model in which CheY switches in a non-concerted, segmental fashion. Local sites may occupy their active conformations at different times utilizing a previously undescribed signaling network consisting of A88 with T87 and the quartet of W58, M85, E89, and Y106.

## Results

### Microsecond-millisecond motions are along the allosteric path

Large scale conformational changes in proteins frequently occur on the slow, or  $\mu$ s-ms, timescale (Henzler-Wildman and Kern, 2007; Kleckner and Foster, 2011). In an attempt to measure the dynamic switching between the I and A conformations of unphosphorylated CheY, we used  $^{15}N$  Carr-Purcell-Meiboom-Gill (CPMG) relaxation dispersion experiments to measure motions on this timescale (Palmer et al., 2001). These experiments elucidate processes on the  $\mu$ s-ms timescale that contribute to broadening (or “width”) of NMR peaks. Specifically, the NMR line width is proportional to the transverse relaxation rate,  $R_2$ , which is comprised of an “intrinsic” rate,  $R_2^\circ$ , and a rate that arises from conformational exchange processes on the  $\mu$ s-ms timescale,  $R_{ex}$ :

$$R_2 = R_2^\circ + R_{ex}. \quad (1)$$

The CPMG relaxation dispersion experiment measures the suppression of  $R_{ex}$  contributions to  $R_2$  as a function of spacing between  $180^\circ$  pulses in the CPMG train,  $\tau_{cp}$  (Loria et al., 1999). For a two-state exchange process in the limit of fast exchange on the NMR timescale,

$$R_{2,eff} \left( \frac{1}{\tau_{cp}} \right) = R_2^\circ + \left( \frac{p_I p_A \Delta\omega^2}{k_{ex}} \right) \left[ 1 - \frac{2 \tanh \left( \frac{k_{ex} \tau_{cp}}{2} \right)}{k_{ex} \tau_{cp}} \right], \quad (2)$$

where  $p_I$  and  $p_A$  are populations of the major and minor states,  $\Delta\omega$  is the difference in chemical shift between the two states, and  $k_{ex}$  is the rate of exchange between I and A ( $= k_1 + k_{-1}$ ). A longer expression (the Carver-Richards equation) exists for the general case (Palmer et al., 2001). A particular residue’s  $R_{ex}$  value is considered to be non-zero (i.e. affected by conformational exchange) if  $R_{ex} > 2 \text{ s}^{-1}$ , as described in the Experimental Procedures.

Pilot relaxation dispersion experiments indicated that the dispersion curves were more pronounced at lower temperatures, and  $15^\circ \text{C}$  was determined to be a good compromise between pronounced dispersion curves and signal-to-noise. Under the standard conditions of 10 mM  $Mg^{2+}$ ,  $^{15}N$  CPMG relaxation dispersion data were collected at 500, 600, and 700

MHz for CheY. Residues with non-zero  $R_{ex}$  values localized to the active site and FlIM binding interface, which includes residues T87, A88, and E89 (in the  $\beta$ 4- $\alpha$ 4 loop), as well as Y106 and V107 (Figure 1A). This was not surprising since all substantial conformational differences between the I and A states are limited to this region (Lee et al., 2001a). There were no non-zero  $R_{ex}$  values elsewhere in the protein, detected either by CPMG relaxation dispersion or by model-free analysis of  $T_1$ ,  $T_2$ , and  $\{^1H\}$ - $^{15}N$  NOE data. Therefore, at a qualitative level, these data are consistent with motion corresponding to the I-to-A transition.

### Motions measured by relaxation dispersion do not result from a single conformational switch event

The relaxation dispersion curves (Figure 1B) were quantitatively analyzed using the general Carver-Richards equation (Palmer et al., 2001) to fit each residue individually (i.e., “local fits”, see Experimental Procedures). A total of 11 residues could be fit with reasonable precision. If CheY switches concertedly in a simple two-state manner between I and A conformations, we would expect to measure the same exchange rate and populations for all residues. However, the fits yielded a range of kinetic and thermodynamic parameters (Table 1). Exchange rates varied from slow ( $\sim 1300\text{ s}^{-1}$  for residues 12 and 38) to fast ( $\sim 3000\text{ s}^{-1}$  for residues 106 and 107) and populations varied from 90% – 99% of I. The different exchange rates were also evident in the raw dispersion curves (compare rows of Figure 1B). In addition, two-state or three-state concerted switching could not be rationalized based on attempts to group fit the data (see Experimental Procedures for details).

Yet another test for two-state behavior is to compare the dispersion-based differences in chemical shifts to known differences in chemical shifts for two defined structural states. Accordingly, we compared  $\Delta\omega$  from the local fits of relaxation dispersions to  $\Delta\delta$  from the chemical shift perturbations between unphosphorylated and  $\text{BeF}_x$ -bound (phosphoryl mimic (Yan et al., 1999)) CheY. A poor correlation between  $\Delta\omega$  and  $\Delta\delta$  was obtained (Figure 1C), further indicating that the I-to-A transition cannot be described by a simple two-state transition.

Taken together, the non-uniform values of the individual exchange parameters, the inability to group fit, and poor correlation of  $\Delta\omega$  and  $\Delta\delta$  indicates that CheY does not undergo concerted, two-state switching in the presence of 10 mM  $\text{Mg}^{2+}$  at 15 °C.

### Removal of $\text{Mg}^{2+}$ alters dynamics of unphosphorylated CheY

Another possible source of non-zero  $R_{ex}$  values in CheY and other RRs is the reversible binding of  $\text{Mg}^{2+}$ . A divalent metal ion is necessary for all RR phosphorylation and dephosphorylation (Lukat et al., 1990). In the experiments described in the previous section, near physiological concentrations of  $\text{Mg}^{2+}$  (10 mM) were used to characterize CheY in an environment similar to inside of cells and to have consistency with previous biochemical and NMR work on CheY (Moy et al., 1994; Hubbard et al., 2003). Because the binding affinity of  $\text{Mg}^{2+}$  to CheY is  $1.5 \pm 0.3\text{ mM}$  under our conditions (determined by NMR  $^1H$ - $^{15}N$  HSQC peak shifts, Figure S1), at 10 mM  $\text{Mg}^{2+}$  and 1 mM CheY, 86% of CheY is bound by the ion. Thus, if the  $\text{Mg}^{2+}$  binding kinetics are on the appropriate timescale, the binding of  $\text{Mg}^{2+}$  and associated side-chain rearrangements could be detected by relaxation dispersion and this could complicate the interpretation in terms of conformational exchange.

To separate motions associated with  $\text{Mg}^{2+}$  binding and release from those intrinsic to CheY, we carried out relaxation dispersion experiments without  $\text{Mg}^{2+}$  present (i.e., with 1 mM ethylenediaminetetraacetic acid, EDTA). In the absence of  $\text{Mg}^{2+}$ , 13 residues displayed non-

zero  $R_{\text{ex}}$  values for CheY, compared to 11 when  $\text{Mg}^{2+}$  is present at 10 mM. These residues are localized to the FliM binding interface or the active site (Figure 2A) just as when  $\text{Mg}^{2+}$  is present (Figure 1A). The locally fit parameters, when assuming a two-state mechanism, are diverse (Figure 2B and Table 2). In comparison with dynamics in the presence of  $\text{Mg}^{2+}$ , many of the same residues display significant dispersion, including T87 and V107. The locally fit parameters of the common residues differ significantly between the  $\text{Mg}^{2+}$ -free and  $\text{Mg}^{2+}$ -present conditions (Tables 1 and 2). Additionally,  $\Delta\omega$  and  $\Delta\delta$  do not correlate (Figure 2C), as was the case with  $\text{Mg}^{2+}$  (Figure 1C). Therefore, even without the potentially complicating effects of  $\text{Mg}^{2+}$ , CheY appears not to undergo two-state switching.

It is possible that bound  $\text{Mg}^{2+}$  induces a conformational change in CheY. The crystal structure of  $\text{Mg}^{2+}$ -bound *E. coli* CheY reveals a conformational difference from  $\text{Mg}^{2+}$ -free CheY that localizes to  $\alpha 4$  and the  $\beta 4$ - $\alpha 4$  loop (Bellolell et al., 1994). However, chemical shift perturbations of CheY upon the addition of  $\text{Mg}^{2+}$  are strongly correlated with closeness to the ion and in general are small in regions of allosteric conformational change (Figure S2), suggesting no significant structural rearrangements. In addition, the crystal structure of *Salmonella typhimurium* CheY (Stock et al., 1993) (which differs by 3 amino acids from *E. coli*) and the NMR structure of *E. coli* CheY (Moy et al., 1994), both with  $\text{Mg}^{2+}$  present, have no indication of any large structural rearrangement. For these reasons, the different CheY dynamics observed with or without  $\text{Mg}^{2+}$  do not appear to be the direct result of  $\text{Mg}^{2+}$ -induced conformational change.

### Physiological $\text{Mg}^{2+}$ concentrations enhance allosteric dynamics

To gain greater insight into the effect of  $\text{Mg}^{2+}$  on CheY dynamics, we obtained estimates of  $R_{\text{ex}}$  at additional  $\text{Mg}^{2+}$  concentrations of 1 mM and 75 mM (Figure 3). Thus, assuming a single binding site (see Supplemental Discussion), 1 mM CheY was calculated (based on the  $K_d$  of 1.5 mM) to be bound by  $\text{Mg}^{2+}$  at a level of 0%, 31%, 86%, and 98% for 0 mM, 1 mM, 10 mM, and 75 mM concentrations of  $\text{Mg}^{2+}$ , respectively. For all concentrations of  $\text{Mg}^{2+}$ , estimates of  $R_{\text{ex}}$  were obtained as the difference of  $R_{2,\text{eff}}$  values at the lowest and highest values of  $\tau_{\text{cp}}$  used for dispersion measurements. Overall, the dependence of  $R_{\text{ex}}$  on the concentration of  $\text{Mg}^{2+}$  is immediately apparent:  $R_{\text{ex}}$  in the presence of no  $\text{Mg}^{2+}$  or very high concentration of  $\text{Mg}^{2+}$  (75 mM) is relatively low, and  $R_{\text{ex}}$  at intermediate levels of  $\text{Mg}^{2+}$  (1 mM) is very high (Figure 3). Furthermore, peak broadening was evident in many residues at intermediate  $\text{Mg}^{2+}$  concentrations, implying increased motion. While high  $R_{\text{ex}}$  levels at intermediate concentrations of  $\text{Mg}^{2+}$  are expected for residues in close proximity to the bound ion (assuming appropriate line-broadening kinetics), high  $R_{\text{ex}}$  levels are not expected at residues distal to  $\text{Mg}^{2+}$  that are not structurally affected by the ion, such as on the FliM binding interface. Yet,  $R_{\text{ex}}$  at V107 is quite pronounced ( $28 \text{ s}^{-1}$ ) at 1 mM  $\text{Mg}^{2+}$ . We note that  $\text{Mg}^{2+}$  binding induces a small chemical shift change of 0.52 ppm for the amide nitrogen of V107 (Figure S2), especially compared to the large chemical shift change of 8.6 ppm brought about by binding of  $\text{BeF}_x$ . Furthermore, A88 experiences extreme line broadening at intermediate concentrations of  $\text{Mg}^{2+}$ . The dispersion experienced at this position is larger than expected when compared to neighboring residues. Thus, there is something particular about partial  $\text{Mg}^{2+}$  saturation that increases conformational dynamics at A88 and even further at the FliM binding interface more than  $10 \text{ \AA}$  away. The difference in  $R_{\text{ex}}$  far from  $\text{Mg}^{2+}$  is due to a change in the I-to-A transition (see Supplemental Discussion and Figures S3 and S4). Importantly, this occurs without significant changes to the structure of the ground state (see above), even though phosphorylation – which occurs at nearly the same location as  $\text{Mg}^{2+}$  binding – induces a large structural change. The  $\text{Mg}^{2+}$  concentration's effect on the I-to-A transition is likely biologically relevant since  $R_{\text{ex}}$  is largest at intermediate concentrations of  $\text{Mg}^{2+}$  (1 mM); this concentration is close to the  $K_d$  for  $\text{Mg}^{2+}$

and CheY (1.5 mM), and also close to the concentration of  $Mg^{2+}$  in *E. coli* (~1–2 mM  $Mg^{2+}$  (Alatossava et al., 1985)).

### Key residues display linear chemical shifts in response to $Mg^{2+}$ binding

We observed that, for Y106, V107, and E89, the NH chemical shifts move along a line for various states of CheY, which include apo,  $BeF_x$ -bound, inactive mutant T87I CheY (Ganguli et al., 1995), and activated mutant A113P CheY (Smith et al., 2004) (Figures 4A, 4B, and 4C). Linear chemical shifts are indicative of a classic, two-state, I-to-A switch, as was described for the receiver domain of the NtrC RR (Volkman et al., 2001). To test whether linear shifts were observed throughout CheY, we plotted all  $Mg^{2+}$ -induced CSPs on a normalized vector scale (Figure 4D). If  $Mg^{2+}$  shifts this equilibrium, all chemical shifts should fall along the dotted line from inactive mutant T87I to  $BeF_x$ -bound wild-type CheY. However, with a few exceptions, the CSPs distribute essentially randomly, with no apparent preferred “direction”. Thus, it appears that, overall,  $Mg^{2+}$  binding elicits either a new conformational state(s) that has only subtle differences from the I state and bears no resemblance to A-like chemical shifts, or  $Mg^{2+}$  has essentially no significant effect on CheY conformation. This is further supported by no significant change in the affinity of binding of FliM to unphosphorylated CheY with the addition of  $Mg^{2+}$  (Figure S5). Nevertheless, there are a few key residues (e.g. Y106, E89) that display the linear chemical shifts suggestive of an equilibrium shift (Figure 4). Interestingly, additional residues that fall near the line (Figure 4D, M85, A97 and A98) lie in the pocket surrounding W58. These residues, along with Y106 and E89, may experience a shift in conformation due to the presence of  $Mg^{2+}$  (discussed below).

## Discussion

For half a century now, allosteric conformational change has been revealed to be an extremely successful mechanism for regulating protein function. Yet, key specific questions regarding allostery have yet to be answered. Is the conformational change concerted? Are there structural intermediates? What are those intermediates? What is the dynamic process that allows conversion between the end conformations and how is it triggered? In summary, we still have a poor understanding of how allosteric conformational changes actually happen. By conducting in-depth NMR experiments on the small allosteric protein CheY, we hoped to gain new insights into the mechanism(s) of allostery. CheY is an excellent system for probing allostery because numerous crystal structures have been solved in a variety of liganded states and it is highly amenable to NMR relaxation experiments. It also exhibits a rich complement of residues with  $\mu$ s-ms motion at residues that connect the physically separated upstream and downstream effector sites.

### Allostery in CheY does not operate by a simple shift of a two-state equilibrium

From previous crystallographic, NMR, and functional data (Birck et al., 1999; Feher and Cavanagh, 1999; Volkman et al., 2001; Gardino et al., 2009), the evidence suggests that receiver domains primarily adopt two conformations and hence allostery is achieved by shifting a pre-existing dynamic equilibrium between these two states, consistent with the MWC model of allostery. Based on this previous work on NtrC, FixJ, and Spo0F, and our initial observation here of  $R_{ex}$  in the allosteric network of residues, the simplest expectation would be that CheY undergoes a concerted switch between I (inactive) and A (active) conformations. Here, we show using NMR relaxation dispersion that intrinsic allostery in CheY does *not* appear to operate by a simple shift of a two-state equilibrium. This conclusion rests, in part, on the ability of relaxation dispersion experiments to accurately identify conformational exchange rates and populations.

It is possible that unphosphorylated CheY does not completely switch to the A state but rather samples a conformational intermediate between the I and A states, and this may account for the lack of correlation in  $\Delta\delta$  and  $\Delta\omega$  values. However, this cannot be confirmed experimentally without having accurate chemical shifts for the intermediate state, which has yet to be isolated in solution. Alternatively, if a structural model for the intermediate exists, the chemical shifts can be calculated using SHIFTX (Neal et al., 2003); then  $\Delta\delta$  can be calculated between the inactive and intermediate states and compared with  $\Delta\omega$  from relaxation dispersion fits. This was carried out using the partially switched structure of CheY bound to FliM without phosphorylation (2B1J (Dyer and Dahlquist, 2006)) as the intermediate, but the  $\Delta\omega$  values did not correlate. In any case, if the observed  $R_{ex}$  values result only from unphosphorylated CheY switching concertedly to an intermediate state, one would still expect  $k_{ex}$  and populations to be uniform, yet they are not.

### Distal quartet of residues directs allosteric conformational change

In order for CheY to undergo its allosteric transition it must primarily do two things: (1) reposition the  $\beta 4$ - $\alpha 4$  loop, and (2) swing Y106 into the pocket just under the  $\beta 4$ - $\alpha 4$  loop. We measured three residues, Y106, V107, and E89, that appeared to approach two-state behavior, although this is apparent only in the chemical shift data (Figure 4) and not in the relaxation dispersion data (Figures 1 and 2). One possible explanation for the linear chemical shifts is that because Y106 is at the solvent interface and presumably under less structural constraint, its position is to a reasonable approximation either “out” or “in”, yielding a bimodal distribution of possible chemical shift values for the amide of Y106. The strong ring currents dominate the chemical shifts of both Y106 and V107 amides. The amide chemical shift of E89, located in the  $\beta 4$ - $\alpha 4$  loop, likely monitors the remodeling of this loop. We utilized SHIFTX (Neal et al., 2003) to calculate the effect on chemical shift from a conformational change. Y106 and V107 amide chemical shifts are very sensitive to the rotation of Y106, with minimal effects from the location of the  $\beta 4$ - $\alpha 4$  loop. By contrast, the amide shift of E89 is very sensitive to movement of the  $\beta 4$ - $\alpha 4$  loop, but not to Y106. Therefore, E89 monitors the  $\beta 4$ - $\alpha 4$  loop motion, and Y106 and V107 likely monitor the rotation of Y106. Taken together, the chemical shift data reveal a localized shift to active-like states.

While E89, Y106, and V107 have the most linear chemical shifts, a few residues are nearly linear (M85, A97, and A98) while all others show no semblance of linearity. M85, A97, and A98 have not been discussed previously as part of the allosteric mechanism. Yet, these residues line the pocket occupied by W58 and are proximal to the  $\beta 4$ - $\alpha 4$  loop and Y106. Examination of CheY in its inactive (3CHY (Volz and Matsumura, 1991)), active (1FQW (Lee et al., 2001a)), and intermediate conformations (2B1J (Dyer and Dahlquist, 2006)) reveals an extended set of conformational interactions. Upon activation, the hydroxyl of Y106 hydrogen bonds to the backbone carbonyl of E89, which is allowed by movement of the  $\beta 4$ - $\alpha 4$  loop. What appears to further stabilize this interaction is a large movement of the E89 side chain towards the side chain of W58, resulting in hydrophobic contacts between the methylene groups of E89 and the tryptophan indole. Interestingly, superimposition of inactive and active CheY structures shows that the tryptophan indole group reorients by  $\sim 55^\circ$  and the intermediate state (when CheY is bound only to FliM peptide) shows the indole at an intermediate angle. Although tryptophan at position 58 is not conserved amongst RRs, the receiver domain from PhoB displays a similar rotation of  $\sim 60^\circ$  upon binding  $\text{BeF}_x$  (Sola et al., 1999; Bachhawat et al., 2005). When analyzing a range of crystal structures, we find that the angle of the W58 indole correlates with the pseudodihedral angle 87:88:89:90 (Figure 5A). This pseudodihedral angle has been used to indicate the location of the  $\beta 4$ - $\alpha 4$  loop (Dyer and Dahlquist, 2006). In order to contact E89, W58 increases its distance from M85, which is associated with W58 in the inactive conformation. Thus, M85

and E89 appear to be in competition for W58, with E89 making contact only in the active conformation. In further support of this extended mechanism, we note that in 3CHY, M85 was found to occupy two conformations (similar to Y106), of which only one is adopted in the intermediate and active conformations and the other is adopted in other crystal structures of the inactive state (Figure 5B). Additional crystal structures reveal that M85 often occupies the “active” rotamer even though W58 has not rotated to the active state (Figure 5B). This indicates that W58 is not as robustly coupled to M85 as the  $\beta 4$ - $\alpha 4$  loop. However, one should be cautioned not to overinterpret the lack of coupling between M85 and W58; since the two M85 side-chain conformations have similar electron density, this leads to the possibility that the crystal structures did not accurately account for multiple conformations. We also note that the amide of W58 could not be assigned, which is consistent with its indole sampling at least two conformations prior to full activation. We propose here that the quartet of W58, M85, E89, and Y106 are coupled, as a key component of the allosteric signaling process from phosphorylation.

### A88 links $Mg^{2+}$ binding to the allosteric quartet

Upon  $Mg^{2+}$  binding to CheY, intrinsic fluorescence of W58 is quenched (Lukat et al., 1990), indicating that  $Mg^{2+}$  may play an additional role in this signaling cascade. T87 is considered to be the initial “sensor” of phosphorylation and triggers the allosteric response (Appleby and Bourret, 1998). However, it is unclear whether additional residues play a crucial role and which residues may sense  $Mg^{2+}$ . A88, in addition to T87, also hydrogen bonds with the phosphoryl group via its backbone amide. Perhaps most importantly, from the relaxation dispersion data in 10 mM  $Mg^{2+}$ , A88 fit to parameters unlike any other residue. While most residues fit to populations ( $p_i$ ) of 0.98–0.99, A88 fit to 0.91. Furthermore, its dispersion curve visually resembles residues 12 and 38, which are near to the bound  $Mg^{2+}$  ion, more than the allosteric residues 87, 89, 106, and 107 (Figure 1B). Careful inspection of the fit to A88 indicates that its  $R_2^\circ$  values are anomalously high and the quality of the fits is significantly lower than the other residues (Table 1). Thus, we suspect that in addition to the motion that gives rise to the main part of A88’s dispersion curve, there is a separate, faster motion on the  $\mu s$  timescale. Only A88 appears to have this more complex combination of motions. Furthermore, A88 experiences line broadening severely affected by the concentration of  $Mg^{2+}$ . At 1 mM  $Mg^{2+}$ , A88 is completely broadened away in the  $^1H$ - $^{15}N$  HSQC. Given its behavior at different  $Mg^{2+}$  concentrations, A88 may be especially sensitive to  $Mg^{2+}$  binding and release.

We therefore hypothesize that A88 acts as a crucial monitor of ion binding and/or phosphorylation at the active site. It is important to note that, even though they are adjacent, residues 86–89 have different exchange parameters. Therefore, at face value these residues appear largely independent. Nevertheless, because these residues lie within the allosteric pathway and all show motion on the  $\mu s$ -ms timescale, we suggest that in addition to T87, A88 provides a linkage to E89, and that the recruitment of W58 by E89 represents a key step in switching the  $\beta 4$ - $\alpha 4$  loop. Given these considerations, we propose that the allosteric network in CheY is facilitated by several distinct behaviors that extend from D57, to T87 and A88, and ultimately to the distal quartet of W58, M85, E89, and Y106 (Figure 6).

### Evidence for segmental allosteric dynamics

Considering the chemical shift and relaxation dispersion data together leads inexorably to a paradox. How can there be linear chemical shifts in key allosteric residues when switching appears not to be concerted by relaxation dispersion? A distribution of 2D chemical shifts along a line is typically considered to indicate a rapid exchange between two states (Volkman et al., 2001), with the spread of positions resulting from a shifting of the equilibrium. While this is indeed highly likely to be the case at the spin probes’ “site”



(amide H and N could be considered as one site), if two probes show this behavior at significant distances from one another, they need not result from the same conformational exchange process. Thus, we consider if the observed linearity can result from *local* two-state switching. Specifically, the switching of W58 can be sensed by residues 85, 97 and 98, but the switching of Y106 appears independent of W58. It is possible that some groups of residues switch concertedly, yet with some independence from other switching events. Even though such a group appears to switch independently, there could be coupling with a time lag or interference effects. This situation could give rise to linear chemical shifts, yet show dispersion curves that cannot be fitted together, as was the case for CheY here. Thus, we propose that CheY's allosteric switching is segmental in nature. It is not truly two-state, but there are several groupings of coupled residues that move with some independence and perhaps in a preferred sequential order. This rough model also suggests that CheY shares features of both conformational selection and induced fit. It does not appear that the A-state is fully formed in the absence of phosphorylation, since there is a poor correlation of  $\Delta\omega$  and  $\Delta\delta$ . Yet, some residues may sample "active-like" *local* conformations. We speculate that binding and release of  $Mg^{2+}$  actuates conformational transitions along the allosteric pathway, from T87 and A88 to the W58-M85-E89-Y106 quartet (Figure 6). This is highlighted by maximal  $R_{ex}$  at Y106 occurring at the intermediate  $Mg^{2+}$  concentration of 1 mM.

The segmental nature of CheY's allosteric dynamics is perhaps best supported by the behavior of Y106. The rate of dynamic switching as observed by CPMG relaxation dispersion is clearly faster than the other switching residues, as evidenced by visual inspection of the dispersion curve (Figure 1B), as well as the fitted  $k_{ex}$  of  $3100\text{ s}^{-1}$  (Table 1). Assuming that this observed motion corresponds to the allosteric conformational change, Y106 must be moving with at least partial independence from the other residues. That Y106 switches rapidly and with relative ease is supported by the significantly populated active-like inward position in 3CHY (Volz and Matsumura, 1991), as well as its low activation barrier for switching, as determined from a computational study (Ma and Cui, 2007). Even from the linear chemical shift data (Figure 4), it is seen that in 10 mM  $Mg^{2+}$ , Y106 is only ~20% toward the active state chemical shift, whereas the other linear shift data appear to be in a range of ~20–60% toward the active state (Figure 4D). Thus, neither relaxation dispersion nor chemical shift data fully support a model for pre-existing, concerted, two-state switching dynamics.

### Comparison to other systems

Most of our knowledge of allosteric dynamics in RRs derives from NMR work on the nitrogen regulatory protein C (NtrC) receiver domain. A combination of chemical shift analysis (particularly of mutants) and transverse relaxation measurements led to the proposal of a two-state dynamic equilibrium between I and A conformations in NtrC (Volkman et al., 2001). The recent CPMG relaxation dispersion study on NtrC reported that those data are consistent with a two-state pre-existing equilibrium and that unphosphorylated NtrC populates the active state at the ~14% level (Gardino et al., 2009); interestingly, it appears that that study was carried out in the absence of  $Mg^{2+}$ . By contrast, because CheY does not appear two-state (even in the absence of  $Mg^{2+}$ ) we cannot currently estimate the population of the active state. Even when considering residues locally, the populations of "excited states" from relaxation dispersion (Table 1) are typically much lower than those deduced from chemical shift plots (Figure 4), assuming that T87I populates only the inactive state. Unfortunately, relaxation dispersion data have not been reported for any other RR to the best of our knowledge. CPMG relaxation dispersion is an excellent approach to studying allosteric mechanisms if the conformational switching occurs on the timescale between a few hundred  $s^{-1}$  to  $\sim 3000\text{ s}^{-1}$ . Application of this method to other allosteric proteins holds

the promise of revealing allosteric mechanisms. This has been done successfully for the KIX domain of the CREB binding protein (Bruschweiler et al., 2009), catabolite activator protein (CAP) (Tzeng and Kalodimos, 2009), and the PBX1 homeodomain (Farber and Mittermaier, 2011), each of which exhibited two-state allosteric switching.

Additionally, a study on the *trp* RNA-binding attenuation protein (TRAP) revealed allosteric motions as a series of localized conformational changes (Kleckner et al., 2012). TRAP's behavior is similar to CheY's segmental dynamics and shed light on the diverse mechanisms proteins may employ to achieve allostery. These results imply that a protein may not solely employ a MWC or Koshland-Nemethy-Filmer (KNF) allosteric mechanism, but may combine them to form a complicated switching event.

## Experimental Procedures

### Protein expression and purification

*E. coli* CheY DNA (provided by Dr. Bob Bourret, University of North Carolina at Chapel Hill) was subcloned into the pET28a plasmid (Novagen). Mutants were made by site-directed mutagenesis PCR. The CheY vector was transformed into BL21 Star (DE3) cells (Invitrogen) and grown in minimal media with the appropriate isotope(s):  $^{15}\text{NH}_4\text{Cl}$  (99%) and/or D-glucose ( $\text{U-}^{13}\text{C}_6$ -99%) as the sole nitrogen and carbon sources, respectively). Samples for relaxation dispersion were grown to yield high (>80%)  $^2\text{H}$  incorporation. Cells were grown at 37 °C until  $A_{600}$  reached 0.6. Isopropyl 1-thio- $\beta$ -D-galactopyranoside was added to a final concentration of 1 mM and cells were grown for an additional 22–26 hours (32–36 hours if  $^2\text{H}_2\text{O}$  used) at 20 °C. The cells were harvested by centrifugation, resuspended in buffer A (25 mM Tris, 10 mM  $\text{MgCl}_2$ , pH 8.0), and stored at –20 °C until needed.

Resuspended cells underwent three freeze/thaw cycles and sonication. The lysate was then centrifuged at 6000 rpm and dialyzed overnight into buffer A at 4 °C. The protein was purified on a Q Sepharose Fast Flow column equilibrated with buffer A. CheY eluted approximately halfway along a gradient of buffer B (buffer A with the addition of 1.5 M NaCl) from 10 to 55% buffer B over 135 mL. The eluent was passed over either a G-50 or G-75 gel-filtration column equilibrated with NMR buffer (50 mM  $\text{NaP}_i$ , 0.02%  $\text{NaN}_3$ , pH 7.0 and an appropriate amount of  $\text{MgCl}_2$  and/or EDTA). Pure CheY was concentrated to between 1.2 and 3 mM and stored at 4 °C. Concentrations were determined using the extinction coefficient of  $10.3\text{ cm}^{-1}/\text{mM}$  (Silversmith et al., 2001) and  $14.3\text{ cm}^{-1}/\text{mM}$  for wild-type and Y106W CheY, respectively.

### NMR Spectroscopy

All NMR spectra were collected on 1 mM CheY with NMR buffer and 10%  $^2\text{H}_2\text{O}$ . The appropriate amount of  $\text{MgCl}_2$  or EDTA was present in the NMR buffer, depending on the experiment. NMR experiments were collected at 15 °C on Varian INOVA spectrometers equipped with room-temperature (500 and 600 MHz) or cryogenic (700 MHz) probes. Assignments for wild-type CheY in the presence of 10 mM  $\text{Mg}^{2+}$  were made from triple resonance experiments, supplemented by previously published assignments (Moy et al., 1994). We made use of both triple resonance experiments and assignments of wild-type CheY to assign mutant and  $\text{BeF}_x$ -bound wild-type CheY (both in the presence of 10 mM  $\text{Mg}^{2+}$ ). Assignments of CheY in the presence of 1 mM EDTA were made based on following peak shifts during titration with  $\text{Mg}^{2+}$ . All NMR data were processed using NMRPipe (Delaglio et al., 1995) and visualized with NMRDraw and NMRView (Johnson, 2004).

## Preparation of BeF<sub>x</sub>-bound CheY

NMR samples of BeF<sub>x</sub>-bound CheY were prepared with high enough concentrations of BeCl<sub>2</sub> and NaF to ensure saturation of CheY and in an ideal ratio for a BeF<sub>3</sub><sup>-</sup> complex. The samples were prepared by first combining all components of a typical unphosphorylated CheY sample (e.g. 1 mM CheY protein, 10 mM MgCl<sub>2</sub>, 10% D<sub>2</sub>O). Then BeCl<sub>2</sub> and lastly NaF were added to a final concentration of 4 mM and 25 mM, respectively. The samples were then left at room temperature overnight and centrifuged briefly before transferring to an NMR tube.

## <sup>15</sup>N CPMG Relaxation Dispersion

Relaxation dispersion experiments were carried out using relaxation-compensated CPMG experiments (Loria et al., 1999). Most experiments were collected with a 40 ms total CPMG period and one with 20 ms. 12 or 13  $\tau_{cp}$  values with 2 duplicates and a reference experiment were collected interleaved.  $R_{2,eff}$  was calculated from peak intensities as previously described (Mulder et al., 2001).

NMR relaxation dispersion experiments measuring motions on the  $\mu$ s-ms timescale were fit assuming a two-state model. Residues were considered to have significant motions and worth analyzing if  $R_{ex} > 2 \text{ s}^{-1}$ . These residues were analyzed as previously described (Mauldin et al., 2009). Briefly, using an F-test ( $\alpha_{critical} = 0.01$ ), the data were fit to both a model that assumes no exchange and a simple two-state model to determine which residues exhibited a significant dependence of  $R_{2,eff}$  on  $1/\tau_{cp}$ . These residues were then fit to the Carver-Richards equation (Palmer et al., 2001). By fitting the dependence of  $R_{2,eff}$  on  $1/\tau_{cp}$ , the exchange rate ( $k_{ex}$ ), populations ( $p_I$ ,  $p_A$ ), and chemical shifts ( $\Delta\omega$ ) were determined. Data from two and three field strengths were fit simultaneously for 1 mM EDTA and 10 mM Mg<sup>2+</sup>, respectively, using the in-house program *exrate2.0*. Errors were estimated by Monte Carlo simulation.

In the limit of fast exchange, eq. (2) can be used to fit dispersion curves (Palmer et al., 2001). Residues were also fit to this equation to assess appropriateness of the Carver-Richards equation. For many of the slower residues, the simplified fast equation did not give the same fit as the Carver-Richards equation and therefore we decided the latter was needed to distinguish the best parameters.

Fitting all of the residues together to the same  $k_{ex}$  and  $p_I$  or “group fitting” was done in the same manner as previously described (Mauldin et al., 2009; Mulder et al., 2001). In order to justify group fitting,  $\chi^2_{group} / \chi^2_{local} < 2$  where  $\chi^2_{local}$  is the  $\chi^2$  when a residue is locally fit and  $\chi^2_{group}$  is the  $\chi^2$  of the same residue when group fit. If necessary, the residue with the highest ratio is removed and the group fitting is attempted again. For CheY, the attempt to group fit resulted in elimination of all residues partly due to the high quality of data.

After all fitting was complete, two residues (G39 in 10 mM Mg<sup>2+</sup> and D13 in 1 mM EDTA) were determined to have errors that were too large to be accurately fit. Visual inspection revealed both of the peaks to be extremely broad in the <sup>1</sup>H-<sup>15</sup>N HSQC. Therefore, these residues undergo motions but are unable to be fit due to high error in the intensity measurements. These residues were removed from any further quantitative interpretation of the data.

Utilizing *cpmg\_fit* (provided by Dmitry Korzhnev and Lewis Kay), curves from each residue were fit to a three-state model; however, the fits were not robust due to the limited data. These preliminary fits yielded individual three-state parameters that were diverse. Nevertheless, using an F-test ( $\alpha_{critical} = 0.05$ ), three residues (A36, T87, and E89) in 10 mM Mg<sup>2+</sup> fit better to a three-state model than a two-state model. In 1mM EDTA, no residues

showed improvement with a three-state fit. Additionally, group fitting to a three-state model did not provide significant statistical improvement. This suggests that CheY is not functioning as a three-state concerted switch.

The sign of  $\Delta\omega$  was determined from comparison of peak positions in the HSQC and HMQC spectra (Skrynnikov et al., 2002).

### Determination of approximate $R_{ex}$ by relaxation dispersion

An estimate of  $R_{ex}$  can be obtained by using a very low and a very high  $1/\tau_{cp}$ , as long as the rate of exchange is less than  $\sim 5000\text{ s}^{-1}$ . We used the same experimental set-up as used for measuring full relaxation dispersion curves (as described above) except only two  $\tau_{cp}$  values, 10.0 ms and 0.556 ms were employed. A 40 ms total relaxation time (T) was used and all  $R_{ex}$  estimates were made at 15 °C.  $R_{ex}$  was calculated from the peak intensities for each plane using

$$R_{ex} = \frac{1}{T} \ln \frac{I_{\tau_{cp}=0.556\text{ms}}}{I_{\tau_{cp}=10\text{ms}}}, \quad (2)$$

in which I is the peak intensity for the given  $\tau_{cp}$  value.

### Supplementary Material

Refer to Web version on PubMed Central for supplementary material.

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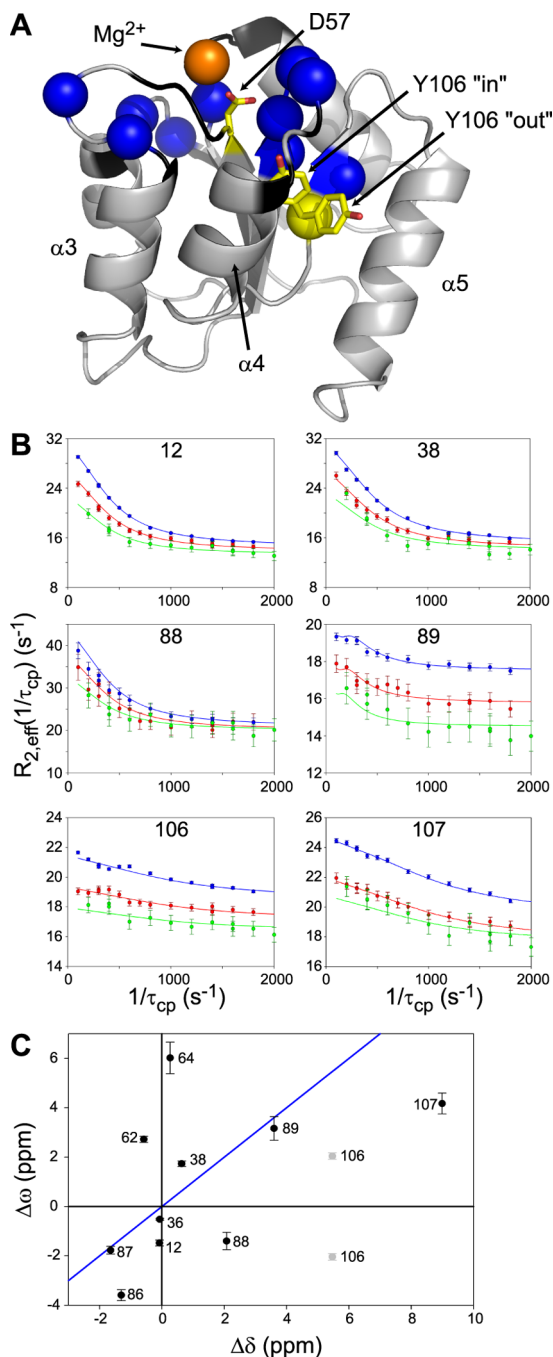
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### Highlights

- NMR dynamics data reveal CheY does not undergo concerted, two-state switching
- Asynchronous, local switching of a network of residues facilitates CheY activation
- $Mg^{2+}$  enhances the dynamics associated with the inactive-to-active transition





**Figure 1.  $^{15}\text{N}$  CPMG Relaxation dispersion data from unphosphorylated CheY in the presence of 10 mM  $\text{Mg}^{2+}$**

(A) Structure of unphosphorylated CheY using 3CHY with the side-chain orientation of D57 and  $\text{Mg}^{2+}$  location from 2CHE. Residues with non-zero  $R_{ex}$  values at 10 mM  $\text{Mg}^{2+}$  are displayed as blue spheres. Black indicates residues that are too broad to be measured. Highlighted in yellow are Y106 and the site of phosphorylation, D57. For Y106, the “in” rotamer indicates active and “out” indicates inactive CheY.

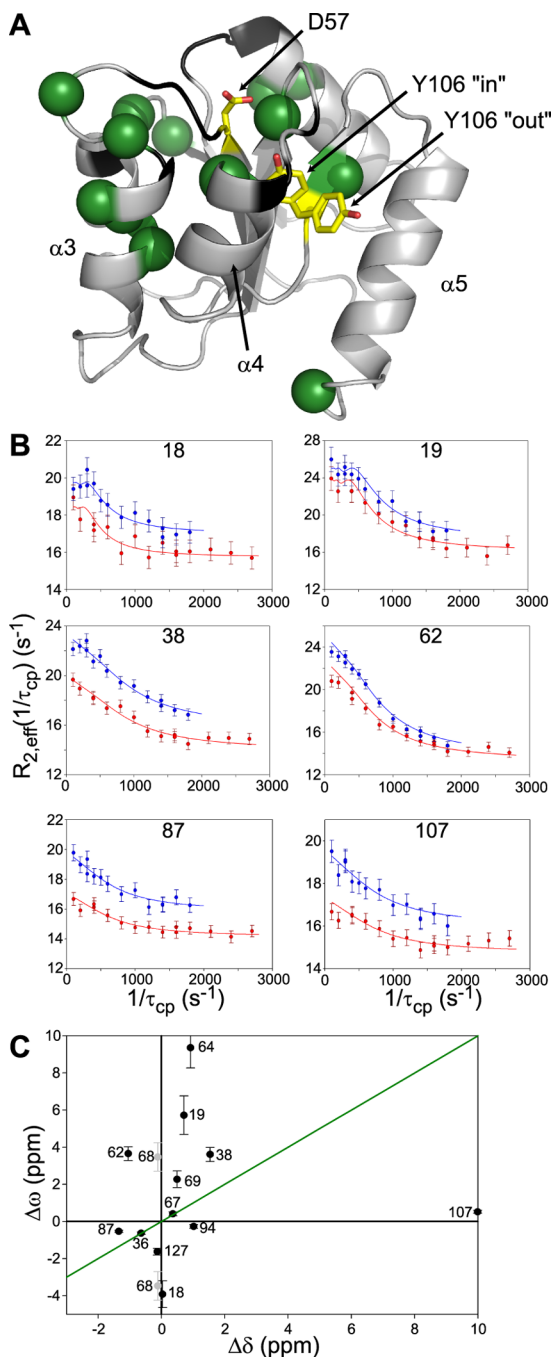
(B) Raw CPMG dispersion curves are displayed for 6 example residues at 700 MHz (blue), 600 MHz (red) and 500 MHz (green). Lines are local fits to the Carver-Richards equation with parameters given in Table 1.

(C)  $\Delta\omega$  values from local fits of relaxation dispersion are plotted against  $\Delta\delta$  using the difference between unphosphorylated CheY and  $\text{BeF}_x$ -bound CheY both with 10 mM  $\text{Mg}^{2+}$ .  $\Delta\delta$  was determined using  $\Delta\delta' = N^{10\text{mM Mg}^{2+}} - N^{\text{BeF}_x}$ , where to get the final  $\Delta\delta$  value we made an additional adjustment. Since unphosphorylated CheY is not completely inactive and  $\text{BeF}_x$ -bound CheY is not completely active, we used  $\Delta\delta = 1.05\Delta\delta'$ . The blue line has a slope of 1. See also Figures S1 and S2.

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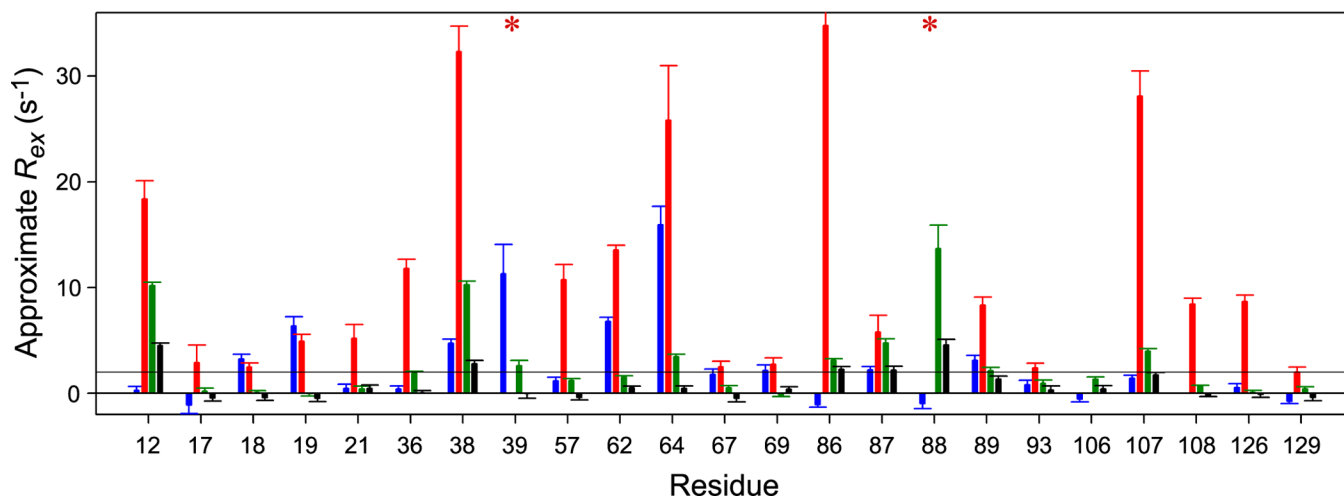
**Figure 2.  $^{15}\text{N}$  CPMG Relaxation dispersion data from unphosphorylated CheY in the absence of  $\text{Mg}^{2+}$  and presence of 1 mM EDTA**

(A) Similar to Figure 1, residues with non-zero  $R_{ex}$  values at 1 mM EDTA are displayed as green spheres.

(B) At 700 MHz (blue) and 600 MHz (red), data are displayed with local fits using the Carver-Richards equation with parameters given in Table 2.

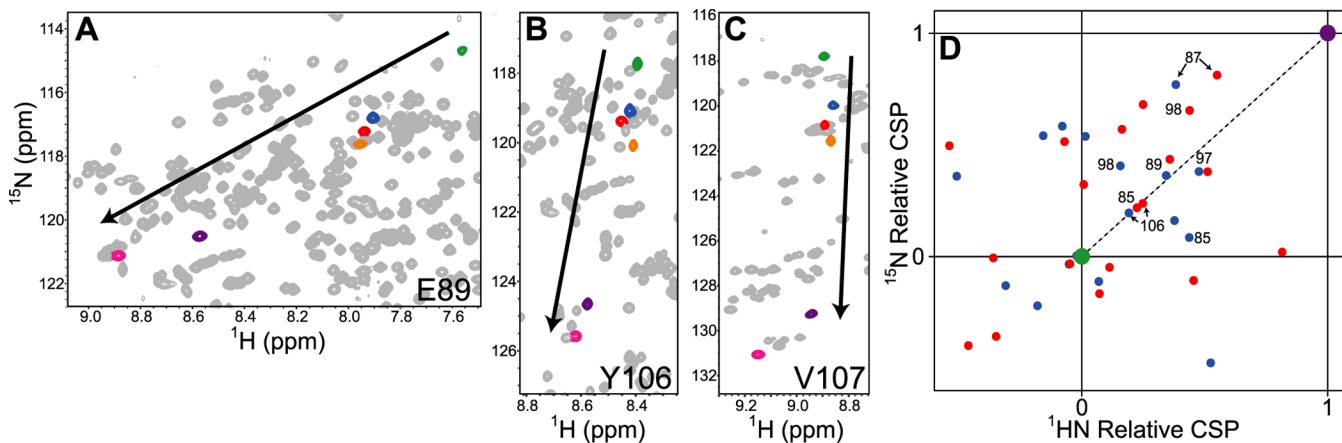
(C) Plot of  $\Delta\omega$  vs.  $\Delta\delta$  using  $\Delta\delta$  between CheY with 1 mM EDTA and CheY with saturating  $\text{BeF}_x$  and 10 mM  $\text{Mg}^{2+}$ .

See also Figures S1 and S2.



**Figure 3. Effect of Mg<sup>2+</sup> concentration on  $R_{ex}$**

$R_{ex}$  is approximated by the change in  $R_{2,eff}$  of two CPMG relaxation dispersion planes. Approximate  $R_{ex}$  is shown for CheY in the presence of 1 mM EDTA (blue), 1 mM Mg<sup>2+</sup> (red), 10 mM Mg<sup>2+</sup> (green), and 75 mM Mg<sup>2+</sup> (black) for all residues with non-zero  $R_{ex}$  values at 1 mM Mg<sup>2+</sup>. Residues G39 and A88 are broadened away at 1 mM Mg<sup>2+</sup> presumably from exchange, indicated by \*. Also, there are no data for Y106 at 1 mM Mg<sup>2+</sup> and V108 at 1 mM EDTA because of peak overlap. See also Figures S1, S2, S3 and S4.



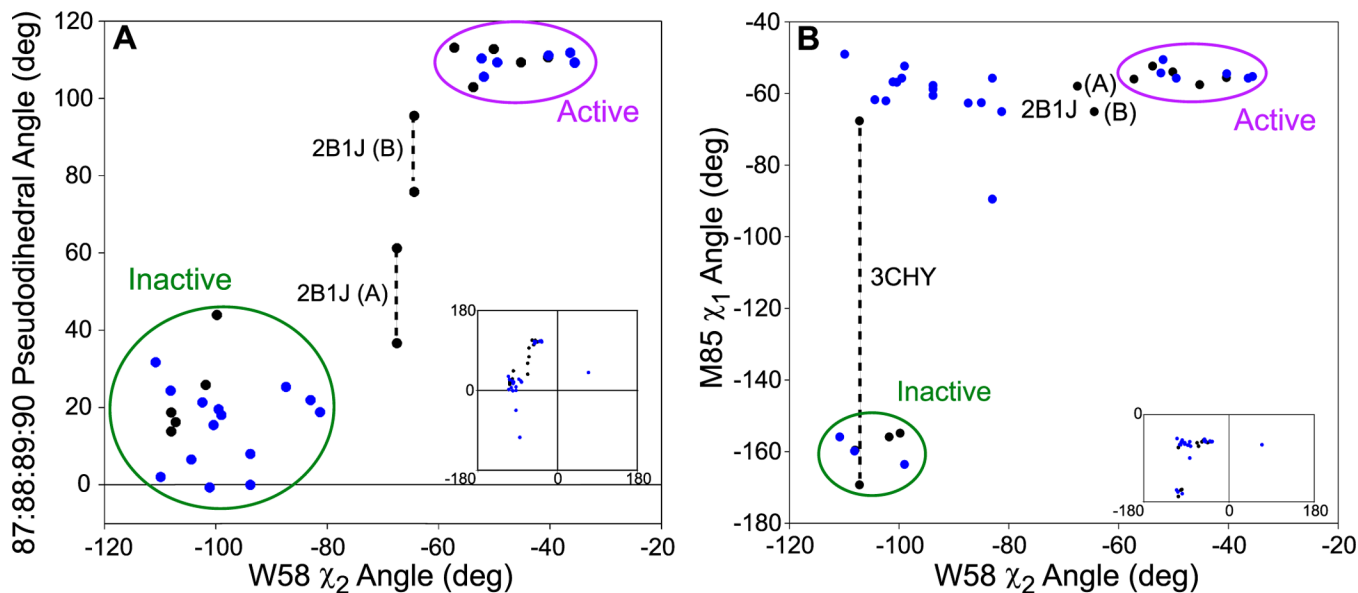
**Figure 4. Chemical shift perturbations (CSP) from inactive to BeF<sub>x</sub> activated CheY**

Overlaid <sup>1</sup>H-<sup>15</sup>N HSQCs are shown with peaks highlighted for (A) E89, (B) Y106, and (C) V107 from various conditions and mutants of CheY: inactivating mutant T87I with 10 mM Mg<sup>2+</sup> (green), wild-type with 1 mM EDTA (blue), wild-type with 10 mM Mg<sup>2+</sup> (red), activating mutant A113P with 10 mM Mg<sup>2+</sup> (orange), BeF<sub>x</sub>-bound wild-type with 10 mM Mg<sup>2+</sup> (purple) and BeF<sub>x</sub>-bound A113P with 10 mM Mg<sup>2+</sup> (pink).

(D) To display all residues, CSP is normalized according to inactivating mutant T87I with 10mM Mg<sup>2+</sup> (green) and BeF<sub>x</sub>-bound wild-type CheY (purple). Each circle represents the normalized CSP of each residue of wild-type CheY in the presence of 10 mM Mg<sup>2+</sup> (red)

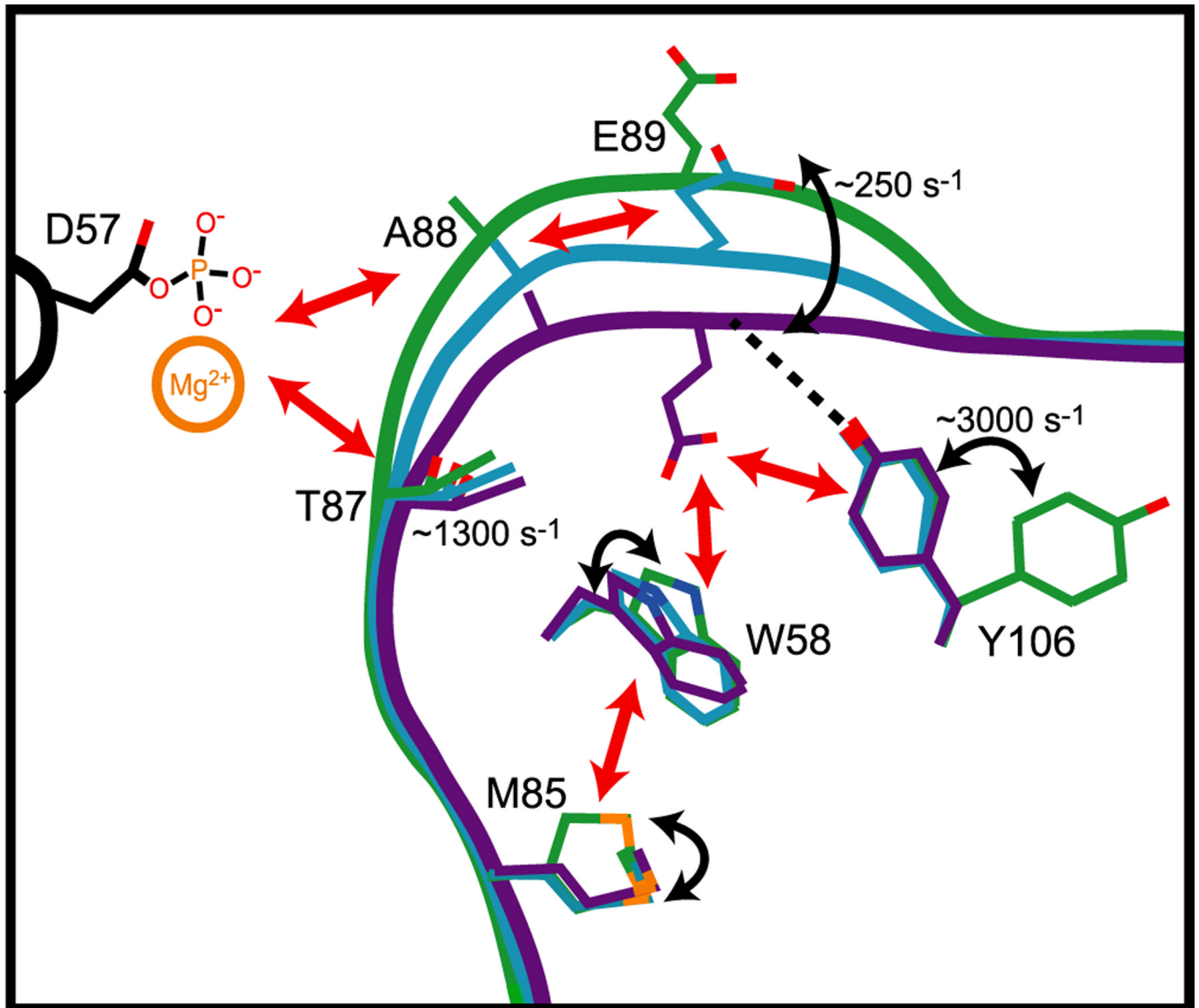
and 1mM EDTA (blue). Residues are only shown if  $CSP = \sqrt{\Delta HN^2 + \left(\frac{\Delta N}{5}\right)^2}$  was greater than 0.20 where  $\Delta HN$  and  $\Delta N$  are between T87I and wild-type BeF<sub>x</sub>-bound CheY. If  $\Delta HN < 0.06$  ppm or  $\Delta N < 0.25$  ppm, the chemical shifts are considered within error in that dimension and the residues were removed. Therefore, while V107 appears linear, it is not shown in (D) since the proton shift is  $< 0.06$  ppm.

See also Figure S5.



**Figure 5. Comparison of dihedral angles in crystal structures of CheY**

(A) Comparison of the 87:88:89:90 pseudodihedral angle with the rotation of W58 and (B) comparison of M85  $\chi_1$  angle with the rotation of W58. The pseudodihedral angle was measured using  $C_\alpha$  coordinates. In black are wild-type CheY crystal structures (1FQW, 1F4V, 2B1J, 1JBE, 3CHY, 2CHE, and 1C4W) and in blue are mutant or CheA-bound CheY crystal structures (only single site mutations and  $>2.5$  Å resolution, full list is provided in Table S1). Insets are the full range of angles occupied while large plots are zoomed on the majority of crystal structures. See Table S1.



**Figure 6. Cartoon model of the allosteric signaling mechanism in CheY**  
 Representations of inactive (green), intermediate (blue) and active (purple) CheY with important signaling residues displayed.

Table 1

Local fits of  $^{15}\text{N}$  CPMG relaxation dispersion for CheY in the presence of 10 mM  $\text{Mg}^{2+}$ .

Residue	$k_{\text{ex}}$ ( $\text{s}^{-1}$ )	$\Delta\omega$ (ppm)	Pt	$R_2^\circ$ ( $\text{s}^{-1}$ ) 500 MHz	$R_2^\circ$ ( $\text{s}^{-1}$ ) 600 MHz	$R_2^\circ$ ( $\text{s}^{-1}$ ) 700 MHz	$\chi^2$
12	$1220 \pm 50$	$1.5 \pm 0.1$	$0.94 \pm 0.01$	$13.3 \pm 0.3$	$13.9 \pm 0.2$	$14.6 \pm 0.1$	4.0
36	$1910 \pm 300$	$0.52 \pm 0.04$	$0.91 \pm 0.02$	$13.1 \pm 0.2$	$13.2 \pm 0.1$	$14.1 \pm 0.1$	2.6
38	$1370 \pm 70$	$1.7 \pm 0.1$	$0.95 \pm 0.01$	$14.0 \pm 0.4$	$14.3 \pm 0.2$	$15.1 \pm 0.2$	46
62	$2060 \pm 180$	$2.7 \pm 0.1$	$0.99 \pm 0.01$	$12.0 \pm 0.1$	$12.2 \pm 0.1$	$13.1 \pm 0.1$	0.75
64	$3620 \pm 460$	$6.0 \pm 0.6$	$0.99 \pm 0.01$	$13.7 \pm 0.5$	$13.0 \pm 0.6$	$13.8 \pm 0.7$	5.3
86	$2100 \pm 330$	$3.6 \pm 0.2$	$0.99 \pm 0.01$	$15.1 \pm 0.3$	$15.5 \pm 0.2$	$16.9 \pm 0.2$	13
87	$1370 \pm 180$	$1.8 \pm 0.2$	$0.98 \pm 0.01$	$15.3 \pm 0.4$	$14.8 \pm 0.3$	$16.6 \pm 0.2$	3.6
88	$1230 \pm 280$	$1.4 \pm 0.4$	$0.91 \pm 0.04$	$20.1 \pm 1$	$20.2 \pm 0.9$	$21.0 \pm 0.8$	14
89	$246 \pm 31$	$3.2 \pm 0.5$	$0.99 \pm 0.01$	$14.5 \pm 0.3$	$15.8 \pm 0.2$	$17.5 \pm 0.1$	2.9
106	$3100 \pm 410$	$2.0 \pm 0.1$	$0.99 \pm 0.01$	$16.4 \pm 0.2$	$17.1 \pm 0.2$	$18.5 \pm 0.2$	4.5
107	$2840 \pm 380$	$4.2 \pm 0.4$	$0.99 \pm 0.01$	$17.6 \pm 0.3$	$17.7 \pm 0.3$	$19.3 \pm 0.4$	5.5



Table 2

Local fits of  $^{15}\text{N}$  CPMG relaxation dispersion for CheY in the presence of 1 mM EDTA.

Residue	$k_{ex}$ ( $s^{-1}$ )	$\Delta\omega$ (ppm)	Pt	$R_2^o$ ( $s^{-1}$ ) 600 MHz	$R_2^o$ ( $s^{-1}$ ) 700 MHz	$\chi^2$
18	$288 \pm 57$	$3.9 \pm 0.7$	$0.99 \pm 0.01$	$15.7 \pm 0.3$	$17.0 \pm 0.3$	2.9
19	$180 \pm 30$	$5.7 \pm 1$	$0.96 \pm 0.01$	$16.1 \pm 0.7$	$17.4 \pm 0.9$	7.6
36	$2400 \pm 500$	$0.6 \pm 0.1$	$0.86 \pm 0.04$	$13.3 \pm 0.2$	$14.1 \pm 0.3$	2.6
38	$2360 \pm 640$	$3.6 \pm 0.4$	$0.99 \pm 0.01$	$14.0 \pm 0.5$	$15.8 \pm 0.7$	4.3
62	$1820 \pm 520$	$3.7 \pm 0.4$	$0.99 \pm 0.01$	$13.4 \pm 0.4$	$13.8 \pm 0.7$	11
64	$96 \pm 16$	$9.4 \pm 1$	$0.80 \pm 0.04$	$20.0 \pm 1$	$22.2 \pm 2$	13
67	$1860 \pm 920$	$0.4 \pm 0.1$	$0.73 \pm 0.1$	$15.2 \pm 0.4$	$16.3 \pm 0.6$	3.5
68	$1400 \pm 790$	$3.5 \pm 0.8$	$0.99 \pm 0.01$	$16.1 \pm 0.4$	$15.9 \pm 0.6$	1.5
69	$2400 \pm 1100$	$2.3 \pm 0.5$	$0.99 \pm 0.01$	$15.0 \pm 0.5$	$15.0 \pm 0.8$	1.1
87	$2130 \pm 540$	$0.53 \pm 0.1$	$0.81 \pm 0.05$	$14.2 \pm 0.3$	$15.9 \pm 0.4$	2.3
94	$1330 \pm 800$	$0.27 \pm 0.1$	$0.67 \pm 0.2$	$15.4 \pm 0.2$	$16.1 \pm 0.3$	3.3
107	$2440 \pm 570$	$0.52 \pm 0.1$	$0.80 \pm 0.08$	$14.8 \pm 0.2$	$16.1 \pm 0.4$	2.6
127	$1980 \pm 550$	$1.6 \pm 0.2$	$0.98 \pm 0.01$	$15.8 \pm 0.3$	$16.5 \pm 0.5$	2.5