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# **ORIGINAL ARTICLE**

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# Characteristics of Eddy Current Attenuation and Thickness Measurement of Metallic Plate



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# Abstract

In eddy current testing, the law of attenuation of eddy current (EC) is of great concern. In conductive half space under the excitation of uniform magnetic field, the EC density decreases exponentially in the depth direction. However, in conductor with finite thickness tested by coil, the distribution of EC in the depth direction is more complicated. This paper studies the characteristics of EC attenuation in metallic plate of finite thickness. Simulation results show that there is an EC reflection at the bottom of plate, which changes the law of EC attenuation. A new concept, namely the equivalent attenuation coefficient, is proposed to quantify the speed of EC attenuation. The characteristics of EC attenuation are utilized to explain the nonmonotonic relation between coil voltage and plate thickness. Procedure of selecting frequency is discussed. Thereafter, measurement of plate thickness is carried out and accurate result is obtained.

Keywords: Eddy current testing, Attenuation of eddy current, Thickness measurement, Metallic plate

### **1** Introduction

Eddy current testing (ECT) which is based on the principle of electromagnetic induction is regarded as an effective nondestructive testing method [1]. It can be used to detect surface [2, 3] and subsurface [4, 5] defects in conductive materials and measure thicknesses [6, 7] of metallic plates.

It is well known that eddy current (EC) density is maximal on conductor surface and decreases exponentially along the depth direction, which is called skin effect [8, 9]. Figure 1(a) illustrates the skin effect, in which the dashed circles represent the EC flows at various depths and larger thickness of circle means larger amplitude of EC density. The capability of EC penetration is determined by the speed of attenuation of EC [10, 11]. The speed of attenuation of EC is related to the reduction of the normalized EC density per unit depth rather than referring to the absolute value of EC density, in which the normalized EC density is computed by dividing the amplitude of EC density by the value on the conductor

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surface. Thus, the speed of attenuation of EC is independent of the intensity of the excitation current. The depth at which the normalized EC density equals 1/e (36.8%) is referred to as skin depth and is illustrated in Figure 1(b). A prior knowledge of skin depth is of great concern when detecting subsurface defects and measuring thicknesses of plates.

In isotropic conductive half space under the excitation of uniform magnetic field, the relation between normalized EC density  $J / J_0$  and depth d is

$$I/J_0 = e^{-\alpha d},\tag{1}$$

where *a* is the attenuation coefficient. The skin depth denoted as  $\delta_S$ , is given by

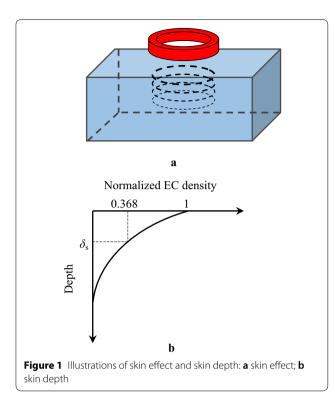
$$\delta_{\rm S} = \frac{1}{\alpha} = \frac{1}{\sqrt{\pi f \,\mu\,\sigma}},\tag{2}$$

where f is the excitation frequency,  $\mu$  is the permeability,  $\sigma$  is the conductivity. This skin depth is referred to as the standard skin depth.

It is found that the skin depth of EC in conductive plate with infinite thickness under the excitation of pancake coil can have a large difference from the standard skin depth [10]. The discrepancy is attributed to the diffusion

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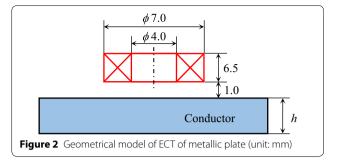
effect and the combined cancellation/diffusion effect of EC and is greatly dependent on the geometric parameters of pancake coil, as explained in Ref. [11].

In the available references discussing skin effect in ECT [10, 11], the test samples are assumed to have large thicknesses such that the EC densities at the bottom surfaces are very small and can be neglected. In many applications of ECT, however, the test samples have finite thicknesses and the EC densities at the bottom surfaces cannot be neglected. Study on the characteristics of EC attenuation in flat plate of finite thickness is not found. This paper attempts to fill in the blank. We utilizes the finite element analysis software AxisymMag [12] developed in lab for solving axisymmetric magnetic field problems to compute EC densities. The software has been verified by comparing computation results with analytical solutions [12].

The  $A_{\phi}$  formulation [13] is used in the paper where  $A_{\phi}$  is the circumferential component of the modified magnetic vector potential. The governing equation is

$$-\nu \frac{\partial^2 A_{\phi}}{\partial z^2} + \nu \frac{A_{\phi}}{\rho^2} - \frac{\nu}{\rho} \frac{A_{\phi}}{\partial \rho} - \nu \frac{\partial^2 A_{\phi}}{\partial \rho^2} + \sigma \frac{\partial A_{\phi}}{\partial t} = J_{s_{\phi}},$$
(3)

where  $\nu$  is the reluctivity of the material,  $\sigma$  is the conductivity, and  $J_{s\phi}$  is the circumferential component of the excitation current density. Eq. (3) is solved by the Galerkin method.



The paper then studies the measurement of plate thickness. Thickness is an important variable that plays an essential role in controlling product performance and cost [14, 15]. Methods for thickness measurement have been presented in literature. Model based approaches seek for plate thickness by comparing measured signal with model predicted signal and updating plate thickness in the model by certain optimization methods [16, 17]. The approaches have been applied in measuring thicknesses of conducting plates [18, 19], coatings [20, 21] and air gaps in aeronautical multilayered structures [22, 23]. Feature based methods measure plate thickness based on the analysis of signal characteristics. Signal features can be signal phase [24, 25] and amplitude [26, 27] in sinusoidal ECT or peak value [28, 29] and rise time [30, 31] in pulsed ECT. Other features include the no-reference-needed feature [32] and the frequency-tozero feature [33]. This paper uses sinusoidal excitation and selects voltage amplitude as the feature for thickness measurement.

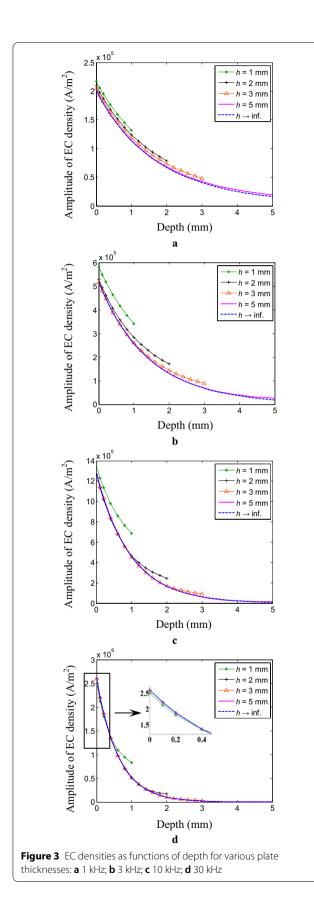
In short, the paper investigates the characteristics of attenuation of EC in the depth direction in metallic plate of finite thickness tested by pancake coil. Based on this, measurement of plate thickness is studied.

## 2 Characteristics of EC Attenuation in Metallic Plate

Figure 2 shows the geometrical model of ECT of metallic plate. The plate is assumed to be infinitely large in the directions parallel to the conductor surface. Therefore the model can be regarded as axisymmetric. The dimensions of the pancake coil are also shown in the figure. The density of the excitation current is  $3.077 \times 10^6$  A/m<sup>2</sup>. The material of the test sample has a conductivity of  $2.0 \times 10^7$ S/m and a relative permeability of 1. The thickness of the plate and the working frequency are variable.

### 2.1 Analysis of EC Density

Figures 3(a)–(d) show the amplitudes of EC densities as functions of depth for various plate thicknesses at 1 kHz, 3 kHz, 10 kHz, and 30 kHz, respectively. For comparison,



the EC densities in conductive half space are also shown. Two characteristics are found.

(1) The EC densities at the bottoms of the plates are increased comparing with those at the same depths in conductive half space. It means that there is reflection of EC at the bottoms of the plates. To explain the changes of EC densities, we plot the incidence EC densities  $J_i$  and the reflection EC densities  $J_r$  as functions of depth at frequencies of 3 kHz and 30 kHz, respectively as shown in Figures 4 and 5. The incidence EC densities are actually the EC densities in conductive half space. The reflection EC densities are obtained by subtracting the incidence EC densities from the total EC densities. It is seen that with reducing plate thickness the reflection EC density is increased and the reflection EC densities are larger at the bottoms of the plates and smaller at the tops of the plates. At the bottoms of the plates at either frequency, the phase differences between the reflection EC densities and the corresponding incidence EC densities are smaller than 90°; therefore the reflection of EC increases the total EC densities at the bottoms of the plates.

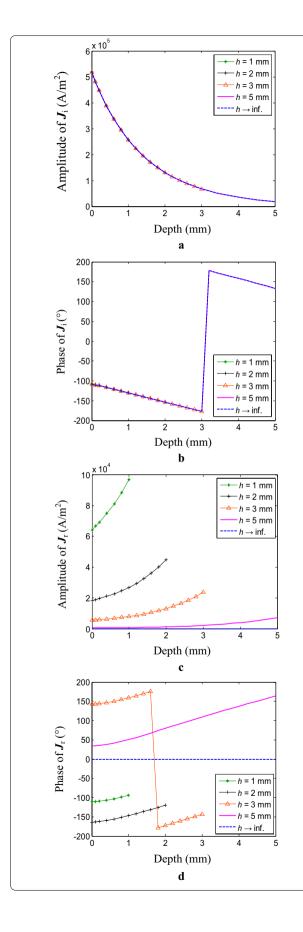
(2) The EC density at the top of the plate of finite thickness can be larger (e.g., when the frequency is 3 kHz and the plate thickness is 1 mm) or smaller (e.g., when the frequency is 30 kHz and the plate thickness is 1 mm) than that at the top of the conductive half space. This is because the phase difference between the reflection EC density and the incidence EC density at the top of the plate can be acute or obtuse. When the frequency is 30 kHz and the plate thickness is 1 mm, the phases of the reflection EC density and the incidence EC density at the top surface are 89.44° and -123.6°, respectively. Their angle is larger than 90°. Therefore the EC density at the top of the plate in this case is reduced.

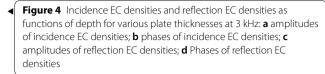
#### 2.2 Analysis of Normalized EC Density

Figures 6(a)–(d) show the normalized EC densities as functions of depth for various plate thicknesses at 1 kHz, 3 kHz, 10 kHz, and 30 kHz, respectively. The following characteristics are found.

(1) Within certain ranges of working frequency and plate thickness, the speed of attenuation of EC in plate of finite thickness is much slower than that in conductive half space. Two cases are analyzed. One is that, both the EC densities at the top surface and the bottom surface are increased and the increment of EC density at the bottom of the plate is larger than that at the top of the plate (e.g., when the frequency is 3 kHz and the plate thickness is 1 mm). The other case is that the EC density at the bottom of the plate is increased whereas the EC density at the top of the plate is increased whereas the EC density at the top of the plate is decreased (e.g., when the fre-







quency is 30 kHz and the plate thickness is 1 mm). In both cases, the speed of attenuation of EC is decreased.

- (2) When the frequency is low (e.g., 1 kHz) or the thickness of the plate is large (e.g., 5 mm), the difference between the speed of attenuation of EC in plate of finite thickness and that in conductive half space is inconspicuous. At low frequency, the variation of the reflection EC density in the range of plate thickness is slow. With large plate thickness, the incidence EC density at the bottom of the plate is small, which results in small reflection of EC. In both cases, the influence of EC reflection on the speed of attenuation of EC is not obvious.
- (3) With fixed frequency, the attenuation of EC in thinner plate is slower than that in thicker plate. The reason is that compared with the EC density in thicker plate, the increase of EC density at the bottom of the plate is larger than that at the top of the plate.

## 2.3 Equivalent Attenuation Coefficient

When the plate thickness is small, the ratio of EC density at the bottom of the plate to that at the top of the plate may be larger than 1/e. In this case, skin depth cannot be obtained. Furthermore, the relation between EC density and depth in the plate no longer possesses an exponential function owing to the nonuniform excitation field and the EC reflection at the bottom of the plate. Thus the attenuation coefficient in Eq. (1) is not applicable. In this scenario, we propose the concept of equivalent attenuation coefficient to take place of skin depth and attenuation coefficient to quantify the speed of attenuation of EC. The equivalent attenuation coefficient is defined in such a way that the curve of normalized EC density vs. depth is fitted by an exponential function in the form of

$$J/J_0 = e^{-\tilde{\alpha}d},\tag{4}$$

and the coefficient of the exponential term is the equivalent attenuation coefficient.

Table 1 shows the values of equivalent attenuation coefficient of EC in the test samples at different



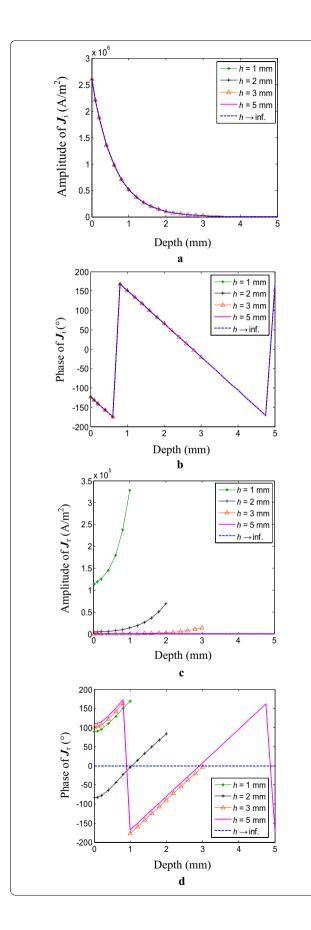


 Figure 5 Incidence EC densities and reflection EC densities as functions of depth for various plate thicknesses at 30 kHz; a amplitudes of incidence EC densities; b phases of incidence EC densities; c amplitudes of reflection EC densities; d phases of reflection EC densities

frequencies. It can be seen that at a certain frequency, the equivalent attenuation coefficient is smaller for thinner plate, which means that the EC reflection at the bottom of thinner plate is stronger, and the attenuation of EC in the plate becomes slower. When the plate thickness is fixed, the equivalent attenuation coefficient increases with increasing excitation frequency.

# **3** Measurement of Plate Thickness

This section studies the measurement of plate thickness. The relation between coil voltage and plate thickness is obtained by numerical simulation. The characteristics of EC attenuation obtained in Section 2 are utilized to explain such relation. Thereafter, the working frequency is selected and the experiment of thickness measurement is performed.

The method of measuring plate thickness is as follows. First, the relation between coil voltage and plate thickness is established by numerical simulation. Then we choose working frequency such that the range of plate thickness of interest is coved by the linear part of the curve of voltage-thickness relation. Given measured voltage, the plate thickness is found by interpolating the curve.

#### 3.1 Relation between Coil Voltage and Plate Thickness

Figure 7 shows the simulation result of coil voltage due to EC vs. plate thickness obtained using the model and parameters shown in Figure 2. As the plate thickness increases, the coil voltage due to EC increases dramatically, then decreases slowly, and finally approaches constant.

The relation between coil voltage due to EC and plate thickness is analyzed as follows. As illustrated in Figure 8, with increasing plate thickness, the volume of EC distribution becomes larger, which has an effect of increasing the coil voltage. At the same time, according to the analysis in Section 2, the increase of plate thickness reduces the EC reflection at the bottom of the plate and the attenuation of EC becomes faster, which has a trend of reducing the coil voltage. When the plate is thin, the voltage increases dramatically as the plate thickness increases, indicating that the increasing effect is dominant over the decreasing effect. As the plate thickness increases, the distance between EC in the range of  $\Delta d$  and coil becomes larger, resulting in weaker influence of this part of EC on the coil. When the plate thickness reaches a certain value, the increasing effect and the decreasing effect are



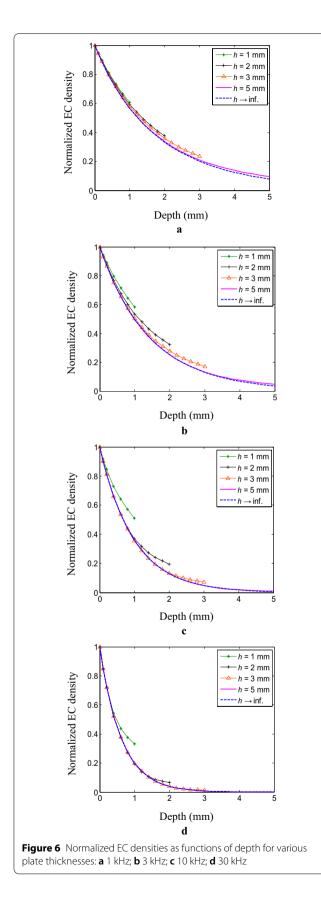
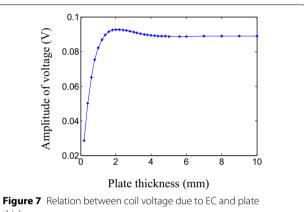


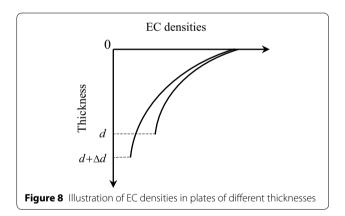
 Table 1 Equivalent attenuation coefficient of EC in the test

 samples of various thicknesses at different frequencies

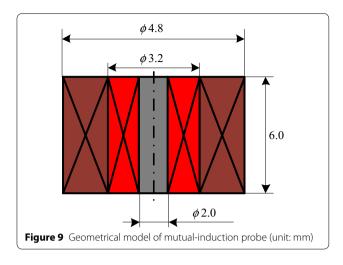
Plate thickness (mm)	Equivalent attenuation coefficient $\widetilde{lpha}$				
	1 kHz	3 kHz	10 kHz	30 kHz	
1	0.4995	0.5409	0.6684	1.2400	
2	0.5011	0.5886	0.9211	1.6070	
3	0.5054	0.6342	1.0080	1.6170	
5	0.5191	0.6776	1.0200	1.6180	
inf.	0.5283	0.6810	1.0200	1.6180	



thickness



balanced and the coil voltage reaches the maximum. When the plate is thicker, the decreasing effect is greater than the increasing effect. Accordingly, the coil voltage reduces when the plate thickness increases. When the plate thickness is large, there is almost no EC reflection at the bottom of the plate, and the increased volume of EC has little effect on the coil. Both of the increasing effect and the decreasing effect are negligible, leading to a constant voltage.

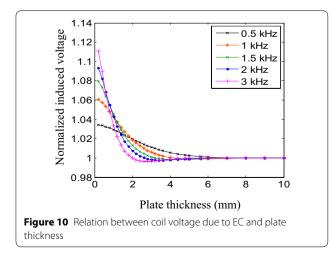


#### 3.2 Selection of Frequency

In thickness measurement, the working frequency is a critical parameter. This subsection studies the selection of frequency.

A probe is designed for measuring thickness of metallic plate. As shown in Figure 9, the probe consists of an excitation coil in the inside and a reception coil in the outside. The liftoff is 0.5 mm. A magnetic core with a relative permeability of 650 is placed at the center of the probe to improve detection sensitivity. The plate being tested has an electrical conductivity of  $3.5 \times 10^7$  S/m and a relative permeability of 1. The thickness of the plate ranges from 0.2 mm to 10.0 mm.

Figure 10 shows the relations between normalized output voltage and plate thickness at different frequencies. For the convenience of comparing the voltages under different frequencies, at each frequency, the voltages for various plate thicknesses have been divided by the voltage for conductive half space. Each of the curves in Figure 10 has a change rule that is reverse to that in Figure 7,



because Figure 7 only shows the coil voltages attributed to EC and each point in Figure 10 is the total coil voltage which is composed of the voltage caused by EC and the voltage generated by the mutual-induction between the excitation coil and the reception coil.

The frequency is selected by the following process. First of all, the range of measurement of plate thickness is specified. Then fit each curve in Figure 10 in the range of measurement by a straight line and compute the corresponding fitting correction coefficient  $R^2$  and the slope of the line. Finally, select the frequency such that the value of  $R^2$  meets requirement and the slope is the maximum. In this example, the range of measurement is [0, 2] mm. In this range, the values of  $R^2$  of linear fitting of the curves at 0.5 kHz, 1 kHz, 1.5 kHz, 2 kHz, and 3 kHz are 0.99, 0.99, 0.97, 0.94, and 0.86, respectively. The requirement of  $R^2$  for a good linear fitting in this case is that  $R^2$  should be larger than 0.9. Among the curves meeting this requirement, the curve at 2 kHz has the largest slope. Thereby, the frequency is selected as 2 kHz.

#### 3.3 Experiment of Thickness Measurement

After selecting working frequency, the experiment of measuring plate thickness is conducted. The experimental system includes the mutual-induction probe, phase-locked amplifier 7280 and signal generator DG3061. The signal generator supplies a sinusoidal signal with frequency of 2 kHz and peak-to-peak value of 1.5 V. Each of the aluminum plates has a square surface with side length of 100 mm. The plate thicknesses range from 0.482 mm to 1.968 mm.

When doing experiment, we put the probe on the plates in succession from the smallest thickness to the largest thickness and read the voltages on the reception coil. This process is then repeated twice. Table 2 lists the results of three measurements at the same position of each plate and the corresponding average values. Excellent repeatability of measurement is seen from the figure.

Finally, another plate with known thickness of 1.194 mm is used to verify the accuracy of thickness measurement. The average voltage of three measurements is 40.627 mV. Lagrangian interpolation using this voltage and the data in Table 2 gives a thickness of 1.189 mm. The absolute error is 0.005 mm and the relative error is 0.42%.

## 4 Conclusions

The characteristics of variation of EC in the depth direction in flat plates of finite thicknesses are systematically studied. More specifically, EC densities in metallic plates excited by pancake coils are computed using the finite element method and then the distributions of EC densities in the depth direction are plotted and analyzed. The

Plate thickness (mm)	1st measurement (mV)	2nd measurement (mV)	3rd measurement (mV)	Average (mV)
0.482	42.110	42.112	42.112	42.111
0.755	41.430	41.432	41.428	41.430
0.970	40.954	40.954	40.954	40.954
1.465	40.404	40.406	40.406	40.405
1.968	40.170	40.172	40.170	40.171

Table 2 The output voltages for various plate thicknesses

results show that the speed of attenuation of EC in metallic plate of finite thickness is slower than that in conductive half space. The thinner the plate is, the slower the EC attenuates. This fact indicates that reflection of EC happens at the bottom of the plate. The speed of EC attenuation in metallic plate of finite thickness is quantified by equivalent attenuation coefficient, a new concept proposed in this paper.

Based on the analysis of the behavior of EC in metallic plate of finite thickness, we elaborate the relation between coil voltage and plate thickness. Thereby, selection of working frequency is discussed. Finally, the experiment of thickness measurement is carried out and plate thickness is found by interpolating the voltage-thickness relation given measured coil voltage.

#### Authors' Contributions

ZZ was in charge of the whole trial; PD and JL wrote the manuscript; JL and SJ conducted the simulation and experiment; JL and YD assisted with sampling and laboratory analyses. All authors read and approved the final manuscript.

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#### Competing Interests

The authors declare that they have no competing interests.

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