

Analysis of the Accuracy of Actuation Electronics in the Laser Interferometer Space Antenna Pathfinder

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The Laser Interferometer Space Antenna Pathfinder (LPF) main observable, labelled Δg , is the differential force per unit mass acting on the two test masses in free fall condition, after the contribution of all non-gravitational forces have been compensated. At low frequencies the differential force is compensated by an applied electrostatic actuation force, which then must be subtracted from the measured acceleration to obtain Δg . Any inaccuracy in the actuation force contaminates the residual acceleration. This study investigates the accuracy of the electrostatic actuation system and its impact on the LPF main observable. It is shown that the inaccuracy is mainly caused by the rounding errors in the waveform processing and also by the random error caused by the analog to digital converter (ADC) random noise in the control loop. Both errors are one order of magnitude smaller than the resolution of the commanded voltages. We developed a simulator based on the LPF design to compute the close-to-reality actuation voltages and consequently, the resulting actuation forces. The simulator is applied in post-processing the LPF data.

I. INTRODUCTION

The Laser Interferometer Space Antenna (LISA) aims to detect gravitational waves in space¹. The mission was approved by the European Space Agency (ESA) in June 2017

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and it is currently in the first planning stages. The technology requirements for LISA were successfully tested by the LISA Pathfinder²⁻⁷ satellite (LPF), which flew from December 2015 to July 2017. In the LPF mission, two test masses (TMs) were put in nearly perfect free fall condition and their residual acceleration was measured precisely. Identical TMs will be used in the LISA satellites. Each TM is a gold-platinum cube with 46 mm edge length, surrounded by a set of electrodes in an electrode housing (EH), as shown in figure 1. Each electrode creates a capacitor with the TM surface. The electrodes are used to simultaneously sense (electrostatic sensing) and, for all degrees of freedom except the sensitive x interferometer axis, control (electrostatic actuation) the TM position with respect to the centre of their EH. Applying a voltage to each electrode induces an electric field and consequently an electrostatic control force/torque on the TM. The desired magnitude of the force/torque, that should be applied on the TM, is calculated by the drag-free attitude control system (DFACS). DFACS also calculates the analytical conversion of force/torque to voltage. The set of TM, EH, charge management system and electronics is called the gravitational reference sensor (GRS). The analog circuits of the GRS front-end electronics (FEE) and the digital code of the field programmable gate array (FPGA) chip inside FEE are responsible for applying the commanded voltages to the corresponding electrodes. The applied voltages are not exactly equal to the commanded voltages, because of hardware limitations. The difference between commanded and applied voltages induces a consequent difference between commanded and applied forces/torques, resulting in loss of actuation accuracy. Quantifying and characterizing this inaccuracy (actuation error) is crucial to understand how this affects the main mission observable of both LPF and LISA.

In the LPF mission two GRS are located in the one spacecraft. In science mode, the spacecraft uses micro-newton thrusters to follow one of the TMs (labeled TM1) along its natural geodesic trajectory. The other TM (labeled TM2) is actuated, such that it follows the spacecraft and TM1. The main LPF observable, Δg , is the total residual acceleration of the two TMs in free fall condition, along the main measurement axis x , and it is calculated by^{6,7}

$$\Delta g(t) = \ddot{x}_{12}(t) + \omega_2^2 x_{12}(t) + (\omega_2^2 - \omega_1^2) x_1(t) - g_c(t) - g_\Omega(t) \quad (1)$$

\ddot{x}_{12} is the second derivative of the relative displacement (relative acceleration) of the TMs, which is measured by a laser interferometer. x_1 is the absolute displacement of TM1 with respect to the spacecraft and it is measured by a dedicated interferometer. ω_1^2 and ω_2^2 are parasitic stiffnesses and they can be considered as spring constants per unit mass for oscillatory like force couplings between the TMs and the electrode housing. $g_\Omega(t)$ is the centrifugal force per unit mass and $g_c(t)$ is the actuation force per unit mass applied to TM2 to follow TM1. According to equation 1, error in the actuation force along the x -axis contaminates the total residual acceleration, Δg . Thus, it is important to quantify the actuation force error and consider its effect on the calculation of Δg for LPF.

The residual acceleration in LPF is measured with a precision of $1.74 \pm 0.01 \text{ fm s}^{-2}/\sqrt{\text{Hz}}$ above 2 mHz and $60 \pm$

$10 \text{ fm s}^{-2}/\sqrt{\text{Hz}}$ at $20 \mu\text{Hz}$ ⁷. This was achieved, among others, by a more accurate calculation of the electrostatic actuation force based on the method presented in this paper. This paper addresses systematic inaccuracies, and an effective correction for these effects, in the actuation waveforms. A brief description of the FEE that is important for this study is provided in section II. More details on FEE and how it simultaneously serves for actuation and sensing is given in references⁸⁻¹¹. In order to study the actuation accuracy, we have implemented a simulator that reproduces the "close-to-reality" behavior of the FPGA code and the actuation electronics. The critical elements that cause actuation errors are elaborated in section III and the simulator is described in IV. As shown later in the paper, the impact of actuation accuracy on the residual acceleration is especially important at frequencies below 1 mHz. Analysis of the actuation error and the impact on LPF data is discussed in section V. A separate article, (to be published) will address low frequency force noise from slow random fluctuations in the actuators¹².

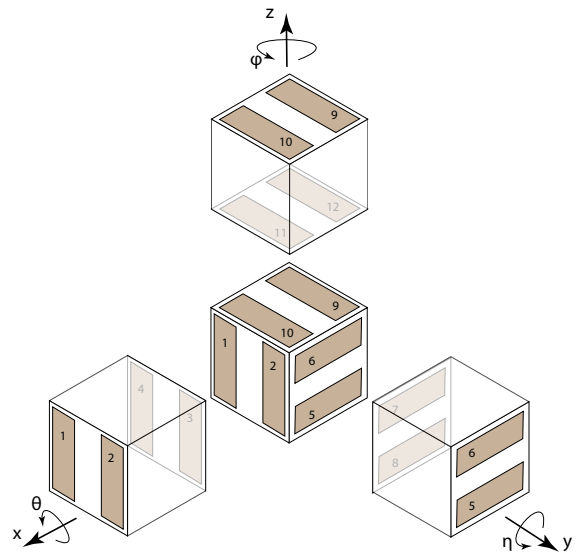


FIG. 1. Electrode arrangement around the TM. The illustrated electrodes are used for TM actuation and position sensing.

II. ACTUATION ELECTRONICS

TM actuation is possible in all degrees of freedom (DoF) by applying commanded force and/or torque to the corresponding electrode. Each electrode is connected to the actuation electronics, which is responsible for exactly two DoF: one translation and one rotation. Considering figure 1, electrodes 1 to 4 are applied for actuation along x and ϕ , electrodes 5 to 8 for y and θ and electrodes 9 to 12 for z and η .

Figure 2 illustrates the block diagram for actuation along x and ϕ . Actuation along other DoF is similar. Given the commanded voltage amplitudes for translation and rotation by the DFACS, two orthogonal sinusoidal waveforms (AC signals)

are generated, using a common lookup table (LuT). The LuT contains a quarter sinusoid composed of 100 samples and it is used to make samples at a rate of 12 kHz using the waveforms. The FPGA logic uses that LuT to generate waveforms at a frequency of $n \times 30$ Hz. For the waveforms at those frequencies, each sample can be calculated by using an entry of the same LuT, where the frequency depends on how the entries are picked. The integer n is set differently for each DoF, such that the generated waveforms are all orthogonal to each other and no cross-talk occurs. This yields 60 Hz and 270 Hz for actuation along x and ϕ , 90 Hz and 240 Hz for y and θ , and 120 Hz and 180 Hz for z and η . The amplitude of the generated sinusoidal waveforms, labelled V_{cmd} , are commanded by the DFACS at 10 Hz frequency. Thus, the waveforms are applied for 100 ms, that form integer multiplications of the waveform periods. The commanded voltage amplitudes are discretized with a resolution of $153 \mu\text{V}$, described in this paper as the least significant bit (LSB) voltage. In addition to these AC signals, DC voltages, labelled $V_{\text{cmd, DC}}$, can also be commanded by DFACS to compensate for parasitic electrode voltages, and to aid TM charge measurement and discharge. The AC and DC signals are combined as given in equation 2.

After combining the waveforms, the combined signal is mapped from the LSB value used for the commanded voltage ($\pm 10 \text{ V}/2^{17} \approx 153 \mu\text{V}$) to the one used for the proportional, integral and derivative (PID) controller ($\pm 14.5 \text{ V}/2^{18} \approx 111 \mu\text{V}$). Thus, the LSB value of the control loop is a factor of 1.379 smaller. The PID controller was designed for a different LSB value in order to make use of the full conversion range of the analog to digital converter (ADC) used for its feedback, i.e. the signal is scaled to use the full range of the ADC output codes. This was done to minimize the noise contribution of the ADC, which has an input referred noise of $0.72 \text{ LSB}_{\text{RMS}}$, according to the data sheet for the ADC component LTC1604¹³. To maintain the same amplitude, the digital values are multiplied by the reciprocal factor before entering the loop. For that, they are multiplied by the factor 1.375, which is the closest value to 1.379 that can be represented with 3 fractional bits and it is well within the tolerance of the analog gain, for which a calibration is performed. This is shown in figures 2 and 3 by the "Scale 1.375" sub-blocks. As discussed later in this paper, scaling the waveform – as it is currently designed in LPF – causes an actuation error, which cannot be neglected. The current electronics design requires several optimizations, before it can be integrated in LISA. These are proposed in a separate study¹⁴.

In order to increase the signal resolution and reduce actuation error, the 12 kHz sampled signal is up-sampled to 96 kHz by linear interpolation and then fed to a sigma-delta control loop. The control loop consists of a PID controller with the derivative block configured feed-forward and the proportional and integral blocks acting on the PID error signal. To generate an analog signal at the electrode, a digital to analog converter (DAC) is used, followed by an analog integrator, which has an integration gain that fits with the feed-forward gain of the digital part. Finally, to provide the feedback for the PID controller, the signal is digitized by an ADC.

The signals injected to electrodes 1 to 4 are then given by

$$\begin{aligned} V_1(t) &= +V_{\text{cmd},1x} \sin(2\pi 60 t) + V_{\text{cmd},1\phi} \sin(2\pi 270 t) + V_{\text{cmd, DC}1} \\ V_2(t) &= -V_{\text{cmd},1x} \sin(2\pi 60 t) + V_{\text{cmd},2\phi} \cos(2\pi 270 t) + V_{\text{cmd, DC}2} \\ V_3(t) &= +V_{\text{cmd},2x} \cos(2\pi 60 t) - V_{\text{cmd},1\phi} \sin(2\pi 270 t) + V_{\text{cmd, DC}3} \\ V_4(t) &= -V_{\text{cmd},2x} \cos(2\pi 60 t) - V_{\text{cmd},2\phi} \cos(2\pi 270 t) + V_{\text{cmd, DC}4} \end{aligned} \quad (2)$$

The time-averaged x -force and ϕ -torque applied to the TM are calculated by

$$\begin{aligned} F_x &= \frac{1}{2} \left(\left| \frac{\partial C_{EL, TM}}{\partial x} \right| - \left| \frac{\partial C_{EL, H}}{\partial x} \right| \right) (\langle V_1^2 \rangle + \langle V_2^2 \rangle - \langle V_3^2 \rangle - \langle V_4^2 \rangle) \\ T_\phi &= \frac{1}{2} \left(\left| \frac{\partial C_{EL, TM}}{\partial \phi} \right| - \left| \frac{\partial C_{EL, H}}{\partial \phi} \right| \right) (\langle V_1^2 \rangle - \langle V_2^2 \rangle + \langle V_3^2 \rangle - \langle V_4^2 \rangle) \end{aligned} \quad (3)$$

with $C_{EL, TM}$ being the capacitance between each electrode and the TM, and $C_{EL, H}$ the one between the electrode and the housing, respectively.

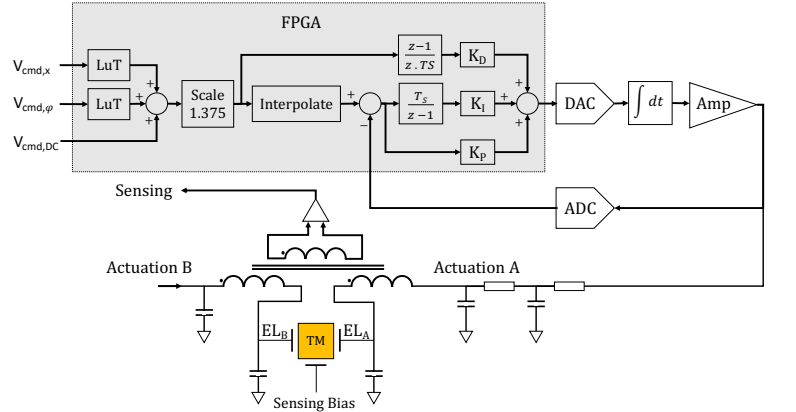


FIG. 2. Circuit block diagram featuring actuation elements for two opposing electrodes. Here, the commanded voltages correspond to x/ϕ actuation. Actuation along other DoF is similar.

III. ACTUATION ACCURACY

The accuracy of the actuation system is mainly affected by rounding errors and also by the random noise caused by ADC in the control loop. These are elaborated below.

• Number of available bits in waveform generation

The AC commanded voltages from the DFACS are 16 bits unsigned integers with $153 \mu\text{V}$ resolution. These are multiplied by LuT words, which are 16 bits integers. The multiplication would result in 32 bits integers, whereas only 16 bits are used in the FPGA code. Therefore, the lower significant 16 bits are discarded and the generated waveform is rounded to 16 bits and this introduces a rounding error.

While a change of 1 LSB in the commanded AC voltage will indeed change the amplitude of the sinusoidal waveform (peak voltage) by 1 LSB, the rest of the

points in the waveform will change by 0 or 1 LSB, depending on the rounding. This results in changing the effective voltage on the electrode (see equation 4) by an amount around, but not exactly, 1 LSB, and it directly influences the applied force/torque, according to equation 3.

- **Scaling after waveform generation**

After scaling by the non-integer factor 1.375, the fractional part is removed. Similar to the last step, this also results in a rounding error. Data resolution at this stage is $153 \mu\text{V}/1.375 = 111 \mu\text{V}$ ⁸.

- **Error caused by ADC random noise**

Among the electronic elements used in the actuation system, DAC and ADC produce the major noise. Although the noise entering the feedback loop is compensated, the noise caused by the ADC is critical, because it directly affects the behavior of the loop. Simulation of the ADC noise is elaborated in the appendix A.

The first two points describe the rounding error that happens on single samples of a waveform. This error is one order of magnitude smaller than the resolution of the commanded voltages, as elaborated in section VA and illustrated in figure 5. The random error caused by ADC random noise is also one order of magnitude smaller than the resolution of the commanded voltages. This is discussed in section VB and illustrated in figure 6.

IV. SIMULATION

We implemented a simulation of the actuation system based on the LPF design, in order to reproduce the close-to-reality actuation voltages and consequently, actuation forces/torques. The critical elements affecting the actuation accuracy were mentioned in the previous section. They are implemented in the simulation in order to quantify their impact. The simulation is divided into two main parts: The first simulates only the waveform generation and scaling and essentially considers the systematic error between the perfect analog sinusoidal wave and the digitized generated waveform. We refer to this part as partial model (PM). The second part simulates the entire circuitry, which includes the waveform generation, scaling, digital controller and the feedback loop. We call this part entire model (EM). EM allows us to investigate the effect of ADC random noise on actuation accuracy and also the systematic error of the control loop.

The simulation inputs are the commanded (AC) voltages V_{cmd} , which are the peak amplitudes of the generated waveforms, and also the commanded DC voltage $V_{\text{cmd,DC}}$. The sinusoidal waveforms $V_i(t)$ applied to each electrode i are implemented according to equation 2, with i denoting electrodes 1 to 4, which are responsible for x/ϕ actuation. The effective voltage on each electrode is the root mean square (RMS) of $V_i(t)$. The rate of DFACS commanded voltages is 10 Hz and the waveforms are 100 ms long. The simulation outputs $V_{\text{out},x,i}$ and $V_{\text{out},\phi,i}$ are the amplitudes of the AC signal (peak

voltages) at 60 Hz and 270 Hz frequencies, which are numerically evaluated in the simulation by the fast Fourier transform of the waveform $V_i(t)$ applied to electrode i . The effective voltage on each electrode is related to the peak voltages as follows

$$\text{RMS}(V_i(t)) = \sqrt{\frac{1}{2} (V_{\text{out},x,i}^2 + V_{\text{out},\phi,i}^2)}, \quad i = 1, \dots, 4 \quad (4)$$

For each input command voltage, $V_{\text{out},x,i}$ and $V_{\text{out},\phi,i}$ are calculated by both PM and EM. For sake of simplicity, only the actuation of TM2 along x is discussed further in this paper and the applied voltages to electrode 1 are illustrated. Therefore, the indices x and $i = 1$ are removed and the simulation output is labelled V_{PM} or V_{EM} , with respect to PM and EM. The analysis for other electrodes and also other DoF provide similar results. Figure 3 illustrates the simulation scheme.

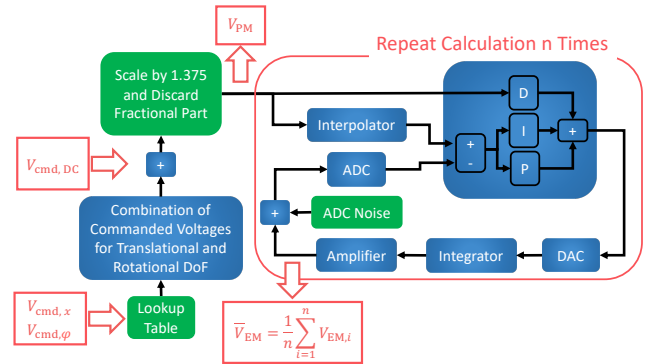


FIG. 3. Simulation scheme. V_{cmd} and $V_{\text{cmd,DC}}$ are the inputs. V_{PM} and V_{EM} are output of PM and EM, respectively.

V. ANALYSIS AND RESULTS

In this section, we elaborate on the analyses related to quantifying the systematic error of the waveform generation, the systematic error of the control loop and also the random error of the control loop. The systematic error of the waveform generation is studied by comparing V_{PM} to the commanded voltages and the systematic and also the random error of the control loop are investigated by comparing V_{EM} to V_{PM} .

For these analyses, the simulation input is an integer multiplication of the LSB value, given by $V_{\text{cmd}} = 153 \mu\text{V} \cdot \mathcal{D}_{\text{cmd}}$, whereby \mathcal{D}_{cmd} denotes the integer value of the digital code resulting from a commanded voltage. $V_{\text{cmd,DC}}$ is zero. We consider V_{cmd} from 153 mV to 535 mV, which correspond to the digital values of \mathcal{D}_{cmd} from 1000 to 3500. Voltages around 500 mV are especially interesting, because they correspond to the LPF commands during the acceleration measurements.

At the end of this section, we analyze the effect of actuation error on the LPF data. For this analysis, the simulation inputs are the commanded voltages for several LPF acceleration measurements.

A. Systematic error of the waveform generation

In figure 4, V_{cmd} is compared to V_{PM} . The resolution of V_{cmd} (vertical difference between adjacent steps) is equal to $153 \mu\text{V}$. For V_{PM} however, the resolution varies for different digital commands. This variation is due to the rounding error after waveform generation and also after the waveform scaling. As shown in figure 5, $\Delta V = V_{PM} - V_{cmd}$ is typically of the order of $10 \mu\text{V}$, which is one order of magnitude smaller than the input resolution. Considering equation 3, ΔV in average introduces an erroneous induced field and results in a force noise.

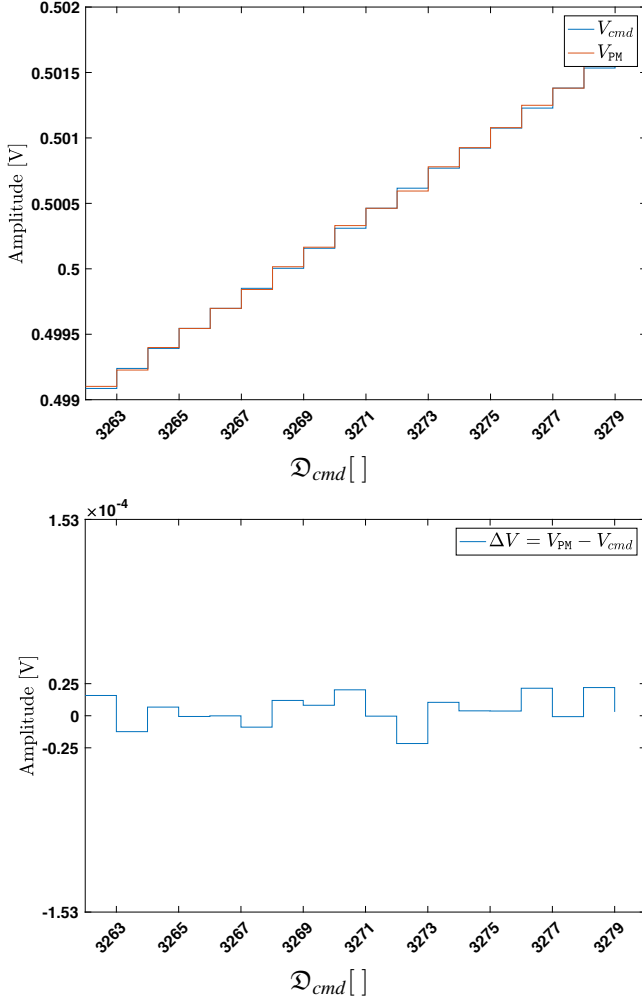


FIG. 4. Comparing V_{cmd} to V_{PM} shows that the resolution of V_{PM} varies for different input commands. The variation is one order of magnitude smaller than V_{cmd} resolution, $153 \mu\text{V}$.

B. Error of the control loop

The ADC random noise is a white noise and it dominates the systematic error of the control loop. Therefore, the simulation is implemented such that it allows suppressing the ran-

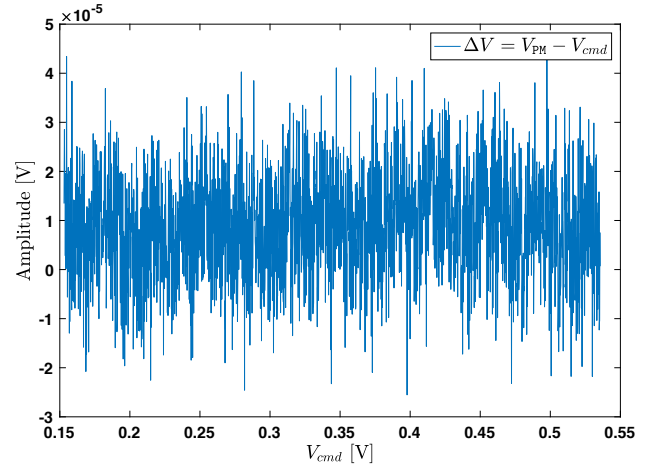


FIG. 5. The variation of V_{PM} resolution is one order of magnitude smaller than the resolution of commanded voltages, $153 \mu\text{V}$. Thus typical amplitude errors are of order 0.1 LSB.

dom noise, in order to quantify the systematic error of the control loop. To this purpose, the simulator repeats the calculation of V_{EM} (usually $n = 10^4, 10^5$ or 10^6 times) and outputs the averaged value $\bar{V}_{EM} = \frac{1}{n} \sum_{i=1}^n V_{EM,i}$.

A fit of the data from simulation outputs \bar{V}_{EM} and V_{PM} demonstrates a relation $\bar{V}_{EM} = \alpha V_{PM} + \beta$. The values α and β are different for various n . The error of the control loop is specified by $\Delta V_n = \bar{V}_{EM} - (\alpha V_{PM} + \beta)$.

The random error is determined by $\Delta V_{n=1}$ and it is illustrated in figure 6. The systematic error is verified by $\Delta V_{n=10^6}$ and it is shown in figure 7. In appendix B we elaborate why $n = 10^6$ is sufficient for analysing the systematic error. The systematic error is of the order of 10^{-8}V and its effect on the accuracy is negligible. It is at worst four orders of magnitude smaller than the resolution of the commanded voltages, and thus, 1000 times smaller than the systematic rounding error (see figures 4 and 5).

C. Effect of actuation error on the LPF data

The LPF main observable, Δg , is the residual acceleration of the two test masses and it is calculated by equation 1. According to this equation, the error of the actuation force per unit mass, $g_c(t)$, contaminates Δg , especially at frequencies below 1 mHz ⁶. In order to investigate the effect of actuation inaccuracy on the LPF data, we first calculated the actuation force by V_{cmd} , V_{PM} and \bar{V}_{EM} , respectively and afterwards, we considered the amplitude spectral density of the residual acceleration, $S_{\Delta g}^{1/2}(f)$, at low frequency, i.e. $f = 0.1 \text{ mHz}$, for the LPF acceleration measurements. As shown in figure 8, calculation of the actuation force by V_{PM} (red curve) results in lower amplitude spectral density, $S_{\Delta g}^{1/2}(f = 0.1 \text{ mHz})$, compared to the case in which the actuation force is calculated by V_{cmd} (blue curve).

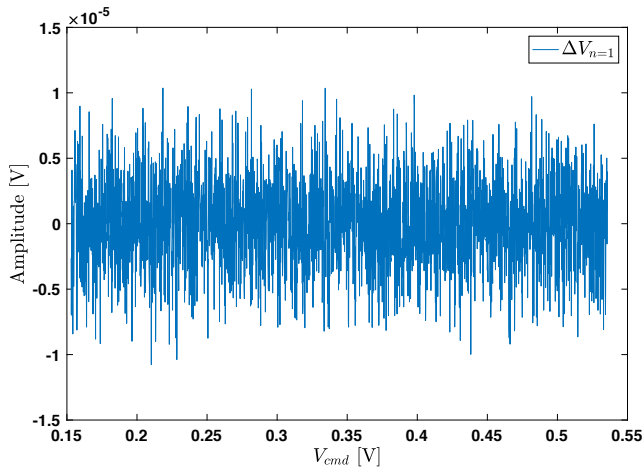


FIG. 6. The figure illustrates $\Delta V_{n=1} = V_{EM} - (0.73V_{PM} - 9 \text{ nV})$. Note that here $\bar{V}_{EM} = V_{EM}$, as $n = 1$. The random error caused by the ADC random noise is of the order of $10 \mu\text{V}$.

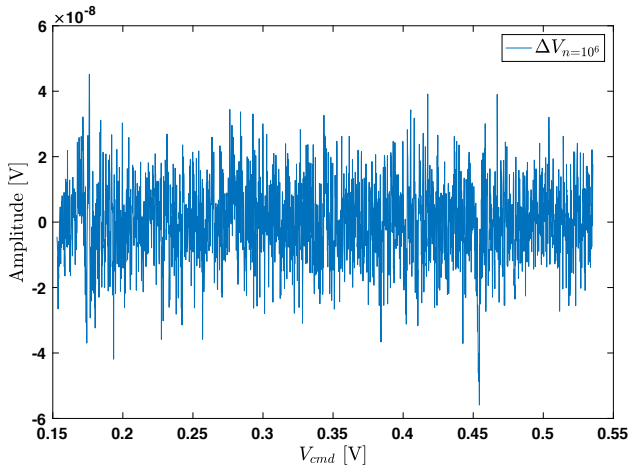


FIG. 7. The error of the control loop $\Delta V_{n=10^6} = \bar{V}_{EM} - (1.03V_{PM} - 2.7 \text{ nV})$ is four orders of magnitude smaller than the resolution of commanded voltages, $153 \mu\text{V}$.

Comparing the data calculated by V_{PM} (red curve) to the ones calculated by \bar{V}_{EM} (yellow curve) indicates that the governing actuation error refers to the error caused by rounding and scaling the waveform and not the error of the control loop. This is also illustrated in figure 9, where the amplitude spectral density of Δg is calculated for the LPF acceleration measurement in June 19th. As shown in the figure, considering the inaccuracy in calculation of Δg is especially important at low frequencies.

VI. CONCLUSION

We have identified, how the systematic rounding errors have a relevant, detectable impact on the AC actuation amplitudes and developed a simulator to verify the expected round-

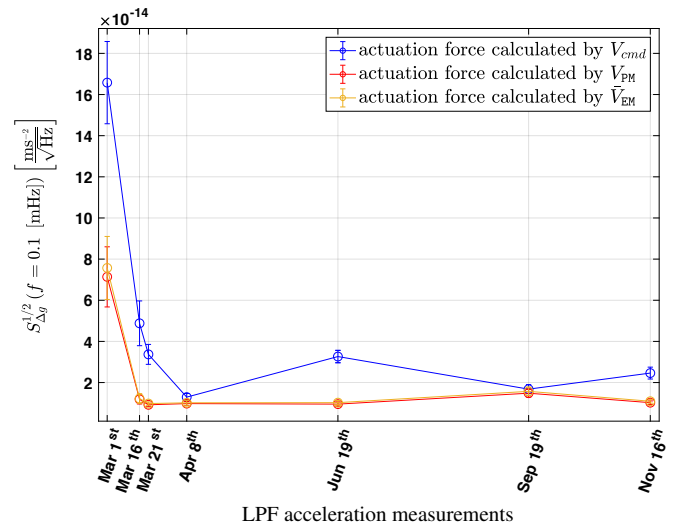


FIG. 8. Amplitude spectral density of the LPF residual test-mass acceleration at 0.1 mHz for different segments of acceleration measurement in 2016. Considering the actuation inaccuracy in calculation of Δg results in stationary amplitude spectral density $S_{\Delta g}^{1/2}$ at $f = 0.1 \text{ mHz}$ frequency.

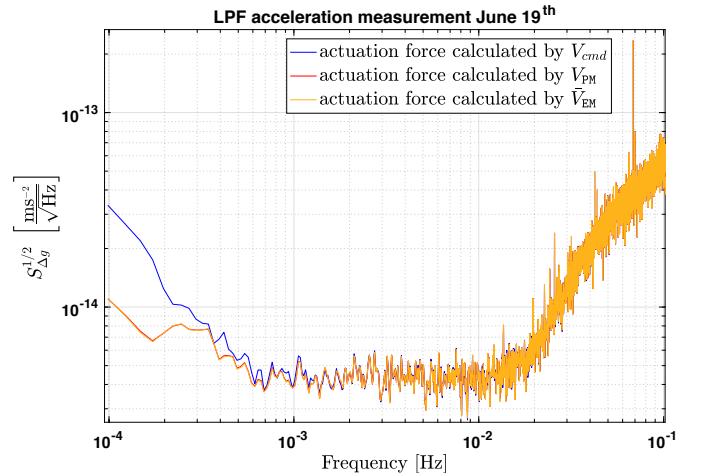


FIG. 9. Amplitude spectral density of the LPF residual test-mass acceleration for the measurement in June 2016.

ing behavior. Using the simulation results, we have also quantified the random error of the control loop, which is caused by the ADC random noise. Both of these errors are of the order of $10 \mu\text{V}$, which is one order of magnitude smaller than the resolution of the input commanded voltages ($153 \mu\text{V}$). We also analyzed the systematic error of the control loop, by suppressing the ADC random noise by three orders of magnitude in the simulation. It was shown, that this error is four orders of magnitude smaller than the resolution of the commanded voltages and thus, negligible.

Considering the systematic rounding errors, our simulator enables us to estimate the close-to-reality value of the actuation forces and more accurately calculate the residual accel-

ation of the test masses, which is the LPF main observable, at low frequencies. It is worth mentioning that actuation forces are calculated by our simulator in post-processing the LPF data.

Contrary to LPF, there is no actuation along the main measurement axis in LISA. However, optimizing the electronics to avoid the systematic error, specially in waveform generation, is relevant for accurate actuation of the TMs in other DoF. A proposal for optimizing the waveform generation for LISA will be provided in a separate study¹⁴.

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Appendix A: ADC Noise

The ADC noise density is calculated from the signal to noise ratio (SNR), which is 90.1 dB according to the data sheet for the ADC component LTC1604¹³.

$$\text{SNR} = 90.1 \text{ dB} = 20 \log \left(\frac{\text{RMS}_{\text{signal}}}{\text{RMS}_{\text{noise}}} \right) \quad (\text{A1})$$

The RMS value for the signal at the ADC input is calculated from the amplitude of the generated sinusoidal signal (A_{signal}) and it is given by

$$\text{RMS}_{\text{signal}} = \frac{A_{\text{signal}}}{\sqrt{2}} = \frac{20 \text{ V}/1.375}{\sqrt{2}} = 10.2852 \text{ V}_{\text{RMS}} \quad (\text{A2})$$

20 V refers to the maximum commanded AC amplitude for each actuation channel and the scaling factor is elaborated in section II. Therefore, $\text{RMS}_{\text{noise}} = 321.52 \mu\text{V}_{\text{RMS}}$ is obtained by substituting equation A2 in A1 and it corresponds to $\text{RMS}_{\text{noise}} = 0.72 \text{ LSB}$ given in the data sheet. The noise density is then given by

$$S_{\text{ADC}}^{1/2} = \frac{\text{RMS}_{\text{noise}}}{\sqrt{f_s}} = 1.0377 \frac{\mu\text{V}}{\sqrt{\text{Hz}}}, \quad (\text{A3})$$

whereby $f_s = 96 \text{ kHz}$ is the sampling frequency of the signal. The implemented noise in the simulation is a pseudo random white noise drawn from normal distribution with zero mean and standard deviation $\sigma_{\text{ADC}} = \text{RMS}_{\text{noise}} = 321.52 \mu\text{V}$.

For big number of iterations, such as $n = 10^5$ and 10^6 , the output of EM is parallel computed. In order to avoid the same random numbers generated by different processor cores, the random number generation is seeded, using the Matlab function `rng(seed)`¹⁵.

Appendix B: cross correlation

The random error of the control loop is labelled ΔV_1 , and its standard deviation is $\sigma_{\Delta V_1} = 3.5 \mu\text{V}$ for the commanded voltages 0.153 V to 0.536 V. These data are shown in figure 6.

According to the central limit theorem, we expect that the random error is reduced by $\sigma_{\Delta V_1} / \sqrt{10^6} = 3.5 \text{ nV}$ for $n = 10^6$ iterations. However, the standard deviation of ΔV_{10^6} obtained by the simulation is $\sigma_{\Delta V_{10^6}} = 11.8 \text{ nV}$ (ΔV_{10^6} data are illustrated in figure 7). It means that $\sigma_{\Delta V_{10^6}}$ derived from the simulated data is 3.37 times bigger than the value expected from the central limit theorem. Therefore, ΔV_{10^6} is not purely random, and it is dominated by the systematic error. This is also analysed by sample cross correlation function, as described below.

For the same input voltages, we calculated $\Delta V_{n=10^6}$ three times and labelled each calculation a *run*. The index $n = 10^6$ indicates that each \bar{V}_{EM} is averaged over 10^6 iterations. As shown in figure 10, the three runs follow similar trajectories. Furthermore, the sample cross correlation coefficient of combination of pairs of runs at lag zero is close to one, as illustrated in figure 11. This proves the dependency of $\Delta V_{n=10^6}$ on the commanded voltages and indicates a systematic error in the control loop.

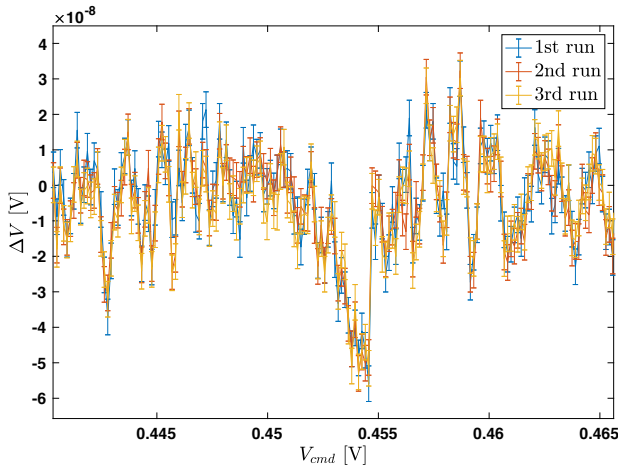


FIG. 10. $\Delta V = \vec{V}_{EM} - V_{PM}$, fitted is calculated three times. The parallel trajectories indicate a systematic error in the control loop.

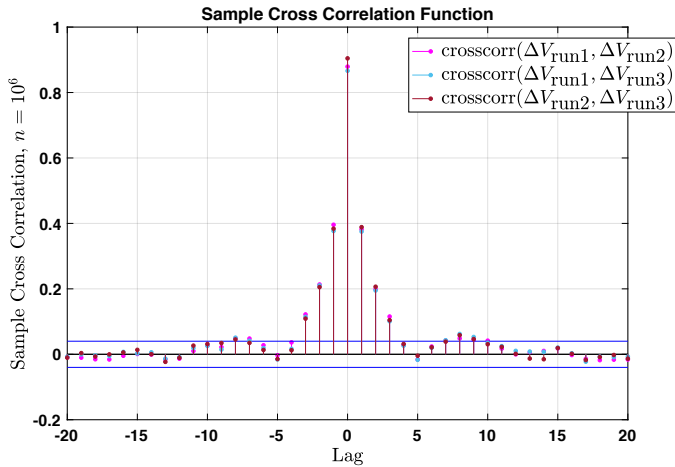


FIG. 11. The sample cross correlation of each two runs at lag zero proves the systematic error in the control loop.

¹P. Amaro-Seoane, H. Audley, S. Babak, J. Baker, E. Barausse, P. Bender, E. Berti, P. Binetruy, M. Born, D. Bortoluzzi, J. Camp, C. Caprini, V. Cardoso, M. Colpi, J. Conklin, N. Cornish, C. Cutler, K. Danzmann, R. Dolesi, L. Ferraioli, V. Ferroni, E. Fitzsimons, J. Gair, L. G. Bote, D. Giardini, F. Gibert, C. Grimani, H. Halloin, G. Heinzel, T. Hertog, M. Hewitson, K. Holley-Bockelmann, D. Hollington, M. Hueller, H. Inchauspe, P. Jetzer, N. Karnesis, C. Killow, A. Klein, B. Klipstein, N. Korsakova, S. L. Larson, J. Livas, I. Lloro, N. Man, D. Mance, J. Martino, I. Mateos, K. McKenzie, S. T. McWilliams, C. Miller, G. Mueller, G. Nardini, G. Nelemans, M. Nofrarias, A. Petiteau, P. Pivato, E. Plagnol, E. Porter, J. Reiche, D. Robertson, N. Robertson, E. Rossi, G. Rusano, B. Schutz, A. Sesana, D. Shoemaker, J. Slutsky, C. F. Sopuerta, T. Sumner, N. Tamanini, I. Thorpe, M. Troebs, M. Vallisneri, A. Vecchio, D. Vetrugno, S. Vitale, M. Volonteri, G. Wanner, H. Ward, P. Wass, W. Weber, J. Ziemer, and P. Zweifel, “Laser Interferometer Space Antenna,” (2017), [arXiv:1702.00786](https://arxiv.org/abs/1702.00786).

²P. McNamara, S. Vitale, and K. Danzmann, “LISA Pathfinder,” *Classical and Quantum Gravity* **25**, 114034 (2008).

³F. Antonucci, M. Armano, H. Audley, G. Auger, M. Benedetti, P. Binetruy, C. Boatella, J. Bogenstahl, D. Bortoluzzi, P. Bosetti, N. Brandt, M. Caleno, A. Cavalleri, M. Cesa, M. Chmeissani, G. Ciani, A. Conchillo, G. Congedo, I. Cristofolini, M. Cruise, K. Danzmann, F. De Marchi, M. Diaz-

Aguilo, I. Diepholz, G. Dixon, R. Dolesi, N. Dunbar, J. Fauste, L. Ferraioli, D. Fertin, W. Fichter, E. Fitzsimons, M. Freschi, A. Garcia Marin, C. G. Marirrodriga, R. Gerndt, L. Gesa, D. Giardini, F. Gibert, C. Grimani, A. Grynagier, B. Guillaume, F. Guzmán, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, J. Hough, D. Hoyland, M. Hueller, J. Huesler, O. Jennrich, O. Jennrich, P. Jetzer, B. Johlander, C. Killow, X. Llamas, I. Lloro, A. Lobo, R. Maarschalkerweerd, S. Madden, D. Mance, I. Mateos, P. W. McNamara, J. Mendes, E. Mitchell, A. Monsky, D. Nicolini, D. Nicolodi, M. Nofrarias, F. Pedersen, M. Perreur-Lloyd, A. Perreca, E. Plagnol, P. Prat, G. D. Racca, B. Rais, J. Ramos-Castro, J. Reiche, J. A. R. Perez, D. Robertson, H. Rozemeijer, J. Sanjuan, A. Schleicher, M. Schulte, D. Shaul, L. Stagnaro, S. Strandmoe, F. Steier, T. J. Sumner, A. Taylor, D. Texier, C. Trenkel, D. Tombolato, S. Vitale, G. Wanner, H. Ward, S. Waschke, P. Wass, W. J. Weber, and P. Zweifel, “From laboratory experiments to LISA Pathfinder: achieving LISA geodesic motion,” *Classical and Quantum Gravity* **28**, 094002 (2011).

⁴F. Antonucci, M. Armano, H. Audley, G. Auger, M. Benedetti, P. Binetruy, C. Boatella, J. Bogenstahl, D. Bortoluzzi, P. Bosetti, M. Caleno, A. Cavalleri, M. Cesa, M. Chmeissani, G. Ciani, A. Conchillo, G. Congedo, I. Cristofolini, M. Cruise, K. Danzmann, F. De Marchi, M. Diaz-Aguilo, I. Diepholz, G. Dixon, R. Dolesi, N. Dunbar, J. Fauste, L. Ferraioli, D. Fertin, W. Fichter, E. Fitzsimons, M. Freschi, A. G. Marin, C. G. Marirrodriga, R. Gerndt, L. Gesa, F. Gilbert, D. Giardini, C. Grimani, A. Grynagier, B. Guillaume, F. Guzmán, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, J. Hough, D. Hoyland, M. Hueller, J. Huesler, O. Jennrich, O. Jennrich, P. Jetzer, B. Johlander, C. Killow, X. Llamas, I. Lloro, A. Lobo, R. Maarschalkerweerd, S. Madden, D. Mance, I. Mateos, P. W. McNamara, J. Mendes, E. Mitchell, A. Monsky, D. Nicolini, D. Nicolodi, M. Nofrarias, F. Pedersen, M. Perreur-Lloyd, A. Perreca, E. Plagnol, P. Prat, G. D. Racca, B. Rais, J. Ramos-Castro, J. Reiche, J. A. R. Perez, D. Robertson, H. Rozemeijer, J. Sanjuan, A. Schleicher, M. Schulte, D. Shaul, L. Stagnaro, S. Strandmoe, F. Steier, T. J. Sumner, A. Taylor, D. Texier, C. Trenkel, D. Tombolato, S. Vitale, G. Wanner, H. Ward, S. Waschke, P. Wass, W. J. Weber, and P. Zweifel, “LISA Pathfinder: mission and status,” *Classical and Quantum Gravity* **28**, 094001 (2011).

⁵F. Antonucci, M. Armano, H. Audley, G. Auger, M. Benedetti, P. Binetruy, J. Bogenstahl, D. Bortoluzzi, P. Bosetti, N. Brandt, M. Caleno, P. Cañizares, A. Cavalleri, M. Cesa, M. Chmeissani, A. Conchillo, G. Congedo, I. Cristofolini, M. Cruise, K. Danzmann, F. De Marchi, M. Diaz-Aguilo, I. Diepholz, G. Dixon, R. Dolesi, N. Dunbar, J. Fauste, L. Ferraioli, V. Ferrone, W. Fichter, E. Fitzsimons, M. Freschi, A. G. Marin, C. G. Marirrodriga, R. Gerndt, L. Gesa, F. Gilbert, D. Giardini, C. Grimani, A. Grynagier, B. Guillaume, F. Guzmán, I. Harrison, G. Heinzel, V. Hernández, M. Hewitson, D. Hollington, J. Hough, D. Hoyland, M. Hueller, J. Huesler, O. Jennrich, P. Jetzer, B. Johlander, N. Karnesis, C. Killow, X. Llamas, I. Lloro, A. Lobo, R. Maarschalkerweerd, S. Madden, D. Mance, I. Mateos, P. W. McNamara, J. Mendes, E. Mitchell, A. Monsky, D. Nicolini, D. Nicolodi, M. Nofrarias, F. Pedersen, M. Perreur-Lloyd, E. Plagnol, P. Prat, G. D. Racca, J. Ramos-Castro, J. Reiche, J. A. R. Perez, D. Robertson, H. Rozemeijer, J. Sanjuan, A. Schleicher, M. Schulte, D. Shaul, L. Stagnaro, S. Strandmoe, F. Steier, T. J. Sumner, A. Taylor, D. Texier, C. Trenkel, H.-B. Tu, S. Vitale, G. Wanner, H. Ward, S. Waschke, P. Wass, W. J. Weber, T. Ziegler, and P. Zweifel, “The LISA Pathfinder mission,” *Classical and Quantum Gravity* **29**, 124014 (2012).

⁶M. Armano, H. Audley, G. Auger, J. Baird, M. Bassan, P. Binetruy, M. Born, D. Bortoluzzi, N. Brandt, M. Caleno, L. Carbone, A. Cavalleri, A. Cesarini, G. Ciani, G. Congedo, A. Cruise, K. Danzmann, M. de Deus Silva, R. De Rosa, M. Diaz-Aguilo, L. Di Fiore, I. Diepholz, G. Dixon, R. Dolesi, N. Dunbar, L. Ferraioli, V. Ferroni, W. Fichter, E. Fitzsimons, R. Flatscher, M. Freschi, A. García Marín, C. García Marirrodriga, R. Gerndt, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, F. Guzmán, A. Grado, C. Grimani, A. Grynagier, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, V. Ferraioli, W. Fuchler, H. Inchauspe, O. Jennrich, P. Jetzer, U. Johann, B. Johlander, N. Karnesis, B. Kaune, N. Korsakova, C. Killow, J. Lobo, I. Lloro, L. Liu, J. López-Zaragoza, R. Maarschalkerweerd, D. Mance, V. Martín, L. Martín-Polo, J. Martino, F. Martin-Portuqueras, S. Madden, I. Mateos, P. McNamara, J. Mendes, L. Mendes, A. Monsky, D. Nicolodi, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, P. Prat, U. Ragnit, B. Raïs, J. Ramos-Castro, J. Reiche, D. Robertson, H. Roze-

- meijer, F. Rivas, G. Russano, J. Sanjuán, P. Sarra, A. Schleicher, D. Shaul, J. Slutsky, C. Sopaerta, R. Stanga, F. Steier, T. Sumner, D. Texier, J. Thorpe, C. Trenkel, M. Tröbs, H. Tu, D. Vetrugno, S. Vitale, V. Wand, G. Wanner, H. Ward, C. Warren, P. Wass, D. Wealthy, W. Weber, L. Wissel, A. Wittchen, A. Zambotti, C. Zanoni, T. Ziegler, and P. Zweifel, “Sub-Femto- g Free Fall for Space-Based Gravitational Wave Observatories: LISA Pathfinder Results,” *Physical Review Letters* **116**, 231101 (2016).
- ⁷M. Armano, H. Audley, J. Baird, P. Binetruy, M. Born, D. Bortoluzzi, E. Castelli, A. Cavalleri, A. Cesarini, A. Cruise, K. Danzmann, M. de Deus Silva, I. Diepholz, G. Dixon, R. Dolesi, L. Ferraioli, V. Ferroni, E. Fitzsimons, M. Freschi, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, C. Grimaldi, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, D. Hoyland, M. Hueller, H. Inchauspé, O. Jennrich, P. Jetzer, N. Karnesis, B. Kaune, N. Korsakova, C. Killow, J. Lobo, I. Lloro, L. Liu, J. López-Zaragoza, R. Maarschalkerweerd, D. Mance, N. Meshksar, V. Martín, L. Martin-Polo, J. Martino, F. Martin-Porqueras, I. Mateos, P. McNamara, J. Mendes, L. Mendes, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, J. Ramos-Castro, J. Reiche, D. Robertson, F. Rivas, G. Russano, J. Slutsky, C. Sopaerta, T. Sumner, D. Texier, J. Thorpe, D. Vetrugno, S. Vitale, G. Wanner, H. Ward, P. Wass, W. Weber, L. Wissel, A. Wittchen, and P. Zweifel, “Beyond the Required LISA Free-Fall Performance: New LISA Pathfinder Results down to 20μ Hz,” *Physical Review Letters* **120**, 061101 (2018).
- ⁸D. Mance, “Development of electronic system for sensing and actuation of test mass of the inertial sensor LISA,” (2012).
- ⁹M. Armano, H. Audley, G. Auger, J. Baird, M. Bassan, P. Binetruy, M. Born, D. Bortoluzzi, N. Brandt, M. Caleno, A. Cavalleri, A. Cesarini, A. Cruise, K. Danzmann, M. de Deus Silva, R. De Rosa, L. Di Fiore, I. Diepholz, G. Dixon, R. Dolesi, N. Dunbar, L. Ferraioli, V. Ferroni, E. Fitzsimons, R. Flatscher, M. Freschi, C. García Marirrodiga, R. Gerndt, L. Gesa, F. Gibert, D. Giardini, R. Giusteri, A. Grado, C. Grimaldi, J. Grzymisch, I. Harrison, G. Heinzel, M. Hewitson, D. Hollington, D. Hoyland, M. Hueller, H. Inchauspé, O. Jennrich, P. Jetzer, B. Johlander, N. Karnesis, B. Kaune, N. Korsakova, C. Killow, J. Lobo, I. Lloro, L. Liu, J. López-Zaragoza, R. Maarschalkerweerd, D. Mance, V. Martín, L. Martin-Polo, J. Martino, F. Martin-Porqueras, S. Madden, I. Mateos, P. McNamara, J. Mendes, L. Mendes, N. Meshksar, M. Nofrarias, S. Paczkowski, M. Perreur-Lloyd, A. Petiteau, P. Pivato, E. Plagnol, P. Prat, U. Ragnit, J. Ramos-Castro, J. Reiche, D. Robertson, H. Rozemeijer, F. Rivas, G. Russano, P. Sarra, A. Schleicher, J. Slutsky, C. Sopaerta, R. Stanga, T. Sumner, D. Texier, J. Thorpe, C. Trenkel, M. Tröbs, D. Vetrugno, S. Vitale, G. Wanner, H. Ward, P. Wass, D. Wealthy, W. Weber, L. Wissel, A. Wittchen, A. Zambotti, C. Zanoni, T. Ziegler, and P. Zweifel, “Capacitive sensing of test mass motion with nanometer precision over millimeter-wide sensing gaps for space-borne gravitational reference sensors,” *Physical Review D* **96**, 062004 (2017).
- ¹⁰W. J. Weber, D. Bortoluzzi, A. Cavalleri, L. Carbone, M. D. Lio, R. Dolesi, G. Fontana, C. D. Hoyle, M. Hueller, and S. Vitale, “Position sensors for flight testing of LISA drag-free control,” in *Gravitational-Wave Detection*, Vol. 4856, edited by P. Saulson and A. M. Cruise, International Society for Optics and Photonics (SPIE, 2003) pp. 31 – 42.
- ¹¹R. Dolesi, D. Bortoluzzi, P. Bosetti, L. Carbone, A. Cavalleri, I. Cristofolini, M. DaLio, G. Fontana, V. Fontanari, B. Foulon, C. D. Hoyle, M. Hueller, F. Nappo, P. Sarra, D. N. A. Shaul, T. Sumner, W. J. Weber, and S. Vitale, “Gravitational sensor for LISA and its technology demonstration mission,” *Classical and Quantum Gravity* **20**, S99–S108 (2003).
- ¹²W. J. Weber, “Nano-Newton electrostatic actuators for femto-Newton small-force measurements: performance of the LISA Pathfinder electrostatic force actuation system,” (To be published).
- ¹³“Analog Devices,” www.analog.com (2000 - 2007).
- ¹⁴N. Meshksar, L. Ferraioli, D. Mance, and J. ten Pierick, “Electrostatic Actuation Accuracy for the Laser Interferometer Space Antenna (LISA),” (To be published).
- ¹⁵MATLAB, “version 9.3.0.713579 (R2017b),” Natick, Massachusetts (2017).