# "Faster GOW0 implementation for more accurate material design" 

Laflamme Janssen, Jonathan ; Rousseau, Bruno ; Côté, Michel


#### Abstract

Density-functional theory (DFT) is currently the ab initio method most widely used to predict electronic energy levels of new materials. However, approxima- tions intrinsic to the theory limit the accuracy of calculated energy levels to about 0.5 eV . The GOWO approach is an alternate ab initio method that provides an enhanced precision (about 0.05 eV ). However, its computational cost is currently prohibitive for systems with more than a few tens of electrons, thus limiting its use in the simulation and design of technologically relevant materials. This limitation of current GOWO implementations can be traced to two bottle- necks : the need to invert a large matrix (the dielectric matrix) and the need to carry out summations over a large number of electronic states (conduction states). The first bottleneck is caused by the choice of the basis in which the dielectric matrix is represented : traditional GOW0 implementations use a plane wave basis, which needs to be relatively large to $\mathrm{p} .$.


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## Faster $\mathrm{G}_{0} \mathrm{~W}_{0}$ implementation for more accurate material design

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## Design of polymers



## Design of polymers



- Gap: optimal for solar spectrum


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## - HOMO :

- higher than $\mathrm{C}_{60}$
(for good charge transfer)


## Design of polymers



- Gap: optimal for solar spectrum


## - HOMO :

- higher than $\mathrm{C}_{60}$
(for good charge transfer)
- low enough for good Voc and air stability


## DFT to the rescue?

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## DFT to the rescue?

HOMO


- Allows to select the top 10-20\% candidates polymers


## The $\mathrm{G}_{0} \mathrm{~W}_{0}$ : performance

## $\mathrm{G}_{0} \mathrm{~W}_{0}$ :

$\left(T+V_{e x t}+V_{H}+V_{x c}\right)\left|\phi_{n}\right\rangle=\varepsilon_{n}\left|\phi_{n}\right\rangle$
$G=\sum_{n=1}^{\infty} \frac{\left|\phi_{n}\right\rangle\left\langle\phi_{n}\right|}{\omega-\varepsilon_{n}}$
$P=-i G G$
$\epsilon=1-v P$
$W=\epsilon^{-1} v$

$\Sigma=i G W$

## LDA and $\mathrm{G}_{0} \mathrm{~W}_{0}$ : precision

$\epsilon_{n} \approx \varepsilon_{n}+\left\langle\phi_{n}\right| \Sigma-V_{x c}\left|\phi_{n}\right\rangle$

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## LDA and $\mathrm{G}_{0} \mathrm{~W}_{0}$ : precision

## The $\mathrm{G}_{0} \mathrm{~W}_{0}$ : bottlenecks

- Why $\mathrm{G}_{0} \mathrm{~W}_{0}$ so expensive ?

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## - $\mathrm{N}_{\mathrm{c}} \sim 10 \mathrm{~N}_{\mathrm{v}}$ to $100 \mathrm{~N}_{\mathrm{v}}$ for $\epsilon_{\mathrm{n}}$ at $\pm 0.05 \mathrm{eV}$

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- $N_{c} \sim 10 N_{v}$ to $100 N_{v}$ for $\epsilon_{n}$ at $\pm 0.05 \mathrm{eV}$
- inversion of $\epsilon \Rightarrow N^{3}$ operation ( $N=$ basis size)


## The case of antracene



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$$
N_{v}=33
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- $\epsilon$ matrix $\sim 7000 \times 7000$ planewaves


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- $\epsilon$ matrix $\sim 7000 \times 7000$ planewaves
$\Rightarrow 1 \mathrm{~Gb}$ of RAM usage to store $\epsilon$


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$\Rightarrow 10 \mathrm{~Gb}$ of RAM usage to store $\{|c\rangle\}$
$\Rightarrow$ 100's hours of CPU time to obtain $\{|c\rangle\}$
- $\epsilon$ matrix $\sim 7000 \times 7000$ planewaves
$\Rightarrow 1 \mathrm{~Gb}$ of RAM usage to store $\epsilon$
$\Rightarrow$ 10's hours of CPU time to $\epsilon^{-1}$


## The plan

$$
\begin{aligned}
P & =-i G G \\
\epsilon & =1-v P \\
W & =\epsilon^{-1} v \\
\Sigma & =i G W
\end{aligned}
$$

$N_{c}$ : number of conduction states

$$
G=\sum_{n=1}^{m-v_{0}} \frac{|n\rangle\langle n|}{\omega-\varepsilon_{n}}
$$

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- summations $\rightarrow$ Sternheimer's equations


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- summations $\rightarrow$ Sternheimer's equations


## The plan

$$
\begin{array}{rlrl}
\text { 1) } P & =-i G G \\
\epsilon & =1-v P \\
W & =\epsilon^{-1}, & G=\sum_{n=1} \frac{|n\rangle\langle n|}{\omega-\varepsilon_{n}} \\
\text { 2) } \begin{aligned}
& \text { number of conduction states } \\
& \sigma_{\sim} \sim N_{c} \\
&=i G W
\end{aligned} &
\end{array}
$$

- summations $\rightarrow$ Sternheimer's equations


## The plan

$$
\begin{aligned}
\text { 1) } \begin{aligned}
P & =-i G G \\
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\text { 3) } W & =\epsilon^{-1} v \\
\text { 2) } \Sigma & =i G W
\end{aligned} \quad G=\sum_{n=1} \frac{}{\text { No : number of conduction states }} \\
\omega-\varepsilon_{n}
\end{aligned}
$$

- summations $\rightarrow$ Sternheimer's equations
- planewaves basis $\rightarrow$ Lanczos basis


## 1) $P=-i G G$ <br> Sternheimer's equation

$$
\text { - } P|\psi\rangle=\sum_{v}|\nu\rangle\left(\sum_{c}|c\rangle \frac{1}{\omega-\left(\varepsilon_{c}-\varepsilon_{v}\right)}-\frac{1}{\omega+\left(\varepsilon_{c}-\varepsilon_{v}\right)}\langle c|\langle\nu|\right)|\psi\rangle
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- We define :

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\left|\phi_{v}^{-}\right\rangle \equiv \sum_{c} \frac{|c\rangle\langle c|}{\omega-\left(\varepsilon_{c}-\varepsilon_{v}\right)}|\nu\rangle|\psi\rangle
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$$

## 1) $P=-i G G \quad$ Sternheimer's equation

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{ }_{\uparrow}\left|\phi_{v}^{-}\right\rangle \equiv \sum_{c} \frac{|c\rangle\langle c|}{\omega-\left(\varepsilon_{c}-\varepsilon_{v}\right)}|v\rangle|\psi\rangle=\frac{1}{\omega-(H)+\varepsilon_{v}} \widehat{P}|v\rangle|\psi\rangle
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$\left(H-\varepsilon_{v}-\omega\right)\left|\phi_{v}^{-}\right\rangle=-\mathcal{P}_{c}|v\rangle|\psi\rangle$ Sternheimer equation

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$$

$\left(H-\varepsilon_{v}-\omega\right)\left|\phi_{v}^{-}\right\rangle=-\mathcal{P}_{c}|v\rangle|\psi\rangle$ Sternheimer equation

$$
\longrightarrow \text { Solving } A|x\rangle=|b\rangle
$$

## Sternheimer's equation

## 2) $\Sigma=i G W$

 .$\qquad$

2 , $-2$

## 2）$\Sigma=i G W$ <br> Sternheimer＇s equation

－We want：$\epsilon_{m} \approx \varepsilon_{m}+\langle m| \Sigma_{x}+\Sigma_{c}-V_{x c}|m\rangle$

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\langle m| \Sigma_{c}|m\rangle=\langle m| i G W_{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{n} \frac{\left\langle n^{*}\right| \Phi_{m}^{\dagger}\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle}{\omega-\left(\varepsilon_{n}-\varepsilon_{m}\right)}
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\begin{aligned}
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& \omega-\left(\varepsilon_{n}-\varepsilon_{m}\right) \\
&=(\ldots) \sum_{n q} \frac{\left\langle n^{*}\right| \Phi_{m}^{*}|q\rangle\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}\left|n^{*}\right\rangle}{\omega-\left(\varepsilon_{n}-\varepsilon_{m}\right)}
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& =(\ldots) \sum_{\text {(1) }} \frac{\left\langle n^{*}\right| \Phi_{m}^{\dagger}(q\rangle\langle Q)\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle}{\omega-\left(\varepsilon_{n}-\varepsilon_{m}\right)}
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& =(\ldots) \sum_{n q} \frac{\left\langle q^{*}\right| \Phi_{m}^{\dagger}|n\rangle\langle n| \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle}{\omega-\left(\varepsilon_{n}-\varepsilon_{m}\right)}
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\end{aligned}
$$

$$
\begin{aligned}
& =(\ldots) \sum_{n q} \frac{\left\langle q^{*}\right| \Phi_{m}^{\dagger}|n\rangle\langle n| \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle}{\omega-\left(\varepsilon_{0}\right)^{\left.-\varepsilon_{m}\right)}}
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

- We want : $\epsilon_{m} \approx \varepsilon_{m}+\langle m| \Sigma_{x}+\Sigma_{c}-V_{x c}|m\rangle$

$$
\begin{aligned}
&\langle m| \Sigma_{c}|m\rangle=\langle m| i G W_{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{\omega}^{D} \frac{\left\langle n^{*}\right| \Phi_{m}^{\dagger}\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle}{\omega-\left(\varepsilon_{n}-\varepsilon_{m}\right)} \\
&=(\ldots) \sum_{(9)}^{\left\langle n^{*}\right| \Phi_{m}^{\dagger}(q\rangle\langle Q|\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle} \\
& \omega-\left(\varepsilon_{n}-\varepsilon_{m}\right) \\
&=(\ldots) \sum_{n q} \frac{\left\langle q^{*}\right| \Phi_{m}^{\dagger}(n\rangle\langle n) \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle}{\left.\omega-\left(\mathcal{E}_{n}\right)-\varepsilon_{m}\right)} \\
&=(\ldots) \sum_{q}\left\langle q^{*}\right| \Phi_{m}^{\dagger} \frac{1}{\omega-\left(H-\varepsilon_{m}\right)} \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

- We want : $\epsilon_{m} \approx \varepsilon_{m}+\langle m| \Sigma_{x}+\Sigma_{c}-V_{x c}|m\rangle$

$$
\begin{aligned}
&\langle m| \Sigma_{c}|m\rangle=\langle m| i G W_{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{\omega}^{( } \frac{\left\langle n^{*}\right| \Phi_{m}^{\dagger}\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle}{\omega-\left(\varepsilon_{n}-\varepsilon_{m}\right)} \\
&=(\ldots) \sum_{(9)}^{\left\langle n^{*}\right| \Phi_{m}^{\dagger}(q\rangle\langle Q|\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle} \\
& \omega-\left(\varepsilon_{n}-\varepsilon_{m}\right) \\
&=(\ldots) \sum_{n q} \frac{\left\langle q^{*}\right| \Phi_{m}^{\dagger}(n\rangle\langle n) \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle}{\left.\omega-\left(\mathcal{E}_{n}\right)-\varepsilon_{m}\right)} \\
&=(\ldots) \sum_{q}\left\langle q^{*}\right| \Phi_{m}^{\dagger} \frac{1}{\left.\omega-(H)-\varepsilon_{m}\right)} \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

- We want: $\epsilon_{m} \approx \varepsilon_{m}+\langle m| \Sigma_{x}+\Sigma_{c}-V_{x c}|m\rangle$

$$
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&=(\ldots) \sum_{(9)}^{\left\langle n^{*}\right| \Phi_{m}^{\dagger}(q\rangle\langle Q|\left[\epsilon^{-1}-1\right] v \Phi_{m}\left|n^{*}\right\rangle} \\
& \omega-\left(\varepsilon_{n}-\varepsilon_{m}\right) \\
&=(\ldots) \sum_{n q} \frac{\left\langle q^{*}\right| \Phi_{m}^{\dagger}(n\rangle\langle n) \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle}{\left.\omega-\left(\varepsilon_{n}\right)-\varepsilon_{m}\right)} \\
&=(\ldots) \sum_{q}\left\langle q^{*}\right| \Phi_{m}^{\dagger} \frac{1}{\left.\omega-(H)-\varepsilon_{m}\right)} \Phi_{m} v^{T}\left[\epsilon^{-T}-1\right](\omega)\left|q^{*}\right\rangle
\end{aligned}
$$

## Sternheimer's equation

$\square$

## 2) $\Sigma=i G W$

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## 2) $\Sigma=i G W$ <br> Sternheimer's equation

$$
\langle m| \Sigma_{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}^{T} \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{+T}|q\rangle
$$

## 2) $\Sigma=i G W$ <br> Sternheimer's equation

$$
\begin{aligned}
\langle m| \Sigma_{c}|m\rangle & =\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \boldsymbol{\Phi}_{m}^{T} \frac{1}{\omega-H^{T}+\varepsilon_{m}} \boldsymbol{\Phi}_{m}^{+T}|q\rangle \\
& =\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \Phi_{m}^{T}\left|x_{q}\right\rangle
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

$$
\begin{aligned}
\langle m| \Sigma_{c}|m\rangle & =\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \Phi_{m}^{T} \frac{1}{\omega-H^{T}+\varepsilon_{m}} \boldsymbol{\Phi}_{m}^{i T}|q\rangle \\
& =\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \Phi_{m}^{T}(\mathbb{Q})
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

$$
\begin{gathered}
\langle m| \Sigma_{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}^{T} \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{+T}|q\rangle \\
=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}^{T}|Q\rangle \\
\left|x_{q}\right\rangle \equiv \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{+T}|q\rangle
\end{gathered}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

$$
\begin{aligned}
\langle m| \Sigma_{c}|m\rangle= & \frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \Phi_{m}^{T} \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{\leftarrow T}|q\rangle \\
= & \left.\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \Phi_{m}^{T} \right\rvert\,(Q) \\
& \uparrow \quad\left|x_{q}\right\rangle \equiv \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{+T}|q\rangle
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

$$
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= & \left.\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}^{T} \right\rvert\,(Q) \\
& \uparrow \quad\left|x_{q}\right\rangle \equiv \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{+T}|q\rangle
\end{aligned}
$$

$$
\left(\omega-H^{T}+\varepsilon_{m}\right)\left|x_{q}\right\rangle=\Phi_{m}^{\dagger T}|q\rangle \quad \text { Sternheimer equation }
$$

## 2) $\Sigma=i G W \quad$ Sternheimer's equation

$$
\begin{aligned}
\langle m| \Sigma_{c}|m\rangle= & \frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] v \Phi_{m}^{T} \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{\dagger T}|q\rangle \\
= & \left.\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}^{T} \right\rvert\,(Q) \\
& \uparrow \quad\left|x_{q}\right\rangle \equiv \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{\dagger T}|q\rangle
\end{aligned}
$$

$$
\left(\omega-H^{T}+\varepsilon_{m}\right)\left|x_{q}\right\rangle=\Phi_{m}^{\dagger T}|q\rangle \quad \text { Sternheimer equation }
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## 2) $\Sigma=i G W \quad$ Sternheimer's equation

$$
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= & \left.\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right] \nu \Phi_{m}^{T} \right\rvert\,(Q) \\
& \uparrow \quad\left|x_{q}\right\rangle \equiv \frac{1}{\omega-H^{T}+\varepsilon_{m}} \Phi_{m}^{+T}|q\rangle
\end{aligned}
$$

$$
\left(\omega-H^{T}+\varepsilon_{m}\right)\left|x_{q}\right\rangle=\Phi_{m}^{\dagger T}|q\rangle \quad \text { Sternheimer equation }
$$

 Solving $A|x\rangle=|b\rangle$

## 2）$\Sigma=i G W \quad$ Lanczos algorithm

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$$
\square
$$

$\qquad$
$\qquad$
$\qquad$
$\qquad$
$\qquad$

$$
\begin{aligned}
& \langle m| \Sigma|m\rangle=\frac{i}{2 \pi} \oint^{*} d \omega \sum_{q}\left\langle q \mid\left[E^{-1}-1\right](\ldots) q\right\rangle \\
& \text { 號 } \\
& \square
\end{aligned}
$$

## 2) $\Sigma=i G W \quad$ Lanczos algorithm

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right](\ldots)|q\rangle
$$

- Need a basis: $\{|q\rangle\}$


## 2) $\Sigma=i G W \quad$ Lanczos algorithm

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right](\ldots)|q\rangle
$$

- Need a basis: $\{|q\rangle\}$
- The ideal basis:

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right](\ldots)|q\rangle
$$

- Need a basis: $\{|q\rangle\}$
- The ideal basis:
- Is small, e.g.
$\Rightarrow$ NOT planewaves


## 2) $\Sigma=i G W$

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right](\ldots)|q\rangle
$$

- Need a basis: $\{|q\rangle\}$
- The ideal basis:
- Is small, e.g.
$\Rightarrow$ NOT planewaves
- Is easy to compute
$\Rightarrow$ NOT $\{|n\rangle\}$

2) $\Sigma=i G W$

## Lanczos algorithm

5
2) $\Sigma=i G W$

Lanczos algorithm

- Idea :


## 2) $\Sigma=i G W$

- Idea :
- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$


## 2) $\Sigma=i G W$

-Idea :

- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$
-Take (any) vector: $|\psi\rangle$


## 2) $\Sigma=i G W$

- Idea :
- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$
-Take (any) vector: $|\psi\rangle$
- Build : $\left\{|\psi\rangle,(v P)|\psi\rangle,(v P)^{2}|\psi\rangle, \ldots\right\}$
$\rightarrow$ biggest eigenvalues pop out
-Idea :
- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$
-Take (any) vector: $|\psi\rangle$
- Build : $\left\{|\psi\rangle,(v P)|\psi\rangle,(v P)^{2}|\psi\rangle, \ldots\right\}$
$\rightarrow$ biggest eigenvalues pop out
- Orthonormalize : $\{|q\rangle\}$
- Idea :
- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$
-Take (any) vector: $|\psi\rangle$
- Build : $\left\{|\psi\rangle,(v P)|\psi\rangle,(v P)^{2}|\psi\rangle, \ldots\right\}$
$\rightarrow$ biggest eigenvalues pop out
- Orthonormalize : $\{|q\rangle\}$
- Lanczos procedure: same $\{|q\rangle\}$
- Idea :
- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$
-Take (any) vector: $|\psi\rangle$
- Build : $\left\{|\psi\rangle,(v P)|\psi\rangle,(v P)^{2}|\psi\rangle, \ldots\right\}$
$\rightarrow$ biggest eigenvalues pop out
- Orthonormalize : $\{|q\rangle\}$
- Lanczos procedure: same $\{|q\rangle\}$
- Don't pay all orthogonalization
-Idea :
- efficient sampling of big eigenvalues of : $\epsilon^{-1}-1 \approx 1-\epsilon=v P$
-Take (any) vector: $|\psi\rangle$
- Build : $\left\{|\psi\rangle,(v P)|\psi\rangle,(v P)^{2}|\psi\rangle, \ldots\right\}$
$\rightarrow$ biggest eigenvalues pop out
- Orthonormalize : $\{|q\rangle\}$
- Lanczos procedure: same $\{|q\rangle\}$
- Don't pay all orthogonalization
- Obtain $\epsilon$ for free



## 3) $W=\epsilon^{-1} v$ <br> Lanczos algorithm

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right](\ldots)|q\rangle
$$

## 3) $W=\epsilon^{-1} v$ <br> Lanczos algorithm

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|\left[\epsilon^{-1}-1\right](\ldots)
$$

## 3) $W=\epsilon^{-1} v$ <br> Lanczos algorithm

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}(q)\left[e^{\Theta 1}-1\right](\ldots)(q)
$$

## 3) $W=\epsilon^{-1} v \quad$ Lanczos algorithm

$$
\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\left(q \mid\left[e^{\Theta}-1\right](\ldots)(q)\right.
$$

- We have a small $\{|q\rangle\}$


## 3) $W=\epsilon^{-1} v \quad$ Lanczos algorithm

$\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}(q)\left[\right.$ e- $\left.^{(1)}-1\right](\ldots)(q)$ Benzene


## 3) $W=\epsilon^{-1} v \quad$ Lanczos algorithm

$$
\left.\langle m| \Sigma^{c}|m\rangle=\frac{i}{2 \pi} \int_{-\infty}^{+\infty} d \omega \sum_{q}\langle q|[-1]-1\right](\ldots) \text { Benzene }
$$



## Performance

## Performance



$$
A|x\rangle=|b\rangle
$$

$A|x\rangle=|b\rangle$
[197]

$$
A|x\rangle=|b\rangle
$$

$$
\{|q\rangle\}
$$

Faster?

## Performance



## Performance



## Performance



## Preliminary results

## Preliminary results

- Working in reals systems?


## Preliminary results

- Working in reals systems?

> | > | HOMO (eV) |  |  |
| :--- | :---: | :---: | :---: |
| > | LDA | GoW0 | Exp. |
| >  Benzene | -6.51 | -9.22 | -9.30 |
| > Thiophene | -6.05 | -8.94 | -8.85 > |



## Conclusion

## Conclusion

- Bottleneck assessed :


## Conclusion

- Bottleneck assessed :
- no knowledge of conduction states required


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- no inversion of $\epsilon$ in cumbersome basis


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## Conclusion

- Bottleneck assessed :
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- Future work :
- refine DFT calculations for candidate polymers


## Conclusion

- Bottleneck assessed :
- no knowledge of conduction states required
- no inversion of $\epsilon$ in cumbersome basis
- Future work :
- refine DFT calculations for candidate polymers
- interface states: what do they look like?



## Thank you!

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## Thank you!

Gabriel Antonius


Simon Blackburn Hélène Antaya Vincent Gosselin

## Acknowledgments

Michel Côté's group


Bruno Rousseau


## Calcul Québec

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## Thank you!

Nicolas Bérubé



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Bruno Rousseau


## Calcul Q̨uébec

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## 

## Thank you!

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Vincent Gosselin


Photovoltaic ${ }^{\circ}$ Innovation Network We get the Sun

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