



"Accurate effective masses from first principles"

Laflamme Janssen, Jonathan ; Gonze, Xavier

Abstract

The accurate ab initio description of effective masses is of key interest in the design of materials with high mobility. However, up to now, they have been calculated using finite-difference estimation of density functional theory (DFT) electronic band curvatures. To eliminate the numerical noise inherent to finite-difference and obtain an approach that is more suitable for material design using high throughput computing, we develop a method allowing to obtain the curvature of DFT bands using Density-Functional Perturbation Theory (DFPT), taking a change of wavevector as a perturbation. Also, the inclusion of G\$_0\$W\$_0\$ corrections to DFT bands in our method will be presented.

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Accurate effective masses from first principles

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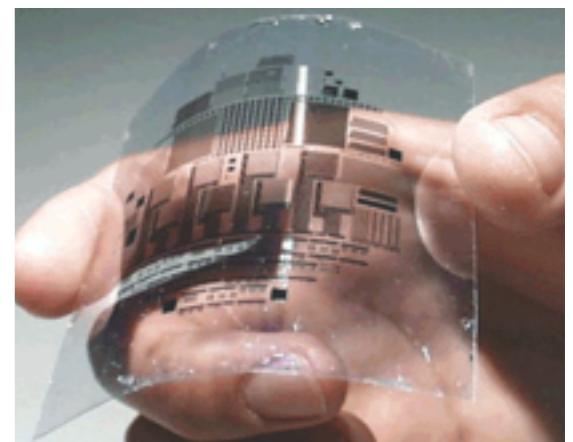
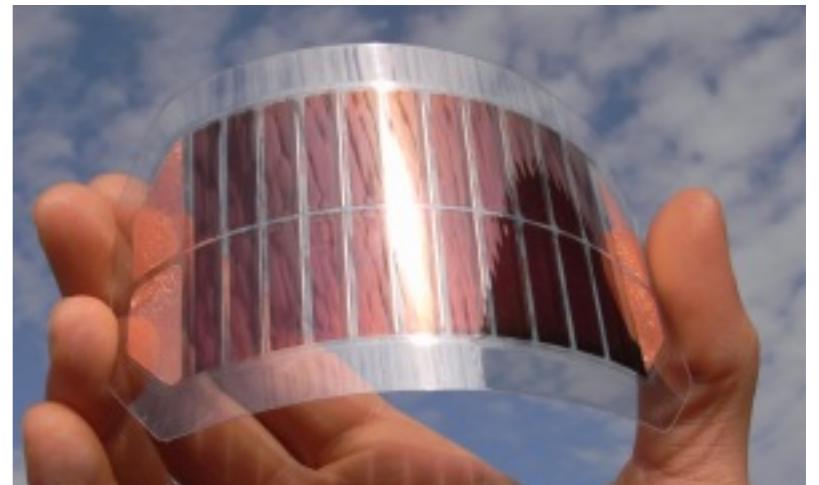
Effective masses

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- design materials with better mobility for electronic applications

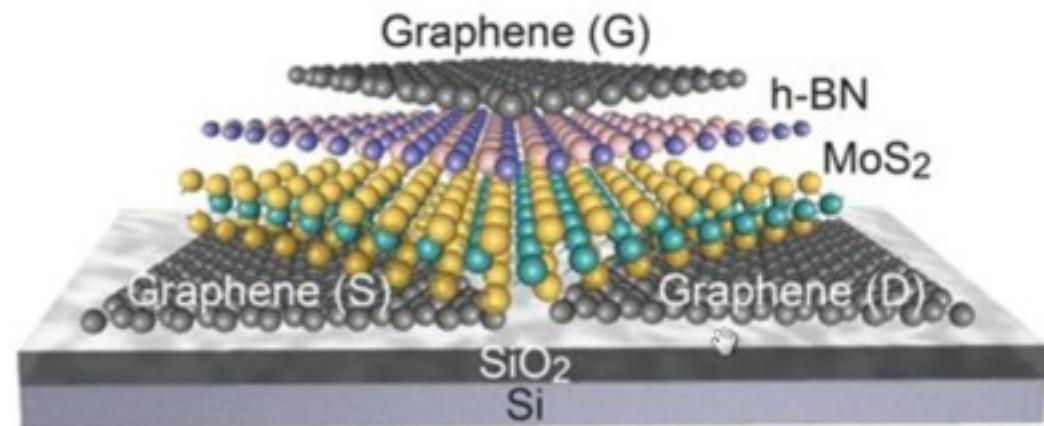
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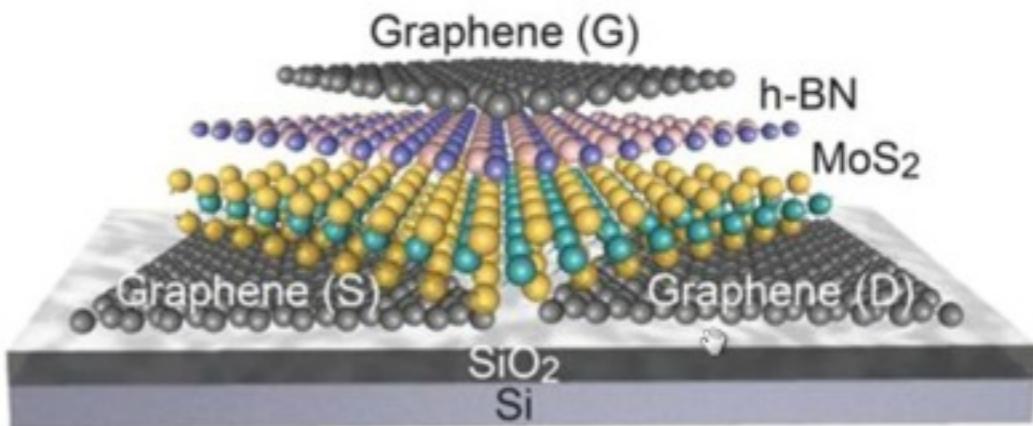
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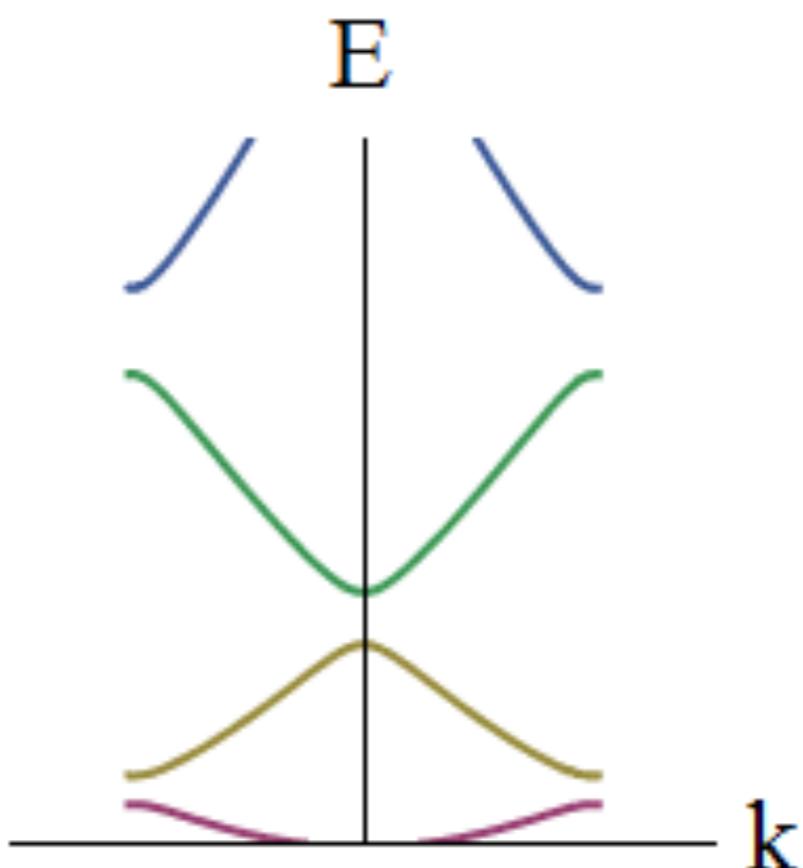
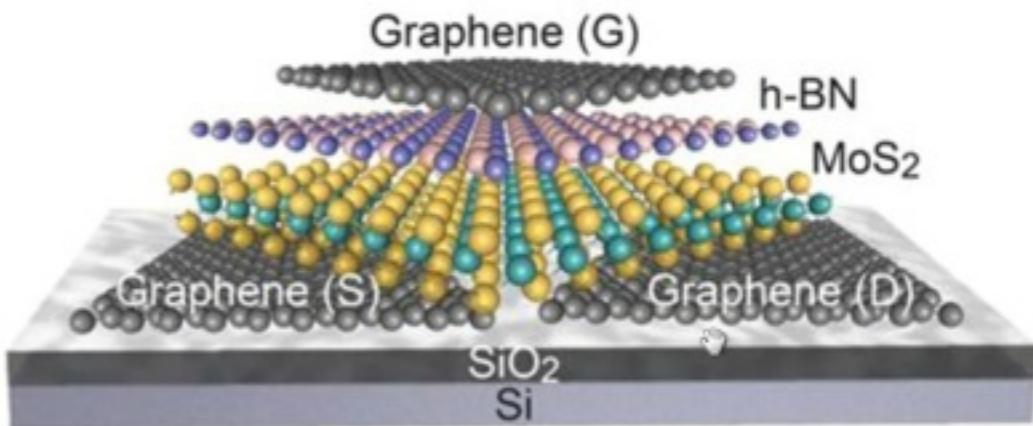
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better μ : both higher conductivity and transparency



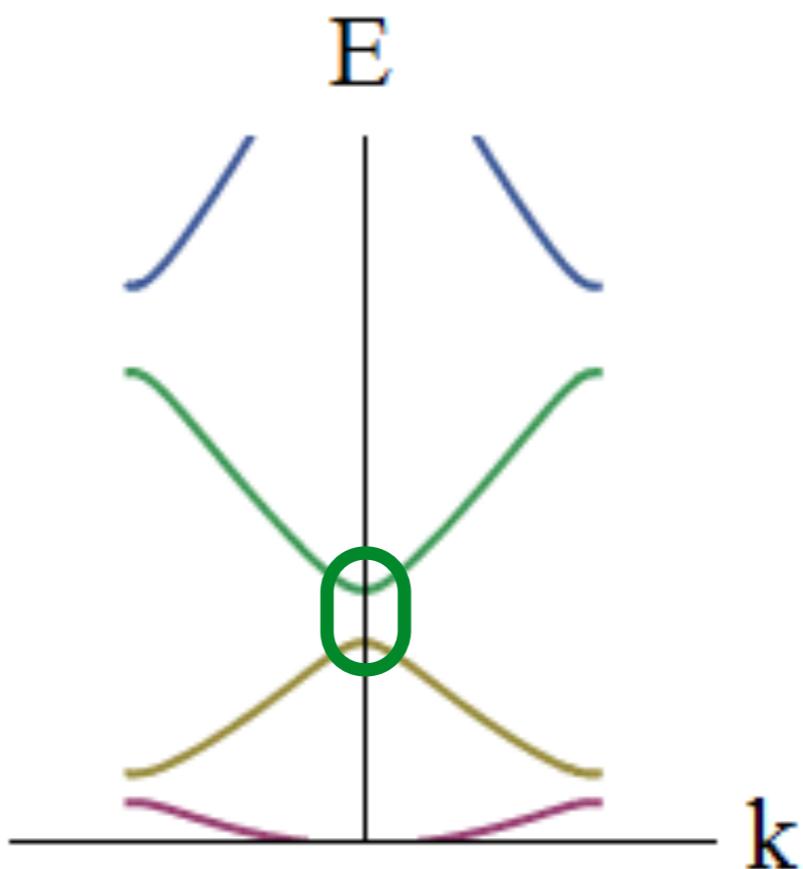
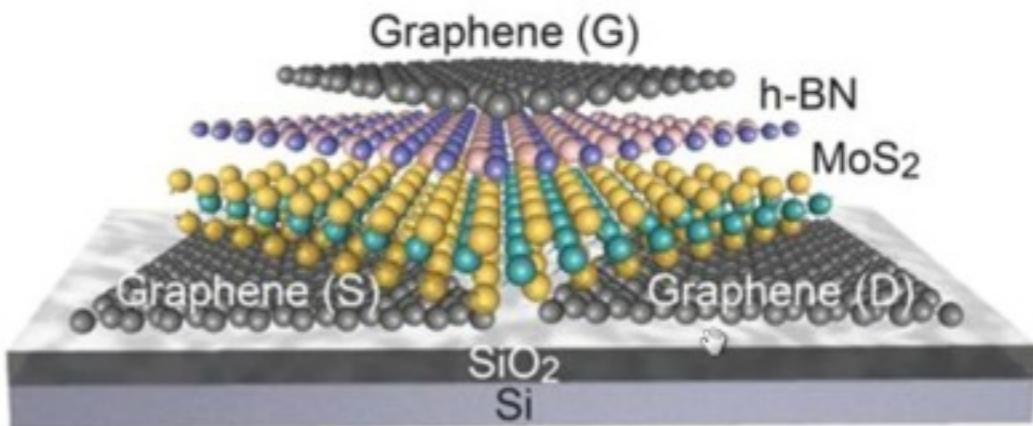
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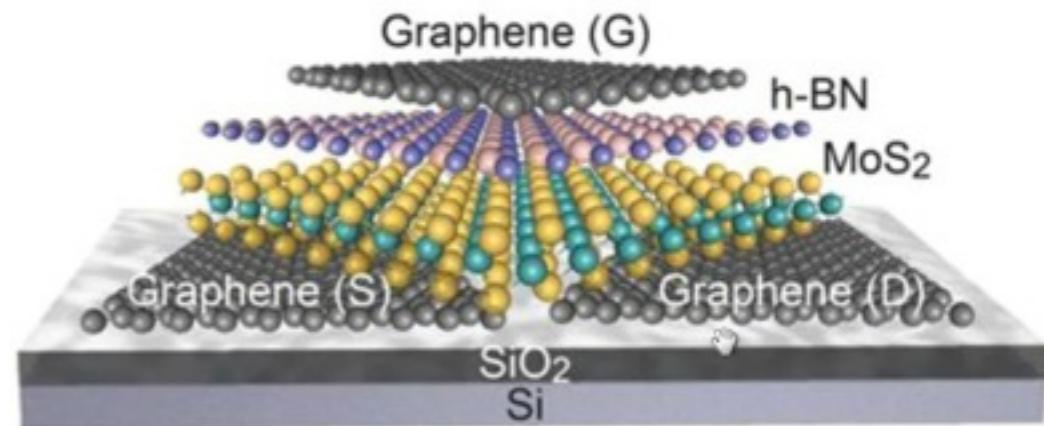
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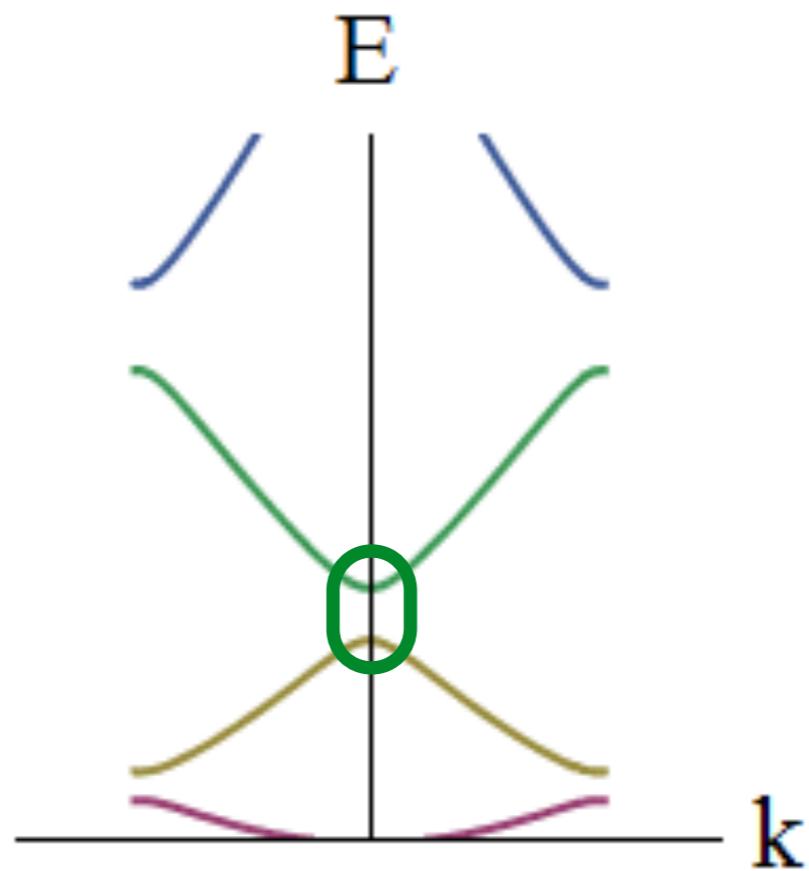


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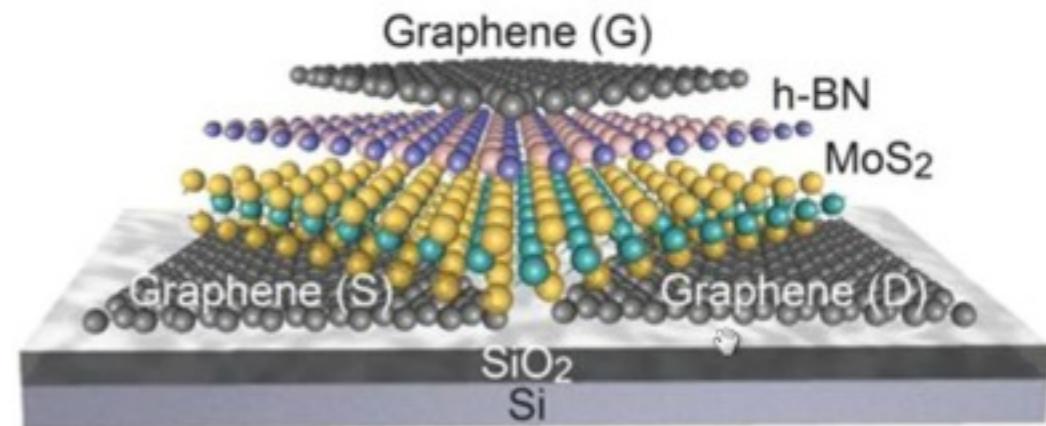


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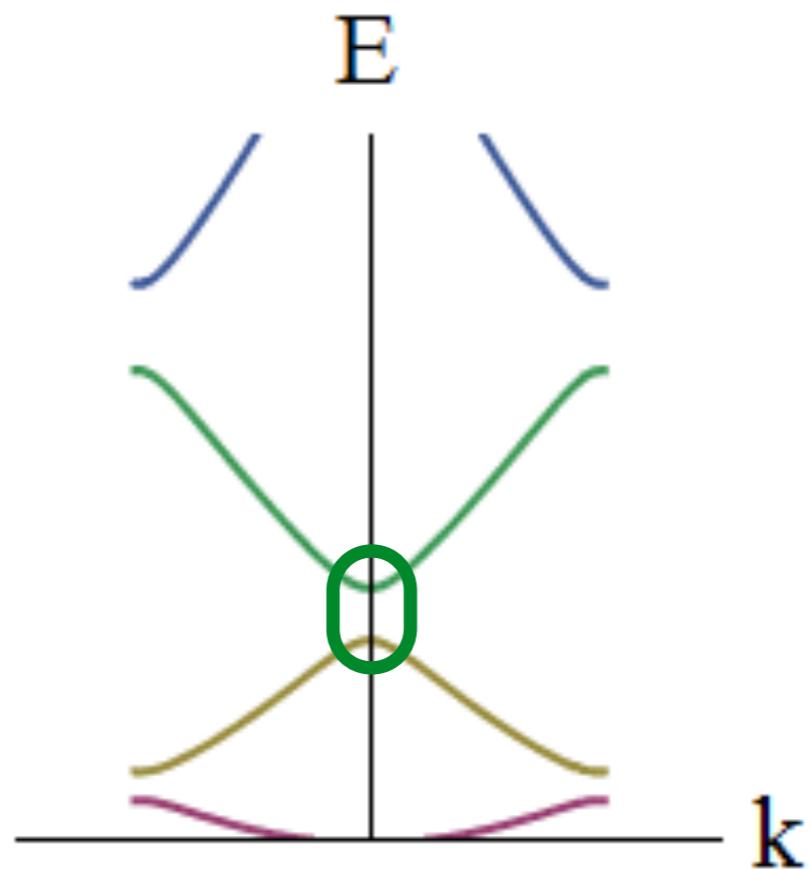


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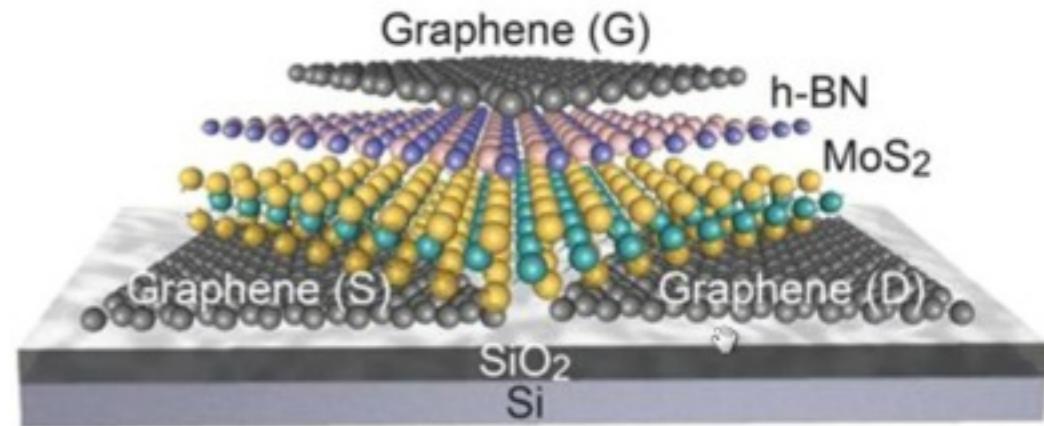


$$\epsilon_{kn} \approx \epsilon_{\Gamma n} + \frac{k^2}{2M} \quad M \neq 1$$



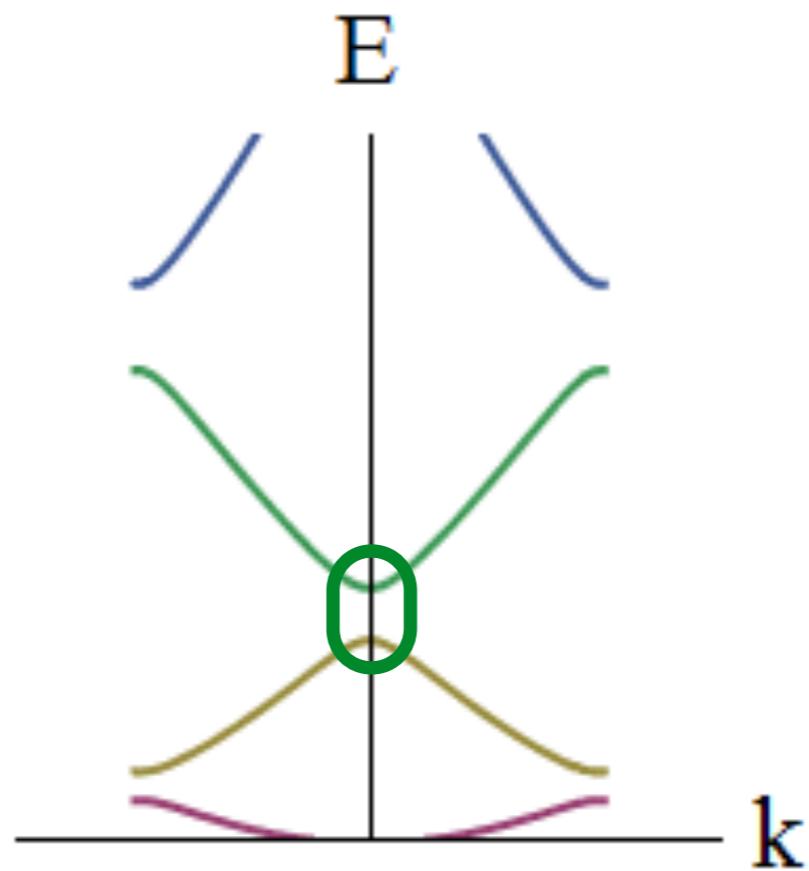
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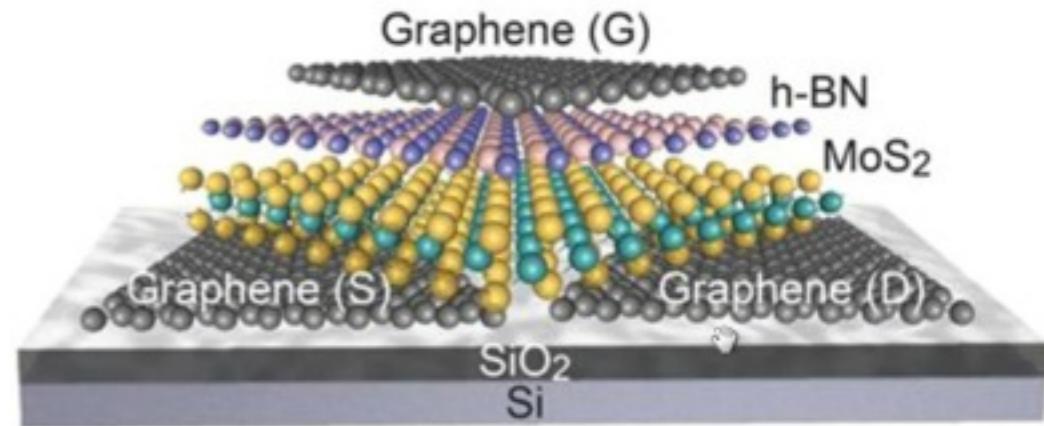
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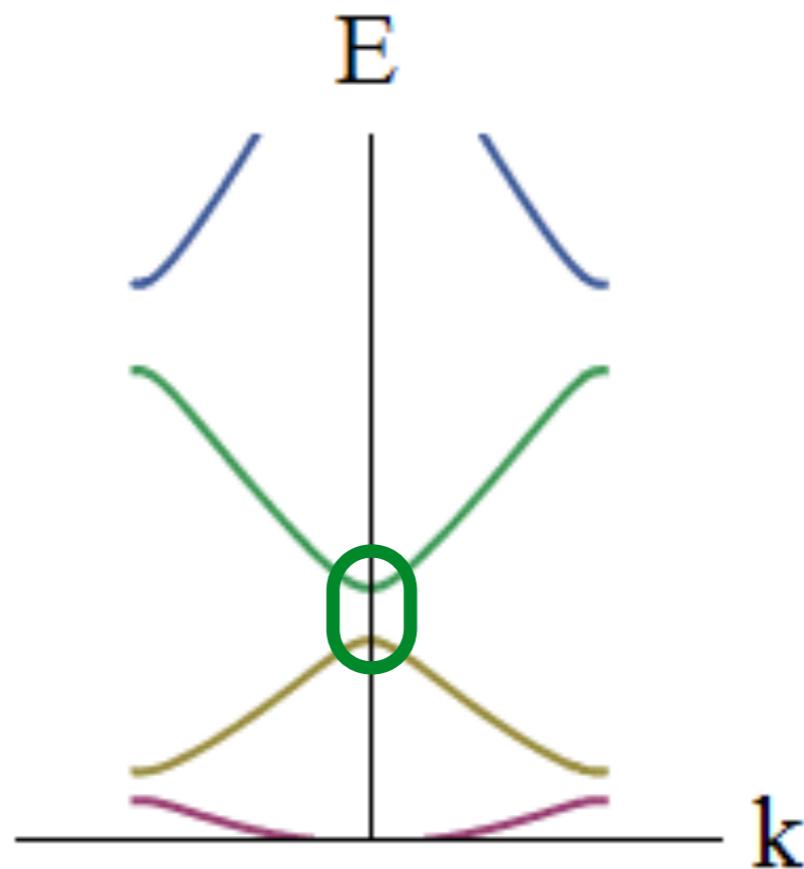
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Effective masses calculation

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DFPT

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Derivatives of \hat{H} up to 2nd order

DFPT

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Derivatives of \hat{H}

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Mecholsky *et al.*, PRB **89** 155131 (2014)

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Results

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Si, 1st band, Γ

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Finite

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0.000	1.161	-0.000
0.000	-0.000	1.161

DFPT

1.161	0.000	0.000
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$$\overline{\Delta} = 10^{-4}$$

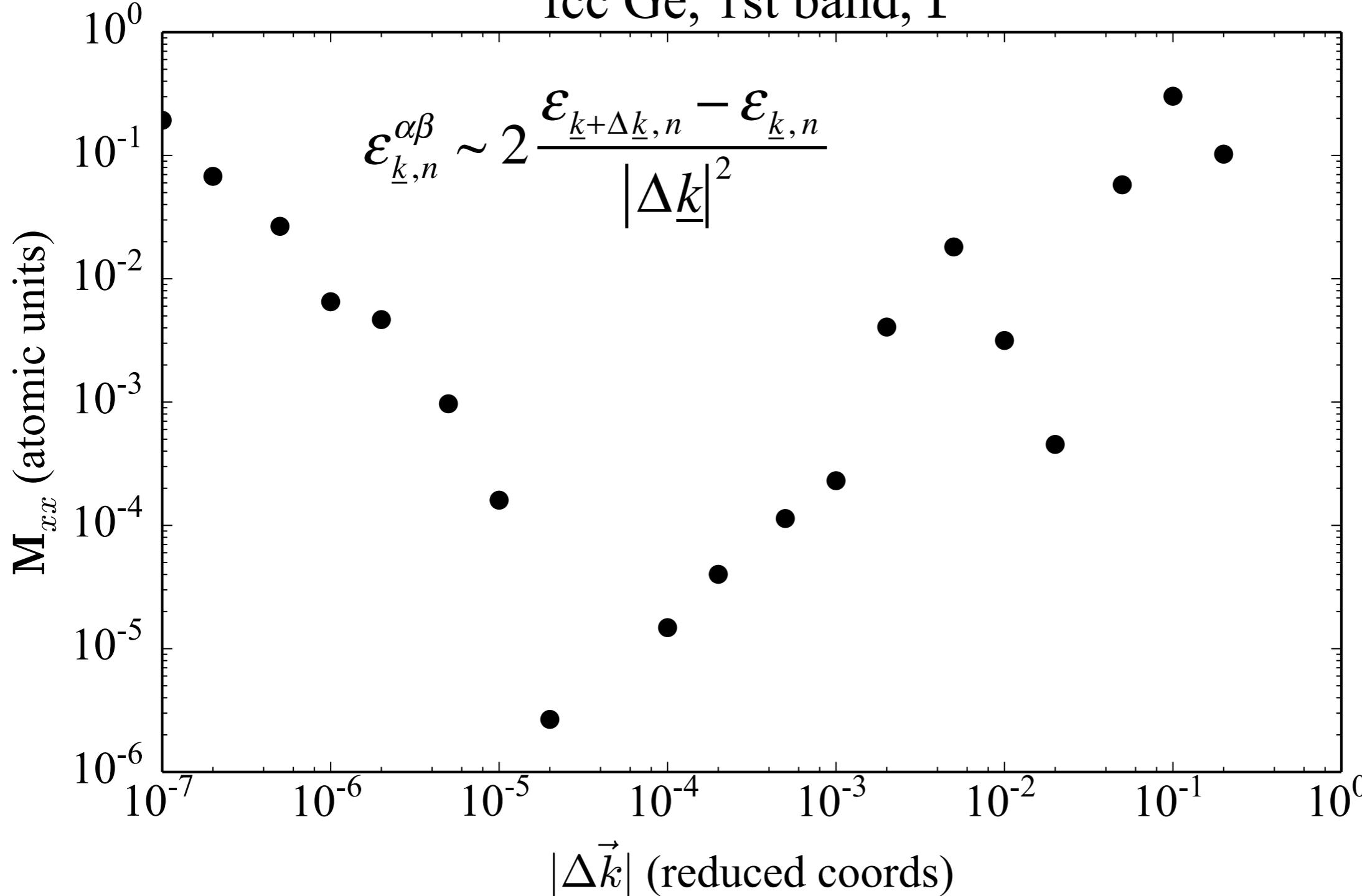
Numerical noise : finite differences

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$$\epsilon_{\underline{k},n}^{\alpha\beta} \sim 2 \frac{\epsilon_{\underline{k}+\Delta\underline{k},n} - \epsilon_{\underline{k},n}}{|\Delta\underline{k}|^2}$$

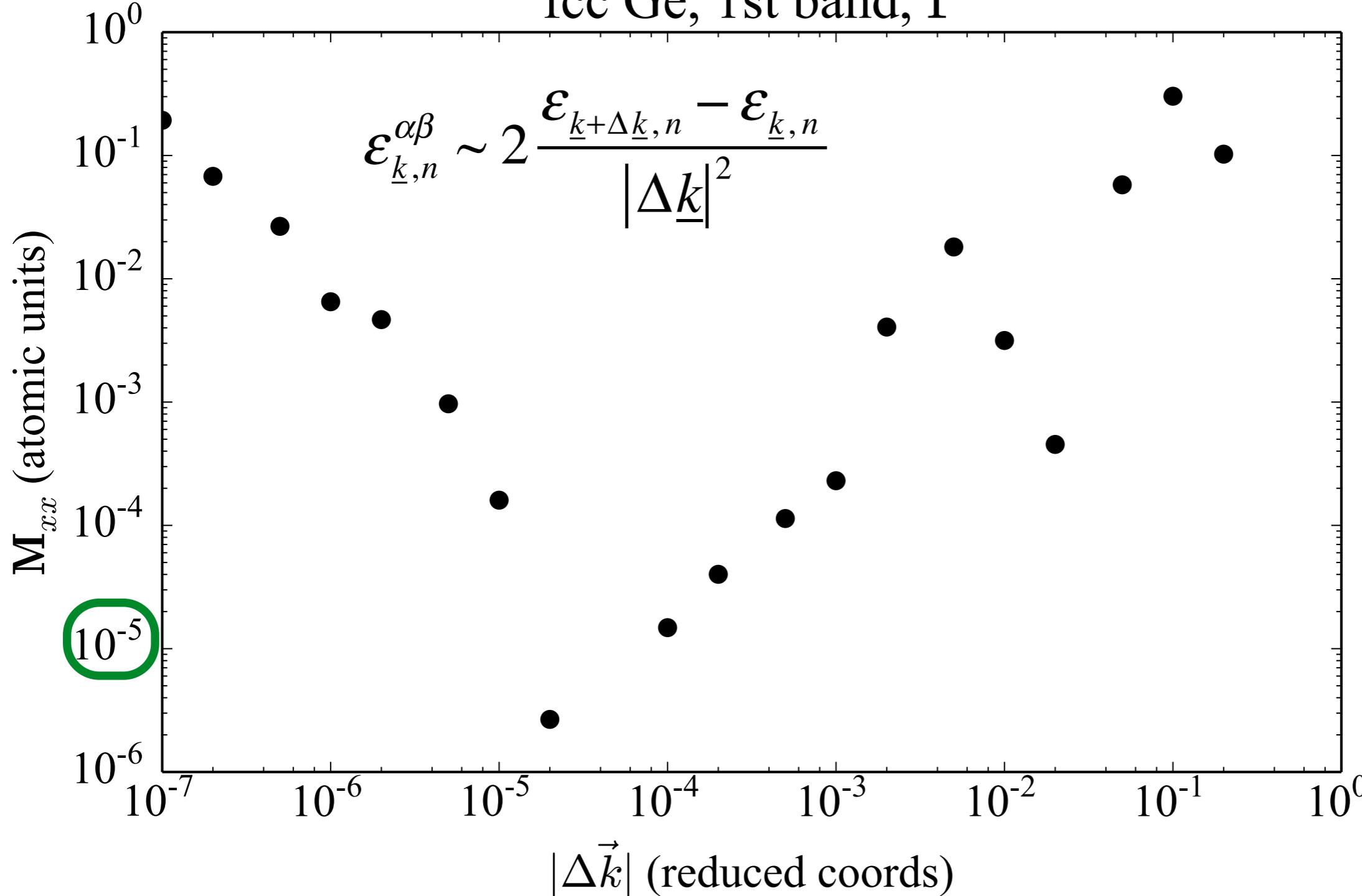
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Convergence study for the effective mass fcc Ge, 1st band, Γ



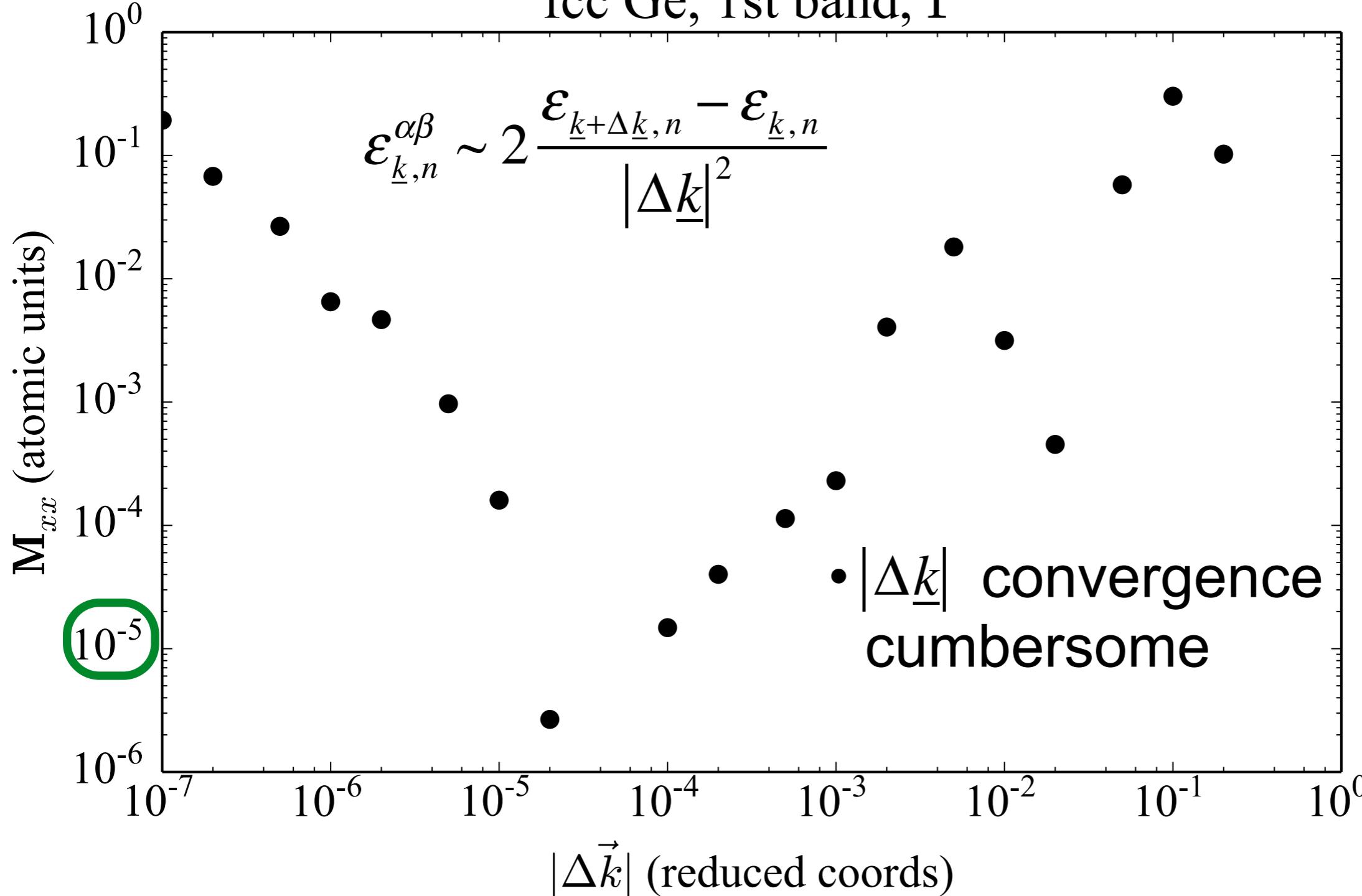
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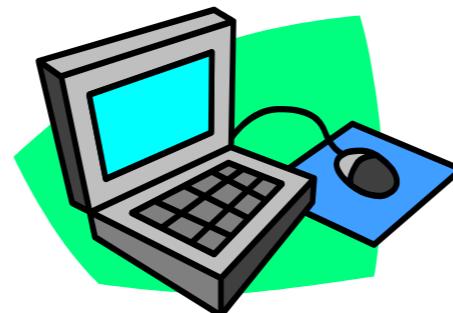
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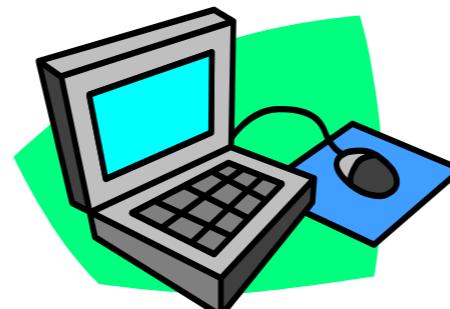
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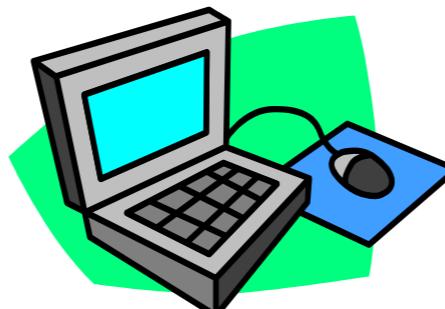


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Thank you!



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$$\frac{M^{G_0W_0}}{M^{\text{DFT}}} = \frac{\cancel{\frac{\partial^2 \varepsilon^{\text{DFT}}}{\partial k^2}}}{\cancel{\frac{\partial^2 \varepsilon^{G_0W_0}}{\partial k^2}}} = \frac{\partial \varepsilon^{\text{DFT}}}{\partial \varepsilon^{G_0W_0}}$$

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$$\frac{M^{G_0W_0}}{M^{\text{DFT}}} = \frac{\cancel{\frac{\partial^2 \varepsilon^{\text{DFT}}}{\partial k^2}}}{\cancel{\frac{\partial^2 \varepsilon^{G_0W_0}}{\partial k^2}}} = \frac{\partial \varepsilon^{\text{DFT}}}{\partial \varepsilon^{G_0W_0}}$$

$$\partial \varepsilon^{G_0W_0} = \partial \varepsilon^{\text{DFT}} + \frac{\partial \langle \hat{\Sigma}(\omega) \rangle}{\partial k} \delta k + \frac{\partial \langle \hat{\Sigma}(\omega) \rangle}{\partial \omega} \delta \varepsilon^{G_0W_0}$$

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Oshikiri et al. PRB **66** 125204 (2002)

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- How general is this apprx?