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# Atomic Layer Deposition of PbS Thin Films at Low Temperatures

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deposition temperatures enable direct ALD of PbS onto a halide perovskite CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> (MAPI) without its decomposition. The stability of MAPI in ambient air is greatly improved by capping with ALD PbS. More generally, these new processes offer valuable alternatives for PbS-based devices, and we hope that this study will inspire more studies on ALD of non-oxides on halide perovskites.

# 1. INTRODUCTION

Lead(II) sulfide (PbS) is one of the oldest known and most wellstudied semiconductors. Depending on whether excess lead or sulfur is present, PbS exhibits either n-type or p-type semiconductivity.<sup>1</sup> The band gap of PbS is narrow (0.4 eV); therefore, PbS absorbs visible and near-mid-infrared light. The exciton Bohr radius in PbS is large (ca. 20 nm) compared to most semiconductors which enables quantum confinement effects to manifest at relatively large dimensions, causing widening of the band gap.<sup>2</sup> These unique properties enable application of PbS thin films in photovoltaics,<sup>3,4</sup> photodetectors,<sup>5,6</sup> transistors,<sup>7</sup> and gas sensors.<sup>8,9</sup> PbS quantum dots demonstrate an even wider range of applications.<sup>7</sup> Atomic layer deposition (ALD) is capable of producing both PbS thin films<sup>10</sup> and quantum dots.<sup>11</sup> Our main interest in ALD of PbS is in using PbS thin films together with halide perovskites.

Halide perovskites are a class of materials with remarkable optoelectronic properties. Perovskites have radically changed the landscape in the photovoltaics research<sup>12</sup> and are studied in a variety of other devices such as X-ray scintillators,<sup>13</sup> photodetectors,<sup>14,15</sup> light-emitting diodes (LEDs),<sup>16</sup> and transistor<sup>17</sup> as well as memory components.<sup>18,19</sup> However, halide perovskites also face stability and scalability issues that hinder their commercialization. Perovskites decompose in ambient air, and scalable deposition methods must be developed for all layers in perovskite-based devices for their production. One approach to tackle these issues is to use ALD for thin film deposition while making these devices. ALD is compatible with roll-to-roll processing and is inherently scalable to large batches and deposition areas. ALD offers uniform and pinhole-free films and ability to conformally coat even the most demanding 3D structures with high aspect ratios. When deposited on top of perovskite, such films encapsulate the perovskite and improve its stability in ambient air.<sup>20</sup> Because of these characteristics, ALD has recently made its way into perovskite photovoltaics.<sup>21</sup> Zardetto et al.,<sup>22</sup> Seo et al.,<sup>23</sup> Brinkmann et al.,<sup>24</sup> and Raiford et al.<sup>25</sup> have given comprehensive reviews on this topic.

A great number of different materials can be deposited with ALD: at least 501 unique compositions according to the database of ALD processes.<sup>26</sup> However, in context of ALD on perovskites, previous studies have focused on a few selected ALD oxides such as  $Al_2O_3$  and  $SnO_x^{22-24}$  This disparity is caused by the relatively high deposition temperatures (100-400 °C) that most ALD processes require. These deposition

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**Figure 1.** (a) Lead precursors used in this study. (b) Overview of ALD processes for PbS with  $H_2S$ . For more details on  $Pb_4O(OBu^{\dagger})_{6}$ ,  $Pb(thd)_{2}$ , and  $Pb(dedtc)_2$  refer to Nykänen et al.<sup>10</sup> and on  $PbEt_4$  to Yun et al.<sup>28</sup> Growth per cycle (GPC) of the PbS film on silicon as (c) a function of deposition temperature, (d) Pb precursor pulse duration, and (e)  $H_2S$  pulse duration. (f) PbS film thickness on silicon as a function of applied deposition cycles. Unless otherwise evident, data are from films deposited with 1000 cycles, 1.0 s precursor pulses, and purge durations.

temperatures are incompatible with the low thermal budget (ca.  $\leq$ 100 °C) of the perovskites. Another incompatibility issue arises from chemical reactions between the ALD precursors and the perovskites.<sup>27</sup> In this work we show that, in addition to the handful of ALD oxide processes, also ALD sulfide processes can be compatible with the perovskites. In particular, PbS can be deposited directly on the CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> (MAPI) perovskite, improves the stability of the perovskite, and has valuable functional properties. We chose PbS because it can exhibit ptype semiconductivity which is scarce among perovskite compatible ALD materials and because PbS electrical as well as optical properties can be tuned. We aimed to achieve proper protective and carrier transport properties in one ALD layer that is compatible with the perovskites. Combining two functions into one ALD layer provides the advantage of simplifying the device structure and its deposition procedure.<sup>29</sup> Very few ALD processes are capable of depositing films with such properties, and p-type semiconductors suitable for hole transporting layers are even more scarce.<sup>22,25</sup> Previous studies have already demonstrated the applicability of PbS films in perovskite photovoltaics.<sup>30,31</sup> We believe that protective, p-type PbS can be used in other p-i-n junction based devices such as LEDs,<sup>1</sup> memory components,<sup>19</sup> and photodetectors.<sup>14,15</sup> Devices based on other structures can still exploit the protective property of PbS provided that the conductivity of the protective layer is not detrimental. PbS films also have another useful property: a series of chemical treatments convert them into MAPI films while retaining the typical advantages of ALD, as demonstrated by Sutherland et al.<sup>32,33</sup> in one of the first studies involving both perovskites and ALD.

Several ALD processes already exist for PbS (Figure 1b).<sup>10,28</sup> These processes operate in a temperature range of 130–400 °C. Such temperatures are incompatible with the halide perovskites. For example, MAPI, the most well-known halide perovskite, decomposes completely at 130 °C in our ALD reactors. Lowering the deposition temperature below 130 °C with the earlier ALD PbS processes is impossible because the lead precursors are either not sufficiently volatile or reactive at lower temperatures. Hence, we decided to look into developing new ALD processes for PbS that would work at lower temperatures. In our previous work on ALD of  $PbI_2^{34}$  we had identified one lead precursor that is volatile and reactive at low temperatures: lead(II) bis[bis(trimethylsilyl)amide], Pb(btsa)<sub>2</sub> (Figure 1a). Another potentially suitable precursor, also a diamide, lead(II)  $rac-N^2,N^3$ -di-tert-butylbutane-2,3-diamide, Pb(dbda) (Figure 1a), was recently reported by Bačić et al.<sup>35</sup> In this work we describe two new ALD processes for PbS using Pb(btsa)<sub>2</sub> and Pb(dbda) as lead precursors and H<sub>2</sub>S as a sulfur precursor. We also demonstrate that these PbS processes are compatible with MAPI films and discuss challenges that we encountered during the scale-up of the Pb(dbda) synthesis.

# 2. RESULTS AND DISCUSSION

2.1. Precursor Characteristics, Synthesis, and Scale-Up. Pb(btsa)<sub>2</sub> and Pb(dbda) (Figure 1a) are diamidoplumbylenes with low melting points (32 and 53 °C, respectively) and high volatility (60 and 40 °C precursor source temperatures in the F120 ALD reactor, respectively). While Pb(btsa)<sub>2</sub> was readily scalable through a facile salt-metathesis reaction, Pb(dbda) had only been synthesized on a preparative scale by transamination of Pb(btsa)<sub>2</sub>.<sup>35</sup> ALD process development required larger quantities of Pb(dbda), and scale-up uncovered deficiencies in the original procedure. We found that residual Pb(btsa)<sub>2</sub> as low as 5% in the crude product mixture after the transamination lowered the melting point mixture to far below room temperature. All attempts to separate Pb(dbda) and Pb(btsa)<sub>2</sub> by recrystallization failed, and fractional sublimation was complicated by their similar volatility. While Pb(dbda) sublimed first, the partial pressure of Pb(btsa)<sub>2</sub> rose when Pb(dbda) was depleted in the sublimation pot; thus,  $Pb(btsa)_2$ began to sublime and mix with crystalline Pb(dbda) on the coldfinger. This caused the crystalline Pb(dbda) to become a



Figure 2. TOF-ERDA composition of ca. 100 nm thick PbS films deposited at different temperatures with (a) Pb(btsa), and (c) Pb(dbda). For more details on impurity contents see Figure S5. Grazing incidence XRD (GIXRD) patterns of PbS films deposited at different temperatures with (b) Pb(btsa)2 and (d) Pb(dbda). Dashed black and dotted gray lines mark PbS and substrate (Si) reflection positions, respectively. 135 °C capped refers to a PbS film capped with ca. 3 nm ALD Al<sub>2</sub>O<sub>3</sub> and measured immediately after the deposition. PbS films were deposited with 1.0 s precursor pulses and purge durations. In (b) and (d) films were deposited with 1000 cycles, except for the 135 °C capped film that was deposited with 2000 cycles.

molten mixture that would drip back into the pot indefinitely. Thus, all  $Pb(btsa)_2$  must either be converted to Pb(dbda) or destroyed by thermolysis for effective purification. Reaction monitoring by <sup>1</sup>H NMR spectroscopy allowed us to scale-up and hasten the preparation: we reduced the reaction time (from ca. 24 to 3 h) by increasing the reaction temperature (ca. 110  $^{\circ}$ C) and using toluene as solvent. Ensuring that no  $Pb(btsa)_2$ remained by reaction monitoring enabled simple Pb(dbda) sublimation. We obtained reduced yield (51% on a 10 g scale), but larger quantities of pure Pb(dbda) could be prepared in a timely manner. Optimization of the kinetic conditions for this reaction currently lies outside the scope of this work but is the subject of ongoing research.

**2.2. PbS Film Growth.** Both Pb(btsa), and Pb(dbda), in combination with H<sub>2</sub>S, produced PbS films at low deposition temperatures (Figure 1c and Figure S1). The deposition temperature ranges that we explored were limited by the precursor source temperatures on the lower end and by the precursor decomposition temperatures on the higher end. The  $Pb(btsa)_2-H_2S$  process has two regions where the GPC is independent of the deposition temperature. In this process the GPC decreases toward higher temperatures. We attribute the GPC decrease to the surface passivation by -SiMe<sub>3</sub> groups, forming from the protonated btsa ligand, that becomes more

effective at higher temperatures.<sup>36</sup> In the Pb(dbda)-H<sub>2</sub>S process, at first, the GPC also decreases with increasing deposition temperature, but at temperatures higher than 95 °C the GPC starts to increase again. This increase is related to the changes in film composition that will be discussed later.

We studied the influence of the process parameters on the film deposition in more detail at the deposition temperatures of 75 and 60 °C with Pb(btsa)<sub>2</sub> and Pb(dbda), respectively. At these temperatures, in both processes, the GPC saturates with respect to the Pb precursor as well as H<sub>2</sub>S pulse durations (Figure 1d,e) and is unaffected by the duration of the nitrogen purge (Figure S2). The thickness of the film is straightforwardly adjustable due to its linear dependence on the number of deposition cycles on both silicon (Figure 1f) and sapphire (Figure S3). The films are surprisingly uniform for such low temperatures and short pulse and purge durations. The thickness nonuniformities over the 5  $\times$  $5 \text{ cm}^2$  substrate are 3.9% and 4.4% with Pb(btsa)<sub>2</sub> and Pb(dbda), respectively (Figure S4). These values are similar to those of the TMA-H<sub>2</sub>O alumina process (2.4%, Figure S4c) and other metal sulfide processes such as  $SnS_2^{37}$  ZrS<sub>2</sub> and HfS<sub>2</sub><sup>38</sup> in the same cross-flow F120 ALD reactor but at much higher temperatures.

**2.3. Film Composition and Crystallinity.** The  $Pb(btsa)_2$ -H<sub>2</sub>S process produces crystalline and nearly stoichiometric lead sulfide films in a wide temperature range (Figure 2a,b). The films are the purest at the lowest deposition temperatures with a total impurity content of ca. 2 at. % that increases slightly toward higher deposition temperatures (Figure S5a). Typical time-offlight elastic recoil detection analysis (TOF-ERDA) elemental depth profiles are shown in Figure S5c,d. Also, the Pb(dbda)–  $H_2S$  process deposits the purest films at the lowest deposition temperatures (Figure S5b), but the stoichiometry changes significantly above 95 °C where the films contain more lead than sulfur (Figure 2c). XRD patterns of these lead-rich films expose reflections of impurity phases: PbO and Pb<sub>3</sub>(CO<sub>3</sub>)<sub>2</sub>(OH)<sub>2</sub> (Figure 2d).

**2.4. Formation of Pb Using Pb(dbda).** The presence of oxygen-containing phases in the XRD patterns was startling because the precursors did not contain oxygen. We speculated that at higher temperatures Pb(dbda) decomposes into elemental lead which would explain the changes in the stoichiometry of the films. When these films are exposed to ambient air, the elemental lead can react with  $O_2$ ,  $CO_2$ , or  $H_2O$  in air, producing the oxygen-containing phases observed in the XRD patterns. To test this hypothesis, we capped a film deposited at 135 °C with ca. 3 nm ALD  $Al_2O_3$  (TMA- $H_2O$  process) without breaking the vacuum and measured the XRD pattern immediately after removing the film from the reactor. Indeed, the XRD pattern of the capped film lacked reflections of the oxygen-containing phases and instead had reflections of elemental lead (Figure 2d).

Even though the presence of elemental lead supports our hypothesis on decomposition of Pb(dbda), it is still inconsistent with the higher thermal decomposition temperature (ca. 160 °C) reported for Pb(dbda) in the literature<sup>35</sup>—our own observation that signs of Pb(dbda) thermal decomposition (blackening of the hot end of the precursor source tube) in the F120 ALD reactor appear only at ≥155 °C—and an experiment where at 135 °C we pulsed only Pb(dbda) without H<sub>2</sub>S that yielded no deposit at all. Additionally, a set of depositions at 135 °C revealed that the process still works in an ALD regime with GPC saturating with respect to precursor pulse durations and films being visually uniform (Figures S6 and S7e). These observations suggest that elemental lead is formed by a mechanism that involves H2S rather than thermal decomposition of Pb(dbda). In Pb(dbda) lead has an oxidation state of +2, and therefore lead ion must be reduced and another species must be oxidized We speculated that the the oxidized species is sulfide  $(S^{2-})$  and that it oxidizes to elemental sulfur, which is volatile enough to be removed during the deposition. It is not uncommon for H<sub>2</sub>S to act as a reducing agent in ALD of transition metal dichalcogenides and reduce the metal ions in metal pentachlorides such as MoCl<sub>5</sub><sup>39</sup> and ReCl<sub>5</sub><sup>40</sup> to a lower oxidation state. This speculation is further supported by the fact that sulfur is sufficiently volatile at the temperatures and conditions being discussed; for example, Tripathi et al. used elemental sulfur as a precursor at a source temperature of 115 °C in ALD of copper sulfide,<sup>41</sup> in the same F120 ALD reactor as used here.

While reduction by  $H_2S$  accounts for some of our empirical observations, it does not explain the difference between the two lead precursors. To understand the cause of that difference, we evaluated the thermodynamics of the reduction with gas-phase density functional theory (DFT, def2-QZVPP// $\omega$ B97X-D3(BJ)/def2-SVP), correcting for the gas-to-solid transition of Pb and PbS with literature thermochemical data. These calculations indicate that reduction with  $H_2S$  with concomitant elimination of elemental sulfur is nonspontaneous, in direct

contradiction to the initially proposed mechanism (Table S1). To elucidate a plausible mechanism for the elemental lead formation, we then considered several other reactions such as film etching by the free ligand and formation of sulfur diamides. Of these candidates, only one fit with our empirical observations by requiring H<sub>2</sub>S and being spontaneous for Pb(dbda) but not Pb(btsa)<sub>2</sub>. That reaction is similar to the precursor thermolysis that yields elemental lead and *N-tert*-butylacetaldimine (non-spontaneous,  $\Delta G = 56$  kJ/mol); however, the consequent thiolysis of the imine to produce thioacetaldehyde which further oligomerizes to a heterocyclic trithiane makes this pathway sponatenous ( $\Delta G = -10$  kJ/mol, Scheme 1). While this

Scheme 1. Proposed Pathway for Deposition of Elemental Lead with Pb(dbda)

$$+ 2 H_2 S \rightarrow Pb + 2/3 \qquad S + 2 H_2 N$$

speculative mechanism fulfills all the requirements, we emphasize that it, or any alternative, would require a dedicated in situ mechanism study for verification. We believe that such a study, while outside of the scope of this work, would improve understanding of processes occurring in low-temperature deposition of lead materials with amide precursors.

The possibility of depositing elemental Pb with ALD at low temperatures without using a dedicated reducing agent is intriguing. We briefly explored whether that is possible by adjusting the process parameters. The PbS to Pb intensity ratios were unaffected by changes in the H<sub>2</sub>S pulse or purge durations (Figure S7a,b). In contrast, increasing the Pb(dbda) pulse duration increased the intensity from the Pb phase (Figure S7c,d). However, even with 15 s pulse duration some PbS was still present. The film with the highest Pb content had a low resistivity of 1.0 m $\Omega$  cm (Pb bulk resistivity 0.02 m $\Omega$  cm), resembled rough metallic films in appearance, and consisted of large micrometer-sized grains connected by smaller submicrometer-sized grains (Figure S7e,f).

**2.5. Film Morphology.** Formation of elemental lead at higher temperatures in the Pb(dbda)–H<sub>2</sub>S process affects also the morphology of the films. Films deposited at lower temperatures by both processes consist of small cube-shaped grains (Figure 3a) as is characteristic to PbS. The grains increase in size with increasing deposition temperature in the Pb(btsa)<sub>2</sub>–H<sub>2</sub>S process (Figure 3a and Figure S8), which linearly increases the roughness of the films with the deposition temperature (Figure 3b). In the Pb(dbda) H<sub>2</sub>S process, in addition to the cubic grains growing in size, irregularly shaped lumps start to appear (Figure 3a and Figure S9) at higher deposition temperatures. These lumps increase in size more rapidly with the deposition temperature than the cubic grains, and this in turn causes an exponential increase in the roughness of the films (Figure 3b).

**2.6. Electrical Properties.** Hall effect measurements revealed that both processes yield conductive p-type PbS films (Figure 3c). The hole mobilities are  $10-70 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$  (Figure 3d) and are comparable to or exceed the mobilities of semiconductors typically used in perovskite based devices, both organic such as PEDOT:PSS ( $2 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ )<sup>42</sup> and spiro-MeOTAD ( $2 \times 10^{-3} \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ )<sup>43</sup> as well as inorganic such as NiO ( $10 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$ ).<sup>44</sup>



**Figure 3.** (a) FESEM images of PbS films deposited at different temperatures (see also Figures S8 and S9). (b) AFM roughness  $R_{qr}$  (c) conductivity, (d) hole mobility, and (e) carrier concentration of the PbS films as a function of deposition temperature. The substrate was silicon in (a, b) and sapphire in (c, e). Films were ca. 100 nm thick and were deposited with 1.0 s precursor pulses and purge durations.



**Figure 4.** (a) Transmittance of PbS films deposited on sapphire with Pb(btsa)<sub>2</sub> and different number of deposition cycles at 75 °C. The dashed gray line separates wavelength ranges measured with different instruments. (b) Photograph of the films in (a) viewed against an office lamp. (c) Crosssection TEM image of a PbS film deposited with Pb(btsa)<sub>2</sub> at 75 °C with 2000 cycles. Some grain boundaries are marked with dashed lines. (d) Out-ofplane and in-plane crystallite sizes, as measured with XRD, in PbS films deposited with Pb(btsa)<sub>2</sub> at 75 °C with different number of cycles. (e) Transition energies defined from log(1/*T*) plots (Figure S14a,c), including the band gap  $E_0$ . Dashed gray lines are reference energy values by Cardona et al.<sup>66</sup> All films were deposited with 1.0 s pulse and purge durations.

2.6.1.  $Pb(btsa)_2$ . In the PbS films deposited with  $Pb(btsa)_2$ , hole mobility, conductivity, and carrier concentration increase at higher deposition temperatures (Figure 3c-e). We attribute the changes to the enlargement of grains (Figure 3a) and to the increase of sulfur and oxygen concentrations in the films (Figure 2a, Figure S5a, and Table S2). Enlargement of grains reduces grain boundary density, which in turn increases mobility.<sup>45,46</sup> whereas the excess of sulfur and oxygen result in defects that increase hole concentrations.<sup>1,47</sup> All chalcogens, when in excess, are known to be p-type dopants in lead chalcogenides.<sup>48,49</sup> The observed changes in composition are on the order of 10<sup>20</sup> atoms cm<sup>-3</sup> and are larger by a factor of 100 compared to the observed carrier concentrations. We speculate that this difference could be caused by carbon and hydrogen impurities because defects created by these impurities can decrease hole concentration. Overall, the observed carrier concentrations are a result of all defects created by nonstoichiometry and impurities.<sup>50</sup> However, describing the interactions of these defects and their effect on carrier concentration is beyond this study.

2.6.2. Pb(dbda). The presence of several phases in the films deposited with Pb(dbda) at higher temperatures (Figure 2d) prevents drawing unambiguous conclusions about their electrical properties. At lower deposition temperatures the changes in grain sizes and composition are not significant enough (Figure 2c and Figure S9a-c) to cause major differences in electrical properties of these films (Figure 3c-e). Compared to Pb(btsa)<sub>2</sub>, Pb(dbda) yields films with a conductivity larger by a factor of 100 (Figure 3c). The difference in conductivity arises from a difference in carrier concentration (Figure 3e). We attribute these differences to the excess of sulfur and lower impurity contents in the films made with Pb(dbda) at lower temperatures (Figure 2c and Figure S5b). The differences in the electrical properties between the two ALD processes offers flexibility for the design of devices based on semiconductor junctions. Various postdeposition annealing treatments can provide additional control over the electrical properties of PbS.5

**2.7. Optical Properties.** We expected that the large exciton radius (ca. 20 nm) of PbS, its narrow band gap, and the dependence of the grain size on the deposition temperature in our processes would allow us to deposit films with tailorable optical properties. To explore this hypothesis, we measured transmittance of PbS films deposited on sapphire in the 190-8000 nm range (Figure 4a,b and Figure S10a-c) and determined their optical band gaps. Optical band gaps are commonly determined from Tauc plots constructed from transmittance data, and for PbS thin films many examples of this approach are present in the literature (Table S3). $^{6,52-6}$ However, we found optical band gap extraction from Tauc plots to be problematic in the case of PbS. The values extracted from Tauc plots of ~100 nm thick PbS films deposited at different temperatures were unreasonably large ( $\geq 2.5$  eV, Figure S10df) for the observed grain sizes (Figures S8, S9, and S11). The grains could also be confined in the film thickness direction, which is not visible in top-down SEM. We initially assumed that the grains are columnar and extend through the entire film thickness; however, TEM revealed that this is not the case, and there are also grain boundaries through film thickness (Figure 4c)—despite that the grains are still too large to explain the Tauc values.

The grains observed with microscopy methods can consist of smaller crystallites, which in turn if sufficiently small would explain the Tauc values. We calculated crystallite sizes using the Scherrer equation and XRD data measured out-of-plane  $(\theta - 2\theta)$ and in-plane  $(2\theta\chi)$ , respectively (Figure S12), for the films deposited with Pb(btsa)<sub>2</sub>. In the 100 nm thick films deposited at different temperatures the in-plane crystallite size increased significantly with increasing deposition temperature, from  $47 \pm$ 3 nm at 65 °C to 135  $\pm$  33 nm at 135 °C, whereas the out-ofplane crystallite size remained almost unchanged at ca. 56 nm (Figure S13). The dependence of the in-plane crystallite size on the deposition temperature is in good agreement with the observed morphology (Figure S8) and trends in electrical properties (Figure 3). The out-of-plane XRD patterns show that in these films most crystallites are {200} oriented with some minor  $\{111\}$  and  $\{220\}$  components. Correspondingly, the most prominent reflections in the in-plane patterns of these films are also {200} (Figure S12c). In the films deposited at 75 °C with different number of cycles both the out-of-plane and inplane crystallite sizes increase rapidly with the increasing number of cycles until 1500 cycles (Figure 4d). After 1500 cycles, the growth of the in-plane crystallite size slows down considerably, whereas the growth of the out-of-plane crystallite size continues and approaches the limit that can be measured from peak broadening. The XRD patterns also reveal changes in the film texture. In the out-of-plane direction {200} dominates up to 2000 cycles and  $\{111\}$  afterward, whereas in the in-plane direction also {200} dominates up to 2000 cycles, then {111} at 3000 cycles, and {220} afterward.

The crystallite dimensions we observed imply no quantum confinement effects in the films deposited with 500 cycles or more and are therefore in disagreement with the Tauc values. Others have also noted a similar discrepancy between the Tauc values and crystallite sizes but provided no explanation.<sup>52,53</sup> Additionally, there is an inconsistency between the reported Tauc values and electronic band gaps. Dasgupta et al. determined the electronic band gap of ALD PbS films using scanning tunneling microscopy (STM).<sup>67</sup> These measurements showed that the electronic band gap decays non-linearly with increasing film thickness and approaches the bulk value when the film thickness is 16 nm, which is in line with the exciton Bohr radius of PbS (ca. 20 nm).

The cause of the conflicting evidence is elucidated in the earlier works of Gibson,<sup>68,69</sup> Scanlon,<sup>70</sup> and Cardona et al.<sup>66</sup> on the optical and electronic properties of PbS. Examination of PbS transmittance or reflectance in a wide energy range (0.2-20 eV)reveals several band-to-band transitions in addition to the lowest energy transition, i.e., the band gap or  $E_0$ . Cardona et al. give a comprehensive list of these transitions  $(E_0 \text{ to } E_6)$  as well as estimate some of the energies of these transitions using reflectance and transmittance spectra from single crystals and epitaxial thin films. In the case of thin film transmittance spectra, the transition energies were estimated based on the position of the features in the  $\log(1/T)$  plot. In fact,  $\log(1/T)$  plots of our films are strikingly similar to Cardona et al. (Figure S14a-d) and contain four features corresponding to the interband transitions  $E_0 - E_3$ . The optical band gap of thicker films (>500 deposition cycles) defined from log(1/T) plots is ca. 0.5 eV (Figure 4e and Figure S14e). This value is in excellent agreement with the reported bulk band gap and electronic band gap of PbS, whereas the increase of the interband transition energies in very thin films (<500 deposition cycles) is in excellent agreement with decrease of crystallite size below the exciton radius (20 nm) in these films (Figure 4d,e).

To summarize, a Tauc plot is not an ideal method for determining the optical band gap of PbS. The Tauc plot

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**Figure 5.** (a) Cross-section SEM image of a PbS-capped MAPI used for stability study in (d). PbS was deposited with Pb(dbda). To examine the cross section, the sample was broken in half after being stored in air for 413 days. For EDS line scans across the cross section see Figure S19. Relative reflection intensities (area) of crystalline phases observed in GIXRD patterns of samples where PbS was deposited on MAPI films with (b) Pb(btsa)<sub>2</sub> and (c) Pb(dbda) at different temperatures. (b) and (c) are derived from GIXRD patterns in Figures S16 and S17. MAPI film reflection intensity (d) as a function of storage time in ambient air (open system) with and without different capping films derived from GIXRD patterns in Figures S18–S22. PbS capping film was deposited at 45 °C with Pb(dbda) and at 65 °C with Pb(btsa)<sub>2</sub>. The Al<sub>2</sub>O<sub>3</sub> capping film was deposited at 65 °C with 800 cycles of the TMA–H<sub>2</sub>O process using a 1.0 s precursor pulse and 4.0 s purge durations. The capping film thicknesses were 66, 50, and 52 nm for PbS deposited with Pb(dbda), PbS deposited with Pb(btsa)<sub>2</sub>, and Al<sub>2</sub>O<sub>3</sub>, respectively. PbS films in (b), (c), and (d) were deposited with 1000 cycles, 1.0 s precursor pulses, and purge durations. (e) Schematic of the changes occurring on the MAPI film in the ALD reactor when deposition temperatures higher than 65 °C are used.

guidelines available in the literature suggest deriving the band gap from the most noticeable Tauc property  $((\alpha h\nu)^2)$  change,<sup>7</sup> which in our case yielded unreasonably large values. However, in optical spectra the most noticeable feature does not necessarily correspond to the band gap  $(E_0)$ . For PbS the most noticeable feature is  $E_2$  (Figures S10d,e and S14a-d,f).<sup>66</sup> The Tauc property itself provides an additional challenge for narrow gap materials as it accentuates higher energy transitions and masks lower energy transitions (Figure S14f-h). Logarithmic plots of Tauc property can reveal these hidden transitions and enable extraction of their energies. In the case of our Tauc plots in addition to  $E_2$  we were able to find also  $E_0$  (ca. 0.5 eV) and  $E_1$ (ca. 1.6 eV) by extensive rescaling (Figure S14g,h). We note the similarity between the  $E_1$  and  $E_2$  transitions energies extracted from Tauc plots and the band gaps reported in the literature for PbS thin films that are also extracted from Tauc plots (Table S3). An accurate distinction of the band gap transition from the other higher energy interband transitions from optical data requires examination of said data in a wide energy range that also covers the region of the bulk band gap of PbS (0.2-1.0 eV). In our case log(1/T) plots were more robust than the Tauc plots

for this purpose. The transitions observed in optical data are best understood in terms of the band structure of PbS, for which we refer the reader to Albanesi et al.<sup>72</sup>

**2.8. Deposition of PbS onto Halide Perovskites.** We investigated whether the ALD processes developed and described here are compatible with halide perovskites. When sulfides are deposited onto perovskites, there is a possibility of an anion exchange between the perovskite and  $H_2S$  (eq 1).

$$CH_{3}NH_{3}PbI_{3}(s) + H_{2}S(g) \rightarrow CH_{3}NH_{3}I(g) + 2HI(g) + PbS(s)$$
(1)

However, based on thermodynamic values reported in the literature,  $^{73-75}$  that reaction is endothermic ( $\Delta H = 146 \text{ kJ/mol}$ ) and unfavorable ( $\Delta G = 121 \text{ kJ/mol}$  at RT). In addition, it was experimentally verified that the MAPI film was unaffected when H<sub>2</sub>S was pulsed onto it at 50 °C (Figure S15). We deposited PbS films at different temperatures on MAPI films (Figure 5a) and used XRD to examine the state of the MAPI film after the deposition. In the case of the Pb(btsa)<sub>2</sub>-H<sub>2</sub>S process only the lowest deposition temperature of 65 °C left the MAPI film unchanged (Figure 5b and Figure S16). At deposition

temperatures higher than 65 °C PbI<sub>2</sub> reflections appeared and increased in intensity, whereas MAPI reflections decreased in intensity. This is caused by the decomposition of MAPI into PbI<sub>2</sub> and volatile CH<sub>3</sub>NH<sub>3</sub>I upon heating to the target deposition temperature before PbS deposition begins (Figure 5e). When the PbS deposition temperature was above 95 °C, also the intensity of PbI<sub>2</sub> reflections decreased. This decrease is caused by the evaporation of PbI<sub>2</sub> before the PbS deposition.

MAPI films remained unchanged when PbS was deposited on them with Pb(dbda) in the temperature range 45–60 °C (Figure 5c and Figure S17). At the deposition temperature of 75 °C and higher, MAPI expectedly decomposed into CH<sub>3</sub>NH<sub>3</sub>I and PbI<sub>2</sub>, as in the case when Pb(btsa)<sub>2</sub> was used. At the deposition temperature of 95 °C and higher reflections of Pb<sub>5</sub>S<sub>2</sub>I<sub>6</sub> and Pb<sub>3</sub>(CO<sub>3</sub>)<sub>2</sub>(OH)<sub>2</sub> phases appeared (Figure S17). The appearance of these phases coincides with the onset of formation of elemental Pb. Notably, no Pb<sub>5</sub>S<sub>2</sub>I<sub>6</sub> was observed with Pb(btsa)<sub>2</sub>. These results imply that the species involved in the formation of elemental Pb can react with the PbI<sub>2</sub> remaining from the decomposition of the MAPI film to yield Pb<sub>5</sub>S<sub>2</sub>I<sub>6</sub>. The Pb<sub>3</sub>(CO<sub>3</sub>)<sub>2</sub>(OH)<sub>2</sub> phase is apparently formed upon exposure of Pb<sub>5</sub>S<sub>2</sub>I<sub>6</sub> or Pb to the ambient air.

In the case of PbS deposition on MAPI the Pb(dbda)–H<sub>2</sub>S process is superior to the Pb(btsa)<sub>2</sub>–H<sub>2</sub>S. Pb(dbda) offers a 20 °C wide temperature window where PbS can be deposited on MAPI, whereas Pb(btsa)<sub>2</sub> only provides a single useful deposition temperature of 65 °C. The benefit of the wider temperature window comes with the cost of more complicated synthesis. Therefore, Pb(btsa)<sub>2</sub> can be a more convenient choice for perovskites that can withstand higher deposition temperatures such as  $FA_xMA_{1-x}Pb(Br_yI_{1-y})_3$  (up to 120 °C),<sup>76</sup>  $Cs_xFA_{1-x}Pb(Br_yI_{1-y})_3$  (up to 110 °C),<sup>77</sup> and others.<sup>24</sup> On the contrary, the deposition of Pb makes Pb(dbda) less relevant at higher deposition temperatures. Ultimately, the need of certain electrical properties can dictate which of the two processes is more suitable in a specific device design.

**2.9. Application of PbS as a Protective Layer.** Finally, to determine whether encapsulation with an ALD PbS film can protect halide perovskites from the ambient air, we followed the state of PbS-capped MAPI films with XRD. Without any encapsulation, the MAPI film completely decomposed into PbI<sub>2</sub> in a week, and after 5 months, PbI<sub>2</sub> had further reacted into PbIOH (Figure 5d and Figure S18). On the other hand, MAPI films were stable for at least 5 months in air when encapsulated with PbS films deposited either from Pb(btsa)<sub>2</sub> or Pb(dbda) (Figures S20 and S21). For comparison, we also evaluated the stability of a MAPI film encapsulated with ALD Al<sub>2</sub>O<sub>3</sub> film deposited with the TMA–H<sub>2</sub>O process (Figure S22). The improvement in the MAPI stability provided by the PbS encapsulation equals that provided by an Al<sub>2</sub>O<sub>3</sub> film of similar thickness (Figure Sc).

We attempted to further characterize the barrier properties of our films by depositing them on 50  $\mu$ m thick biaxially oriented poly(ethylene terephthalate) (BoPET) and measuring water vapor (WVTR) and oxygen transmission rates (OTR) of these samples. However, film deposition on BoPET was challenging. Relatively long pulse (6 s) and purge times (3 s) were required for visually uniform films on BoPET, and the GPC was significantly lower than on rigid substrates: 0.20 Å with Pb(btsa)<sub>2</sub> at 65 °C and 0.25 Å with Pb(dbda) at 45 °C. With the PbS capping the OTR decreased by ca. 90% with respect to the bare BoPET, however, the WVTR decreased only slightly (Figure S23). We also observed a significant decrease in the OTR and only a slight decrease in the WVTR of the reference BoPET samples coated with ALD  $Al_2O_3$ . These results indicate that ALD PbS and  $Al_2O_3$  coatings perform similarly in these barrier tests, despite the polycrystalline structure of the PbS. In general, it is thought that amorphous films will serve as better barriers because they lack grain boundaries along which diffusion can be faster.<sup>78</sup> A thin amorphous passivating PbSO<sub>3</sub>/PbSO<sub>4</sub> layer that forms on the PbS surface<sup>79</sup> can account for the  $Al_2O_3$  comparable barrier performance of our crystalline PbS films.

While low deposition temperatures enable the use of polymers as substrates, the ALD of barrier films on polymers is not straightforward. This is partially because of the engineering challenges associated with the handling of flexible substrates (see Figure S24) but even more because of the complex mechanism of ALD film nucleation and growth on polymers.<sup>80</sup> Barrier performance of an ALD film is affected by many factors such as porosity, roughness, and surface structure of the substrate polymer as well as crystallinity, impurity content, and thickness of the coating itself.<sup>81</sup> For instance, a thicker coating could block gases more efficiently and be more mechanically durable but could also crack more easily, whereas a thinner coating would not necessarily form a continuous and tight layer.<sup>80–82</sup> In the literature WVTR values as low as  $10^{-1}$  g m<sup>-2</sup> day<sup>-1</sup> were reported for 50 nm thick thermal ALD Al<sub>2</sub>O<sub>3</sub> films deposited at 65 °C on thicker (125  $\mu$ m) PET substrates<sup>83</sup> and even lower ones for ALD Al<sub>2</sub>O<sub>3</sub> on rigid substrates measured with electrical calcium corrosion tests.<sup>84</sup> On the basis of the literature results and our own observations, we anticipate that further optimization of the pretreatment, process, and handling parameters will enable lower WVTR values for ALD PbS coatings on polymers. However, the extent of optimizing these parameters is a topic of a dedicated study and outside of the scope of this work which aimed to demonstrate the ability of ALD PbS films to protect halide perovskite films on rigid substrates.

### 3. CONCLUSIONS

We described two new ALD processes for deposition of PbS films using Pb(btsa)<sub>2</sub>, Pb(dbda), and H<sub>2</sub>S as precursors. The scale-up of Pb(dbda) synthesis afforded reasonable precursor purity and quantity for process development. Yet challenges encountered during the scale-up highlight the importance of further studies on the synthetic chemistry of volatile Pb complexes. Both Pb precursors enable deposition of high quality PbS films at low temperatures of 65-155 °C with Pb(btsa)<sub>2</sub> and 45-75 °C with Pb(dbda). The PbS films are uniform, crystalline, and of high purity as well as conductive and p-type. We demonstrated that at deposition temperatures of  $\leq 65$  °C these ALD processes are compatible with the perovskite halide (CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>) films. Pb(dbda) is more suitable for capping the CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> perovskite because it is applicable in a range of temperatures versus only one temperature with Pb(btsa)<sub>2</sub>.  $Pb(btsa)_{2}$ , on the other hand, is a more suitable choice for perovskites that withstand higher temperatures such as those based on cesium or formamidinium as well as in cases where high carrier mobility is sought. Capping the CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub> films with ALD PbS prevents the perovskite from decomposing in the ambient air. The encapsulation performance of ALD PbS is comparable to ALD Al<sub>2</sub>O<sub>3</sub>. Compatibility with various substrates such as polymers provided by the low deposition temperatures and the ability to control electrical and optical properties of the PbS films offer flexibility for design of device

structures based on halide perovskites or PbS. We hope that this work will encourage other studies on the use of non-oxide ALD films together with halide perovskites and that the new ALD PbS processes are viable options for PbS based devices.

## 4. EXPERIMENTAL SECTION

**4.1. Precursor Synthesis.** Lead(II) bis[bis(trimethylsilyl) amide] (Pb(btsa)<sub>2</sub>) was synthesized from lead(II) chloride (Alfa Aesar, 99.999%) and lithium bis(trimethylsilyl)amide (Li(btsa), Aldrich, 95%) as previously described.<sup>34,85</sup> Lead(II) *rac-N<sup>2</sup>,N<sup>3</sup>*-di-*tert*-butylbutane-2,3-diamide (Pb(dbda)) was prepared by a modified literature procedure.<sup>35</sup> Pb(btsa)<sub>2</sub> (13.160 g, 24.923 mmol) was dissolved in 100 mL of toluene at ambient temperature under nitrogen in a 200 mL Schlenk flask to yield a clear orange solution. Then, free rac-N,N'-ditert-butyl-2,3-diaminobutane (H2dbda, 5.000 g, 24.94 mmol) diluted in 10 mL of toluene was added all at once via cannula. Some decomposition was evident as the mixture turned dark green-gray with the precipitation of metallic Pb. The mixture was then immediately heated to reflux in an oil bath to facilitate the transamination reaction. Aliquots were taken by syringe and placed into an NMR tube under nitrogen; the majority of solvent was removed under vacuum; then C<sub>6</sub>D<sub>6</sub> was added to prepare the sample for spectroscopy. In this way, the reaction was monitored by ex situ <sup>1</sup>H NMR spectroscopy until the signal for  $Pb(btsa)_2$  had disappeared (ca. 3 h). The reaction mixture was then cooled to ambient temperature, and volatiles (i.e., free bis(trimethylsilyl)amine and toluene) were removed under high vacuum (ca.  $10^{-2}$  Torr) until a blood-red oil remained. Most of the unreacted H<sub>2</sub>dbda was removed by distillation through a 90° glass elbow under high vacuum (40  $^{\circ}C/10^{-2}$  Torr) into a dry ice-acetone cooled (-78 °C) receiving Schlenk flask. The almost complete removal of H<sub>2</sub>dbda was conveniently indicated by solidification of the mixture in the pot. After the receiving flask was replaced with a clean Schlenk flask, the crude product was purified by sublimation (60 °C,  $10^{-2}$  Torr) into the air-cooled elbow to yield Pb(dbda) (5.1 g, 51%, based on Pb(btsa)<sub>2</sub>) as a blood-red crystalline solid. Remaining ligands and solvent were condensed into the dry ice-acetone cooled receiving flask and removed mechanically. Samples prepared in this way had identical spectroscopic and physical properties to those previously reported.<sup>35</sup>

**4.2. Film Deposition.** Caution!  $H_2S$  is a flammable, corrosive, and highly toxic gas. The ALD reactor used in this work was in a dedicated laboratory equipped with an  $H_2S$  sensor and where gas masks were readily available. The  $H_2S$  bottle was kept in a gas cabinet with isolated ventilation. EPDM (ethylene propylene diene monomer) rubber seals were used. The exhaust was bubbled though an aqueous solution of  $Cu(NO_3)_2$ . The authors urge the readers to acquaint themselves with the article by Dasgupta et al.<sup>86</sup> before working with  $H_2S$ .

The films were deposited in a commercial, hot-wall, flow-type F120 ALD reactor (ASM Microchemistry Ltd.) operated under reduced pressure of ca. 10 mbar. A continuous flow (400 sccm) of nitrogen (AGA, 99.999%), introduced into the reactor through a gas purifier (SAES Pure Gas, MC1-902F, H<sub>2</sub>O, O<sub>2</sub>, CO, CO<sub>2</sub>, H<sub>2</sub>, NMHC removal <1 ppb), was used as a carrier and purge gas. Pb(btsa)<sub>2</sub> and Pb(dbda) were sublimed from open glass boats held inside the reactor at 60 and 40 °C, respectively. For both precursors the first signs of decomposition were observed at deposition temperatures of 155 °C. H<sub>2</sub>S (Linde, 99.5%) was introduced into the reactor through a gas purifier (SAES Pure Gas, MC1-302F, H<sub>2</sub>O removal <1 ppb) and a needle valve, at a flow rate of 14 sccm. Al<sub>2</sub>O<sub>3</sub> was deposited by using the trimethylaluminum (TMA)– $H_2O$  process. The films were deposited on 5 × 5 cm<sup>2</sup> silicon (100), soda lime glass, and biaxially oriented poly(ethylene terephthalate) (BoPET, GoodFellow, 50  $\mu$ m thick) substrates as well as on 2 in. sapphire wafers (*c*-plane off to *M*-plane by  $0.2 \pm 0.1^{\circ}$ , double side polished, University Wafer). For the deposition on BoPET a special adapter was machined from Al sheets to prevent the movement of the polymer substrate by pressing it at the edges, without the adapter the movement of the polymer occasionally blocked the gas flows in the deposition chamber (Figure S24).

 $CH_3NH_3PbI_3$  films were made by a conversion of  $PbI_2$  films. First  $PbI_2$  films were deposited with ALD at 75 °C with the  $Pb(btsa)_2-SnI_4$ 

process (Acros Organics, >99%, -10 mesh).<sup>34</sup> Then a TLC plate (0.75 mL/g pore volume silica gel on Al foil, Sigma-Aldrich) was submerged in a 500 mg/mL solution of CH<sub>3</sub>NH<sub>3</sub>I (>99.0%, TCI) in dimethylformamide ( $\geq$ 99.9%, Sigma-Aldrich) for 30 min. The TLC plate was then dried in air by keeping it on a 60 °C hot plate for 30 min. The dry plate was placed on the bottom of a steel cylinder (V = 0.5 L) with the MAI permeated SiO<sub>2</sub> side facing up. An aluminum support holding the PbI<sub>2</sub> film face down was then placed on the TLC plate. The design of the aluminum support is such that the distance between the MAI permeated SiO<sub>2</sub> face of the TLC plate and the surface of the PbI<sub>2</sub> film was ca. 1 mm. After the loading the cylinder was cycled with N<sub>2</sub> (AGA, 99.999%)/vacuum three times and isolated with vacuum inside (ca. 1 mbar). The evacuated cylinder was heated in an air oven at 140 °C for 24 h to complete the conversion of PbI<sub>2</sub> to CH<sub>3</sub>NH<sub>3</sub>PbI<sub>3</sub>.

**4.3. Film Characterization.** Prior to characterization, the films were stored in closed containers filled with ambient air unless indicated otherwise. No significant oxidation is expected for PbS stored in air,<sup>79,87,88</sup> but instead a thin (ca. 0.5 nm) passivating PbSO<sub>3</sub>/PbSO<sub>4</sub> layer can form upon prolonged exposure.<sup>79</sup>

Film thicknesses were measured with X-ray reflectivity (XRR, Rigaku SmartLab diffractometer). To measure XRR thickness maps, slits were used to limit the irradiated area<sup>89</sup> to  $\leq$ 5 × 4 mm<sup>2</sup> (at the angle range of interest), and an automated XY stage was used to measure a 7 × 7 matrix across the 5 × 5 cm<sup>2</sup> samples with 0.5 cm edge exclusion. Grazing incidence XRD (GIXRD) patterns were measured at 1° incidence angle. Grain sizes were calculated by using the Scherrer equation and full width at half-maximum of the most intense reflection. Scherrer constants from Langford and Wilson et al.<sup>90</sup> were used for the respective lattice planes assuming cubic crystallites. The instrumental broadening was defined by measuring PbS powder (99.9%, Sigma-Aldrich) out-of-plane and in-plane with the same optics as used to measure PbS thin films.

The morphology of the films was studied with a Hitachi S-4800 field emission scanning electron microscope (FESEM) and with atomic force microscopy (AFM) by using a Veeco Multimode V instrument. AFM tapping mode images were captured in air using silicon probes with a nominal tip radius of 10 nm and a nominal spring constant of 3 N/m (NFESP from Bruker). Images were flattened to remove artifacts caused by a sample tilt and a scanner bow. Resistivities were determined from sheet resistance measured with four-point probe (CPS Probe Station, Cascade Microtech combined with a Keithley 2400 SourceMeter) from films deposited on soda lime glass.

Hall effect measurements of films on sapphire substrates were made at room temperature with an Ecopia Van der Pauw HMS-5000 measurement system. Time-of-flight elastic recoil detection analysis was performed using a 40 MeV <sup>127</sup>I<sup>7+</sup> beam. The incident angle of the beam was 16°, and the scattering angle was 40°. An FEI Quanta 3D 200i dual-beam microscope and standard lift-out procedures were used to prepare TEM cross-section specimens. An FEI Tecnai F20 microscope operated at 200 kV was used for bright-field TEM imaging. An Oxford INCA 350 energy spectrometer connected to the FESEM was used for energy dispersive X-ray (EDS) measurements.

Transmittance of the films on sapphire substrates in the 190-3000 nm wavelength range was measured with an Agilent Cary 5000 spectrophotometer equipped with an internal diffuse reflectance measurement accessory (DRA-2500). The internal DRA consists of a 110 mm diameter integrating sphere with input and output apertures to perform reflectance and transmittance measurements. Sample of interest was mounted at the input aperture and the output aperture of the integrating sphere was covered with a standard reference material in the transmittance mode measurement. A Labsphere Spectralon SRM-99 diffuse reflectance standard was used for this purpose. Reference light was passed through the sample of interest when sample spectra were collected and the reference light intensity, in this case the emission light of Cary 5000, was measured without a sample. The scan speed was set to 10 nm/s and step size to 1 nm. Transmittance in the 3000-8000 nm wavelength range was measured with a Bruker Alpha II FTIR spectrometer. For Tauc plot construction direct allowed band gap was assumed for PbS.

Water vapor transmission (WVTR) tests were performed, and results were calculated according to the ASTM standard E 96/E 96M-05 desiccant method with modifications in custom-built chambers that comply with the standard requirements. The relative humidity (RH) gradient used in these measurements was 0/75% instead of the 0/50%in the standard. The exposed sample surface area was limited to 5 cm<sup>2</sup> with an adhesive PTFE/aluminum masking. The temperature was 23 °C. The temperature and RH of the cabinet were measured by using a Rotronic RH meter (Bassersdorf, Switzerland) before each weighing. The experiments were performed by using three replicates. The WVTR was calculated from the weight gain via equation WVTR =  $(\Delta W/t)/A$ , where  $\Delta W$  is the weight gain, *t* the time, and *A* the exposed surface area of the sample. The oxygen transmission rate and permeability were tested according to the ASTM standard D 3985-05 by using a Systech Illinois 8001 oxygen permeation analyzer (Johnsburg, IL). A mask (A Systech Illinois 8001 accessory) was used to limit the surface area of the samples to 5 cm<sup>2</sup>. The  $O_2$  partial pressure gradient was 1 atm with  $N_2$  on the receiving side, the temperature was set to 23 °C, and the relative humidity of both gases (N2 and O2) was 50%. Two replicates of each sample were measured.

4.4. Computational Details. Reaction energies were calculated with density functional theory (DFT) in ORCA v4.2.1.92 Geometries were optimized by using Head-Gordon's range-separated hybrid functional modified by Najibi and Goerigk<sup>93</sup> to include Grimme's dispersion correction with Becke–Johnson damping<sup>94,95</sup> ( $\omega$ B97X-D3(BJ)), Ahlrich's double- $\zeta$  split-valence quality basis set on all atoms (def2-SVP),<sup>96</sup> and the Stuttgart ECP60MDF pseudopotential on Pb.<sup>9</sup> The resolution of the identity approximation for the Coulomb integrals with the chain of spheres approximation for the exchange term (RIJCOSX)<sup>98</sup> with the def2/J auxiliary basis set<sup>99</sup> was used to speed up the self-consistent field (SCF) step of all calculations. Frequency analysis was used to confirm there were no imaginary vibrational modes, and the geometries were minima on the potential energy surface. Numerical frequency calculations were performed to work around a bug in the ORCA code involving analytical Hessian calculations together with the RIJCOSX approximation. Single-point calculations were then performed with a quadruple- $\zeta$  basis set (def2-QZVPP) to obtain more accurate electronic energies to use with the thermochemical corrections from the previous frequency analysis. Enthalpy, entropy, and free energy changes of reaction were calculated by using Hess's law, with gas-to-solid transition enthalpy and entropy changes for Pb and PbS taken from the NIST Standard Reference Database Number 69.7

# ASSOCIATED CONTENT

# **③** Supporting Information

The Supporting Information is available free of charge at https://pubs.acs.org/doi/10.1021/acs.chemmater.0c01887.

Energetics of reactions considered in the discussion of the mechanism of elemental lead formation; element concentrations measured with TOF-ERDA in PbS films; band gaps reported for PbS films in the literature; photographs of PbS films; GPC as a function of purge duration; thickness as a function of deposition cycles; thickness maps; impurity content as a function of deposition temperature; elemental depth profiles of the films; ALD study of Pb(dbda)-H<sub>2</sub>S process at 135 °C; FESEM images of PbS films deposited at different temperature and with different number of deposition cycles; optical properties of PbS films; in-plane and outof-plane XRD measurements of PbS films; Scherrer crystallite size as a function of temperature; XRD measurements of various samples with MAPI perovskite; transmission rates of PbS and Al<sub>2</sub>O<sub>3</sub> covered BOPET; description of the adapter used for BOPET depositions (PDF)

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#### Notes

The authors declare no competing financial interest.

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