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Working Paper

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A GENETIC MODEL WITH LIMIT CYCLE BEHAVIOR**

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**International Institute for Applied Systems Analysis
A-2361 Laxenburg, Austria**

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Abstract

A two-locus genetic model, based on Dawkin's "sex war" game, with the fitness of the genotypes at each locus depending on the gene frequencies at the other, is shown to give rise to a stable limit cycle. The mathematical analysis involves averaging techniques and elliptic integrals.

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**THE "BATTLE OF THE SEXES":
A GENETIC MODEL WITH LIMIT CYCLE BEHAVIOR**

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1. Introduction

This note describes a simple two-locus genetic model that gives rise to a stable limit cycle. The origin of the model is Dawkins' (1976) "sex war" game, a simple evolutionary game describing the parental investment conflict. In this game males have two possible strategies, faithful (E_1) or philanderer (E_2), and females may play the strategies fast (F_1) or coy (F_2). Assuming +15 as the reward for having a child (to both parents), -20 as the total cost of raising the offspring, and -3 for the prolonged courtship coy females insist on, Dawkins obtained the following pay-off matrices A for males and B for females:

$$A, B = \begin{bmatrix} 5,5 & 2,2 \\ 15,-5 & 0,0 \end{bmatrix}. \quad (1)$$

This bimatrix game has a totally mixed Nash equilibrium at $(\frac{5}{8}, \frac{3}{8})$ and $(\frac{1}{6}, \frac{5}{6})$ (see Dawkins (1976), Maynard Smith (1982), Schuster and Sigmund (1981)). It is not an evolutionarily stable state (ESS), however, since the second (stability) condition is not satisfied. In fact, Selten (1980) has shown in general that in asymmetric conflicts, only pure strategies can be evolutionarily stable. Hence the static, game theoretic approach is not very satisfactory for this particular example, and it seems reasonable to introduce some dynamics.

It turns out, however, that the standard continuous time game dynamics, as studied e.g. by Zeeman (1980), may be generalized in (at least) two different ways to asymmetric conflicts: the one, see Schuster and Sigmund (1981), gives rise to conservative oscillations, the other, see Maynard Smith (1982), Appendix J, makes

the equilibrium asymptotically stable. With discrete time, however, the stationary solution would always be unstable (Eshel and Akin, 1983).

The first type of dynamics is also relevant for the analysis of our new model in § 2: With x_i, y_j denoting the relative frequencies of the strategies E_i, F_j respectively and the basic assumption

$$\frac{\dot{x}_i}{x_i} = \text{growth rate of } E_i = \text{average payoff for } E_i \quad (2)$$

$$\text{—average payoff for males} = (A\underline{y})_i - \underline{x}^t A \underline{y} .$$

Schuster and Sigmund (1981) derive the equations

$$\dot{x} = x(1-x)(2-12y) \quad (3)$$

$$\dot{y} = y(1-y)(-5+8x)$$

for $x = x_1 = 1-x_2$, $y = y_1 = 1-y_2$. They show that

$$V = x^5(1-x)^3 y^2(1-y)^{10} \quad (4)$$

is a constant of motion, so that all solutions oscillate around the Nash equilibrium, which also represents the time-average of any orbit. (Note the strong analogy to the Lotka-Volterra predator-prey equations.)

This model has one main drawback, however: It neglects the genetic structure. The derivation (2) assumes implicitly that individuals reproduce asexually. There is thus a need for a genetic model which takes into account the sexual interaction between males and females. One such diploid model was studied by Bomze et al. (1983), again using continuous time and leading to similar conservative oscillations. As this lacks structural stability too, a more natural model would be to consider a diploid model with separated generations.

2. The Model

We consider two loci, with two alleles at each locus. The A-locus, with alleles A and a , regulates male behavior: $A \rightarrow$ strategy E_1 (= faithful), $a \rightarrow E_2$ (= philanderer), and is not expressed in females. More precisely, we assume that homozygotes AA and aa play the pure strategies E_1 and E_2 respectively, and that the heterozygote Aa plays the mixed strategy $\frac{1}{2}(E_1 + E_2)$. In a similar way the B-locus determines female behavior: $B \rightarrow F_1 =$ fast, $b \rightarrow F_2 =$ coy. Let p_A, q_A be the fre-

quencies of allele A in adult males and females respectively, and p_B, q_B be the corresponding frequencies of allele B .

Assuming linkage equilibrium, the frequencies of AA, Aa, aa in the next generation are $p_A q_A, p_A(1-q_A) + q_{sub}A(1-p_A) = p_A + q_A - 2p_A q_A, (1-p_A)(1-q_A)$ respectively. The marginal frequency of gene A , which also measures the probability of strategy E_1 among males in the new generation, is then given by $p_A q_A + \frac{1}{2}(p_A + q_A - 2p_A q_A) = \frac{1}{2}(p_A + q_A)$. Similar results hold for the B-locus.

In order to simplify the mathematical analysis, we replace Dawkins' original pay-off matrix (1) by the following more symmetric one (this does not change the situation qualitatively):

$$\begin{bmatrix} 1-K, & 1+L & 1+K, & 1-L \\ 1+K, & 1-L & 1-K, & 1+L \end{bmatrix} \quad \text{with } 0 < K, L < 1. \quad (5)$$

This gives the following fitness values for males, depending on the frequency of females:

$$\begin{aligned} w_{AA} &= (1-K)\frac{p_B+q_B}{2} + (1+K)\left(1-\frac{p_B+q_B}{2}\right) = 1+K(1-p_B-q_B) \\ w_{aa} &= (1+K)\frac{p_B+q_B}{2} + (1-K)\left(1-\frac{p_B+q_B}{2}\right) = 1-K(1-p_B-q_B) \\ w_{Aa} &= \frac{1}{2}(w_{AA}+w_{aa}) = 1, \end{aligned} \quad (6)$$

with similar results for females.

Hence writing p'_A , etc. for the allelic frequencies in the adults of the new generation,

$$\begin{aligned} p'_A &= \frac{w_{AA} p_A q_A + \frac{1}{2} w_{Aa} (p_A + q_A - 2p_A q_A)}{\bar{w}_A} \\ q'_A &= \frac{p_A q_A + \frac{1}{2} (p_A + q_A - 2p_A q_A)}{1} = \frac{p_A + q_A}{2}, \text{ etc.} \end{aligned}$$

The mean fitness of males at the A-locus, w_A , is given by

$$\begin{aligned} \bar{w}_A &= 1 + K[p_A q_A (1-p_B-q_B) - (1-p_A)(1-q_A)(1-p_B-q_B)] \\ &= 1 - K(1-p_A-q_A)(1-p_B-q_B) . \end{aligned}$$

This leads finally to the equations

$$\left. \begin{aligned}
 p'_A &= \frac{\frac{1}{2}(p_A + q_A) + K(1 - p_B - q_B)p_A q_A}{1 - K(1 - p_A - q_A)(1 - p_B - q_B)} \\
 q'_A &= \frac{1}{2}(p_A + q_A) \\
 p'_B &= \frac{1}{2}(p_B + q_B) \\
 q'_B &= \frac{\frac{1}{2}(p_B + q_B) - L(1 - p_A - q_A)p_B q_B}{1 + L(1 - p_A - q_A)(1 - p_B - q_B)}
 \end{aligned} \right\} \quad (7)$$

Since $0 \leq K, L < 1$ this is a well-defined transformation on the state space $[0,1]^4$.

Obviously $p_A = q_A = p_B = q_B = \frac{1}{2}$ is a stationary solution of (7) with two zero eigenvalues (corresponding to a very quick convergence to a sex ratio 1:1) and a pair of complex conjugate eigenvalues with real part $+1$ (due to the cyclic structure of Dawkins' game) which make this solution unstable.

On the other hand, using the method of "average Ljapunov functions", developed by Hofbauer and Sigmund (1984) and Hutson and Moran (1982), one can analyze the local behavior near the boundary of the state space, showing that this boundary is repelling, and system (7) exhibits "permanent coexistence" (see Appendix A). The attractors of (7) thus have to lie in the interior of $[0,1]^4$, and since the only stationary solution there is unstable, this attractor has to be a more complicated set.

Numerical simulations show that this attractor is a "limit cycle", i.e. a closed invariant attracting curve surrounding the stationary solution (see Fig. 1). In the next section we will give a mathematical proof of this observation for small values of L, K (this is the biological most relevant case) and determine the explicit equation of this limit cycle in the limit $K = L \rightarrow 0$. This is done by averaging techniques, as, for example, when treating Van der Pol's equation as a perturbation of the harmonic oscillator.

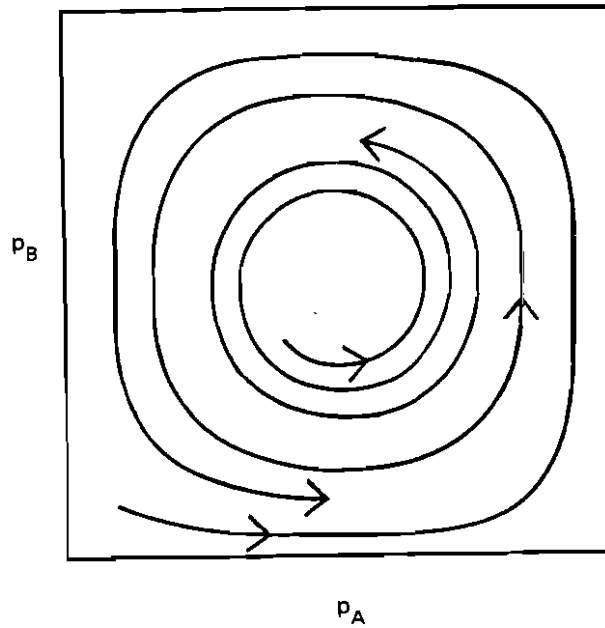


FIGURE 1. The limit cycle.

3. The Limit Cycle

We begin with a linear change of variables which transforms the stationary solution to the origin:

$$x = 1 - p_A - q_A, y = 1 - p_B - q_B, u = p_A - q_A, v = p_B - q_B. \quad (8)$$

Then $(x, y, u, v) \in [-1, +1]^4$.

Furthermore we assume $K = L$ to make the analysis simpler and set $\varepsilon = K/4$. Then (7) is transformed into

$$\begin{aligned} x' &= x - \varepsilon y \frac{1-x^2-u^2}{1-4\varepsilon xy} & y' &= y + \varepsilon x \frac{1-y^2-v^2}{1+4\varepsilon xy} \\ u' &= \varepsilon y \frac{1-x^2-u^2}{1-4\varepsilon xy} & v' &= -\varepsilon x \frac{1-y^2-v^2}{1+4\varepsilon xy}. \end{aligned} \quad (9)$$

Now $|u'|, |v'| \leq \frac{\varepsilon}{1+4\varepsilon}$, and so u and v will be of order $O(\varepsilon)$ for $\varepsilon \rightarrow 0$, uniformly on the whole state-space, in all subsequent generations. (This corresponds to a rough equilibration of the sex-ratio.) Thus the four-dimensional system (9) reduces for small ε to a two-dimensional one:

$$\begin{aligned}x' &= x + \varepsilon f(x, y) + \varepsilon^2 \tilde{f}(x, y) + O(\varepsilon^3) \\y' &= y + \varepsilon g(x, y) + \varepsilon^2 \tilde{g}(x, y) + O(\varepsilon^3)\end{aligned}\tag{10}$$

with

$$\begin{aligned}f(x, y) &= -y(1 - x^2) & \tilde{f}(x, y) &= -4xy^2(1 - x^2) \\g(x, y) &= x(1 - y^2) & \tilde{g}(x, y) &= -4x^2y(1 - y^2).\end{aligned}\tag{11}$$

For $\varepsilon \rightarrow 0$ we obtain, as a first-order approximation, the differential equation

$$\begin{aligned}\dot{x} &= f(x, y) = -y(1 - x^2) \\ \dot{y} &= g(x, y) = x(1 - y^2).\end{aligned}\tag{12}$$

These are - up to a translation of the Nash solution to the origin - just the Schuster-Sigmund equations (3) for Dawkins' game, see also Maynard Smith (1982), Appendix J.

These equations are integrable with

$$I = R^2 = x^2 + y^2 - x^2y^2\tag{13}$$

as a constant of motion. Therefore we change to "canonical coordinates" (R, φ) such that (12) appears in the normal form

$$\dot{z} = iz\omega(|z|^2) = i(a_0z + a_1z|z|^2 + a_2z|z|^4 + \dots)$$

with $z = Re^{i\varphi}$, or equivalently

$$\dot{R} = 0, \quad \dot{\varphi} = \omega(R^2).\tag{14}$$

The canonical angle variable φ and the angle velocity $\omega(R^2)$ are evaluated in Appendix B. They can be expressed only in terms of elliptic integrals,

Now we shall express our difference equation (10) in terms of the canonical variables $I = R^2$ and φ :

$$\begin{aligned}I' &= I(x', y') = I(x + \varepsilon f + \varepsilon^2 \tilde{f} + \dots, y + \varepsilon g + \varepsilon^2 \tilde{g} + \dots) \\ &= I(x, y) + (\varepsilon f + \varepsilon^2 \tilde{f})I_x + (\varepsilon g + \varepsilon^2 \tilde{g})I_y + \frac{1}{2}[(\varepsilon f + \dots)^2 I_{xx} \\ &\quad + 2(\varepsilon f + \dots)(\varepsilon g + \dots)I_{xy} + (\varepsilon g + \dots)^2 I_{yy}] + O(\varepsilon^3).\end{aligned}$$

Now the linear terms in ε vanish since (14) implies $\dot{I} = I_x f + I_y g = 0$. For the ε^2

-terms we obtain from (13)

$$I_x = 2x(1-y^2), \quad I_{xx} = 2(1-y^2), \quad I_{xy} = -4xy, \dots$$

Therefore

$$I' = I + \varepsilon^2 F(x, y) + O(\varepsilon^3) \quad (15)$$

with

$$\begin{aligned} F(x, y) &= \tilde{f}I_x + \tilde{g}I_y + \frac{1}{2}(f^2 I_{xx} + 2fg I_{xy} + g^2 I_{yy}) \\ &= -4xy^2(1-x^2)2x(1-y^2) - 4x^2y(1-y^2)2y(1-x^2) \\ &\quad + y^2(1-x^2)^2(1-y^2) + xy(1-x^2)(1-y^2)4xy + x^2(1-y^2)^2(1-x^2) \\ &= (1-x^2)(1-y^2)[-16x^2y^2 + y^2(1-x^2) + 4x^2y^2 + x^2(1-y^2)] \\ &= (1-R^2)(x^2 + y^2 - 14x^2y^2) \\ &= \frac{1-R^2}{1-x^2}[R^2 - 14R^2x^2 + 13x^4]. \end{aligned} \quad (16)$$

Similarly, using $\dot{\varphi} = \varphi_x f + \varphi_y g = \omega(I)$, we obtain

$$\begin{aligned} \varphi' &= \varphi(x', y') = \varphi(x + \varepsilon f + \dots, y + \varepsilon g + \dots) \\ &= \varphi(x, y) + \varepsilon f \varphi_x + \varepsilon g \varphi_y + O(\varepsilon^2) \\ &= \varphi + \varepsilon \omega(R^2) + O(\varepsilon^2). \end{aligned} \quad (17)$$

So we end up with

$$\begin{aligned} I' &= I + \varepsilon^2 F(I, \varphi) + O(\varepsilon^3) \\ \varphi' &= \varphi + \varepsilon \omega(I) + O(\varepsilon^2). \end{aligned} \quad (18)$$

The next step is to average out the φ -dependence in the ε^2 -term of (18) by the Ansatz $P = I + \varepsilon h(I, \varphi)$. Then

$$\begin{aligned} P' &= I + \varepsilon h(I', \varphi') = I + \varepsilon^2 F(I, \varphi) + \varepsilon h(I + \varepsilon^2 F + \dots, \varphi + \varepsilon \omega(I) + \dots) + \dots \\ &= I + \varepsilon h(I, \varphi) + \varepsilon^2 [F(I, \varphi) + \frac{\partial h}{\partial \varphi} \omega(I)] + O(\varepsilon^3). \end{aligned} \quad (19)$$

Now, following the usual averaging procedure, compare e.g. Arnold (1983), p. 147

ff, we separate F into its mean and its oscillating part,

$$F(I, \varphi) = G(I) + \tilde{F}(I, \varphi)$$

with

$$G(I) = \frac{1}{2\pi} \int_0^{2\pi} F(I, \varphi) d\varphi. \quad (20)$$

Inserting this into (19) we may choose h in such a way that

$$\tilde{F}(I, \varphi) + \frac{\partial h}{\partial \varphi} \omega(I) = 0$$

or

$$\frac{\partial h}{\partial \varphi} = - \frac{\tilde{F}(I, \varphi)}{\omega(I)}. \quad (21)$$

(Since h is a 2π -periodic function in φ , the mean of its derivative, $\frac{1}{2\pi} \int_0^{2\pi} \frac{\partial h}{\partial \varphi} d\varphi$, equals zero. Since this will not in general be the case for the mean of $F(I, \varphi)$ we can average out only the oscillating part $\tilde{F}(I, \varphi)$.)

So with our new variable P instead of I we have simplified (18) to

$$\begin{aligned} P' &= P + \varepsilon^2 G(P) + O(\varepsilon^3) \\ \varphi' &= \varphi + \varepsilon \omega(P) + O(\varepsilon^2). \end{aligned} \quad (22)$$

The averaged function $G(P)$ is computed in Appendix B (in terms of elliptic integrals again), where we also prove the existence of a number P_0 such that $G(0) = G(1) = G(P_0) = 0$

$$\begin{aligned} G(P) &> 0 \quad \text{for } 0 < P < P_0 \\ G(P) &< 0 \quad \text{for } P_0 < P < 1. \end{aligned} \quad (23)$$

Thus the graph of G is similar to that shown in Fig. 2.

If we neglect O -terms in (22), then $P = P_0$ is obviously the equation of an invariant globally attracting circle. That this circle also persists under the O -perturbations for small $\varepsilon > 0$, can be shown e.g. by the technique developed by Iooss (1979) for proving the Hopf bifurcation theorem for maps.

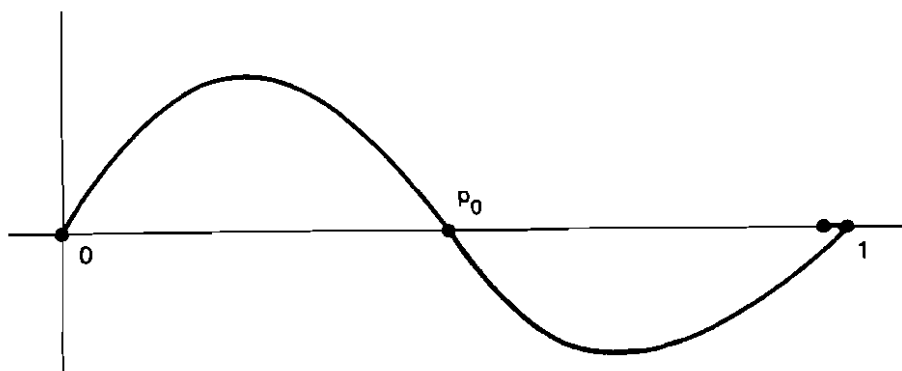


FIGURE 2. The graph of the function $G(P)$.

In particular it is found that for small $\varepsilon > 0$ (10) has an attracting limit cycle of the form

$$P = I + \varepsilon h(I, \varphi) = P_0 + O(\varepsilon^{1/2}). \quad (24)$$

So one particular periodic orbit from the Hamiltonian system (12), i.e. the Schuster-Sigmund dynamics for Dawkins' game, $R^2 = x^2 + y^2 - x^2 y^2 = P_0$, serves as the limiting case $\varepsilon \rightarrow 0$ for the limit cycles of our more refined genetic model of the "battle of the sexes". The constant P_0 which determines the position of this cycle is the zero of the function $G(P)$ and is computed in Appendix B.

4. Discussion

We have described limit cycle behavior in a simple genetic model, with two alleles at each of two loci, with sex-dependence and fitnesses depending on the frequencies of the other sex. The model is based on Dawkins' (1976) "battle of the sexes". For small ε , measuring the intensity of frequency dependent selection, we have proved the existence and stability of a limit cycle. The amplitude of this limit cycle is rather insensitive to changes in the selection intensity, while the "period" is approximately inversely proportional to the selection intensity. This is in contrast to limit cycle behavior constructed by standard Hopf bifurcation techniques where the amplitude is small and the period is approximately constant.

In the limit $\epsilon \rightarrow 0$, this limit cycle tends to one particular cycle of the simple Hamiltonian system identified by Schuster and Sigmund (1981) as the first dynamic model for Dawkins' game. In particular, at least for $\epsilon \rightarrow 0$, the time average over the limit cycle represents the interior stationary solution predicted by game theory. Thus this solution, though unstable both dynamically and evolutionarily, still has some biological relevance. This shows once more that despite their simplicity such haploid models already carry the basic ingredients.

Numerical simulations show that this limit cycle behavior is rather robust. It does not depend on the assumption that fitnesses are additive, or that the selection intensity is the same in the two sexes (i.e. $K = L$). Limit cycle behavior is still observed if there is some directional, non-frequency-dependent selection, provided that this directional selection is less intense than the frequency-dependent selection. Also, the conclusion does not depend on the assumption that the genes are sex-limited in their effects. A similar limit cycle arises if the fitnesses of the genotypes at the A locus, in both sexes, depend additively on the frequencies of the alleles at the B locus, and vice-versa.

All these observations could be proved using the same method of averaging. We have confined ourselves to the simplest possible case, however, in order to keep the mathematical analysis tractable. In more general models difficulties could arise since the nice elliptic integrals have to be replaced by more general integrals where formulas like (B5-7) are probably not available.

That limit cycle behavior is possible in rather simple genetic systems has been shown in the last few years. The pioneering work in this direction is due to Akin (1979). Other papers dealing with this topic are Akin (1982, 1983), Hastings (1981), Hofbauer and Iooss (1984), Hofbauer (1984, 1985), Hunt (1982), Koth and Kemler (1985), Selgrade and Namkoong (1984).

Appendix A: Behaviour Near the Boundary

First it is easy to check that one of p'_A, q'_A, p'_B, q'_B is zero if and only if one of the four alleles A, a, B, b is missing in both sexes. In the new variables (8): $x = \pm 1$ or $y = \pm 1$. Thus the set

$$M = \{ |x| = 1 \text{ or } |y| = 1 \},$$

which is a union of four squares, is the maximal invariant subset of the boundary of the state space: Starting on the boundary outside M , the next generation and all subsequent ones will be in the interior of $[0,1]^4$, i.e. all four alleles will be present in both sexes. Thus in order to show that the boundary is a repelling set, it is sufficient to prove that M is repelling.

For this we first need the behavior on M itself. Since $x' \leq x$ iff $y \geq 0$, it is easy to see that on the part where $y = -1$, x converges monotonically, increasing from -1 to 1 . Where $x = 1$, y goes from -1 to 1 ; where $y = 1$, x goes from 1 back to -1 , and finally where $x = -1$, y decreases from 1 to -1 , and the cycle is closed. Thus the ω - limit set of M consists only of the four fixed points where $|x| = |y| = 1$ (and $u = v = 0$).

Now consider the function

$$P = (1 - x^2)(1 - y^2).$$

Then $P \geq 0$ holds, with $P = 0$ on M . Furthermore $\frac{P'}{P} > \delta > 0$ on the whole state space and

$$\frac{P'}{P} = \left[\frac{1 - 4\epsilon^2}{1 - 16\epsilon^2} \right]^2 > 1$$

at the four boundary fixed points in M . Thus $\frac{P'}{P} > 1$ on the whole ω -limit set of M , and P is an "average Lyapunov function" for M . According to Hutson and Moran (1982), M is a repeller.

Appendix B: Elliptic Integrals in the Battle of the Sexes

We start with the normal form (14)

$$\dot{R} = 0, \quad \dot{\varphi} = \omega(R^2) \tag{B1}$$

of the Schuster-Sigmund differential equation (12) for the battle of the sexes. Since, by (13), we have

$$R^2 = x^2 + y^2 - x^2 y^2, \tag{B2}$$

the state space $\{(x, y) \in \mathbb{R}^2 : |x| < 1\}$ for (12) transforms into $\{(R, \varphi) \in \mathbb{R} \times \mathbb{R} / 2\pi\mathbb{Z} : 0 \leq R < 1\}$. In order to determine the canonical angle φ and the angle velocity $\omega(R^2)$ we treat φ as a function of x and $R: \varphi = \varphi(x, R)$. Then $\dot{\varphi} = \varphi_x \dot{x} + \varphi_R \dot{R}$ implies

$$\omega(R^2) = -\varphi_x y (1 - x^2).$$

Eliminating y from (B2),

$$\frac{d\varphi}{dx} = -\omega(R^2)(1-x^2)^{-\frac{1}{2}}(R^2-x^2)^{-\frac{1}{2}}. \tag{B3}$$

It is well-known that this leads to an elliptic integral. In particular we obtain

$$-\int_{-R}^{+R} \varphi_x dx = -\varphi \Big|_{-R}^{+R} = \pi = \omega(R^2) \int_{-R}^{+R} (1-x^2)^{-\frac{1}{2}}(R^2-x^2)^{-\frac{1}{2}} dx. \tag{B4}$$

Now recall Legendre's formulas for complete elliptic integrals (see e.g. Groebner and Hofreiter (1950), p. 39):

$$\begin{aligned} \int_0^R (1-x^2)^{-\frac{1}{2}}(R^2-x^2)^{-\frac{1}{2}} dx &= \int_0^1 (1-R^2 x^2)^{-\frac{1}{2}} (1-x^2)^{-\frac{1}{2}} dx \\ &= K(R) = \frac{\pi}{2} \sum_{n=0}^{\infty} \left[\frac{\binom{1}{2}_n}{n!} \right]^2 R^{2n} \end{aligned} \tag{B5}$$

$$\int_0^R (1-x^2)^{\frac{1}{2}} (R^2-x^2)^{-\frac{1}{2}} dx = \int_0^1 (1-R^2x^2)^{\frac{1}{2}} (1-x^2)^{-\frac{1}{2}} dx$$

$$= E(R) = \frac{\pi}{2} \sum_{n=0}^{\infty} \frac{\left(\frac{1}{2}\right)_n \left(\frac{1}{2}\right)_n}{n! n!} R^{2n} \quad (B6)$$

$$\int_0^R (1-x^2)^{-3/2} (R^2-x^2)^{-\frac{1}{2}} dx = \int_0^1 (1-R^2x^2)^{-3/2} (1-x^2)^{-\frac{1}{2}} dx$$

$$= (1-R^2)^{-1} E(R) \quad (B7)$$

where $(\alpha)_n = \alpha(\alpha-1)\cdots(\alpha-n+1)$. These series converge for $|R| < 1$. (B4) and (B5) imply

$$\omega(R^2) = \frac{\pi}{2K(R)} = \left[\sum_{n=0}^{\infty} \left[\frac{\left(\frac{1}{2}\right)_n}{n!} \right]^2 R^{2n} \right]^{-1} \quad (B8)$$

The next problem is to evaluate the averaged function (20):

$$G(R^2) = \frac{1}{2\pi} \int_0^{2\pi} F(R^2, \varphi) d\varphi \quad (B9)$$

Since we do not have an explicit expression for φ , we have to transform the integral. Using (B2) and splitting the circle into 4 equal parts we obtain

$$G(R^2) = \frac{2\omega(R^2)}{\pi} \int_0^R \frac{F(R^2, x)}{y(1-x^2)} dx = \frac{2}{\pi} \omega(R^2) \int_0^R F(R^2, x) (R^2-x^2)^{-\frac{1}{2}} (1-x^2)^{-\frac{1}{2}} dx \quad (B10)$$

So again we have to calculate an elliptic integral. From (16) we obtain

$$(1-R^2)^{-1} \int_0^R F(R^2, x) (1-x^2)^{-\frac{1}{2}} (R^2-x^2)^{-\frac{1}{2}} dx = \quad (B11)$$

$$= \int_0^R (R^2-14R^2x^2+13x^4) (1-x^2)^{-3/2} (R^2-x^2)^{-\frac{1}{2}} dx =$$

$$13 \int_0^R (1-x^2)^{\frac{1}{2}} (R^2-x^2)^{-\frac{1}{2}} dx + (14R^2-26) \int_0^R (1-x^2)^{-\frac{1}{2}} (R^2-x^2)^{\frac{1}{2}} dx +$$

$$+ 13(1-R^2) \int_0^R (1-x^2)^{-3/2} (R^2-x^2)^{-\frac{1}{2}} dx =$$

$$= 13E(R) + (14R^2-26)K(R) + 13E(R) =$$

$$= 26E(R) + (14R^2 - 26)K(R) .$$

So (B10) together with (B8) gives an "explicit" formula for G :

$$G(R^2) = (1 - R^2)(14R^2 - 26 + 26 \frac{E(R)}{K(R)}) . \quad (B12)$$

Inserting the series expansions (B5), (B6) into (B11) we obtain after some calculation:

$$\begin{aligned} G(R^2) &= \omega(R^2)(1 - R^2) \sum_{n=1}^{\infty} [14n - 13(2n - 3)^2] \left(\frac{1}{2}\right)_{n-1}^2 \frac{R^{2n}}{n!(n-1)!} = \quad (B13) \\ &= \omega(R^2)(1 - R^2) \left[R^2 + \frac{15}{8}R^4 - \frac{25}{64}R^6 - \dots \right] . \end{aligned}$$

Obviously all term (up to the first two) of the sum have a negative coefficient. Since, furthermore, $G(R^2)(1 - R^2)^{-1}|_{R=1} = -12 < 0$ from (B12), the function $G(P)$ has a unique zero P_0 in the interval $(0,1)$ and is therefore of the desired form (23). The numerical value of P_0 can be found approximately from a table of elliptic functions: $P_0 \approx 0.46$.

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