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Lagos, F., Schreiber, M.R., Parsons, S.G. orcid.org/0000-0002-2695-2654 et al. (2 more authors) (2020) Most EL CVn systems are inner binaries of hierarchical triples. Monthly Notices of the Royal Astronomical Society: Letters. slaa164. ISSN 1745-3925
https://doi.org/10.1093/mnrasl/slaa164

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# Most EL CVn systems are inner binaries of hierarchical triples 

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Accepted XXX. Received YYY; in original form ZZZ


#### Abstract

In spite of their importance for modern astronomy, we do not fully understand how close binary stars containing at least one white dwarf form from main sequence binary stars. The discovery of ELCVn binaries, close pre-white dwarfs with A/F main sequence star companions, offers now the unique possibility to test models of close compact binary star formation. Binary evolution theories predict that these ELCVn stars descend from very close main sequence binaries with orbital periods shorter than 3 days. If this is correct, nearly all ELCVn stars should be inner binaries of hierarchical triples because more than 95 per cent of very close main sequence binaries (the alleged progenitor systems) are found to be hierarchical triples. We here present SPHERE/IRDIS observations of five ELCVn binaries, finding in all of them tertiary objects, as predicted. We conclude that ELCVn systems are inner binaries of hierarchical triples and indeed descend from very close main sequence binaries that experience stable mass transfer.


Key words: binaries: close - white dwarfs - stars: evolution

## 1 INTRODUCTION

A large variety of close white dwarf binaries exist, including objects as interesting as super soft X-ray sources, catalcysmic variables (CVs) and their detached progenitors (e.g. Schreiber \& Gänsicke 2003) or double white dwarf binaries (Han 1998). For all these close white dwarf binaries, theories for their formation and evolution, in most cases involving common envelope evolution (Webbink 1984; Zorotovic et al. 2010), have been developed.

However, as our prescriptions for mass transfer interactions and angular momentum loss are relatively simple conservation equations containing a number of neither theoretically nor observationally well constrained parameters, binary population simulations are unable to reliably predict detailed characteristics of most white dwarf binary populations. In fact, in several cases, the predictions of theoretical binary population models strongly disagree with the observations. To provide just three examples, current models are unable to produce double degenerate systems via a combination of two common envelope phases (Nelemans \& Tout 2005) and require either stable mass transfer or additional energy sources contributing during common envelope evolution (Webbink 2008). No model for supernovae Ia, the thermonuclear explosion of a white dwarf growing in mass through binary star interactions, is able to reproduce the observed delay time distribution (e.g. Yungelson \& Kuranov

[^0]2017). Even worse, in CVs several predictions of the standard model for their evolution drastically disagree with the observations (e.g. Zorotovic et al. 2011). While some of these problems are solved in a recently suggested revision of the model (Schreiber et al. 2016), others remain (Pala et al. 2017; Belloni et al. 2020).

To test and eventually calibrate theoretical models of close white dwarf binary formation it would be ideal to have reliable and concrete information about the main sequence binary progenitors for a given type of white dwarf binary. In general, this is illusory. However, here we present, to the best of our knowledge, the first case where such information is indeed available: ELCVn binaries.

ELCVn-type binaries are eclipsing binaries that contain an A- or F-type star and a very low mass (typically 0.2 solar masses) helium white dwarf precursor (pre-Helium WD). The orbital periods of ELCVn stars are typically very short, i.e. 1-3. days. A number of these compact binary stars have been discovered in the Wide Angle Search for Planets, Kepler photometric surveys, or with the Palomar Transient Factory (Maxted et al. 2014a; van Roestel et al. 2018), and have been found to show different types of pulsations (e.g. Maxted et al. 2014b).

What makes ELCVn binaries special for close white dwarf binary evolution theories is the following. While most white dwarfs and pre-white dwarfs in close binaries are assumed to have formed through common envelope evolution, according to current theories, ELCVn stars must form from dynamically stable mass transfer when the more massive star of the initial main sequence binary was
at the end of its main sequence lifetime or just entered the sub-giant branch (Chen et al. 2017). This condition is required as otherwise the core of the more massive star will grow to masses exceeding measurements of pre-white dwarfs in ELCVn binaries. Such an early start of mass transfer implies that the orbital period of the progenitor binary star system must have been shorter than $\sim 3$ days, as the radius of the primary star was still close to that of a main sequence star when mass transfer began. For longer orbital periods, the sub-giant star is simply too small to fill its Roche-lobe.

The prediction of very short orbital periods for the progenitor main sequence binary stars has fascinating consequences. Virtually all close main sequence stars with orbital periods below $\sim 3$ days are known to be the inner binaries of hierarchical triple systems (Tokovinin et al. 2006). Even if the mass transfer of the inner binary was not conservative and a certain amount of mass was expelled from the inner binary (thereby increasing the orbit of the tertiary), the third companions should still be there after the mass transfer phase. This implies that, if current theories for the formation of ELCVn stars are correct, virtually all ELCVn stars must be the inner binaries of hierarchical triples.

We here present SPHERE (Beuzit et al. 2019) observations of five ELCVn binaries and indeed find strong candidates for tertiary objects in all of them. We conclude that ELCVn binaries form from the inner and very close binaries of hierarchical triple systems in perfect agreement with the formation channels described in Chen et al. (2017).

## 2 SPHERE OBSERVATIONS AND DATA REDUCTION

The separations of the potential tertiary companions to the inner (ELCVn) binary are expected to be in the rage from 10 to 10000 au (Tokovinin et al. 2006). With the spatial resolution of roughly 100 mas and the field of view of SPHERE/IRDIS (11 x 11 $\operatorname{arcsec}^{2}$, Dohlen et al. 2008), our targets must be at distances between $\approx 200-1800 \mathrm{pc}$ in order to cover a large range of the predicted separations for all our targets. We selected five ELCVn binaries from the sample of Maxted et al. (2014a) with distances between $420-1600 \mathrm{pc}$ and magnitudes in the 2MASS $H$ band between 9.7 and 11.9. We used Gaia distances except for TYC 6736. For the latter, the Gaia measurement implies an unrealistically large distance ( 2080 pc ). In fact, the poorly constrained paralax ( $0.46 \pm 0.23 \mathrm{mas}$ ) derived by Gaia may already indicate the presence of a companion. Instead of the Gaia value, we used the distance from RAVE data release 5 (Kunder et al. 2017) for TYC 6736. Table 1 summarizes the parameters of our targets. Thanks to the extreme contrast of 8-14 mags that can be reached with SPHERE our observations are even sensitive to low-mass $\left(\approx 0.1 \mathrm{M}_{\odot}\right)$ stellar companions to our most distant objects.

The five objects that we selected were observed with the high contrast imager SPHERE between 2019 April 9 and August 13. Acquisition of direct imaging was made with IRDIS in the dual band imaging mode (Vigan et al. 2010) using the broad band $H$ filter $\left(\lambda_{\mathrm{H}}=1625 \pm 290 \mu \mathrm{~m}\right)$. Furthermore, the pupil tracking mode was implemented in order to perform angular differential imaging (ADI, Marois et al. 2006). We used the N-ALC-YJH-S coronagraph.

At the beginning and end of each observing sequence, we obtained flux (with the neutral density filter ND 1 to avoid saturation) and star centre calibration images, adding one more exposure for the latter in the middle of the sequence. A total of 112 science images were taken with an exposure time of 32 s for each target except TYC 6736 for which only 66 science images were taken due to bad

Table 1. Distances, 2MASS $H$ photometry, orbital period, mass ratios and proper motion of all ELCVn stars studied in this work. Orbital periods and mass ratios are obtained from Maxted et al. (2014a). Distances are obtained from Bailer-Jones et al. (2018), except for TYC 6736, whose distance has been obtained from data release 5 of the RAVE survey (Kunder et al. 2017).

| Target | Distance [pc] <br> $\pm$ | $H_{2 \text { MASS }}$ <br> $\pm$ | $\mathrm{P}_{\text {orb }}$ <br> $[\mathrm{d}]$ | $q$ <br> $\pm$ | Proper Mot. <br> $[\mathrm{mas} / \mathrm{yr}]$ |
| :--- | :--- | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |
| TYC 5204-1575-1 | 745.37 | 11.308 | 1.29 | 0.130 | 12.8 |
|  | 46 | 0.022 |  | 0.026 |  |
| TYC 9337-2511-1 | 1572.55 | 11.874 | 1.162 | 0.148 | 9.5 |
|  | 66 | 0.023 |  | 0.026 |  |
| TYC 6736-69-1 | 424 | 9.924 | 2.173 | 0.136 | 13.7 |
|  | 85 | 0.021 |  | 0.041 |  |
| TYC 6631-538-1 | 540.07 | 10.432 | 0.901 | 0.143 | 8.9 |
|  | 13 | 0.027 |  | 0.027 |  |
| TYC 5450-1192-1 | 447.37 | 9.748 | 0.793 | 0.176 | 12.8 |
|  | 8.5 | 0.022 |  | 0.012 |  |

weather. The IRDIS data were first pre-processed (sky background subtraction, flat-fielding, bad-pixels correction) with the vLT/SPHERE python package ${ }^{1}$. The frames were recentred based on star centre exposures using the four satellite spots. After pre-processing, and without any post-processing technique to remove speckle patterns produced by the coronograph, we detected at least one potential companion in all five ELCVn binaries. We used the principal component analysis (PCA) algorithm available in the Vortex Image Processing (vip, Gonzalez et al. 2017) python package to look for fainter companions, only finding one extra candidate in TYC 5450 (TYC 5450 d, see Fig. 1).

Making use of the strong signal of our detections, we derotated and median-collapsed the science images to perform aperture photometry and obtained their relative magnitudes $\Delta H_{\text {mag }}$ with respect to the central binary. Since the object detected in TYC 6736 surpassed the IRDIS saturation threshold, flux calibration images were used to obtain $\Delta H_{\text {mag. }}$. As a complementary approach we also used the negative fake companion technique coupled with a Markov chain Monte Carlo algorithm available in vip to calculate $\Delta H_{\mathrm{mag}}$, without finding any discrepancy between both methods.

## 3 THE TERTIARIES TO EL CVN STARS

In all five cases, our SPHERE observations revealed potential tertiary companions (Fig. 1). In what follows we show that the probability for these detections to be background objects is very small and estimate the possible nature of the companions.

### 3.1 Excluding background contamination

We calculated the probability $\mathrm{P}(\Theta, m)$ for each detection to be a chance alignment with a background source within an angular distance $\Theta$ following Brandner et al. (2000) as

$$
\begin{equation*}
P\left(\Theta, m_{\lim }\right)=1-e^{-\pi \Theta^{2} \rho\left(m_{\lim }\right)} \tag{1}
\end{equation*}
$$

where $\rho(m)$ is the cumulative surface density of background sources down to a limiting magnitude $m_{\text {lim }}$ (i.e. the magnitude of the detections). In order to calculate $\rho\left(m_{\lim }\right)$, we used the Besançon galaxy model (Czekaj, M. A. et al. 2014) to generate a synthetic 2MASS $H$ photometric catalogue of point sources within 1 square degree, centred on the coordinates of each target. To calculate $m_{\text {lim }}$, we

[^1]

Figure 1. SPHERE/IRDIS images of the five ELCVn stars that we observed. In all cases we find at least one potential tertiary. The magnitude difference obtained for the possible companions extend from 1.3 to 14 mags. The probabilities for being background objects are small in all cases with values in the range of $0.0002-3 \%$ (except for TYC 5450 d and e).

Table 2. Measured separations, magnitude differences, apparent magnitudes, projected separations, probabilities for chance alignment with a background source, mass and effective temperature for the seven detections.

| Object | $\Theta$ | $\Delta H_{\text {mag }}$ <br>  <br>  <br> $[\operatorname{arcsec}]$ | Proj. Sep. Hau$]$ <br> 2MASS | $P\left(\Theta, m_{\text {lim }}\right)$ | Mass $\left[\mathrm{M}_{\odot}\right]$ <br> $\mathrm{T}_{\text {eff }}[\mathrm{K}]$ |
| :--- | :--- | :---: | :---: | :---: | :---: |
| TYC 5204 c | 4.32 | $4.988 \pm .018$ | 3193 | 0.023 | $0.3-0.4$ |
|  |  | $16.29 \pm .08$ | 198 |  | $3400-3500$ |
| TYC 9337 c | 3.96 | $6.946 \pm .016$ | 6227 | 0.031 | $0.2-0.3$ |
|  |  | $18.82 \pm .08$ | 261 |  | $3200-3400$ |
| TYC 6736 c | 0.2 | $1.330 \pm .030$ | 88 | $2 \cdot 10^{-6}$ | $\approx 0.9$ |
|  |  | $11.53 \pm .09$ | 18 |  | $5300-5500$ |
| TYC 6631 c | 1.0 | $3.595 \pm .008$ | 544 | $4 \cdot 10^{-4}$ | $\approx 0.6$ |
|  |  | $14.03 \pm .07$ | 13 |  | $3700-4000$ |
| TYC 5450 c | 0.53 | $4.701 \pm .014$ | 235 | $1.6 \cdot 10^{-4}$ | $\approx 0.5$ |
|  |  | $14.45 \pm .07$ | 5 |  | $\approx 3700$ |
| TYC 5450 d | 3.36 | $13.86 \pm .15$ | 1503 | 0.089 | $<0.07$ |
|  |  | $23.6 \pm .2$ | 28 |  | $<1600$ |
| TYC 5450 e | 5.58 | $10.402 \pm .120$ | 2497 | 0.15 | $<0.07$ |
|  |  | $20.15 \pm .18$ | 47 |  | $<1600$ |

assumed that both SPHERE and 2MASS $H$ filters are identical, so that the relation between the SPHERE $\Delta H_{\text {mag }}$ of our candidates and their apparent 2MASS $H$ magnitudes ( $H_{2 \text { MASS }}$ ) is given by
$H_{2 \mathrm{MASS}}=H_{\text {arch }, 2 \mathrm{MASS}}+2.5 \log (1+\alpha)$,
where $H_{\text {arch,2MASS }}$ is the 2MASS archival magnitude of the ELCVn and $\alpha=10^{0.4 \Delta H_{\text {mag }}}$ the value of $\Delta H_{\text {mag }}$ expressed in terms of the counts ratio between the EL CVn and the candidate. To avoid underestimating the value of $\mathrm{P}\left(\Theta, m_{\mathrm{lim}}\right)$, we added 0.5 magnitudes to $m_{\mathrm{lim}}$. We note that discrepancies between 2MASS and SPHERE $H$ bands are clearly irrelevant. We used VOSA's (Bayo et al. 2008) synthetic photometry tool coupled with the BT-NextGen spectral library (Allard et al. 2012), finding differences of 0.06 magnitudes at most. Table 2 lists the final apparent magnitudes and probabilities for being a background object for each candidate. Combining the obtained probabilities we derive a probability for all five ELCVn binaries to be triple systems or higher multiples of 0.95 , i.e. with $2 \sigma$ significance we measured that five out of five ELCVn stars are triples or higher multiples.

For any assumed true triple fraction of ELCVn stars we can now calculate the significance for rejecting this hypothesis. Given
an assumed triple fraction $p$, the probability of measuring five out of five is simply $p^{5}$. This implies that we can say with one (two) sigma significance that the true fraction of ELCVn stars exceeds 80 (55) per cent.

In Fig. 2 we compare our result with the triple fraction of sunlike stars as a function of orbital period. Apparently, the measured high fraction of tertiaries in ELCVn binaries corresponds to those of close binary stars with periods below 3 days, in perfect agreement with theoretical predictions.

In general one could further strengthen this result by confirming common proper motions of the candidates with respect to the central ELCVn binary. According to the proper motions listed in Table 2 and the IRDIS pixel scale ( 12.25 mas), the time between observations would need to be at least 3 yr. However, such second epoch observations would only slightly improve an already highly significant result given that for two objects independent evidence for the companion nature of our candidates is already provided by Gaia. As mentioned previously, the large uncertainty and unrealistic value of the Gaia paralax of TYC 6736 is consistent with the presence of a companion. In the case of TYC 5204, Gaia even detected the companion and, despite relatively large uncertainties, the Gaia astrometric solution is consistent with common proper motion. Given this independent evidence, and the very low background probabilities, our result would remain nearly unchanged even if the Galaxy model we used was significantly underestimating the background object probability.

### 3.2 Potential nature of the companions

The optical and IR emission of ELCVn binaries is dominated by their A (or F) star. If our detections are physically bound to the EL CVn binary, we would expect they were later G, K or M type dwarfs. Although a full characterization of the potential companions is impossible given the limited data currently available, we can derive a rough estimate of their masses and temperatures by using isochrones for main sequence low-mass stars (Baraffe et al. 2015).

Chen et al. (2017) showed that EL CVn binaries with orbital periods shorter than 2.2 days descend from main sequence binaries
where the most likely mass for the pre-Helium white dwarf progenitor is in the range $1.15-1.2 \mathrm{M}_{\odot}$. As the mass transfer that led to the formation of the ELCVn system started after/at the end of main sequence evolution, the age of our systems should be at least $\approx 5-7$ Gyr. By comparing the derived distance-modulus corrected $H_{\text {2MASS }}$ magnitude with those values in the isochrones between 5 and 10 Gyr , the companions for TYC 5204, TYC 9337, TYC 5450 and TYC 6631 are consistent with with M and K type stars, while for TYC 6736 the closest companion is likely a G type star (Table 2).

## 4 DISCUSSION

Binary evolution theory predicts that ELCVn stars descend from very close main sequence binary stars. Multiplicity surveys show that practically all such close main sequence binaries host a distant tertiary, i.e. they are the inner binaries of triple systems. Our observations with SPHERE/IRDIS of five EL CVn confirm this prediction: we find strong companion candidates in all five systems. In what follows we discuss the implications of our finding on the evolutionary history and the future of ELCVn stars.

### 4.1 The evolutionary history of EL CVn triples

ELCVn binaries descend from hierarchical main sequence triple star systems where the inner binary has a very short orbital period of $\lesssim 3$ days. The formation of such close binaries has been a mystery for decades as they cannot form directly via fragmentation of molecular clouds or protostellar discs. Fabrycky \& Tremaine (2007) and Naoz \& Fabrycky (2014) proposed that Kozai-Lidov oscillations coupled with tidal friction could be responsible for the formation of these close binaries. However, Moe \& Kratter (2018) showed that the close binary fractions of very young stars is consistent with those of field stars and that Kozai-Lidov oscillations are too slow to reproduce the observations. Instead, they suggest that energy dissipation due to interactions with the primordial disc during the pre-main sequence is the main mechanism driving the formation of inner binaries with periods less than 10 days.

Key information concerning the origin of very close binary stars may come from combining statistics of stellar multiples with those of the most massive planets and brown dwarfs. Recently, Fontanive et al. (2019) showed that the fraction of tertiaries in systems with close substellar companions is even higher than in the case of stellar binaries. We compare the triple fraction of ELCVn stars with both samples as a function of orbital period in Fig.2. As also noted by Moe \& Kratter (2019), the similar large fractions of systems with a distant tertiary indicate a similar formation mechanism which is most likely disc fragmentation followed by disc migration. Why this disc fragmentation and subsequent migration occurs nearly only in hierarchical triples is not entirely clear. Either the total initial mass required for this to work is large enough that it is virtually always accompanied by core fragmentation, or a close companion formed through core fragmentation is triggering disc fragmentation at an early stage. In any case, we conclude that according to current theories of close binary formation, ELCVn systems are triple systems with a close inner binary formed through disk fragmentation and subsequent migration. After the main sequence triple has formed, it is assumed to evolve quietly until the more massive star of the inner binary moves off the main sequence and crosses the Hertzsprung gap. This star then fills its Roche-lobe before entering the first giant branch and thermal time scale mass transfer starts. It depends then on the thermal time scales and on


Figure 2. Triple fraction for stellar binaries (blue) and hot Jupiter/Brown dwarfs (green) from Tokovinin et al. (2006) and Fontanive et al. (2019). In both cases the fraction of systems with tertiary companions is very large, suggesting that close binaries (including those with massive Jupiter or brown dwarf companion) form through disc fragmentation (see also Moe \& Kratter 2019). The teal and red shaded regions represent the $1 \sigma$ and $2 \sigma$ significance level for rejecting the hypothesis that the tertiary fraction of ELCVn stars is less than 80 and 55 per cent respectively. The large triple fraction clearly implies that ELCVn stars descend mostly from close inner binaries of hierarchical triples.
the masses of both stars that to what degree this dynamically stable mass transfer remains conservative and how much the orbital period is increased during the mass transfer phase. Mass transfer stops when the envelope of the initially more massive star has been fully stripped off. The result is an inner binary consisting of the exposed low-mass core and a secondary that is more massive than it initially was (Chen et al. 2017). The tertiary is affected little by the evolution of the inner binary. In most cases it simply spirals out depending on the total mass lost by the inner binary: an ELCVn triple system, as the ones we observed, is born.

### 4.2 The future of EL CVn triples

The evolutionary phase the triple system appears as an ELCVn binary with a distant tertiary is probably relatively short. Given the short periods of the ELCVn binaries and in particular of the sample observed by us (see Table 1), a second phase of mass transfer is unavoidable. For shorter periods the main sequence star either fills its Roche-lobe while still on the main sequence or in the Hertzsprung gap. In both cases it will not have a deep convective envelope and mass transfer will be stable for mass ratios as large as $3-5$ (see Ge et al. 2015, their Fig. 8 and 9). Nevertheless, the mass ratios of the ELCVn stars we observed exceed this limit (table 1 , using $1 / q$ for a proper comparison) and it is therefore likely that the second phase of mass transfer will be dynamically unstable and lead to common envelope evolution. As the orbital energy that is available will not be sufficient to expel the entire envelope of the main sequence star, both stars will merge, hydrogen burning will start around the core until the envelope will be expelled. In other words, the triple systems we currently observe will evolve into wide binary systems with one component being a white dwarf. If a significant part of the secondary is expelled during this merger process, the white dwarf may not significantly grow in mass and remain in the lowmass white dwarf regime. Interestingly, such a system has recently been observed (Vos et al. 2018) and the authors concluded that it
most likely descended from a triple, in perfect agreement with our prediction for the future of close ELCVn binaries.

ELCVn stars with slightly smaller mass ratios than the systems we observed will start thermal time scale mass transfer and evolve into double degenerate stars consisting of two low-mass white dwarfs. These ELCVn stars might therefore be the progenitors of observed double white dwarfs (Bours et al. 2014; Parsons et al. 2020). EL CVn stars with significantly longer periods have not been observed but predicted to exist (though in small numbers). These objects will start mass transfer when the main sequence star evolved into a giant star which will lead to common envelope evolution. Due to the larger initial orbital period, these systems may survive the common envelope phase and form very close double degenerates. For both scenarios, the amount of mass lost during mass transfer or common envelope evolution would be, in most cases, not enough to unbound the third object, even in an impulsive massloss regime (Veras et al. 2011), i.e. the distant tertiary will move to a larger separation but remains bound to the system. As ELCVn systems are likely to evolve into systems containing extremely lowmass (ELM) white dwarfs, the sample established by the ELM survey (e.g. Brown et al. 2010) might contain a significant number of triple systems. The outlined future evolution of ELCVn stars and the origin of ELM white dwarfs can be tested by searching for these tertiaries. To predict the precise relative number of such systems, triple population synthesis models are required which we intend to present in an upcoming publication.

## 5 CONCLUSIONS

The characteristic configuration of ELCVn binaries, close A-F type dwarfs with a very low-mass pre-white dwarf evolving towards higher effective temperatures at nearly constant luminosity, offers a unique way to test the latest white dwarf binary formation theories. The narrow parameter space predicted for the progenitor systems, i.e. orbital periods below $\sim 3$ days and white dwarf progenitors with masses $\simeq 1.15-1.20 \mathrm{M}_{\odot}$, implies that (if theories are correct) nearly all ELCVn binaries must be the inner binary stars of hierarchical triples because virtually all main sequence binary stars with periods shorter than 3 days are known to be the inner binaries of such triples.

We performed SPHERE/IRDIS observations of five ELCVn stars and indeed found very strong companion candidates in all five systems. Our results represent a unique and independent confirmation of the predictions of formation theories for ELCVn stars. Discussing the future of ELCVn binaries we found that ELCVn stars and their tertiaries either evolve into wide binaries with a low-mass white dwarf component or into triples with inner binaries consisting of at least one extremely low-mass (ELM) white dwarf. EL CVn triples therefore represent a link between hierarchical main sequence triples with very close inner binary stars and a sub-population of systems containing ELM white dwarfs.

## ACKNOWLEDGEMENTS

FL, MRS, and NG thank for support from the ICM Millennium Nucleus for Planet Formation, NPF. FL is supported by an ESO studentship and MRS by FONDECYT (1181404). SGP acknowledges support from the STFC Ernest Rutherford Fellowship. BTG was supported by the UK STFC grant ST/P000495. NG acknowledges support from project CONICYT-PFCHA/Doctorado Nacional/2017
folio 21170650 . The data presented in this work have been obtained through ESO program 0103.D-0346(A).

## DATA AVAILABILITY

The data underlying this letter will be shared on reasonable request to the corresponding author.

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