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# Multifractal structure of the monthly rainfall regime in Catalonia (NE Spain): Evaluation of the non-linear structural complexity

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## ABSTRACT

The complex non-linear regime of the monthly rainfall in Catalonia (NE Spain) is analyzed by means of the reconstruction fractal theorem and the multifractal detrended fluctuation analysis algorithm. Areas with a notable degree of complex physical mechanisms are detected by using the concepts of persistence (Hurst exponent), complexity (embedding dimension), predictive uncertainty (Lyapunov exponents), loss of memory of the mechanism (Kolmogorov exponent), and the set of multifractal parameters (Hölder exponents, spectral asymmetry, spectral width, and complexity index). Besides these analyses permitting a detailed description of monthly rainfall pattern characteristics, the obtained results should also be relevant for new research studies concerning monthly amounts forecasting at a monthly scale. On one hand, the number of necessary monthly data for autoregressive processes could change with the complexity of the multifractal structure of the monthly rainfall regime. On the other hand, the discrepancies between real monthly amounts and those generated by some autoregressive algorithms could be related to some parameters of the reconstruction fractal theorem, such as the Lyapunov and Kolmogorov exponents.

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The monthly rainfall regime in Catalonia, NE Spain, is analyzed by means of the fractal theory with the aim of improving the knowledge about its complex physical mechanism.

(1998), Talkner and Weber (2000), García-Marín *et al.* (2013; 2019), Rodríguez *et al.* (2013), Burgueño *et al.* (2014), Lana *et al.* (2015; 2016), and Herrera-Grimaldi *et al.* (2019), among others, can be cited.

By focusing the analysis on the characteristics of a rainfall regime, a detailed analysis of its complexity could permit a better detection of regions where rainfall amount forecasting is quite easy or, alternatively, difficult, and of high uncertainty. The complexity of a rainfall regime could be quantitatively evaluated by fractal and multifractal theories. The application of the reconstruction fractal theorem (RFT) (Diks, 1999) to monthly amount series permits one to quantify several concepts such as persistence or randomness of the time series (Hurst exponent), predictive uncertainty (Lyapunov exponents), complexity of the physical process (embedding dimension), and loss of memory of the physical system (Kolmogorov entropy). The multifractal behavior, determined by the multifractal detrended analysis (MDFA) (Kantelhardt *et al.*, 2002), quantifies the complexity of the physical mechanism bearing in mind the central, maximum, and minimum Hölder exponents; spectral amplitude and asymmetry; as well as a complexity index that summarizes these multifractal parameters.

## I. INTRODUCTION

The fractal structure of nature phenomena can be successfully analyzed by taking Mandelbrot (1983) as a reference. More concretely, the physical mechanisms of these phenomena can be studied by taking into account concepts such as predictive instability, degree of complexity, and loss of memory of the physical mechanism (Turcotte, 1997; Diks, 1999; and Dimr, 2005). Another relevant point of view is the analyses of these phenomena by means of the multifractal theory (Goltz, 1997) and, more concretely, by means of specific algorithms applied to time series (Kantelhardt *et al.*, 2002). Some examples of these fractal and multifractal analyses on solid Earth sciences (seismology and tectonics) are Hirabayashi *et al.* (1992), Godano *et al.* (1996), Enescu *et al.* (2005), and Ozturk (2012), among others. With respect to dynamic atmospheric and climatology, Koscielny-Bunde *et al.*

59 The main objective of this study is a detailed analysis of the spatial  
60 distribution of fractal and multifractal parameters characterizing  
61 the complexity of the monthly pluviometric regime of a Mediter-  
62 ranean area, Catalonia, where a homogeneous *Mediterranean cli-*  
63 *mate* is not expected due to its varied and complex topography. As  
64 a consequence, advantages and shortcomings concerning monthly  
65 amounts forecasting and possible time trends on monthly amounts  
66 could be relevant. On one hand, different levels of the monthly  
67 regime complexity could be detected for the different climatic areas  
68 of Catalonia. On the other hand, monthly rainfall forecasting uncer-  
69 tainties could be expected with a certain degree of veracity. The  
70 randomness or persistence of the monthly pluviometric series is  
71 quantified by the Hurst exponent. The complexity of the physical  
72 process governing the pluviometry of the different climatic areas  
73 is characterized by the embedding dimension concept (number of  
74 non-linear equations to describe the physical mechanism), and the  
75 uncertainty degree on forecasting monthly amounts is assumed to  
76 be strongly related to the Lyapunov exponents. Additionally, the  
77 Kolmogorv entropy permits quantification of the loss of memory  
78 of the physical mechanism, this question being a relevant factor  
79 if autoregressive processes are applied to forecast monthly rainfall  
80 amounts. The complexity of the rainfall regime is described from  
81 different points of view based on the multifractal spectrum (maxi-  
82 mum, minimum, and central Hölder exponents; spectral amplitude;  
83 and spectral asymmetry) all of them being summarized by a com-  
84 plexity index. The results obtained in these fractal and multifractal  
85 analyses permit a detailed description, at a local scale, of the phys-  
86 ical mechanisms complexity governing the pluviometry of an area  
87 of varied topography. In short, improvements on monthly rainfall  
88 forecasting and validation of monthly amount time trends should  
89 be expected.

90 The contents of this paper are organized as follows. The quality  
91 of the dataset and the recording continuity are discussed in Sec. II.  
92 The RFT theorem and the MDFA algorithm are described with  
93 detail, being also explained the meaning of the corresponding frac-  
94 tals parameters in terms of monthly rainfall regimes, in Sec. III. The  
95 obtained results are introduced in detail in Sec. IV and the rele-  
96 vance of them, with possible improvements on rainfall forecasting at  
97 a monthly scale and verification of possible time trends on monthly  
98 amounts, are summarized in Sec. V.

## 99 II. STUDY AREA AND DATABASE

100 The area of the Iberian Peninsula corresponding to Catalonia  
101 (NE Spain) should be identified as belonging to a Mediterranean cli-  
102 mate in agreement with the Köppen–Geiger classification (Mc Night  
103 and Hess, 2000), bearing in mind its latitude on the Northern Hemi-  
104 sphere and the characteristics of several climatic variables. Neverthe-  
105 less, some relevant differences on thermometric and pluviometric  
106 regimes within the territory have to be taken into consideration due  
107 to the relatively complex orography (Fig. 1) and the temperate effects  
108 of the Mediterranean Sea on areas close to the Mediterranean coast,  
109 whereas the littoral fringe is characterized by moderate low tempera-  
110 tures in winter and notable hot episodes in summer, the inner areas,  
111 delimited by the Pre-Littoral chains, the Eastern Pyrenees, and the  
112 Ebro valley, are characterized by cold winters and very hot sum-  
113 mers. Annual average temperatures vary from 0°C (Pyrenees) to

114 17°C (inner areas), and the extreme temperature records are within  
115 a very wide range of 75° (maximum 43°C for inner Catalonia and  
116 minimum −32° for the Pyrenees). With respect to the pluviometric  
117 regime, the complex topography also contributes to different spatial  
118 distributions of rainfall amounts, and three different domains could  
119 be assumed, whereas Pyrenees and Pre-Pyrenees areas are character-  
120 ized by average annual amounts close to or exceeding 1000 mm/yr,  
121 sometimes with copious snow episodes (especially for the north  
122 face of the Pyrenees), some places of the Pre-Littoral chain achieve  
123 amounts close to or slightly exceeding 700 mm/yr, and the records  
124 of the rest of territory (inner Catalonia and Littoral fringe) range  
125 from 300 to 700 mm/yr. Consequently, the expected equinoctial  
126 thermometric and rainfall regimes are only explicitly accomplished  
127 in the Littoral fringe and some inner domains, being a relevant  
128 factor for the rest of the territory the Pre-Littoral and Eastern Pyre-  
129 nees chains. In short, bearing in mind rainfall temperatures and  
130 wind regimes, five different climate domains, all of them within  
131 a generic *Mediterranean climate*, can be established for Catalonia.  
132 These five domains would be that corresponding to Littoral fringe,  
133 Inner territory, Mediterranean-mountain, transition from Mediter-  
134 ranean to Atlantic domain, and Atlantic Mountain (the last one a  
135 small domain in the north face of the Pyrenees). Previous recent  
136 analyses of pluviometry and thermometric regimes in Catalonia  
137 can be found in Burgueño *et al.* (2014), Lana *et al.* (2016), and  
138 Casas-Castillo *et al.* (2018), among others, who have recently ana-  
139 lyzed them at local and regional scales. A very detailed description  
140 of thermometric and pluviometric regimes can be also found in  
141 Clavero *et al.* (1996).

142 Due to the relative complex orography of Catalonia (Fig. 1),  
143 with altitudes about sea level achieving 2900 m in the Pyrenees,

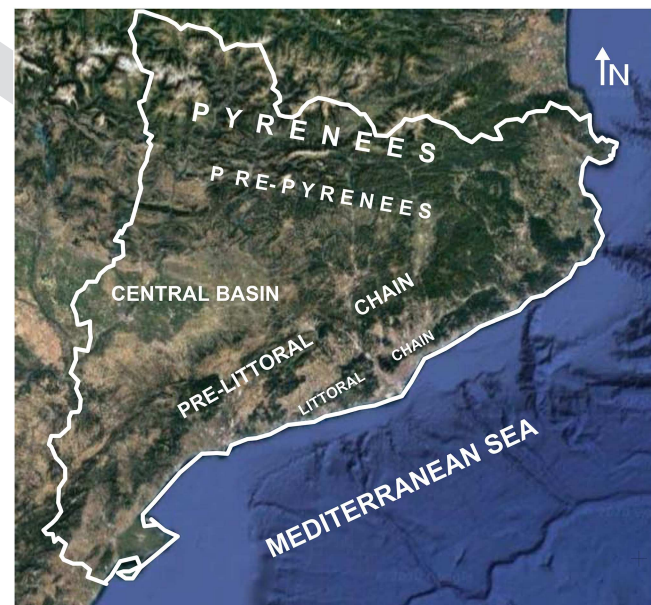


FIG. 1. Topographic image with the most relevant orographic elements (Eastern Pyrenees, Central Basin, and Littoral and Pre-Littoral chains).

144 2000 m on the Pre-Pyrenees, varying from 900 to 1500 m in the Pre-  
 145 Littoral chain, close to 500–600 m in the Central Basin, and only a  
 146 few hundred meters on the Littoral chain, a pluviometric network  
 147 as dense as possible is necessary. Monthly rainfall amount series  
 148 have been collected from Servei Meteorològic de Catalunya (SMC,  
 149 [www.meteo.cat](http://www.meteo.cat)), the meteorological agency of Catalonia, including  
 150 records of two other organisms: Fabra Observatory (Reial Acadèmia  
 151 de Ciències i Arts, RACA, Barcelona) and Ebre Observatory (Ramon  
 152 Llull University, URL, and CSIC, Consejo Superior de Investiga-  
 153 ciones Científicas, CSIC, Spanish Government). A set of 96 monthly  
 154 amount series [Fig. 2(a)] have been selected, offering the emplace-  
 155 ment of the rain gauges a relatively dense spatial distribution and,  
 156 at the same time, accomplishing relevant conditions. First, several  
 157 tests of homogeneity and data quality were applied at daily scale by  
 158 [Llabrés-Brustenga et al. \(2019\)](#) to distinguish the acceptable rain-  
 159 fall series. Second, for very short lags of few days without records,  
 160 unknown daily rain amounts have been substituted by those gener-  
 161 ated by a krigging process ([Stein, 1999](#) and [Press et al., 2007](#)),  
 162 taking into account daily records close to the gauge without data and  
 163 the topography of the area including these emplacements have been  
 164 taken in this krigging process. Third, an appropriate length (40 years  
 165 interval) for a right computation of fractal parameters has been con-  
 166 sidered. The results of applying these constraints are illustrated in  
 167 Fig. 2(b), where the chosen 96 rainfall records represent 80%–100%  
 168 of the available data for 1960–2000.

169 A first description of the rainfall regime is shown in Figs. 3(a)  
 170 and 3(b), where the spatial distribution of the average annual  
 171 rain amounts and the standard deviations are represented. The  
 172 differences on the average amounts at annual scale are quite  
 173 evident by comparing Littoral and Pre-Littoral emplacements  
 174 (500–800 mm/yr), areas of the Central basin (300–500 mm/yr), as  
 175 well as emplacements close to the Eastern Pyrenees and the Pyre-  
 176 nees chain itself (800–1000 mm/yr and 1100 mm/yr in a few places).  
 177 It is also remarkable the high standard deviations for many areas,  
 178 including the Pyrenees area, being this fact in agreement with the  
 179 expected irregularity of a Mediterranean rainfall regime. In fact, the  
 180 spatial patterns of the geographical distribution of average and stan-  
 181 dard deviation are quite similar, with a quite evident tendency to  
 182 increase the standard deviations with the average amounts.

### 183 III. FRACTAL AND MULTIFRACTAL THEORY

#### 184 A. Reconstruction theorem

##### 185 1. Rescaled-range analysis

186 A first step on a fractal analysis of time series could be the pro-  
 187 cess designed as the rescaled-range analysis ([Korvin, 1992](#)), which is  
 188 quantified by the power law,

$$\frac{R(\tau)}{S(\tau)} \propto \tau^H, \quad (1)$$

189 with  $R(\tau)$  and  $S(\tau)$  being the range of variation and standard deviation  
 190 respectively, of segments of length  $\tau$  and  $H$ , the Hurst exponent  
 191 of Eq. (1). If this power-law equation is well accomplished, the time  
 192 series could be qualified as random ( $H$  very close to 0.5), persis-  
 193 tent ( $H$  notably exceeding 0.5), or anti-persistent ( $H$  clearly lowering

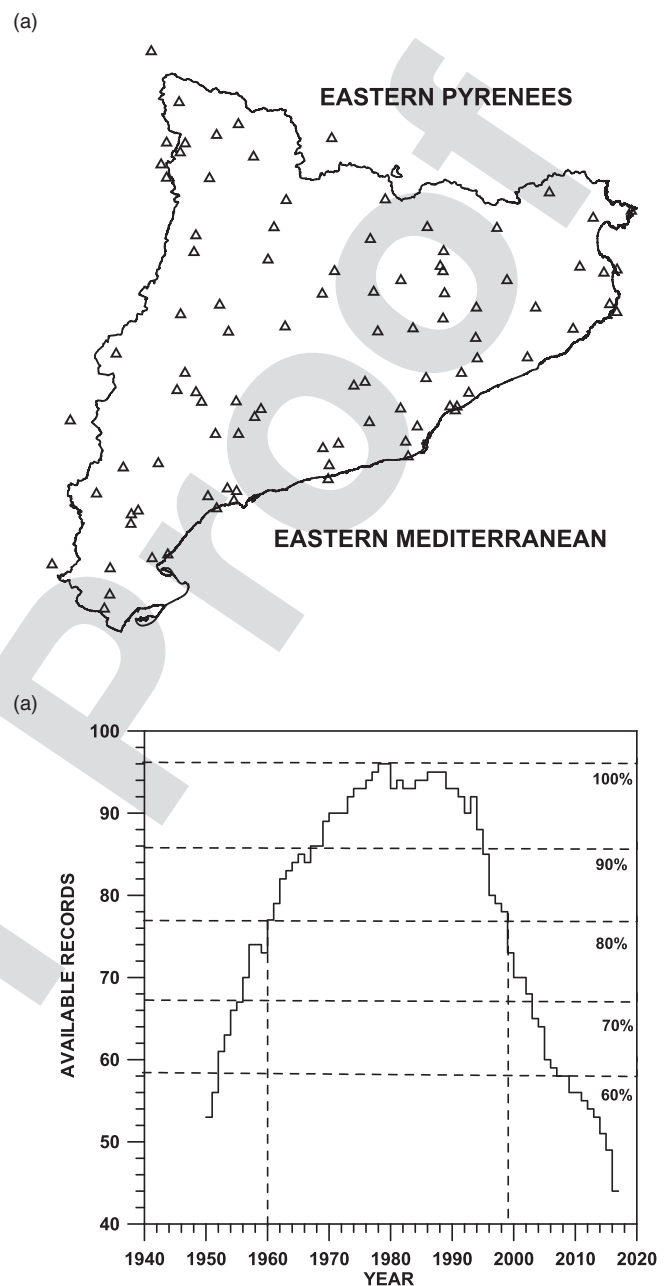


FIG. 2. (a) Pluviometric network with the location of the 96 rain gauges (open triangles). (b) Annual evolution of the number of available gauges. Vertical dashed lines limit the chosen 40 years interval for the analysis.

0.5). Besides a possible way of self-similar/affine character verifica- 194  
 tion of the series ([Turcotte, 1997](#)), the numeric values of  $H$  permit to 195  
 decide if the randomness of the series discourages forecasting strate- 196  
 gies or the persistent/anti-persistent character facilitates successful 197  
 application of forecasting algorithms. 198



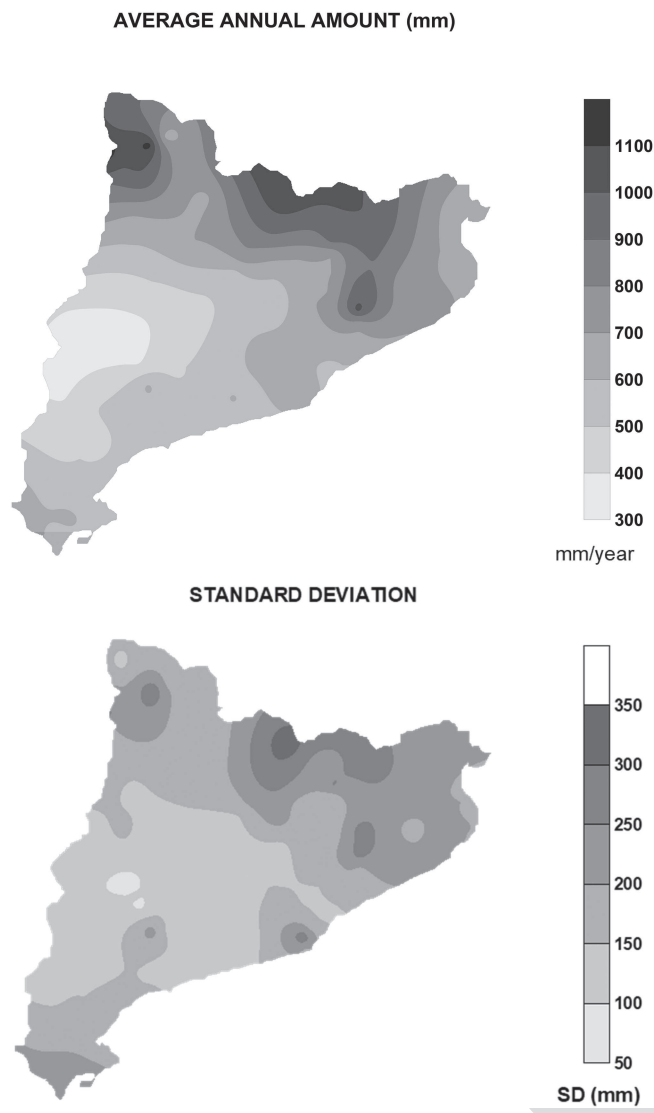


FIG. 3. Spatial distribution of the average annual rainfall amounts and their standard deviations.

199 **2. Reconstruction fractal theorem**

200 The reconstruction theorem (Diks, 1999) permits one to quan-  
 201 tify the complexity and predictive instability of a physical process.  
 202 With respect to the complexity and possible chaotic behaviors, the  
 203 fundamental parameters are the minimum number of nonlinear  
 204 equations associated with a physical mechanism, the embedding  
 205 dimension,  $d_E$ , necessary to obtain an asymptotic value of the cor-  
 206 relation dimension  $\mu$  and the Kolmogorov entropy,  $\kappa$ , which is a  
 207 measure of the loss of memory of the physical mechanism along the  
 208 process. The mathematical process of the reconstruction theorem is

based on the generation of  $m$ -dimensional space vectors,

$$Z(i) = \{x(i), x(i + 1), \dots, x(i + m - 1)\}; \quad i = 1, \dots, N, \quad (2)$$

with  $\{x(k)\}$  being the set of  $n$  elements of a empiric data series and  $N = n - m + 1$ . In terms of  $Z(i)$ , the correlation integral, following the Grassberger–Procaccia formulation (Grassberger and Procaccia, 1983a; 1983b), is written as

$$C(m, r) = \lim_{N \rightarrow \infty} \frac{1}{N^2} \sum_{i,j=1}^N H\{r - \|Z(i) - Z(j)\|\}, \quad (3)$$

with  $r$  being an Euclidean distance in the  $m$ -dimensional space and  $H\{\cdot\}$ , the Heaviside function. Equation (3) is the starting point to obtain numeric estimations of parameters  $\mu$  and  $\kappa$ . Assuming that the correlation integral can be expressed as  $C(m, r) = A_m e^{-\mu r}$ , with  $A_m$  being the correlation amplitude for the reconstruction dimension  $m$  and plotting this correlation integral in terms of  $r$  on log–log scales,

$$\log\{C(m, r)\} = \log(A_m) - \mu r + \mu(m)\log(r), \quad (4)$$

whose slope  $\mu(m)$  is straightforwardly obtained for every reconstruction dimension  $m$  by means of linear regression on log–log scales. As mentioned in other papers related to this concept of correlation dimension (Burgueño et al., 2014), two factors which could lead to wrong estimations of  $\mu(m)$  have to be carefully revised. On one hand, the phenomenon of lacunarity (Turcotte, 1997), many times detected for small values of  $r$ ; on the other hand, the saturation of  $C(m, r)$  for high values of  $r$ . These shortcomings concerning  $\mu(m)$  can be easily solved by searching for the  $r$  interval for which  $\log\{C(m, r)\}$  is strongly linear dependent on  $\log(r)$  and, at the same time, the highest square regression coefficient is detected. In this way, the slope of the log–log linear dependence becomes a good estimation of  $\mu(m)$ .

After a right quantification of the correlation dimension for every reconstruction dimension  $m$ , it is relevant to observe that  $\mu(m)$  will tend asymptotically to a value,  $\mu^*$ , which is interpreted as the minimum number of non-linear equations describing the mechanism that governs empiric data  $\{x\}$ . In this way,  $\mu^*$  becomes a first evaluation of the complexity of the analyzed time series. It is also relevant that a high maximum dimension  $m$  or, in other words, a necessary high embedding dimension  $d_E$  for achieving the asymptotic value of  $\mu(m)$ , is usually associated with a random behavior of the analyzed time series.

Together with the degree of complexity and a possible random component of the physical mechanism, the loss of memory of the physical system is another valuable parameter. This question is quantified by the Kolmogorov entropy exponent,  $\kappa$ , which can be estimated as follows. According to Eq. (4) and Lana et al. (2010), grouping  $\log\{C(m, r)\} - \mu(m)\log(r)$ , as  $\alpha(m)$ ,

$$\alpha(m) = \log(A_m) - \mu r. \quad (5)$$

This last equation permits a fast and accurate estimation of  $\kappa$  by a least square regression of empiric  $\alpha(m)$  in terms of  $m$ , provided that  $\log(A_m)$  should be a constant. Empirical data confirm this behavior only with  $m$  tending to  $\infty$ , being then assumed that  $A_{m+1}/A_m$  tends to 1.0 for high reconstruction dimensions  $m$ . In other words, correlation amplitudes tend to be very similar for high reconstruction dimensions.

Another relevant application of the reconstruction theorem is the quantification of the predictive instability. After generating  $m$ -dimensional vectors according to Eq. (2), in agreement with Wiggins (2003), the Lyapunov exponents,  $\lambda_j$  ( $j = 1, 2, \dots, m$ ), can be computed according to the algorithms proposed by Eckmann *et al.* (1986) and Stopp and Meier (1988). Assuming that the addition of all the  $m$  Lyapunov exponents is negative, the trajectory in the  $m$ -dimensional space is described by aperiodic orbits around a strange attractor defined by the Kaplan–Yorke dimension, DKY (Kaplan and Yorke, 1979). Quantitatively,

$$D_{KY} = c + \frac{1}{|\lambda_{c+1}|} \sum_{j=1}^c \lambda_j, \quad (6)$$

with  $c$  being the maximum number of positive and negative Lyapunov exponents in a decreasing order, and accomplishing  $\lambda_1 + \lambda_2 + \dots + \lambda_c \geq 0$ .

## B. Multifractal spectrum

The concept of multifractal spectrum offers a new viewpoint of non-linear systems' complexity. At the same time, establishes the multifractal characteristics (central Hölder exponent, extreme Hölder exponents, spectral amplitude and spectral asymmetry), as well as a synthesis of the physical mechanism degree complexity by taking advantage of a combination of Hölder exponents.

### 1. MDFA algorithm

The multifractal detrended fluctuation analysis (MDFA) represents a reliable way of characterization of multifractal nonstationary and stationary time series (Kantelhardt *et al.*, 2002), the main objective being quantification of the rescaled structure of the  $q$ -order moments of the analyzed series, which is a process inherent to a multifractal structure analysis. It is worth mentioning that the MDFA offers higher quality and simplicity in comparison with other algorithms with similar objectives (Feder, 1988; Muzy *et al.*, 1994) and it has been applied to very different scientific fields such as human health, biology or Earth sciences, among others.

The MDFA algorithm can be summarized as the following five steps:

- Computation of the time series profile obtaining residuals from the average of the whole series.
- Segmentation of the profile.  $N_s$  non-overlapping segments of the profile, with equal length  $s$ , are generated. With the aim of improving computational accuracy, this second step is usually repeated starting from both extremes of the profile. In this way,  $2N_s$  segments are available.
- Computation of the local variance for each one of the  $2N_s$  segments. A least-square polynomial fit is computed for every segment and the residual variance of this fitting process is quantified for every segment, the non-stationarity of the series being then removed. In agreement with Koscielny-Bunde *et al.* (2006), polynomial degrees varying from 2 to 5 could be convenient.

- Computation of the  $q$ th-order fluctuation function,  $F(s)_q$ . This function is defined by Eqs. (7a) and (7b),

$$F(s)_q = \left\{ \frac{1}{2N_s} \sum_1^{2N_s} [F^2(s, \nu)]^{q/2} \right\}^{1/q}; \quad q \neq 0; -\infty < q < +\infty, \quad (7a)$$

$$F(s)_0 = \left\{ \frac{1}{4N_s} \sum_1^{2N_s} \ln[F^2(s, \nu)] \right\}; \quad q = 0. \quad (7b)$$

$F^2(s, \nu)$  is the local variance for every one of the  $2N_s$  segments and parameter  $q$  is chosen varying within the  $+15, -15$  range. Steps 2–4 have to be repeated for several segment lengths  $s$ . In agreement with Kantelhardt *et al.* (2002) useful values of  $s$  would be in the  $m + 2 \leq s < N/4$  interval, where  $m$  is the selected order for polynomial fits of step 3.

• *The scaling behavior of  $F(s)_q$ .* By assuming the hypothesis that the analyzed series are long-range power-law correlated, log–log plots of  $F(s)_q$  vs  $s$  for each value of  $q$  accomplish the power-law given by Eq. (8),

$$F(s)_q \approx s^{h(q)}, \quad (8)$$

with the exponent  $h(q)$  depending on  $q$ . If the analyzed series is non-stationary or noisy (fractal Brownian signals, for example), the exponent  $h(q = 2)$  will be equal to  $H + 1$ , with  $H$  being the above introduced Hurst exponent. Conversely, for the stationary time series, as daily extreme temperatures or monthly atmospheric circulation indices, the exponent  $h(q = 2)$  is exactly the Hurst exponent  $H$ . The exponent  $h(q)$  is many times also cited as the generalized Hurst exponent.

Differences between monofractal and multifractal behaviors have to be considered. Monofractal structures do not contribute very relevantly to obtain more information about the structure of the mechanism governing the analyzed series. Nevertheless, details offered by multifractal structures are much more complete. In the first case, variances  $F^2(s, \nu)$  are identical or very similar for all segments  $s$  and the generalized Hurst exponent is reduced to the Hurst exponent  $H$ . Conversely, for multifractal behavior, segments with large variance dominate the  $q$ -order fluctuation function for positive  $q$ . Thus, the generalized Hurst exponent,  $h(q)$ , describes the scaling behavior of the segments with large fluctuations. Alternatively, for negative values of  $q$ , small fluctuations govern the  $q$ -order fluctuation function and then  $h(q)$  describes the scaling behavior of the segments with small fluctuations (Movahed and Hermanis, 2008).

### 2. Singularity spectrum

The singularity spectrum  $f(\alpha)$  [Eq. (9a)] is closely related to the  $q$ -order fluctuation function,  $F(s)_q$ , in terms of the generalized Hurst exponent,  $h(q)$ , and the Legendre transform

$$f(\alpha) = q\{\alpha - h(q)\} + 1 \quad (9a)$$

being also related the Hölder exponent  $\alpha$  to the generalized Hurst exponent by Eq. (9b),

$$\alpha = h(q) + q \frac{dh(q)}{dq}, \quad (9b)$$

345 where  $\alpha$  is also identified as the singularity strength or the Hölder  
 346 exponent and  $f(\alpha)$  represents the fractal dimension of the differ-  
 347 ent subset of the series. The multifractal scaling exponent,  $\tau(q)$ , is  
 348 defined as

$$\tau(q) = qh(q) - 1 \tag{10}$$

349 and the Hölder exponent  $\alpha$ , in agreement with Eqs. (9b) and (10),

$$\alpha(q) = d\tau/dq. \tag{11}$$

350 The singularity spectrum provides new viewpoints of the mul-  
 351 tifractal structure of a series, given that  $f(\alpha)$  quantifies the fractal  
 352 dimensions of subsets of the series associated with the same sin-  
 353 gularity strength  $\alpha$ . Several relevant fractal parameters related to  
 354 the singularity spectrum have to be considered. One of them is  
 355 the critical (central) Hölder exponent  $\alpha_0$ , which corresponds to  
 356 the maximum of  $f(\alpha)$ . A small value of  $\alpha_0$  implies that the “fine-  
 357 structure” of the physical mechanism cannot be analyzed from  
 358 empiric data. Conversely, a large value of  $\alpha_0$  strongly suggests the  
 359 recovering of the “fine-structure.” It is worth mentioning that the  
 360 Hurst exponent,  $H = h(q = 2)$ , and  $\alpha_0$  show a clear linear relation-  
 361 ship, confirmed, for instance, by [Burgueño et al. \(2014\)](#). The other  
 362 two relevant parameters, spectral asymmetry and spectral width, are  
 363 related to the mathematical structure of  $f(\alpha)$ . The shape of this func-  
 364 tion is expected well fitted to a quadratic function [Eq. (12)] around  
 365 the position  $\alpha_0$ ,

$$f(\alpha) = A(\alpha - \alpha_0)^2 + B(\alpha - \alpha_0) + C, \tag{12}$$

366 where  $C$  is an additive constant theoretically equal to 1, given that  
 367  $f(\alpha)$  is defined as a normalized function with a maximum equal  
 368 to 1.0, and  $B$  quantifies the asymmetry of the spectrum:  $B = 0$  for  
 369 a symmetric spectrum;  $B > 0$  (“fine-structure”) for a right-skewed  
 370 spectrum;  $B < 0$  (“smooth-structure”) for a left-skewed spectrum.  
 371 The spectrum width,  $W$ , defined as the range of  $\alpha$  is defined as  
 372  $W = \alpha_{\max} - \alpha_{\min}$ , with  $f(\alpha_{\max}) = f(\alpha_{\min}) = 0$ . Bearing in mind that  
 373  $\alpha_{\max}$  and  $\alpha_{\min}$  are theoretically obtained by tending  $q$  to  $\pm\infty$ , these  
 374 extreme Hölder exponents are estimated by extrapolating the fitted  
 375 curve of Eq. (12) to zero. In short, the wider the range of the  
 376 Hölder exponent, the stronger is the multifractality. Similarly, the  
 377 wider the range of  $\alpha$ , the “richer” is the structure of the physical  
 378 process. As a summary, a series with a high value of  $\alpha_0$ , a wide  
 379 range of fractal exponents and a right-skewed shape is more com-  
 380 plex than one with the opposite characteristics ([Shimizu et al., 2002](#)).  
 381 In terms of physical mechanism, a fine-structure could be analyzed  
 382 provided that parameters  $\alpha_0$ ,  $B$ , and  $W$  confirm high complexity.  
 383 On the contrary, only a smooth-structure of the physical mechanism  
 384 would be obtained if the values of these three parameters suggest low  
 385 complexity.

386 The complexity of the monthly rainfall multifractal structure  
 387 can be summarized by the complexity index, CI, proposed by  
 388 [Shimizu et al. \(2002\)](#). This global coefficient of complexity is defined  
 389 by the addition of three normalized multifractal parameters, the cen-  
 390 tral Hölder exponent,  $\alpha_0$ , the multifractal amplitude,  $(\alpha_{\max} - \alpha_{\min})$ ,  
 391 and a new quantification of the asymmetry  $(\alpha_{\max} - \alpha_0)/(\alpha_0 - \alpha_{\min})$ .  
 392 In agreement with this definition of the asymmetry, a quotient very  
 393 close (or equal) to 1.0 will imply high (absolute) symmetry. Con-  
 394 versely, a left asymmetry will be characterized by a quotient lower  
 395 than 1.0 and a right asymmetry by a quotient higher than 1.0. The

TABLE I. Minimum, maximum, mean, standard deviation, SD, and skewness, Sk, of the RFT parameters for the 96 records.

	Min	Max	Mean	SD	Sk
$H$	0.37	0.70	0.55	0.06	0.338
$\lambda_1$	0.11	0.22	0.16	0.02	0.316
$D_{KY}$	11.99	13.62	12.66	0.38	0.286
$\mu^*$	6.71	9.75	8.18	0.60	0.205
$K$	0.54	2.31	1.64	0.31	-0.643

396 asymmetry could be also directly represented by the coefficient  $B$  of  
 397 Eq. (12). Nevertheless, this coefficient could be sometimes affected  
 398 by computational uncertainties due to a relative bad fit of empirical  
 399 data to the mentioned Eq. (12).

400 In short, whereas high positive values of CI would imply a  
 401 notable complexity on the physical mechanisms governing the ana-  
 402 lyzed phenomena, negative and lower positive CI's would be associ-  
 403 ated with more simple physical mechanisms. In consequence, high  
 404 positive values of CI would be associated with difficult success pre-  
 405 dictability. Conversely, negative or low positive values of CI would  
 406 suggest an easier predictability.

407 Finally, it is straightforward to conclude that the monofractal-  
 408 ity will be characterized by a singularity spectrum  $f(\alpha) = 1.0$  of null  
 409 width. This fact would imply that the dependence of  $h(q)$  on  $q$  disap-  
 410 pears, and it is reduced to the Hurst exponent  $H$ , which at the same  
 411 time is coincident with the single Hölder exponent  $\alpha$ . In short, it is  
 412 evident that a monofractal structure has to be assumed notably less  
 413 complex than a multifractal structure.

414 **IV. RESULTS**

415 A relatively complex spatial distribution of the fractal parame-  
 416 ters has to be expected due to, as mentioned in Sec. II, the complex  
 417 orography including several mountain chains, a central basin, and  
 418 proximity to the Mediterranean coast. As a consequence, a revision  
 419 of several characteristics concerning RFT and MDFA parameters is  
 420 necessary to detect places where the physical processes governing  
 421 the monthly rainfall would be relatively easy or complex. First of

TABLE II. Cross-correlation coefficients between different pairs of FRT parameters. Bold types correspond to coefficients very close to 1.0 and exceeding 95% statistical confidence.

$H$	$\lambda_1$	-0.10
$H$	$D_{KY}$	-0.19
$H$	$\mu^*$	0.03
$H$	$K$	-0.03
$\lambda_1$	$D_{KY}$	<b>0.86</b>
$\lambda_1$	$\mu^*$	<b>0.89</b>
$\lambda_1$	$K$	0.05
$D_{KY}$	$\mu^*$	-0.53
$D_{KY}$	$K$	0.13
$\mu^*$	$K$	0.29

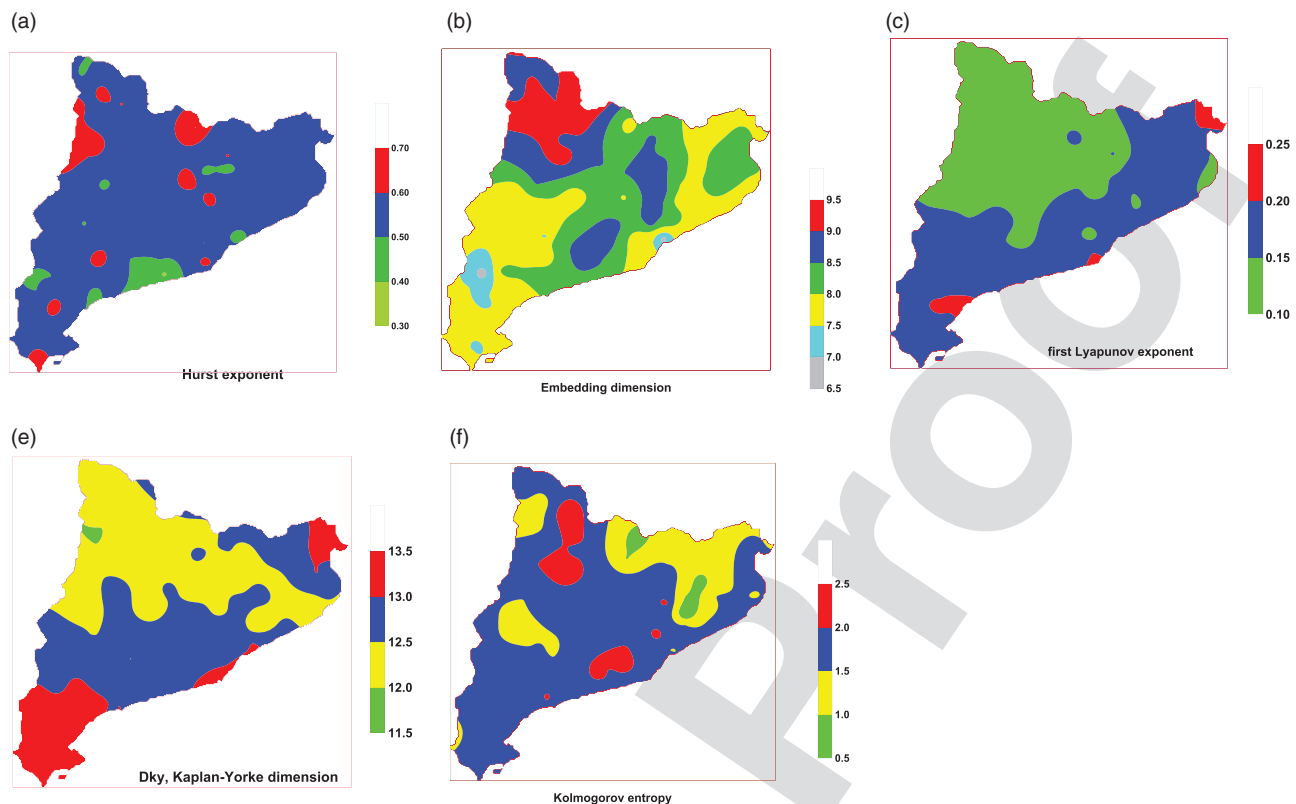


FIG. 4. Spatial distribution of (a) Hurst exponents, (b) embedding dimensions, (c) first Lyapunov exponents, (d) Kaplan–Yorke dimension, and (e) Kolmogorov entropy.

all, the predominant and extreme parameter values are introduced and their meaning, with respect to the complexity of the physical mechanism, are discussed. After that, more details are given with respect to the spatial distribution of these parameters and possible relationships to orography and proximity to the Mediterranean Sea.

#### A. Reconstruction fractal theorem

(a) *The Hurst exponent:* With respect to parameters obtained from the RFT (Table I), the Hurst exponent is characterized by predominant values slightly exceeding 0.5 (showing signs of persistence), manifested by a skewness of 0.34, average and standard deviation of 0.55 and 0.06, respectively, and only a few cases with low and high Hurst exponents, implying anti-persistence and persistence, respectively. The spatial distribution of this exponent (Fig. 4) is characterized by a domain of quasi-randomness behavior ( $H$  varying from 0.4 to 0.6) covering almost all the analyzed area, and some isolated emplacement of clear anti-persistence ( $H < 0.4$ ) and persistence ( $H > 0.6$ ). In short, the monthly rainfall regime could be assumed with a few notable signs of persistence and anti-persistence and predominant behavior close to randomness. This predominance would be an impediment on autoregressive processes leading to compute forthcoming monthly amounts.

(b) *The embedding dimension:* The minimum number of necessary non-linear equations, represented by  $\mu^*$ , to quantify the monthly rainfall series is characterized by a relatively long interval (6.71–9.75), with average and standard deviations of 8.18 and 0.60, respectively, and a low skewness (Table I). A high ratio of embedding dimensions is within the interval (7.25–9.0). Consequently, due to these high values of  $\mu^*$ , the autoregressive process to quantify forthcoming monthly amounts is expected to be complex. The spatial distribution of  $\mu^*$  (Fig. 4) detects the Central Basin and some isolated emplacement of the Mediterranean coast with relatively low values of  $\mu^*$ . It is also relevant signs of an increasing tendency toward the North-eastern and, especially, toward the North-western (Pyrenees domain). In short, the monthly rainfall regime would be characterized by more complex mechanisms (systems of non-linear equations), especially in these two just mentioned areas.

(c) *The predictive instability:* The instability, in other words, the uncertainty on forthcoming predicted monthly amounts, is notably governed by the first Lyapunov exponent. Its values are delimited by a short range from 0.11 to 0.22, with a very small standard deviation (0.02) and a notable number of samples exceeding an average of 0.16 (Table I). A high ratio of monthly rainfall amounts are characterized by  $\lambda_1$  varying from 0.125 to 0.175. In agreement with Fig. 4, besides two small areas at the



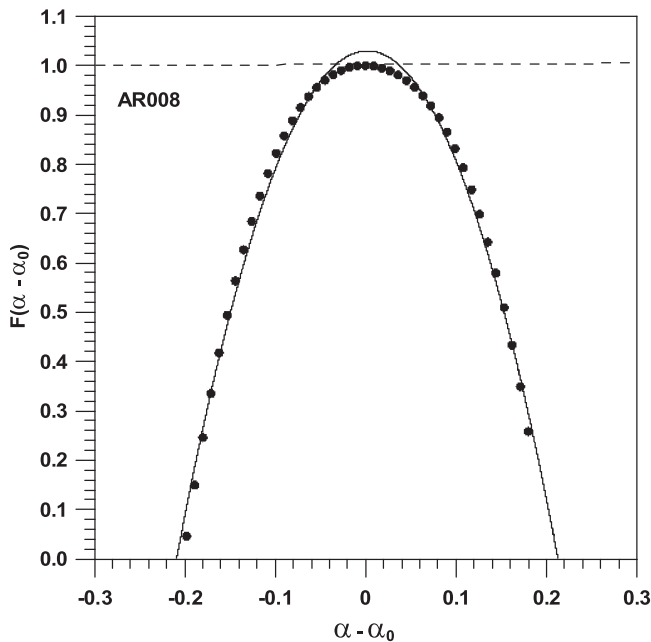


FIG. 5. An example of multifractal spectrum corresponding to a gauge in North-Western Catalonia. Solid points represent empirical samples of multifractality for a wide range of  $\alpha(q)$  parameters. The continuous line describes the theoretical multifractal spectrum, given by a second order polynomial.

- 468 North-east and the South-west with the highest values, this first  
 469 Lyapunov exponent is characterized by two well defined areas;  
 470 one of them approximately covering the Mediterranean coast,  
 471 Littoral, and Pre-littoral chains and the beginning of the East-  
 472 ern Pyrenees; the other corresponding to the rest of the analyzed  
 473 domain, including a good part of the North-western Catalonia  
 474 and most of the Eastern Pyrenees chain.  
 475 (d) *Kaplan–Yorke dimension*: The Kaplan–Yorke dimension is  
 476 characterized by a narrow range (11.99–13.62), an average of  
 477 12.66, a moderate standard deviation (0.38), and skewness (0.29)  
 478 (Table I), and a high ratio of monthly series with  $D_{KY}$  ranging  
 479 from 12.25 to 13.25. Due to its expected relation to the Lyapunov  
 480 exponents (specially the first exponent), a spatial structure, sim-  
 481 ilar to that observed for  $\lambda_1$ , should be expected, excepting for  
 482 the range of (12.0–12.5) of Kaplan–Yorke dimensions expand-  
 483 ing along a narrow fringe up to the Northern Mediterranean  
 484 coast (Fig. 4). In short, higher predictive instabilities have to be  
 485 expected toward the South and North-east of the Catalonia.  
 486 (e) *The loss of memory of the physical system*: The measure of  
 487 the loss of memory of the physical system (the Kolmogorov  
 488 exponent  $K$ ), is characterized by a wide range (0.54–2.31) and  
 489 standard deviations and average values of 0.31 and 1.64 (Table  
 490 I). Nevertheless, a high number of cases are within the (1.5–2.0)  
 491 interval. Consequently, most of the monthly rainfall patterns  
 492 would be characterized by outstanding loss of memory. Then,  
 493 the necessary number of monthly amounts to obtain a reliable  
 494 estimation of the next amount would be high, this fact making

TABLE III. Minimum, maximum, mean, standard deviation, SD, and skewness, Sk, of the MDFA parameters.

	Min	Max	Mean	SD	Sk
$\alpha_0$	0.36	0.66	0.50	0.05	0.101
$\alpha_{max}$	0.55	0.91	0.72	0.07	0.184
$\alpha_{min}$	0.02	0.55	0.27	0.09	-1.471
$W$	0.18	1.01	0.44	0.12	1.530
$\Gamma$	-4.45	2.50	-0.05	1.18	-0.404
CI	-2.83	2.52	0.00	1.00	-0.071

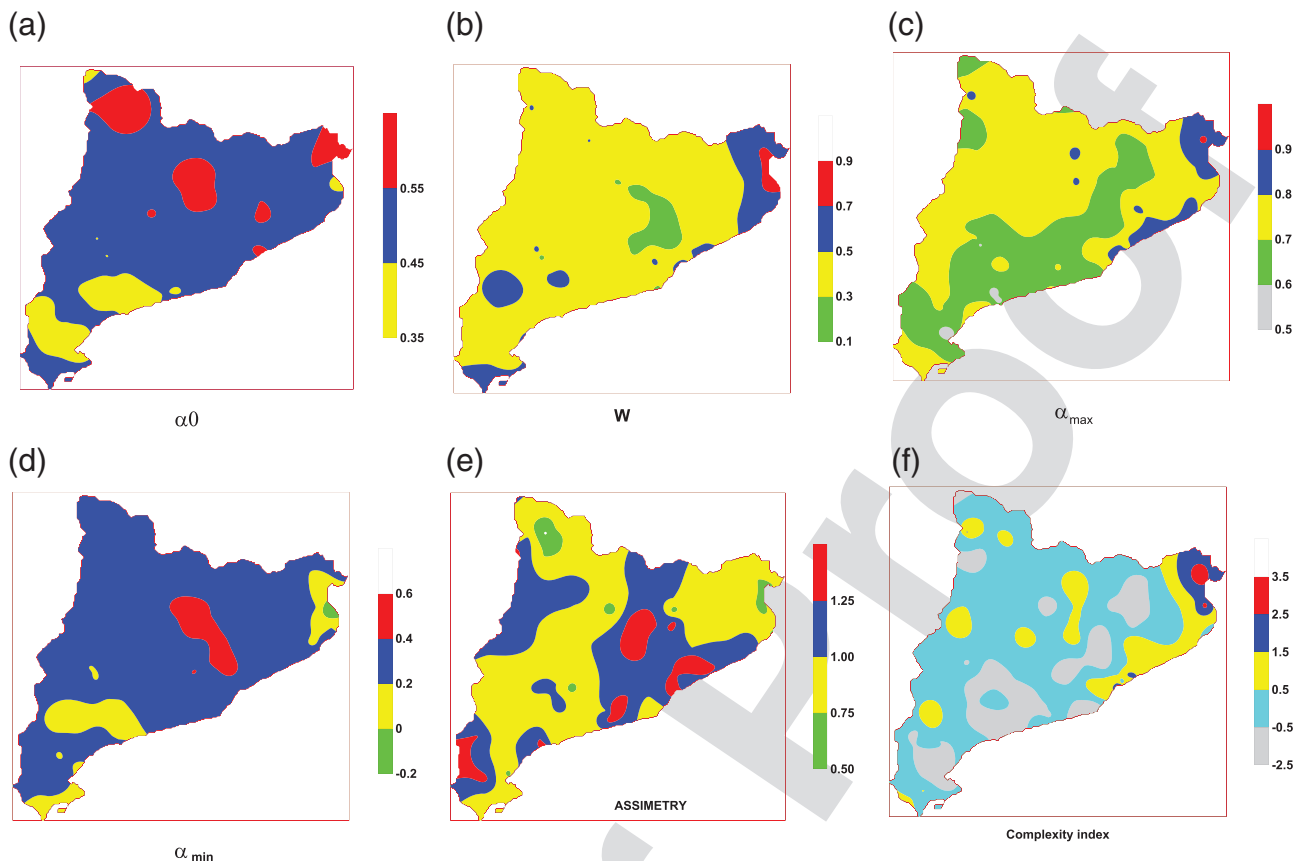
495 complex an autoregressive process. The spatial distribution of  
 496  $K$  is not very homogeneous (Fig. 4) and distinction has to be  
 497 made between a clear dominant area with  $K$  ranging from 1.5  
 498 to 2.5 and another lowering 1.5. In consequence, two regions  
 499 are detected. On one hand, an area representing a high ratio of  
 500 the domain where the loss of memory is relevant; on the other  
 501 hand, a more reduced area (North-eastern and some Western  
 502 emplacements), where the loss of memory of the physical system  
 503 is notably smaller.

504 Table II also summarizes possible relationships among the dif-  
 505 ferent RFT parameters. The notable cross-correlation between the  
 506 first Lyapunov exponent and the Kaplan–Yorke dimension  
 507 could be expected bearing in mind Eq. (6), where the dependence  
 508 of  $D_{KY}$  on the Lyapunov exponents is quite evident. Beside this  
 509 first high cross-correlation, it is also very relevant the correlation  
 510 between the first Lyapunov exponent,  $\lambda_1$ , and the dimension  $\mu^*$ .  
 511 In agreement with the meaning of these two fractal parameters,  
 512 the predictive instability would increase with the number of non-  
 513 linear equations describing the physical mechanism of the successive  
 514 monthly rainfall amounts.

### B. Multifractal detrended fluctuation analysis

515  
 516 An example of multifractal spectrum is shown in Fig. 5. This  
 517 spectrum has been obtained from monthly rainfall records belong-  
 518 ing to a rain gauge emplaced in the Pyrenees. Solid points are the  
 519 empiric samples of multifractal spectrum for a wide range of  $q$   
 520 exponents, within the  $\pm\infty$  interval. The solid line is the quadratic  
 521 polynomial given by Eq. (12), and the dashed horizontal line delimit-  
 522 its the maximum expected multifractal amplitude  $f(\alpha - \alpha_0) = 1.0$ . In  
 523 spite of the good fit of empiric multifractality to Eq. (12), the the-  
 524 oretical maximum is slightly higher than 1.0, which is a, sometimes  
 525 detected, not very relevant shortcoming. The example of this figure  
 526 is also a good example of almost null asymmetry ( $\gamma \approx 0$ ), given that  
 527 the extrapolated values to determine the extremes ( $\alpha_{max} - \alpha_0$ ) and  
 528 ( $\alpha_{min} - \alpha_0$ ) are very close to  $\pm 0.21$ , leading to a spectral amplitude  
 529 close to 0.42. In agreement with Table III, these last two values are  
 530 very close to the average of  $W$  and  $\gamma$  when the whole sample of 96  
 531 parameters is analyzed.

- 532 (a) *The central  $\alpha_0$  Hölder exponent*: This multifractal parameter,  
 533 characterizing the complexity of the physical process, is iden-  
 534 tified by extreme values of 0.36 (low complexity) and 0.66 (high  
 535 complexity). Bearing in mind that its average and standard



**FIG. 6.** Spatial distribution of (a) the central Hölder exponent, (b) spectral amplitude, (c) maximum and (d) minimum Hölder exponents, (e) spectral asymmetry, and (f) complexity index.

deviation are 0.50 and 0.05, respectively, and the low skewness (0.101) (Table III), most of the  $\alpha_0$  values are detected within the interval (0.35–0.45) which would represent physical processes of moderate complexity. Nevertheless, bearing in mind the contribution of the other two parameters, multifractal spectral amplitude  $W$  and asymmetry  $\gamma$ , the degree of physical complexity will be finally quantified by means of the index CI. With respect to the spatial distribution of the central Hölder exponent (Fig. 6), it is evident that a great part of Catalonia is characterized by  $\alpha_0$  within the 0.45–0.55 interval. Only toward the south, for some areas,  $\alpha_0$  is less than 0.45 (minor complexity). Opposite to this, for some places on the North, not necessarily emplaced in the Pyrenees, the parameter exceeds 0.55, describing higher complexity. At least for this multifractal parameter, the proximity to the Mediterranean coast and the Pyrenees, or the emplacement on the Central Basin, are factors not conditioning  $\alpha_0$ .

(b) *The maximum,  $\alpha_{max}$ , and minimum,  $\alpha_{min}$ , Hölder exponents:* In agreement with Table III,  $\alpha_{max}$  and  $\alpha_{min}$  are characterized by small standard deviations (0.09 and 0.05, respectively) and a moderate skewness on  $\alpha_{max}$ . Most of the values of these

parameters are detected within narrow ranges (0.65–0.80 for  $\alpha_{max}$  and 0.20–0.35 for  $\alpha_{min}$ ). From the spatial point of view (Fig. 6),  $\alpha_{max}$  with values from 0.7 to 0.8 covers a great area, being also notable the domain characterized by values within the 0.6–0.7 interval. Are also outstanding some small isolated domains along the Mediterranean coast where  $\alpha_{max}$  exceeds 0.8.  $\alpha_{min}$  varying from 0.2 to 0.4 covers a great portion of the map, being also worthy of mention two small nuclei (0.0–0.2 and 0.4–0.6). Although both extreme Hölder parameters contribute to the quantification of the complexity, emplacements with similar behaviors on  $\alpha_{max}$  and  $\alpha_{min}$  are difficult to detect. For this reason, for a better characterization of the complexity, the spectral amplitude  $W$  can be considered, ranging from 0.18 to 1.01, an average and standard deviation of 0.44 and 0.12, a very relevant skewness of 1.53 (Table III), and with many of the  $W$  samples within the (0.30–0.55) interval. Bearing in mind the notable value of the skewness, a quite heterogeneous spatial distribution of this last index could be expected. Nevertheless, Fig. 6 describes a great portion of Catalonia associated with spectral amplitudes from 0.3 to 0.5, being also detected the widest amplitudes on two reduced domains of the Mediterranean coast

**TABLE IV.** Cross-correlation, CCor, for pairs of MDFA parameters. Bold types correspond to coefficients exceeding 95% statistical confidence level.

Pairs of MDFA	CCor
$\alpha_0, \alpha_{\max}$	<b>0.60</b>
$\alpha_0, W$	-0.12
$\alpha_0, \gamma$	-0.04
$\alpha_0, CI$	0.52
$\alpha_{\max}, W$	<b>0.64</b>
$\alpha_{\max}, \gamma$	0.07
$\alpha_{\max}, CI$	<b>0.89</b>
$W, \gamma$	-0.24
$W, CI$	0.48
$\gamma, CI$	0.42

(North-east and South-west extremes) and the narrowest ones on a small area in the Pre-Littoral chain.

(c) *The asymmetry*: Another point of view of the complexity is offered by the asymmetry  $\gamma$ , with extreme values varying from -4.45 (strong left asymmetry) to 2.50 (strong right asymmetry). The average of all the 96 asymmetry samples is almost null (-0.05) and the standard deviation is 1.18 (Table III). In agreement with Fig. 6, the most common range of asymmetry is from 0.75 to 1.0 and from 1.0 to 1.25. Both positive ranges would be associated with right asymmetries of the multifractal spectra, characterizing signs of complexity. Exceptionally, the highest right asymmetries ( $\gamma$  exceeding 1.25) are detected in reduced areas on the Central Basin, the Mediterranean coast, and the inner Catalonia. Conversely, the left asymmetries ( $\gamma$  lowering 0.75 appear in reduced areas and, especially, on the north face of the Eastern Pyrenes).

(d) *The complexity index*: Taking into account the possibility of monthly rainfall series associated with very different values of  $\alpha_0$ ,  $W$ , and  $\gamma$ , the index CI is finally chosen as the best parameter to define the complexity of every monthly rainfall series. This index is characterized (Table III) by extremes of -2.83 (low complexity) and 2.52 (high complexity), with moderate symmetry (skewness equal to 0.071) and a clear predominance of monthly series, 71 out of 96 cases, within the (-1.0, +1.0) interval. The spatial distribution of CI can be described (Fig. 6) by a few places with low values ( $< -0.5$ ), being an example the mouth of the Ebro river and neighboring domains, at the southern extreme, prevalence of a very moderate complexity (-0.5, 0.5) throughout the country and outstanding values of CI (strong complexity) toward the North-eastern extreme.

Possible correlations among the different multifractal parameters are summarized in Table IV. The cross-correlation coefficients certainty has been quantified, in agreement with Hirsch *et al.* (1992), being obtained in all cases cross-correlations exceeding 90%–95% significant levels. It is noticeable the strong positive correlation between  $\alpha_{\max}$  and CI as well as the notable correlations between  $\alpha_{\max}$  and  $W$  and  $\alpha_{\max}$  and  $\alpha_0$ . It is also worth mentioning that the weight of the central Hölder exponent  $\alpha_0$ , the multifractal spectral amplitude  $W$ , and the asymmetry  $\gamma$  on the parameter CI is

quite similar by revising the corresponding cross-correlation coefficients. Consequently, besides the strong correlation between  $\alpha_{\max}$  and CI, the complexity of the physical mechanism governing the monthly rainfall regime is due to a similar weighted contribution of the central Hölder exponent, the multifractal spectral amplitude, and asymmetry.

## V. CONCLUSIONS

The analysis of the fractal/multifractal structure of the monthly rainfall regime in Catalonia permits to confirm from a new viewpoint the heterogeneous spatial distribution of monthly rainfall patterns. This heterogeneous distribution, in spite of a relatively reduced geographic domain, would be the consequence of a complex orography, with a notable range of altitudes from zero meters above sea level (Mediterranean coast) up to 3000 m (Pyrenees) and the different mountain chains (Eastern Pyrenees, Littoral, and Pre-Littoral chains) and a wide Central Basin. Bearing in mind the predominant NW frontal passages from the Atlantic Ocean and the Eastern advections on the Mediterranean Sea, these one especially in autumn, the effects of the topographic barriers created by the mountain chains on local rainfall regimes are expected. Additionally, the proximity or remoteness to the Mediterranean coast is another factor to be considered.

Concepts as randomness/persistence, uncertainties on forecasting processes, the complexity degree of the non-linear equations describing the physical process, its loss of memory, and a combined set of multifractal parameters, leading to quantifying the complexity of the monthly rainfall regime, have permitted to accomplish two relevant objectives. First, a new and complete classification of rainfall pattern areas, bearing in mind the spatial distribution of the different fractal/multifractal parameters describing properties of the rainfall regime. Second, the detection of rainfall regimes needing a high, low, or moderate number of consecutive monthly amounts previous to the forthcoming amount obtained by autoregressive processes. Additionally, places where the uncertainty and the discrepancy between forecasted and really recorded amounts would be high, low, or moderate are also detected. These results, concerning number of necessary monthly amounts and discrepancies between real and forecasted amounts, are relevant bearing in mind that the analyzed Mediterranean region is characterized by long dry spells. Then, accurate autoregressive processes are necessary to forecast forthcoming monthly amounts preventing or mitigating the effects of these long dry spells. Consequently, the next step on the increasing knowledge about the monthly rainfall regime in Catalonia should be the application of autoregressive processes, bearing in mind results obtained in this paper from the viewpoint of the fractal/multifractal theory. Additionally, validation of possible time trends on monthly amounts should be also considered.

## DATA AVAILABILITY

The monthly scale data that support the findings of this research are available from the corresponding author upon reasonable requests. Daily scale data are also available from the meteorological agency of Catalonia (SMC, [www.meteo.cat](http://www.meteo.cat)), Ref. ■; Fabra Observatory (*Reial Acadèmia de Ciències i Arts, RACA*, Barcelona);

670 and Ebre Observatory (Ramon Llull University, ■, and Consejo  
671 Superior de Investigaciones Científicas, Spanish Government).

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