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Behaviour of the Mg II 2798Å Line Over the Full Range of AGN Variability

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ABSTRACT

We investigate the responsiveness of the 2798Å Mg II broad emission line in AGN on timescales of several years. Our study is based on a sample of extremely variable AGN as well as a broad population sample. The observed response of the line in previous studies has been mixed. By focussing on extreme variability ($|\Delta g| > 1$) we find that Mg II clearly does respond to the continuum. However, the degree of responsiveness varies strikingly from one object to another: we see cases of Mg II changing by as much as the continuum, more than the continuum, or very little at all. In 74% of the highly variable sample the behaviour of Mg II corresponds with that of H β , with 30% of the objects showing large variations in both lines. We do not detect any change in the line width that would correspond to Broad Line Region ‘breathing’, in accordance with results from the literature. Some of the objects in our highly variable sample show a clear asymmetry in the Mg II profile. This skewness can be both to the blue and the red of the line centre. Results from our broad population sample show that highly variable quasars have lower Eddington ratios. This result holds for the variability of the continuum, but the correlation is significantly reduced for the variability of the Mg II line. For the first time, we present an overview of the value of the intrinsic Baldwin Effect for Mg II in a large sample.

Key words: galaxies:active – quasars:emission lines – quasars:general

1 INTRODUCTION

Variability is inherent to AGN emission and affects all parts of the electromagnetic spectrum. The timescales and amplitude of fluctuations differ across the spectral range. Changes in the UV-optical continuum output are accompanied by changes in broad emission lines (BELs), often at a smaller amplitude than the continuum (Peterson 2008; Bentz et al. 2013; Sun et al. 2015). Average optical continuum variability on a timescale of months-years is ~10-20% (e.g. MacLeod et al. 2010). However a subset of Extremely Variable Quasars (EVQs) can, on the same timescales, show variability in the order of several (Lawrence et al. 2016; Rumbaugh et al. 2018; Graham et al. 2020). Extreme continuum variability is closely linked to the concept of changing-look (CL) AGN (Tohline & Osterbrock 1976; Shappee et al. 2014; LaMassa et al. 2015; MacLeod et al. 2016), in which the change in broad H β is such that the AGN changes between a type 1

and a type 2 classification (Khachikian & Weedman 1971; Osterbrock 1981). The extreme variability of continuum and BELs in these objects provides the opportunity for greater insight into AGN evolution and structure.

The focus of this work is on the response of the broad Mg II λ 2798 emission line. This line differs from most other emission lines in its formation mechanism, because it is collisional (MacAlpine 1972; Netzer 1980; Guo et al. 2019)¹. Results from previous studies with regard to the response of Mg II to the continuum have been mixed. Some studies have found a response (Clavel et al. 1991; Reichert et al. 1994; Metzroth et al. 2006; Hryniewicz et al. 2014), although the variable fraction of the flux is usually only a few percent. An exception to this is the quasar LBQS 2113-4538, studied by Hryniewicz et al. (2014), which shows a fluctuation in Mg II Equivalent Width (EW) of ~25% over several hundred days.

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¹ Recombination of Mg III is expected to occur at a negligible rate (e.g. Guo et al. 2019)

A similar number of studies finds no correlation between continuum changes and the variability of Mg II (Trevese et al. 2007; Cackett et al. 2015).

This complex behaviour has also been studied using larger samples consisting of survey data. Woo (2008) finds peak-to-peak changes of up to 25% on a timescale of approximately a year, indicating the potential for large flux changes. In a study of $\sim 9,000$ Stripe 82 quasars Kokubo et al. (2014) find that both Mg II and the Fe II complex in the same wavelength range vary only slightly, when compared to the Balmer lines. If Mg II does respond to the continuum, this line could be a promising candidate for Reverberation Mapping (RM) studies at high z . Early results of the SDSS-IV RM programme showed a credible Mg II lag in six quasars ($0.3 < z < 0.8$) out of a sample of ~ 100 objects (Shen et al. 2016a). A different study, also based on the SDSS-RM data, finds that Mg II is only slightly less variable than H β (Sun et al. 2015).

In comparisons from object to object, the line flux of Mg II is found to correlate with the optical continuum at 5100 Å (Shen & Liu 2012) and the width of Mg II is found to correlate strongly with the widths of H α and H β (Shen & Liu 2012; Wang et al. 2019). These correlations suggest that Mg II and the Balmer lines are formed in the same, or a similar, region in the BLR (McLure & Dunlop 2004). This would broadly agree with photoionisation models (Korista & Goad 2004; Baskin et al. 2014). Dong et al. (2009) find that the Equivalent Width (EW) of Mg II is strongly anti-correlated with the Eddington ratio, but not with the continuum luminosity at 3000Å. These authors argue that the lack of correlation with the luminosity underlies the Mg II ensemble Baldwin Effect.

An investigation into epoch to epoch Mg II variability in a sample of SDSS quasars is presented in Zhu et al. (2017), who find that the Mg II line tracks the continuum and that the responsivity of the line decreases with a higher initial continuum state and with a higher Eddington ratio. Both conclusions are supported by the results of Sun et al. (2015). Yang et al. (2020) report a weak correlation between the change in the 3000Å continuum and the change in Mg II luminosity, based on spectroscopic follow up of 16 objects. The variability of Mg II is of smaller amplitude than that of the UV continuum, but does appear to track it. A suggested explanation is that only part of the line responds to the ionising flux (cf. Korista & Goad 2004).

In contrast to the flux changes, the Full Width at Half Maximum (FWHM) of Mg II does not change with the fluctuating continuum. The lack of change in line width suggests that the radius to luminosity relationship (on which single epoch BH mass estimates are based) does not hold for Mg II. This marks a particular difference with broad H β . Interestingly, Roig et al. (2014) found a class of quasars characterised by a strongly suppressed broad H α and H β emission and a very prominent Mg II line, without signs of reddening. This population fits well in the evolutionary sequence presented in Guo et al. (2020): Based on a photoionisation model Guo et al. (2020) find that Mg II responds consistently to continuum changes, but at a much diminished level compared to the Balmer lines.

In this paper we aim to show that the Mg II line does respond to continuum variability. The response is complex and we detect a range of behaviours of the Mg II line, both

in flux and in line profile. We consider the variability of the broad Mg II line in two samples. The first sample consists of 43 highly variable quasars. 40 of these are part of the CLQ candidate observations described in MacLeod et al. (2019, hereafter referred to as MCL19). This paper aims to present a complementary analysis of this data-set. The remaining three objects are part of the SDSS CLQ search presented in MacLeod et al. (2016). The second sample is a broad population sample, based on the SDSS quasar catalogue.

The structure of this paper is as follows. In Section 2 we describe our observations and the selection of our data-sets. Mg II line characteristics are calculated from new fits to the spectra in both samples. Section 3 presents the fitting method used, as well as several methods we will employ to quantify the changes in the spectra. In Section 4 we discuss the connection between Mg II and H β variability. The changes in the flux of the Mg II line and its relation to the continuum flux are discussed in Section 5. Following the discussion of the flux, we will consider the line profile and the possibility of extracting information about the kinematics of the line forming regions in Section 6. The connection to physical parameters is discussed in Section 7. We discuss the implications of our findings in Section 8 and summarise our conclusions in Section 9. In Appendix A we provide additional details on our statistical methods. Where necessary we assume a standard flat cosmology with $\Omega_\Lambda = 0.7$, $\Omega_m = 0.3$, and $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$.

2 MG II SAMPLES

This study is based on two samples: a set of 43 highly variable objects, 40 of which have new observations (the **Supervariable sample**), and the set of SDSS DR14 quasars for which we have repeat spectra available that cover Mg II (the **Full Population sample**). The bases of the samples are the quasar catalogues for DR7 (Schneider et al. 2010) and DR14 (Pâris et al. 2014) respectively, referred to here as DR7Q and DR14Q.

2.1 Observations

All observations for SDSS were made using the 2.5m Sloan telescope at Apache Point (Gunn et al. 2006), as part of the SDSS-I/II, SDSS-III/BOSS, and SDSS-IV/eBOSS observation campaigns. The spectroscopic reduction pipelines for SDSS I/II and BOSS respectively are presented in Stoughton et al. (2002) and Bolton et al. (2012). Of particular importance to this study is that the DR14 reduction pipeline for quasar spectra includes a correction for an error in the flux calibration due to atmospheric differential refraction which is present in SDSS III/BOSS spectra (Margala et al. 2016; Harris et al. 2016).

The new observations were all part of the study discussed in MCL19. Here we provide a brief overview, for a full description please see MCL19. The observations were made using the Intermediate Spectrograph and Imaging System (ISIS) on the the 4.2m WHT in La Palma, with the Blue Channel Spectrograph on the MMT on the 6.5m telescope (Mt. Hopkins, Arizona) and the Low Dispersion Survey Spectrograph 3 (LDSS3)-C on the 6.5m Magellan telescope.

Table 1. Overview of the selection of CLQ candidates for observations, as presented in MCL19. This table is based on Table 1 in MCL19. The Supervariable sample consist of these 40 objects plus 3 known CLQs from previous studies.

Criterion	N_{obj}
DR7 catalogue	105,783
No BOSS spectrum	79,838
max $ \Delta g > 1$, $ \Delta r > 0.5$ & $z < 0.83$	1,727
No radio detection	1,403
Recent $ \Delta g > 1$	262
Observed	130
Mg II in the spectrum	40

All observations were reduced using standard long slit spectroscopy methods. The flux calibration of the MMT and Magellan spectra has been scaled to that of SDSS by matching the [OIII]4959 and [OIII]5007 fluxes. For the WHT spectra we deemed the flux calibration of our spectra to be of sufficient quality not to require this additional grey scaling.

2.2 Sample Selection

2.2.1 The Supervariable Sample

The Supervariable sample is a combination of two data-sets. The first is the CLQ sample presented in MacLeod et al. (2016) (MCL16) and the second is the CLQ candidate sample presented in MCL19. The CLQs discussed in MCL16 were found by using Pan-STARRS (PS1) and SDSS light curve data and subsequently comparing BOSS and SDSS-I/II spectra. From both data-sets the objects that have Mg II in their spectra are included. From the MCL16 sample 3 objects qualified.

The selection criteria, discussed in detail in MCL19, are listed below:

- (i) Quasars included in the DR7 quasar catalogue (Schneider et al. 2010)
- (ii) And with no BOSS spectrum following the initial SDSS observation
- (iii) And a large magnitude change in PS1 and SDSS photometry: max $|\Delta g| > 1$ and $|\Delta r| > 0.5$
- (iv) $z < 0.83$, so that H β is in the spectral range (for WHT, MMT and Magellan)
- (v) There is no associated radio source (this excludes e.g. blazars)
- (vi) The most recent photometry (2013) indicates a change in g of at least 1 mag compared to the SDSS epoch

Table 1 summarises these selections. A total of 130 candidates were observed and 40 of these objects had Mg II in the spectrum, qualifying for the Supervariable sample. In combination with the three MCL16 objects, this gives a total of 43 objects. It should be emphasised the Supervariable sample consists of both confirmed CLQs and observed CLQ candidates. All objects are characterised by strong variation in the optical flux and can be considered EVQs. The sample is listed in Table 4.

Table 2. The selection criteria for the Full Population sample, where the number in bold is the total number of quasars included in this sample. The table also shows the number of objects in a subsample of objects with high quality spectra available.

Criterion	N_{obj}
DR14 catalogue	523,356
$N_{spec} \geq 2$	72,747
$\Delta t > 30$ days	61,147
not BALQ	56,724
z-warning/exposure/coordinates	55,028
Mg II in range	15,824
Successful pipeline fit	15,101

2.2.2 The Full Population Sample

For each object in DR14Q with more than one spectrum available the pair with the *largest difference* between observation dates was selected, with a required minimum of 30 days. For the included objects, only the pair of spectra with the largest time difference is used from here on. The motivation for selecting the objects with the largest temporal baseline was to select for fitting the objects most likely to show large flux changes. This is based on the assumption that large continuum changes are more likely the longer we observe a quasar.

We exclude BALQs from the sample as these objects could represent an irregular subset of quasars, with strong outflow signatures in their spectra (e.g. Hamann et al. 2019). Further checks remove observations with an SDSS `z_warning` (indicative of a redshift problem, Boller et al. (2007)) and observations for which the exposure time is not registered in the file. The Mg II λ 2798 line is required to be in the spectral range for both the SDSS I/II and BOSS spectrographs. Assuming that for a proper fit of the Mg II line it is necessary to have flux data at ± 100 Å about the Mg II line centre, and using the smaller λ -range available to the SDSS spectrograph (3800–9200 Å), we are able to include quasars with redshift $0.41 < z < 2.17$. The Full Population sample consists of 15,824 objects; see Table 2.

The distributions of the rest-frame timespans between epochs are displayed in Figure 1. For QSOs in the Supervariable sample with more than two spectra available, Δt is calculated between spectra 1–2, 2–3, up to (N–1)–N, in chronological order, giving a total of 108 spectral pairs for the 43 objects. The distributions clearly indicate the wide range in intervals covered by our study, extending into a decade-long baseline for assessing spectral evolution.

3 TRACKING THE CHANGES IN MG II

The spectra of the Supervariable and the Full Population samples are fitted using the `lmfit` package. The pipeline written for this purpose uses an iterative process to improve the fit quality. In brief overview the stages of the pipeline are:

- (i) Load spectrum and create mask for low quality flux bins;
- (ii) Correct for Galactic dust extinction;
- (iii) Trim spectrum to fitting range;
- (iv) Fit the relevant spectral component;
- (v) Use fitting results for the next iteration of fitting;

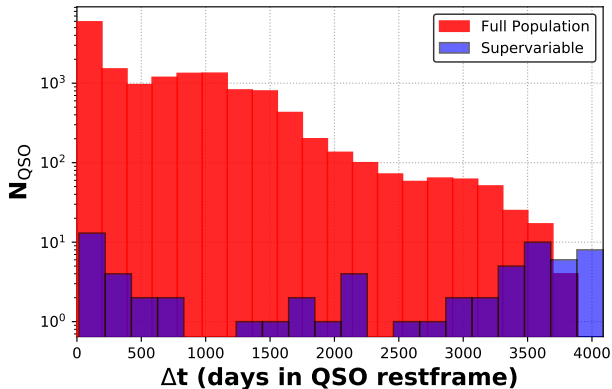


Figure 1. Histogram of the elapsed time between spectral epochs for the Supervariable sample (*blue*) and the Full Population sample (*red*). Both samples cover a large range in Δt and could therefore contain changes in Mg II responsivity as a result of two different underlying effects: structural changes in the BLR, possible only for $\Delta t \sim$ years, and large fluctuations of the continuum. Δt is defined as the time between two sequential spectra.

The steps (iv)-(v) are iterated for each spectral component, to improve the quality of the fit. In step (v) we use the fitting results as the new initial values and for sigma clipping.

The components of the model used to fit the Mg II line are the continuum power law, the Fe II template, and either one or two Gaussian profiles for the Mg II line itself. The number of iterations and the sigma-clipping limit are free parameters in the procedure, which were calibrated using the Supervariable sample. The power law continuum is defined as $A\nu^{-\alpha}$. The wavelength range for fitting this component is selected to contain a good representation the near-UV continuum, ranging from 2300Å to 3088Å. This range is also used to jointly fit the power law and the Fe II template. The Fe II template used in this step is that presented by Vestergaard & Wilkes (2001). The best-fit continuum+Fe II model is subtracted from the spectra.

The final component of the fit is the Mg II line itself, which is calculated in the window $2700\text{Å} < \lambda < 2900\text{Å}$. The line is represented with one or two Gaussian profiles. We limited the number of possible Gaussians to two, as the primary interest of this study is in flux changes: a large number of Gaussians increases the chance of spurious flux measurements, in particular due to remnant Fe II contamination in the line wings. The second Gaussian, although narrower than the other component, does not constitute a ‘narrow’ line in the same sense as e.g. the [OIII] narrow lines. The Akaike information criterion is calculated for a single-Gaussian and a two-Gaussian fit, the line model with the lowest AIC is selected. A table with an overview of all fitting parameters is included in the Online Supplement and the pipeline can be found here². An example fit to a spectrum is shown in Figure 2.

For each epoch the spectral fitting provides: the continuum flux level at the line centre (f_{2798}), a Mg II line flux (f_{MgII}), the Mg II line width (σ_{MgII}), and the central wave-

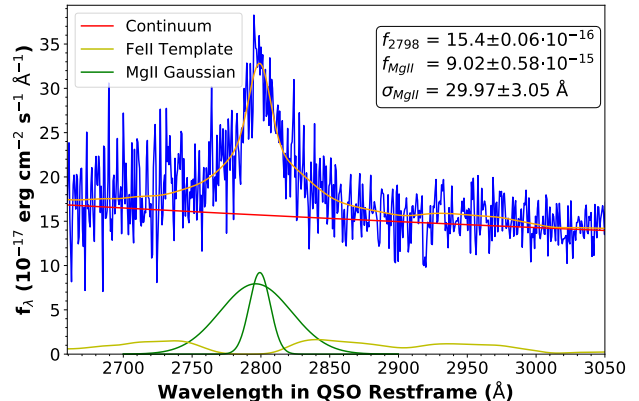


Figure 2. Example of the result of a fit by the pipeline, for the SDSS spectrum of J212436 taken on MJD 52200. The spectrum is shown in blue and the fit components are: *red*: power law continuum; *green*: Gaussian fit to Mg II (this can be one or two Gaussians, both of which constitute broad emission lines); *yellow*: Fe II template by Vestergaard et al., smoothed with a Gaussian filter; *orange*: the combined fit to the spectrum. Note in particular the presence of Fe II emission in both the red and the blue wing of the Mg II line. The most important fitted parameters are listed in the top right corner.

length of the Mg II line. In the case where two Gaussians are fit, their combined flux constitutes the broad Mg II emission. From the Mg II line centre we derive the QSO redshift. This value of z , rather than the SDSS value, will be used for all results presented here. In the case of a two-Gaussian line model, σ_{MgII} is derived from the broadest Gaussian, as this most accurately represents the full width of the line (see e.g. Figure 2). In cases of low S/N, not all spectra can be fit successfully by the pipeline. Objects with one or both spectra for which no fit could be found are excluded. This issue only applies to a small fraction of the Full Population sample: 1446 spectra ($\sim 5\%$). In the results presented further on, the number of objects in the Full Population sample is 15,101 (Table 2).

The measure of emission strength for the main results is a flux rather than the EW. This is similar to the method used in e.g. Zhu et al. (2017) and Yang et al. (2020). Although a study of EW provides important information about the behaviour of broad emission lines (Shen et al. 2011; Dong et al. 2009; Korista & Goad 2004), considering the line and continuum flux as separate parameters can allow for more insight into the behaviour of each component. For the purpose of comparison to other studies, the EW is used in Section 5.3 to calculate the intrinsic Baldwin Effect.

We use several measures to quantify Mg II variability. Different methods of normalisation highlight different characteristics of the data. The first category of normalisations pertains to the fluxes only, and is with respect to one spectral epoch. We consider both a normalisation to the high and the low state. For normalisation to the high state, the spectral epoch with the highest *continuum* flux is identified for each object. The values f_{MgII} and f_{2798} in all other epochs are then divided by f_{MgII} and f_{2798} of the high state spectrum. The emphasis in this normalisation is on small flux changes, and the associated changes in line flux. The case

² <https://github.com/dshoman/MgII>

of normalisation to the low state, using the epoch of lowest f_{2798} , allows for a clearer view of whether the line flux is able to track an increase in the continuum. The second category is a normalisation of the changes in f_{2798} , f_{MgII} , and σ_{MgII} . The quantities Δf_{2798} and Δf_{MgII} are defined as the change in flux, from epoch to epoch, divided by the flux of the second epoch.

To quantify the responsivity in the normalised flux samples, as well as any differences in behaviour, we fit two empirical functions to the normalised data. The functions are fit to the epoch-normalised flux data. The first function is a simple linear response of the line to changes in the continuum, and the second function has two linear components. The fitting is iterative, which improves the assessment of the uncertainties associated with the fit. The measurement errors on the fluxes (propagated from the errors on the fits) appear to underestimate the uncertainties in the data. The latter are indicated by the scatter. Our uncertainties are therefore estimated using an initial fit to the data, after which the standard deviation of the residuals is set to be the uniform error. The next iteration of the fit provides the parameters presented in this study.

We use a number of methods to test the line responsivity (see Appendix A for more detail). In addition to well known statistical tests, we define a *Sequence Test*. This is a test to discern which of the two fitting functions is most appropriate to describe the normalised data sets. The method measures the distribution of data points above or below the fitted functions and compares it with the binomial distribution. We also define a *Responsivity Measure*. This metric is based on a simple criterion: the perpendicular distance in the normalised flux plane from the line that would indicate a 1:1 correspondence between line and continuum flux. The responsivity measure will be referred to as α_{rm} .

4 MG II AND CHANGING-LOOK AGN

There is a wide range of variability among the objects in the Supervariable sample, see e.g. Table 4 column 7 and the spectra shown in Figure 3. The four objects included in this figure were selected as examples of four categories of Mg II variability: the top two panels show objects which both have a dimming continuum: the Mg II line changes only in one of the two cases. The bottom two panels show a similar difference in response, but for a rising continuum. Of particular interest is J002311 (panel c), for which the range of observations covers the object brightening and re-dimming. The response from Mg II as the continuum rises (between MJD 51900 and 55480) is very limited, but once the continuum diminishes again (MJD 58037) the reaction is more pronounced. The responsivity of the line therefore not only differs from object to object, *but also from epoch to epoch*.

The bottom three objects in Figure 3 can be classified as CLQs, as only the narrow line component of H β remains in the low state spectra. The classification is summarised in Table 3. The categories in Table 3 are based on two classifications: variability in H β is based on whether the object was classified as a CLQ in MCL19 and the classification of Mg II variability requires $|\Delta f_{\text{MgII}}| \geq 0.5$. The category ‘no Mg II’ covers a stable line flux, as the continuum changes. Although the relation between variability in Mg II and in

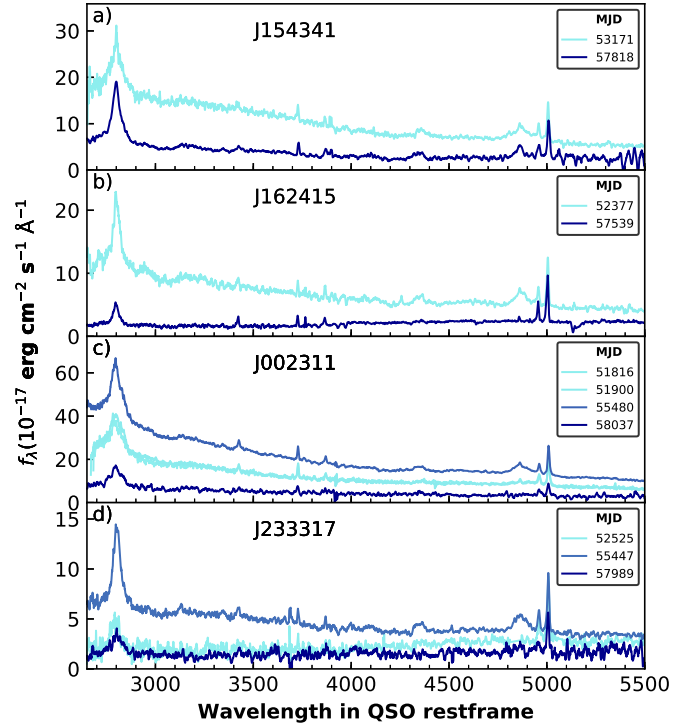


Figure 3. Examples of the wide range of behaviours detected in the Supervariable sample. *a)*: J154341 displays a diminishing continuum, with the Mg II line flux remaining stable. *b)*: J162415, the Mg II line flux tracks the lowering continuum quite closely. Note that this is also a CLQ. *c)*: J002311, where we see a rising continuum, without a strong response in the Mg II flux. We have enough spectra to see the spectrum both rise and fall: after the lack of response to a rising continuum, the Mg II tracks the continuum more closely after it dims. *d)*: J233317, which shows the Mg II flux responding to a rising continuum, and returning to a lower state subsequently.

Table 3. Four categories of variability displayed in the Supervariable sample based on H β and Mg II line flux changes. The H β variability is based on a visual identification as a CLQ in MCL19 and the Mg II variability requires $|\Delta f_{\text{MgII}}| \geq 0.5$. The second column lists the number of objects that fall in each category and the third column the average of time spans of the spectral epochs for the included objects. The value in parentheses is the standard deviation of Δt .

Type of Variability	N_{obj}	avg. Δt_{spec} in days
H β , Mg II	13	3,335 (772)
H β , no Mg II	3	3,871 (198)
no H β , Mg II	8	3,375 (527)
no H β , no Mg II	19	3,379 (769)

H β is not one-to-one, there does appear to be a connection: objects with changing H β are more likely to show variability in Mg II than not.

The categorisation in Table 3 indicates that there is no direct correlation between the elapsed time between spectral epochs and the type of variability detected. There is a large scatter in the range of Δt_{spec} (see also Figure 1) and this scatter is present in all four categories. This potentially indicates that the change in ionising flux, rather than the

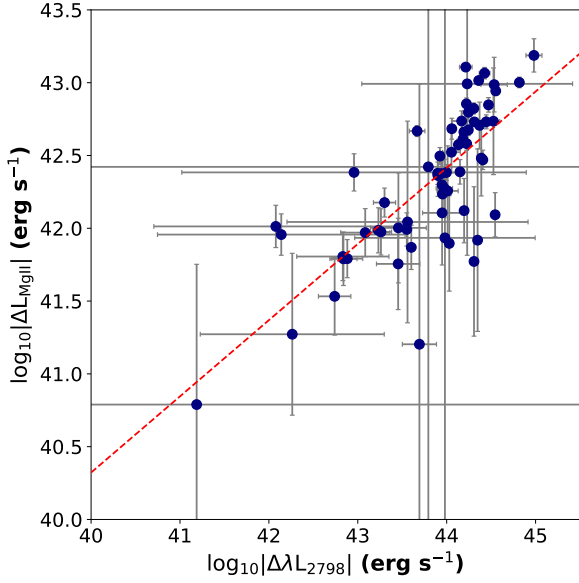


Figure 4. There is a positive correlation between the changes in line (L_{MgII}) and continuum (L_{2798}) luminosity for the Supersample: a linear regression of $\log L_{\text{MgII}}$ on $\log L_{2798}$ (red) gives a positive slope of 0.52. The plot displays the changes in luminosity from epoch to epoch, including the information from all spectra available for every objects. The horizontal axis represents the absolute value of the change in continuum luminosity at 2798\AA and the vertical axis the absolute value of the change in the Mg II line luminosity.

elapsed time is the driving force behind the variability of Mg II as well as of $H\beta$.

5 HOW VARIABLE CAN MG II FLUX BE?

Using the methods outlined in section 3 we fit and normalise the data in our two samples. The results of the spectral fitting of the Supersample are summarised in Table 4. The parameters calculated for the Supersample represent the behaviour of the Mg II line in a sample with a large range in time-spans between observations.

5.1 Luminosities

An overview of the change in line and continuum luminosities of the Supersample, using the absolute values of the changes, is shown in Figure 4. The data in this figure concern the change from epoch to epoch in total energy emitted per second. A linear regression of the change in line flux on the change in continuum gives a slope of 0.52. This indicates a positive but non-linear correlation between the two parameters. The plot also illustrates the large range in luminosities covered by the Supersample.

Table 4. Overview of the results of the spectral fitting for the Supervariable sample, listed by SDSS in column 1. Column 2 indicates whether $|\Delta f_{\text{MgII}}| \geq 0.5$ (\checkmark), is smaller than 0.5 (\times), or the continuum is stable (-). Column 3 lists whether the object was identified as a CLQ in MCL19. Column 4 is the time span between the spectral epochs, in the AGN restframe. Columns 5 and 6 list the changes in line and continuum luminosity, columns 7 and 8 are the fractional changes in line flux, and column 9 the change in line width. For objects with more than two spectra available the values in columns 5–9 are between the two spectra with the largest $|\Delta L_{2798}|$. Columns 10 list the measured QSO redshift, and columns 11–13 in turn give \log_{10} of the the bolometric luminosity, the black hole mass and the Eddington ratio, as provided by Shen et al. (2011). Column 14 provides the maximum absolute change in g-band photometry over the period between spectral observations, based on SDSS, PS1 (Heinis et al. 2016) and Catalina Real Time Survey (CRTS; Drake et al. 2009) data. The final column lists the telescope with which the latest spectral observation was made. A machine-readable version of this table is available in the Online Supplement.

SDSS name	Mg II	CLQ	max Δt_{spec} (RF days)	ΔL_{2798} (10^{42} erg s^{-1})	ΔL_{MgII} (10^{42} erg s^{-1})	Δf_{2798}	Δf_{MgII}	$\Delta \sigma$ (\AA)	z	L_{Bol}	M_{BH}	$\frac{L_{\text{bol}}}{L_{\text{Edd}}}$	max $ \Delta g $	Telescope	
1	001113.46-110023.5	\checkmark	\checkmark	3906	39.5 ± 3.7	0.9 ± 0.2	0.69	0.5	4.2 ± 3.6	0.495	45.2	8.2	-1.1	1.2	MMT
2	001206.25-094536.3	\checkmark	\checkmark	3484	135.6 ± 21.9	2.2 ± 1.8	0.87	0.54	17.1 ± 11.1	0.566	45.5	8.5	-1.1	2.0	Magellan
3	002311.06+003517.5	\checkmark	\checkmark	4378	319.9 ± 10.8	4.9 ± 0.4	0.91	0.55	-4.9 ± 13.2	0.422	45.5	9.2	-1.8	1.0	MMT
4	002450.50+003447.7	\times	\times	3795	110.4 ± 9.2	1.3 ± 0.2	0.69	0.33	-10.7 ± 6.7	0.524	45.5	8.7	-1.3	1.1	MMT
5	002714.21+001203.7	-	\times	4086	-17.2 ± 53.6	-0.5 ± 0.8	-0.08	-0.17	0.5 ± 1.4	0.454	45.5	8.2	-0.8	0.5	MMT
6	010529.61-001104.0	\times	\times	3353	29.4 ± 15.0	0.4 ± 0.4	0.52	0.12	-3.6 ± 6.8	0.738	45.2	8.3	-1.2	0.8	MMT
7	013458.36-091435.4	\times	\checkmark	4041	168.0 ± 7.6	0.6 ± 0.2	0.59	0.1	-8.4 ± 7.1	0.443	45.6	8.5	-1.0	1.4	MMT
8	022556.07+003026.7	\checkmark	\checkmark	2920	68.9 ± 2.3	2.0 ± 0.1	0.88	0.82	-9.9 ± 11.3	0.504	45.4	8.4	-1.1	1.7	SDSS
9	022652.24-003916.5	\checkmark	\checkmark	2206	30.0 ± 2.5	0.9 ± 0.3	0.75	0.5	3.0 ± 16.6	0.625	45.3	7.5	-0.3	1.4	SDSS
10	034144.72-002854.3	\checkmark	\times	3419	-100.8 ± 13.9	-4.4 ± 0.4	-1.98	-1.8	5.8 ± 11.5	0.623	45.4	9.9	-2.6	1.2	MMT
11	034330.48-073703.1	\times	\times	2404	97.0 ± 10.7	2.3 ± 0.3	0.71	0.46	20.1 ± 5.6	0.742	45.5	8.7	-1.3	1.0	Magellan
12	074502.90+374947.0	\checkmark	\times	3895	66.2 ± 3.5	1.5 ± 0.2	0.69	0.77	21.3 ± 8.2	0.593	45.3	8.4	-1.2	1.3	MMT
13	074719.65+215245.9	\times	\times	3674	39.7 ± 4.2	1.5 ± 0.2	0.62	0.36	-4.9 ± 8.5	0.462	45.2	8.7	-1.6	1.2	MMT
14	075228.65+405931.7	\times	\times	4349	48.9 ± 4.9	1.2 ± 0.1	0.34	0.19	1.1 ± 1.2	0.423	45.4	8.2	-0.9	1.3	MMT
15	081916.15+345050.3	\times	\times	767	19.7 ± 2.1	0.4 ± 0.1	0.43	0.17	7.0 ± 5.4	1.500	45.9	9.3	-1.4	2.1	WHT
16	090104.75+345524.2	\checkmark	\checkmark	2855	140.4 ± 48.0	4.0 ± 1.7	0.78	0.73	-0.3 ± 21.2	0.564	45.7	9.2	-1.6	1.8	MMT
17	092223.17+444629.0	\checkmark	\times	3322	74.5 ± 419.0	4.3 ± 24.0	0.76	0.76	11.5 ± 8.6	0.536	45.5	8.9	-1.5	1.4	MMT
18	094132.89+000731.1	\times	\times	3911	28.2 ± 325.2	1.2 ± 13.2	0.37	0.37	-0.2 ± 4.8	0.489	45.5	8.0	-0.7	1.2	MMT
19	102016.04+014231.7	\times	\times	3746	-54.2 ± 10.9	-1.6 ± 0.4	-0.79	-0.44	-5.6 ± 8.0	0.448	45.3	8.1	-0.9	1.4	MMT
20	111348.64+494522.4	\times	\times	3071	88.3 ± 10.0	1.1 ± 0.7	0.44	0.12	-2.8 ± 5.6	0.659	45.7	9.2	-1.6	1.5	WHT
21	123228.08+141558.7	\times	\times	3084	53.1 ± 5.2	-0.4 ± 0.3	0.62	-0.2	-2.5 ± 9.0	0.427	45.3	8.6	-1.4	1.2	MMT
22	132815.50+223823.8	\checkmark	\times	2149	78.6 ± 5.8	2.1 ± 0.2	0.76	0.63	-2.8 ± 7.7	0.486	45.4	8.8	-1.5	1.4	WHT
23	145519.88+364800.4	-	\times	3043	3.9 ± 19.2	-1.0 ± 0.4	0.03	-0.32	9.8 ± 7.8	0.525	45.5	8.6	-1.1	1.3	WHT
24	152044.63+394930.3	-	\times	3581	23.9 ± 10.4	-0.1 ± 0.3	0.38	-0.04	4.7 ± 7.0	0.439	45.2	8.5	-1.4	1.2	MMT
25	153734.06+461358.9	\checkmark	\checkmark	3704	92.2 ± 17.6	3.3 ± 0.6	0.76	0.79	-24.2 ± 3.2	0.378	45.3	8.5	-1.3	1.4	MMT
26	153912.76+524540.2	\checkmark	\checkmark	3896	43.1 ± 8.9	1.2 ± 0.3	0.66	0.78	-10.7 ± 4.8	0.415	45.3	8.3	-1.2	1.1	MMT
27	154341.64+385319.3	\times	\times	3241	109.3 ± 14.6	-0.4 ± 1.1	0.67	-0.11	-3.4 ± 7.4	0.432	45.4	9.1	-1.8	1.2	MMT
28	160226.01+323019.1	\times	\times	3350	83.6 ± 8.8	0.2 ± 0.3	0.4	0.07	7.0 ± 6.9	0.557	45.7	8.8	-1.1	1.1	MMT
29	162415.02+455130.0	\checkmark	\checkmark	3479	93.9 ± 5.7	2.5 ± 0.1	0.83	0.73	6.4 ± 5.9	0.481	45.4	8.5	-1.2	1.2	WHT
30	163031.47+410145.8	\times	\times	3603	88.2 ± 5.9	2.9 ± 0.3	0.51	0.49	-17.5 ± 10.5	0.531	45.5	8.8	-1.4	1.3	MMT
31	164053.05+451033.7	\times	\times	3766	47.7 ± 197.6	0.4 ± 11.3	0.63	0.23	-7.2 ± 4.3	0.422	45.4	8.9	-1.6	2.6	MMT
32	212436.64-065722.1	\times	\times	3734	76.7 ± 6.5	0.6 ± 0.3	0.51	0.2	-8.1 ± 4.2	0.430	45.4	8.3	-1.0	2.1	MMT
33	214613.31+000930.8	\checkmark	\checkmark	1547	-43.5 ± 4.2	-1.8 ± 0.3	-4.1	-1.44	13.5 ± 24.2	0.621	45.2	8.9	-1.8	0.5	SDSS
34	223133.89+003312.7	\times	\checkmark	3917	41.4 ± 4.3	0.8 ± 0.2	0.45	0.29	-12.3 ± 6.1	0.476	45.4	8.6	-1.3	1.2	MMT
35	223500.50-004940.7	\checkmark	\times	3559	-86.4 ± 36.7	-1.9 ± 1.3	-1.45	-0.71	-4.5 ± 9.0	0.642	45.1	8.5	-1.5	0.5	MMT
36	224017.28-011442.8	\checkmark	\times	3396	-62.2 ± 6.1	-1.1 ± 0.2	-1.23	-0.96	-20.1 ± 4.9	0.502	45.1	8.3	-1.3	1.2	MMT
37	224829.47+144418.0	\checkmark	\checkmark	3634	173.2 ± 4.4	4.3 ± 0.3	0.87	0.66	25.6 ± 7.7	0.424	45.5	8.9	-1.5	1.5	MMT
38	225240.37+010958.7	\checkmark	\checkmark	3529	98.7 ± 5.5	4.5 ± 0.3	0.8	0.63	5.7 ± 7.5	0.534	45.3	8.9	-1.7	1.7	MMT
39	230427.86-000803.2	\checkmark	\times	3632	-63.3 ± 16.1	-2.3 ± 0.4	-1.07	-1.46	-8.0 ± 6.4	0.526	45.2	8.7	-1.6	1.1	MMT
40	231552.95+011406.1	\times	\times	3643	-17.9 ± 14.0	-0.1 ± 0.3	-0.51	-0.12	-5.3 ± 9.1	0.521	45.1	8.5	-1.6	1.1	MMT
41	231742.68+011425.2	\times	\checkmark	3654	40.6 ± 11.9	0.8 ± 0.5	0.53	0.43	-1.9 ± 17.7	0.518	45.3	9.0	-1.8	1.6	MMT
42	233317.38-002303.4	\checkmark	\checkmark	3821	66.8 ± 2.1	1.8 ± 0.1	0.79	0.64	-0.2 ± 7.9	0.513	45.3	10.2	-3.0	1.7	Magellan
43	234623.42+010918.1	\checkmark	\times	3635	-71.5 ± 11.7	-5.6 ± 0.2	-2.9	-3.64	-7.0 ± 9.7	0.509	45.3	8.9	-1.7	2.0	MMT

5.2 Flux Normalisations

The results of the normalisation to the epoch of the highest continuum are presented in Figure 5. All fluxes in this figure are fractional changes in flux: the further away from unity the normalised value is, the greater the relative change. The data for both samples show a strong scatter, but there is a pattern to the relation between line and continuum fluxes. Large *relative* changes in continuum are matched by large relative changes in line emission. For smaller changes in continuum flux the scatter in line fluxes becomes greater. The distribution of fluxes in the Full Population sample is concentrated in the region where normalised line and continuum fluxes are ~ 1 . These are spectral pairs where the overall flux change has been small, likely relating to quasars that were not varying significantly over the period of observation. The population does show a clear linear correlation, however.

For the Full Population sample we have visually examined the quality of the data and fits for 236 of the strongest outliers in the normalised sample. In a small fraction of the cases ($\sim 5\%$) large changes were associated with changes in data quality between observations, such as missing data in the spectrum that partially covered the Mg II range in one of the epochs. A relatively low S/N in both spectra may have affected the reliability of the fit in $\sim 30\%$ of the spectra. However, for most objects the measured flux change is visible in the spectrum. Given the small fraction of outliers overall, we are therefore confident in the identification of the trend visible in the right panel of Figure 5.

The data in Figure 5 are fit with the response functions defined in Section 3. A summary of results of the relevant statistical tests is given in Table 5. For the Supervariable sample the results of the one-component fit are included in red in the plot. The right bottom panel shows the associated residuals, as well as the χ_{red}^2 . The sequence test statistic, S (section 3), for this fit is 31 corresponding to a p-value (p_s) of 0.1. Fitting the same functions to the Full Population sample, we find a slight preference for the two-component function (Table 5).

For the normalisation to the minimum epoch (Figure 6) the scatter in both samples is clearly visible. The distribution of the data of the Supervariable sample suggests a pattern different from simple linear correspondence. There is a sharp initial rise, followed by a possible flattening of the responsivity. Note that the scatter for $f_c \gtrsim 6$ is particularly large. This is reflected in a lower Spearman correlation coefficient than for the normalisation to the maximum epoch, 0.67 versus 0.78 (Table 5). The tentative pattern in the data can be tested using the two responsivity functions. The results of both fits are included in Figure 6. The residuals are shown in the bottom two panels. The reduced χ^2 value is lower for the two component fit, but the fact that the value is less than one indicates this function is likely overfitting the data. The results of the sequence test (listed in Table 5) indicate that neither fit describes the data particularly well.

For the Full Population sample (right panel in Figure 6), the coefficients (Table 5) clearly indicate a correlation. Evidence of levelling off is not as clear as it was for the Supervariable sample. This is possibly due to the strong concentration of the distribution around objects with relatively small flux and continuum changes. A firm conclusion is that the Mg II line and UV continuum flux are clearly

Table 5. Correlation coefficients and the fitting statistics for the two normalisations, for the Supervariable and Full Population samples. The top row lists the Spearman (r_s) and Pearson (r) correlation coefficients. The bottom rows list the χ_{red}^2 and the p-value associated with the sequence test statistic S, for the one and two component linear fits.

	Maximum		Minimum	
	Supervar.	Full Pop.	Supervar.	Full Pop.
r_s	0.78	0.32	0.67	0.30
r	0.76	0.30	0.46	0.32
1C: χ_{red}^2	1.01	1.15	1.39	1.28
p_s	0.10	~ 1	0.94	~ 1
2C: χ_{red}^2	0.92	1.02	0.77	0.93
p_s	0.47	~ 1	0.97	0.98

Table 6. A comparison of various metrics of responsivity between the Supervariable and the Full Population sample for both normalisations. The included values are the slope of a linear fit to the data, the Pearson coefficient (r) and the Spearman coefficient (r_s).

	Maximum			Minimum		
	Slope	r	r_s	Slope	r	r_s
Supervariable	1.39 ± 0.03	0.76	0.78	0.58 ± 0.03	0.46	0.67
Full Population	1.13 ± 0.00	0.30	0.32	0.74 ± 0.01	0.32	0.30

correlated and scale linearly. This is an ensemble effect: the large scatter suggests that individual objects can show a different correspondence between line and continuum. The difference from object to object will be a combination of intrinsic differences between the objects, and an artefact of the cadence of observations. The latter is caused by the fact that we only have two snapshots of the full line response, which means we will observe the line in different stages of its response for every target (in agreement with the evolutionary ‘sequence’ proposed in Guo et al. (2020)).

The metrics of correlation and responsivity for the two samples are summarised in Table 6. In all cases the Pearson and Spearman coefficients show a strong correlation between the line and the continuum flux. The columns labelled ‘slope’ list the gradient of the one-component linear fit to the data. The error on the slope was calculated in a Monte Carlo simulation (N=10,000) by removing a random selection of data points each iteration, 10 for the Supervariable sample and 3,000 for the Full Population Sample, and refitting. The slopes for the two data-sets differ significantly. For both normalisations the gradient for the Supervariable sample indicates a stronger response of the line to the continuum than in the Full Population sample. This means that for a sample dominated by large continuum flux changes the line responsivity differs from a sample dominated by quasars with a more tranquil continuum.

5.2.1 The Responsivity Measure

The distribution of the responsivity measure, α_{rm} (section 3), in the two samples is shown in Figure 7. The responsivity measure was chosen to easily identify over-responsive and under-responsive objects in the data set. Objects with $\alpha_{rm} > 0$ have a larger relative change in line flux than in

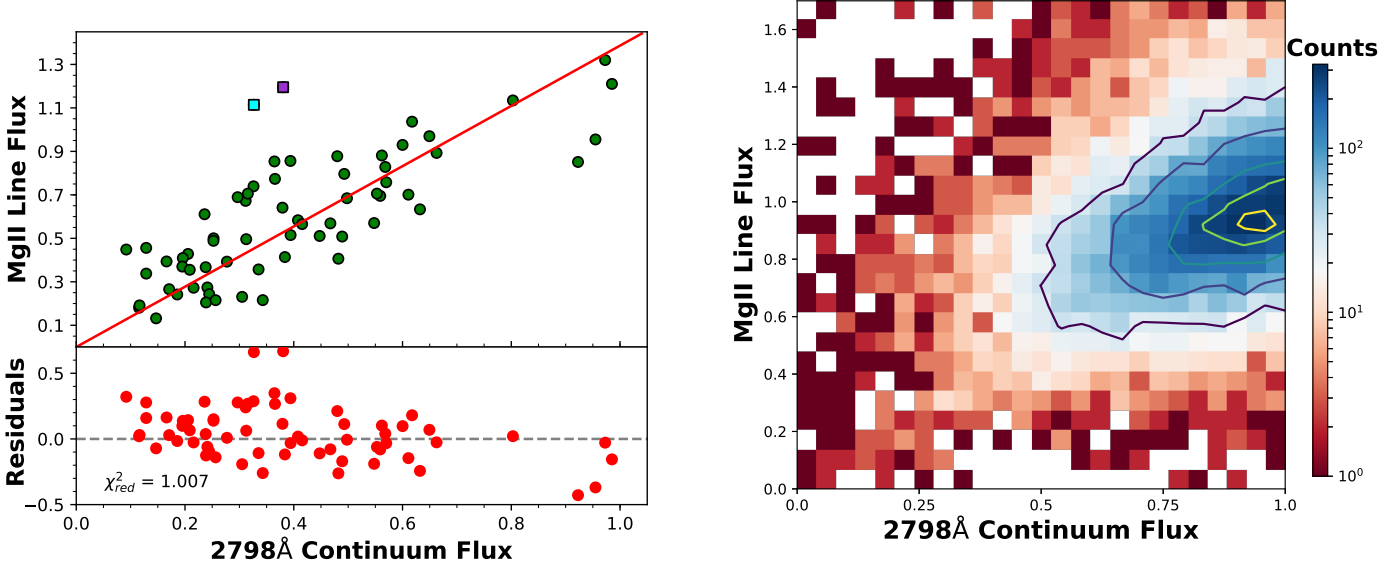


Figure 5. *left,top*: Fluxes in the Supervariable sample, normalised to the epoch of *maximum* continuum. There is a correlation between large continuum changes and large line flux changes. The plot shows the line flux plotted against the continuum as well as the fitted linear responsivity function. Two outlying data points have been marked with squares: J154341 on MJD 57818 (cyan) and J123228 on MJD 57845 (purple). *left,bottom*: The residuals and χ^2_{red} for the response function. *right*: The Mg II line and 2798Å continuum fluxes for the Full Population sample, normalised to the epoch of maximum continuum. The data show a clear correlation. Most objects in the sample are concentrated on the right side of the plot, exhibiting only small changes in both line and continuum. The colours in this 2D histogram correspond to the number of objects. The contours are placed at 10, 25, 50, 75, and 90% of the maximum count number for the sample.

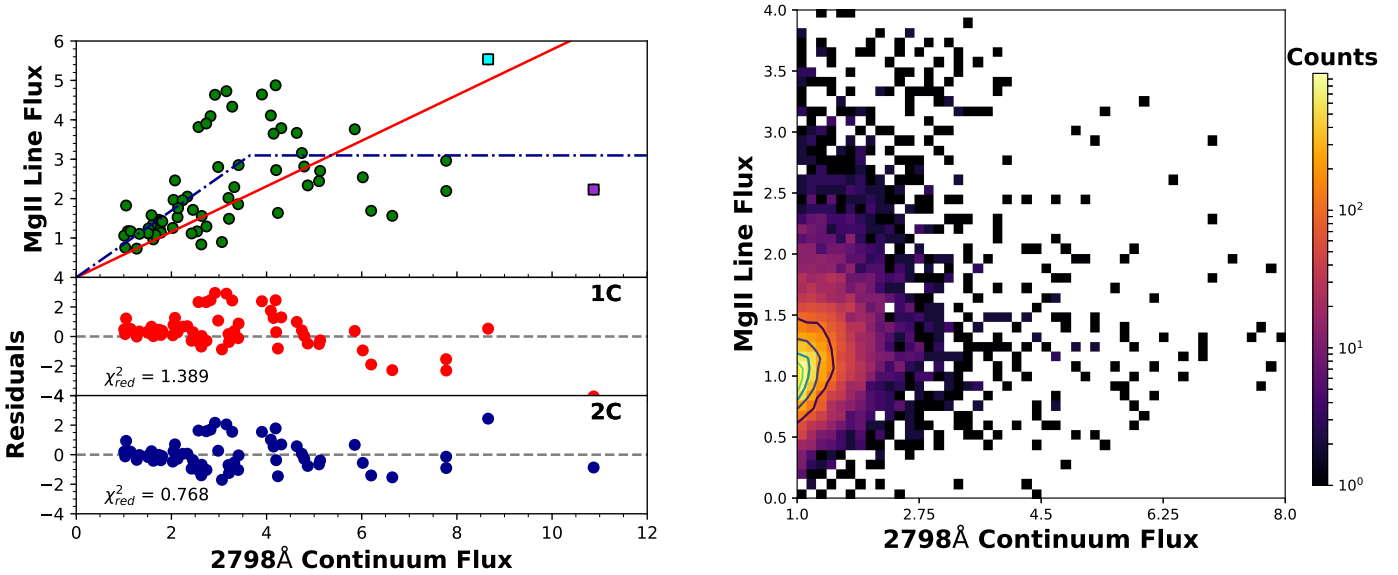


Figure 6. *left,top*: Normalised to the epoch of *minimum* continuum level, the data for the Supervariable quasar sample show a significant scatter. This is particularly clear for the largest fractional flux changes. The panel also shows the two fitted responsivity functions. Two outliers are marked with squares: J022556 on MJD (cyan) and J002311 on MJD 55480 (purple). These objects are discussed in more detail at the end of section 5. *right,bottom*: The residuals from the two fitted functions as well as the χ^2_{red} for the fits. The two-component fit (lower) results in a smaller reduced χ^2 value. As $\chi^2_{red} < 1$, this could be a result over overfitting. *left*: The distribution in the $f_{2798}-f_{MgII}$ parameter space of the Full Population sample, when normalised to the epoch of minimum continuum. The distribution shows a linear correlation. The layout of this figure is the same as that of the right panel in Figure 5.

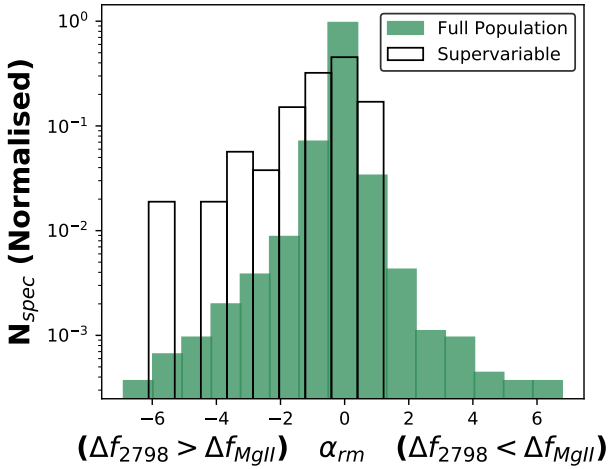


Figure 7. A comparison of the distribution of α_{rm} in the two Mg II samples. Both the Supervariable and the Full Population sample appear skewed to negative values of α_{rm} , although this trend is more prominent for the Supervariable sample. Negative values are associated with an under-responsive Mg II line flux and positive values with over-responsivity.

the continuum and can be considered over-responsive. Conversely, $\alpha_{rm} < 0$ indicates an under-responsive object. The distribution of the Supervariable sample is clearly skewed to the negative values. The same skewness can be seen in the distribution of α_{rm} for the Full Population sample, although in this case the distribution is strongly dominated by objects with $\alpha_{rm} \sim 0$. The α_{rm} distribution therefore shows the same pattern as observed in the normalisation plots. Most objects show only limited variability of both line and continuum, resulting in a responsivity measure close to zero.

There are a number of spectra on the positive side of the distribution, representing strong responses in the Mg II line. For the objects that show a response greater than 1:1, the extremes in α_{rm} are not as large as for the low-responsivity cases. This implies that the difference between fractional flux changes in continuum and line flux for objects with $\alpha_{rm} > 0$ is overall smaller than for the less responsive objects. It should be noted, however, that if the average response function is shallower than a 1:1 response the objects with $\alpha_{rm} > 0$ are truly exceptionally responsive.

The responsivity parameter can also be used to quantify the link between Mg II and H β variability. As seen in Table 3, the correspondence between CLQ behaviour and Mg II variability varies. Figure 8 shows the distribution of α_{rm} for CLQ and non-CLQ objects in the Supervariable sample. For objects with more than one spectrum the average of the responsivity measures is used. The distribution for the CLQs is clearly broader than for the non-CLQ objects and stretches to more negative values. For both distributions the centre lies slightly below 0, corresponding to objects where Mg II only slightly tracks the continuum. We therefore again see a range of behaviours: some CLQs also show strong variability in Mg II, whereas for some CLQs Mg II barely responds at all to the changing continuum.

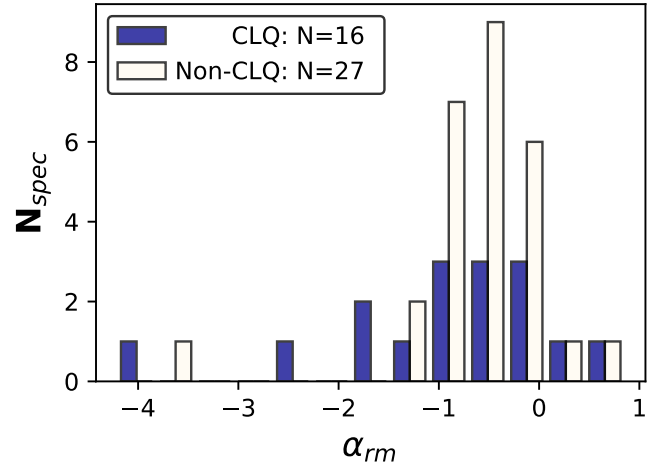


Figure 8. The distribution of the responsivity measure among the CLQs and non-CLQs in the Supervariable sample. The average value of α_{rm} is used for objects with more than two spectra available.

5.2.2 Interesting Outliers

The responsivity functions and α_{rm} are used to define the response of the Supervariable sample. Compared to these measures there are individual objects that stand out. The first of these are identified in Figure 5, using the normalisation to the epoch of maximum continuum. The data points are found at the top left of the plot. These objects showed a large decrease of the continuum associated with an increase in the line flux. The spectra for the two objects, J123228 and J154341, are shown in Figure 9. Neither of the objects lose their broad H β flux. This could indicate that the response of the BLR to the drop in continuum has not occurred yet. The elapsed rest-frame time in both cases is close to two decades. An alternative hypothesis is that the observed line fluxes are in response to a flare in the continuum that occurred in the intervening years. The increasing line flux in J154341 seems coupled to a narrowing of the line, indicating a change in the physical conditions in the line forming region.

The clearest outlier with respect to the normalisation to the lowest continuum (Figure 6; purple square), shows a continuum change of a factor ~ 10 coupled to a line flux change of a factor ~ 2 . The object associated with this spectrum is J002311. The spectra for this object are shown in panel *c* of Figure 3. The spectrum marked in the right panel of Figure 6 was taken on MJD 55480. The high state continuum was associated with a strong broad H β line, which disappears in the lower state spectra. The cadence of observations for J002311 allows for a more detailed view of the evolution. The object is varying rapidly: it doubles its continuum flux on a restframe timescale of ~ 14 years, and drops it by a factor 10 in the following decade. It is plausible the observation was made before the line had had enough time to respond in full to the continuum changes.

Another interesting spectrum is marked in cyan in Figure 6. This epoch shows a line flux change greater than a factor 5. The spectrum belongs to J022556, for which

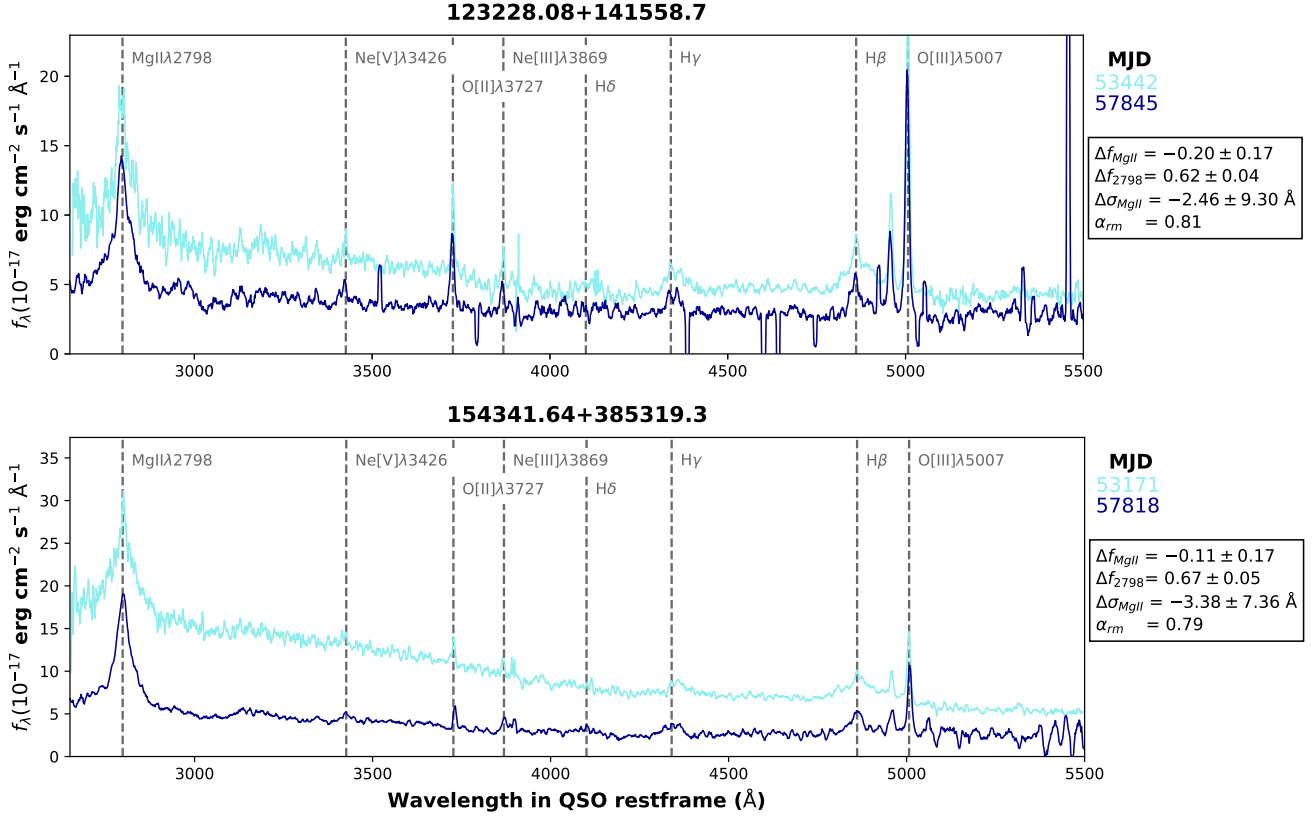


Figure 9. The spectra for J123228 (*top*) and J154341 (*bottom*), the two outliers marked in the left panel of Figure 5. Both objects show large continuum changes, but a limited response of the Mg II flux. The change in H β is also limited for both objects.

we have a high number of spectra available because the object was part of extra deep BOSS plates³. This QSO was identified as a CLQ in MCL16, where its strong Mg II variability was also noted. The high cadence of observations allows us to see there is evidence of varying responsivity for this object.

Figure 10 shows a close-up of the changing Mg II line for J022556 (*left*) as well as the line and continuum fluxes using the normalisation to the maximum continuum (*right*). The spectra cover ~ 8 years in the QSO restframe. The shading of the markers indicates the progression of time. The first epoch available (SDSS, MJD 52200) is represented by the light blue marker in the centre of the plot. The initial rise in the continuum is to the epoch of maximum continuum (top right). The gradient of this response is approximately unity. In the next spectrum the continuum has collapsed, and the line has also almost disappeared (bottom left corner). When the continuum rises again, the gradient is considerably steeper.

The epoch dependent responsivity visible in J022556 (and J002311) indicates that the scatter around the responsivity functions (in Figures 5 and 6) could be the result of sampling the varying objects at different stages during the BLR response to the continuum. The response of the line to the continuum does not follow a single gradient for *the same*

object, therefore we can expect considerable scatter around any linear response function fitted to a large ensemble.

5.3 The Intrinsic Baldwin Effect

A final measure for the responsivity is the intrinsic Baldwin Effect (iBE), which measures the change in EW compared to the change in continuum luminosity for individual objects. For the Full Population sample the iBE is calculated as $\Delta \text{EW} / \Delta L_{2798}$ for each spectral pair. For the Supervariable sample the same is done for each chronologically sequential pair. The results for the Full Population sample are shown in Figure 11.

The distribution is largely symmetric and centred around a slope of -0.5 . For the Supervariable sample the average slope is -0.54 . Unfortunately there are no literature values of the Mg II iBE to compare these results to. For other UV lines such as CIV $\lambda 1549$ and for the Balmer lines the slope of the intrinsic BE is significantly steeper than for the ensemble BE (Osmer & Shields 1999; Rakić et al. 2017). We find that the same holds for Mg II: the eBE of the DR7 sample has a slope of -0.16 , shallower than the slope of the iBE.

³ https://www.sdss.org/dr12/spectro/special_plates/

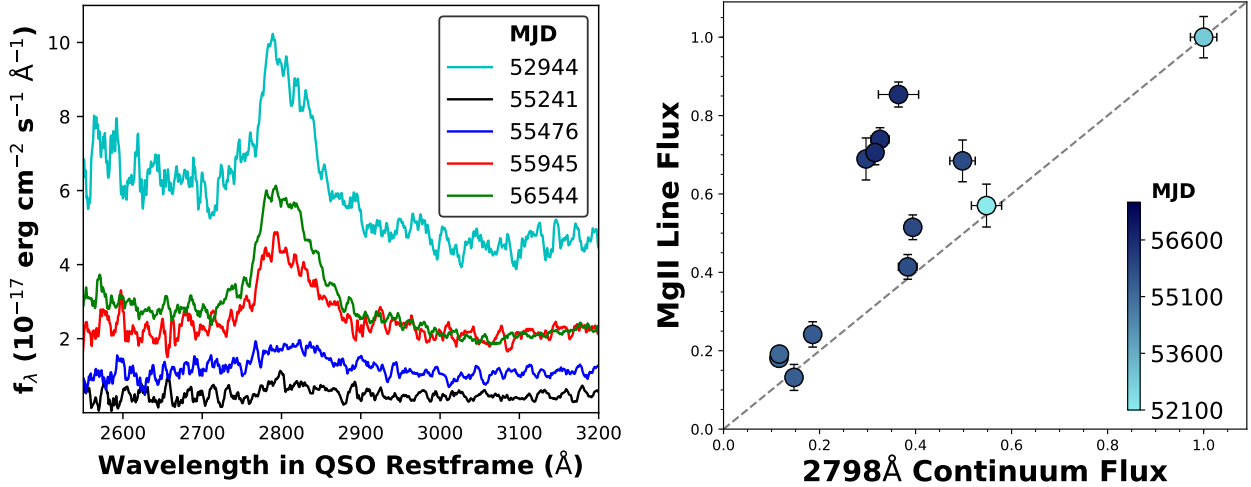


Figure 10. *left:* The changing Mg II line flux in J022556. The high cadence of observations for this object allows us to track the line flux as the continuum falls and rises over the course of ~ 8 years in the QSO restframe. *right:* The Mg II line fluxes normalised to the maximum continuum. The epochs of the spectra are indicated with the colour, increasing chronologically from light to dark. This plot is based on the normalisation of the fluxes to the epoch of maximum continuum. A 1:1 response is indicated with the dashed grey line. The initial increase in line flux is matched approximately 1:1 by the Mg II line. Following a collapse in line and continuum, however, the gradient of the line responsivity is considerably steeper as the continuum recovers.

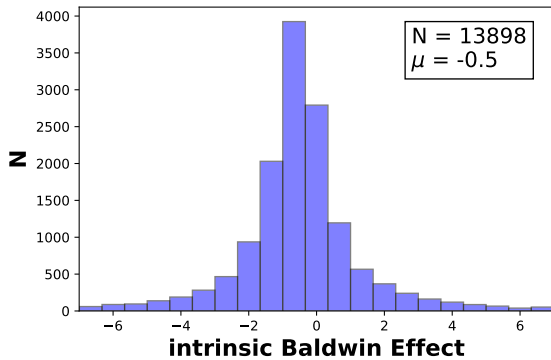


Figure 11. A histogram of the intrinsic Baldwin Effect for the Full Population sample. The number of spectral pairs, N , and the sample average, μ , are listed in the top right corner. The iBE is calculated as $\Delta EW / \Delta L_{2798}$ for each pair. Although there are some strong outliers, the distribution is clearly centred around -0.5 . This is significantly lower than the slope of the ensemble BE of -0.16 calculated for the data from [Shen et al. \(2011\)](#).

5.4 Line Responsivity over Time

If the line responsivity changes over time, this should manifest itself as a dependency on Δt , the elapsed time between the spectral measurements for each pair. The Full Population sample is large enough to investigate the effect of Δt in more detail. The first test is to divide the data-set into Δt bins and calculate the normalised fluxes in the subsets as in the previous section. The results are shown in Figure 12. To emphasise the visualisation of the outliers, the normalisation to the epoch of minimum continuum is used. To ensure a good sample size in each Δt bin, the bin size increases for longer elapsed times. The time bins represent re-observations of the object within one year, between one and two years, between two and three-and-a-half years, and

three-and-a-half up to 10 years respectively. All counts have been normalised by dividing by the maximum per plot.

Two changes take place in the Δt bins in Figure 12: the shape of the contours changes and the gradients of the line responsivity appear to decrease. The first effect is visible as an elongation of the contours as Δt increases. This implies that large changes in the continuum are more common for large Δt . Changes in flux that are very unlikely to occur within the first year become more likely as Δt increases. A longer waiting time means a greater chance of observing a large change in UV continuum. This agrees with established results on quasar continuum variability ([MacLeod et al. 2010](#); [Welsh et al. 2011](#)). The change in shape also implies that for greater Δt the distribution has a larger spread.

The second change visible in Figure 12 is in the slope of the linear correlation and can be seen in the plots as a gradual tilting of the contours. The contrast is most pronounced between the shortest and longest time differences: as a visual aid the contours from the top left plot have been included as dotted lines in the bottom right plot. Although the effect is subtle, the slope decreases over time. This effect can be further quantified by a linear fit to the data. This is the same one-component fitting function used in the previous section. The fitted slopes and the errors are listed in Table 7. The results show a steady decline with Δt in the slope of the responsivity. For a larger elapsed time between observations the response of the line to continuum changes therefore appears to be smaller.

A different approach to investigating the effect of elapsed time on the line responsivity is illustrated in Figure 13. Here we compare the distribution of the normalised change in line flux (Δf_{MgII}) for several Δt subsets. The top panel of Figure 13 shows the distribution of Δf_{MgII} against Δt . The Full Population sample is skewed toward objects with a relatively small Δt , as is also evident in Figure 1. The sections in

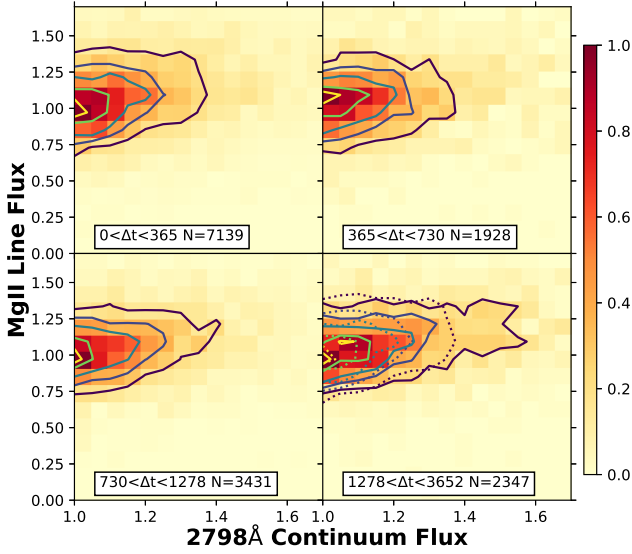


Figure 12. The Full Population sample normalised to the minimum continuum state, binned by Δt between observations, in days in the QSO restframe. The contours are set to 10, 25, 50, 75, and 90% of the maximum value. To observe any difference in response over time the contours from the shortest elapsed time bin (*top left*) are included in the panel for the longest elapsed time (*bottom right*). The maximum number of counts per bin in each panel has been normalised to one, to allow for comparison between the four subsets.

Table 7. Slopes of the linear correlation between normalised f_{MgII} and f_{2798} for the Δt subsets shown in Figure 12. The slope shows a slight but steady evolution with Δt : the line responsivity decreases for longer time intervals between measurements. The error on the slopes is calculated from the covariance matrix produced by the `scipy` fitting routine. The error on the measurements was estimated based on the scatter, as described in section 3. Δt is given in the QSO restframe.

RF Δt (days)	Slope	Uncertainty
$0 < \Delta t < 365$	0.746	0.006
$365 < \Delta t < 730$	0.745	0.011
$730 < \Delta t < 1278$	0.738	0.007
$1278 < \Delta t < 3652$	0.724	0.007

Δt have been chosen to cover both the full range of elapsed times and include enough objects in each subsample to allow for meaningful statistical analysis. The Δt bins are listed in the bottom three panels and the associated sample sizes are included in Table 8.

If there is no pattern to the flux changes, we expect the data in each bin to follow a normal distribution. For each of the three subsets the distribution of Δf_{MgII} is therefore compared with a normal distribution using the Lilliefors test. The bottom three panels of Figure 13 show the cumulative distributions of the three subsets, as well as the CDF of the normal distribution for a visual comparison. The results of the test are quite clear, rejecting the hypothesis the data are drawn from the normal distribution. The Lilliefors test provides a $p\text{-value} < 10^{-10}$ in each case.

The distributions of Δf_{MgII} in the Δt bins indicated in Fig-

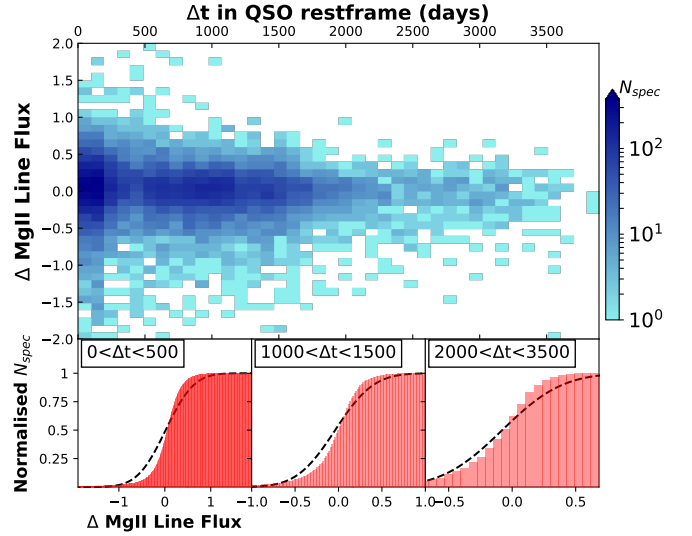


Figure 13. *top:* The normalised change in line fluxes (Δf_{MgII}) plotted against the elapsed time between observations, in days in the QSO restframe. *bottom:* Testing for normalcy on three different Δt cuts in the data shows a significant deviation from a normal distribution. The three cumulative histograms represent the sub-samples, marginalised over Δt . The normal CDF, calculated with the sample average and sample standard deviation, is plotted as the black dashed line as a visual aid for the comparison.

ure 13 and their deviation from a normal distribution can be quantified by calculating the sample moments. These are presented in Table 8. The moments of the distribution do change with time. The mean shifts to negative values as Δt increases. This is likely the result of a selection bias: the spectral pairs in the lowest Δt bin are more likely to have both observations made during the BOSS survey. For sources in the higher Δt bins the probability is larger that the first of the spectra is from SDSS I/II. Sources that were in a relatively bright state during SDSS I/II were more likely to be included in DR7Q and therefore to be included in the Full Population sample. This makes it more likely to see transitions to dimmer states. We conducted a similar analysis for f_{2798} and included the results in Table 8. The continuum flux distributions show the same pattern as those of the line flux.

The more interesting effects occur in the third and fourth moments of the distribution. The skewness indicates the distributions are skewed towards more positive values. This effect seems to decrease with Δt , however, particularly in the case of Δf_{MgII} . A similar trend with Δt can be seen in the kurtosis, which is a measure of the weight of the tails of a distribution. The two lower Δt bins have quite prominent tails. This is also evident from the scatter in the top panel of Figure 13. The moment is significantly smaller for the higher bins, and more so for the line flux than for the continuum flux. The reduced number of strong outliers could indicate a reduced responsivity for objects in this bin. The data in Table 8 indicate $|\Delta f_{2798}|$ and Δt are therefore correlated parameters, making it difficult to see which of the two parameters most strongly affects the observed changes in

the Mg II line flux. However, we do note that large line flux changes are observed on all time scales. This is an indication that large continuum changes are an important driving force for changes in the line flux.

6 THE MG II LINE PROFILE

6.1 Changes in the Line Profile

Line profiles provide insight into the kinematics of the line forming regions. The measure used to quantify the Mg II profile is the width (σ) of the broadest Gaussian fit to the line. The widths are defined in the QSO restframe. Table 4 lists the change in line width for the spectral pair with the largest change in continuum ($\Delta\sigma_{\text{MgII}}$) for all targets in the Supervariable sample. The distribution of $\Delta\sigma_{\text{MgII}}$ is shown in Figure 14. The distribution shows that most objects have relatively small changes in line width (< 10). However, there is also an extended tail towards larger changes. The large spread in the sample is another indication of the broad range of Mg II variability among extremely variable quasars.

The largest values of $\Delta\sigma_{\text{MgII}}$ are $\sim 20\text{\AA}$ for the Supervariable sample. Under the assumption of Keplerian rotation of the BLR gas this corresponds to a change in rotational speed (along our line of sight) of approximately 2,100 km/s. For Mg II the Keplerian speeds deduced from the FWHM are in the order $10^{3.6}$ km/s (Shen et al. 2011). The extremes in the observed changes in line width therefore imply large transitions.

For the Full Population sample the distribution is strongly peaked around $\Delta\sigma_{\text{MgII}} = 0$, with a slight asymmetry towards decreasing line widths. Large changes in σ are possible, but most objects exhibit only limited variability. This is likely the same effect as the observed concentration of flux changes around small values (Figure 7): most quasars in the sample have changed only slightly between epochs. For the Full Population sample we have visually inspected the spectral fits of 164 outliers in $|\Delta\sigma_{\text{MgII}}|$ to exclude the possibility of issues with the fit or the data quality as a cause for the behaviour.

Figure 15 shows the change in line width plotted against the change in Mg II flux, where both quantities have been normalised by dividing by the value from the first epoch of the pair. The two parameters represent a fractional change over time. The figure includes the result of a linear regression of $\Delta\sigma_{\text{MgII}}$ on Δf_{MgII} for the Supervariable sample (red). The slope of the line is 0.07 ± 0.05 . Although the scatter in the data is strong, it is possible that there is a positive correlation between the magnitude *and* direction of change of the line flux and the line width. For the Full Population sample a linear regression of $\Delta\sigma_{\text{MgII}}$ on the change in line flux has a gradient of 0.09 ± 0.01 .

For the Full Population sample we show a comparison of σ_{MgII} with the fitted line luminosities in Figure 16. For the bulk of the spectra the line width appears unrelated to either the line or the continuum luminosity. This would be in agreement with results from Shen et al. (2011), who find no correlation between the FWHM of Mg II and the 3000 \AA luminosity. This is in contradiction with the expectation based on BLR breathing and could be an indication that the Mg II line width is not entirely set by the gravitational potential

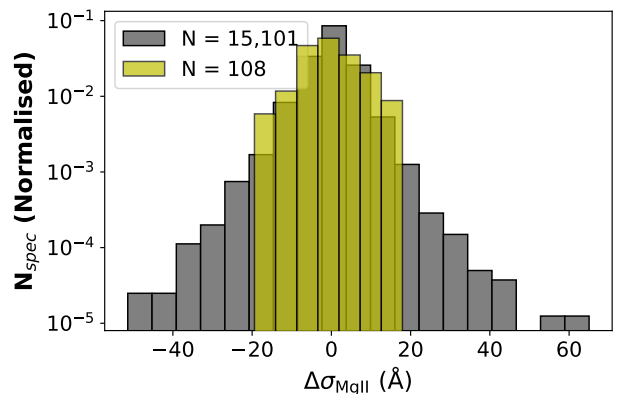


Figure 14. The distribution of $\Delta\sigma_{\text{MgII}}$ for the Full Population sample (grey) and the Supervariable sample (yellow). Each distribution has been normalised such that the surface area under the histogram equals one. N_{spec} is the number of spectra in each bin. The distributions are strongly centred around objects with no observable change in line width.

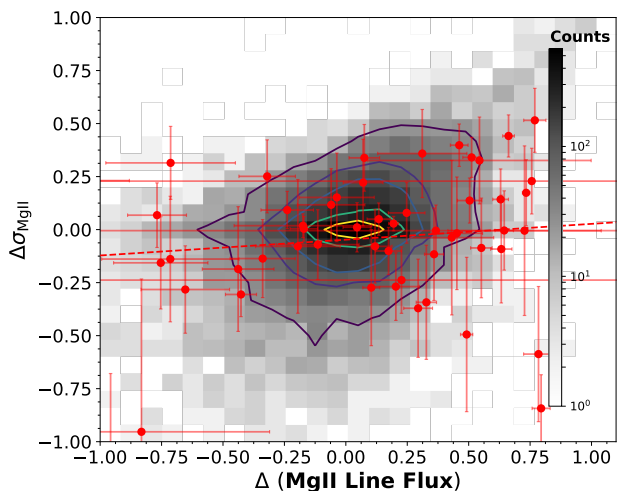


Figure 15. The change in line width ($\Delta\sigma_{\text{MgII}}$) plotted against the change in line flux. The values have been normalised to the preceding epoch and therefore represent a fractional change with respect to the older spectrum. The data for the Supervariable sample is shown in red, including the result of a linear regression of $\Delta\sigma_{\text{MgII}}$ on Δf_{MgII} . The Full Population sample is represented by the grey histogram. The contours indicate 5, 15, 25, 50 and 75% of the peak value of the histogram. For most objects in the Full Population sample the two parameters do not seem to correspond at all. For large changes in line flux, associated with more extreme variability, a connection is possibly present.

or that the location of the line forming region does not vary strongly with the continuum flux.

6.2 Skewed Profiles

There are five objects in the Supervariable sample that display a skewed line profile: J022556, J034144, J111348, J225240, and J234623. Of these J034144 and J111348 show

Table 8. The sample moments of the Δf_{MgII} and Δf_{2798} distributions for the Δt subsamples shown in Figure 13. Δt is in days in the restframe. There is an indication of change in the sample distributions over time. For the calculation of the moments the distributions have been filtered requiring $|\Delta f| < 7$ (i.e. no flux changes greater than a factor 7 between epochs).

	Δf_{MgII}			Δf_{2798}		
	$0 < \Delta t < 500$	$1000 < \Delta t < 1500$	$2000 < \Delta t < 3500$	$0 < \Delta t < 500$	$1000 < \Delta t < 1500$	$2000 < \Delta t < 3500$
N	6282	2472	502	6273	2475	499
μ	0.02	-0.04	-0.08	0.10	-0.08	-0.26
σ	0.48	0.35	0.37	0.39	0.42	0.78
skewness	-2.71	-3.16	0.04	-4.13	-3.54	-3.35
kurtosis	41.9	23.9	12.7	60.4	32.0	17.2

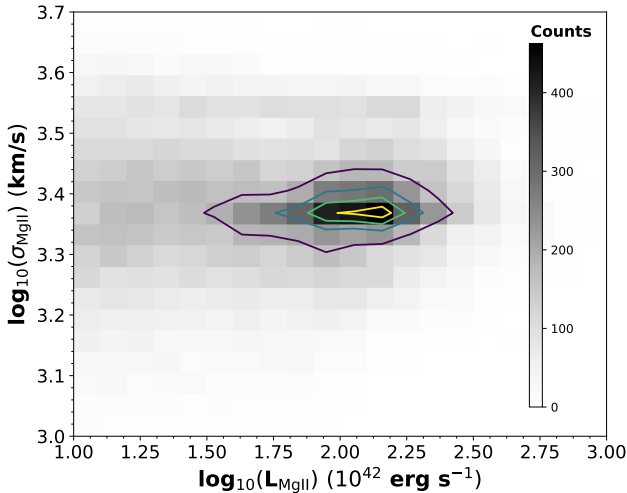


Figure 16. The line width (σ_{MgII}) plotted against the line luminosity for the Full Population sample. The contours represent 40, 60, 75, and 90% of the peak numbers of counts per bin. A correlation is possibly present for lower luminosities, however most objects do not show such a correlation. A lack of correlation between line width and continuum luminosity would be in agreement with the findings in Shen et al. (2011).

a skewness to the red, the others to the blue. There are more spectra that can tentatively be identified as asymmetric, but the data quality does not allow for a definitive identification. Two examples of skewed line profiles in our sample are shown in Figure 17.

Only one of the five objects, J111348, shows an indication of a variable line profile over time. Between the epochs, the maximum of the line, possibly associated with a distinct component, appears to shift from the centre to the red. The two epochs are shown in Figure 18. The change to a more asymmetric profile was associated with a large drop in flux, where the 2798Å continuum dropped by a factor of approximately 4.

7 CONNECTION TO PHYSICAL PARAMETERS

To investigate whether the differences in spectral variability are connected to physical parameters, we will compare different subsets of the samples. The most important parameters to consider are the bolometric luminosity (L_{bol}), the black

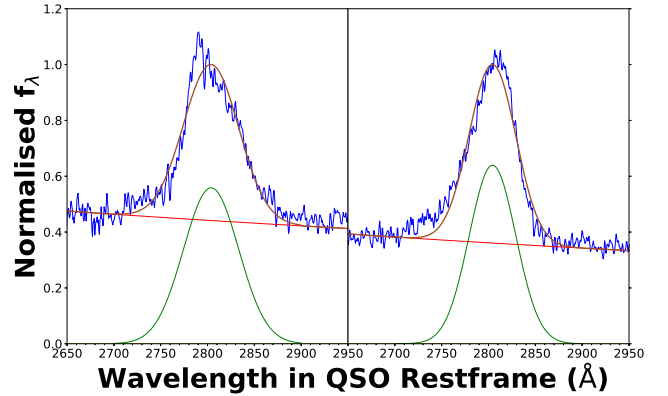


Figure 17. Examples of asymmetry in the Mg II line: *left* J225240 observed at MJD 55500 (SDSS), *right* J111348 observed at 57539 MJD (WHT). The fluxes have been normalised to the maximum flux of the fitted functions. In both objects the line appears skewed with respect to the symmetric Gaussian. The skewness of the lines is to different sides: to the blue for J225240 and to the red for J111348. A possible explanation for this effect is an offset between different line components.

hole mass (M_{BH}), and the combination of these two values into the Eddington ratio (λ_{Edd}). The Eddington ratio can be used as a measure of the accretion rate and EVQs appear to have relatively low Eddington rates on average (Rumbaugh et al. 2018, MCL19). The influence of this parameter on the variability of Mg II is therefore worth investigating. We use the catalogue provided by Shen et al. (2011) for our values of L_{bol} , M_{BH} , and λ_{Edd} , which means the quasars included in this part of the analysis are objects that were included in DR7Q. The number of objects included in this part of the analysis is 3,494. The black hole masses are the fiducial estimates, based on the single epoch virial method (Shen et al. 2011).

The combined distributions, for both samples, of $\log L_{\text{bol}}$, $\log \lambda_{\text{Edd}}$, and the redshift are shown in Figure 19. A number of selection effects distinguishes the Supersample from the general quasar population. The Eddington ratio for the strongly variable objects is on average lower than that of the Full Population sample. This is in agreement with the results presented in Figure 6 of MCL19 and the EVQ sample discussed by Rumbaugh et al. (2018). However, the strong selection biases that separate the samples imply that a direct comparison of the two samples should be treated with caution.

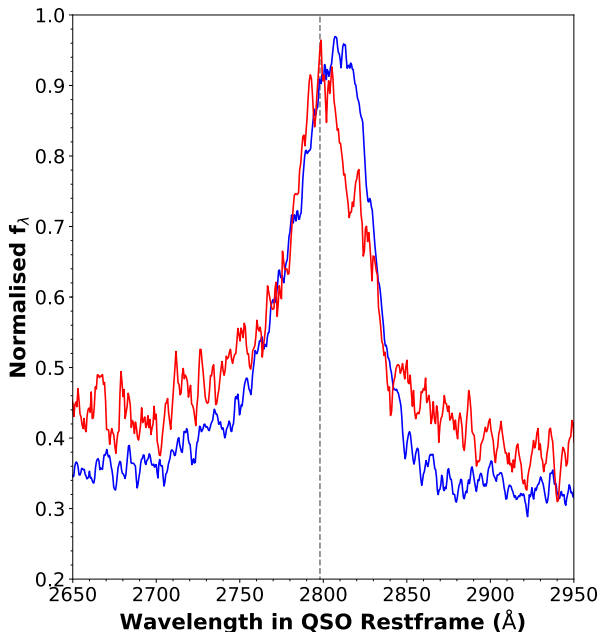


Figure 18. The profile of the Mg II $\lambda 2798$ line in J111348 as observed on MJD 52438 (*red*) and 57539 (*blue*). The fluxes have been normalised to the maximum of the Mg II line flux. Over the period between the two epochs the line flux dropped by a factor ~ 4 . In the more recent spectrum, associated with the lower emission state, the line has acquired a skewness to the red. The vertical line marks 2798Å in the QSO restframe.

7.1 Data Subsets Based on Normalised Flux

By making selections in the normalised flux space, it is possible to create subsets of the samples based on spectral variability. Figure 20 illustrates a comparison of physical parameters associated with different regimes of spectral variability. We consider the λ_{Edd} distribution in more detail in Table 9.

Figure 20 is based on the normalisation to the minimum continuum epoch. The three sections have been selected to represent different types of responsivity. The sections have been made as follows: section *a* targets highly variable and highly responsive objects, section *b* selects a general quiescent sample, and section *c* targets objects where the line is relatively more variable than the continuum. Section *c* overlaps with *b* in the range of continuum fluxes, and with *a* in the range of line fluxes.

The distributions in the right panels of Figure 20 show that the subset *a* is a clearly distinct selection of quasars from those in *b* and *c*. The most variable objects (*a*) are characterised by lower luminosities, lower masses, and lower Eddington ratios. A correlation between continuum variability and Eddington ratio is therefore evident. Considering the Mg II line variability, the λ_{Edd} distributions of the objects in *b* and *c* indicate that these populations are quite similar, suggesting strong Mg II variability can occur in both high and low Eddington systems. However, the averages and means for the distributions (Table 9) do show an offset: the objects with greater line variability have on average somewhat smaller Eddington ratios. We find an average of the $\log_{10}\lambda_{\text{Edd}}$ distribution of -1.02 for section *a* and -0.75 for section *b*, compared to -0.62 for section *c*. The connection of

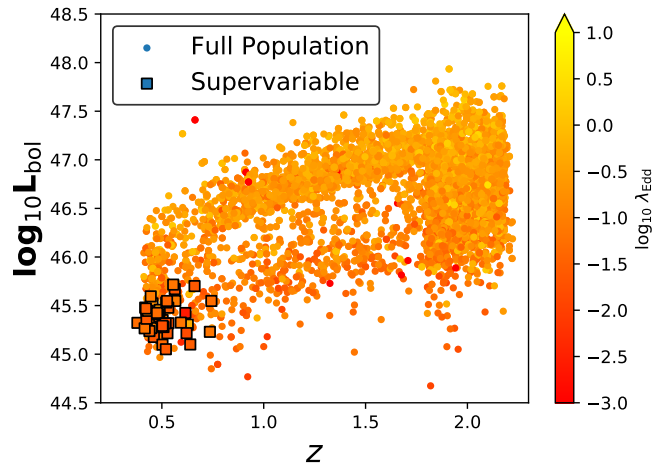


Figure 19. The bolometric luminosity plotted against the redshift for the two QSO samples; colour indicates the Eddington ratio. The Supervariable sample is on average intrinsically dimmer as well as having a lower Eddington ratio. The requirements for the Supervariable sample (table 1) and the SDSS quasar criteria introduce a number of selection effects evident in the distributions of the samples.

Table 9. Summary statistics of the distributions of $\log_{10}\lambda_{\text{Edd}}$ distributions for the data subsets defined in Figure 20. The results for the Full Population and Supervariable sample are included. The sections in the normalised flux space that define the subset were the same for each sample. The columns represent, from left to right, the sample size, median, sample average, and standard deviation.

	Full Population				Supervariable			
	N	\bar{X}	\bar{X}	σ	N	\bar{X}	\bar{X}	σ
a	25	-1.14	-1.02	0.45	24	-1.23	-1.48	0.62
b	1403	-0.61	-0.62	0.33	2	-	-0.90	0.13
c	33	-0.73	-0.75	0.44	0	-	-	-

Mg II flux variability with the Eddington ratio is therefore noticeably weaker than the connection between λ_{Edd} and the continuum flux.

For the Supervariable sample we make the same selections in flux-flux parameter space. We characterise the distributions of $\log_{10}\lambda_{\text{Edd}}$, calculating the median, sample average and standard deviation of the data subsets in Table 9. For the Supervariable sample the comparison shows that the dependence of the magnitude of relative continuum changes on the Eddington ratio also applies. The average value of $\log_{10}\lambda_{\text{Edd}}$ for the highest continuum variability subset *a* is -1.48 and for *b* $\log_{10}\lambda_{\text{Edd}}=-0.90$. The number of objects in the subsets of the Supervariable sample is low, making a reliable comparison difficult. No objects matched the requirements for subset *c*, most likely because the Supervariable sample was selected based on strong continuum variability.

7.2 Considering the Responsivity Measure

We now consider an analysis of sample subsets based on α_{rm} . The first selection simply splits the samples between positive and negative α_{rm} , i.e. the over and under responsive subsets

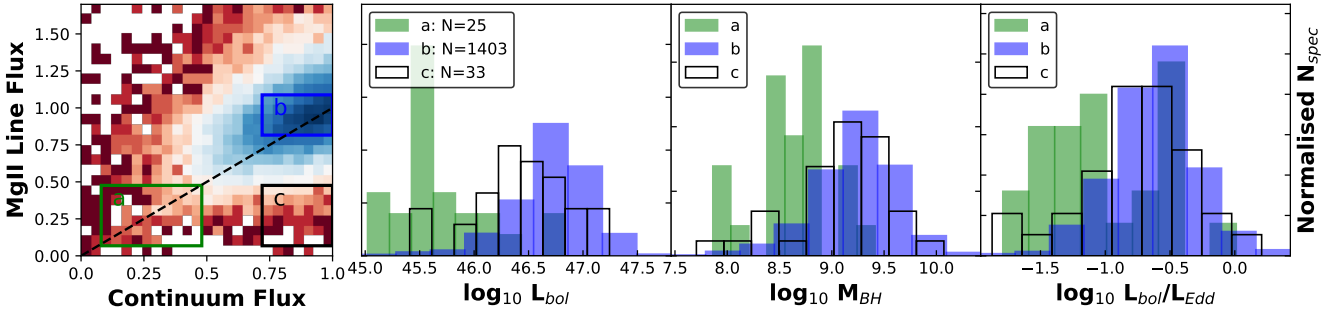


Figure 20. Quasar properties for the subsets of the Full Population sample, normalised to the epoch of highest continuum emission. *left:* The selection in the normalised flux space. A 1:1 ratio of line and continuum flux is indicated by the dashed line. *right:* Three histograms displaying (from left to right) the normalised distribution of the bolometric luminosity, the fiducial black hole mass, and the Eddington ratio. The values for the distributions are from Shen et al. (2011).

respectively. A second selection focuses on the more extreme cases by requiring $|\alpha_{rm}| > 0.75$. This selection will function as a check on the first results: if there is a distinction in physical parameters based on α_{rm} , the distinction should be more prominent among the outliers.

The results for the selection $|\alpha_{rm}| > 0.75$ are shown in Figure 21. For none of the four parameters considered (L_{bol} , M_{BH} , λ_{Edd} , and z) do the distributions of the two sets of parameters differ visibly. For the distributions of λ_{Edd} summary statistics are listed in Table 10. However, the medians and averages in Table 10 do indicate a distinction between the two $|\alpha_{rm}| > 0.75$ samples: the over-responsive objects have a slightly higher λ_{Edd} . For the Supervariable sample a similarly small distinction based on α_{rm} can be detected (Table 10).

A lower value of the Eddington ratio for objects with a negative responsivity measure would be in agreement with the results based on the subsets in flux. The correlation between the continuum flux and λ_{Edd} is stronger than the correlation of the line flux and λ_{Edd} . Therefore objects in which the continuum has relatively changed the most ($\alpha_{rm} < 0$) can be expected to have lower Eddington ratios. The difference between the two distributions is slight, making a firmer conclusion difficult.

8 INTERPRETATION

8.1 Flux Variability

The broad range in Mg II line flux variability noted in both samples agrees with previous studies. Differences from object to object can be seen among RM studies (e.g. Clavel et al. 1991; Hryniewicz et al. 2014; Trevese et al. 2007; Cackett et al. 2015) and are a likely contribution to the strong scatter found in Mg II variability in larger samples of quasars (Kokubo et al. 2014; Sun et al. 2015; Zhu et al. 2017), as well as in the EVQ sample presented by Yang et al. (2020). Even for very large variations in the (ionising) continuum, the response can differ considerably. As was evident from the example of J022556, the behaviour does not only differ among objects, but evolves for individual quasars over time. The range in behaviours illustrated in Figure 3 appears representative of the EVQ population, when compared to other objects in the literature (Ross et al. 2018; Yang et al. 2020; Kynoch et al. 2019; Guo et al. 2019; Dexter et al. 2019):

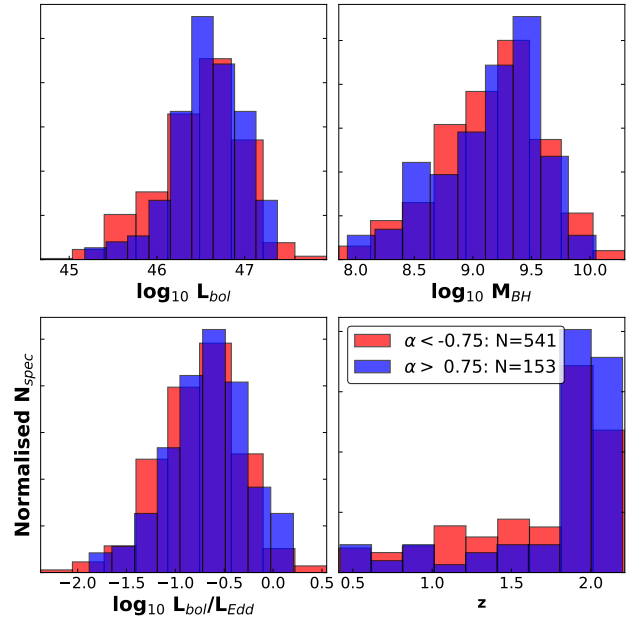


Figure 21. A comparison of the distribution of AGN properties among subsections of the Full Population sample based on the responsivity measure. The properties are: bolometric luminosity (*top left*), black hole mass (*top right*), Eddington ratio (*bottom left*), and redshift (*bottom right*). No clear distinction is visible between the subsets, despite a cut based on a relatively large value of $|\alpha_{rm}|$.

strong changes in the optical or UV continuum can be followed by dramatic changes in the Mg II flux, but this is not always the case .

The primary behaviour observed in the Mg II response is that the line tracks the continuum. This is true for both the Supervariable and the Full Population sample. The correlation between line and continuum has been found in other studies (Sun et al. 2015; Zhu et al. 2017; Yang et al. 2020), although not all authors agree on the strength of the response (Kokubo et al. 2014). Although a more complex flux response (e.g. a saturation effect, Figure 6) is certainly possible, the approximation to first order is the clearest effect in the data. To compare the gradient in our sample with that presented in Yang et al. (2020) we

Table 10. Summary statistics for the distributions of the Eddington ratio in subsets of the Full Population and Supervariable samples based on α_{rm} . The statistics are, from left to right: the sample size, median, sample average, and standard deviation. The same statistics for the samples in their entirety are listed in the bottom row. J081916 was not part of DR7, reducing the total size of the Supervariable sample from 43 to 42.

Sample	Full Population				Supervariable			
	N	Median	\bar{X}	σ	N	Median	\bar{X}	σ
$\alpha_{rm} < 0$	5374	-0.64	-0.65	0.37	35	-1.38	-1.40	0.48
$\alpha_{rm} > 0$	1979	-0.63	-0.65	0.36	7	-1.29	-1.33	0.31
$\alpha_{rm} < 0.75$	541	-0.72	-0.76	0.42	18	-1.34	-1.38	0.28
$\alpha_{rm} > 0.75$	153	-0.65	-0.67	0.41	2	-	-1.16	0.00
Full Sample	7407	-0.63	-0.65	0.36	42	-1.28	-1.36	0.47

calculate $\Delta \log_{10} L_{\text{MgII}} / \Delta \log_{10} L_{2798}$. We find $\Delta \log_{10} L_{\text{MgII}} = (0.58 \pm 0.05) \Delta \log_{10} L_{2798}$ for the Supervariable sample, and 0.35 ± 0.01 for the Full Population sample. Yang et al. (2020) find a value of 0.47 ± 0.05 for $\Delta \log_{10} L_{\text{MgII}} / \Delta \log_{10} L_{3000}$ in an EVQ sample including SDSS quasars. The results are similar and, furthermore, emphasise the stronger overall response observed in the Supervariable sample.

The timescales for strong Mg II variability are \sim years (Ross et al. (2018); Guo et al. (2019), but see Hryniewicz et al. (2014)). For the objects in the Supervariable sample as well as for a large fraction of the objects in the Full Population sample (see Figure 1) the timescales are of the same order of magnitude. On shorter timescales the linear correlation between line and continuum is not as clear, although it can be detected for a fraction of the population (Shen et al. 2016b). Interestingly, a timescale of years corresponds with a dynamical timescale (τ_{dyn}) in the BLR. The fact that a consistent positive correlation (albeit with a large scatter in the amplitude of response) is clearer on longer timescales, could indicate that the ability of the Mg II line to reprocess the ionising flux represents a physical change in the BLR.

A difficulty in linking the evolution of the line response over time to τ_{dyn} is that a larger Δt is on average related to a larger change in continuum. The entanglement of Δf_{MgII} and Δf_{2798} makes a firm conclusion about the effect problematic. The moments of the distributions of Δf_{MgII} and Δf_{2798} in different Δt bins, listed in Table 8, show that the distributions evolve in a similar way. Disentangling the two parameters in this manner therefore does not appear possible. An indicator that Δt is of significant importance is that the gradient of the responsivity is larger in the Full Population sample than in the Supervariable sample (Table 6). It should be stressed that this response is to *fractional changes*: the Full Population sample is dominated by objects with only small changes in both continuum and Mg II. For the most dramatic changes in Mg II flux a large continuum flare or drop might still be required.

An established population effect for quasars is the anti-correlation of continuum variability with the Eddington ratio (MacLeod et al. 2012; Rumbaugh et al. 2018; Graham et al. 2020). Objects with larger Eddington ratios, linked to higher accretion rates, on average show a smaller amplitude in variability. This effect is evident in both samples, as can be seen in Figure 20 and Table 9. The data also show a tentative correlation between the change in Mg II line flux and λ_{Edd} . As the line flux tracks the continuum, a correspondence between the line flux and λ_{Edd} can be expected, based on the correlation between the continuum variability

and this parameter. A similar relation is reported by Dong et al. (2009), who find that the EW of Mg II has an inverse correlation with λ_{Edd} .

Both Sun et al. (2015) and Zhu et al. (2017) find that the response of the Mg II line depends on L_{bol} . Sun et al. (2015) report a decrease in the line variability with increasing L_{bol} . Zhu et al. (2017) find that more luminous objects have a shallower slope in the linear response. The connection between L_{bol} and the Mg II line change is not clear in the data for the Supervariable sample, although the effect is visible for the continuum variability (Figure 20). The decrease of the continuum variability with bolometric luminosity is also found in DRW studies (MacLeod et al. 2010).

The connection between H β and Mg II variability can provide information about the BLR processes that govern the emission. The line forming regions for these two species are thought to be in close proximity (Woo 2008; Shen & Liu 2012; Korista & Goad 2004). The variability of the two lines in relation to each other can show a wide range, as seen in Table 3. A class of objects in which Mg II and H β vary independently would agree with the BOSS objects presented in Roig et al. (2014). The authors of this study reveal the presence of a class of objects in the BOSS sample for which Mg II is significantly stronger than the Balmer emission lines, but which lack signs of reddening. Kokubo et al. (2014) argue that Mg II is intrinsically less variable than H β in the entire quasar population. Although this is in agreement with RM studies, the research presented here indicates that the Mg II behaviour is more complex than a suppressed response compared to the Balmer lines. The observed pattern of behaviour agrees with the CL ‘sequence’ presented in Guo et al. (2020), although we would argue that the marked difference in possible responses among the AGN population suggests that the response is not determined by changes in the luminosity alone. Rather, a range of conditions in the BLR likely conspire to create a specific responsivity (Homan et al. 2020).

Several causes for the difference between H β and Mg II are discussed in the literature. The first possibility is that the material emitting H β has a different geometric distribution than the Mg II emitting material (Sun et al. 2015). The ‘BLR’ that forms H β would then be physically distinct from the one that forms Mg II. For a more extended Mg II region, any change in incident flux would be geometrically diluted, illuminating different radii at different times. This could explain the lack of effect in some of the RM studies, although it is more difficult to explain the strong instantaneous response measured by e.g. Hryniewicz et al.

(2014) in this scenario.

Another issue with an interpretation based on differing BLR geometries is the physical conditions in which the two lines are formed: the Mg II population is likely present in the cooler, self-shielded ‘back’ of the cloud (e.g. Baskin et al. 2014). This region largely overlaps with the parts of the cloud where H is neutral. As the collisional excitation of Mg II depends on the the electron density (n_e) and the temperature (Guo et al. 2020), the formation region of the 2798Å line is likely associated with the transition region between the HII and HI regions of BLR clouds, where H β is also formed (MacAlpine 1972; Netzer 1980; Collin-Souffrin et al. 1986; Krolik 1999). If the line forming regions of H β and Mg II are so intimately linked, a very different geometry for the two appears less probable.

A second possible explanation is based on the difference in optical depths between the two lines (Korista & Goad 2004; Sun et al. 2015; Yang et al. 2020). Mg II is a resonance line, whereas H β is not. Mg II can therefore be expected to have a greater optical depth, increasing the number of absorptions and reemissions of any line photon before it escapes the cloud. This could also dilute the response to a change in incident flux, spreading it out over time. A possible problem with this interpretation is that although Mg II is a resonance line, the population of ground state atoms is itself subject to a collisional equilibrium which could suppress the number of absorbing atoms. Photoionisation modelling does indicate the optical depth effect could be significant (Korista & Goad 2004; Baskin et al. 2014). If optical depth for the lines is a critical factor, the observed evolution with λ_{Edd} could be the result of a changing ionisation structure in the cloud, as a result of a changing shape of the SED. A difference in temperature or n_e could affect the collisional excitation of Mg II in BLR clouds.

8.2 Line Profile

The Supervariable sample shows a range of line profiles. If the skewness seen in some spectra results from an offset between broad and narrow components, a radial outflow or inflow would be a plausible explanation. Another option is to associate the different line components with a blueshifted and a redshifted component. These components could be the result of rotation of the BLR or possibly even the outer regions of the disc, where the line could be formed. Such a scenario would be akin to the interpretation for quasars with double peaked emission lines (e.g. Eracleous et al. 2009). It is impossible to discern the different dynamic components in either of the two samples, however.

To quantify the behaviour of the line profile, we have made use of the line width (σ) of the broadest Gaussian fit to the line. For both the Supervariable and the Full Population sample σ lacks a correlation with the luminosity. We compare our results to those of Yang et al. (2020). We find a slope of $\Delta \log_{10} \text{FWHM}_{\text{MgII}} / \Delta \log_{10} L_{2798}$ of 0.013 ± 0.035 for the Supervariable sample and 0.008 ± 0.066 for the Full Population sample, compared to $\Delta \log_{10} \text{FWHM}_{\text{MgII}} = 0.012 \pm 0.012 \Delta \log_{10} L_{3000}$. These values are in agreement with other results (Woo 2008; Shen et al. 2011) and indicate that the Mg II line does not ‘breathe’.

The lack of breathing presents another marked differ-

ence with the behaviour of H β (Bentz et al. 2013). A possible explanation is that the Mg II emission region is confined to a smaller set of radii (Yang et al. 2020), perhaps as the result of the small geometrical depth of the Mg II line emitting region in each cloud.

9 CONCLUSIONS

The Mg II line exhibits a complex behaviour. The analysis presented in this work falls into two broad categories. The first is flux responsivity: the changing levels of the line flux under a varying continuum provide information about the incident ionising flux and the physical conditions in the line forming region. The second category is the line profile: changes in line width and skewness can provide information about the kinematics of the gas creating the Mg II emission. Below is a brief summary of the results.

9.1 Flux Responsivity

- i) There is a broad range of behaviour in Mg II variability. In some objects the line changes more than the continuum, whereas in others the line shows barely any response at all. Based on the responsivity measure, it is clear that the objects where the line varies less than the continuum form the largest part of the quasar population.
- ii) $|\Delta L_{\text{MgII}}|$ and $|\Delta \lambda L_{2798}|$ show a linear correlation, visible for the Supervariable sample (Figure 4). The relation holds over five orders of magnitude.
- iii) The connection between Mg II variability and CLQ behaviour is not one-to-one (Table 3). Considering H β and Mg II in the same object, the Supervariable sample contains objects where either, both and neither of the lines vary. Joint and separate Mg II and H β variability occur in approximately equal measure.
- iv) The normalised fluxes show a linear correlation between the Mg II line and the 2798Å continuum. This holds for both normalisations, to the minimum and the maximum state (Figures 5 and 6; Table 6). The scatter in the data is considerable. Although the scatter is partially the result of measurement and fitting uncertainties, it is indicative of the range in Mg II variability.
- v) There is clear evidence of a changing responsivity over time in the spectra of J022556 (Figure 10), for which relatively high cadence observations are available. The response of Mg II not only changes from object to object, but also from epoch to epoch.
- vi) A change in the linear responsivity of Mg II with time can be observed in the Full Population sample. The responsivity appears to decrease for longer elapsed times between spectral epochs (Figure 12). When modelled as a simple linear response, the change in gradient is modest, but steady over the Δt bins (Table 7).
- vii) The distributions of Δf_{MgII} and Δf_{2798} change with Δt (Figure 13 and Table 8). The averages of the distributions become more negative over time and the skewness increases. The change in the average is likely the result of a selection bias (see the discussion in section 5.4).
- viii) The Eddington ratio is connected to the magnitude of changes in the continuum. Objects with larger continuum

changes on average have a lower λ_{Edd} . The difference is visible in the offset between the λ_{Edd} distributions of the Supersample and the Full Population sample (Figure 20 and 21 and Tables 9 and 10). For changes in the line flux this connection possibly also exists, but is not as strong as for the continuum.

9.2 Kinematics

- i) Some of the objects in the Supersample show a clear asymmetry in the Mg II profile. This asymmetry manifests itself as a skewness of the lines. The skewness can be both to the blue and to the red of the line centre (Figure 17).
- ii) Asymmetry of the Mg II line can evolve over time, as observed for J111348 (Figure 18). The evolution of the line profile could be due to a change in the relative strength of different line components. These could be the broad and narrow components, but it is also possible they represent a more dynamic structure.
- iii) The Mg II line width lacks a correlation with either the 2798Å continuum or the line flux (Figure 16). There is some tentative evidence of a correspondence for the lowest luminosities, however the bulk of the distribution shows a flat response. The Mg II shows no clear signs on breathing. This is in agreement with previous results (e.g. Shen et al. 2011).

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APPENDIX A: STATISTICAL TESTS

χ^2 Tests

We use χ^2 as the goodness-of-fit parameter when fitting responsivity functions to the normalised flux data. The one-component and two-component responsivity functions are nested functions. In the case of nested functions, the difference between the two χ^2 ($\Delta\chi^2$) of the fits is distributed according to an M d.o.f. χ^2 distribution, where M is the difference in number of parameters between the functions.

In the case of our study $M=1$, therefore $P(\chi_1^2 > \Delta\chi_f^2) = 1 - \text{CDF}(\Delta\chi_f^2)$, CDF is the cumulative distribution function associated with χ_1^2 . If this p -value is small, this indicates that the two-component fit is the preferred model.

Sequence Test

The ‘sequence test’ is intended to check the distribution of the data points around the fitted responsivity functions. If the function truly represents the behaviour of the fluxes, the scatter of the data points should be random, as it would result only from the uncertainties in the measurements. This can be pictured by considering the residuals (see Figure 5). Any data point is just as likely to be above the line of the function as it is to be below it.

The sequence test is performed on the residuals of the fit, and counts the number of data points that have the same sign (\pm) as the preceding data point. The test statistic is labelled S , and represents the number of data points that have the same sign as the preceding point. The sampling distribution for this statistic is the binomial distribution for $N-1$ trials, where N is the sample size. For a perfect fit each data point (after the first one) has a $\frac{1}{2}$ probability of lying on the same side of the response function. S is normalised by dividing by $N-1$. The expectation value for S is therefore $\frac{1}{2}$ and the variance $\frac{1}{4}$. Using the sample distribution a p -value can be calculated for a given S statistic (p_s). This value represents the probability that the S statistic is not drawn from the binomial distribution. A low value of p_s suggests a good fit. We verified the robustness of our test with a bootstrap MC ($N=10^4$) simulation, using the spectra of the Supervariable sample.

Responsivity Measure: α_{rm}

The responsivity measure is constructed using the normalisation to the minimum continuum state. This metric is based on the perpendicular distance in the normalised flux plane from the line that would indicate a 1:1 correspondence between line and continuum flux.

Different types of behaviour can be associated with different values of α_{rm} : the tracking of the line by the continuum will result in $\alpha_{rm} \sim 0$, as will a stable line and continuum; a positive responsivity measure is associated with a larger line flux change than continuum flux change; a negative measure corresponds to objects with a continuum flux change that outdoes the line flux change.

Test for Normal Distribution

There are several statistical tests available to test for normality of a distribution. For the purposes of our analysis we have tested the Lilliefors and Shapiro-Wilk tests on their suitability. The Kolmogorov-Smirnov test was not directly applicable, as the normalisation of the data requires using the sample average and sample variance. The Lilliefors test is an adaptation of the KS test, designed to account for the fact that the normalisation uses properties derived from the sample distribution.

Our testing consisted of a bootstrap MC ($N=10^4$) on the Δf_{MgII} of the Full Population Sample as well as on a test set of data drawn from a normal distribution. The Lilliefors test proved the most robust. We have applied the Lilliefors test as implemented in the `statsmodels` package⁴.

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⁴ <https://www.statsmodels.org/dev/>