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# Few-Mode Fibers With Uniform Differential Mode Group Delay for Microwave Photonic Signal Processing

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**ABSTRACT** We present a novel design of few-mode fiber (FMF) with a uniform differential mode group delay (U-DMGD) among propagating modes and apply the FMF to a single-fiber delay line module for implementing microwave photonic finite impulse response (FIR) filters. By optimizing key parameters such as the core grading exponent and the dimension of the trench of FMF, a U-DMGD between adjacent modes among four modes (LP<sub>01</sub>, LP<sub>11</sub>, LP<sub>02</sub> and LP<sub>31</sub>) over the entire C band is achieved. Wavelength dependence is entirely removed. An FIR microwave photonic filter (MPF) implemented using the designed 1-km FMF is investigated through numerical simulations. The free spectral range (FSR) of the MPF is 5.7 GHz, the 3-dB bandwidth is 1.26 GHz, and the main lobe-to-side lobe ratio (MSR) is 10.42 dB. Discussions on fabrication aspects have also been presented. The proposed single-fiber delay line structure based on FMF can significantly reduce the system complexity of microwave photonic signal processing.

**INDEX TERMS** Few-mode fiber, uniform differential mode group delay (U-DMGD), delay line, microwave photonic filter (MPF).

## I. INTRODUCTION

The microwave photonic filter (MPF), which is a pivotal module in microwave photonics, has been applied in various fields such as radar, phased array beamforming and radio over fiber networks etc [1]. As a dominant branch of MPFs, the tapped transverse MPF uses a set of finitely or infinitely weighted and differentially delayed optical signals to carry the same microwave signal. After photodetection, the summation of these weighted and delayed microwave signal copies leads to a required microwave filter transfer function [2]. A typical finite impulse response (FIR) MPF architecture is shown in Fig. 1. Input RF signals are modulated onto optical carrier through the intensity modulator. The modulated signals are divided into N channels, and the signals in each channel passes through attenuators and differentially delayed fibers to adjust the signal weights and delays. The N-channel optical delay lines are recombined using a standard  $N \times 1$  fused

fiber coupler [2]. Finally, the signals are detected by a photodetector (PD). According to the number of light sources, FIR MPFs can be divided into two categories, single-source MPFs and multi-source MPFs. For single-source MPFs, optical interference effects adversely impact the performance of MPFs, where the differential optical delay is smaller than the coherence time of the source. In order to avoid the optical interference effects, MPF source devices, such as fiber Bragg gratings and multi-wavelength sources, are implemented in the incoherent regime [3]. However, these methods will increase the complexity and cost of the fabrication.

On the other hand, the development of spatial division multiplexing (SDM) technologies including multi-core multiplexing (MCM) [4], mode division multiplexing (MDM) [5] and their successful application in optical fiber communication systems have brought new opportunities for microwave photonic subsystems such as MPFs. In the case of MCM, an MPF has been proposed based on the selective inscription of Bragg grating in homogeneous multicore fibers. The multi-core fiber includes N cores engraved with M gratings,

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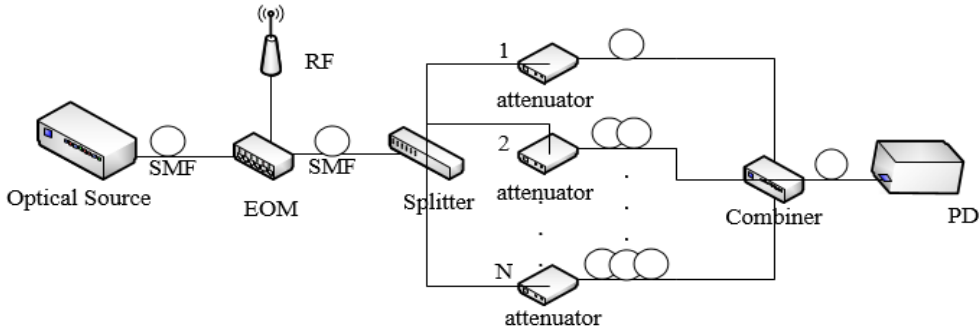


FIGURE 1. The FIR filter architecture with individually weighted and differentially delayed taps.

so this device offers a total of  $N \times M$  delayed signal samples [6]. Meanwhile, an FMF based MPF was also experimentally demonstrated [7]. The MPF utilized DMGDs as optical delay lines, and the amplitudes of the modes excited by the higher-order mode (HOM) excitation method are employed as the tap weights of the filter. In addition, an MPF was realized by determining the wavelength according to the uniform DMGD between adjacent modes and the inscription of the long period grating (LPG) at a specific longitudinal position along the fiber [8], [9]. However, it is crucial to select the specific wavelength or implement LPGs to achieve U-DMGD between adjacent modes. As the number of modes increases, it is more and more difficult and complex to control the operation wavelength.

In this paper, we present a novel optimization scheme of the FMF with U-DMGD over the C band. The proposed FMF supports four modes:  $LP_{01}$ ,  $LP_{11}$ ,  $LP_{02}$  ( $LP_{21}$ ), and  $LP_{31}$  ( $LP_{12}$ ). Compared to delay line elements based on the single-mode fiber (SMF), the delay line elements based on the proposed FMF offers a higher dispersion efficiency, a better power handling capability and less nonlinear effects [7]. Correspondingly, the MPF based on the FMF with U-DMGD is theoretically analyzed, and the spectral response of the MPF is obtained. Compared to the reported designs [8], [9], the MPF can employ the mode dispersion and the chromatic dispersion to provide the uniform time delay difference for signals at different modes and wavelengths, and can thus increase its number of taps. This paper is organized as follows: Section 2 introduces the design strategy of the FMF with U-DMGD; Section 3 introduces the optimization process of the FMF with U-DMGD; Section 4 presents the schematic setup and spectral response of the MPF with U-DMGD FMF; Section 5 draws conclusions of the paper.

## II. DESIGN STRATEGY OF FMF WITH U-DMGD

The definitions of the geometry and parameters in the U-DMGD FMF are described below. The U-DMGD FMF consists of four regions: core (CO), inner cladding (IC), trench (TR) and outer cladding (CL), as shown in Fig. 2. The radii of the core, the inner cladding, the trench and the cladding are represented using  $R_{CO}$ ,  $R_{IC}$ ,  $R_{TR}$ , and  $R_{CL}$ , respectively.

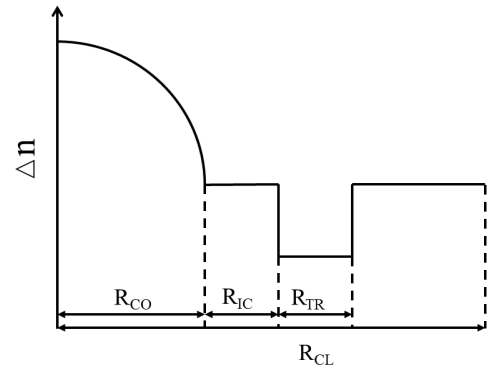


FIGURE 2. The illustrative schematic of the refractive index profile.

The refractive index profile of the U-DMGD FMF, composed by a graded-index core region and a step-index cladding trench region, is described by:

$$n(r) = \begin{cases} n_{CO} [1 - \Delta n_{CO} (\frac{r}{R_{CO}})^\alpha] & 0 < r < R_{CO} \\ n_{CL} & R_{CO} < r < R_{CO} + R_{IC} \\ \frac{n_{CL}}{1 - \Delta n_{TR}} & R_{CO} + R_{IC} < r < R_{CO} + R_{IC} + R_{TR} \\ n_{CL} & R_{CO} + R_{IC} + R_{TR} < r < R_{CL} \end{cases} \quad (1)$$

where  $n(r)$  is the refractive index as a function of the radial distance  $r$ ,  $n_{CL}$  is the cladding refractive index,  $\alpha$  is the core grading exponent, and  $n_{CO}$  and  $\Delta n_{TR}$  are the relative refractive index differences at  $r = 0$  and at the trench, respectively. The relative refractive index difference is defined as  $n(r) = (n(r) - n_{CL})/n(r)$ . The refractive index profile is shown in Fig. 2. Considering the traditional fiber manufacture technology and the influence of the trench on the propagating modes [10], fiber parameters are subject to the following constraints:

$$\begin{cases} 1.5 < \alpha < 2.5 \\ -0.0045 \leq \Delta n_{TR} \leq 0 \\ 0 \leq R_{IC} \leq \frac{1}{2} R_{CO} \\ 0 \leq R_{TR} \leq R_{CO} \end{cases} \quad (2)$$

The range of  $\alpha$  is specified in (2) and  $\alpha = 2$  is typical for FMFs. Since a deeper trench will be difficult to fabricate, (2) constrains  $n_{TR}$  to be greater than  $-0.0045$ . In (2),  $R_{IC}$  is limited to be  $R_{CO}/2$  as a trench further away from the core has much reduced impact on the guided modes. A wider trench will make the fabrication more challenging, so  $R_{TR}$  is generally chosen to be smaller than  $R_{CO}$ . The wavelength range is bounded according to the C-band wavelengths of from 1530 nm to 1565 nm.

The normalized cutoff frequency  $V$  determines the number of modes in a graded index fiber, and can be described by [11]:

$$V = \frac{2\pi R_{CO}}{\lambda} \sqrt{n_{CO}^2 - n_{CL}^2} \sqrt{1 + \frac{2}{\alpha}} \quad (3)$$

In order to support four modes (LP<sub>01</sub>, LP<sub>11</sub>, LP<sub>02</sub> and LP<sub>31</sub>), the normalized cutoff frequency range is  $5.8 \leq V \leq 7.2$ .  $R_{CO}$  should be as large as possible to reduce the nonlinear coefficients for all modes. A core radius of FMF,  $R_{CO} = 19.6 \mu\text{m}$ , is applied. The material of the cladding is set to be SiO<sub>2</sub>. The center of the CO is GeO<sub>2</sub>-doped SiO<sub>2</sub> corresponding to a 4.5% molecular fraction of GeO<sub>2</sub> in SiO<sub>2</sub>. The fiber diameter is designed to be 125  $\mu\text{m}$ , which refers to a typical diameter of the standard single mode fiber.

According to the definition of group delay between adjacent modes,

$$DMGD = \frac{L}{v_{g1}} - \frac{L}{v_{g2}} = \frac{L}{\frac{c}{n_{g1}}} - \frac{L}{\frac{c}{n_{g2}}} = \frac{L}{c} (n_{g1} - n_{g2}) \quad (4)$$

where  $L$  is the transmission distance,  $v_{g1}$  and  $v_{g2}$  are the group velocities of two modes, respectively,  $n_{g1}$  and  $n_{g2}$  are the group refractive indices of corresponding modes, respectively, and  $c$  is the speed of light in vacuum.

From (4) we have,

$$\Delta v_g = \frac{v_{g2} v_{g1}}{c} \Delta n_g \quad (5)$$

where  $v_g = v_{g2} - v_{g1}$  and  $n_g = n_{g1} - n_{g2}$ . Since  $g_1$  and  $g_2$  are very close and both are much larger than  $v_g$ , the product  $v_{g2} \times v_{g1}$  can be almost considered as a constant. Therefore, the group velocity difference  $\Delta v_g$  is proportional to the group refractive index difference  $\Delta n_g$ . To achieve a U-DMGD among adjacent modes, it is necessary to achieve a uniform differential group refractive index  $\Delta n_g$  among adjacent modes.

Simulations of the designed fiber structure are carried out using COMSOL Multiphysics based on the full vector finite element analysis. COMSOL Multiphysics is a general-purpose simulation software for modeling designs, devices, and processes in all fields of engineering, manufacturing, and scientific research. Within the structural parameters, it is found that that the differential effective refractive index between the adjacent modes in different mode groups, such as LP<sub>01</sub>-LP<sub>11</sub>, LP<sub>02</sub>-LP<sub>11</sub> and LP<sub>31</sub>-LP<sub>02</sub>, is  $\sim 10^{-3}$ , while the effective refractive index difference between the

adjacent modes in the same mode group, such as LP<sub>21</sub>-LP<sub>02</sub> and LP<sub>12</sub>-LP<sub>31</sub>, is  $\sim 10^{-5}$ . Subsequent investigations are all performed based on the modes of the different mode groups, unless specified.

### III. PARAMETERS OPTIMIZATION OF FMF WITH U-DMGD

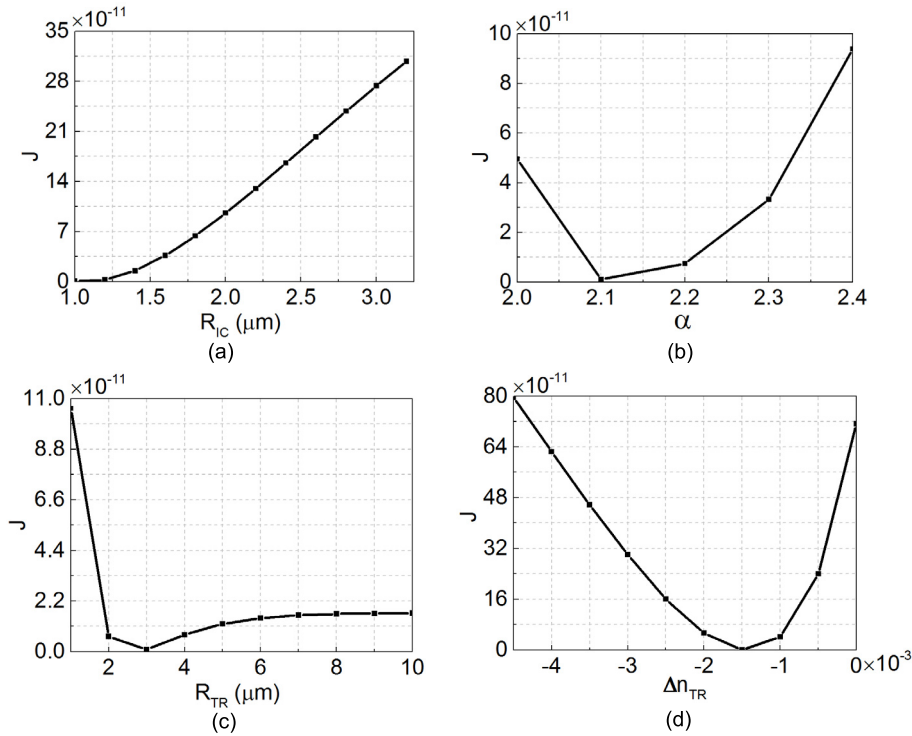
In order to investigate the optimum value of the core grading exponent  $\alpha$  and the dimensions of the trench ( $R_{IC}$ ,  $R_{TR}$ , and  $n_{TR}$ ) that allows U-DMGD among the modes, an optimization function  $J$  is introduced with an objective of minimizing  $\Delta n_g$  at a specific wavelength:

$$J(R_{IC}, \alpha, \Delta n_{TR}, R_{TR}) = (\Delta n_g^{LP02-LP11} - \Delta n_g^{LP11-LP01})^2 + (\Delta n_g^{LP31-LP02} - \Delta n_g^{LP11-LP01})^2 \quad (6)$$

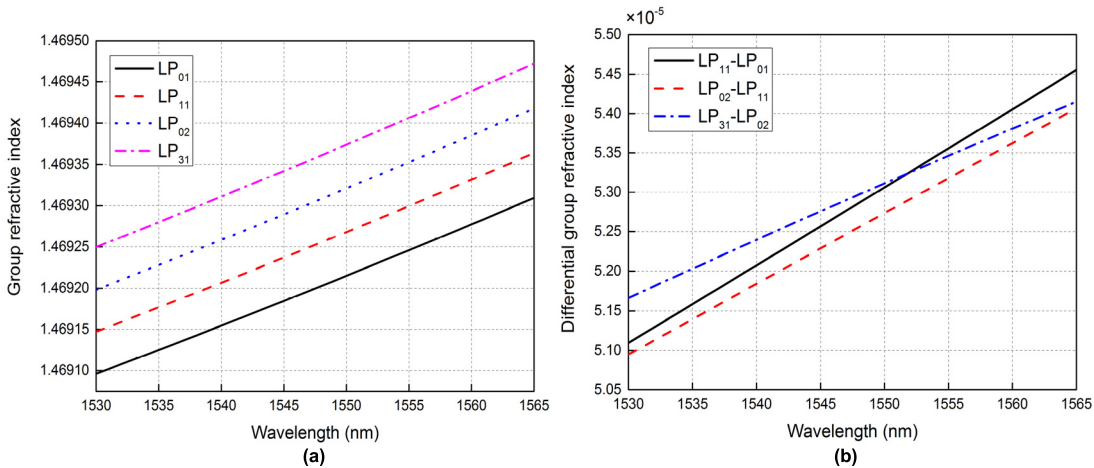
We firstly analyzed the impact of the  $R_{IC}$ ,  $\alpha$ ,  $R_{TR}$ , and  $\Delta n_{TR}$  on  $\Delta n_g$ . A change in the structural parameter  $\alpha$  will result in a change in the refractive index profile of the core, thereby will affect  $n_g$  of all modes. As the LP<sub>01</sub> mode is mainly confined in the core, the other three structural parameters ( $R_{IC}$ ,  $R_{TR}$ , and  $\Delta n_{TR}$ ) have much less influence on the LP<sub>01</sub> mode, and will mainly affect high-order modes.

Based on the explanation presented above, we studied the effect of four parameters ( $R_{IC}$ ,  $\alpha$ ,  $R_{TR}$ , and  $n_{TR}$ ) on  $J$  at a selected wavelength of 1550 nm, in order to realize the design of the FMF with U-DMGD.

In particular, we first fix the parameters with  $\alpha = 2.1$ ,  $R_{TR} = 3 \mu\text{m}$ ,  $\Delta n_{TR} = -1.5 \times 10^{-3}$  and sweep the parameter  $R_{IC}$  (from 1  $\mu\text{m}$  to 3.2  $\mu\text{m}$  with a step of 0.2  $\mu\text{m}$ ). The impact of  $R_{IC}$  on  $J$  at the wavelength of 1550 nm is shown in Fig. 3(a), where the value of  $J$  increases as  $R_{IC}$  increases. The difference of  $\Delta n_g$  between adjacent modes becomes greater as the distance from the core to trench increases, because the influence of the trench on each mode becomes smaller and smaller. We then fix the parameters with  $R_{TR} = 3 \mu\text{m}$ ,  $R_{IC} = 1 \mu\text{m}$ ,  $\Delta n_{TR} = -1.5 \times 10^{-3}$ . The effect of  $\alpha$  (from 2.0 to 2.4 with a spacing of 0.1) on  $J$  is shown in Fig. 3(b). It can be clearly seen that when  $\alpha$  increases, the value of  $J$  first drops and then rises. The curve shows that the value of  $J$  has a minimum value near  $\alpha = 2.1$  for  $\Delta n_{TR} = -1.5 \times 10^{-3}$ ,  $R_{TR} = 3 \mu\text{m}$  and  $R_{IC} = 1 \mu\text{m}$ . Fig. 3(c) shows  $J$  as a function of  $R_{TR}$  (from 1  $\mu\text{m}$  to 10  $\mu\text{m}$  with a step of 1  $\mu\text{m}$ ), at the wavelength of 1550 nm. The value of  $J$  first decreases dramatically and then increases slowly with the increase of  $R_{TR}$  in the scenario of  $R_{IC} = 1 \mu\text{m}$ ,  $\alpha = 2.1$ ,  $\Delta n_{TR} = -1.5 \times 10^{-3}$ . Therefore, it is concluded that when  $R_{TR}$  is around 3  $\mu\text{m}$ , the value of  $J$  reaches its minimum. From Fig. 3(c), we can draw the conclusion that  $R_{TR}$  has marginal influence on the difference of  $\Delta n_g$ , when  $R_{TR}$  exceeds a certain value (around 3  $\mu\text{m}$ ). Fig. 3(d) describes the relationship between  $J$  and  $\Delta n_{TR}$ , in order to explain the relationship between  $\Delta n_g$  and  $\Delta n_{TR}$ . We analyzed  $\Delta n_{TR}$  (from  $-0.0045$  to 0 with a step of 0.0005) and it is found that as  $\Delta n_{TR}$  increases, the  $J$  value decreases first and then increases at the case of  $R_{IC} = 1 \mu\text{m}$ ,  $\alpha = 2.1$ ,



**FIGURE 3.** The optimization function  $J$  as a function of: (a)  $R_{IC}$ , (b)  $\alpha$ , (c)  $R_{TR}$ , and (d)  $\Delta n_{TR}$  at a wavelength of 1550 nm.



**FIGURE 4.** (a)  $n_g$  as a function of wavelength (b)  $\Delta n_g$  as a function of wavelength ranging from 1530 to 1565 nm of the optimized parameters.

$R_{TR} = 3 \mu\text{m}$ . The optimal parameter  $\Delta n_{TR}$  to minimize  $J$  should be around  $-0.0015$ .

Therefore, the value of  $J$  can be minimized with adequate information of trench dimensions ( $R_{IC}$ ,  $R_{TR}$ , and  $\Delta n_{TR}$ ) and at a given wavelength. The structural parameters of the proposed FMF with U-DMGD between adjacent modes at 1550 nm are listed in Table 1.

Fig. 4(a) presents  $n_g$  as a function of the wavelength based on the optimized parameters listed in Table 1 ( $n_{CO}$ ,  $n_{CL}$  and  $\Delta n_{TR}$  follow the Sellmeier Equation [12]). As shown in Fig. 4(a),  $\Delta n_g$  of  $LP_{01}$ ,  $LP_{11}$ ,  $LP_{02}$  and  $LP_{31}$  modes scale linearly with the wavelength ranging from 1530 nm to

**TABLE 1.** The optimized parameters of an FMF to achieve uniform group-delay difference.

$R_{CO}$	19.6 $\mu\text{m}$
$R_{IC}$	1 $\mu\text{m}$
$R_{TR}$	3 $\mu\text{m}$
$R_{CL}$	62.5 $\mu\text{m}$
$n_{CO}$	1.4509 (1550 nm)
$n_{CL}$	1.444 (1550 nm)
$\Delta n_{TR}$	$-1.37 \times 10^{-3}$ (1550 nm)
$\alpha$	2.1

1565 nm. We further analyze  $\Delta n_g$  between adjacent modes. In Fig. 4(b),  $\Delta n_g^{LP_{11}-LP_{01}}$ ,  $\Delta n_g^{LP_{02}-LP_{11}}$  and  $\Delta n_g^{LP_{31}-LP_{02}}$

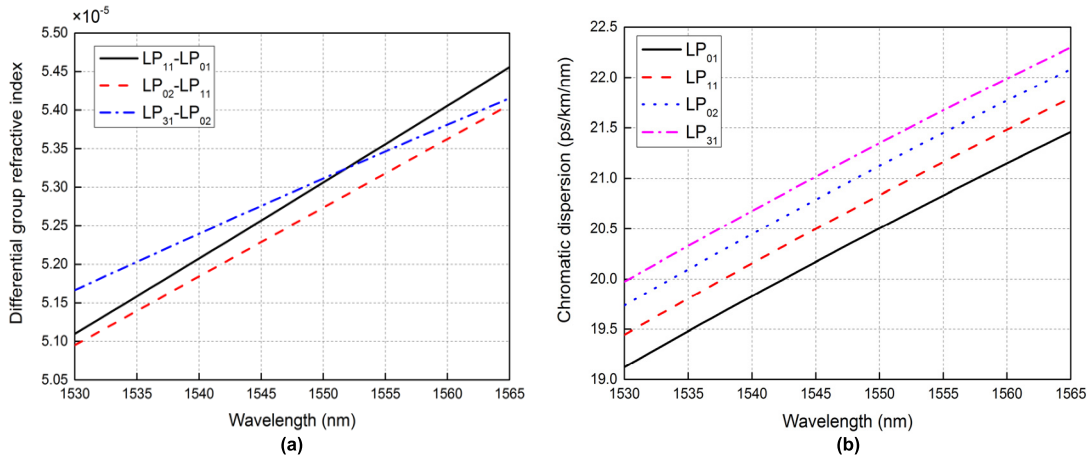


FIGURE 5. (a) DMGD (b)chromatic dispersion as a function of wavelength over the C band for the optimized parameters.

TABLE 2.  $\Delta n_{eff}$  and DMGD between adjacent modes at 1550 nm.

	LP <sub>11</sub> -LP <sub>01</sub>	LP <sub>02</sub> -LP <sub>11</sub>	LP <sub>31</sub> -LP <sub>02</sub>
$\Delta n_{eff} (\times 10^{-3})$	1.18	1.21	1.19
DMGD (ps/m)	0.1769	0.1745	0.1751

increase with the increase of the wavelength, but the curves of  $\Delta n_g^{LP_{11}-LP_{01}}$  and  $\Delta n_g^{LP_{31}-LP_{02}}$  intersect near the wavelength of 1550 nm.

We further analyze the DMGDs per unit length between LP<sub>11</sub> and LP<sub>01</sub>, LP<sub>02</sub> and LP<sub>11</sub>, LP<sub>31</sub> and LP<sub>02</sub>, which are expressed as  $DMGD^{LP_{11}-LP_{01}}$ ,  $DMGD^{LP_{02}-LP_{11}}$  and  $DMGD^{LP_{31}-LP_{02}}$ , respectively. Simulation results are shown in Fig. 5(a). The  $DMGD^{LP_{11}-LP_{01}}$ ,  $DMGD^{LP_{02}-LP_{11}}$  and  $DMGD^{LP_{31}-LP_{02}}$  are 0.1769 ps/m, 0.1745 ps/m, 0.1751 ps/m at 1550 nm, respectively. The chromatic dispersion values of all guided modes over the C band are shown in Fig. 5(b). The  $\Delta n_{eff}$  and the DMGD between adjacent modes at 1550 nm are shown in Table 2.

The mode coupling effect has to be considered when a fiber is designed to provide a significant and exact delay.  $\Delta n_{eff}$  and DMGD between LP<sub>02</sub> and LP<sub>21</sub> are  $1.33 \times 10^{-5}$  ps/m, and 0.0048 ps/m, respectively. The overlap between the LP<sub>02</sub> and the LP<sub>21</sub> modes is 49.34%. The value of  $\Delta n_{eff}$  and DMGD between LP<sub>12</sub> and LP<sub>31</sub> are  $1.63 \times 10^{-5}$ , 0.0085 ps/m, respectively. The overlap between the LP<sub>12</sub> and the LP<sub>31</sub> modes is 39.66%. The mode coupling is related to the overlap integral and the effective refractive index difference [13]. The effect of the mode coupling on the proposed MPF is discussed in Section 4.

The optimized parameters in Table 1 neglect the limited precision margin of the manufacturing control process [14].

Those margins could lead to the deviation of fiber parameters from the optimized values. The typical fabrication tolerances for  $n_{CO}$ ,  $\Delta n_{TR}$ ,  $R_{CO}$ ,  $R_{IC}$ ,  $R_{TR}$  and  $\alpha$  are listed in Table 3 according to [14]. These discrepancies result in corresponding deviations in the DMGD relative to the targeted value. When the deviation of each fiber parameter reaches its maximum, the DMGD of the FMF at 1550nm

TABLE 3. Typical fabrication tolerances of FMF.

$R_{CO}$	$19.6 \pm 0.49 \mu\text{m}$
$R_{IC}$	$1 \pm 0.025 \mu\text{m}$
$R_{TR}$	$3 \pm 0.075 \mu\text{m}$
$R_{CL}$	$62.5 \mu\text{m}$
$n_{CO}$	$1.4509 \pm 0.0001$ (1550 nm)
$n_{TR}$	$-1.37 \times 10^{-3} \pm 0.0001$
$\alpha$	$2.1 \pm 0.01$

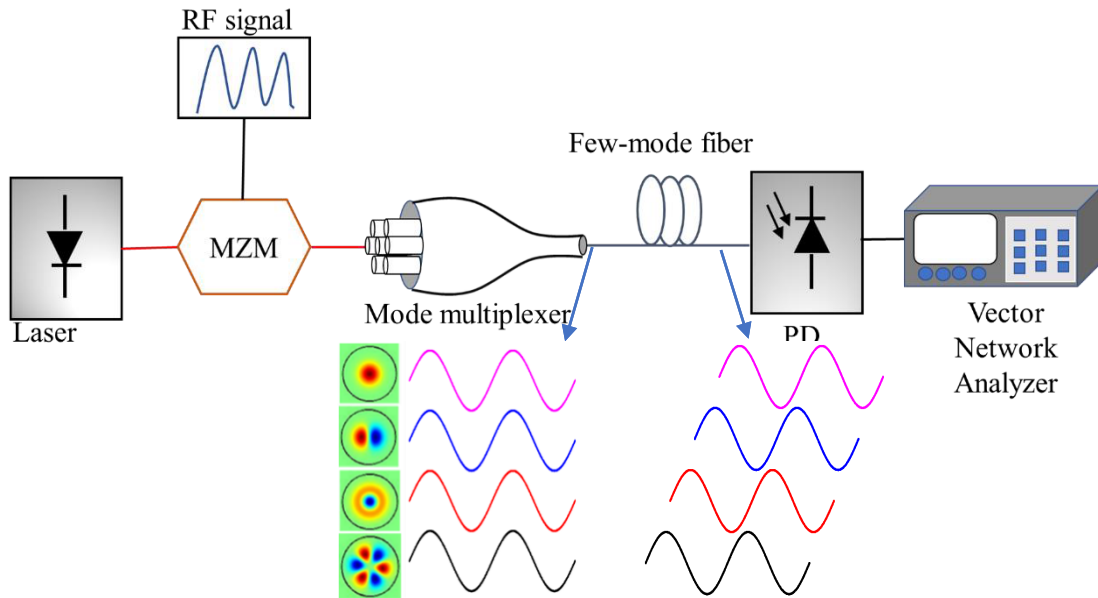
is analyzed. The deviation of DMGD reach the largest value when  $R_{CO} = 20.09 \mu\text{m}$ ,  $n_{CO} = 1.4508$ ,  $\alpha = 2.09$ ,  $R_{IC} = 0.075 \mu\text{m}$ ,  $R_{TR} = 3.075 \mu\text{m}$ ,  $\Delta n_{TR} = -1.38 \times 10^{-3}$  at 1550 nm. The  $DMGD^{LP_{11}-LP_{01}}$ ,  $DMGD^{LP_{02}-LP_{11}}$  and  $DMGD^{LP_{31}-LP_{02}}$  are 0.179 ps/m, 0.182 ps/m, 0.185 ps/m, respectively. The difference of DMGD is 0.003 ps/m.

#### IV. SETUP OF MPF AND PERFORMANCE EVALUATION

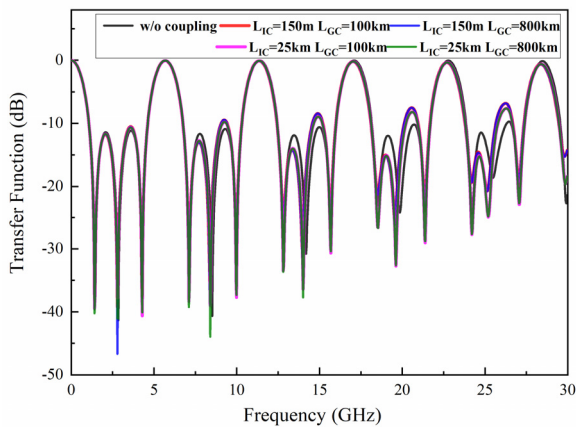
Since we lack of the fiber manufacturing facility, numerical simulations using VPItransmissionMaker software have been implemented to verify the utility of the designed FMF in MPF systems. Realistic fibre parameters have been selected in the simulation and the stated manufacturing tolerances have been considered. Studies in the previous papers have shown that when the bending radius of the FMF is not particularly small, light propagation loss and mode crosstalk can be ignored [15], [16]. The system setup is shown in Fig. 6. The MPF is composed of an RF signal source, a laser, a Mach-Zehnder modulator, a mode multiplexer, a photodetector, a vector network analyzer and 1-km FMF with U-DMGD.

The characteristic parameters of the FMF with U-DMGD are employed into the model. The RF signals are modulated onto the optical carriers through the intensity modulator. All modes are equally excited by the mode multiplexer and are coupled into the FMF with U-DMGD. At the receiver end, the optical signals of all modes are detected by a PD.

The normalized radio-frequency response of the MPF is obtained by the vector network analyzer as shown in Fig 7.

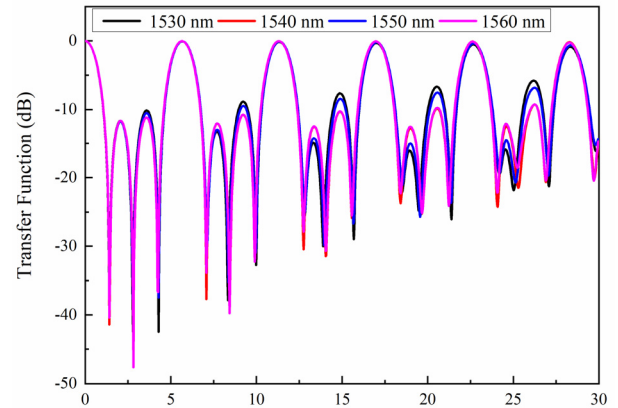


**FIGURE 6.** Setup of the MPF based on the FMF. RF signal: radio-frequency signal; MZM: Mach-Zehnder modulator, PD: photodetector.



**FIGURE 7.** Normalized radio-frequency response of the implemented MPF. Black solid line: without coupling; red solid line:  $L_{IC} = 150$  m,  $L_{GC} = 100$  km; blue solid line:  $L_{IC} = 150$  m,  $L_{GC} = 800$  km; pink solid line:  $L_{IC} = 25$  km,  $L_{GC} = 100$  km; green solid line:  $L_{IC} = 25$  km,  $L_{GC} = 800$  km.

In Fig. 7, the black curve corresponds to the spectral response of the MPF without considering any mode coupling. However, due to the perturbation of the refractive index, the mode coupling along the FMF has to be considered. As mentioned in Section 2, the differential effective refractive index among the mode groups is in the order of  $10^{-3}$ , which can be regarded as a case of weak coupling, while the differential effective refractive index within each mode group is in the order of  $10^{-5}$ , which can be considered as a case of strong coupling [17]. The correlation coupling length is larger than the total system length in the weak coupling regime, while the correlation coupling length is much smaller than the total system length in the strong coupling regime [18]. We studied the impacts of the strong coupling and the weak coupling on the transfer function of the MPF. For comparison, the correlation coupling lengths within the mode group ( $L_{IC}$ ) are selected as 150 m and 25 km,

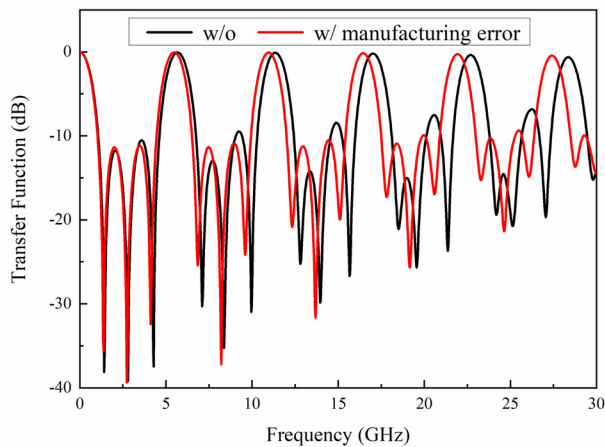


**FIGURE 8.** Normalized radio-frequency response of the MPF implemented at 1530 nm, 1540 nm, 1550 nm, and 1560 nm.

respectively [19]. The correlation coupling lengths between mode groups ( $L_{GC}$ ) are selected as 100 km and 800 km, respectively [20], [21]. As shown in Fig. 7, the FSR of the MPF is 5.7 GHz, the 3-dB bandwidth is 1.26 GHz, and the MSR is 10.42 dB. By comparing the five curves, it is found that the FSR and the 3-dB bandwidth of the MPF remain unchanged, while the MSR of the MPF decreases due to the mode coupling.

The normalized radio-frequency response of the MPF over the C band is also analyzed. The normalized radio-frequency response functions of the MPF at 1530 nm, 1540 nm, 1550 nm, and 1560 nm are shown in Fig. 8, respectively. The  $L_{GC}$  is selected as 100 km and the  $L_{IC}$  is selected as 150 m. As shown in Fig. 8, the transfer function walks off slightly because the DMGD and the dispersion of the FMF differ slightly at different wavelengths.

In fact, the finite precision margins of the manufacturing processes cannot be neglected. Hence, the transfer function will vary with the manufacturing error [14]. The transfer



**FIGURE 9.** The normalized radio-frequency response of the MPF implemented. Black solid line:  $R_{CO} = 19.6 \mu\text{m}$ ,  $n_{CO} = 1.4509$ ,  $\alpha = 2.1$ ,  $R_{IC} = 1 \mu\text{m}$ ,  $R_{TR} = 3 \mu\text{m}$ ,  $\Delta n_{TR} = -1.37 \times 10^{-3}$ ; red solid line:  $R_{CO} = 20.09 \mu\text{m}$ ,  $n_{CO} = 1.4508$ ,  $\alpha = 2.09$ ,  $R_{IC} = 0.075 \mu\text{m}$ ,  $R_{TR} = 3.075 \mu\text{m}$ ,  $\Delta n_{TR} = -1.38 \times 10^{-3}$ .

functions of the MPF without any manufacturing error and with a maximum manufacturing error have been simulated in Fig. 9. The parameters used are listed in Table 3. Due to the influence of the process tolerance, the DMGDs between modes will change, and the FSR and MSR of the MPF will decrease. The FSR of the MPF is reduced by 0.03 GHz.

## V. CONCLUSIONS

In this paper, an FMF with U-DMGD among propagated modes is designed. Through the optimization of structural parameters, the stable excitation and transmission of four modes ( $LP_{01}$ ,  $LP_{11}$ ,  $LP_{02}$  and  $LP_{31}$ ) with U-DMGD among them are achieved over the entire C band. The designed FMF features wavelength independent U-DMGD. As a single-fiber multitap delay line structure, the FMF holds great potential in effectively reducing the complexity of the MPF system in microwave photonic signal processing. Specifically, the deviation of fiber geometric dimensions and the doping concentrations are tolerable and reasonable within the existing fiber manufacture facility restrictions. Thus, our design is highly promising for practical fabrications. The FIR MPFs based on the U-DMGD FMF can employ the mode coupler (such as the photonic lantern) to independently adjust and control the power of each excited mode, and can thus realize the tunability of the filter response. Although the MPF based on the FMF with U-DMGD can only realize a low-pass filter response, equivalent negative or complex tap coefficients can be obtained by deliberately designing non-uniform time delays among filter taps using our proposed method, in order to realize a reconfigurable frequency response.

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