# Parallel Coupled Microstrip Filters With Ground-Plane Aperture for Spurious Band Suppression and Enhanced Coupling

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Abstract—Parallel coupled microstrip sections with a slotted ground plane are proposed as building blocks of coupled-line microstrip filters with enhanced performance. It is shown that, by proper adjustment of the ground-plane slot dimensions, the double frequency spurious band associated with unequal even/odd electrical lengths can be suppressed or meaningfully reduced. As an additional feature, this simple design relaxes tolerances of strip width and spacing in those cases where tightly coupled high-impedance sections are required. A rough preliminary design can be obtained within a few seconds using a fast optimization algorithm based on a quasi-TEM analysis of the coupled sections. Fine tuning is based on the use of a commercial electromagnetic simulator. Finally, experimental check of filter performance is provided.

*Index Terms*—Coupled-line microstrip filters, slotted ground plane, spurious response suppression.

## I. INTRODUCTION

ARALLEL coupled-line microstrip filters (PCMFs) are usual components in microwave integrated circuits owing to their simple design and versatility. However, when implemented in their traditional form on a single-layer substrate in microstrip technology, the filter developed by Cohn in 1958 [1] inherently presents two relatively serious drawbacks. The first one is the existence of a spurious passband at  $2f_0$ (where  $f_0$  is the central frequency of the filter). This undesired passband is associated to the nonhomogeneous nature of the dielectric surrounding the conductors (this causes even and odd modes supported by the coupled line to have different phase velocities). The second problem comes from the difficulty of fabricating wide-band filters because of the weak coupling between the lines in the conventional structure [see Fig. 1(a)]. Strong coupling required by some filter specifications leads to very small values of strip width w and strips spacing s, which cannot sometimes be accurately achieved in practice. Due to the practical importance of the aforementioned problems,

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Fig. 1. Cross sections of basic building blocks of microstrip coupled lines for filter design. (a) Conventional. (b) Modified with ground-plane aperture.

a number of methods have been proposed in the literature to improve the out-of-band performance of PCMFs (apart from the obvious substitution of PCMFs by other types of microstrip filters not exhibiting those intrinsic drawbacks [2], although maybe showing other limitations). Roughly speaking, most of the methods available in the literature to suppress the undesired passband of PCMFs fall into one of the following two categories: those based on lumped load compensation [3]. [4] and those based on modified structures with intrinsically equal modal phase velocities [5], [6]. Recent proposals, which can be considered sophisticated implementations of a mixing of those approaches, are the use of continuously varying (nonuniform) coupled transmission lines [7] or corrugated coupled microstrip lines [8]. Completely different approaches are also available. Thus, extra transmission zeros associated to extra coupling capacitances are used in [9] and [10], the utilization of a two-dimensional (2-D) photonic-bandgap ground plane is proposed in [11], and a periodic (sinusoidal) modulation of the coupled lines width is reported in [12].

Our proposal in this paper is to use the modified structure in Fig. 1(b), which simply incorporates in the well-known microstrip line a centered slot at the ground plane. The slot width  $s_R$  can be adjusted so as to tune the even/odd phase velocities in order to suppress the spurious band. As an additional advantage, this structure provides tight coupling (in comparison with conventional microstrip lines), thus, relaxing the requirements on physical dimensions w and s in those cases where tight coupling is necessary [13]. This structure was analyzed by Masot *et al.* in [14] as a suitable configuration for high-directivity and/or

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high-coupling printed directional-coupler applications. In [14], a fast algorithm for perfect phase velocity matching (based on the judicious choice of  $s_R$ ) was proposed in order to improve the directivity of the coupler. This algorithm is also used in our case as a first step in the design process. This step is carried out very quickly since the computer code providing the modal parameters of the basic coupled section was highly optimized [15] (the code is based on an accelerated quasi-static spectral-domain analysis). This is an additional advantage of our proposal when compared with other approaches requiring electromagnetic (EM) simulation since the beginning of the design process. Nevertheless, EM simulations are also necessary in our scheme for final fine tuning. The reason is that the end effects are important and quite different for even and odd modes when the coupled sections in Fig. 1(b) are used as building blocks of a PCMF. Note that overall electrical lengths, rather than modal phase velocities, must be matched for spurious band suppression and that end effects meaningfully contribute to this electrical length. Two different strategies have been employed in this paper to achieve electrical-length equalization. The first one (illustrated by cases A and B in Section II) consists of selecting  $s_R$  so as to compensate differential end effects, i.e.,  $s_R$ is chosen not to achieve identical mode phase velocities, but to equalize overall modal electrical lengths (including end effects). The second strategy (case C) consists in retaining the value of  $s_R$ that provides phase velocity matching along most of the extension of the coupled section, while local perturbations (widening of the slots) at the ends of the coupled sections are introduced to compensate for the difference between equivalent end effect modal electrical lengths. Experimental confirmation of the proposed methods is finally provided.

# **II. FILTER DESIGN**

## A. One-Pole Butterworth Filter of $\Delta = 20\%$

Let us consider the simple case of a Butterworth bandpass filter centered at  $f_0 = 2$  GHz having a bandwidth of  $\Delta = 20\%$ and order N = 1 in order to illustrate the use of the block in Fig. 1(b). Following [1], we obtain for the two identical sections of the filter,  $Z_{\text{even}} = 77.67 \ \Omega$  and  $Z_{\text{odd}} = 38.04 \ \Omega$ . The corresponding physical dimensions of the conventional and slotted ground filters are shown in Table I, as well as the even (e) and odd (o) modal effective permittivities  $\epsilon^e_{ef}$  and  $\epsilon^o_{ef}.$  This design has been carried out in seconds on the basis of the quasi-static approach in [15] and the use of an ad hoc optimization scheme [14]. The optimization strategy takes into account that mode impedances are mainly controlled by strip width w and strips separation s, while phase velocity mismatching is mainly affected by the ground slot width  $s_R$ . Thus, instead of using a general-purpose gradient optimization method, we only change one of the structural variables, each time attending only to the associated error function. If this scheme is repeated a few times, a very good solution is attained. Many more evaluations of the coupled-line parameters are required if the general-purpose optimization algorithm is used. From Table I, it can be observed that, as was predicted, both w and s are larger for the modified geometry than for the conventional one. This feature relax tolerances in the definition of linewidth and line separation when an accurate fabrication technique is not available.

TABLE I DIMENSIONS (IN MILLIMETERS) AND EFFECTIVE PERMITTIVITIES OF CONVENTIONAL AND MODIFIED COUPLED-LINE BANDPASS FILTER A. SUBSTRATE THICKNESS: 0.635 mm; PERMITTIVITY: 10  $\epsilon_0$ 

Top view	Bottom view
CONVENTIONAL	MODIFIED
w = 0.413 s = 0.173 l = 14.69 $\epsilon^{e}_{ef} = 6.96, \ \epsilon^{o}_{ef} = 5.63$	$ \begin{array}{c} w = 0.732 \\ s = 0.361 , s_R = 1.78 \\ l = 15.93 \\ \epsilon^e_{ef} = \epsilon^o_{ef} = 5.54 \end{array} $



Fig. 2. Full-wave simulations obtained with *Ensemble* for the conventional filter of Table I (dashed line), the modified version with equal-mode phase velocities, i.e., with  $s_R = 1.78$  mm (gray line), and the modified version with  $s_R = 1.90$  mm for modal electrical lengths compensation (solid line).



Fig. 3. Simulated and measured results for the final version of the filter design A.

Note that no edge effects were considered in this preliminary design. The physical lengths of the coupled sections were calculated following [16] (conventional filter) or using  $l = \lambda_e/4 = \lambda_o/4$  (modified filter),  $\lambda_e$  and  $\lambda_o$  being the even/odd wavelengths at  $f_0$ . These structures were simulated using the commercial package *Ensemble* and the obtained responses are shown in Fig. 2. As expected, the conventional filter presents spurious transmission at approximately 4 GHz. Since the central frequency matches pretty well the preliminary design prediction, we conclude that end effects are small in this case. However, owing to strong distortion of the surface current at the ground plane, these effects become important for the

TABLE II DIMENSIONS (IN MILLIMETERS) OF CONVENTIONAL AND MODIFIED BANDPASS PCMF B. THE SAME SUBSTRATE AS IN A



Type of design	Sections 1, 3	Section 2
Conventional	w = 0.184, s = 0.0828, l = 14.98	w =0.238, s =0.0918, l =14.90
	$\epsilon^{e}_{ef}$ =6.52, $\epsilon^{o}_{ef}$ =5.53	$\epsilon^e_{ef}$ =6.63, $\epsilon^o_{ef}$ =5.54
Modified with	$w = 0.309, s = 0.145, s_R = 1.29$	$w = 0.404, s = 0.165, s_R = 1.65,$
$\epsilon^{e}_{ef} = \epsilon^{o}_{ef}$	$l = 15.96, \ \epsilon^{e}_{ef} = \epsilon^{o}_{ef} = 5.52$	$l = 15.94, \ \epsilon^{e}_{ef} = \epsilon^{o}_{ef} = 5.52$
Modified with	$w = 0.412, s = 0.194, s_R = 1.80,$	$w = 0.404, s = 0.165, s_R = 1.65,$
spurious band suppression	$l$ =15.85, $\epsilon^{e}_{ef}$ =5.45, $\epsilon^{o}_{ef}$ =5.00	$l = 15.85, \ \epsilon^{e}_{ef} = \epsilon^{o}_{ef} = 5.52$

modified version of the filter, leading to a severe displacement of the central frequency. Moreover, this effect is different for each mode, in such a way that the equivalent excess lengths  $\Delta l_e$  (even mode) and  $\Delta l_o$  (odd mode) are not equal at all. As a consequence, the spurious band is not completely removed. This problem can be corrected by slightly increasing the slot width, as is also shown in Fig. 2 (solid line). The dimensions of w and s must be simultaneously adjusted so as to keep the same modal impedances, but this is not a problem thanks to the use of the fast quasi-static solver described in [15]. When  $s_R = 1.90$  mm, the small difference between  $\epsilon_{ef}^e$  and  $\epsilon_{ef}^o$  (note that the even mode is more affected than the odd mode when the slot width is augmented) compensates for the difference between  $\Delta l_e$  and  $\Delta l_o$ . The spurious band then disappears.

As a final step, the length of the sections are slightly reduced so as to match the desired central frequency (l = 13.83 mm for this filter). The full-wave simulation of the final design is shown in Fig. 3. Measured results are also included in that figure. Discrepancies probably come from lack of accuracy in the fabrication process, but it is clear that the spurious band is meaningfully reduced with respect to the one expected from a conventional design. Interestingly, it should be mentioned that the final dimension of  $s_R$  is not critical, which is in contrast with other approaches [6].

# *B.* Two-Pole Chebyshev Filter of $\Delta = 30\%$ and Ripple 0.1 dB

As a second example to illustrate additional details of the design procedure, we have designed a  $\Delta = 30\%$  bandwidth Chebyshev bandpass filter with  $f_0 = 2$  GHz, order N = 2, and ripple 0.1 dB. Classical design requires  $Z_{\rm even} = 115.3 \ \Omega$  and  $Z_{\rm odd} = 40.57 \ \Omega$  (Sections 1 and 3), and  $Z_{\rm even} = 103.7 \ \Omega$  and  $Z_{\rm odd} = 38.64 \ \Omega$  (Section 2). Dimensions of conventional, modified with identical modal velocities, and modified for end-effect compensation are shown in Table II. In this filter design, a detail deserves to be clarified. Note that, in this case, there are two different coupled sections. To obtain the final dimensions, we have used *Ensemble* to suppress the spurious band in the response of the one-pole filter that could be constructed based on *each section separately*. It can be



Fig. 4. *Ensemble* simulations for the conventional filter of Table II (dashed line) and the modified version (solid line) when  $\epsilon_{ef}^e = \epsilon_{ef}^o$ .



Fig. 5. Comparison between simulated and measured results for the final version of the designed filter B.

observed that no relevant changes are required for the central section of the filter. We would like to emphasize again that the small values of *s* required by the conventional design are not required by the filter based on the slotted structure. This permits to push further on the limitations on specifications associated with the line precision achievable with the available technology.

An important practical issue is that the slot under each section must be separated from the adjacent one (see the figure atop Table II) to avoid slot mode propagation. If the printed metal bridges separating the slots are eliminated, the filter response

TABLE III DIMENSIONS (IN MILLIMETERS) OF CONVENTIONAL AND MODIFIED COUPLED-LINE BANDPASS FILTER C

Type of design	Sections 1, 3	Section 2
Conventional	w = 0.357, s = 0.134,	w = 0.472, s = 0.240,
	l = 14.74	l = 14.63
	$\epsilon^{e}_{ef}$ =6.86, $\epsilon^{o}_{ef}$ =5.60	$\epsilon^e_{ef}$ =7.07, $\epsilon^o_{ef}$ =5.69
Modified	w = 0.620, s = 0.263,	w = 0.865, s = 0.551,
	$s_R=1.59, l=15.93,$	$s_R$ =2.09, $l$ =15.96,
	$\epsilon_{ef}^e = \epsilon_{ef}^o = 5.54$	$\epsilon^{e}_{ef} = \epsilon^{o}_{ef} = 5.52$



Fig. 6. Full-wave simulations obtained from *Ensemble* for the conventional filter of Table III (dashed line) and the modified version (solid line) when  $\epsilon^e_{ef} = \epsilon^e_{ef}$ .

deteriorates drastically (this has been both simulated and measured, although results are not included for the sake of space saving). This phenomenon is closely related to the excitation of the slot mode in the coplanar-waveguide version of parallel-coupled line filters. Air bridges are required to suppress the undesired mode in such a case [17] because of the coplanar nature of the structure, but thanks to our double sided geometry, metal bridges are just a part of the ground-plane pattern. Fig. 4 depicts the simulated responses for the conventional and modified filters (assuming  $\epsilon_{ef}^e = \epsilon_{ef}^o$ ). Important spurious band reduction is clearly observed. In Fig. 5, we compare *Ensemble* results for the final design (row 3 in Table II) and the measurements of the fabricated filter. Although the spurious band is not completely removed (due to imprecision in the filter dimensions associated to our fabrication process), its level is below -25 dB.

# C. Two-Pole Chebyshev Filter of $\Delta = 20\%$ and Ripple 0.5 dB

Here, as an example of the second strategy mentioned in Section I, we have designed a Chebyshev bandpass filter of parameters  $f_0 = 2$  GHz,  $\Delta = 20\%$ , order N = 2, and ripple 0.5 dB. In this case, the classical design provides the values  $Z_{\text{even}} = 84.86 \ \Omega$  and  $Z_{\text{odd}} = 37.54 \ \Omega$  for Sections 1 and 3 and  $Z_{\text{even}} = 70.75 \ \Omega$  and  $Z_{\text{odd}} = 39.20 \ \Omega$  for Section II. Dimensions for a conventional filter and modified filter implemented in the substrate of cases A and B are shown in Table III (no edge effects are accounted for). Fig. 6 shows the simulated filter responses for filters in Table III. As expected, the spurious band of the modified filter is significantly reduced and the central frequency is slightly shifted. To completely remove the spurious band, we have widen the slots only in the neighborhood of the transitions between filter sections, such as shown in Fig. 7 (dimensions of perturbed slots are noted



Fig. 7. Local perturbation of the slots in the ground plane to completely remove the spurious band (case filter C). The sections lengths to achieve the desired central frequency are  $l_1 = 14.83$  mm and  $l_2 = 14.68$  mm. The slots are 1 mm wider along a distance of 2.15 mm.



Fig. 8. Comparison between simulated and measured results of the final version of filter C.

in the caption). Note that, except in this zone, the slot has the dimensions that make the even and odd modal phase velocities equal.

The results of the simulation with *Ensemble* for the final version of the filter with locally perturbed slots are compared with the measurements in Fig. 8. Once again, it can be appreciated that the spurious band practically disappears and that the agreement between experimental and simulated results can be considered acceptable. Note that local perturbation of the slot width is enough for modal electrical lengths compensation because the main cause of mismatching (different even/odd-mode phase velocities) has been already removed by proper adjustment of  $s_R$ . Local slot widening is just a fine adjustment that is not critical. Just a few full-wave simulations are enough to properly tune the dimensions of the local perturbation.

#### **III.** CONCLUSIONS

A simple method to remove the spurious band of conventional PCMFs has been presented. The width of a slot practiced in the ground plane along each coupled section has been adjusted to compensate for unequal modal electrical lengths. The preliminary design neglecting end effects has been based on a fast quasi-TEM approach. These effects are then accounted for by using a full-wave simulator for planar circuits and compensated either by fine tuning of the ground-plane slot widths or by local perturbation of such slots. Experimental verification of the theory is provided. A satisfactory agreement between simulated and measured responses has been found.

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