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# The importance of Rydberg orbitals in dissociative ionization of small hydrocarbon molecules in intense few-cycle laser pulses 

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Much of our intuition about strong-field processes is built upon studies of diatomic molecules, which have electronic states that are typically well separated in energy. In polyatomic molecules, however, the electronic states are closer together, leading to the involvement of more states. We studied a specific process initiated by intense few-cycle laser pulses in $\mathrm{C}_{2} \mathrm{D}_{2}$, $\mathrm{C}_{2} \mathrm{D}_{4}$, and $\mathrm{C}_{2} \mathrm{D}_{6}$, namely single ionization of the parent molecule followed by hydrogen elimination. We measure the momentum of the remaining $\mathrm{C}_{2} \mathrm{D}_{n-1}^{+}$fragment in each case using velocity map imaging (VMI). Our measurements, along with angle-dependent ionization probability calculations [1], demonstrate that Rydberg orbitals are readily populated for certain orientations of the molecules relative to the laser polarization. Moreover, these commonly-neglected Rydberg orbitals make a key contribution to strong-field ionization.

As an example, the comparison between measured and calculated $\mathrm{C}_{2} \mathrm{D}_{3}^{+}$angular distributions from $\mathrm{C}_{2} \mathrm{D}_{4}$, shown in Fig. 1, clearly illustrates the need to include contributions from Rydberg states. The general four-lobed structure of $\mathrm{C}_{2} \mathrm{D}_{3}^{+}$ from dissociative single ionization is independent of pulse intensity and duration (at least up to 45 fs [2]). Calculations that do not include ionization from Rydberg orbitals result in an approximately isotropic angular distribution, as illustrated by Fig. 1(c), while including contributions from Rydberg orbitals yields the four-lobed structure shown in Fig. 1(d), which qualitatively matches the experimental results. This improved understanding of the links between photofragment images and the electronic dynamics in the molecule can improve endeavors towards image-
based strong-field coherent control.


Figure 1. (a) Momentum distribution of $\mathrm{C}_{2} \mathrm{D}_{3}^{+}$. The inner lobes come from dissociative single ionization of $\mathrm{C}_{2} \mathrm{D}_{4}$, and the weaker outer ring is due to dissociative double ionization. (b) Measured yield of the inner lobes as a function of the angle between the $\mathrm{C}_{2} \mathrm{D}_{3}^{+}$and the laser polarization. Calculated angular distribution for $\mathrm{C}_{2} \mathrm{D}_{3}^{+}$(c) without including Rydberg orbitals and (d) with ionization from Rydberg orbitals included.

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