

THEORETICAL REACTOR KINETIC MODELS AND
EXPERIMENTAL VERIFICATION

by

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B. S., Kansas State University, 1959

A THESIS

submitted in partial fulfillment of the

requirements for the degree

MASTER OF SCIENCE

Department of Nuclear Engineering

KANSAS STATE UNIVERSITY
OF AGRICULTURE AND APPLIED SCIENCE

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INTRODUCTION

Nature of the Study

The control of a neutron chain reaction is probably the most important aspect of atomic energy. The ability to control a nuclear reactor permits the usefulness of such a device. As is well known, the controllability of a neutron chain reaction is primarily due to the existence of delayed neutrons, i.e., neutrons which are born a considerable length of time (relative to neutron lifetimes) after the fission process. The time behavior of a nuclear reactor must be well understood in order that the very concept of nuclear energy not become a menace to mankind. The study of reactor kinetics has progressed from early in the Manhattan Project and many attempts have been made to understand and solve reactor kinetics problems. In many instances, in order to solve the kinetics problems, each reactor must be treated individually and even then the problems may be extremely difficult or impossible.

The usefulness of any theory lies in its ability to predict mathematically the characteristic behavior with which it is concerned. Therefore, kinetics theory must be able to describe the time behavior of a neutron chain reaction. This paper investigates the characteristics, limitations, and utility of some of the presently used theories. The theory utilized in this paper is very similar to that discussed in treatises such as Glasstone and Edlund (8). The theory was applied to reactor systems for the prediction of the neutron density time dependence where both

step and finite changes of reactivity were assumed. This paper will demonstrate the agreement found between theoretical and experimental determination of the reactivity worth of reactor control rods, measured both with the positive period method and with the rod drop method.

The particular reactor of interest here is a TRIGA reactor (see PLATE I) (designed and built by the General Atomic Division of General Dynamics Corp.). Since Kansas State University will install a TRIGA Mark II reactor in 1961, it is important to know beforehand the nature of the kinetic analysis that will be required to predict the behavior of such a reactor. Because of the extreme variation in complexity of the equations for the various kinetic models, it is important to realize the simplest model which will suitably describe the time behavior of the reactor.

Since almost all solutions of the various kinetic equations must be solved with the aid of a computer, either digital or analog, it is important to realize an approximate solution which will accurately predict the kinetic behavior of a reactor. Therefore, solutions of the kinetic equations using both three and six delayed neutron groups were developed and solved with the aid of the Kansas State University IBM-650 computer.

Since the TRIGA reactor has a control rod drop time of approximately 0.35 seconds from "full out" to "full in", the step change in reactivity assumption is not really justifiable when considering neutron flux predictions for times comparable to the rod drop time. A theoretical development is given for

EXPLANATION OF PLATE I

General view of the TRIGA reactor, looking into the reactor tank at the core. The ionization chambers are shown just outside of the core reflector.

PLATE I



reactivity insertion as a function of time. The resulting analytical prediction for the neutron flux as a function of time computed from the one group model was compared with the experimentally measured neutron flux after a change in reactivity.

Nomenclature

- n - Neutron density, neutrons/cm³.
- Φ - Thermal neutron flux, neutrons/cm². sec.
- D - Diffusion coefficient, cm.
- ∇^2 - Laplacian operator
- r - Spatial coordinate
- t - Time, seconds
- β - Fraction of total neutrons that are delayed
- β_i - Fraction of neutrons in the i th delayed group
- k_{eff} - Effective multiplication constant
- k - Infinite medium multiplication constant
- λ_i - Decay constant of precursor of the delayed neutrons of the i th group
- C_i - Precursor density of i th group, atomic nuclei/cm³.
- λ - Prompt neutron lifetime
- B^2 - Buckling
- Σ_a - Macroscopic absorption cross section, cm⁻¹.
- τ - Fermi age, cm².
- v - Neutron velocity, cm./sec.
- L - Diffusion length, cm.
- ρ - Reactivity

ω_i - Root of the characteristic equation

ρ - Resonance escape probability

Literature Survey

One of the first nuclear scientists to publish an extensive derivation of the "pile-kinetic" equations was Hurwitz (11). He described the methods used to solve the pile-kinetic equations for step changes in reactivity and when reactivity was a slowly varying function of time.

Kimel, et. al., (14) presented a theoretical development for the time behavior of neutron density as a function of step changes in reactivity. The space independent kinetic equations were based on the Fermi continuous slowing down model and used six groups of delayed neutrons. Experimental measurements were made with the Argonaut Reactor at Argonne National Laboratory. Reactivity worths of the control rods of the Argonaut were measured experimentally using the theoretical development. Both positive period measurements and negative rod drops measurements were made and good agreement was obtained between the two methods, indicating that the theoretical development could be used with the Argonaut reactor.

Smets (17) presented a solution for the one energy group bare reactor kinetic equations when the reactivity was a function of time. The time dependence of the reactivity was assumed to be either linear, exponential, or a reciprocal. Smets combined the basic space independent differential equations in the form

$$D \left[l + \sum_{i=1}^m \frac{\beta_i}{D + \lambda_i} \right] n(t) = \left[1 - \sum_{i=1}^m \frac{\beta_i D}{D + \lambda_i} \right] k_{\text{eff}}(t) \cdot n(t) \quad (1)$$

where

l = prompt neutron lifetime,

β = fraction of total neutrons that are delayed,

β_i = fraction of neutrons in the i th delayed group,

n = neutron density, neutron/cm³,

λ_i = decay constant of precursor of the delayed neutrons of the i th group,

k_{eff} = effective multiplication constant,

m = number of delayed neutron groups,

and D , in this case was the differential operator d/dt .

Solutions of Eq. (1) in the form of Laplace type contour integrals were obtained. Smets also found that the kinetic equations could be integrated for any piecewise analytic function, $k_{\text{eff}}(t)$, which consisted of straight lines, exponentials, and inverse functions but not the sum of any two or more such functions.

Toppel (20) discussed the errors in determining the magnitude of a step change in reactivity by means of measuring the stable reactor period. He found that for certain reactivities the neutron population would not be represented by a single exponential until several minutes after a positive step was made. Toppel also considered a reactor with a source and showed that a longer wait time was required for the true exponential neutron density distribution than with a source free reactor. A typical

point from Toppel's work showed that in the case of a positive insertion of reactivity, $\rho = 5.0 \times 10^{-4}$, and when the initial effective multiplication constant was 1.000, 100 seconds were required for the reactor to be within one per cent of being on a stable period. In contrast, if the initial reactivity was 0.99995 and then the reactivity was inserted, a wait time of approximately 300 seconds was required for the reactor to be within one per cent of being on a stable period.

Friedman (5) discussed the use of an "effective" fraction of delayed neutrons in place of the true fraction of neutrons. He stated that the fraction of delayed neutrons should be increased because the delayed neutrons have a larger nonleakage probability than prompt fission neutrons, thus a delayed neutron is more effective in producing another fission than is a prompt neutron. Friedman went on to say that if U^{238} was present in the reactor fuel, the fraction of delayed neutrons should be increased due to the increased fraction of delayed neutrons from U^{238} . An expression was given for the effective fraction of delayed neutrons for the i th precursor group (5).

$$\beta_{\text{eff};i} = \frac{p_{\text{delayed}}}{p_{\text{prompt}}} \left[\frac{\beta_i^{25}}{\epsilon^{28}} + \frac{\nu^{28}}{\epsilon^{25}} \left(\frac{\delta^{28}}{\epsilon^{28}} \right) \beta_i^{26} \right] \quad (2)$$

Where

$$\beta_i \begin{matrix} (25) \\ (28) \end{matrix} \equiv \text{fraction of neutrons which are from the } i\text{-th precursor due to fission from } U \begin{matrix} (25) \\ (28) \end{matrix}.$$

ρ $\begin{matrix} \text{(delayed)} \\ \text{(prompt)} \end{matrix}$ \equiv nonleakage probability of $\begin{matrix} \text{(delayed)} \\ \text{(prompt)} \end{matrix}$ neutrons

γ $\begin{matrix} \text{(28)} \\ \text{(25)} \end{matrix}$ \equiv average number of neutrons per fission from $\begin{matrix} \text{(28)} \\ \text{(25)} \end{matrix}$
U

$$\delta^{20} \equiv \frac{U^{20} \text{ fissions}}{U^{25} \text{ fissions}}$$

$\epsilon^{20} \equiv$ fast fission factor of U^{238}

Ewen and Wick (4) extended the work of Friedman (5) pertaining to evaluating the effective fractions of delayed neutrons. They gave results of several cores and showed the difference in core reactivity evaluation by using their calculated effective fraction of delayed neutrons. Keepin, et. al., (13) measured the period, relative abundance, and absolute yield of delayed neutrons from "fast" fission for six nuclides including U^{235} and from thermal fission for three nuclides including U^{235} . The "Godiva" reactor was used as the neutron source. Six exponential periods were found to be necessary and sufficient for an optimum least-square fit of the experimental data. Keepin's work on determination of delayed neutron parameter is the latest presented and is probably the most reliable analytical study that has been made on the delayed neutron phenomenon to date. Many persons have experimentally measured sets of delayed neutron data and most of the sets are slightly different. Keepin

(12) presented a list of experimentally measured sets of delayed neutron constants and presented an excellent description of the delayed neutron phenomenon.

Skinner and Cohen (16) showed that a two delayed neutron group model represented, to a fair degree of accuracy, the kinetic behavior of a reactor when a positive insertion of reactivity was considered. They also demonstrated that at least three groups must be employed for the delayed neutron model used to describe the time behavior of neutron density in a reactor when negative reactivities are inserted into the core. The principal advantages given for using the reduced delayed neutron group models were: 1) the number of differential equations to be solved was reduced, 2) the solutions of the kinetic equations required less computer time, 3) control design was greatly simplified. The following is a list of three group constants as given by Skinner and Cohen (16).

Table 1. Reduced delayed neutron group constants for thermal fission of U^{235} .

Group	λ_i (sec ⁻¹)	β_i/β
1	0.01244	0.033
2	0.03694	0.346
3	0.0632	0.621

Most methods of solving the reactor kinetic differential equations include the assumption that the neutron flux can be separated into two variables, space and time. Solutions of the

kinetic equations are difficult to obtain when the separation process is not performed. Garabedian and Leffert (6) solved the kinetic equations where it was assumed that the spatial distribution contained a time function. They found, in the case examined, that the space function reached an asymptotic shape within 0.005 seconds after a flux perturbation was caused by a sudden reactivity change.

THEORETICAL DEVELOPMENT

Since the theory used by this paper is very similar to that found in references (8), (11), and (7), the detailed development and solution of the reactor kinetics equations is found in Appendix - A. The kinetics model and the associated assumptions are given along with the resulting equations.

A reactor system is described by the time dependent thermal diffusion equation and the precursor rate equations. The assumptions used are listed as follows:

- 1) the reactor is homogeneous,
- 2) the reactor is bare, i.e., has no reflector,
- 3) the delayed neutron energy spectrum is the same as that of the prompt fission neutrons,¹
- 4) there is no fission from U^{238} ,
- 5) the spatial distribution of the neutron flux may be represented by the fundamental mode of the wave equation,

¹This assumption need not be made (see Appendix - A).

6) the delayed neutron precursor spatial distribution is proportional to the spatial distribution of the neutron flux,

7) the Fermi age equation may be used as a source of thermal neutrons for the thermal diffusion equations.

The thermal neutron balance written around a volume element of the reactor is

$$D \nabla^2 \Phi(r,t) - \Sigma_a \Phi(r,t) + S(r,t) = \frac{\partial n(r,t)}{\partial t} = \frac{1}{v} \frac{\partial \Phi(r,t)}{\partial t} \quad (3)$$

The precursor rate equation is

$$\frac{\partial C_i(r,t)}{\partial t} = -\lambda_i C_i(r,t) + \frac{\beta_i \Sigma_a \Phi(r,t) k}{\rho} \quad (4)$$

The space and time variables of Eq.'s (3) and (4) are separated and the resulting space independent kinetic equations are

$$\left[(1-\beta) k_{\text{eff}} - 1 \right] \frac{n(t)}{\ell} + \sum_{i=1}^6 \lambda_i H_i(t) = \frac{dn(t)}{dt} \quad (5)$$

and

$$\frac{dH_i(t)}{dt} = -\lambda_i H_i(t) + \frac{k_{\text{eff}} n(t) \beta_i}{\ell} \quad (6)$$

By assuming a solution for the set of seven linear differential equations of the form

$$n(t) = A' e^{wt} \quad (7)$$

and

$$H_i(t) = B'_i e^{\omega t} \quad (8)$$

a characteristic equation may be obtained as

$$\omega = \frac{\rho - \beta}{\ell(1 - \rho)} + \frac{1}{\ell(1 - \rho)} \sum_{i=1}^6 \frac{\beta_i \lambda_i}{\omega + \lambda_i} \quad (9)$$

or

$$\rho = \frac{\ell\omega}{1 + \ell\omega} + \frac{1}{1 + \ell\omega} \sum_{i=1}^6 \frac{\lambda_i \beta_i}{\omega + \lambda_i} \quad (10)$$

The neutron density as a function of time is expressed as

$$\frac{n(t)}{n(0)} = \frac{\phi(t)}{\phi(0)} = \sum_{j=1}^7 A_j e^{\omega_j t} \quad (11)$$

where the coefficients, A_j 's, are

$$A_j = (1 - \rho) \left[\frac{\ell + \sum_{i=1}^6 \frac{\beta_i}{\omega_j + \lambda_i}}{\ell(1 - \rho) + \sum_{i=1}^6 \frac{\beta_i \lambda_i}{(\omega_j + \lambda_i)^2}} \right] \quad (12)$$

Three groups of delayed neutrons are used as an approximation for the usual six groups. This simplifies the solution of

the kinetic differential equations but they are of the same general form as Eq.'s (7) and (8).

When reactivity is a function of time, the differential equations are, in some instances, very difficult to solve; therefore, an approximation may be made where it is assumed that the six groups of delayed neutrons may be represented by one group of delayed neutrons. McPhee (15) developed a solution to the approximated one group kinetic model and a similar model is developed in Appendix - B. McPhee's equations were developed for a linear time dependent reactivity insertion whereas the kinetic development given in Appendix - B, of this paper, is based on the assumption of reactivity insertion rate empiricized to

$$k_{\text{off}}(t) = 1 - A(1 - e^{-bt})^n. \quad (13)$$

A is a parameter determined by the final value of the reactivity change and b and n are constants determined empirically.

COMPUTED DATA

General

The solutions of the differential equations for several kinetics models were programmed for the Kansas State University IBM-650 computer. Several neutron parameter models were used. First, the usual equations using six groups of delayed neutrons were programmed and two sets of neutron parameters were used, Hughes' (10) and Keepin's (13). Next, an approximation of the

six groups of delayed neutrons by three groups was made. The three delayed neutron parameters were chosen such that the approximation would predict the same time dependent neutron flux after a step change in reactivity as did the six group model. Finally, the kinetic equations were approximated with one group of delayed neutrons and were solved for reactivity as a function of time.

The equations of interest (see Appendix - A) for the six group analysis with step inputs of reactivity were the characteristic equation,

$$\rho = \frac{\lambda\omega}{\lambda\omega + 1} + \frac{1}{\lambda\omega + 1} \sum_{i=1}^6 \frac{\lambda_i\omega}{\omega + \lambda_i} \quad (9)$$

and the equation of neutron flux as a function of time normalized to the critical flux just before the step change in reactivity

$$\frac{\phi(t)}{\phi(0)} = \sum_{j=1}^7 A_j e^{\omega_j t} \quad (11)$$

An IBM-650 program was written to compute roots of Eq. (9) and to compute the A_j 's of Eq. (11). The computed data were used to analyze the experimental results from both the TRIGA reactor and the Argonaut reactor; therefore two different prompt neutron lifetimes, 2.0×10^{-4} seconds for the Argonaut and 8.0×10^{-5} seconds for the TRIGA reactor, were used in the

computations. The parameter variation in each set of data was reactivity.

Delayed Neutron Parameters

The two delayed neutron parameter models used in the theoretical computations were those as given by Keepin (13) and by Hughes (10). The two sets of neutron parameters gave slightly different sets of data. The following tables list the delayed neutron constants used in the analytical computations. A list of several other sets of delayed neutron parameters was made by Keepin (12).

Table 2. Keepin's delayed neutron data for thermal fission of ^{235}U .

Group	Half-life (Sec.)		λ_i (Sec. ⁻¹)	β_i/β
1	55.72	1.28	0.01244	0.033 ± 0.003
2	22.72	0.71	0.03051	0.219 ± 0.009
3	6.22	0.23	0.1114	0.196 ± 0.022
4	2.30	0.09	0.3014	0.395 ± 0.011
5	0.61	0.063	1.1363	0.115 ± 0.009
6	0.236	0.025	3.0137	0.042 ± 0.008

$$\beta = 0.0064$$

Table 3. Hughes' delayed neutron data for thermal fission of ^{235}U .

Group	Half-life (Sec.)	λ_i (Sec. ⁻¹)	β_i/β
1	55.6	0.2	0.01246
2	22.0	0.2	0.03149
3	4.51	0.1	0.1537
4	1.52	0.05	0.456
5	0.43	0.05	1.612
6	0.05	0.02	13.86

$$\beta = 0.00755$$

Reactivity Parameter

Computed data were calculated over wide ranges of reactivity, from $-\$0.001$ to $-\$10.0$ for negative reactivities and from $\$0.001$ to $\$1.20$ for positive reactivities (see PLATES II through XXI).

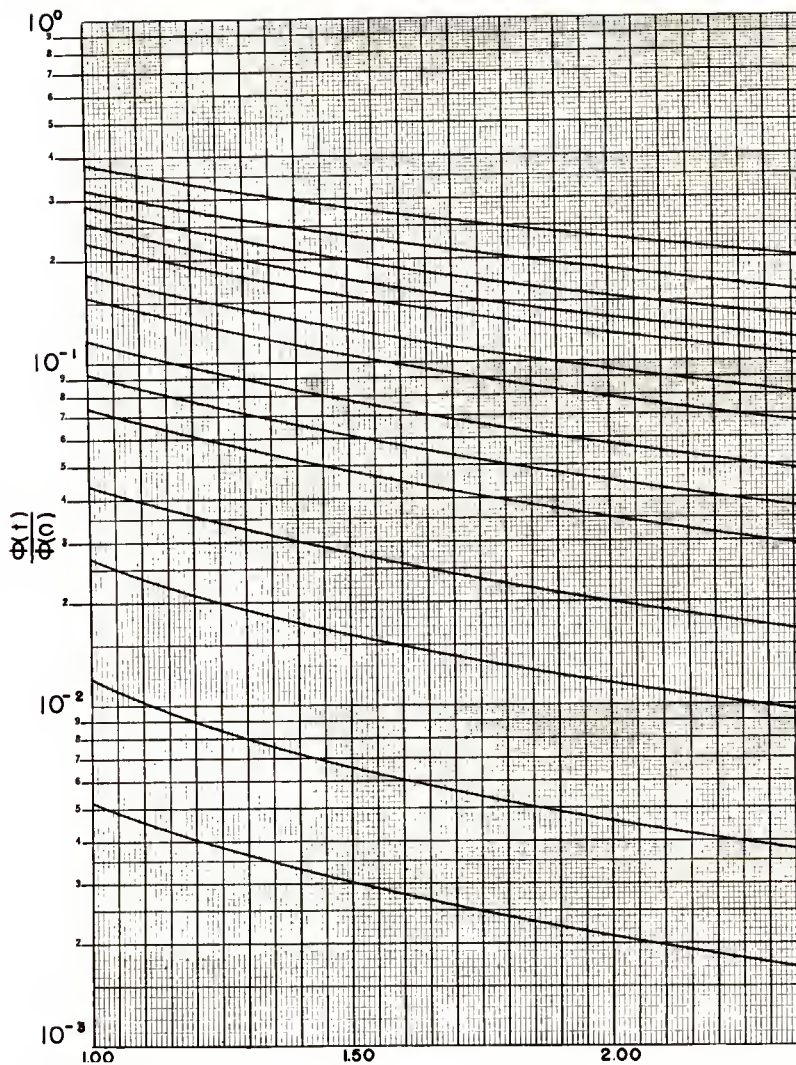
EXPLANATION OF PLATE II

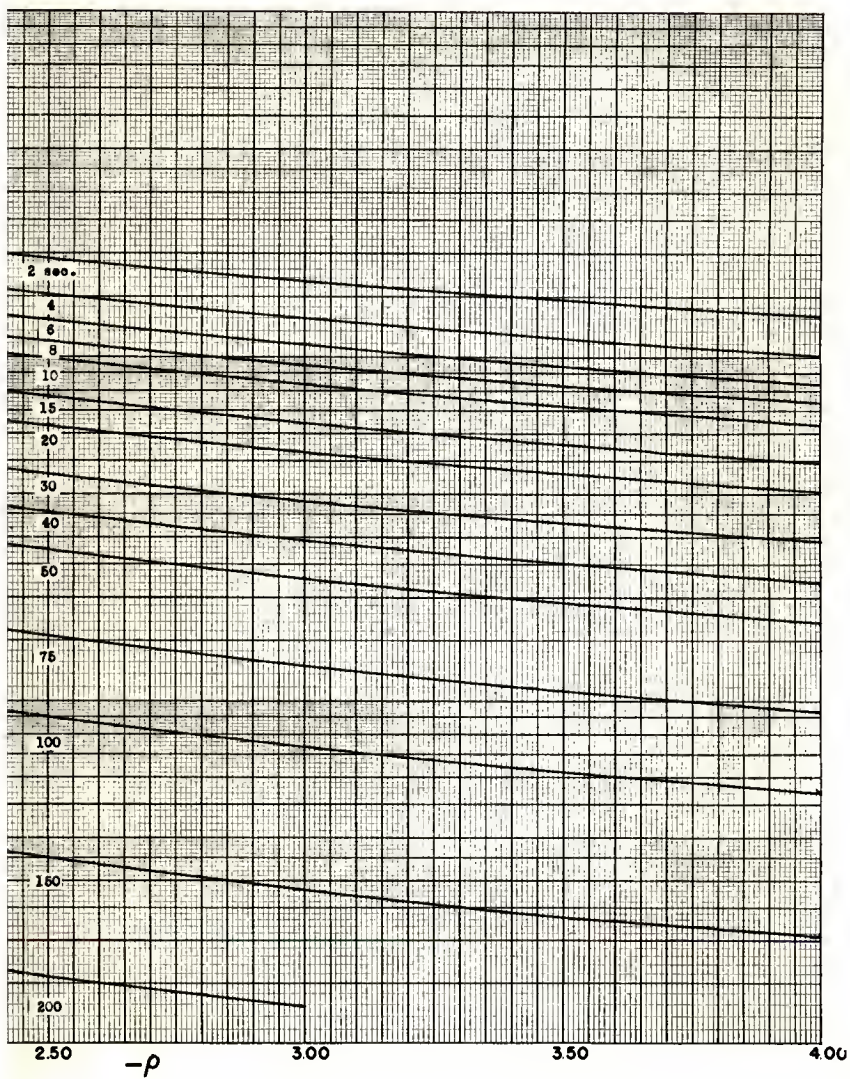
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of negative reactivity, ρ (dollars), for various times, t (seconds), after the reactivity change.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Hughes' delayed neutron parameters.

Data computed from Eq. (11).





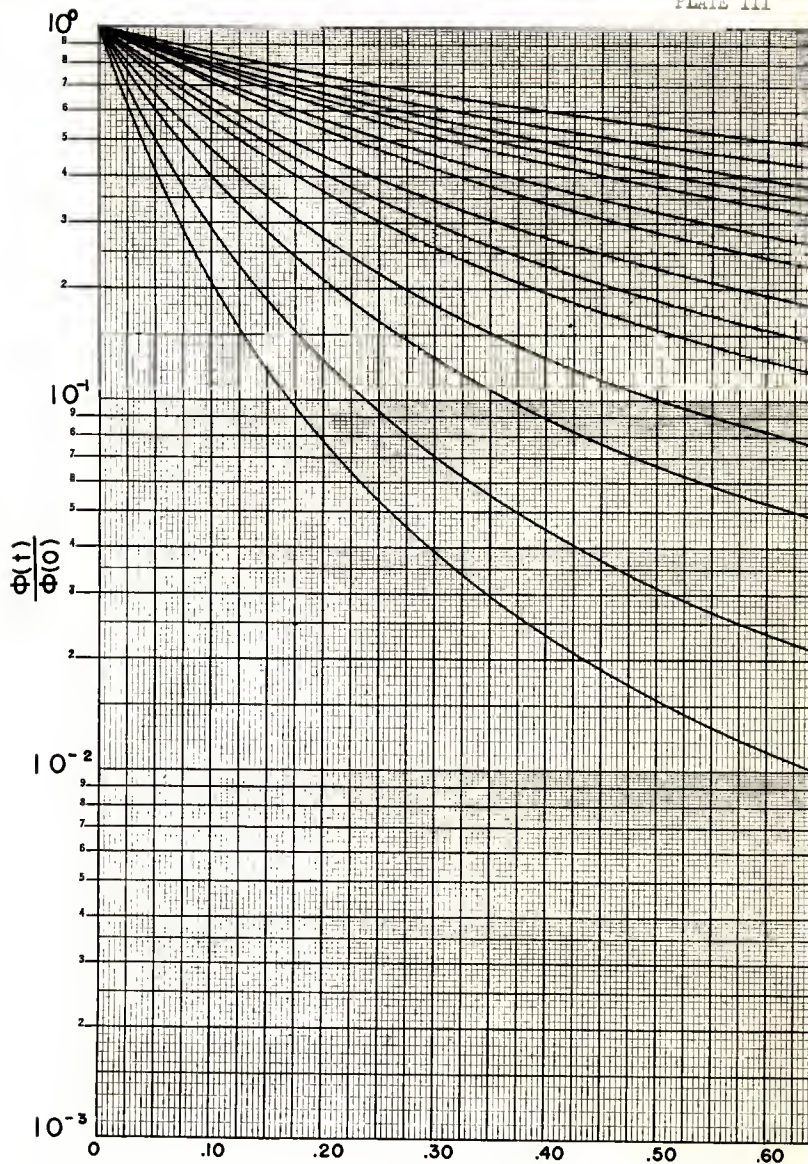
EXPLANATION OF PLATE III

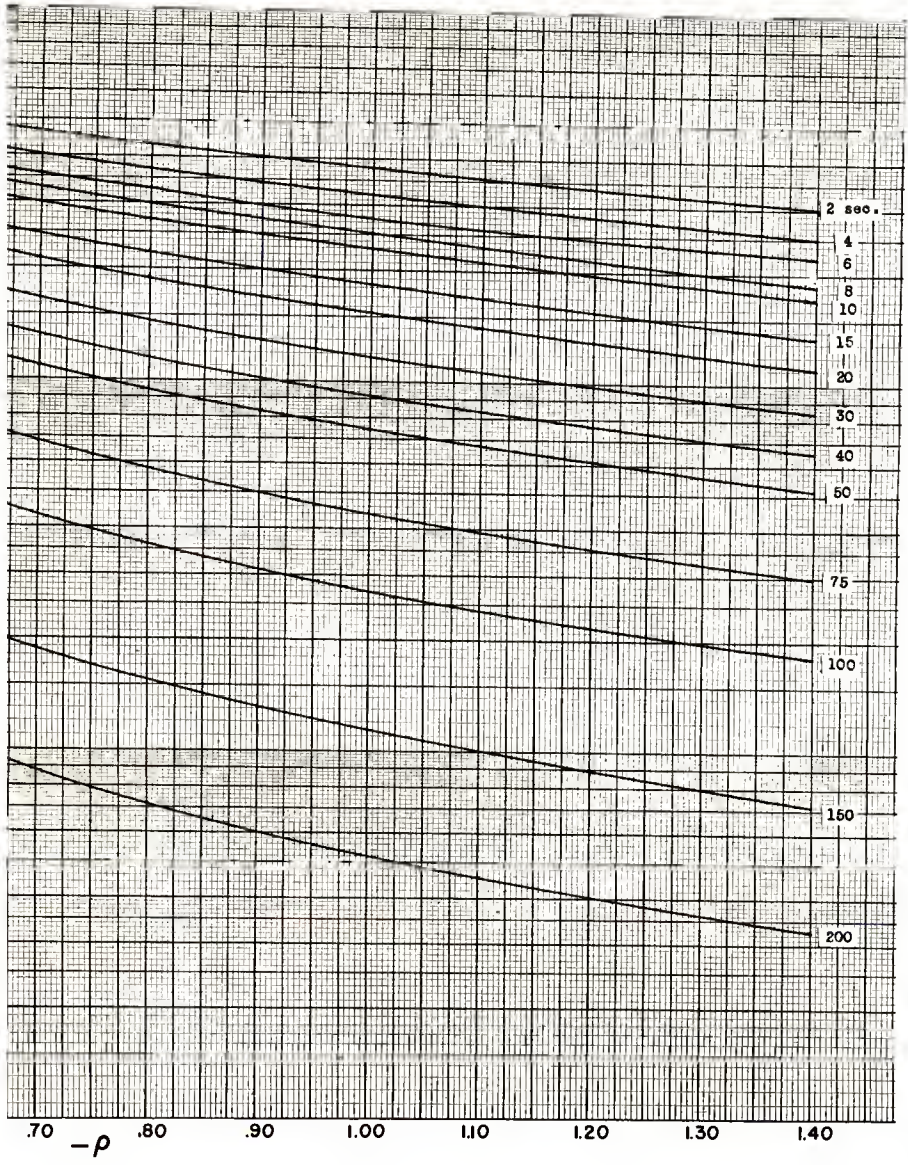
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of negative reactivity, ρ (dollars), for various times, t (seconds), after the reactivity change.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Hughes' delayed neutron parameters.

Data computed from Eq. (11).





EXPLANATION OF PLATE IV

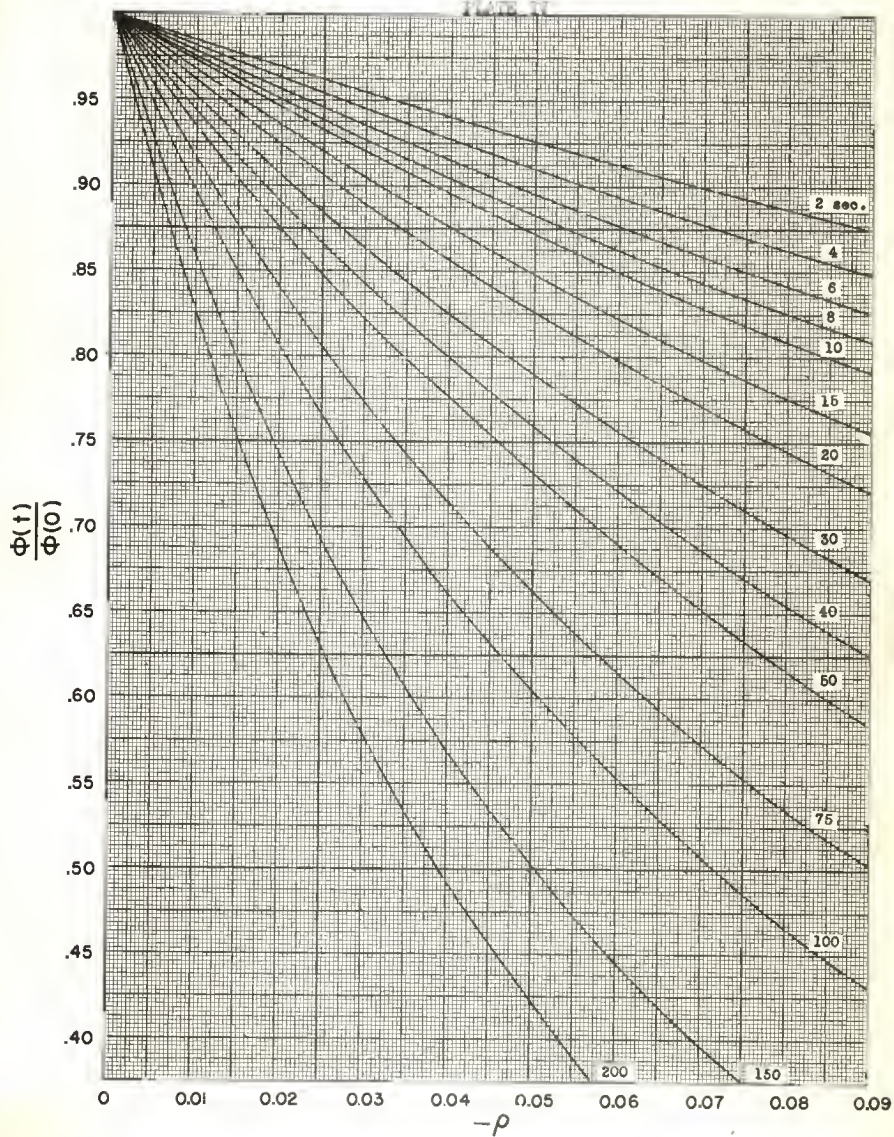
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of negative reactivity, ρ (dollars), for various times, t (seconds), after the reactivity change.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Hughes' delayed neutron parameters.

Data computed from Eq. (11).

TABLE II



EXPLANATION OF PLATE V

Negative reactivity, ρ (dollars), vs. roots, ω (seconds⁻¹), of the characteristic equation for step changes of reactivity.

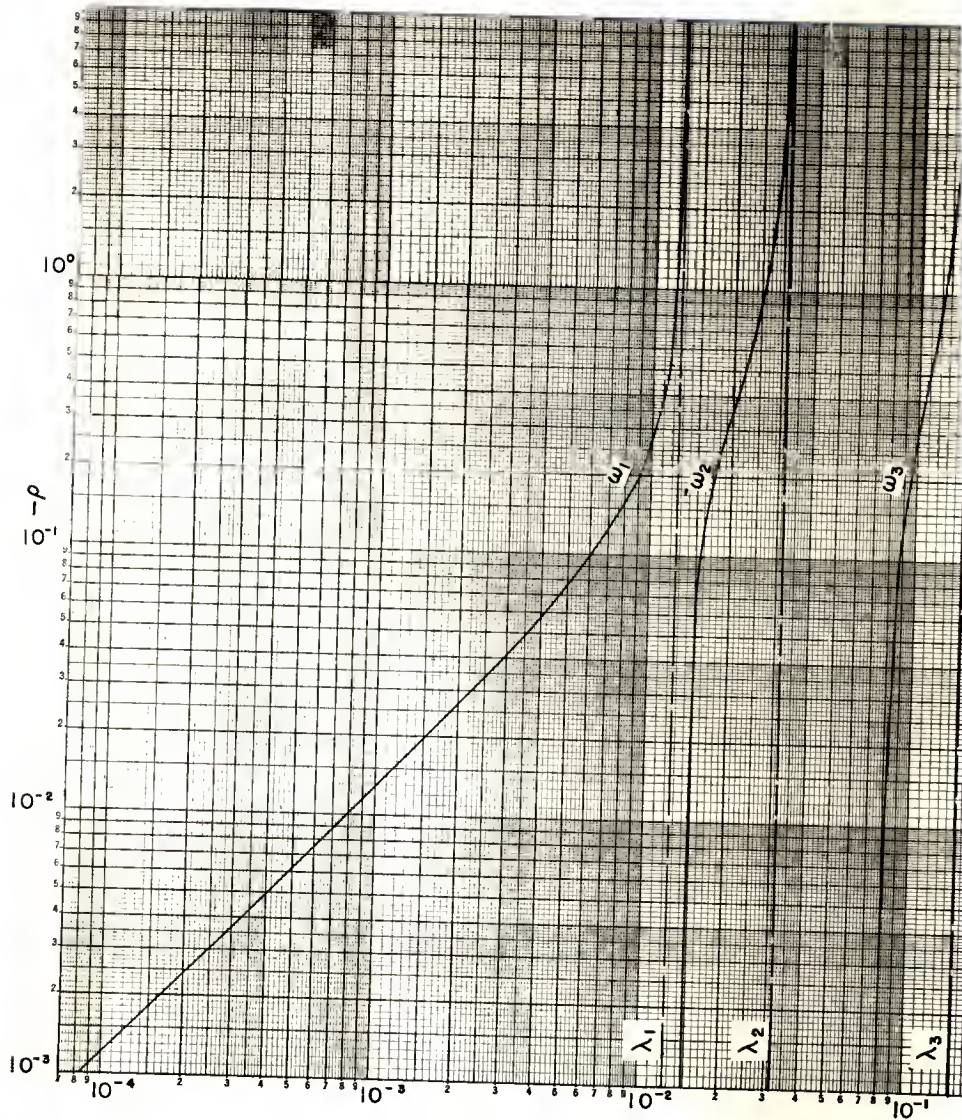
All roots, ω , are negative in value.

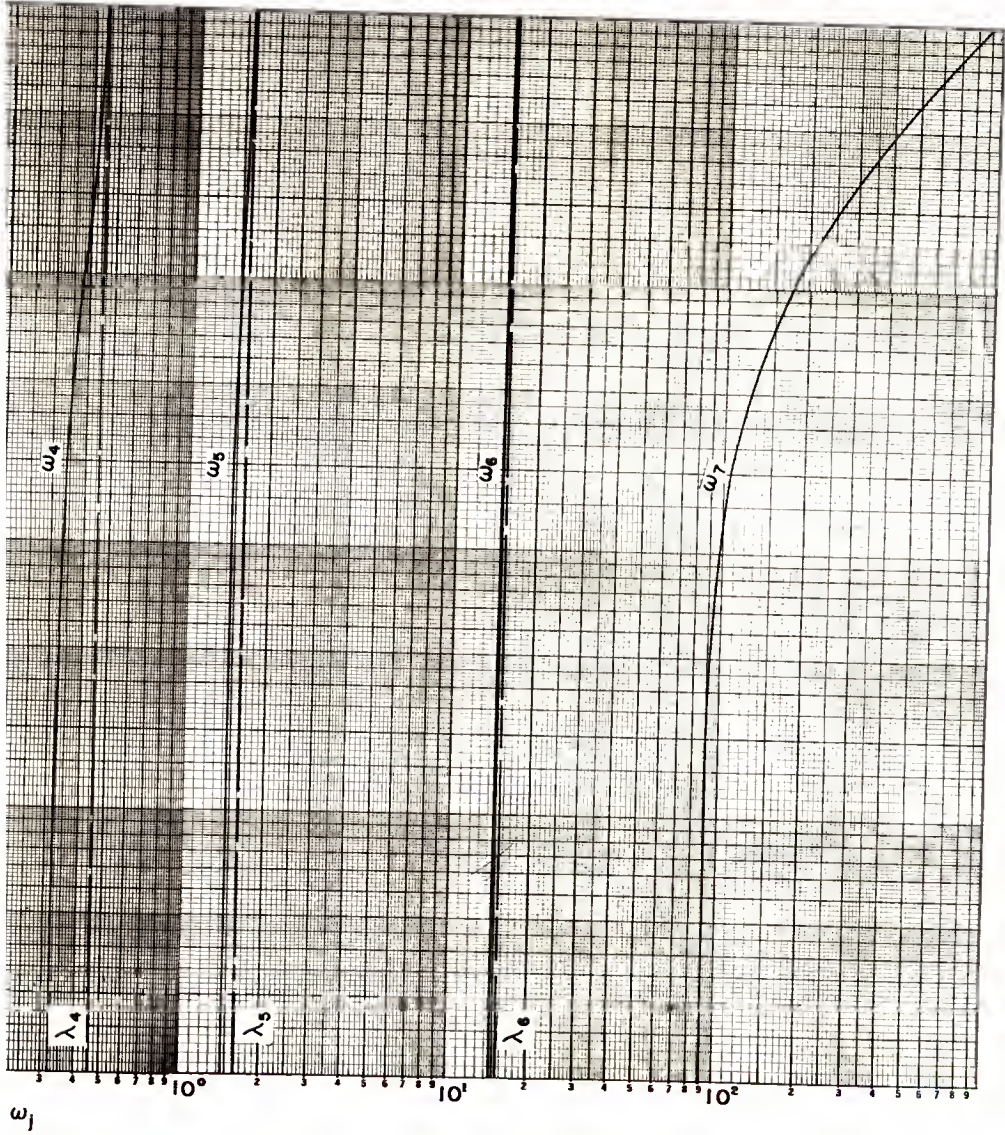
The λ 's indicate the decay constants of the neutron precursor groups.

Prompt neutron lifetime, $l = 8.0 \times 10^{-5}$ second.

Hughes' delayed neutron parameters.

Data computed from Eq. (9).





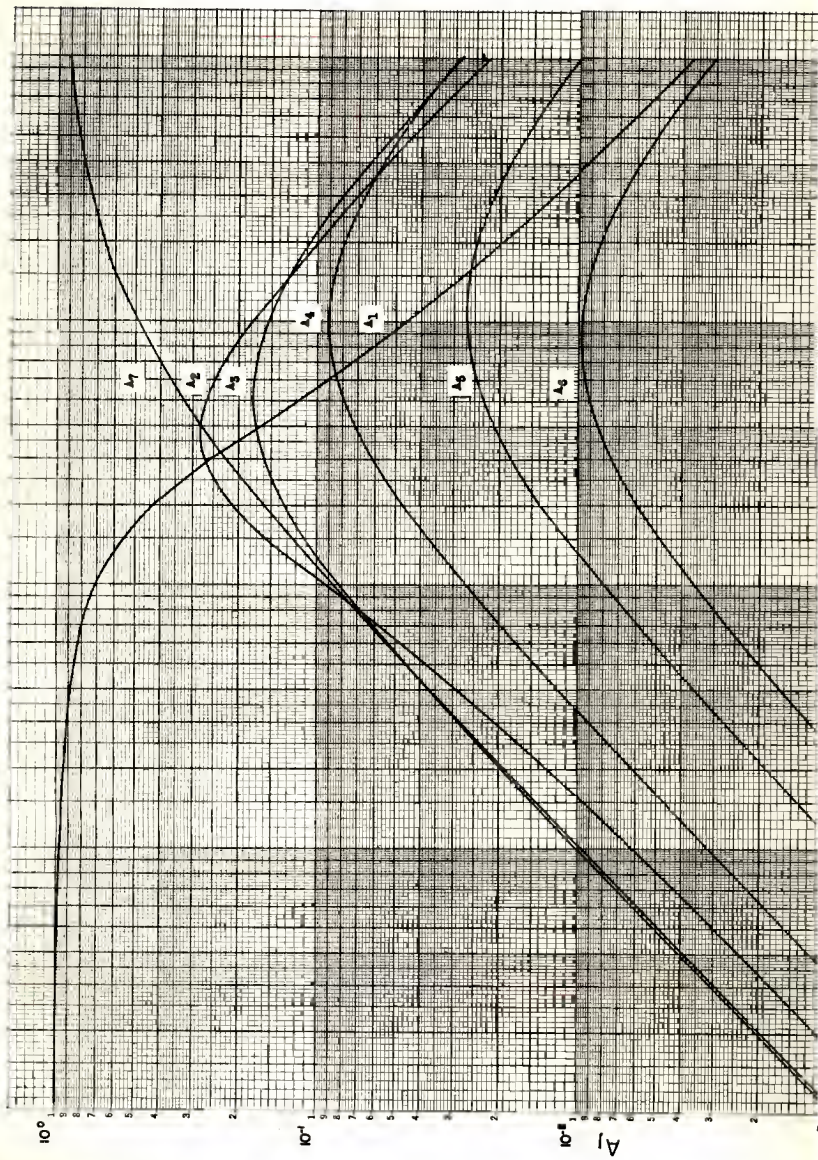
EXPLANATION OF PLATE VI

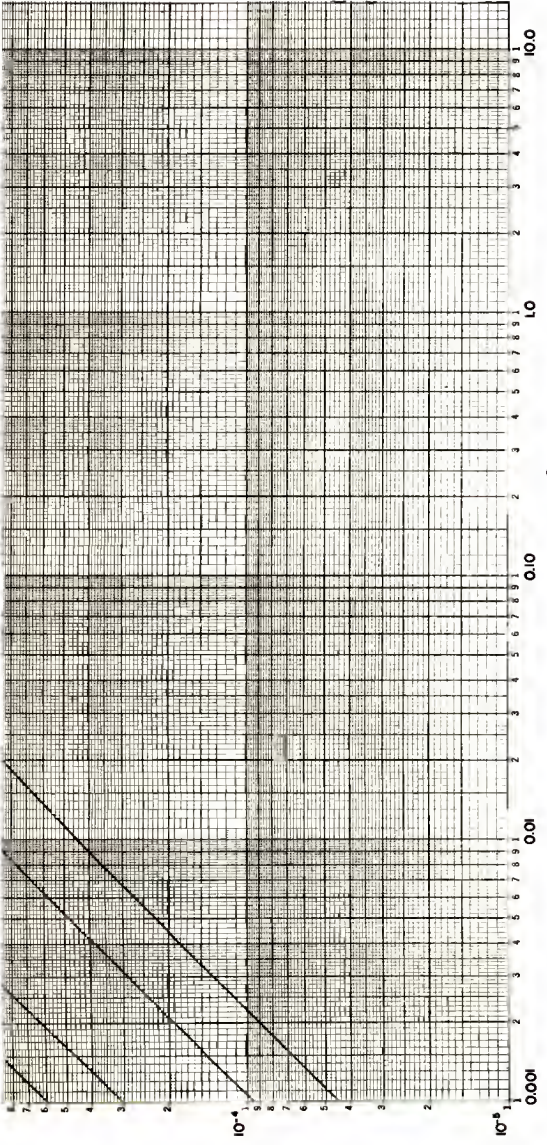
Dimensionless coefficients, A_j 's, of the flux equation
vs. negative reactivity, ρ (dollars), for step
changes in reactivity.

Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.
Hughes' delayed neutron parameters.

Data calculated from Eq. (12).

PLATE VI





EXPLANATION OF PLATE VII

Positive reactivity, ρ (dollars), vs.
positive asymptotic period, τ (seconds).
Prompt neutron lifetime, $\ell = 8.0 \times 10^{-5}$
seconds.
Hughes' delayed neutron parameters.
Data computed from Eq. (9).

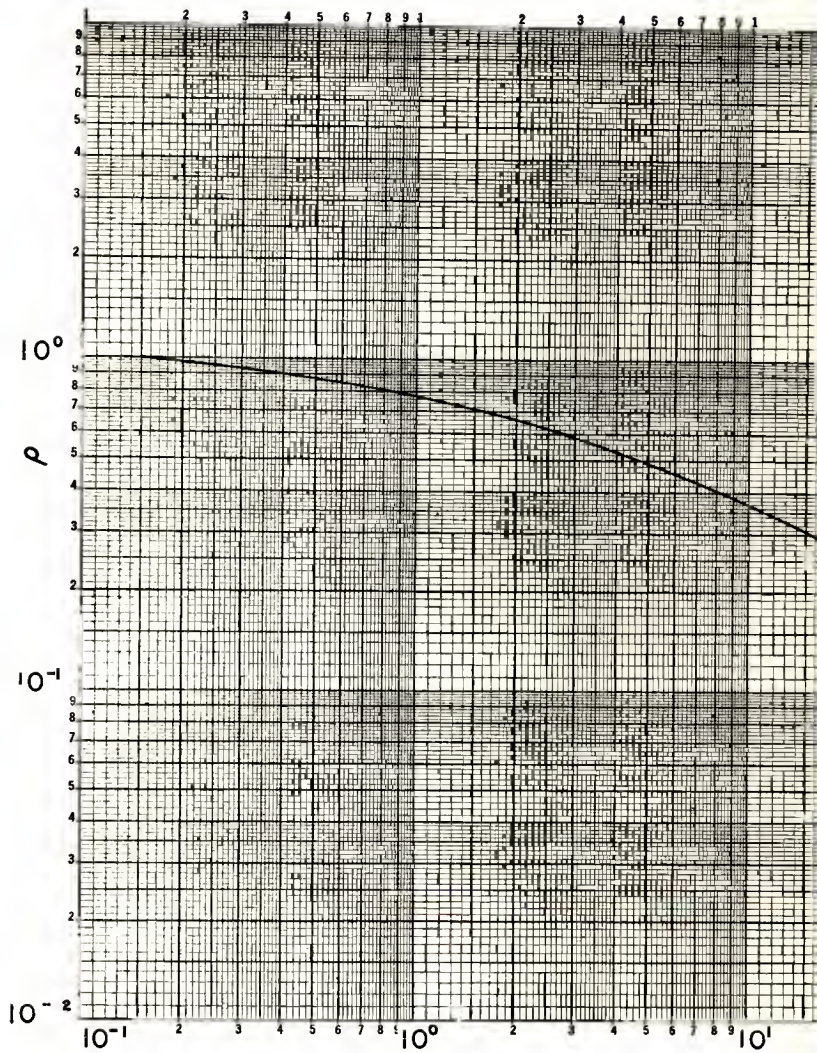
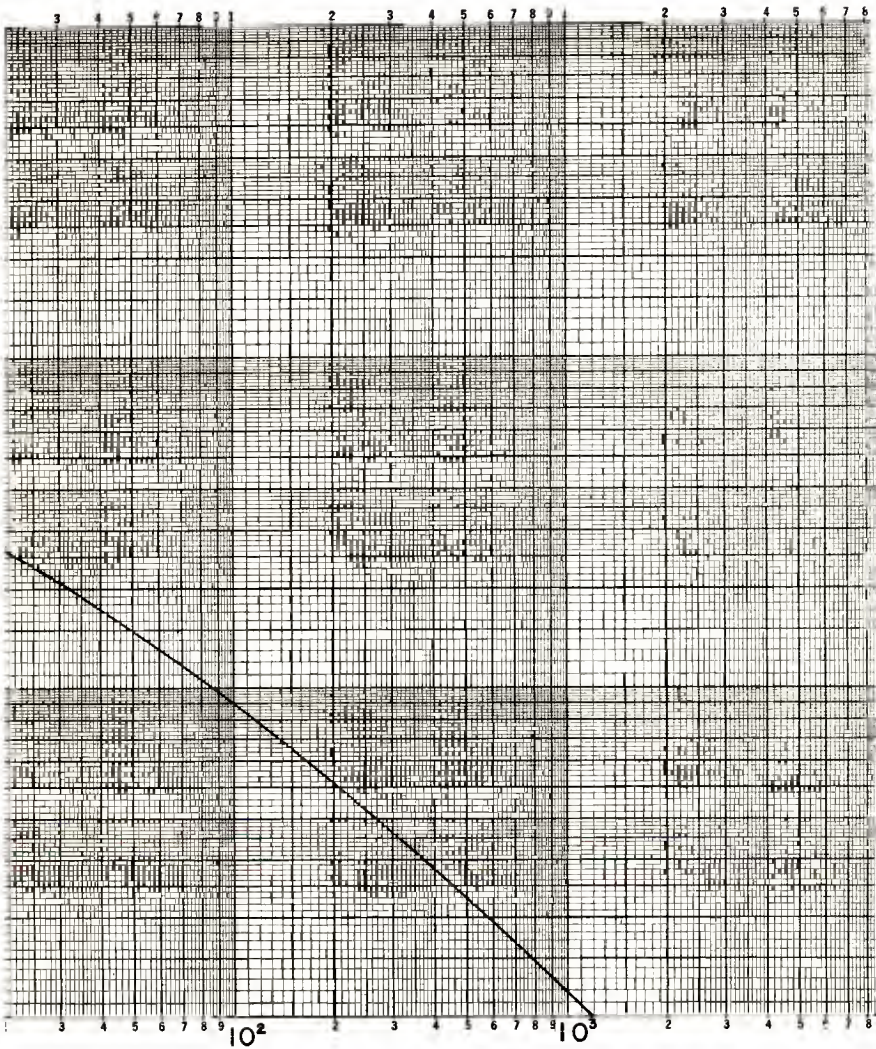


FIGURE VII



τ

EXPLANATION OF PLATE VIII

Positive reactivity, ρ (dollars), vs. roots, ω 's (seconds⁻¹) of the characteristic equation for step changes in reactivity.

All roots are negative except the indicated positive root.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Hughes' delayed neutron parameters.

Data computed from Eq. (9).

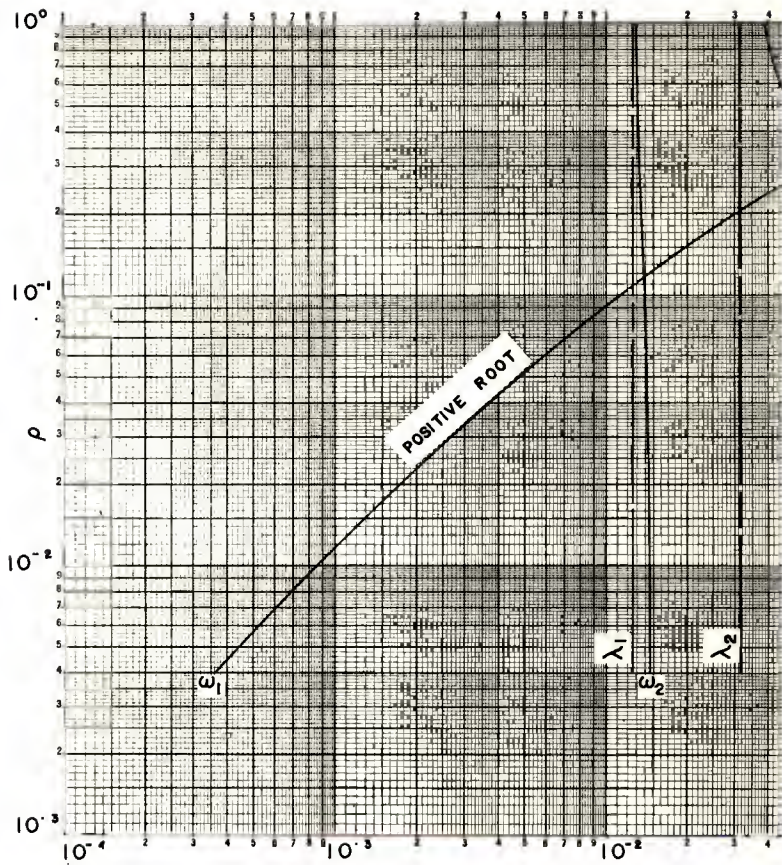
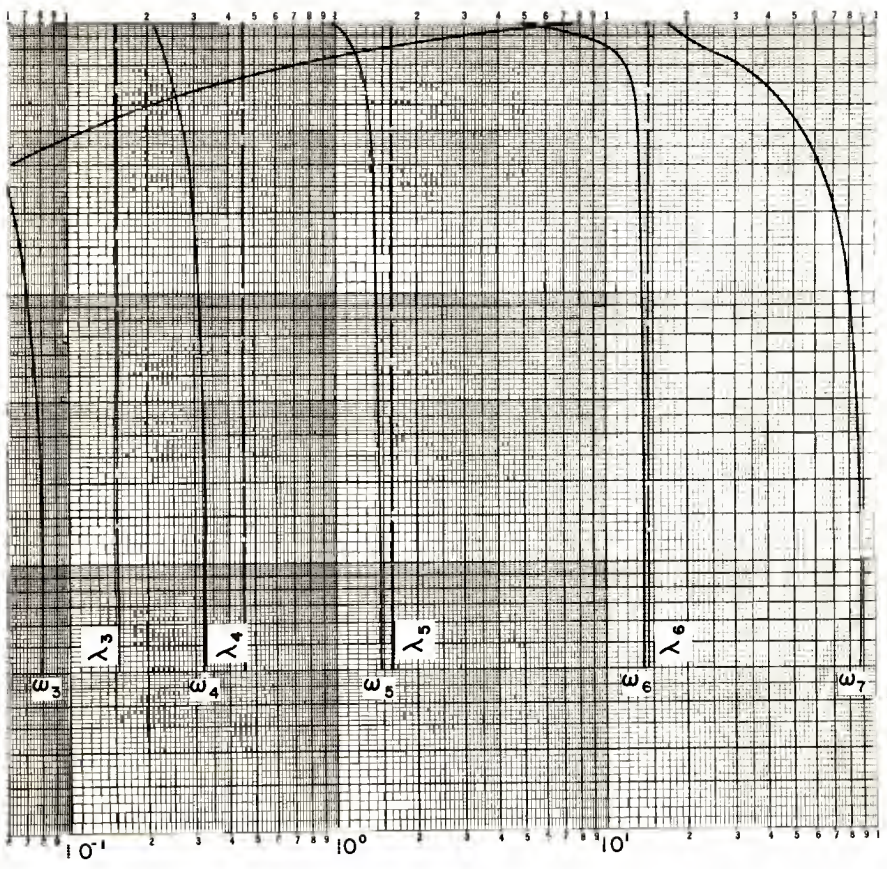


PLATE VIII



ν_j

EXPLANATION OF PLATE IX

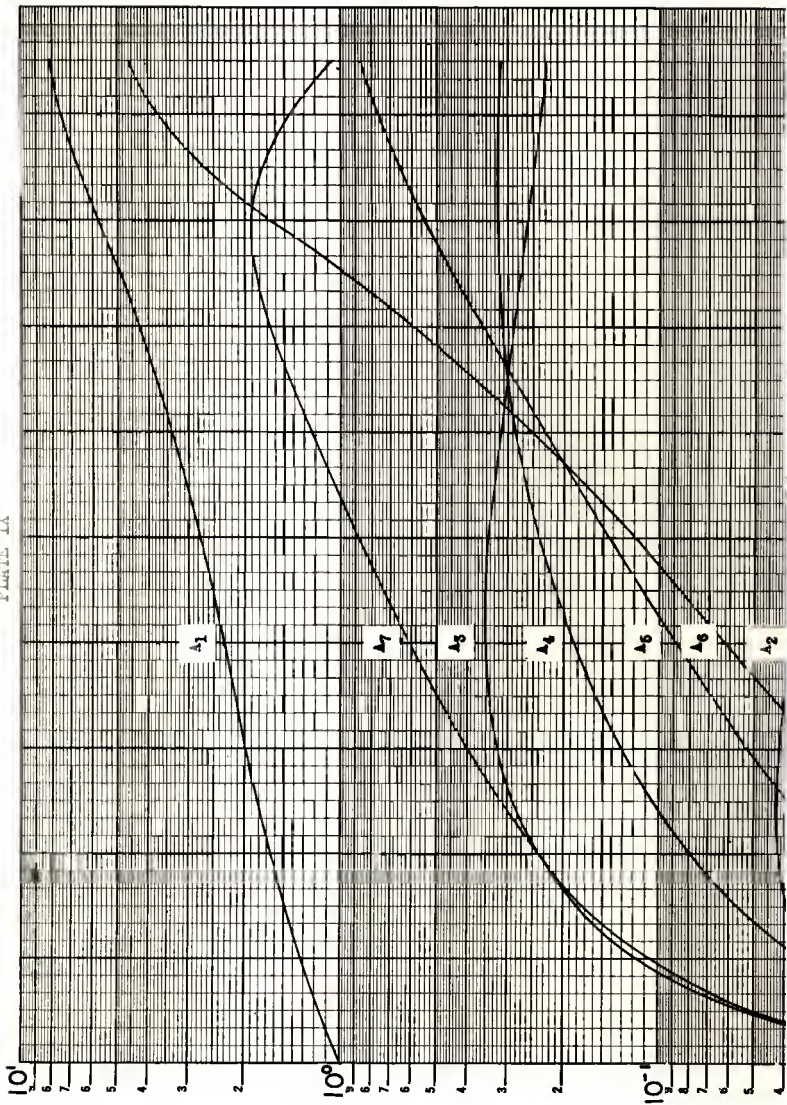
Dimensionless coefficients, A_j 's of the flux equation
vs. positive reactivity, ρ (dollars), for step changes
in reactivity.

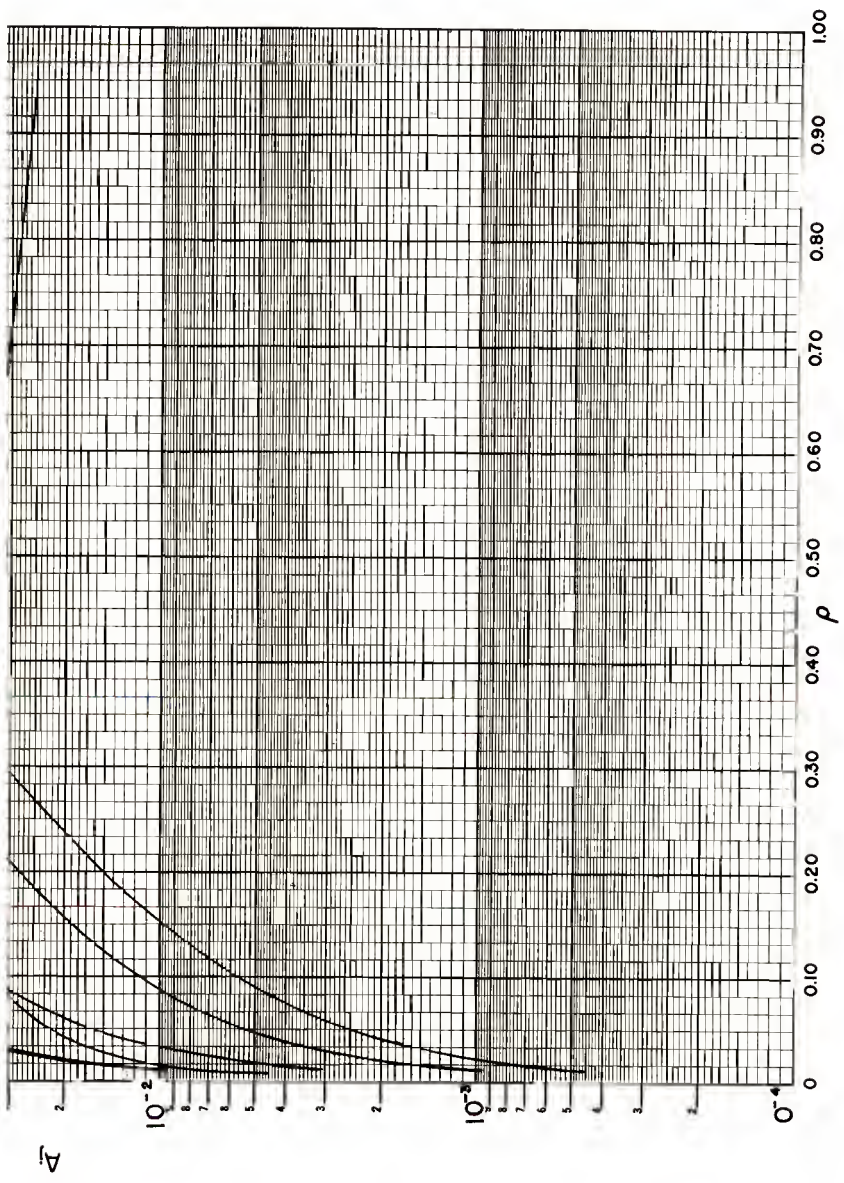
Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Hughes' delayed neutron parameters.

Data computed from Eq. (12).

PLATE IX





EXPLANATION OF PLATE X

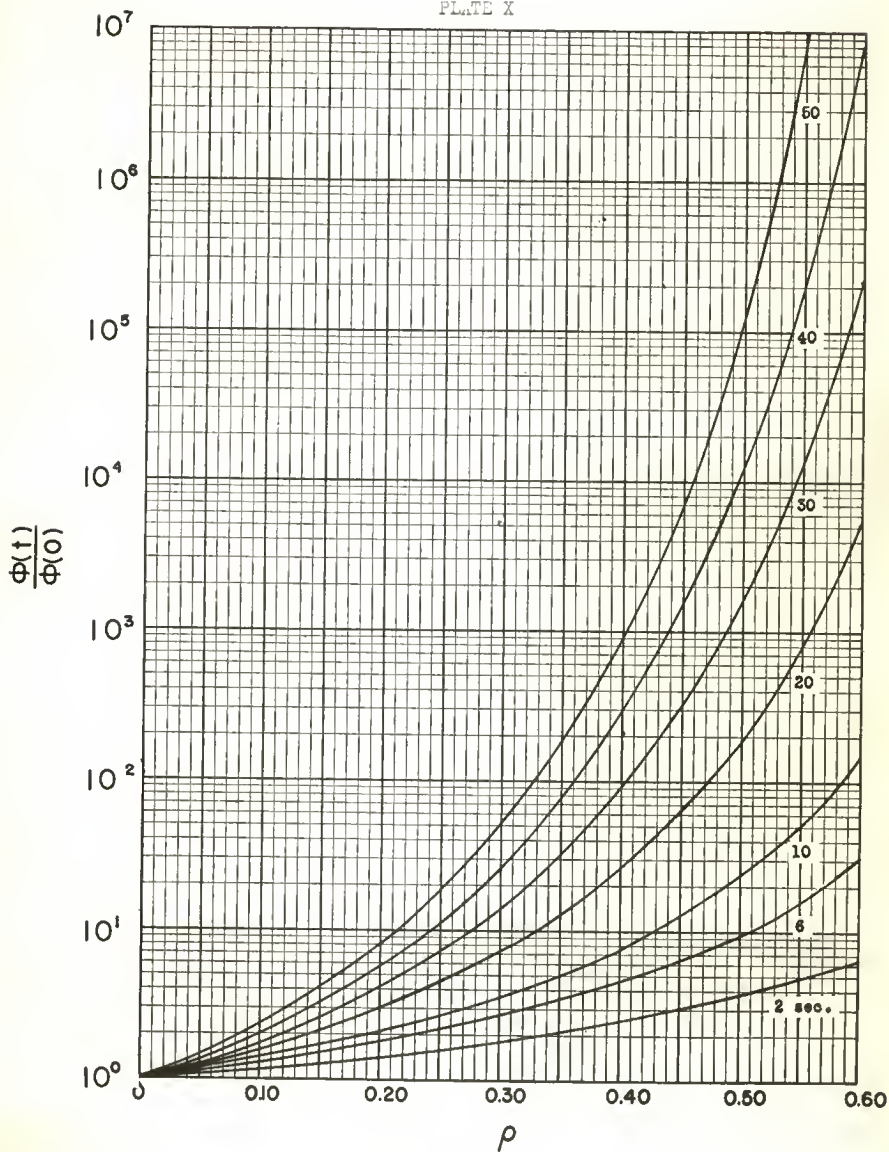
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of positive reactivity, ρ (dollars), for various times, t (seconds), after the reactivity step.

The prompt neutron lifetime, $\lambda \approx 8.0 \times 10^{-5}$ seconds.

Hughes' delayed neutron parameters.

Data computed from Eq. (11).

PLATE X



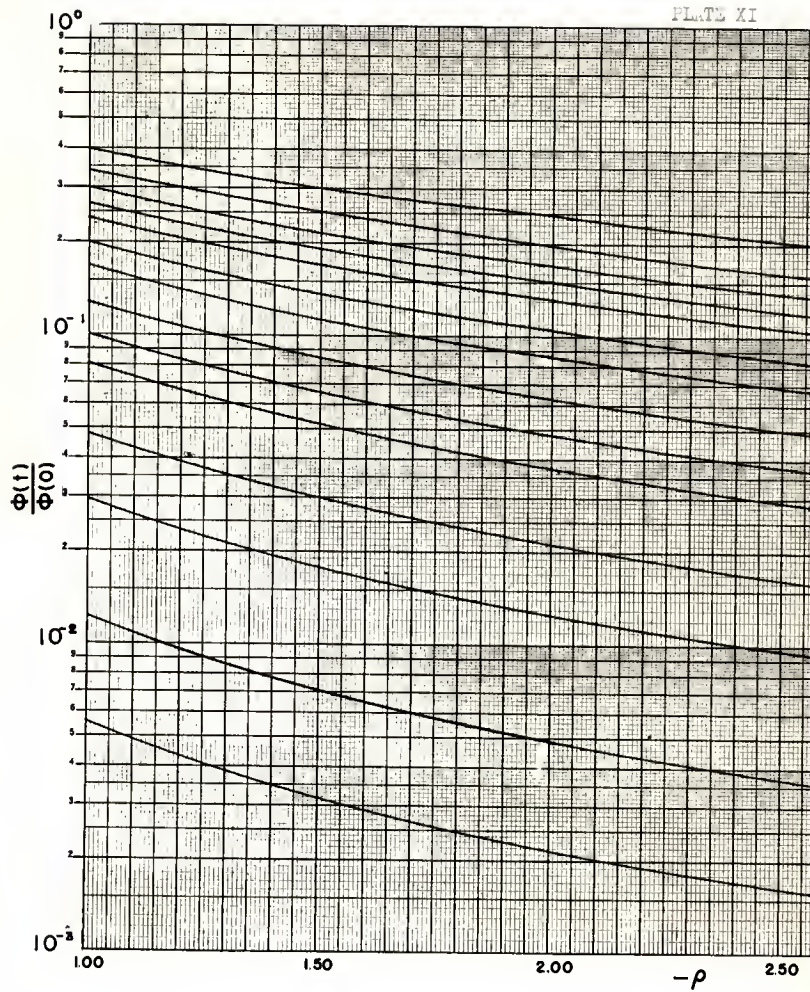
EXPLANATION OF PLATE XI

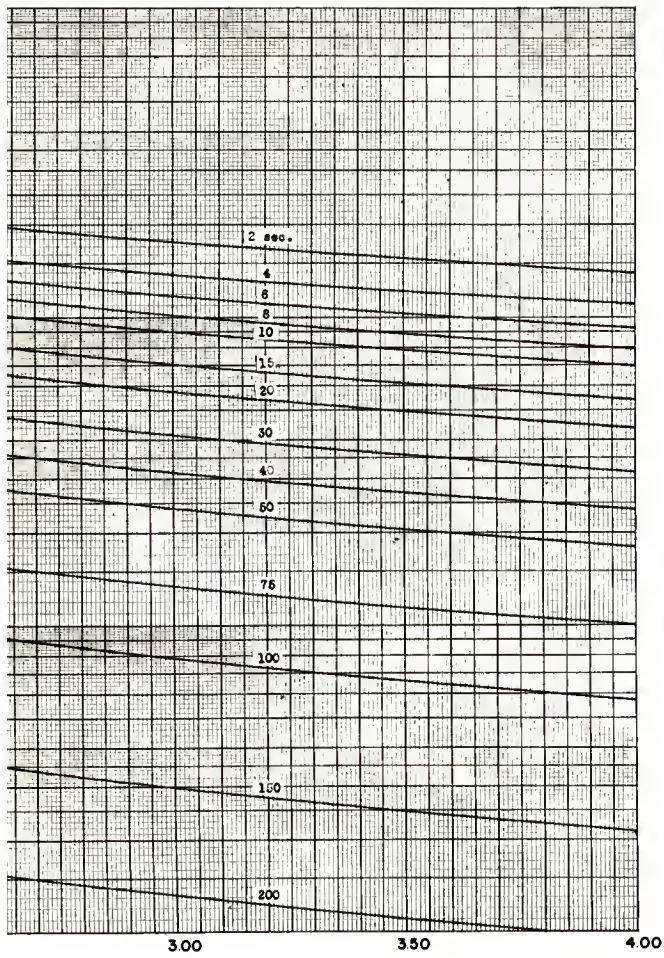
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of negative reactivity, ρ (dollars), for various times, t (seconds), after the reactivity step.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keepin's delayed neutron parameters.

Data computed from Eq. (11).





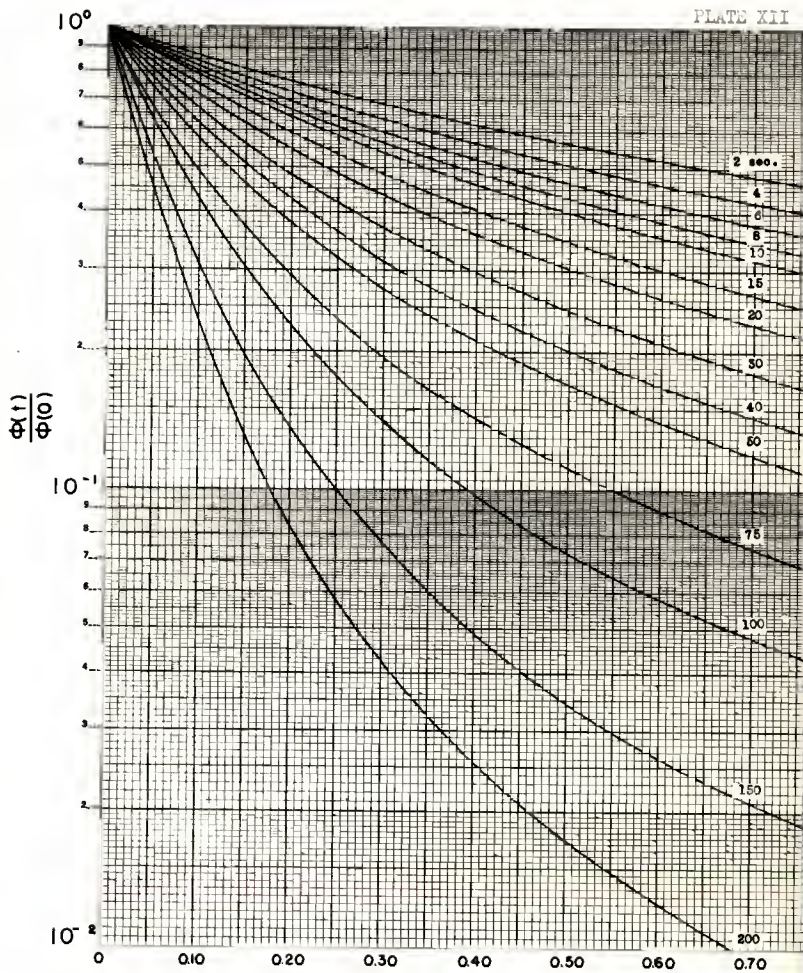
EXPLANATION OF PLATE XII

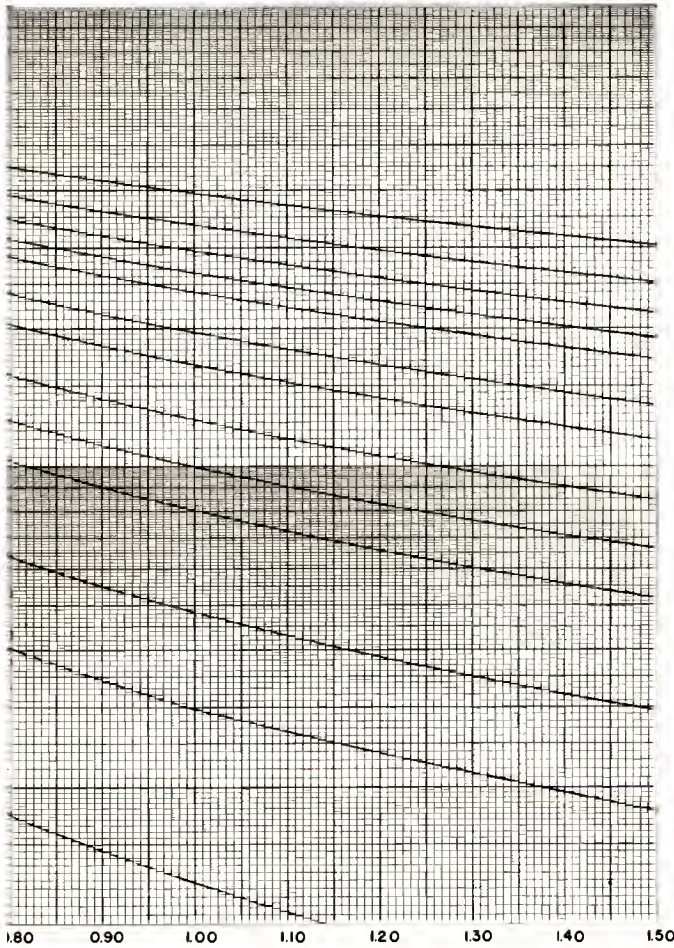
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of negative reactivity, ρ (dollars), for various times, t (seconds), after the reactivity step.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keepin's delayed neutron parameters.

Data computed from Eq. (11).





EXPLANATION OF PLATE XIII

Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of negative reactivity, ρ (dollars), for various times, t (seconds), after the reactivity step.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keepin's delayed neutron parameters.

Data computed from Eq. (11).

EXPLANATION OF PLATE XIV

Negative reactivity, ρ (dollars), vs. roots, ω 's (seconds), of the characteristic equation for step changes of reactivity.

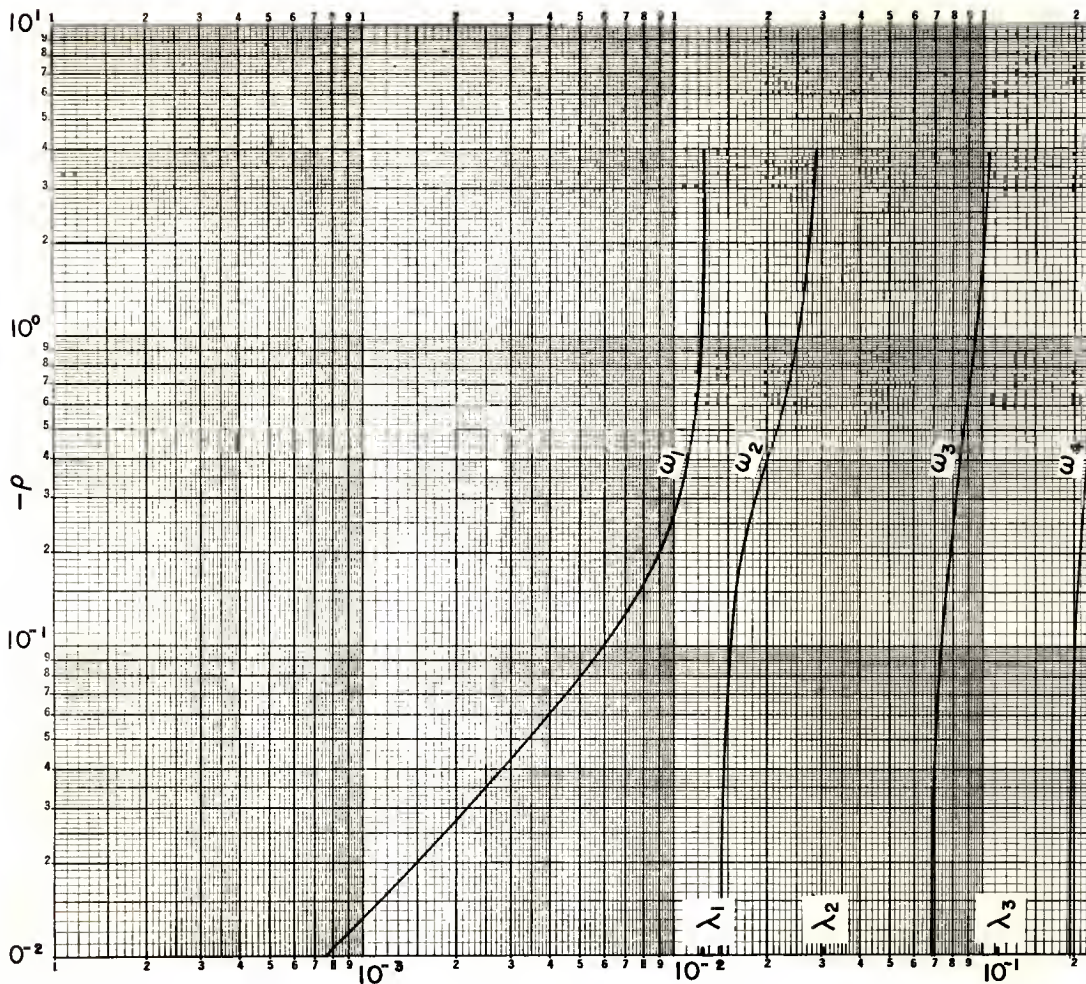
All roots are negative.

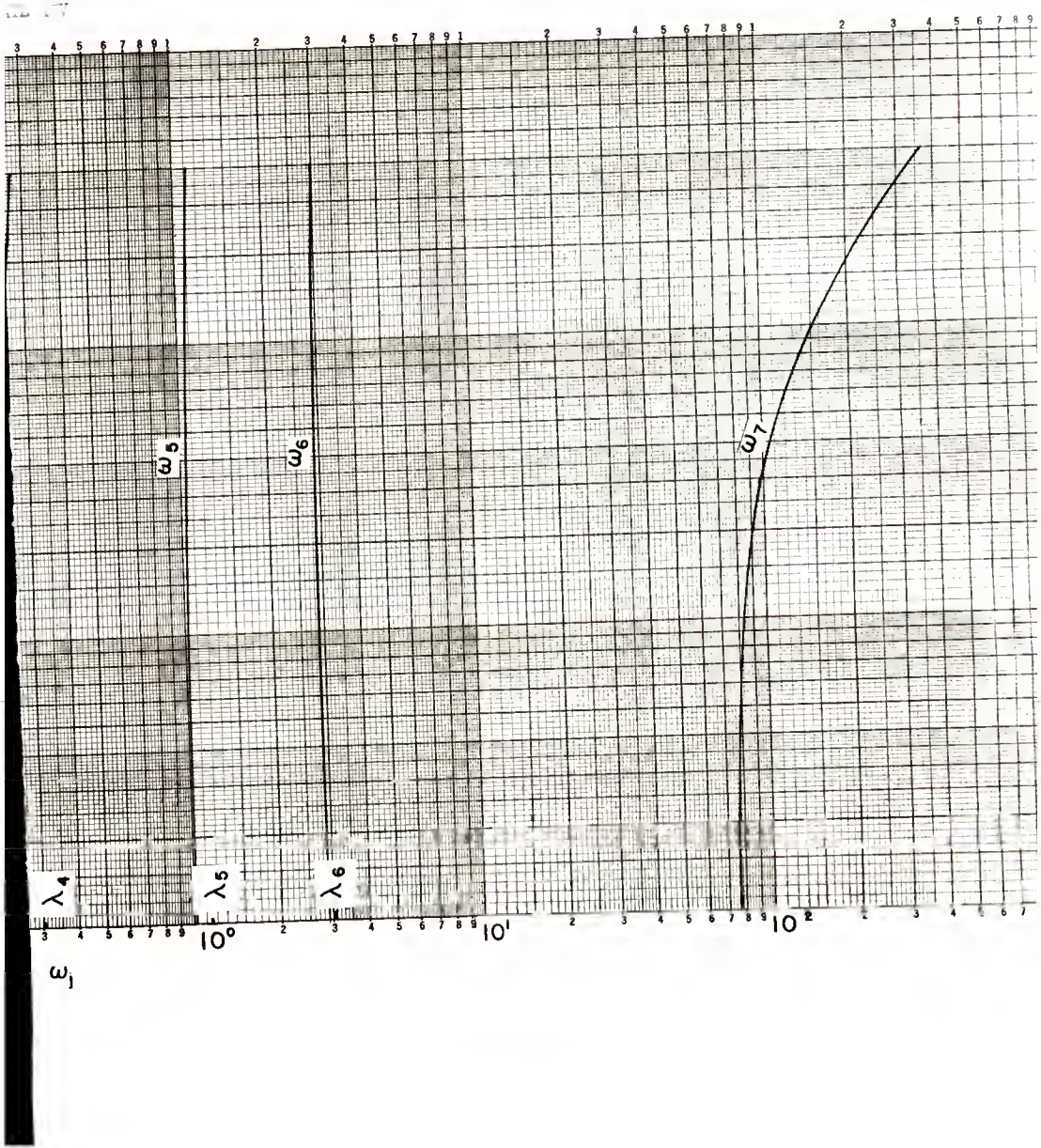
The λ 's indicate the decay constants of the neutron precursor groups.

Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keppin's delayed neutron parameters.

Data computed from Eq. (9).





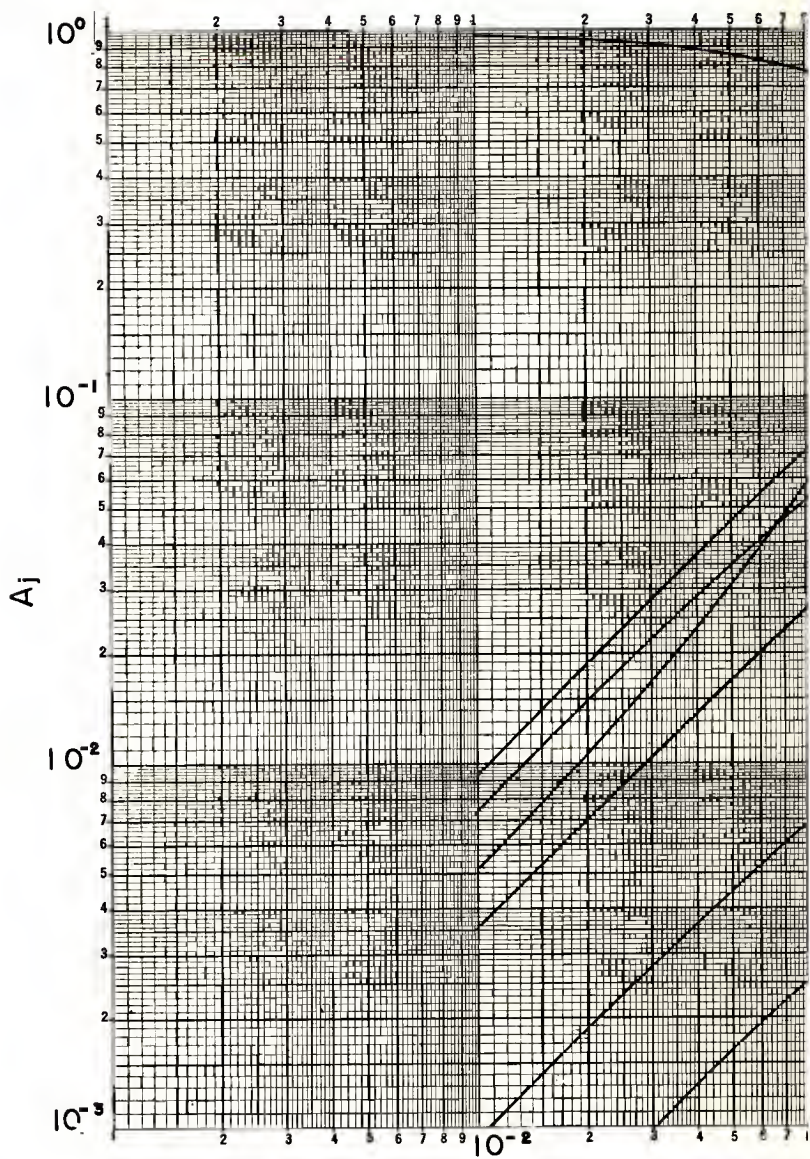
EXPLANATION OF PLATE XV

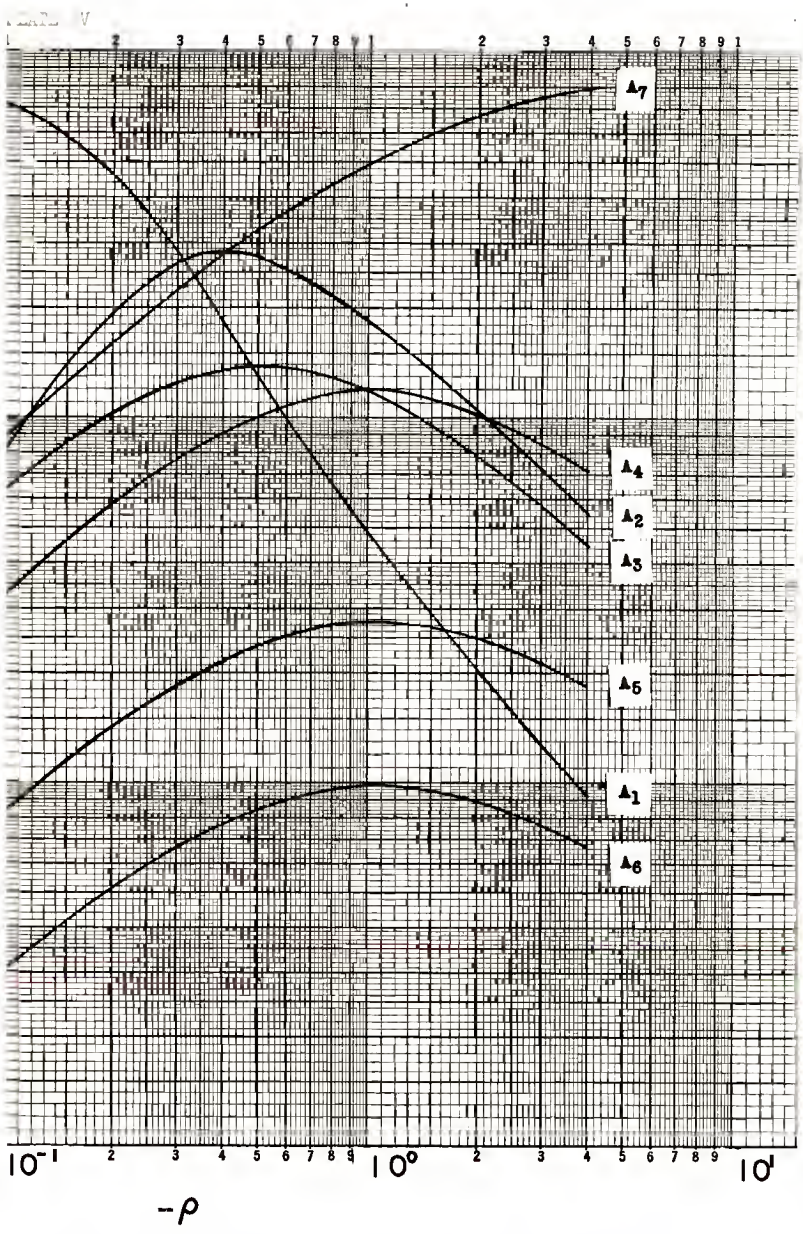
Dimensionless coefficients, A_j 's, of the flux equation vs. negative reactivity, ρ (dollars), for step changes in reactivity.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keepin's delayed neutron data.

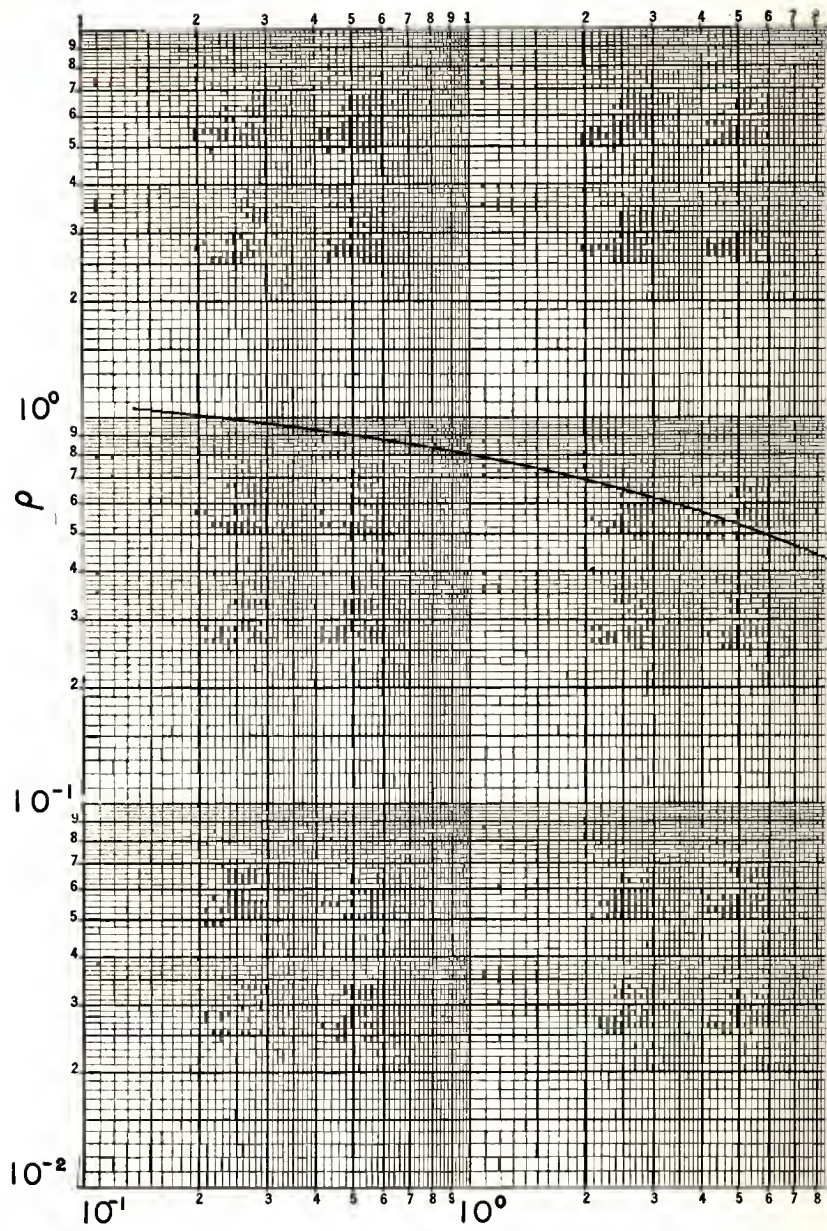
Data computed from Eq. (12).



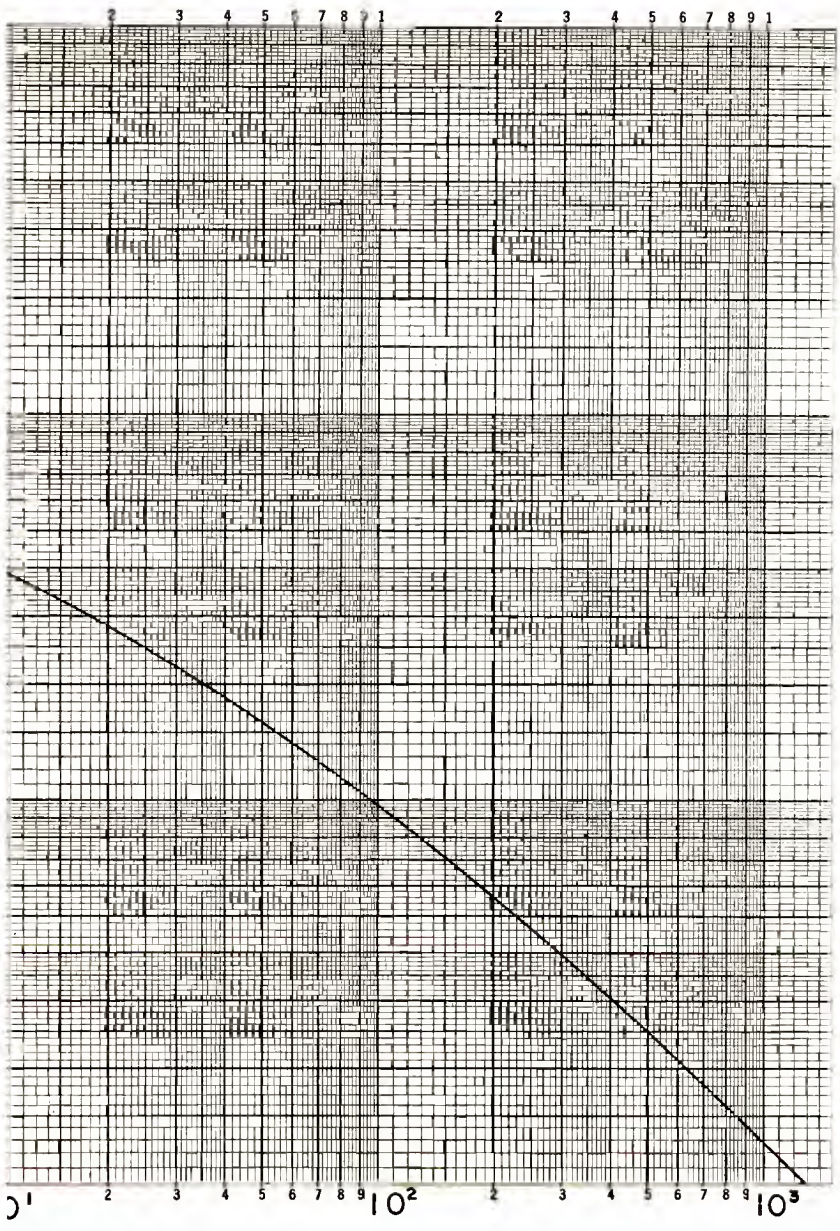


EXPLANATION OF PLATE XVI

Positive reactivity, ρ (dollars), vs.
positive asymptotic period, τ (seconds).
Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$
seconds.
Keepin's delayed neutron parameters.
Data computed from Eq. (9).



LINE 37



τ

EXPLANATION OF PLATE XVII

Positive reactivity, ρ (dollars), vs. roots, ω 's (seconds⁻¹), of the characteristic equation for step changes of reactivity.

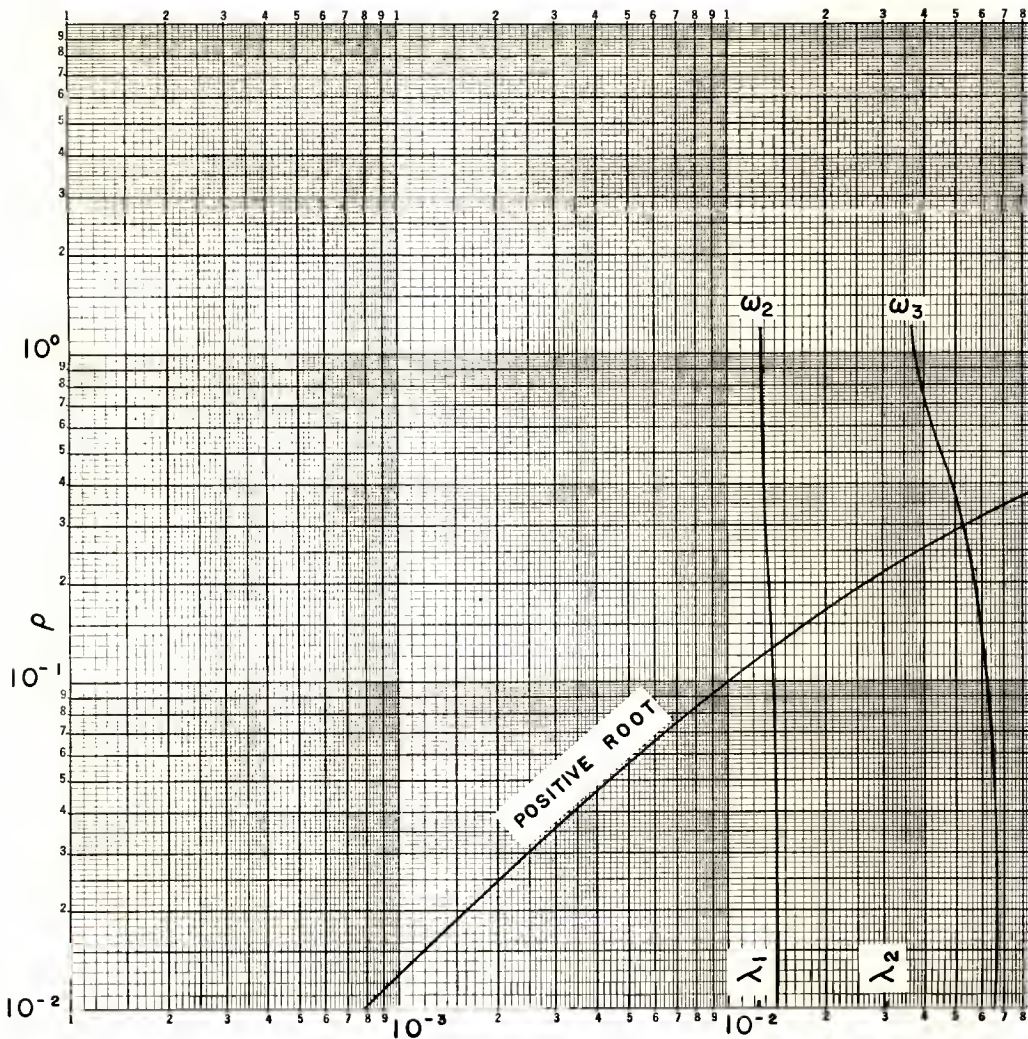
All roots are negative except the indicated positive root.

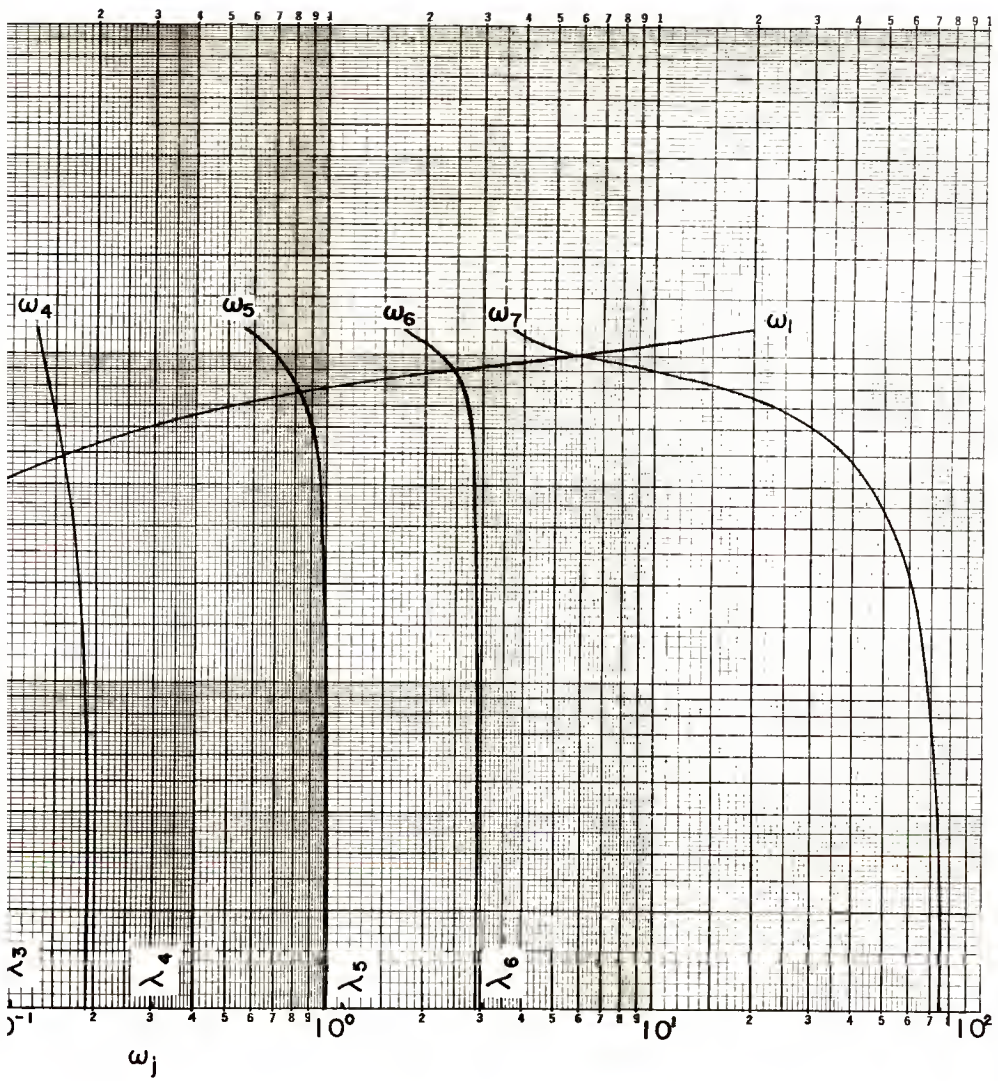
The prompt neutron lifetime, $\ell = 8.0 \times 10^{-5}$ seconds.

The λ 's are the decay constants for the neutron precursor groups.

Keepin's delayed neutron parameters.

Data computed from Eq. (9).





EXPLANATION OF PLATE XVIII

Dimensionless coefficients, A_j 's, of the flux equation
vs. $\rho_{\text{effective}}$ negative reactivity, ρ (dollars), for step
changes in reactivity.

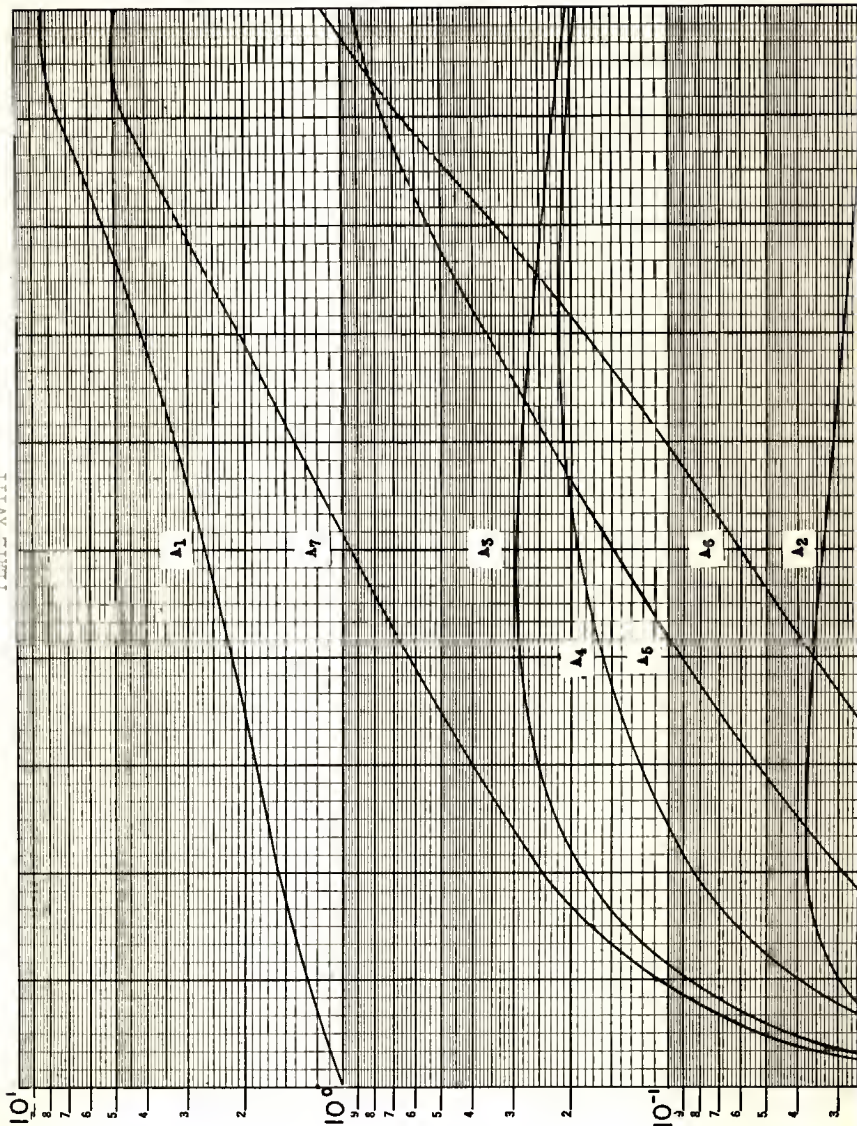
Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keypin's delayed neutron parameters.

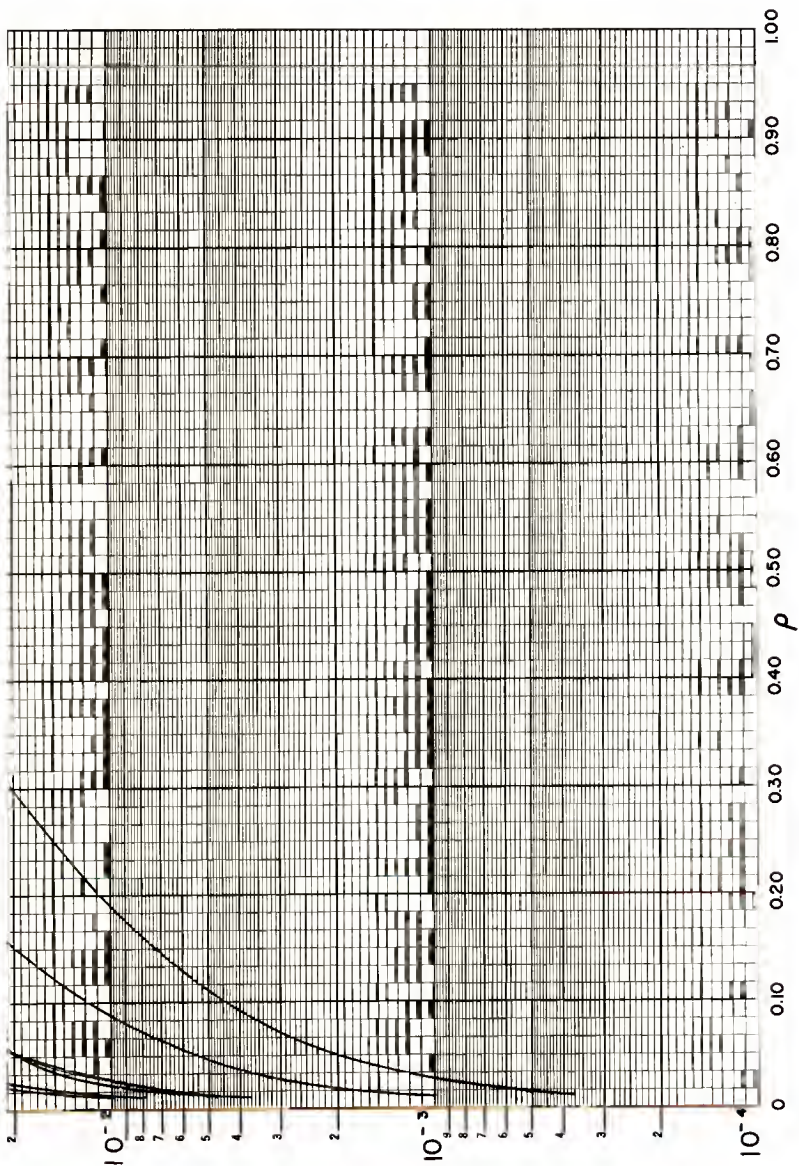
Data computed from Eq. (12).

A

PLATE XVIII



A1



EXPLANATION OF PLATE XIX

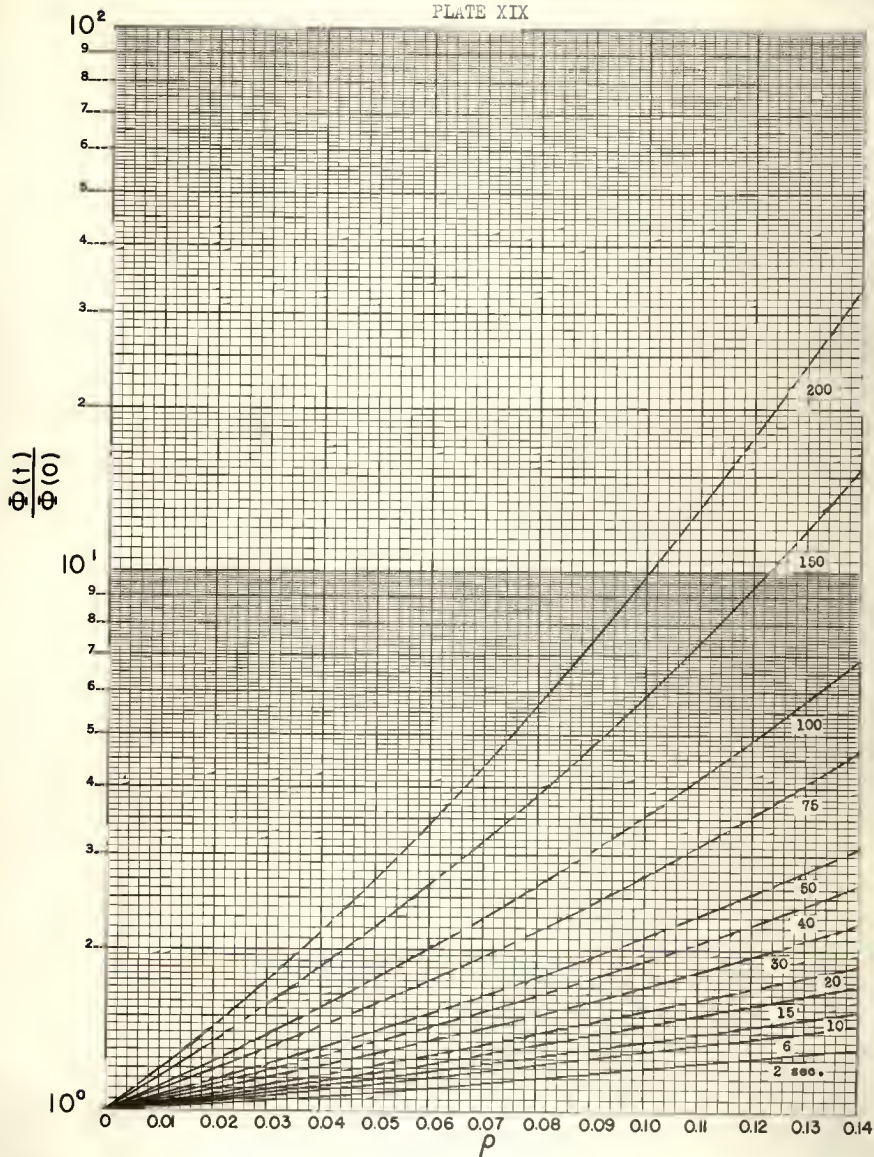
Flux ratio, $\phi(t)/\phi(0)$, vs. step changes of positive reactivity, ρ (dollars), for various times, t (seconds), after the reactivity step.

The prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Keepin's delayed neutron data.

Data computed from Eq. (11).

PLATE XIX



Beta Parameter

The total fraction of delayed neutrons has been shown to vary for different fissionable materials. Keepin (12) published a complete list of the various fissionable materials and their respective fractions of delayed neutrons. Even with the same fissionable material, the effectiveness of the delayed neutrons to promote fission has been found to differ from reactor to reactor, depending upon the size and configuration of each reactor.

Since effective delayed neutron fractions are generally difficult to obtain, it was considered important to evaluate the dependence of the kinetic equations upon the total fraction of delayed neutrons. The ratio, $\phi(t)/\phi(0)$, was calculated for various times after a step insertion of reactivity using different total fractions of delayed neutrons. The total fractions of delayed neutrons, β 's, used in the calculations varied from 0.005 to 0.008, therefore the β 's used cover the range of effective fractions used in most thermal reactors that utilize U^{235} as the primary fuel. The study included several negative and positive reactivities (see Appendix - F).

The same computations, as described above, were recalculated using a constant reactivity of -0.0748 per cent. A plot of the ratio, $\phi(t)/\phi(0)$, vs. the total fraction of delayed neutrons for two different times was made (see PLATE XX). Also shown on the same plot are the flux ratios as a function of β when the reactivity insertion was -0.10. It was found that the computed flux ratios for a particular reactivity (ρ) insertion

EXPLANATION OF PLATE XX

Flux ratio, $\phi(t)/\phi(0)$, vs. the total fraction of delayed neutrons, β .

Flux ratio evaluated at two time points, 10 seconds and 75 seconds after a step change in reactivity.

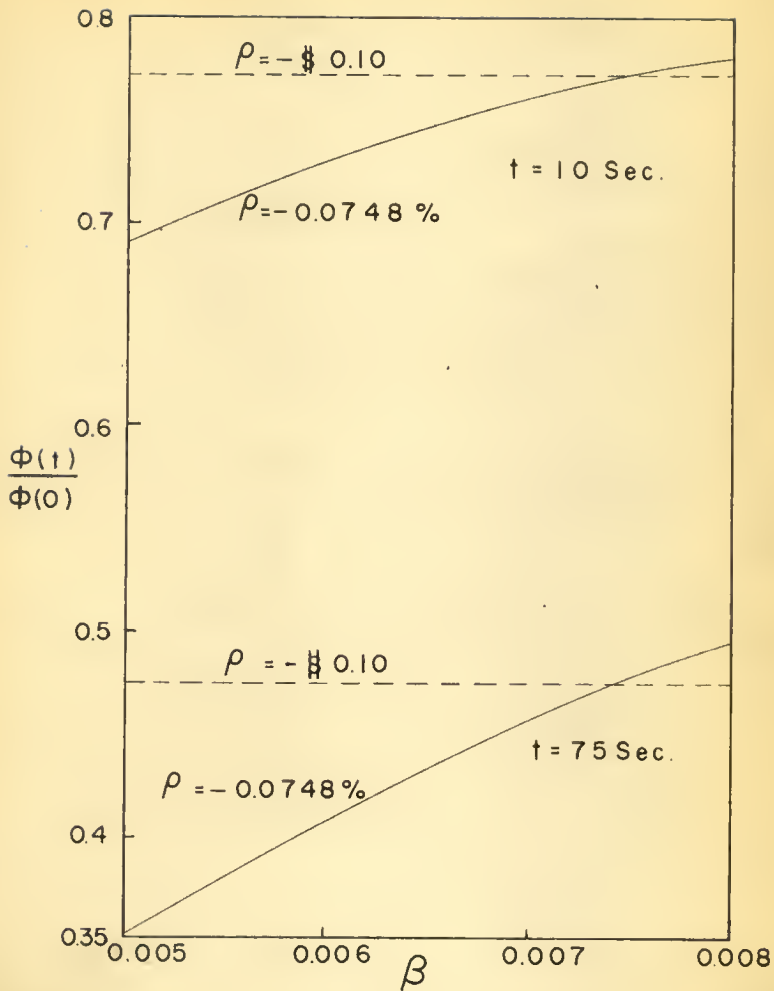
— Reactivity was -0.0748 per cent.

--- Reactivity was -0.10 , as defined by each β .

Hughes' delayed neutron parameters.

Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ sec.

PLATE



was essentially independent of β , whereas the same was not true when the reactivity insertion was given in units of per cent.

Decay Constant Parameter

The decay constants of the neutron precursor groups were measured experimentally by Keepin, who quoted the uncertainty of each of the decay constants. In this paper an evaluation was to determine the effect of varying each decay constant within the experimental uncertainty as reported by Keepin. A set of computed data was made consisting of the ratio, $\phi(t)/\phi(0)$, for various times after step insertions of reactivity using each decay constant modified to its highest value of uncertainty (see Appendix - F).

Prompt Neutron Lifetime Parameter

Since the prompt neutron lifetime varies considerably from reactor to reactor, it is important to compute the effectiveness of the lifetime on the kinetic equations. For thermal reactors, lifetimes are known to vary from approximately 3.0×10^{-5} seconds for certain swimming pool type reactors to approximately 1.0×10^{-3} seconds for large graphite piles.

The flux ratio equation was solved for both positive insertions of reactivity and negative insertions of reactivity using various values of the prompt neutron lifetime. The range of lifetimes investigated was from 1×10^{-5} seconds to 1×10^{-3}

seconds. Plots were made of the flux ratio as a function of time after a step insertion of reactivity for several reactivities and several lifetimes (see PLATES XXI and XXII). The plots indicated that the lifetime parameter is not an important parameter of the kinetic equations except when considering large step changes of positive reactivity, $\rho > \beta 0.30$.

Three Groups of Delayed Neutrons

The reactor kinetic equations were solved using three groups of delayed neutrons. The resulting characteristic equation was programmed and the roots were computed by the IBM-650. It was demonstrated that the six delayed neutron group kinetic equations could be reasonably approximated with three groups of delayed neutrons.

It is difficult to obtain a set of three group constants analytically, either experimentally or theoretically, which will reasonably approximate the kinetic equations using six delayed neutron groups. Therefore, a trial and error method was used to determine the three group constants. Flux ratios had been computed for step reactivity insertions using the six group constants. The three group constants were then obtained such that the flux ratios as given by the three group constants agreed with those as given by the six group constants (see PLATES XXIII and XXIV). A set of neutron constants determined in this manner is given in Table 4.

EXPLANATION OF PLATE XXI

Flux ratio, $\phi(t)/\phi(0)$, vs. time (seconds) after
a negative step change of reactivity, (dollars).

Hughes' delayed neutron parameters.

Two different lifetimes, $\lambda = 1 \times 10^{-5}$ seconds.

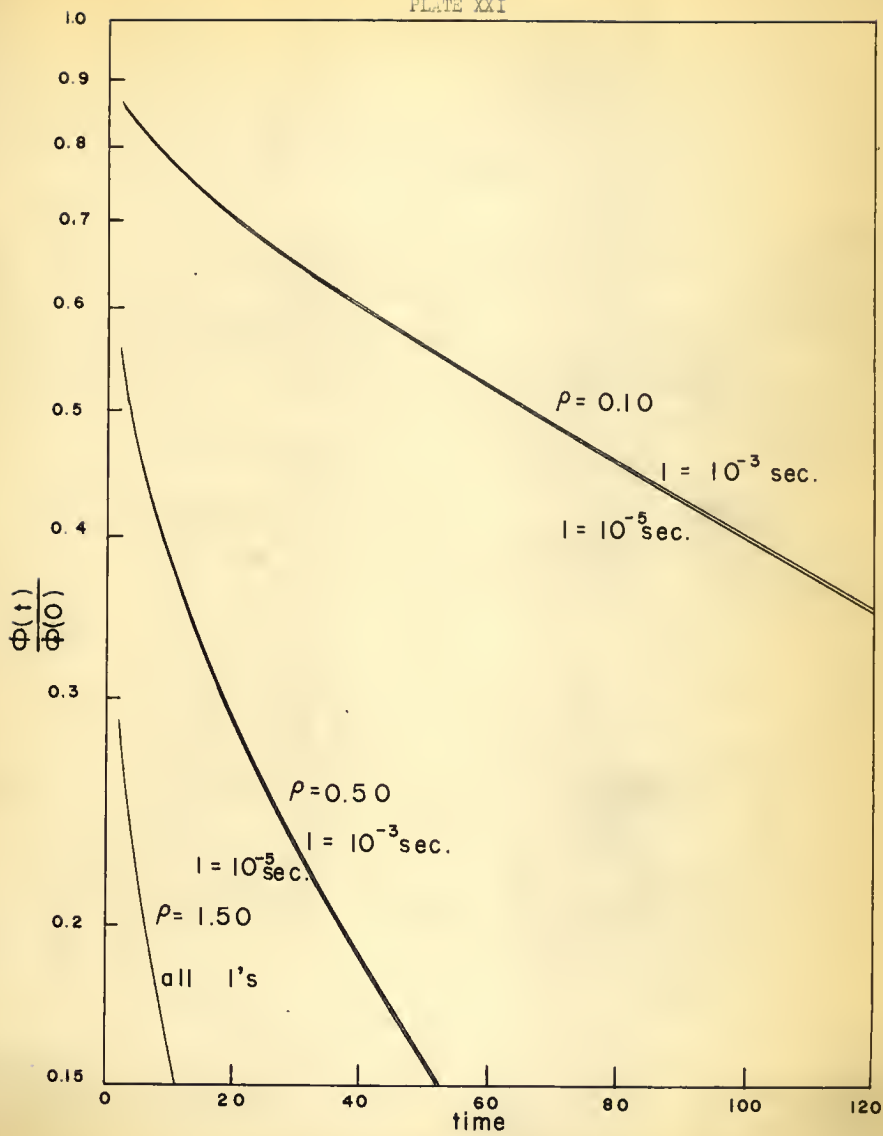
$\lambda = 1 \times 10^{-3}$ seconds.

Three different reactivities, $\rho = -\$0.10$

$\rho = -\$0.50$

$\rho = -\$1.50$

PLATE XXI



EXPLANATION OF PLATE XXII

Flux ratio, $\phi(t)/\phi(0)$, vs. time (seconds) after a positive step change in reactivity, ρ (dollars).

Hughes' delayed neutron parameters.

Two different lifetimes are shown, $\lambda = 1 \times 10^{-5}$ seconds.

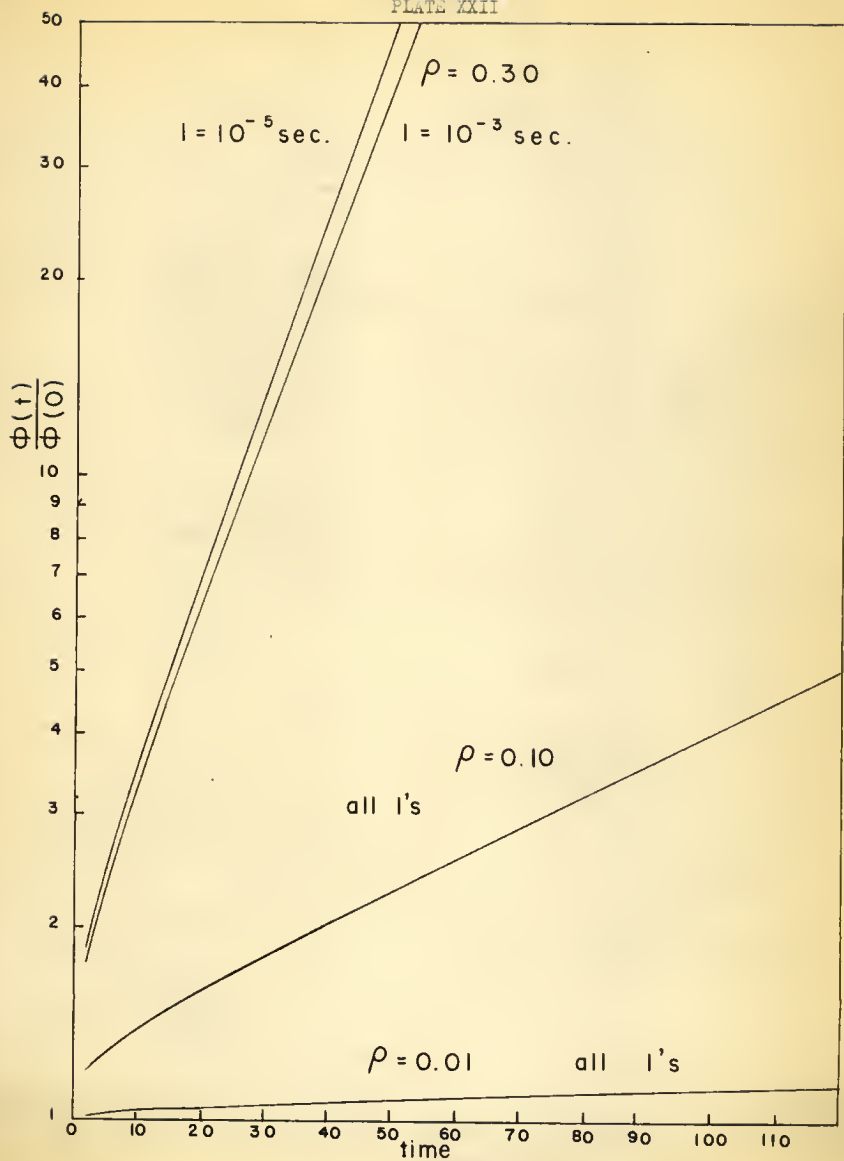
$\lambda = 1 \times 10^{-3}$ seconds.

Three different reactivities, $\rho = \$0.01$

$\rho = \$0.10$

$\rho = \$0.30$

PLATE XXII



EXPLANATION OF PLATE XXIII

Flux ratio, $\phi(t)/\phi(0)$, vs. time (seconds) after a negative step change in reactivity of $-\beta 0.10$.

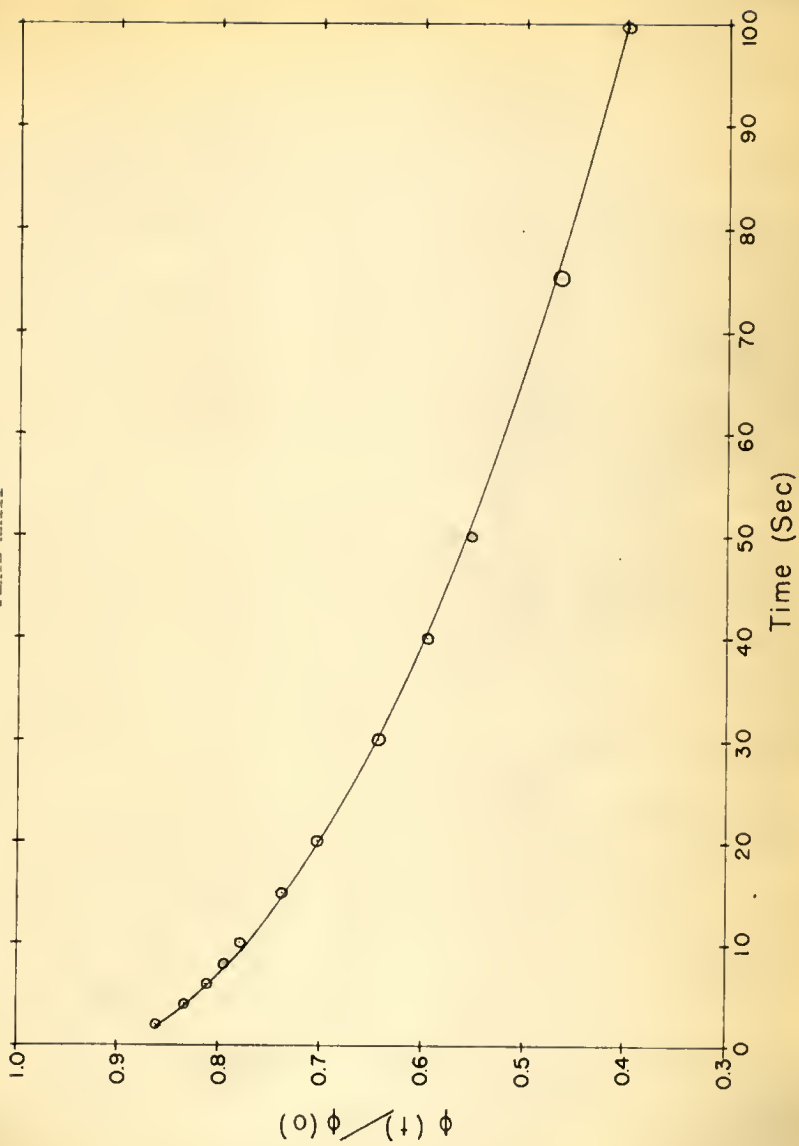
— Theoretically predicted flux ratio using the six group kinetics model.

Hughes' delayed neutron parameters.

⊙ Theoretically predicted flux ratio using the three group kinetics model.

Three group parameters are given in Table 4.

PLATE XXIII



EXPLANATION OF PLATE XXIV

Flux ratio, $\phi(t)/\phi(0)$, vs. time (seconds) after a negative step change in reactivity of -1.42.

— Theoretically predicted flux ratio using the six group kinetics model.

○ Theoretically predicted flux ratio using the three group kinetics model.

Hughes' delayed neutron parameters.

Three group parameters are given in Table 4.

PLATE XXIV

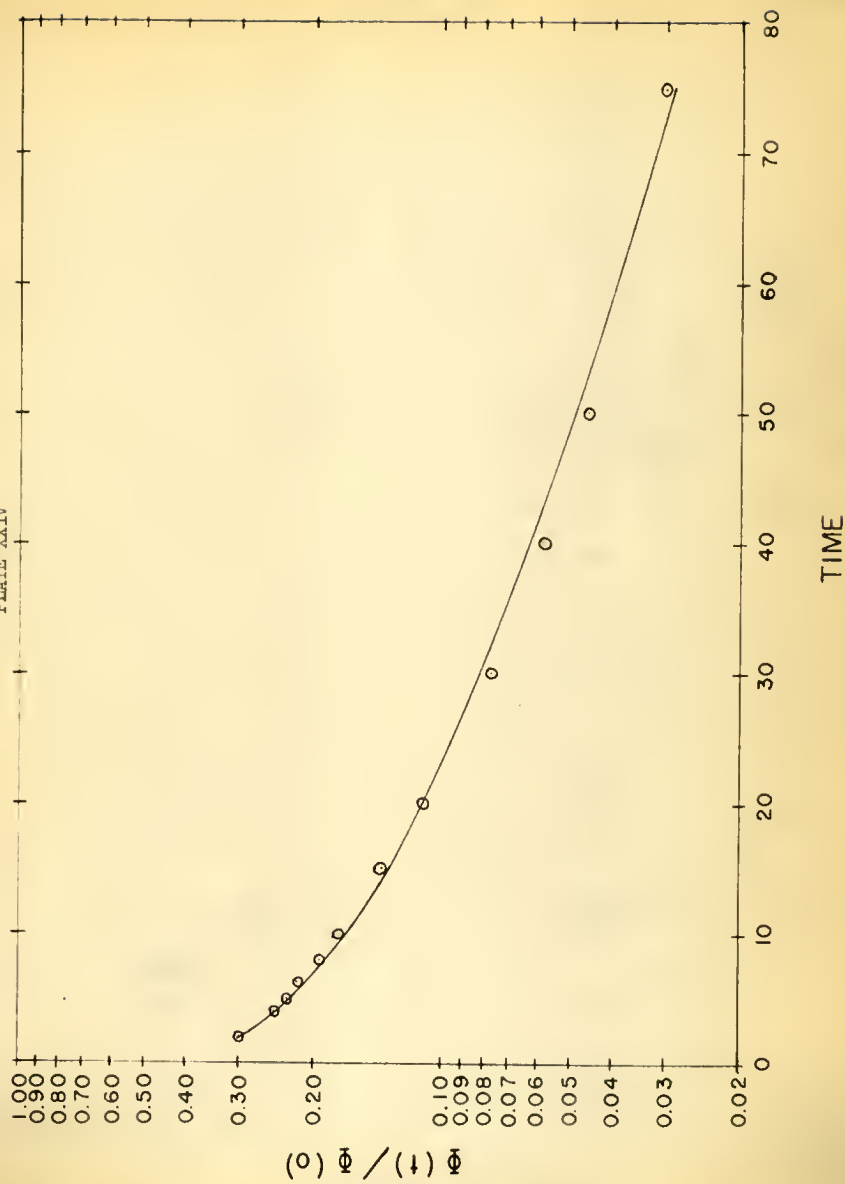


Table 4. Empirically determined three delayed neutron group constants.

Group	λ_i (Sec. ⁻¹)	β_i
1	0.79	0.0033
2	0.091	0.00349
3	0.0117	0.00069

For each reactivity investigated using the three group analysis, the comparisons with the six group analysis gave similar flux ratio deviations as a function of time. At two seconds after the step change in reactivity, the two models, the three group and the six group, gave almost the same flux ratio. For times between six seconds and 20 seconds, the three group model gave slightly higher flux ratios. For times between 20 seconds and 100 seconds, the six group model gave slightly higher flux ratios. For times larger than 100 seconds, the three groups model gave higher flux ratios. The difference between the flux ratios as given by the two models differed significantly for times larger than 150 seconds. As the step input of reactivity increased, the difference in the flux ratios, as given by the two models, increased. The largest reactivity used was $-\$1.42$ (see PLATE XXIV).

The six group code required approximately 100 seconds of IBM-650 time to solve the characteristic equation and calculate flux ratios for 15 time points. The three group code required approximately 25 seconds to solve the characteristic equation

and calculate the flux ratios for 15 time points. Thus, if it becomes necessary to perform a kinetic analysis using certain additional reactivity parameters, the three group analysis will greatly simplify the calculations.

EXPERIMENTAL MEASUREMENTS

General

The experimental data were obtained with the TRIGA reactor located at the John Jay Hopkins Laboratory, San Diego, California and with the Argonaut reactor located at Argonne National Laboratory. The TRIGA data were obtained by the author and the Argonaut data were forwarded to this author by Dr. Bill Sturm, Manager of the Argonaut reactor.

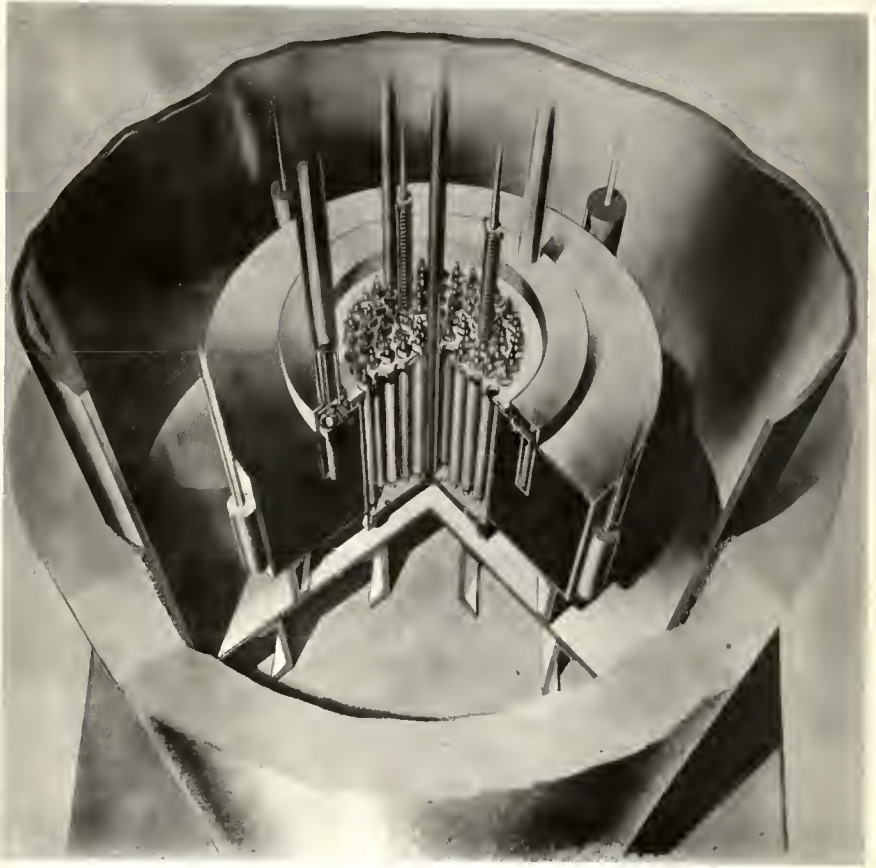
Experimental System

Briefly, the TRIGA reactor consisted of a cylindrical core containing a lattice of cylindrical fuel elements surrounded by a graphite reflector and located at the bottom of a 20 foot aluminum tank (see PLATES I and XXV). The fuel elements contained fissionable material in the form of an alloy of uranium-zirconium hydride composed of eight wt. per cent uranium enriched to 20 per cent in U^{235} . The reactor contained four control rods. The control rods used in the reactivity measurements were the shim rod and the regulating rod. The positions of the rods are shown in PLATE XXVII. More descriptive information of the TRIGA reactor may be found in the references (19), (21), and

EXPLANATION OF PLATE XXV

View of the core of the TRIGA reactor showing the positions of the ionization chambers located just outside the core reflector.

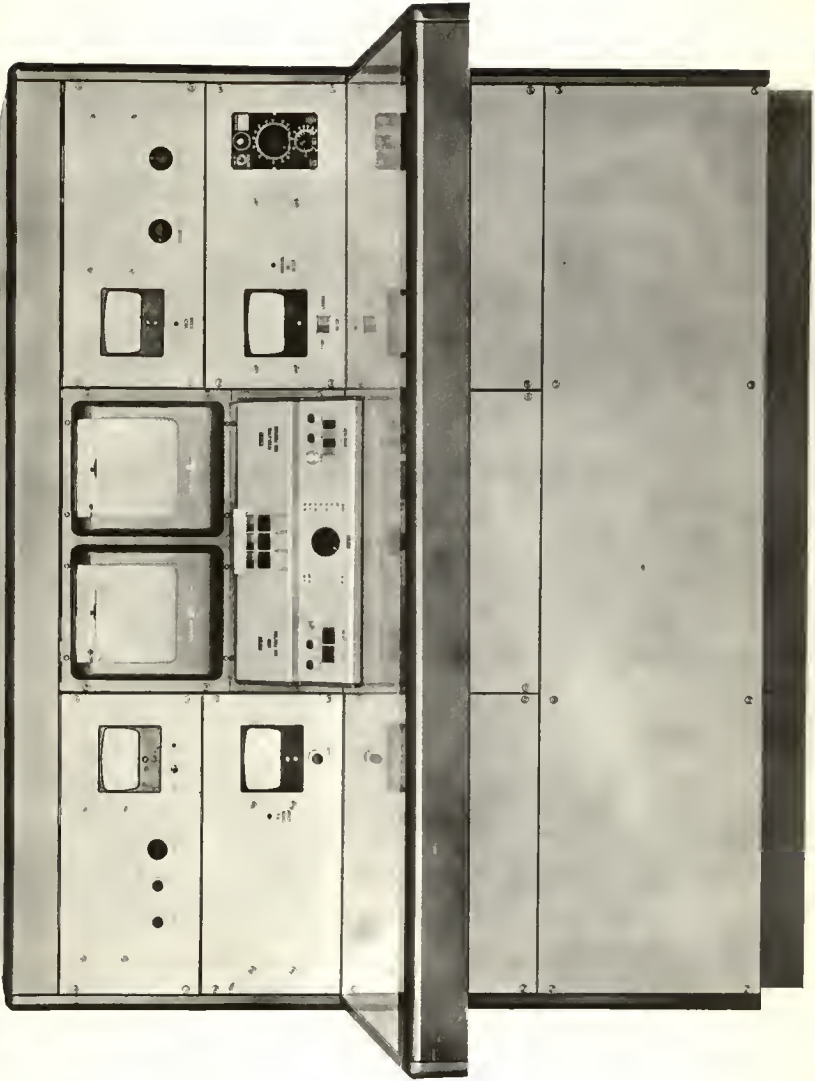
PLATE XXV



EXPLANATION OF PLATE XXVI

View of a TRIGA reactor console.

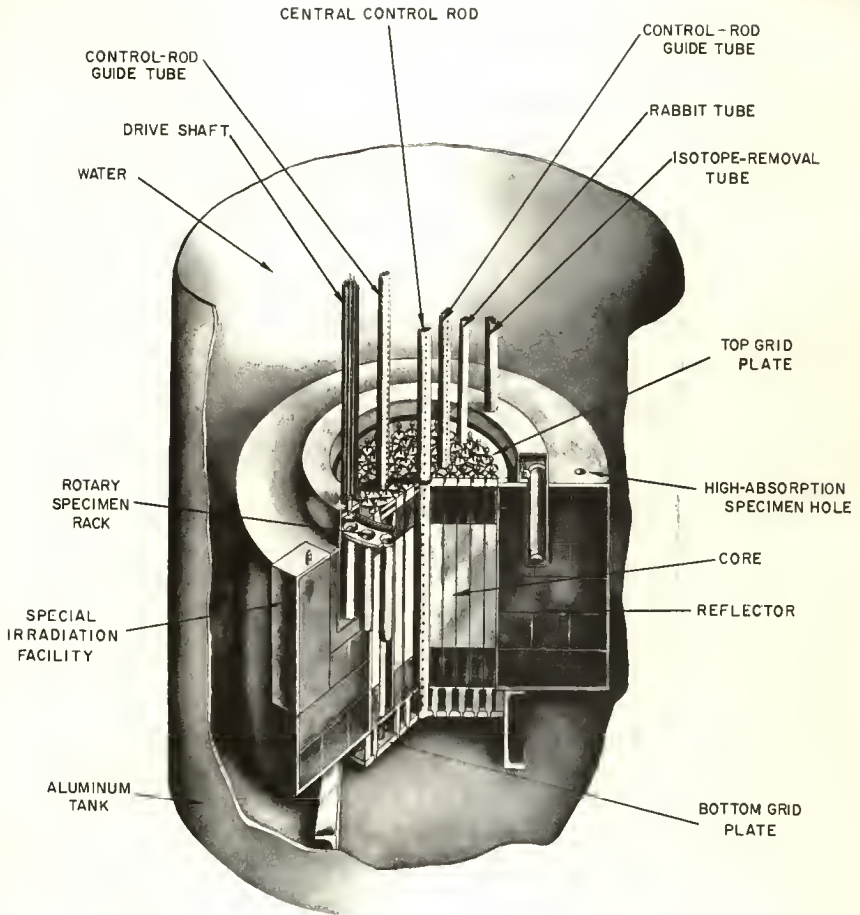
PLATE XXVI



EXPLANATION OF PLATE XXVII

A sketch of the TRIGA reactor showing the
basic core components.

PLATE XXVII



REACTOR CORE AND REFLECTOR ASSEMBLY

(22).

The reactivity insertion was made by changing the position of one of the two control rods. The control rods were composed of boron-carbide canned in aluminum tubes. The control guide tubes may be seen in PLATE XXVIII. The rods had a $15\frac{1}{2}$ inch travel from their "full out" position to "full in". During the rod drop experiments or during a scram, the rods fell under the force of gravity. The time required for the rods to fall from "full out" to "full in" was approximately 0.35 seconds.

The neutron flux level was measured with the use of two Westinghouse WL-6377 compensated ion chambers. The ion chamber sensitivity was 4×10^{-4} amps/nv and were electrically compensated having a compensated gamma sensitivity of about 3×10^{-13} amps/roentgen/hr. The compensated ion chambers are shown in PLATE I.

The current from the compensated ion chambers was monitored with two Keithly micro-micro ammeters. The ammeters are shown in PLATE XXIX just above the reactor operator's head. The signal from the micro-micro ammeters was recorded by a Midwestern model 608 direct-recording, oscillograph. (See PLATE XXX).

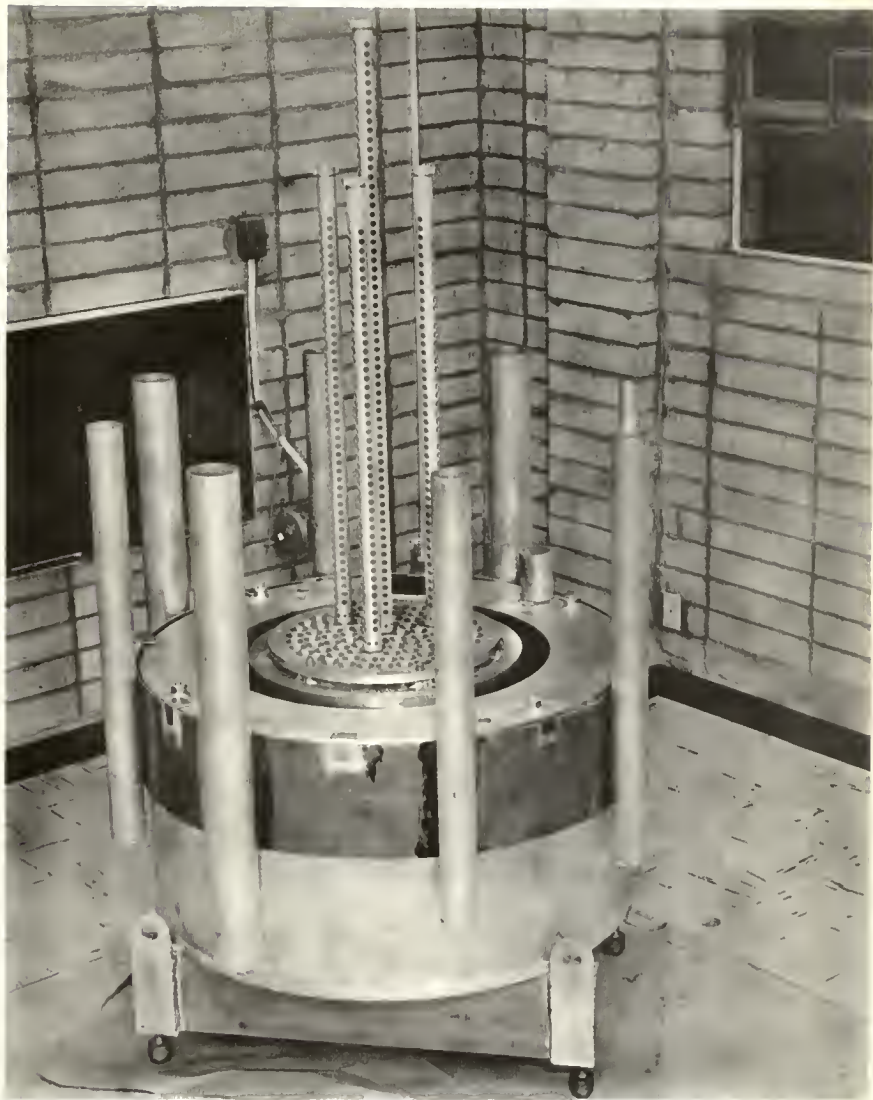
EXPLANATION OF PLATE XXVIII

View of the TRIGA core before positioned into the reactor tank.

Metal tubes at the periphery will hold the ionization chambers.

The perforated tubes are the control rod guide tubes.

PLATE XXVIII

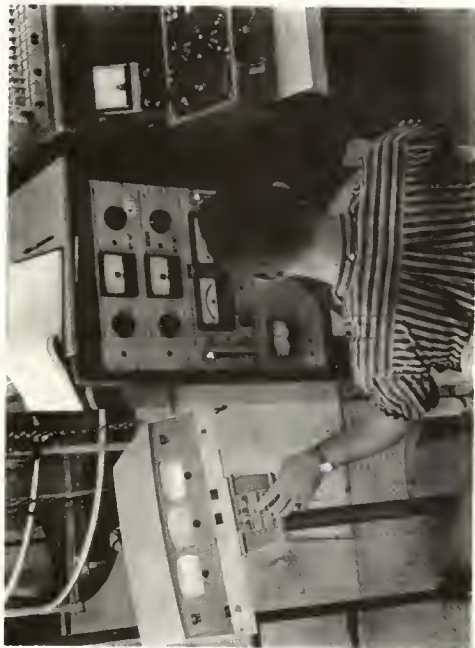


EXPLANATION OF PLATE XXIX

View of the simplified TRIGA reactor console (at left).

The two Keithly micro-micro ammeters are located just above the head of the reactor operator (at right).

PLATE XXIX



EXPLANATION OF PLATE XXX

General layout of the TRIGA building.
Midwestern direct-recording oscillograph is
located on the table (center) and is being
operated by the man who is kneeling beside it.

PLATE XXX



EXPERIMENTAL PROCEDURE

General

Two experimental methods were used to measure the reactivity worth of a control rod. The first method described is by far more commonly used in the nuclear field for reactivity measurement, and this procedure is known as the positive period method. The second method is known as the rod drop method.

The positive period method consisted of several well defined procedural steps. First, the reactor was brought to a critical position at a low power.¹ The reactor was maintained at the critical position for approximately eight minutes to insure that the criterion for delayed criticality² was satisfied. A positive reactivity change was made by partially withdrawing one of the control rods. The neutron population in the core increased as a function of time. One hundred seconds after the change in reactivity the time function of the neutron density was approximately represented by a single exponential function. The time constant of the exponential function was the reciprocal of the stable reactor period. The time required for the neutron density to increase by a factor of "e" was measured and that

¹The TRIGA reactor has a significant temperature coefficient of reactivity and thus any temperature change during reactivity measurements would have surely invalidated the measurements. Other than the heating effect, power was not a factor in the reactivity measurements.

²Delayed criticality is defined as the state of a nuclear reactor in which all the neutrons resulting from fission, both delayed and prompt, are required to maintain a chain reaction.

time was the stable reactor period for that particular reactivity change. From an analytical relationship between the stable reactor period and the magnitude of the step change in reactivity, the reactivity was established for the particular experiment.

Instrumentation Calibration

Frequent checks were made to insure that the instruments used were calibrated and functioning properly. The linearity of the compensated ionization chambers was checked numerous times to insure that the current output of the chambers was proportional to the reactor power level. The micro-micro ammeters were checked several times during each run to insure that the proper zero position was being used. The Midwestern recorder was calibrated several times to insure that the proper calibration was maintained.

It was assumed, in the TRIGA reactor, that the thermal power level was proportional to the neutron density at the position of the detector. The power level calibration was obtained by measuring the average reactor tank water temperature as a function of time while the reactor was operating at a steady power. The temperature measurements were made at several power levels. The calibration indicated that the current output of the ionization chamber was truly proportional to the power level regardless of the detector position.

During a particular run, the micro-micro ammeter measured neutron flux levels over several decades of power. Each decade of the micro-micro ammeter required a slightly different zero

position, thus when the micro-micro ammeter was switched from one decade to a lower value a new zero position was made by an adjustment on the ammeter.

The Midwestern recorder operated on a galvanometer principle, i.e., the current supplied to the recorder galvanometer was proportional to the current fed to the ammeter. The galvanometer reflected ultraviolet light onto a moving photosensitive recording paper. The recorder was calibrated by imposing a false signal, of known value, from the micro-micro ammeter to the galvanometer. The deflection of the false signal line on the recorder paper was then measured as a function of the magnitude of the false signal (see PLATE XXXII).

Positive Period Measurements

The two control rods that were calibrated with the positive period method were the shim rod and the regulating rod (see PLATE XXXI). The safety rod calibration was furnished by General Atomic personnel. The shim and regulating rod calibrations were made using this safety rod calibration.

The reactor was brought to a delayed critical position with the safety rod out of the core, the shim rod completely in the core, and the regulating rod at a position necessary for criticality.¹ The positions of the control rods at the critical position were recorded in the reactor log book. The reactor was

¹The power level at the critical position was generally around 20 watts.

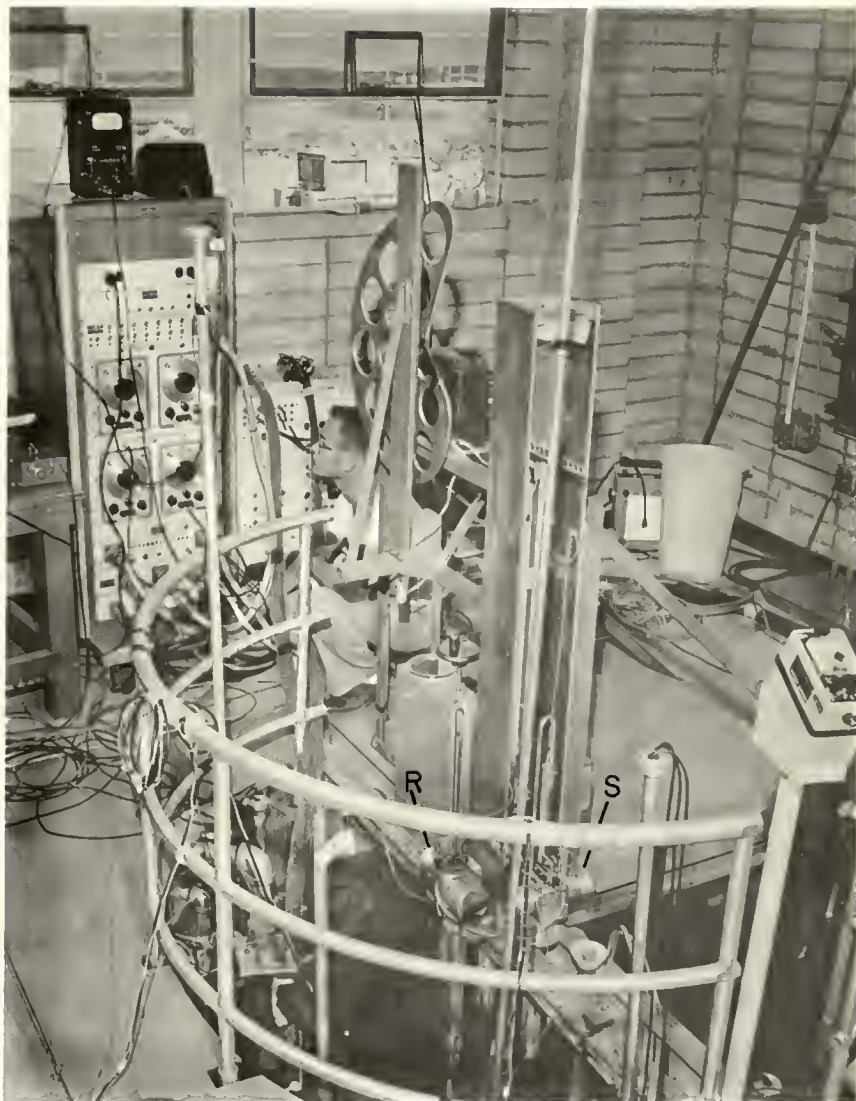
EXPLANATION OF PLATE XXXI

View of the top of the TRIGA reactor showing the control rod drive motors and the transient rod mechanism.

S - shim rod drive mechanism.

R - regulating rod drive mechanism.

PLATE XXXI



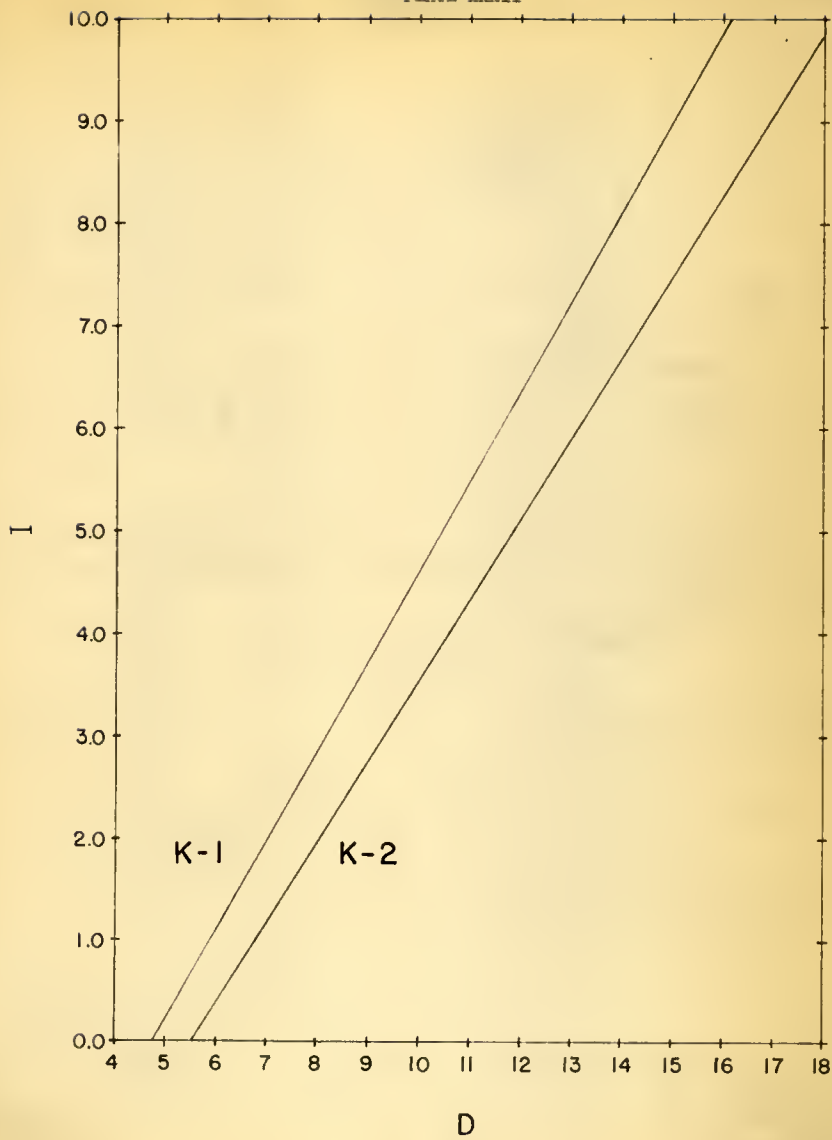
EXPLANATION OF PLATE XXXII

Calibration curve for the Midwestern Model 608
direct-recording oscillograph.

Current, I (arbitrary units) vs. deflection, D
(cm.), of signal line away from the reference
line.

K-1 and K-2 are recording channels.

PLATE XXXII



left in the critical state for approximately eight minutes to insure that a delayed criticality had been obtained. The shim rod was then withdrawn from the core a short distance to make the reactor super-critical. The new position of the shim rod was recorded in the reactor log book. The reactor was left in the supercritical state for at least 100 seconds to insure that the increase in power as a function of time was exponential. A stop watch was used to measure the time required for the reactor power to increase by a factor of 1.5. The time measured was then divided by the natural logarithm of 1.5, 0.4055, to obtain the stable reactor period. Several such measurements were made with each reactivity change. The reactor was brought back to the critical position at the original power level by means of inserting the regulating rod into the core. A technique was used to initially overcompensate the reactivity of the reactor with the regulating rod. It consisted of inserting the regulating rod into the core by an amount greater than necessary to bring the reactor to a critical position at the original power. The power level was allowed to undershoot the original power and then the regulating rod was withdrawn from the core in small steps to obtain the desired criticality. The purpose of the undershoot in power level was to "balance out" the effects of the delayed neutrons such as to obtain a delayed criticality in a much shorter time. The new position of the regulating rod was recorded in the reactor log book. The procedure was repeated by withdrawing the shim rod another short distance. The stepwise

measurements were continued until the regulating rod was almost completely inserted into the reactor core. At that point, the safety rod was slowly inserted into the core while the regulating rod was being withdrawn, to insure that near criticality conditions were being maintained. After the safety rod was fully inserted into the core, the new position of the regulating rod was recorded. The stepwise reactivity measurements were again continued until the shim rod and the regulating rod were completely calibrated.

A theoretical relationship between the reactivity change and the resulting stable period was used to evaluate the reactivity of each of the rod movements. The relationship (see PLATE XXXIII) was derived by General Atomic (19). The reactivity changes were produced by the shim rod, thus the reactivity of each of the rod movements was summed to obtain the integral worth of the shim rod (see PLATE XXXIV). Since the shim rod reactivities were compensated with regulating rod, similar calibration curves were plotted for the regulating rod (see PLATE XXXV).

Rod Drop Measurements

Rod drop measurements were made with two control rods, the shim and the regulating rod. When available, two linear power level channels were used for neutron density measurements. At other times, the second compensated ionization chamber was too far away from the core to accurately measure the low power used in the rod drop measurements.

EXPLANATION OF PLATE XXXIII

Positive reactivity, ρ (dollars), vs.
asymptotic period, τ (seconds).

Analytical data calculated by General Atomic
(19) using Keepin's delayed neutron data.

Total effective fraction of delayed neutrons,
0.0078.

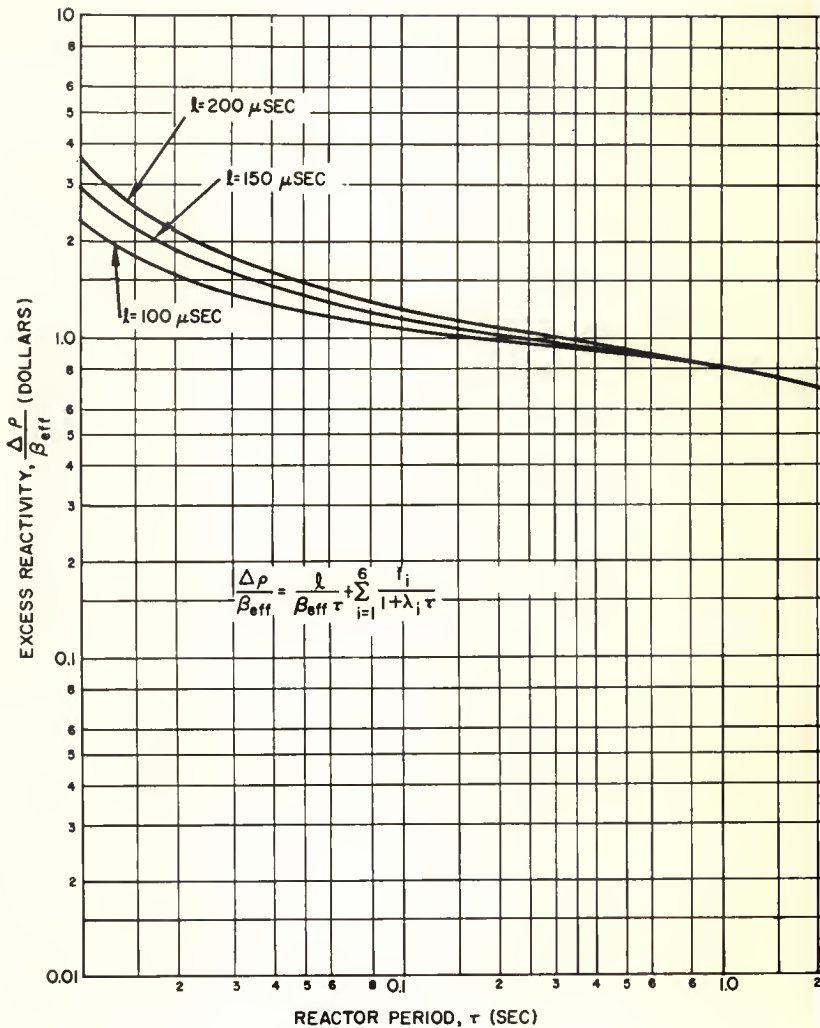
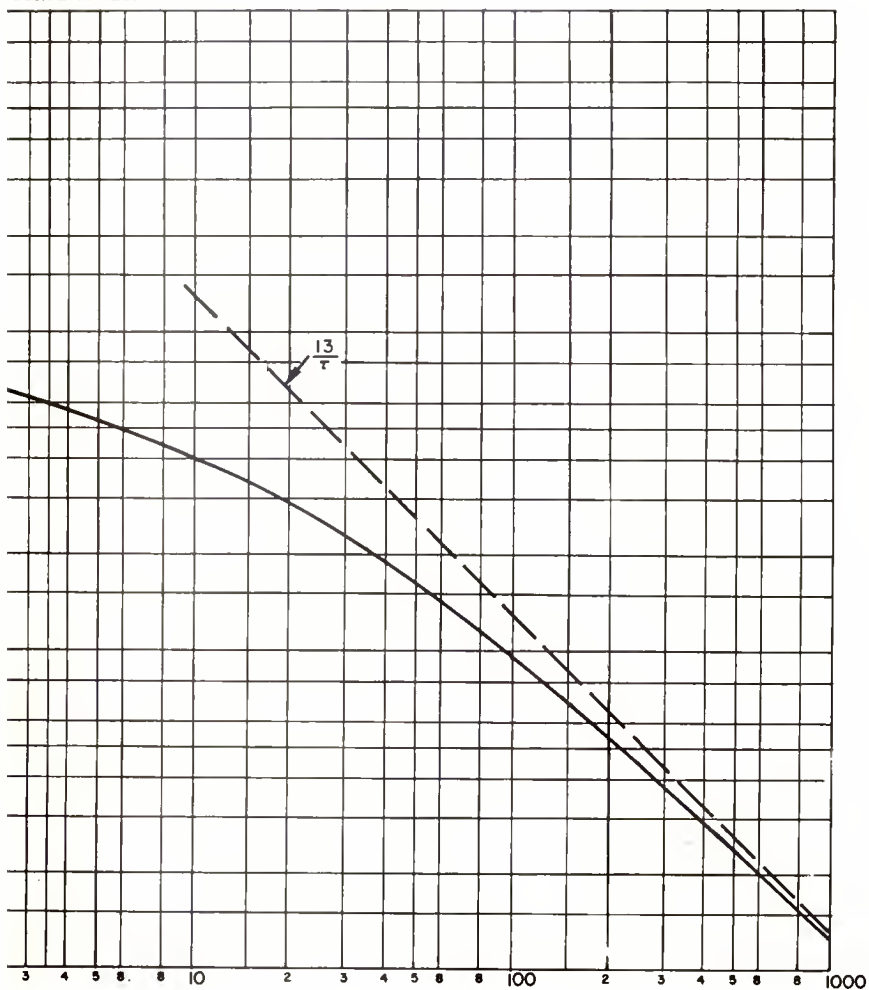


PLATE XXXIII



EXPLANATION OF PLATE XXXIV

Shim control rod calibration for the TRIGA reactor.

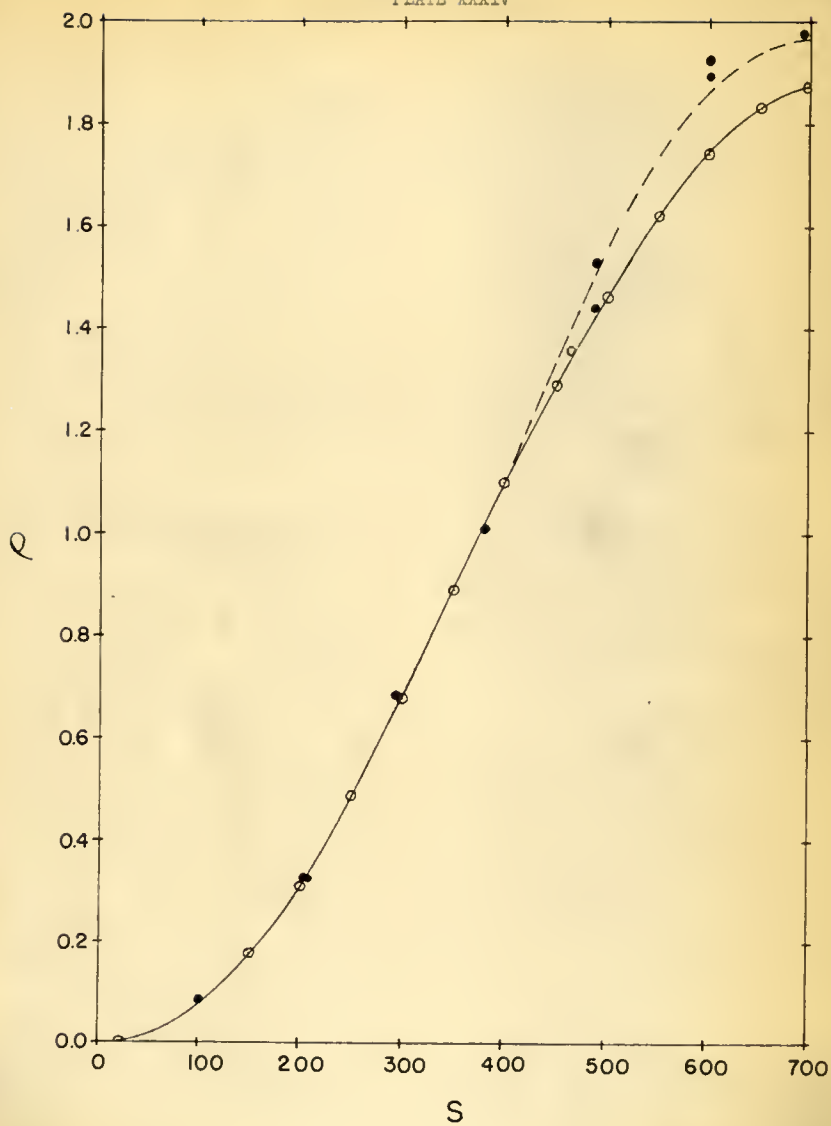
Reactivity, ρ (dollars), vs. control rod position,
S (units).

---◆--- Measured by dropping the shim rod from various
critical positions to "full in".

—○— Measured by the positive period method.

Keepin's delayed neutron parameters.

PLATE XXXIV



EXPLANATION OF PLATE XXXV

Regulating control rod calibration for the
TRIGA REACTOR.

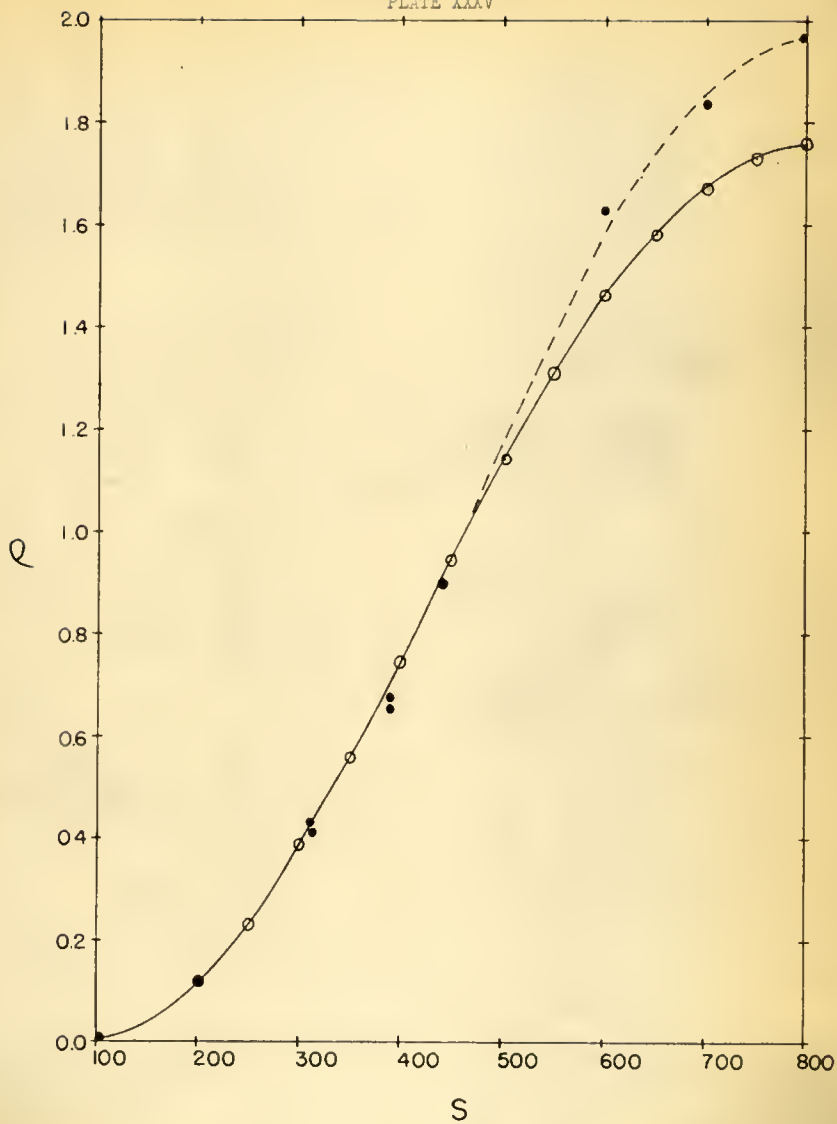
Reactivity, ρ (dollars), vs. control rod
position, S (units).

--♦-- Measured by dropping the regulating
rod from various critical positions
to "full in".

—⊙— Measured by the positive period method
using increments of reactivity.

Keepin's delayed neutron data.

PLATE XXXV



The first step of the rod drop measurements was to bring the reactor to a critical configuration with the rod to be calibrated at the position desired and with the other rods at positions necessary for criticality. The reactor was allowed to remain at the critical positions for a length of time necessary to insure that the criterion for delayed criticality had been satisfied. The Midwestern recorder was turned on to allow for sufficient warm-up time. A zero adjust check was made on the Keithly micro-micro ammeter. The magnet holding the control rod was then de-energized and the rod fell freely into the core. Since there was no means available to stop the rod between the critical position and the "full in" position, all of the rod drop measurements were made from the critical position to the "full in" position.

The scales of the Keithly micro-micro ammeter were 1×10^{-n} and 3×10^{-n} . Each time that the current from the ion chamber became too low, the decade setting on the ammeter was switched to a lower value, therefore a new zero adjust was required. The zero adjust operation was performed quickly and very little recording signal was lost during the adjusting period.

A signal proportional to the power level of the reactor was recorded by the Midwestern recorder. The power level was proportional to the deflection of the recorder signal line from a reference line. Thus, the deflection of the signal line was measured and from the calibration graph (see PLATE XXXII), the power level as a function of time was obtained. The measured power level was then normalized to unity at the critical power

level just before the rod drop, thus the normalized power level gave the ratio, $\phi(t)/\phi(0)$. From the theoretical plots (see PLATES II and III) the reactivity of the rod drop was determined. Since the value of the ratio, $\phi(t)/\phi(0)$, was measured for the various times, two seconds, four seconds, etc., it was possible to obtain the reactivity at the various times. The rod drops were analyzed with two sets of neutron data, Keepin's (13) and Hughes' (10).

After each rod drop was completed, the reactor was brought back to the initial critical power level, thus all of the rod drop measurements started with the same critical power level. A critical configuration was obtained with the rod to be calibrated in a new position and then the rod drop procedure was repeated. Rod drops were made from various critical positions of the rod that was being calibrated, until the reactivity worth of the complete rod was determined. Both the shim and the regulating rods were calibrated in the same manner as described above (see PLATES XXXIV and XXXV).

EXPERIMENTAL ANALYSIS

TRIGA Data

As outlined in the procedure, two methods of reactivity measurements were used to calibrate two control rods of the TRIGA reactor. The control rods that were calibrated were the shim rod and the regulating rod.

When the positive period method was used to measure the

reactivity worth of the shim rod, the rod was calculated to have a complete worth of \$1.87, according to the General Atomic (19) analytical data (see PLATE XXXIV). In these data General Atomic used Keepin's (13) delayed neutron data but instead of the total fraction of delayed neutron as given by Keepin, they calculated a total effective fraction of delayed neutrons. General Atomic used as a total fraction the value of 0.0078 as compared to that given by Keepin of 0.0064, but as was shown in the analytical data of this paper (see Appendix - F), there was no significant difference between the stable period relationship and reactivity when either 0.0064 or 0.0078 was used for the total fraction of delayed neutrons.

Because of the difference between the two sets of neutron constants as given by Keepin as compared to those given by Hughes, there was found to be a slightly lower complete reactivity worth of the control rod when the rod was calibrated by the positive period method using Hughes' delayed neutron data (see PLATE XXXVI). The difference between the two delayed neutron models was most significantly demonstrated when the two plots of reactivity ρ , vs. the stable period, τ , were compared (see PLATE XXXVII). The total worth of the shim rod was calculated to be \$1.75 when Hughes' delayed neutron data were used.

Similarly, the regulating rod was calibrated by the positive period method using Keepin's delayed neutron data (see PLATE XXXV) and Hughes' delayed neutron data (see PLATE XXXVIII). The regulating rod complete reactivity worth was calculated to be \$1.76 and \$1.64 using Keepin's and Hughes' data respectively.

EXPLANATION OF PLATE XXXVI

Shim control rod calibration for the TRIGA reactor.

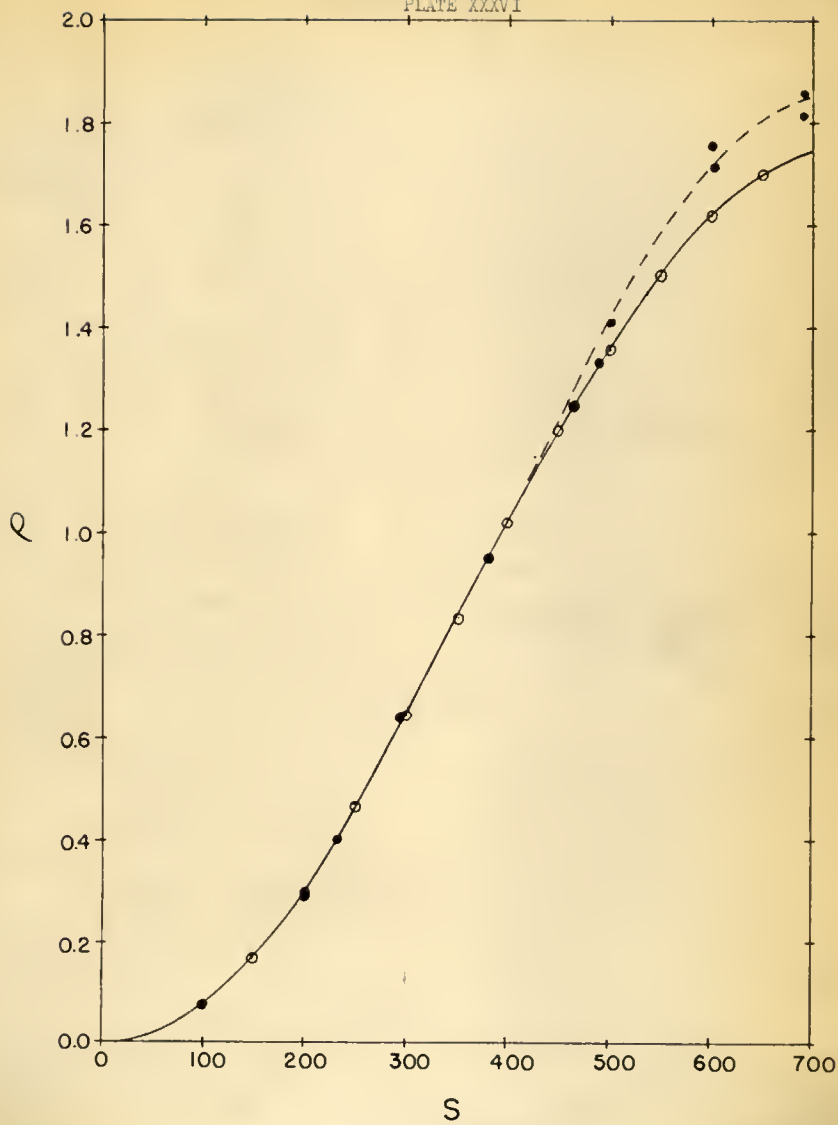
Reactivity, ρ (dollars), vs. control rod position, S (units).

---◆--- Measured by dropping the shim rod from various critical positions to "full in".

—⊖— Measured by the positive period method using increments of reactivity.

Hughes' delayed neutron data.

PLATE XXXVI



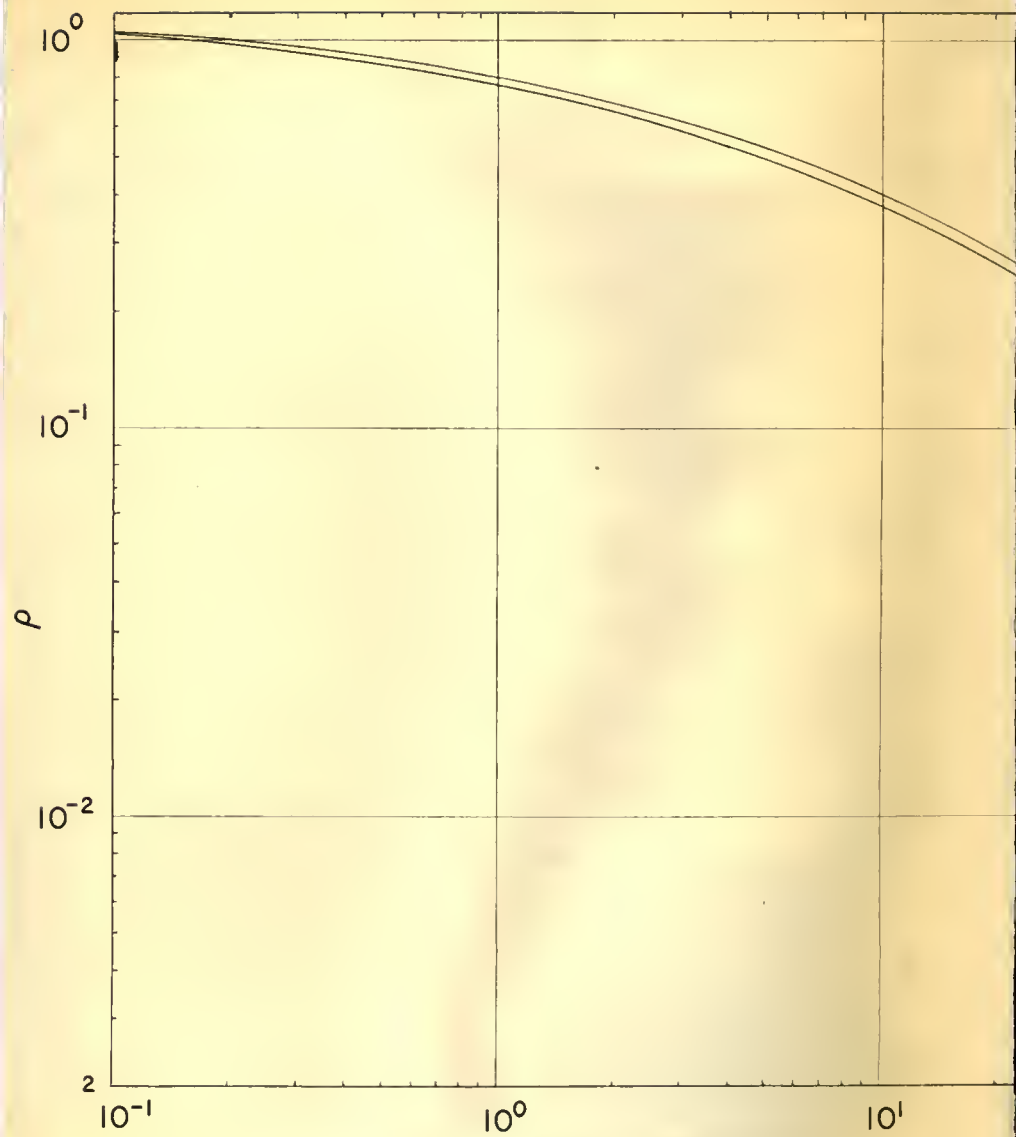
EXPLANATION OF PLATE XXXVII

Positive reactivity, ρ (dollars), vs. positive asymptotic period, τ (seconds).

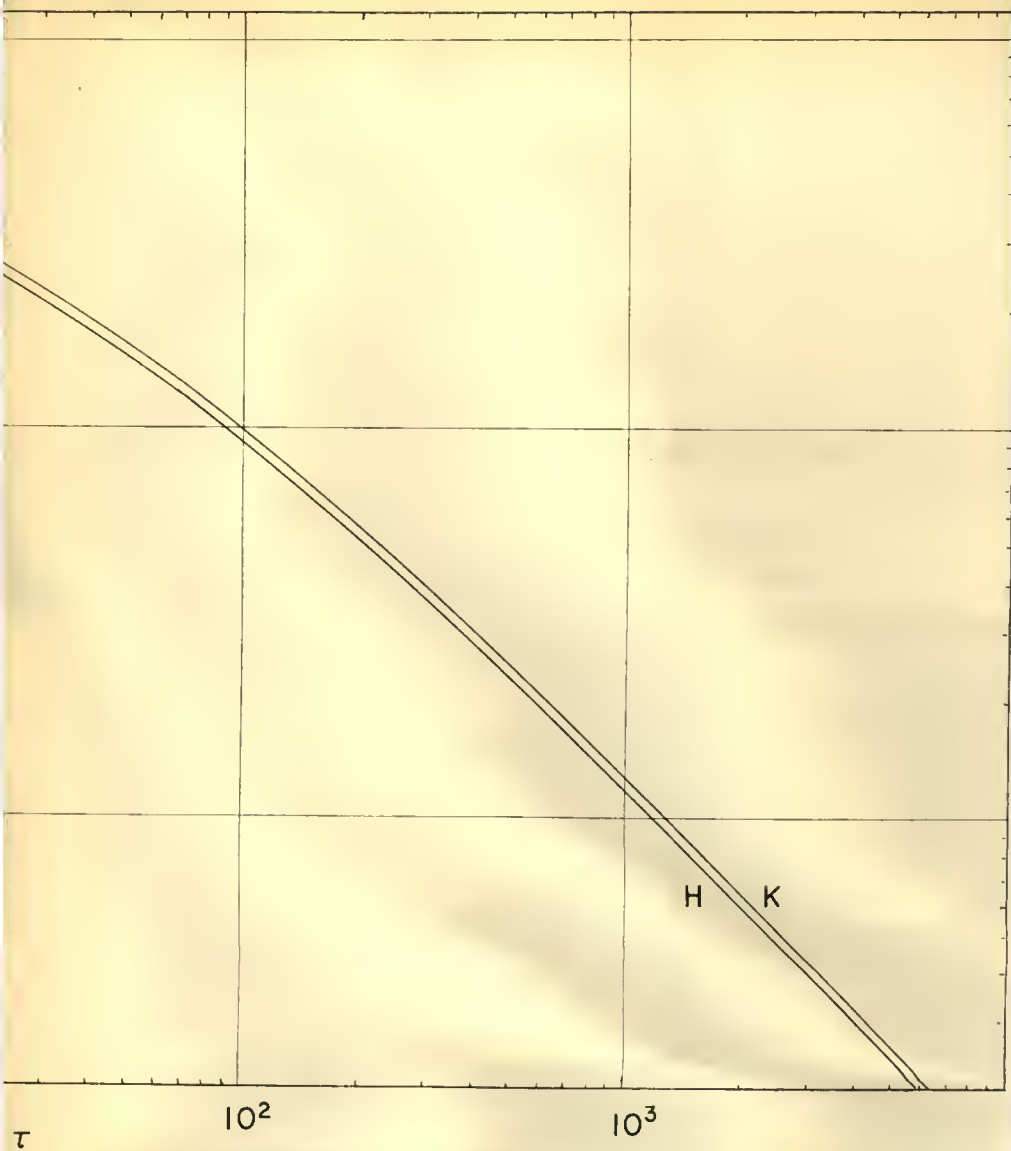
K curve: Keepin's delayed neutron parameters.

H curve: Hughes' delayed neutron parameters.

Neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.



TE XXXVII



H

K

 10^2 10^3 τ

EXPLANATION OF PLATE XXXVIII

Regulating control rod calibration for the TRIGA reactor.

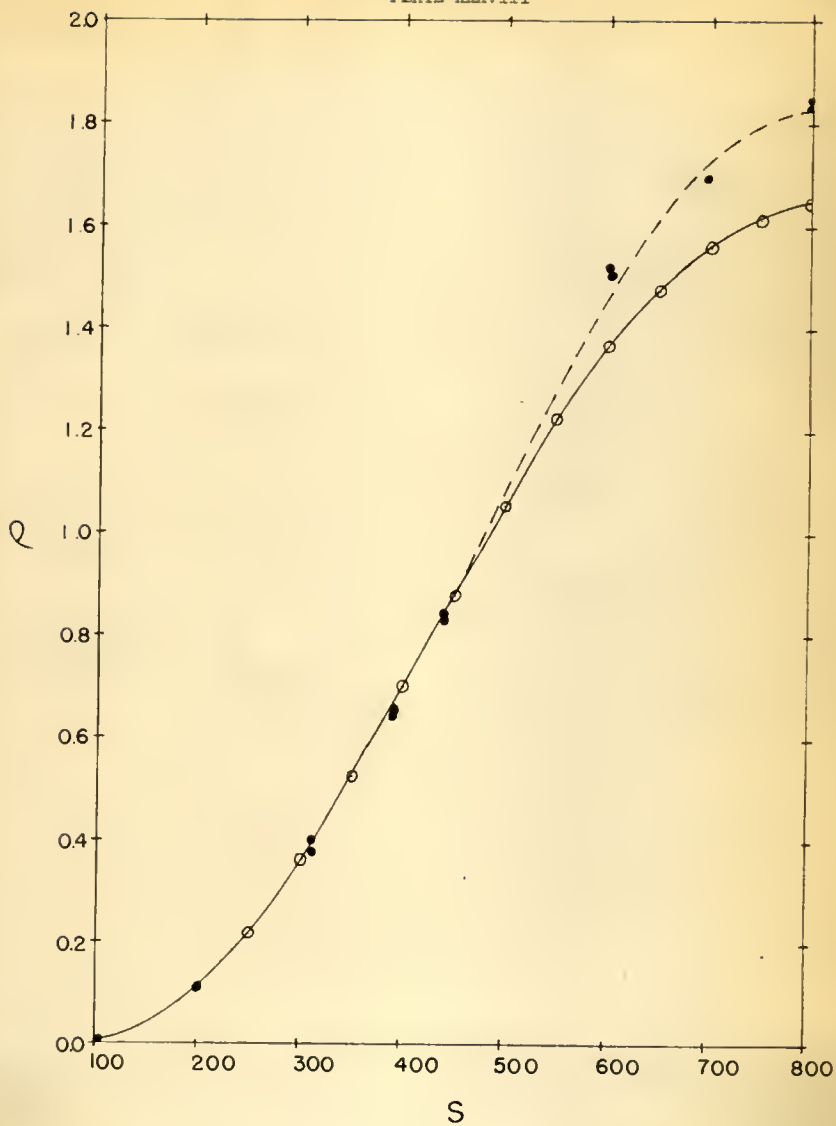
Reactivity, ρ (dollars), vs. control rod positions, S (units).

--- -- Measured by dropping the regulating rod from various critical positions to "full in".

—○— Measured by the positive period method using increments of reactivity.

Hughes' delayed neutron parameters.

PLATE XXXVIII



An interesting point was demonstrated in the control rod calibration by the positive period method. During the process of positive period measurements, the safety rod, a control rod of almost equal magnitude reactivity worth as the shim rod or the regulating rod, was moved from its "full out" position to its "full in" position. Thus, the safety rod was moved from one extreme of its reactivity worth position to the other extreme. The point of interest is that the position of the safety rod had no effect on the calibration of either the shim rod or the regulating rod. If the position of the safety rod had created another parameter, "shadowing effect", in the calibration procedure, a shift in the reactivity calibration curve would have been observed. The shift would have been in the form of a finite jump in the calibration curve.

Forty-three rod drop experiments were performed to measure the reactivity worth of the control rods. The control rods used in the rod drop experiments were the shim rod and the regulating rod. In each case the control rod was dropped from its critical position to its "full in" position; thus the reactivity that was calculated was the difference in reactivity of the two positions. It was assumed that the rod drop represented a negative step change in reactivity. The time required for each control rod to travel from its "full out" to its "full in" position was approximately 0.35 seconds. Thus, at times large compared to 0.35 seconds, the reactivity insertion appeared as a negative step change.

Because of the finite reactivity insertion, it was found

that for times less than eight seconds after the rod drop was initiated the measured neutron flux was higher than given by the theoretically predicted flux. Since a step change in reactivity was assumed in the theory, this observed effect was expected. The discrepancy between the analytically determined flux and the experimentally measured flux was greater, for short times after the rod drop, when larger reactivity insertions were made (see PLATE XLII and PLATE XLIII). Since the reactivities were determined from the ratio, $\phi(t)/\phi(0)$, at various times after the change in reactivity, the apparent reactivity, measured at times just after the rod drop was made, were smaller in magnitude than the true value of the reactivity change. In each case, the apparent reactivity as a function of time would reach a fairly stable value for times larger than ten seconds and then would vary somewhat around the stable value. There was found to be no uniform shift in the reactivity values as given at large times after the rod drop, although the reactivities did tend to drift away from the stable value of reactivity for times larger than 75 seconds. In each run, the reactivities were averaged for times greater than ten seconds and less than 100 seconds and the mean value was reported as the measured reactivity of the rod drop.

Integral control rod calibrations were made for the shim rod and the regulating rod (see PLATES XXXVI and XXXVIII). The total worth of the shim rod was calculated to be \$1.97 and \$1.84 using Keepin's and Hughes' delayed neutron data respectively. The rod drop method gave reactivity worths which were larger in

magnitude for large reactivities than did the positive period method (see PLATE XXXVIII). The total worth of the shim rod was found to be 5.1 per cent higher when measured with the rod drop method than when measured with the positive period method using either Keepin's or Hughes' delayed neutron data. (See PLATES XXXIV and XXXVI). This discrepancy is attributed primarily to the fact that the total reactivity worth of the control rod was measured in one reactivity step when the rod drop method was used whereas when the positive period method was used, the total reactivity worth was determined by 12 incremental steps of reactivity.

Argonaut Data

Through the courtesy of the Argonaut reactor personnel at the International School of Nuclear Science and Engineering, Argonne National Laboratory, experimental data were obtained from the Argonaut reactor. The data consisted of several rod drop experiments with both a slab and an annular core loading.

An analysis of the data gave similarly behaved results as that given by the TRIGA. When the rod drop experiments were analyzed using Keepin's data the evaluated reactivity was always higher than that given by Hughes' data. The same type of disagreement was found between reactivity as measured by the positive period method and reactivity as measured by the rod drop method. At small reactivities, the two methods gave substantially the same results but at larger reactivities, the rod drop method gave higher values. The fine rod of the Argonaut reactor

was calibrated both by dropping the rod from various critical positions to "full in" and then by dropping the rod from a critical position with the fine rod in its "full out" positions to various lower positions. The two methods of rod drops gave approximately the same results (see PLATES XLI and XLII).

Reactivity as a Function of Time

General Atomic (19) measured the time required for a typical TRIGA control rod to fall from its "full out" position to its "full in" position. The time required was approximately 0.35 seconds. The fall time was certainly finite when considering times less than ten seconds after initiating the rod drop. It was considered important to study the problem of predicting the reactivity insertion rate and the resulting time dependence of the neutron flux.

Since the TRIGA core was located in a water bath, the control rods fell in a fairly resistive medium during rod drop measurements (see Experimental System). It was assumed that the control rod motion could be treated as a falling body in a viscous medium where the resistive force is proportional to velocity (3). Therefore, the force, F , acting on the control rod, was assumed to be of the form

$$F = -mg - r \frac{dy}{dt} = m \frac{d^2y}{dt^2} . \quad (14)$$

Where r is the proportionality constant, m is the mass of the control rod, g is the acceleration due to gravity, and $\frac{dy}{dt}$

EXPLANATION OF PLATE XXXIX

Calibration curve for the fine control rod of the Argonaut reactor.

Reactivity, ρ (dollars), vs. control rod position, S.

Upper set of curves: Keepin's delayed neutron data.

—⊖— Measured by dropping rod from various critical positions to full in.

□ Measured by dropping rod from a critical position at 100 to various lower positions.

---+--- Measured by positive period method.

Lower set of curves: Hughes' delayed neutron data.

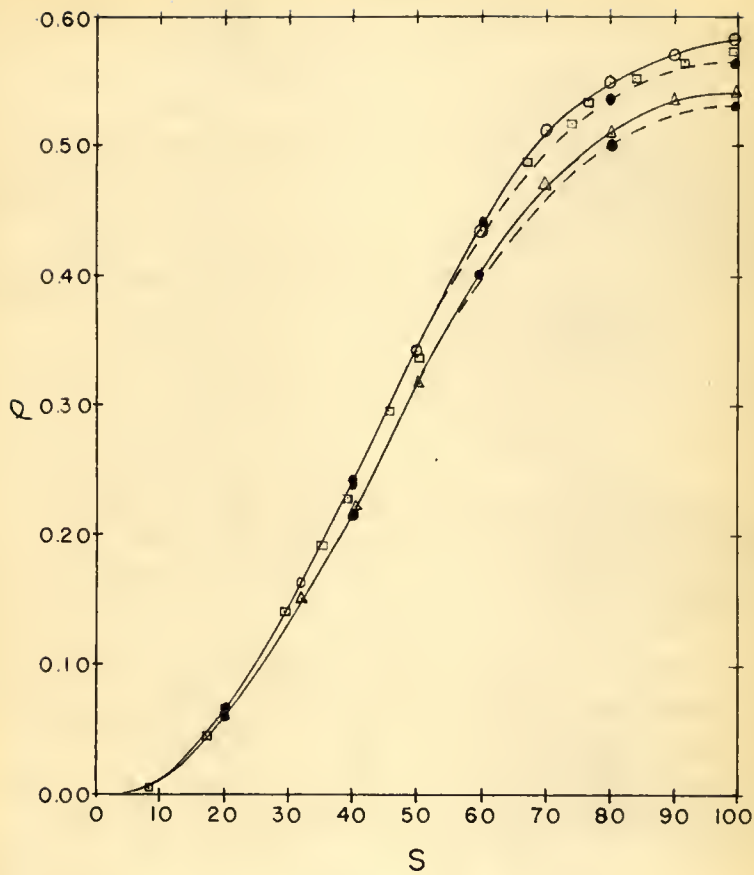
—△— Measured by dropping rod from various critical positions to full in.

---+--- Measured by positive period method.

The prompt neutron lifetime, $\tau = 2.0 \times 10^{-4}$ seconds.

Slab core loading.

PLATE XXXIX



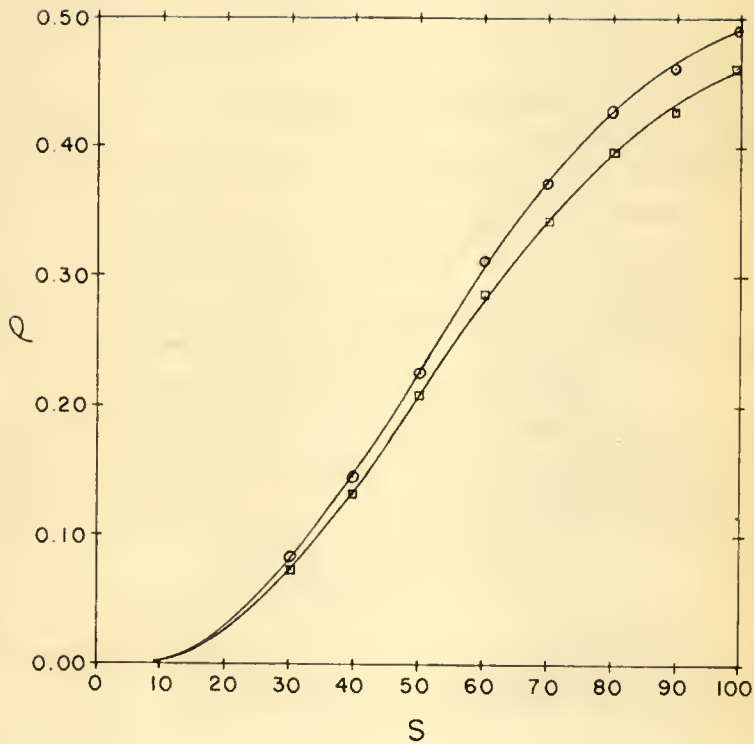
EXPLANATION OF PLATE XL

Shim control rod calibration for the Argonaut reactor, annular loading.

Reactivity, ρ (dollars), vs. control rod position, S (units).

- ⊕— Measured by the rod drop method using Keepin's delayed neutron parameters.
- Measured by the rod drop method using Hughes' delayed neutron parameters.

PLATE XL



is the velocity of the control rod, positive upward. Eq. (14) was solved for the rod displacements, y , as a function of time, i.e.,

$$y = -mgt/r + \frac{m}{r} \left[\frac{dy}{dt}(0) + \frac{mg}{r} \right] (1 - e^{-\frac{r}{m}t}) . \quad (15)$$

Since the control rods were dropped from rest, $\frac{dy}{dt}(0)$ was equal to zero. All of the TRIGA control rods contained approximately the same amount of mass, therefore, the ratio, m/r , was considered to be a constant, c . The rod position as a function of time could then be written as

$$y = cgt - c^2g(1 - e^{-ct}) . \quad (16)$$

The boundary conditions were

$$y(0) = 0$$

$$\text{and } y(0.35 \text{ Sec.}) = -15.25 \text{ inches.}$$

Eq. (16) was solved for the constant, $c = 1.0298$ seconds, by trial and error.

The rod position was then plotted as a function of time (see PLATE XLI). From the rod calibration curve and the rod position curve, a reactivity as a function of time relationship was obtained (see PLATE XLI). It was desirable to obtain an analytic function for reactivity as a function of time because most methods for solving the kinetic equations, where $k_{eff}(t)$ is a function of time, demand well behaved time functions of reactivity (1). The form

$$k_{eff}(t) = 1 - A(1 - e^{-bt^n}) , \quad t \geq 0 \quad (17)$$

EXPLANATION OF PLATE XLI

S curve: Fall of a control rod in the TRIGA reactor, Eq. (16).

The position of a control rod, S (inches), vs. time (seconds).

The zero position of the control rod is the "full out" position.

———— Reactivity curve: The reactivity insertion due to a falling control rod in the TRIGA reactor as predicted by

$$\rho \doteq \frac{\Lambda}{\beta} (1 - e^{-bt^n}) .$$

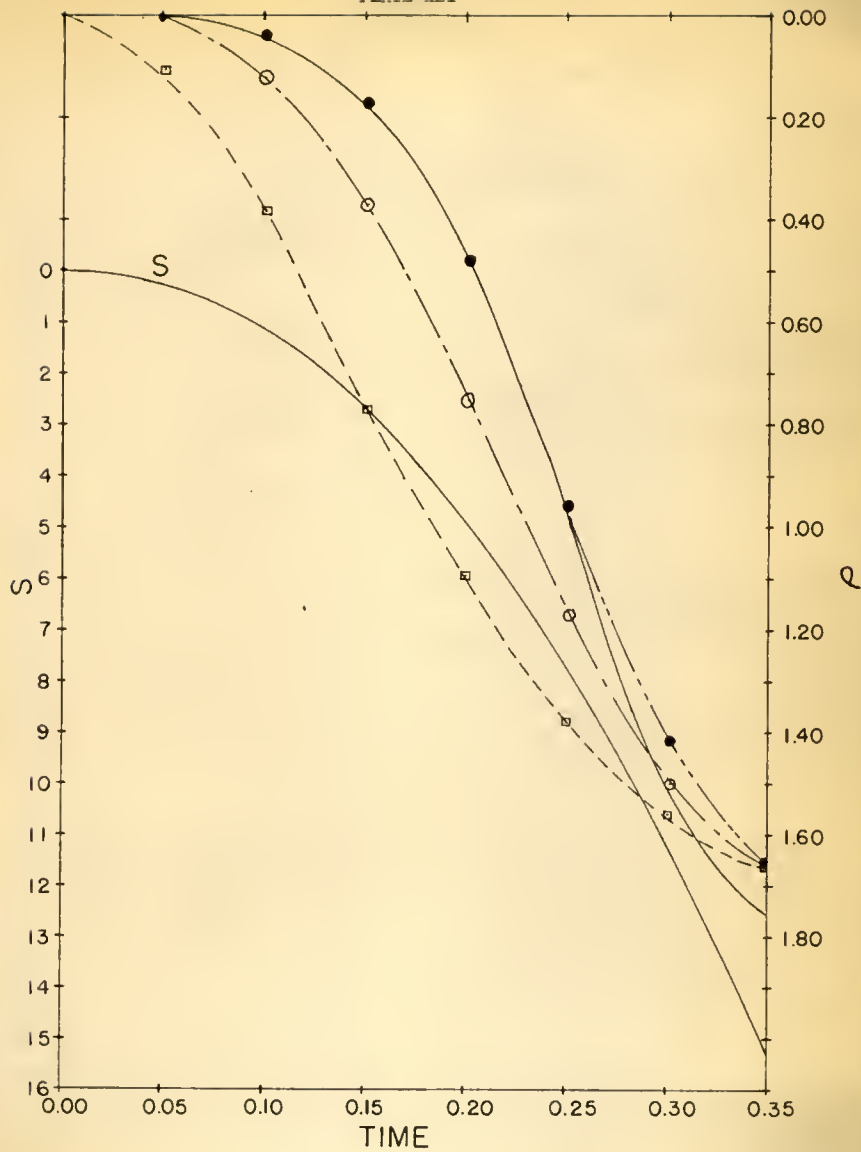
Reactivity, ρ (dollars) vs. time (seconds).

--□-- where n is equal to 2.0.

---○--- where n is equal to 3.0.

----- where n is equal to 4.0.

PLATE XLI



is used here. Note that $k_{\text{eff}}(0)$ is equal to unity and $k_{\text{eff}}(t)$ is approximately equal to $(1.0 - A)$ at large times.

In order to solve for b in Eq. (17), an arbitrary boundary condition was imposed upon Eq. (17) such that the time dependent term of $k_{\text{eff}}(t)$ be 0.95 of its final value at time equal to 0.35 seconds, i.e., $k_{\text{eff}}(0.35) = 1.00 - 0.95A$. Corresponding to various assumed values of n the following values for b were calculated:

$$n = 2 ; b = 24.5$$

$$n = 3 ; b = 69.9$$

$$n = 4 ; b = 300$$

Since reactivity is

$$\rho = \frac{k_{\text{eff}} - 1}{k_{\text{eff}}} = \frac{k_{\text{ex}}}{k_{\text{eff}}} = k_{\text{ex}} \quad (18)$$

and
$$\rho(\beta) = \frac{k_{\text{ex}}}{\beta}$$

then

$$\rho(\beta) = \frac{A}{\beta}(1 - e^{-bt^n}) . \quad (19)$$

Reactivity, ρ , from Eq. (19) was plotted as a function of time (see PLATE XLI) for the various values of n . It was found that n equal to 4.0 gave the best empirical fit to the reactivity curve.

$k_{\text{eff}}(t)$ from Eq. (17) was used in the reactor kinetic equation (see Appendix - B). Several computations were made and the resulting predictions of neutron flux as a function of time

after a change in reactivity were compared to experimentally measured values. The kinetic equations used for the calculations were of the form

$$n(t) = n(0) \left[\frac{a_2 + (\beta - 1)}{a_2 + (\beta - 1) e^{-bt}} \right] e^{-\int_0^t \frac{\lambda dt'}{a_2 + (\beta - 1) e^{-bt'}} + \int_0^t \frac{\lambda dt'}{a_2 e^{-bt'} + (\beta - 1)}} \quad (20)$$

where the decay constant, λ , represents the single delayed neutron group. It was found that the value of n equal to 4.0 gave the best comparison with the experimental data. Even then, although the predicted flux levels from the approximated reactivity equation gave better comparisons for very short times after the initial rod drop than those using the six group step input analysis, the results were in general quite poor. It was then decided to let both b and λ be empirical parameters. The best fit to experimental measurements was obtained when (see PLATE XLII)

$$b = 10.0 \text{ Sec.}^{-1}$$

$$\lambda = 0.20 \text{ Sec.}^{-1}$$

These parameters gave good results over a fairly wide range of reactivities. The largest reactivity insertion used was $-\$1.42$ (see PLATE XLII). The finite reactivity insertion method gave decisively better predictions of the time behavior of the neutron flux for small times, $t < \text{four seconds}$. At larger times the model gave increasingly poorer results due to the assumption of one group of delayed neutrons.

EXPLANATION OF PLATE XLII

Flux ratio, $\phi(t)/\phi(0)$, vs. time (seconds)
after a negative change in reactivity of $-\$0.40$.

○ Experimentally measured flux ratios in the TRIGA reactor after dropping the regulating rod from position 310 to 100.

----- Theoretically predicted flux ratios using six groups of delayed neutrons and assuming a step change in reactivity.

Hughes' delayed neutron parameters.

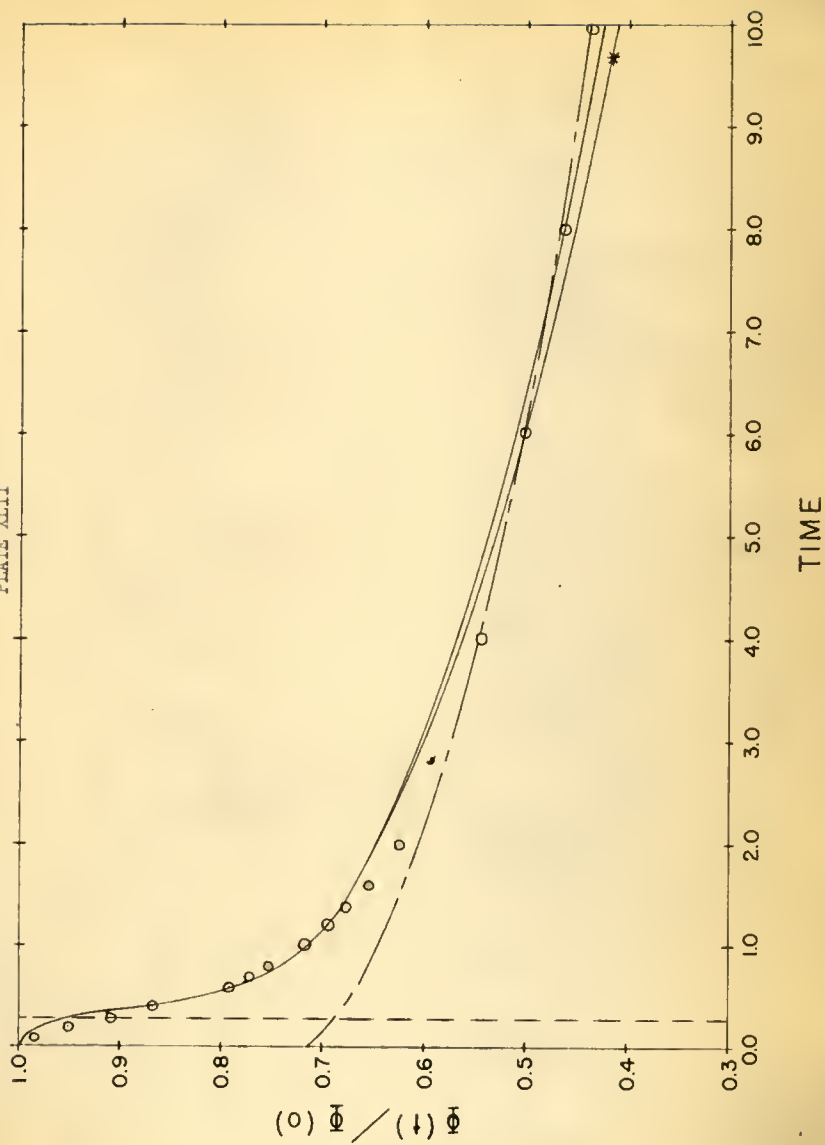
----- Theoretically predicted flux ratios using the reactivity as a function of time, Eq. (16), and assuming one group of delayed neutrons.

b was equal to 10.0 and λ was equal to 0.20 seconds⁻¹.

* In this case λ was equal to 0.22 seconds⁻¹.

----- Total rod drop time.

PLATE XLIII



EXPLANATION OF PLATE XLIII

Flux ratio, $\phi(t)/\phi(0)$, vs. time (seconds) after a negative change in reactivity of -0.142 .

o Experimentally measured flux ratios in the TRIGA reactor after dropping the shim rod from position 500 to 29.

----- Theoretically predicted flux ratios using six groups of delayed neutrons and assuming a step change in reactivity.

Hughes' delayed neutron parameters.

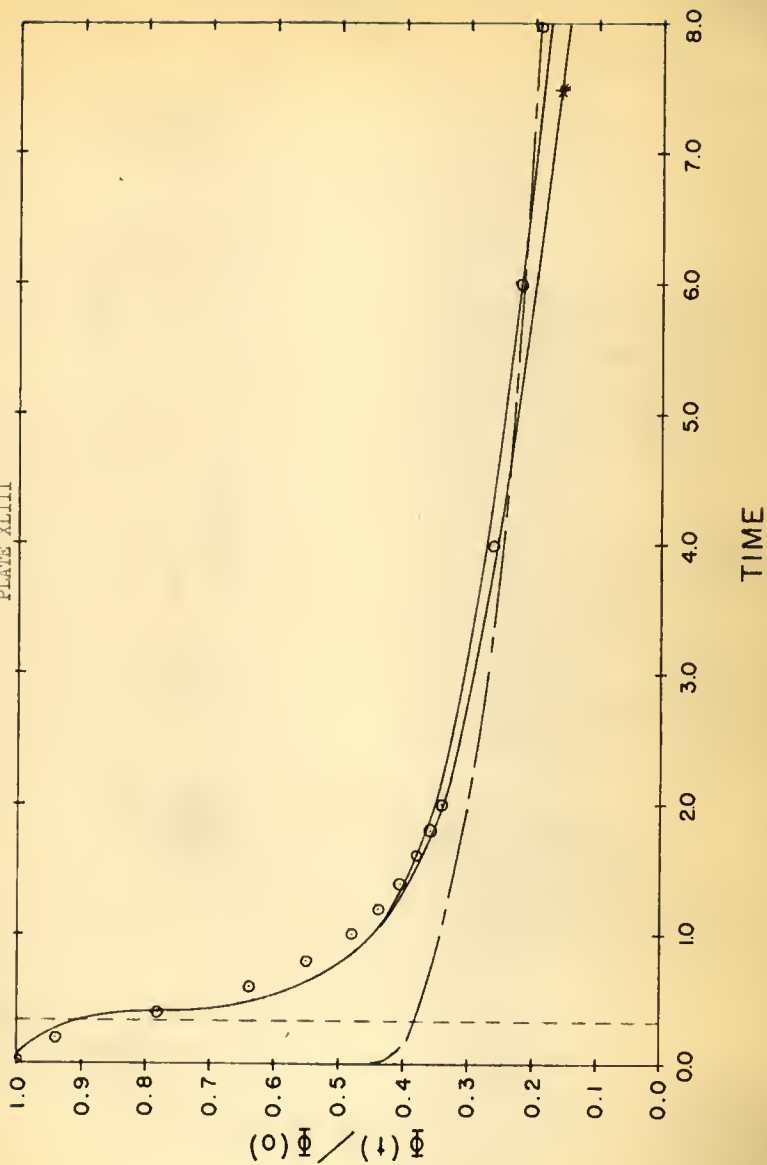
----- Theoretically predicted flux ratios using the reactivity as a function of time, Eq. (16), and assuming one group of delayed neutrons.

b was equal to 10.0 and λ was equal to 0.20 seconds⁻¹.

* In this case λ was equal to 0.22 seconds⁻¹.

----- Total rod drop time.

PLATE XLIII



CONCLUSIONS

The reactor kinetic theory used in this paper is quite satisfactory for use with the TRIGA reactor, at least for reactor operation at low power levels. The conclusion above is based on the agreement between reactivity determination by the method of positive period measurements and by the method of rod drop measurements. The rod drop method depended upon the theoretical prediction of the flux ratio over wide ranges of time. Good agreement was found between the predicted flux ratios and the experimentally measured flux ratios for Keepin's and Hughes' delayed neutron parameters. The agreements between predicted and measured flux ratios were poorer for large than for small reactivities. Agreement was unacceptable at reactivities greater than $\$2.00$. For reactivity measurements greater than $\$1.00$, there was observed a slight scatter of the experimentally determined flux ratios. It is not certain whether the scatter of the observed data was due to the inadequacy of the kinetics model or to the failure of the experimental equipment to accurately measure the rapidly falling flux levels. A possibility that the kinetics model cannot be extended to cover large negative changes in reactivity was indicated by the consistent prediction of larger values of reactivity by the rod drop method as compared with the positive period method. It must be remembered that all of the positive period measurements were composed of small positive changes in reactivity. It is recommended that when using the rod drop method, for evaluating the reactivity

worth of a control rod, the measurements be made using moderately small changes, i.e., that the total reactivity worth of a control rod be measured by a series of incremental rod drops and not be dropping the control rod from "full out" to "full in".

It was found that the two delayed neutron parameter models, Hughes' and Keepin's, gave similar theoretical predictions of the kinetic behavior of a reactor. Keepin's model consistently gave larger evaluations of reactivity, approximately eight per cent, when both the positive period and the rod drop methods were used to calibrate the control rods of the TRIGA reactor and the Argonaut reactor. The two delayed neutron models gave similarly shaped control calibrations for both the TRIGA reactor and the Argonaut reactor, that is, when evaluating large negative changes in reactivity the disagreement between the positive period and the rod drop method evaluations were similar for both reactors and for both delayed neutron models. No conclusion could be made as to which model fit the experimental measurements best, because both models were consistent, within themselves, when considering small changes in reactivity and both models gave similar disagreements for large reactivities. But it is the opinion of this author that either Keepin's or Hughes' model may be used in reactivity evaluation.

It was found, in the Beta Parameter analysis, that if reactivities are quoted in units of dollars there is no confusion as to the response of a reactor to a particular change in reactivity ($\$$), that is, regardless of the total fraction of delayed neutrons, the same reactivity ($\$$) input gave approximately the

same predicted flux ratios.

The kinetic equation parameter analysis revealed that the only significant parameter is reactivity. The other parameters may be varied to any reasonable values and the resulting theoretical change of the kinetic behavior of a reactor will be slight. The parameter study confirms a generalized study by Henry (9).

The three group kinetic equations used in this paper confirm work done by Skinner and Cohen (16), in that the six group equations may be presented by the reduced group equations. The real usefulness of the three group equations probably does not lie in their ability to predict the reactivity worth of reactor control rods, but rather, to predict the general nature of the time dependency of the neutron flux in a reactor where a series of reactor calculations must be made considering many parameters in addition to the usual parameters. The computer code developed for this paper may be altered to consider a wide class of kinetic problems. These computer calculations are more rapid and easier to analyze than the calculations obtained with the six group equations.

When considering rod drop experiments or other reactivity insertion problems, where the reactivity insertion is rapid but still finite, the six group kinetic equations will predict the time dependency of the neutron density when evaluated at times much larger in magnitude than the total reactivity insertion time. But, when neutron densities are evaluated at times comparable to the total insertion time, it is necessary to consider

reactivity as a function of time instead of to assume the step function. A set of equations developed in this paper approximated the reactivity insertion rate of a control rod drop and the resulting neutron flux in the TRIGA reactor. The predicted time dependency of the neutron flux gave values which were in good agreement with the experimentally measured neutron flux. The approximated reactivity equations gave better predictions of the neutron flux at times less than four seconds, after initiating the reactivity change, than did the six group equations using a step input of reactivity. For times greater than four seconds, the neutron flux was predicted to a good degree of accuracy by the kinetic equations using six delayed groups and a step input. The reason for the poor results, using the approximated reactivity equations at large times was, no doubt, due to the one delayed neutron group assumption.

ACKNOWLEDGMENTS

The author takes this opportunity to express his appreciation to the following people whose advice, encouragement, and technical guidance have made the completion of this project possible:

Dr. William R. Kimel, Head of the Department, for suggesting this study and for his excellent technical judgment which greatly aided the progression of this thesis.

Richard C. Bailie, for his excellent guidance and supervision in the use of the IBM-650 computer and for his valuable time that he spent in writing portions of a computer program.

General Atomic, for the use of their TRIGA reactor and the excellent cooperation of the TRIGA reactor staff.

Dr. William Sturm, of Argonne National Laboratory, for the experimental data forwarded to this author.

The author also wishes to express his gratitude to the Kansas State University Engineering Experiment Station for their financial support and to the Atomic Energy Commission for the opportunity to attend graduate school with the aid of an A.E.C. Fellowship.

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APPENDIX

Part - A

Theoretical Development of the Kinetic Equations

The system considered was a bare, homogeneous, thermal reactor. A thermal neutron balance on a volume element of the reactor core gives (8)

$$D\nabla^2\Phi(r,t) - \Sigma_a\Phi(r,t) + S(r,t) = \frac{\partial n(r,t)}{\partial t} = \frac{1}{v} \frac{\partial\Phi(r,t)}{\partial t} \quad (20)$$

The Fermi age equation is used as the source of neutrons for the thermal diffusion equation. The source term for prompt neutrons is

$$S(r,t)_{\text{prompt}} = (1-\beta)k\Sigma_a\Phi(r,t)e^{-B^2\tau} \quad (21)$$

The rate equation for the precursor density is

$$\frac{\partial C_i(r,t)}{\partial t} = -\lambda_i C_i(r,t) + \frac{k\beta_i\Sigma_a\Phi(r,t)}{\rho} \quad (22)$$

where $-\lambda_i C_i$ gives the loss of neutron precursors due to radioactive decay and $\frac{k\beta_i\Sigma_a\Phi(r,t)}{\rho}$ gives the production of precursors from the fission process. Since delayed neutrons are produced following the beta decay of the precursors, the total production rate of all delayed neutrons is $\sum_{i=1}^m \lambda_i C_i$. Since there are six groups of delayed neutrons, m is equal to six in the above summation. Delayed neutrons are born with energies between 0.2 and 0.8 Mev., thus the production rate of delayed neutron term must be multiplied by the nonleakage probability

during slowing down and by the resonance escape probability.

$$S(r,t)_{\text{Delayed}} = \sum_{i=1}^6 p_i e^{-\beta^2 \tau_i} \lambda_i C_i(r,t) \quad (23)$$

The source term for both the prompt neutrons and the delayed neutrons is substituted into the thermal diffusion equations and the resulting equation is

$$\begin{aligned} D \nabla^2 \Phi(r,t) - \Sigma_a \Phi(r,t) + (1-\beta) k \Sigma_a \Phi(r,t) e^{-\beta^2 \tau} \\ + \sum_{i=1}^6 p_i e^{-\beta^2 \tau_i} \lambda_i C_i(r,t) = \frac{1}{v} \cdot \frac{\partial \Phi(r,t)}{\partial t} \end{aligned} \quad (24)$$

The following, well known, relations were substituted into Eq. (24)

$$L^2 = D / \Sigma_a \quad (25)$$

$$l_0 = 1 / \Sigma_a v \quad (26)$$

Eq. (24) was then divided by Σ_a , which gives

$$\begin{aligned} L^2 \nabla^2 \Phi(r,t) + \left[(1-\beta) k e^{-\beta^2 \tau} - 1 \right] \Phi(r,t) \\ \frac{1}{\Sigma_a} \sum_{i=1}^6 p_i e^{-\beta^2 \tau_i} \lambda_i C_i(r,t) = l_0 \frac{\partial \Phi(r,t)}{\partial t} \end{aligned} \quad (27)$$

The space and time variables of the thermal diffusion and the precursor rate equations were assumed to be separable, that is

$$\bar{\Phi}(r,t) = F(r) \phi(t) \quad (28)$$

and

$$C_i(r,t) = G_i(r) H_i(t) . \quad (29)$$

It is assumed that the spatial distribution of the thermal neutrons can be given by the fundamental mode of the wave equation,

$$\nabla^2 F(r) + B^2 F(r) = 0 . \quad (30)$$

Since ∇^2 is a spatial operator, $B^2 F(r)$ is substituted into Eq. (27) after the separation of variables, thus

$$\begin{aligned} -L^2 \phi(t) B^2 F(r) + [(1-\beta) k e^{-B^2 \tau} - 1] F(r) \phi(t) \\ + \sum_{i=1}^6 \frac{\rho_i e^{-B^2 \tau_i}}{\Sigma_a} \lambda_i G_i(r) H_i(t) = l_0 F(r) \frac{d\phi(t)}{dt} . \end{aligned} \quad (31)$$

The definition of the effective multiplication factor and the finite prompt neutron lifetime is used to simplify Eq. (32)

$$k_{\text{eff}} \equiv \frac{k e^{-B^2 \tau}}{1 + L^2 B^2} = \text{effective multiplication factor} \quad (33)$$

$$l \equiv \frac{l_0}{1 + L^2 B^2} = \text{finite prompt neutron lifetime} \quad (34)$$

Eq. (32) is divided by $1 + L^2 B^2$ and then Eq. (33) and Eq. (34) is substituted into Eq. (32).

Eq. (31) is then divided by the spatial flux function to give

$$\frac{\ell d\phi(t)}{dt} = -\phi(t) + (1-\beta)k_{eff}\phi(t) + \frac{1}{\sum_{\alpha} (1 + L^2 B^2) F(\alpha)} \cdot \sum_{i=1}^6 p_i e^{-B^2 \tau_i} \lambda_i G_i(\alpha) H_i(t) \quad (35)$$

Eq. (35) is divided by the neutron velocity, v , and the expression

$$\frac{\phi(t)}{v} = n(t)$$

is substituted into Eq. (35) giving

$$\ell \frac{dn(t)}{dt} = -n(t) + (1-\beta)k_{eff} n(t) + \frac{1}{\sum_{\alpha} v F(\alpha) (1 + L^2 B^2)} \cdot \sum_{i=1}^6 p_i e^{-B^2 \tau_i} \lambda_i G_i(\alpha) H_i(t). \quad (36)$$

The relationship

$$\ell = \frac{1}{\sum_{\alpha} v (1 + L^2 B^2) F(\alpha)} = \text{finite prompt neutron lifetime}$$

is substituted into Eq. (36) to give

$$\left[(1-\beta)k_{eff} - 1 \right] \frac{n(t)}{\ell} + \frac{1}{F(\alpha)} \sum_{i=1}^6 p_i e^{-B^2 \tau_i} \lambda_i G_i(\alpha) H_i(t) = \frac{dn(t)}{dt} \quad (37)$$

$$\frac{\rho_i e^{-\beta^2 \tau_i} G_i(\tau)}{F(\tau)}$$

is defined as the space independent

constant, ϵ_i .

The precursor rate equation is then written in the separated variable form as

$$\frac{dH_i(t)}{dt} G_i(\tau) = -\lambda_i G_i(\tau) H_i(t) + \frac{k \beta_i \sum_a F(\tau) \phi(t)}{\rho} \quad (38)$$

The resulting equation, (39), contains functions of time alone on the left side and functions of space alone on the right side.

$$\frac{dH_i(t)}{dt} \frac{1}{\phi(t)} + \frac{\lambda_i H_i(t)}{\phi(t)} = \frac{k \beta_i \sum_a F(\tau)}{\rho G_i(\tau)} \quad (39)$$

Therefore ϵ_i is a time independent parameter.

The space independent neutron density equation is then written as

$$\left[(1 - \beta) k_{eff} - \rho \right] \frac{n(t)}{\lambda} + \sum_i^6 \epsilon_i \lambda_i H_i(t) = \frac{dn(t)}{dt} \quad (40)$$

The space independent precursor equation is then obtained by dividing Eq. (40) by $G_i(\tau)$, which gives

$$\frac{dH_i(t)}{dt} = -\lambda_i H_i(t) + \frac{k \beta_i \sum_a \phi(t) F(\tau)}{\rho G_i(\tau)} \quad (41)$$

The last term on the right hand side of Eq. (41) is then multiplied and divided by $\nu e^{-\beta^2 \tau} (1 + L^2 B^2)$ to give

$$\frac{dH_i(t)}{dt} = -\lambda_i H_i(t) + \frac{k e^{-B^2 \tau}}{1 + L^2 B^2} \cdot \frac{\phi(t)}{\nu} \beta_i \frac{\sum_{\tau} \nu (1 + L^2 B^2) F(\tau)}{G_i(\tau) p e^{-B^2 \tau}} \quad (42)$$

ϵ'_i then defined as $\frac{F(\tau)}{G_i(\tau) p e^{-B^2 \tau}}$, thus ϵ'_i is similar to ϵ of Eq. (35) and likewise time independent. The resultant space independent precursor rate equation is obtained as

$$\frac{dH_i(t)}{dt} = -\lambda_i H_i(t) + \frac{k_{\text{eff}} n(t) \beta_i \epsilon'_i}{\rho} \quad (43)$$

The usual definition of reactivity, ρ , is

$$\rho = \frac{k_{\text{eff}} - 1}{k_{\text{eff}}} \quad (44)$$

Thus

$$k_{\text{eff}} = \frac{1}{1 - \rho} \quad (45)$$

Eq. (45) is substituted in the space independent kinetic equation to give

$$\frac{\rho - \beta}{\rho(1 - \rho)} n(t) + \sum_{i=1}^6 \epsilon_i \lambda_i H_i(t) = \frac{dn(t)}{dt} \quad (46)$$

The same substitution, as above, into Eq. (43) results in

$$\frac{dH_i(t)}{dt} = -\lambda_i H_i(t) + \frac{\beta_i n(t) \epsilon'_i}{\rho(1 - \rho)} \quad (47)$$

Eq.'s (46) and (47) constitute a set of seven linear differential equations, therefore solutions of the form

$$n(t) = A' e^{\omega t} \quad (48)$$

$$H_i(t) = B_i' e^{\omega t} \quad (49)$$

may be assumed.

Substitutions of Eq.'s (48) and (49) into Eq.'s (46) and (47) give

$$\left(\omega - \frac{\rho - \beta}{\lambda(1-e)} \right) A' + \sum_{i=1}^6 \epsilon_i \lambda_i B_i' = 0 \quad (50)$$

and

$$\frac{\epsilon_i \beta_i A'}{\lambda(1-e)} - (\omega + \lambda_i) B_i' = 0 \quad (51)$$

A solution for the coefficient A' is displayed by Cramer's rule as follows:

$$A = \begin{array}{c}
 \begin{array}{ccccccc}
 0 & \lambda_1 \epsilon_1 & \lambda_2 \epsilon_2 & \lambda_3 \epsilon_3 & \lambda_4 \epsilon_4 & \lambda_5 \epsilon_5 & \lambda_6 \epsilon_6 \\
 0 & -(\omega + \lambda_1) & 0 & 0 & 0 & 0 & 0 \\
 0 & 0 & -(\omega + \lambda_2) & 0 & 0 & 0 & 0 \\
 0 & 0 & 0 & -(\omega + \lambda_3) & 0 & 0 & 0 \\
 0 & 0 & 0 & 0 & -(\omega + \lambda_4) & 0 & 0 \\
 0 & 0 & 0 & 0 & 0 & -(\omega + \lambda_5) & 0 \\
 0 & 0 & 0 & 0 & 0 & 0 & -(\omega + \lambda_6)
 \end{array} \\
 \hline
 \begin{array}{ccccccc}
 \omega - \frac{(\rho - \theta)1}{(1-\rho)\lambda} & \lambda_1 \epsilon_1 & \lambda_2 \epsilon_2 & \lambda_3 \epsilon_3 & \lambda_4 \epsilon_4 & \lambda_5 \epsilon_5 & \lambda_6 \epsilon_6 \\
 \frac{\beta_1 \epsilon_1'}{\lambda(1-\rho)} & -(\omega + \lambda_1) & 0 & 0 & 0 & 0 & 0 \\
 \frac{\beta_2 \epsilon_2'}{\lambda(1-\rho)} & 0 & -(\omega + \lambda_2) & 0 & 0 & 0 & 0 \\
 \frac{\beta_3 \epsilon_3'}{\lambda(1-\rho)} & 0 & 0 & -(\omega + \lambda_3) & 0 & 0 & 0 \\
 \frac{\beta_4 \epsilon_4'}{\lambda(1-\rho)} & 0 & 0 & 0 & -(\omega + \lambda_4) & 0 & 0 \\
 \frac{\beta_5 \epsilon_5'}{\lambda(1-\rho)} & 0 & 0 & 0 & 0 & -(\omega + \lambda_5) & 0 \\
 \frac{\beta_6 \epsilon_6'}{\lambda(1-\rho)} & 0 & 0 & 0 & 0 & 0 & -(\omega + \lambda_6)
 \end{array}
 \end{array} = 0 \quad (52)$$

Since A' is not identically zero except for the trivial solution, the determinant of the coefficients must be identically equal to zero. The determinant of the coefficients was simplified by multiplying the following columns by the indicated products:

column 1 by $(\omega + \lambda_1)(\omega + \lambda_2) \cdots (\omega + \lambda_6)$

column 2 by $(\omega + \lambda_2)(\omega + \lambda_3) \cdots (\omega + \lambda_6) \beta_1 \epsilon'_1 / \ell(1-e)$

column 3 by $(\omega + \lambda_1)(\omega + \lambda_3) \cdots (\omega + \lambda_6) \beta_2 \epsilon'_2 / \ell(1-e)$

column 4 by $(\omega + \lambda_1)(\omega + \lambda_2)(\omega + \lambda_4) \cdots (\omega + \lambda_6) \beta_3 \epsilon'_3 / \ell(1-e)$

column 5 by $(\omega + \lambda_1)(\omega + \lambda_2)(\omega + \lambda_3)(\omega + \lambda_5)(\omega + \lambda_6) \beta_4 \epsilon'_4 / \ell(1-e)$

column 6 by $(\omega + \lambda_1)(\omega + \lambda_2) \cdots (\omega + \lambda_4)(\omega + \lambda_6) \beta_5 \epsilon'_5 / \ell(1-e)$

column 7 by $(\omega + \lambda_1)(\omega + \lambda_2) \cdots (\omega + \lambda_5) \beta_6 \epsilon'_6 / \ell(1-e)$

Each of the columns, except column 1, was added to column 1, thus the determinant is cast into the following form.

A	B	C	D	E	F	G	
0	I	0	0	0	0	0	
0	0	J	0	0	0	0	
0	0	0	K	0	0	0	= 0
0	0	0	0	L	0	0	(53)
0	0	0	0	0	M	0	
0	0	0	0	0	0	N	

The product, $AIJKLMN$, is set equal to zero. Since the terms, I through N, are product terms of the reactor parameters, which are finite, they cannot be equal to zero, hence the A

term must be identically zero. The term, A, in expanded form is

$$A = \left[\omega - \frac{1}{\lambda} \left(\frac{\rho - \beta}{1 - \rho} \right) \right] \left[\prod_i (\omega + \lambda_i) \right] \\ - \frac{1}{\lambda(1-\rho)} \sum_{i=1}^6 \frac{\prod_i (\omega + \lambda_i)}{(\omega + \lambda_i)} \lambda_i \beta_i \epsilon_i \epsilon_i' \quad (54)$$

Equation (54) is divided by $\prod_i (\omega + \lambda_i)$ which gives

$$\lambda(1-\rho) \left[\omega + \frac{1}{\lambda} \left(\frac{\rho - \beta}{1 - \rho} \right) \right] = \sum_{i=1}^6 \frac{\epsilon_i \epsilon_i' \lambda_i \beta_i}{(\omega + \lambda_i)} \quad (55)$$

Since

then, Eq. (55) becomes

$$\lambda\omega - \rho(\lambda\omega + 1) = \sum_{i=1}^6 \epsilon_i \epsilon_i' \left[\frac{\lambda_i \beta_i}{\omega + \lambda_i} - \frac{\beta_i \omega + \beta_i \lambda_i}{\omega + \lambda_i} \right] \quad (56)$$

and

$$\rho = \frac{\lambda w}{\lambda w + 1} + \frac{1}{\lambda w + 1} \sum_{i=1}^6 \frac{\epsilon_i \epsilon'_i \beta_i w}{w + \lambda_i} \quad (57)$$

$\epsilon_i \epsilon'_i \beta_i$ from Eq. (56) is defined as the effective fraction of delayed neutrons, i.e. $\epsilon_i \epsilon'_i$ gives the relative effectiveness of a delayed neutron for producing a succeeding fission as compared to a prompt neutron. $\epsilon_i \epsilon'_i$ is the ratio of the product of the resonance escape probability and the nonleakage probability during slowing down for each of the precursor neutron groups to the product of the resonance escape probability and the nonleakage probability during slowing down for fission spectrum neutrons, that is

$$\epsilon_i \epsilon'_i = \frac{p_i e^{-B^2 \tau_i}}{p e^{-B^2 \tau}} \quad (58)$$

thus

$$\epsilon_i \epsilon'_i \beta_i = \beta_{eff i}$$

$\beta_{eff i}$ was substituted into Eq. (51) and the resulting equation is the characteristic equation for the kinetic differential equations,

$$\rho = \frac{\lambda w}{\lambda w + 1} + \frac{1}{\lambda w + 1} \sum_{i=1}^6 \frac{\beta_{eff i} w}{w + \lambda_i} \quad (59)$$

Equation (59) is a seventh degree polynomial equation in w and thus there are seven roots of Eq. (59). The roots are the

eigen-values, $\omega_1, \omega_2, \dots, \omega_7$, of the characteristic equation. The solution for the neutron density as a function of time is then expressed as

$$n(t) = \sum_{j=1}^7 A_j e^{\omega_j t} \quad (60)$$

and the precursor density is expressed as

$$H_i(t) = \sum_{j=1}^7 B_{ij} e^{\omega_j t} \quad (61)$$

The two coefficients, A_j and B_{ij} , are not independent but may be related by the substitution of Eq.'s (60) and (61) into the rate equation giving

$$B_{ij}(\omega_j + \lambda_i) = \frac{A_j \beta_i \epsilon_i'}{(1-\rho)\lambda}$$

Therefore, (62)

$$B_{ij} = \frac{\epsilon_i' A_j}{(1-\rho)(\omega_j + \lambda_i)}$$

The precursor density equations may then be expressed as

$$H_i(t) = \frac{\beta_i \epsilon_i'}{(1-\rho)} \sum_{j=1}^7 \frac{A_j e^{\omega_j t}}{(\omega_j + \lambda_i)} \quad (63)$$

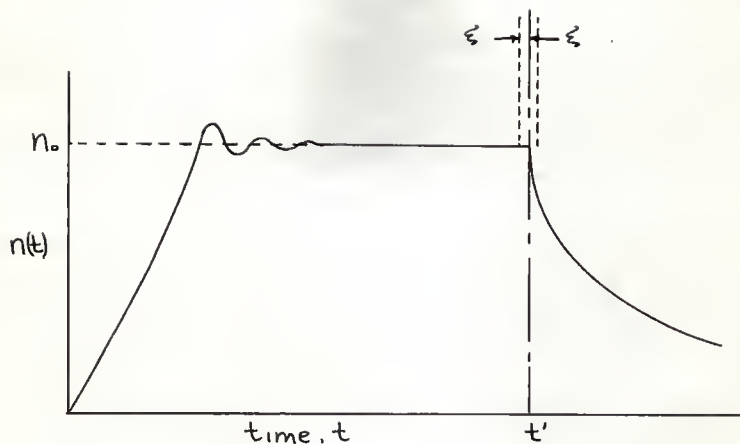


Fig. 1

The neutron and precursor densities are evaluated immediately before and immediately after a step change in reactivity for the following hypothetical case. The reactor is brought to a critical configuration and the neutron density distribution of time is considered as shown in Figure 1. A step change in reactivity occurs at time, t' . At time, $t' - \xi$, the reactivity, ρ , is equal to zero and at time, $t' + \xi$, the reactivity is equal to a constant.

Before the change in reactivity, the neutron density is expressed as

$$n_0 = A_0 e^{\omega_0 t} + A_1 e^{\omega_1 t} + \dots + A_n e^{\omega_n t} \quad (64)$$

When the reactivity is zero, one of the roots, ω_0 , is zero. The fact that one of the roots is exactly equal to zero may be seen by writing the characteristic equation for zero reactivity.

$$\rho = 0 = \frac{\ell\omega}{\ell\omega + 1} + \frac{1}{\ell\omega + 1} \sum_{i=1}^6 \frac{\beta_{eff,i} \omega}{\omega + \lambda_i} \quad (65)$$

The term, $\ell\omega + 1$, may be cancelled from Eq. (65) and the resulting equation is of sixth degree in ω . The remaining roots, all negative, have been called the "natural vibration rates" of the pile by Zoodak (18). Since the reactor is assumed to be critical in the delayed sense, the terms of Eq. (64), $A_{02}e^{\omega_{02}t}$, $A_{03}e^{\omega_{03}t}$, ..., $A_{07}e^{\omega_{07}t}$, truly vanish leaving only the first term, A_{01} . The neutron density is then a single constant at time, $t' - \xi$.

$$n(t' - \xi) = A_{01} = n_0 \quad (66)$$

After the change in reactivity, the neutron density is expressed as

$$n(t' + \xi) = \sum_{j=1}^7 A_j e^{\omega_j(t' + \xi)} \quad (67)$$

The precursor density before the change in reactivity is evaluated from Eq. (62) and Eq. (63) to be

$$H_i(t' - \xi) = \frac{\beta_i A_{01} \epsilon'_i}{\lambda_i} = \frac{\beta_i n(t' - \xi) \epsilon'_i}{\lambda_i} = \frac{\beta_i \epsilon'_i n_0}{\lambda_i} \quad (68)$$

Since $n(t)$ must be continuous at t' , the limit as $\xi \rightarrow 0$ of Eq. (67) is

$$n(t') = \sum_{j=1}^7 A_j e^{\omega_j t'} = n_0 \quad (69)$$

$$H_i(t') = \frac{\beta_i \epsilon'_i}{(1-\epsilon)} \sum_{j=1}^7 \frac{A_j e^{\omega_j t'}}{\omega_j + \lambda_i} = \frac{\beta_i n_0 \epsilon'_i}{\lambda} \quad (70)$$

Since t' is an arbitrary time, a more convenient time was chosen for the step change in reactivity to occur, i.e., $t' = 0$, then Eq. (70) may be written as

$$n(0) = \sum_{j=1}^7 A_j = n_0 \quad (71)$$

The equilibrium precursor density is then

$$H_i(0) = \frac{\beta_i \epsilon_i}{\lambda_i} n(0) \quad (72)$$

The Laplace transform of the spatially independent kinetic equation, Eq. (40), is

$$\Delta \bar{n}(\Delta) - n(0) = \frac{1}{\ell} \left(\frac{\rho - \beta}{1 - \epsilon} \right) \bar{n}(\Delta) + \sum_{i=1}^6 \epsilon_i \lambda_i \bar{H}_i(\Delta) \quad (73)$$

Also, the Laplace transform of the precursor equation is

$$\Delta \bar{H}_i(\Delta) - H_i(0) = -\lambda_i \bar{H}_i(\Delta) + \frac{\beta_i \epsilon'_i \bar{n}(\Delta)}{\ell(1-\epsilon)} \quad (74)$$

Eq. (74) is rewritten as

$$\bar{H}_i(\Delta)(\Delta + \lambda_i) = H_i(0) + \frac{\beta_i \bar{n}(\Delta) \epsilon'_i}{\ell(1-\epsilon)} \quad (75)$$

$\bar{H}_i(\Delta)$ from Eq. (75) is substituted into Eq. (67) and the result is

$$\Delta \bar{n}(\Delta) - n(0+) = \frac{1}{\ell} \left(\frac{\ell - \beta}{1 - \ell} \right) \bar{n}(\Delta) + \sum_{i=1}^6 \epsilon_i \lambda_i \left(\frac{H_i(0+) + \frac{\beta_i \bar{n}(\Delta) \epsilon_i}{\ell(1-\ell)}}{\Delta + \lambda_i} \right) \quad (76)$$

Since $\beta = \sum_{i=1}^6 \beta_i$; $H_i(0+) = \frac{\beta_i n(0) \epsilon_i}{\lambda_i}$

and $n(0+) = n_0$

then Eq. (76) may be written as

$$\bar{n}(\Delta) \left[\Delta + \frac{1}{\ell} \left(\frac{\ell - \beta}{1 - \ell} \right) - \sum_{i=1}^6 \frac{\epsilon_i \epsilon_i' \lambda_i \beta_i n(0)}{(\Delta + \lambda_i)} \right] = n(0) + \sum_{i=1}^6 \frac{\epsilon_i \epsilon_i' \beta_i n(0)}{(\Delta + \lambda_i)} \quad (77)$$

Eq. (77) may be solved for $\bar{n}(\Delta)$,

$$\bar{n}(\Delta) = \frac{n(0) + n(0) \sum_{i=1}^6 \frac{\epsilon_i \epsilon_i' \beta_i}{\Delta + \lambda_i}}{\Delta - \frac{1}{\ell} \left(\frac{\ell - \beta}{1 - \ell} \right) - \sum_{i=1}^6 \frac{\epsilon_i \epsilon_i' \lambda_i \beta_i}{\ell(1-\ell)(\Delta + \lambda_i)}} \quad (78)$$

The characteristic equation may be rewritten as

$$\omega - \frac{\ell - \beta}{\ell(1-\ell)} - \sum_{i=1}^6 \frac{\epsilon_i \epsilon_i' \beta_i \lambda_i}{\ell(1-\ell)(\omega + \lambda_i)} = 0 \quad (79)$$

The denominator of Eq. (78) may be recognized as exactly the same form as Eq. (79) and since the roots, ω , satisfy the characteristic equation the denominator of Eq. (78) is zero. Therefore, an inverse transform of Eq. (78) is of the form

$$\begin{aligned} \mathcal{L}^{-1}[\bar{n}(\Delta)] = n(t) &= \sum_{j=1}^7 A_j(\Delta_j) e^{\Delta_j t} \\ &= \sum_{j=1}^7 A_j(\omega_j) e^{\omega_j t} \end{aligned} \quad (80)$$

A method, as given by Wylie (33), predicts the coefficients, $A_j(\omega_j)$, by using the following form of the inverse transform

$$\mathcal{L}^{-1}[\bar{n}(\Delta)] = \sum_{j=1}^7 \frac{p(\omega_j)}{q'(\omega_j)} e^{\omega_j t} \quad (81)$$

where $q'(\Delta)$ is given by the equation

$$\begin{aligned} q'(\Delta) \Big|_{\Delta=\omega_j} &= \frac{d}{d\Delta} \left\{ \Delta - \frac{1}{\lambda} \left(\frac{\rho - \beta}{1 - \rho} \right) - \sum_{i=1}^6 \frac{\epsilon_i \epsilon'_i \lambda_i \beta_i}{\lambda(1-\rho)(\Delta + \lambda_i)} \right\} \Big|_{\Delta=\omega_j} \\ &= 1 + \frac{1}{\lambda(1-\rho)} \sum_{i=1}^6 \frac{\epsilon_i \epsilon'_i \lambda_i \beta_i}{(\omega_j + \lambda_i)^2} \end{aligned} \quad (82)$$

Using Eq. (82), the coefficients are determined as

$$A_j(\omega_j) = n(0) (1-e) \left[\frac{\rho + \sum_{i=1}^6 \beta_{eff i}}{\rho(1-e) + \sum_{i=1}^6 \beta_{eff i} \lambda_i} \frac{1}{(\omega_j + \lambda_i)^2} \right] \quad (83)$$

Therefore, the neutron density equation is

$$\frac{n(t)}{n(0)} = \frac{\phi(t)}{\phi(0)} = \sum_{j=1}^7 (1-e) \left[\frac{\rho + \sum_{i=1}^6 \beta_{eff i}}{\rho(1-e) + \sum_{i=1}^6 \beta_{eff i} \lambda_i} \right] e^{\omega_j t} \quad (84)$$

which is called the neutron density ratio or flux ratio.

Part - B

Reactivity as a Function of Time

When the coefficient, k_{eff} , in Eq.'s (40) and (43), remains a constant or is a step function of time, the kinetic differential equations are easily solved. When k_{eff} has a time dependence, the kinetic differential equations become difficult to solve exactly (1). The degree of difficulty depends upon the time function which describes $k_{eff}(t)$ and on the number of delayed neutron groups used in the model.

Glasstone (7) presents an approximate solution of the spatially independent kinetic equations where k_{eff} is a linear function of time. One group of delayed neutrons was assumed in

Glasstone's solution.

A solution to the spatially independent kinetic equations is given in this paper in which $k_{\text{eff}}(t)$ is given by Eq. (14). Eq. (14) is an approximation to the reactivity insertion rate for a rod drop in the TRIGA reactor. For simplicity of solution, it is assumed that there is only one group of delayed neutrons. The usual assumption connected with the simplified kinetics model are made (see Appendix - A). Also, in accordance with an assumption made by McPhee (15), in a similar development, the prompt neutron lifetime was assumed to be approximately zero.

The spatially independent kinetic equations are the same as Eq. (40) and Eq. (43), except that here there is assumed to be only one group of delayed neutrons. Thus Eq.'s (40) and (43) become

$$\left[(1-\beta) k_{\text{eff}}(t) - 1 \right] \frac{n(t)}{\lambda} + \epsilon \lambda H(t) = \frac{dn(t)}{dt} \quad (85)$$

and

$$\frac{dH(t)}{dt} = -\lambda H(t) + \frac{k_{\text{eff}}(t) n(t) \beta \epsilon'}{\lambda} \quad (86)$$

$\lambda H(t)$ is substituted from Eq. (86) into Eq. (85) to obtain

$$\left[(1-\beta) k_{\text{eff}}(t) - 1 \right] \frac{n(t)}{\lambda} + \epsilon \left[\frac{k_{\text{eff}}(t) n(t) \beta \epsilon'}{\lambda} - \frac{dH(t)}{dt} \right] = \frac{dn(t)}{dt} \quad (87)$$

It was also assumed that the effective fraction of delayed neutron constant, $\beta\epsilon'$, is equal to unity. Thus, the two terms cancel in Eq. (87).

$$(k_{\text{eff}}(t) - 1) \frac{n(t)}{\lambda} - \epsilon \frac{dH(t)}{dt} = \frac{dn(t)}{dt} \quad (88)$$

Eq. (86) is then differentiated with respect to time to give

$$\frac{d^2n(t)}{dt^2} = (k_{\text{eff}}(t) - 1) \frac{1}{\lambda} \cdot \frac{dn(t)}{dt} + \frac{n(t)}{\lambda} \frac{dk_{\text{eff}}(t)}{dt} - \epsilon \frac{d^2H(t)}{dt^2} \quad (89)$$

Differentiating Eq. (86) gives

$$\frac{d^2H(t)}{dt^2} = -\lambda \frac{dH(t)}{dt} + \frac{k_{\text{eff}}(t)\beta\epsilon'}{\lambda} \cdot \frac{dn(t)}{dt} \quad (90)$$

Eq. (88) is then solved for $\frac{dH(t)}{dt}$.

$$\frac{dH(t)}{dt} = \left[k_{\text{eff}}(t) - 1 \right] \frac{n(t)}{\lambda\epsilon} - \frac{dn(t)}{dt} \cdot \frac{1}{\epsilon} \quad (91)$$

$\frac{dH(t)}{dt}$ is substituted from Eq.(91) into Eq. (90) to give

$$\frac{d^2H(t)}{dt^2} = -\lambda \left[(k_{\text{eff}}(t) - 1) \frac{n(t)}{\lambda\epsilon} - \frac{dn(t)}{dt} \cdot \frac{1}{\epsilon} \right] \quad (92)$$

Then, from Eq.'s (89) and (92), $\frac{d^2n}{dt^2}$ is found to be

$$\begin{aligned} \frac{d^2n(t)}{dt^2} = & (k_{\text{eff}}(t) - 1) \frac{1}{\ell} \cdot \frac{dn(t)}{dt} + \frac{n(t)}{\ell} \frac{dk_{\text{eff}}(t)}{dt} \\ & + \lambda \left(\frac{k_{\text{eff}}(t) - 1}{\ell} \right) n(t) - \lambda \frac{dn(t)}{dt}. \end{aligned} \quad (93)$$

Eq. (93) is simplified into

$$\begin{aligned} \frac{d^2n(t)}{dt^2} = & \frac{dn(t)}{dt} \cdot \frac{1}{\ell} \left[k_{\text{eff}}(t) - 1 - \lambda \ell - k_{\text{eff}}(t) \beta \right] \\ & + \frac{n(t)}{\ell} \left[\frac{dk_{\text{eff}}(t)}{dt} + \lambda (k_{\text{eff}}(t) - 1) - \beta \frac{dk_{\text{eff}}(t)}{dt} \right]. \end{aligned} \quad (94)$$

Following McPhee, it is assumed that the products, $\frac{d^2n(t)}{dt^2} \ell$ and $\lambda \ell$ are approximately zero, therefore Eq. (94) becomes

$$\left[(1 - \beta) k_{\text{eff}}(t) - 1 \right] \frac{dn(t)}{dt} \doteq -n(t) \left[\lambda (k_{\text{eff}}(t) - 1) + (1 - \beta) \frac{dk_{\text{eff}}(t)}{dt} \right] \quad (95)$$

Since $k_{\text{eff}}(t) \doteq 1 - A(1 - e^{-bt^4})$, (96)

therefore $\frac{dk_{\text{eff}}(t)}{dt} = -4bAt^3 e^{-bt^4}$. (97)

Eq. (95) becomes

$$\begin{aligned} \left[(1 - \beta)(1 - A\{1 - e^{-bt^4}\}) - 1 \right] \frac{dn(t)}{dt} = & -n(t) \left[\right. \\ & \left. \lambda(1 - A\{1 - e^{-bt^4}\}) - \lambda + (1 - \beta)(-4bAt^3 e^{-bt^4}) \right]. \end{aligned} \quad (98)$$

Terms in $n(t)$ are collected to give

$$\frac{1}{n(t)} \cdot \frac{dn(t)}{dt} = \frac{[-(1-e^{-bt^4})\lambda + (1-\beta)(-4bt^3 e^{-bt^4})]}{1-\beta + \frac{\beta}{A} - (1-\beta)e^{-bt^4}} \quad (99)$$

Eq. (99) is integrated from zero time to time t .

$$\int_{n(0)}^{n(t)} \frac{1}{n(t)} dn(t) = - \int_0^t \frac{\lambda dt'}{a_2 + (\beta-1)e^{-bt'^4}} - \int_0^t \frac{(\beta-1)(-4bt'^3) e^{-bt'^4} dt'}{a_2 + (\beta-1)e^{-bt'^4}} + \int_0^t \frac{\lambda dt'}{a_2 e^{+bt'^4} + (\beta-1)} \quad (100)$$

Where a_2 is

$$a_2 = 1 - \beta + \frac{\beta}{A} \quad (101)$$

The integration is performed directly with one of the terms of Eq. (100) to give

$$\int_0^t \frac{(\beta-1)(-4bt'^3) e^{-bt'^4} dt'}{a_2 + (\beta-1)e^{-bt'^4}} = \ln \left[\frac{a_2 + (\beta-1)e^{-bt^4}}{a_2 + (\beta-1)} \right]$$

The last two terms of Eq. (100) are integrated numerically with the IBM-650 using Simpson's Rule. The neutron flux is then represented by

$$n(t) = n(0) \left[\frac{a_2 + (\beta - 1)}{a_2 + (\beta - 1)e^{-bt^4}} \right] e^{-\int_0^t \frac{\lambda dt'}{a_2 + (\beta - 1)e^{-bt'^4}} + \int_0^t \frac{\lambda dt'}{a_2 e^{+bt'^4} + (\beta - 1)}} \quad (102)$$

Part - C

Computer Program for the Kinetic Equations With Six Groups of Delayed Neutrons

General. This program solves the reactor kinetic equations for a bare homogeneous thermal reactor assuming that there are six delayed neutron groups (Appendix, Part - A). The code was written in SOAP-II for the IBM-650 computer. The output data from the code produces the following quantities:

- 1) Coefficients of the characteristic polynomial.
- 2) Roots of the characteristic equation.
- 3) Coefficients of the neutron flux ratio equation.
- 4) The ratio, $\phi(t)/\phi(0)$, for various times after a step insertion of reactivity.

The Coefficients of the Characteristic Equations. The expansion of the characteristic equation may be written as

$$\left[\rho(\lambda\omega + 1) - \lambda\omega \right] = \frac{\beta_1\omega}{\omega + \lambda_1} + \frac{\beta_2\omega}{\omega + \lambda_2} + \dots + \frac{\beta_n\omega}{\omega + \lambda_n} \quad (103)$$

By multiplying both sides of Eq. (97) by $\prod_i (\omega + \lambda_i)$ and collecting terms in powers of ω , the following type equation results

$$\sum_{j=0}^7 H_j \omega^j = 0 \quad (104)$$

Roots of the Characteristic Equation. Once the reactor parameters, λ , ρ , λ_i' , and β_i' are known, the numerical coefficients of the characteristic equation may be calculated and thus the particular characteristic equation is determined. Roots of such an equation may be determined in many different fashions. Since all of the roots of the equation are real and the function, $f(\omega)$, is differentiable, the Newton-Raphson method may be used. The Newton-Raphson method is an iterative procedure and convergence to a root of the polynomial equation is fast if the equation is well behaved.

Educated guesses for the roots are aided by the fact that the characteristic equation has six poles, located at the negative value of each of the decay constants of the precursor groups. Test for proper convergence is simple since the roots must lie in well defined positions. One root must be algebraically larger than the negative value of the smallest decay constant and one root must be algebraically smaller than the negative value of the largest decay constant. Each of the other roots lie between respective poles.

For the roots to lie between two poles, one of the poles is tried for a guess of the root. If proper convergence is not obtained, the other pole is tried as a guess for the root. If proper convergence is still not obtained, the midpoint between the two poles is tried as a guess. From experience with the program, the three guesses are sufficient to extract the desired

root when the root lies between two poles. Since the largest negative root is strongly dependent upon the magnitude of the step change in reactivity, a series of guesses are offered. First, the largest negative pole is tried as a guess and if convergence is not obtained to the proper value, a multiple of the pole magnitude is tried. The largest algebraic root is always algebraically larger than the negative value of the smallest decay constant. If the reactivity insertion is negative, this root will be negative and if the reactivity insertion is positive, this root will be positive.

If the reactivity insertion is negative several guesses are tried, i.e., zero, the largest pole, and the midpoint between zero and the largest pole. If the reactivity is positive an approximation, as given by Glasstone (7), is used as a guess for the root.

$$\omega \doteq \frac{\lambda_{av} \rho}{\beta - \rho} \quad (105)$$

Where λ_{av} is the average decay constant and is defined as the inverse of average mean life of the delayed neutrons, i.e.,

$$\lambda_{av} \equiv \frac{\beta}{\sum_{i=1}^6 \beta_i \frac{1}{\lambda_i}} \quad (106)$$

If convergence to the desired root is not obtained, a multiple of the approximation is used.

Coefficients of the Flux Equations. The coefficients, A_j , of the flux equation

$$\frac{\phi(t)}{\phi(0)} = \sum_{j=1}^7 A_j e^{w_j t} \quad (107)$$

as defined by Eq. (94) are computed in a straight forward manner. Since this is very simple, a discussion of the procedure is not given.

Flux Ratios. Since all of the ω 's and A_j 's have been computed, it is a simple matter to calculate the flux ratios.

Program Operation. If one wishes to compute continuously a set of varying parameters, such as reactivity, the following two cards are added to the program just before the parameter data.

Symbolic			Instruction	Op	Data	Address	
	NZC	COM65	COM80	1683	48	1633	1534
COM80	NOP	0000	8000	1543	00	0000	8000

The position, 1534, will then be removed from the table of availability. The various parameter changes may be added to the end of the data and a transfer card positioned between each data card.

The computer will tend to overflow when larger positive reactivities are used, $\rho > \$1.00$. The overflow is due to the unruly nature of the characteristic polynomial at larger reactivities and no special programming precaution has been used to prevent the overflow in this case. In order to prevent large exponential values from causing the computer to overflow, all values "e" to the power ± 135.0 and larger are omitted. At larger reactivities, $\rho > \$0.60$, the computer will give erroneous results for large times after the step change in reactivity.

This is no real handicap since $\frac{\phi(t)}{\phi(0)} \geq 10^{27}$ has no practical value for the usual reactor operation.

Input Data. The nuclear parameters are placed in the code in the following positions.

Position	Symbol	Parameter	
1384	6	λ_6	(smallest decay constant)
1385	5	λ_5	
1386	4	λ_4	
1387	3	λ_3	
1388	2	λ_2	
1389	1	λ_1	
1374	BETA6	β_6	(the fraction associated with λ_6)
1375	BETA5	β_5	
1376	BETA4	β_4	
1377	BETA3	β_3	
1378	BETA2	β_2	
1379	BETA1	β_1	
0148	BETAT	β	
0239	LAMAV	λ_{av}	
0191	RHO	ρ	
0078	L	Prompt Neutron lifetime, l	

Output DATA. The output data, W'_A , H's, A_j 's, and $\frac{\phi(t)}{\phi(0)}$, are given in the following order:

First line of data - all of the decay constants and the average decay constant.

Second line - All of the delayed neutron fractions and the total fraction.

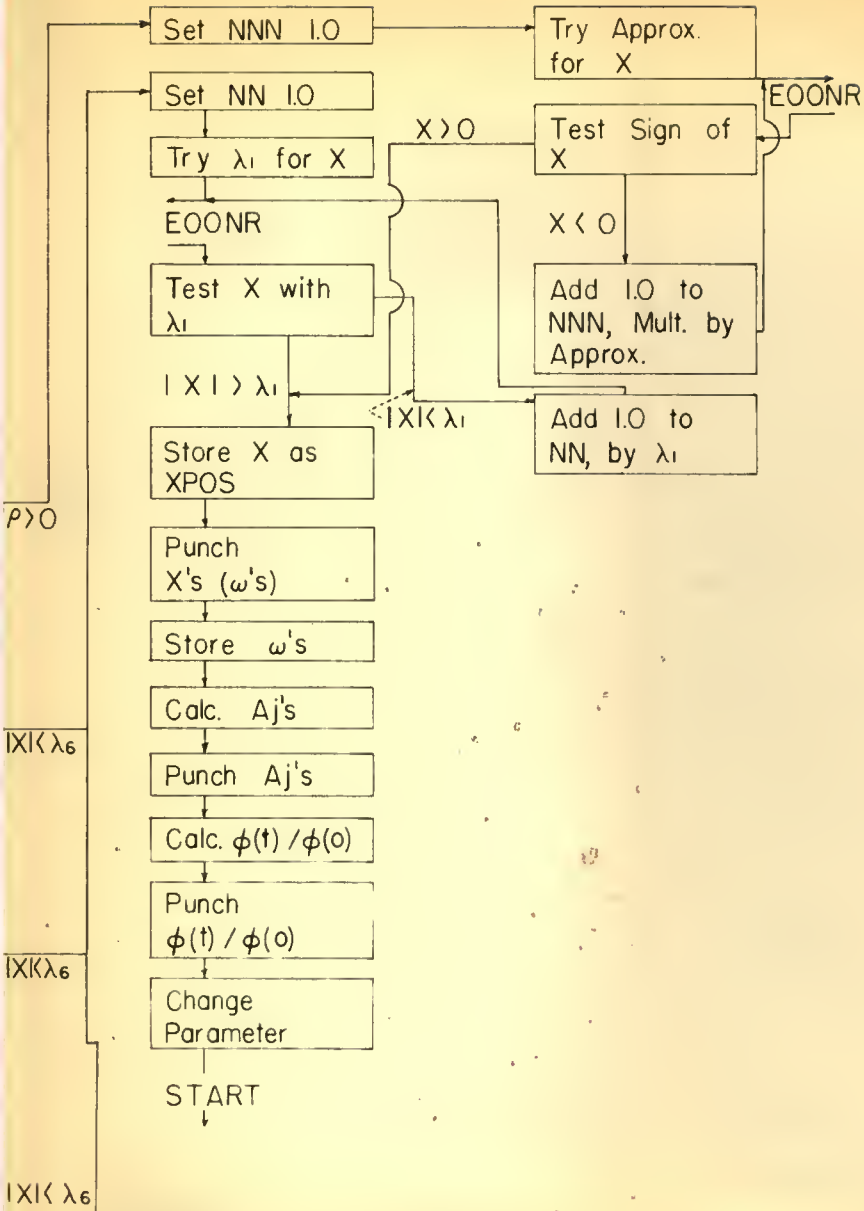
Third line - Reactivity and the prompt neutron lifetime.
Fourth line - The coefficients of the characteristic polynomial.
Fifth line - Roots of the characteristic equation.
Sixth line - Coefficients of the flux equation.
Lines seven through 22 - Flux ratios at the various times.

After the data cards have been punched and positioned in the program deck, the following console settings are used.

Storage Entry Switches	70 1951 1999	plus
Programmed Stop	STOP	
Half Cycle	RUN	
Control	RUN	
Address Selection	Anything	
Display	Anything	
Overflow	SENSE	
ERROR	STOP	

EXPLANATION OF PLATE XLIV

Flow diagram of the computer program
for the solution of the reactor kinetic
equations using six groups of delayed
neutrons.



SIX
 DELAYED
 NEUTRON
 GROUP
 REACTOR
 KINETIC
 EQUATIONS

BLR 0440	0499								
BLR 0500	0599	CORRECT	1	0000	00	0000	0000	0000	
BLR 0800	0899	TRACE	2	0000	00	0000	0000	0000	
BLR 1300	1400		3	0000	00	0000	0000	0000	
BLR 1551	1960	READ	4	0000	00	0000	0000	0000	
BLR 1977	1986	PRINT	5	0000	00	0000	0000	0000	
SYN A	0500		6	0000	00	0000	0000	0000	
SYN AP	0490		7	0000	00	0000	0000	0000	
SYN T	1000		8	0000	00	0000	0000	0000	
SYN T	1350		9	0000	00	0000	0000	0000	
SYN BP J	1360		10	0000					
SYN SJ	1370		11	0000	00	0000	0000	0000	
SYN BETA	1380		12	0000	00	0000	0000	0000	
SYN LAMDA	1390		13	0000	00	0000	0000	0000	
SYN OMEGA	1400		14	0000	00	0000	0000	0000	
SYN START	1999		15	0000	00	0000	0000	0000	
SYN HLT	1111		16	0000	00	0000	0000	0000	
EQOCL	UTO ZZZ1	CONT1	17	0000	00	0000	0000	0000	
CONT1	LBO ZZZ10		18	0000	24	0053	0066		
	STO 1977		19	0006	69	0009	0012		
	STO 1978	READ	20	0012	24	1977	0030		
	STO 1979	DATA	21	0030	24	1978	0031		
	STO 1980	ZERO	22	0031	24	1979	0032		
	STO 1981	SUBROUTINE	23	0032	24	1980	0033		
	STO 1982		24	0033	24	1981	0034		
	STO 1983		25	0034	24	1982	0035		
	STO 1984	ZZZ1	26	0035	24	1983	0036		
EQQNR	STO NRM	CONT2	27	0036	24	1984	0037		
CONT2	REA 0007		28	0050	24	0053	0056		
	RAU A	A CONT3	29	0056	61	0007	0062		
CONT3	FAP X		30	0062	60	2500	0065		
	AXA 0001		31	0005	39	1000	0100		
	NZA LOOP2	CONT4	32	0100	60	0001	0106		
			33	0106	40	0059	0010		

LOOP2	FAO A	A	CON13	34	0059	38	2500	0005	
CONT4	FAD A	A		35	0010	21	0080	0027	
	STU Y			36	0027	22	0082	0085	
	RBB Y	0006		37	0085	83	0006	0041	
	RAU AP		B	CON16	38	0041	60	4490	0045
CON16	FMP X			39	0045	39	1000	0150	
	AXB 0001			40	0130	52	0001	0154	
	NZR L00P3			41	0156	42	0109	0060	
LOOP3	FAO AP		B	CON15	42	0150	32	4490	0045
CONT5	FAO Y P			43	0060	38	0400	0017	
	STU Y P			44	0060	21	0022	0023	
	RSU Y			45	0025	61	0022	0072	
	FOV YP			46	0037	34	0022	0072	
	FAO X			47	0077	32	0000	0077	
	STU XP			48	0077	21	0132	0135	
	RAU X			49	0135	60	1000	0053	
	FOV PREC			50	0055	34	0000	0053	
	STU OIFF			51	0058	21	0112	0015	
	RAU X			52	0058	60	1000	0105	
	FOV PREC			53	0105	34	0008	0108	
	STU OIFF			54	0108	21	0112	0063	
	RAU X			55	0065	60	1000	0059	
	FSP XP			56	0155	33	0132	0159	
	RAM 8003			57	0159	67	8003	0067	
	RAM 8002			58	0067	60	8008	0075	
	F8M OIFF			59	0075	38	0112	0039	
	BMI FIN18			60	0039	46	0042	0043	
ITER	LD0 XIF		ITER	61	0043	69	0132	0185	
	STU X		CONT2	62	0185	24	0103	0056	
	LD0 XIF		NRW	63	0240	65	0132	0053	
FINIB	STL XP			64	0240	20	0011	0014	
EO0EA	STO AAA1		E TO THE X	65	0206	67	0011	0014	
	STL AAA2		SUBROUTINE	66	0114	80	0019	0122	
	RAM AAA2			67	0155	60	0019	0023	
	STL AAA3			68	0128	39	0022	0072	
	RAM AAA3			69	0023	32	0029	0205	
	FMP AAA3			70	0076	39	0019	0069	
	FAO AAA4			71	0205	32	0172	0149	
	FMP AAA4			72	0049	39	0019	0119	
	FAO AAA5			73	0119	38	0022	0039	
	FMP AAA5			74	0099	39	0019	0169	
	FAO AAA6			75	0149	38	0272	0149	
	FMP AAA6			76	0099	39	0019	0169	
	FAO AAA7			77	0149	39	0019	0169	
	FMP AAA7			78	0219	32	0322	0199	
	FAO AAA8			79	0199	39	0019	0169	
	FMP AAA8			80	0249	32	0372	0249	
	FAO AAA9			81	0249	61	0004	0097	
	FMP AAA9			82	0007	39	0004	0057	
	STU AAA9			83	0054	21	0004	0057	
	FMP AAA9			84	0057	24	0004	0104	
	STU AAA9			85	0104	21	0004	0107	
	FMP AAA9			86	0107	60	0021	0165	
AAA5	RAM AAA9		AAA6	87	0165	34	0004	0104	
	FOV AAA9			88	0018	60	0372	0127	
	STU AAA9			89	0127	34	0004	0104	
AAA6	RAM AAA9		AAA1	90	0018	21	0004	0104	
START	RAL AAA4			91	0319	65	0004	0104	
	LD0 1384		LAMBOA 6	92	1999	69	1384	0087	
	LD0 1385		LAMBOA 5	93	0087	24	0040	0093	
	LD0 1386		LAMBOA 4	94	0093	24	0091	0044	
	LD0 1387		LAMBOA 3	95	0038	24	0091	0044	
	LD0 1388		LAMBOA 2	96	0044	24	0092	0095	
	LD0 1389		LAMBOA 1	97	0099	34	0092	0095	
	ST1			98	0095	69	1387	0090	
	ST2			99	0090	69	1388	0141	
	ST3			100	0046	69	1388	0141	
	ST4			101	0141	24	0094	0141	
	ST5			102	0047	69	1389	0142	
	ST6			103	0142	24	0144	0177	
	ST7			104	0048	69	1384	0087	
	ST8			105	0177	24	0080	0083	
	ST9			106	0083	39	1375	0175	
	ST10			107	0028	24	0081	0084	
	ST11			108	0084	69	1376	0079	
	ST12			109	0079	24	0182	0215	
	ST13			110	0215	69	1377	0130	
	ST14			111	0130	24	0130	0081	
	ST15			112	0086	69	1378	0131	
	ST16			113	0131	24	0134	0137	
	ST17			114	0137	69	1379	0132	
	ST18			115	0232	24	0285	0088	
	ST19			116	0088	60	0191	0195	
	ST20			117	0195	33	0098	0125	
	ST21			118	0125	39	0078	0128	
	ST22		H 7	119	0128	21	0282	0135	
	ST23			120	0335	81	0006	0241	
	ST24			121	0241	60	1390	0241	
	ST25			122	0245	50	0001	0001	
	ST26			123	0001	40	0204	0225	
	ST27			124	0204	32	1390	0245	
	ST28			125	0255	21	0110	0013	
	ST29			126	0013	69	0111	0225	
	ST30			127	0225	33	0148	0175	
	ST31			128	0175	21	0180	0183	
	ST32			129	0183	60	0282	0187	
	ST33			130	0187	39	0110	0160	
	ST34			131	0160	24	0180	0177	
	ST35		H 6	132	0157	21	0162	0215	
	ST36			133	0215	60	0048	0073	
	ST37			134	0073	21	0178	0181	
	ST38			135	0181	81	0006	0237	
	ST39			136	0237	60	0237	0145	
	ST40			137	0345	39	1380	0230	
	ST41			138	0230	40	0204	0225	
	ST42			139	0305	21	0178	0231	
	ST43			140	0231	50	0001	0227	
	ST44			141	0227	40	0237	0145	
	ST45			142	0291	60	0069	0133	
	ST46			143	0133	21	0285	0131	
	ST47			144	0281	81	0006	0337	
	ST48			145	0337	39	1390	0140	
	ST49			146	0395	60	1390	0395	
	ST50			147	0140	32	0228	0335	
	ST51			148	0355	21	0228	0335	
	ST52			149	0331	50	0001	0327	
	ST53			150	0327	40	0337	0140	
	ST54			151	0341	60	0110	0245	
	ST55			152	0245	39	0110	0210	
	ST56			153	0210	33	0204	0225	
	ST57			154	0405	34	0158	0208	
	ST58			155	0208	80	008	008	
	ST59			156	0315	60	0282	0437	
	ST60			157	0437	31	0212	0269	
	ST61			158	0269	31	0212	0269	
	ST62			159	0269	60	0100	0365	
	ST63			160	0365	39	0100	0360	
	ST64			161	0280	32	0178	0505	
	ST65			162	0505	21	0216	0193	
	ST66			163	0193	21	0068	0173	
	ST67		H 5	164	0419	60	0068	0173	
	ST68			165	0173	21	0276	0321	

	RSB	0006	CON414	164	0331	H1	0006	0537
CON18	RAU	LAMDA	A	167	0537		60	3390
	FMP	LAMDA	A	168	0545		39	3390
	FMP	LAMDA	A	169	0590		39	3390
	FAD	SUML3		170	0240		32	0278
	STU	SUML3		171	0555		21	0278
	AXA	0001		172	0431		60	0001
	NZA	CON18		173	0567		40	0537
	RAU	ZERO	CON19	174	0391		60	0537
CON19	BTU	SUML2		175	0223		21	0358
	RSB	0006	CON20	176	0376		61	0006
CON20	RAU	LAMDA	A	177	0637		60	3390
	FMP	LAMDA	A	178	0595		39	3390
	FMP	BETA	A	179	0340		39	3390
	FAD	SUML2		180	0330		32	0398
	STU	SUML2		181	0605		21	0398
	AXA	0001		182	0581		60	0001
	NZA	CON20		183	0441		40	0637
	RAU	A02	CON21	184	0687		60	0637
	FBB	SUML2		185	0117		33	0228
	FMP	A02		186	0656		39	0110
	FAD	SUML3		187	0260		32	0278
	FOV	THREE		188	0740		34	0258
	STU	A03		189	0208		39	0312
	RAU	H7		190	0415		60	0282
	FMP	A03		191	0368		39	0312
	STU	H4		192	0368		21	0066
	RAU	A02		193	0259		60	0212
	FMP	RDMLB		194	0167		39	0210
	FAD	H4		195	0380		32	0066
	BTU	H4		196	0243		21	0066
	RAU	A0		197	0569		60	0110
	FMP	SUML1		198	0518		39	0178
	FAD	H4		199	0368		39	0368
	FBB	SUML2		200	0293		33	0368
	STU	H4	H 4	201	0755		39	0166
	RAU	ZERO		202	0619		60	0068
	STU	SUML4		203	0873		60	0068
	RSB	0006	CON22	204	0631		21	0068
CON22	FMP	LAMDA	A	205	0787		60	3390
	FMP	LAMDA	A	206	0645		39	3390
	FMP	LAMDA	A	207	0340		39	3390
	FMP	LAMDA	A	208	0390		39	3390
	FAD	SUML4		209	0540		32	0428
	STU	SUML4		210	1005		21	0428
	AXA	0001		211	0681		60	0001
	NZA	CON22	CON23	212	1037		40	0787
	RAU	ZERO		213	0591		60	0068
	BTU	SUML3		214	0323		21	0368
	RSB	0006	CON24	215	0731		61	0006
CON24	RAU	LAMDA	A	216	1087		60	3390
	FMP	LAMDA	A	217	0695		39	3390
	FMP	LAMDA	A	218	0590		39	3390
	FMP	BETA	A	219	0640		39	3390
	FAD	SUML3		220	0730		32	0528
	BTU	SUML3		221	0058		21	0058
	AXA	0001		222	1438		50	0001
	NZA	CON24		223	1137		40	1087
CON25	RAU	A03	CON25	224	0641		60	0312
	FAD	SUML3		225	0217		32	0878
	FMP	A03		226	0591		39	0110
	FBB	SUML4		227	0310		33	0428
	STU	A04		228	1155		21	0368
	RSB	A02		229	0063		61	0212
	FMP	SUML2		230	0467		39	0228
	FAD	A04		231	0578		32	0368
	FOV	FOUR		232	1187		34	0590
	STU	A04		233	0740		21	0368
	RAU	H7		234	0113		60	0882
	FMP	A04		235	0237		39	0566
	FAD	SUML3		236	0410		32	0568
	STU	H3		237	1205		21	0510
	RAU	A03		238	0163		60	0312
	FMP	RDMLB		239	0317		39	0180
	STU	H3		240	0530		32	0510
	RSB	A02		241	0287		21	0287
	BTU	A03		242	0213		61	0110
	FMP	SUML2		243	0565		39	0565
	FAD	H3		244	0688		32	0510
	STU	H3		245	1437		21	0510
	RAU	A02		246	0263		21	0263
	FMP	SUML1		247	0367		39	0178
	FAD	H3	H 3	248	0678		39	0210
	RAU	ZERO		249	1487		21	0510
	STU	SUML5		250	0313		60	0068
	RSB	0006	CON26	251	0377		61	0068
CON26	RAU	LAMDA	A	252	1031		81	0006
	FMP	LAMDA	A	253	1537		60	3390
	FMP	LAMDA	A	254	0745		39	3390
	FMP	LAMDA	A	255	0790		39	3390
	FMP	LAMDA	A	256	1040		39	3390
	FMP	LAMDA	A	257	1090		39	3390
	FAD	SUML5		258	1268		32	0728
	AXA	0001		259	1255		81	0728
	NZA	CON26		260	1081		50	0001
CON27	RAU	ZERO	CON27	261	0591		50	0001
	STU	SUML4		262	0691		60	0068
	RSB	0006		263	0483		21	0778
CON28	RAU	A02	CON28	264	1131		81	0006
	FMP	LAMDA	A	265	1487		60	3390
	FMP	LAMDA	A	266	1756		39	3390
	FMP	LAMDA	A	267	1190		39	3390
	FMP	LAMDA	A	268	1240		39	3390
	FMP	BETA	A	269	1268		39	3390
	FAD	SUML4		270	1820		39	3380
	STU	SUML4		271	0580		32	0778
	AXA	0001		272	1405		40	0428
	NZA	CON28	CON29	273	1161		50	0001
CON29	RAU	A04		274	0741		60	0360
	FBB	SUML4		275	0615		39	0428
	FMP	A0		276	1455		39	0310
	FAD	SUML5		277	0560		32	0728
	STU	EO		278	0305		21	0305
	RSB	A03		279	0363		61	0312
	FMP	SUML2		280	1280		39	0240
	FAD	EO		281	1028		32	0610
	RAU	A02		282	1737		21	0510
	FMP	SUML3		283	0413		60	0128
	FAD	EO		284	0537		39	0278
	FOV	FIVE		285	1078		32	0128
	STU	EO		286	1797		34	1440
	RAU	H7		287	1490		21	0610
	RSB	0006		288	0133		60	0688
	FMP	SUML4		289	1837		39	0610
	STU	H2		290	0460		33	0778
	FMP	RDMLB		291	1555		21	0710
	RAU	A04		292	0688		39	0360
	FAD	H2		293	0665		39	0360
	STU	H2		294	0630		32	0710
	STU	H2		295	1887		21	0710
	RAU	A03		296	0613		60	0312
	FMP	SUML1		297	0567		39	0178

FAD H2	298	1128	32	0710	1937
FAD H2	299	1937	21	0710	0663
RBU A02	300	0663	61	0212	0617
FMP 0ML82	301	0617	39	0328	1178
FU H2	302	1178	32	0710	1987
RAU A0	303	1987	61	0710	0713
FMP 3ML63	304	0713	60	0110	0715
FAD H2	305	0715	39	0528	1228
FAD H2	306	1228	32	0710	1937
RAU H2	307	1937	21	0710	0763
FMP 2	308	0763	60	0543	0599
FMP 3	309	0599	39	0543	0599
FMP 4	310	0599	39	0143	0343
FMP 5	311	0343	39	0543	0599
FMP 6	312	0192	39	0091	0791
STU H10	313	0791	39	0144	0540
RAU H10	314	1540	21	0100	0763
FDV 1	315	0927	60	0194	0349
STU H11	316	0349	34	0113	0349
RAU H10	317	1045	21	0250	0553
FDV 2	318	0153	60	0194	0349
STU H12	319	0399	34	0094	0244
RAU H10	320	0244	21	0198	0051
FDV 3	321	0051	60	0194	0549
STU H13	322	0549	34	0143	0393
RAU H10	323	0393	21	0248	0212
FDV 4	324	0101	60	0194	0599
STU H14	325	0599	34	0098	0242
RAU H10	326	0242	21	0248	0212
FDV 5	327	0649	60	0194	0599
STU H15	328	0699	34	0091	0241
RAU H10	329	1041	21	0146	0749
FDV 6	330	0749	60	0194	0799
STU H16	331	0799	34	0040	1539
FMP BETA1	332	2590	21	0294	0147
RAU H11	333	0147	30	0285	0139
FMP BETA2	334	0139	39	0250	0300
RAU H11	335	0300	21	0554	0207
FMP BETA3	336	0207	60	0114	0189
STU BH12	337	0189	39	0198	0296
FMP BETA4	338	1605	31	0296	1638
RAU H13	339	1605	60	0133	0188
STU BH13	340	0188	39	0248	0346
FMP BETA5	341	0348	21	0100	0763
RAU H14	342	1655	60	0282	0238
FMP BETA6	343	0238	39	0096	0242
STU BH15	344	0196	21	0350	0203
FMP BETA7	345	0203	60	0248	0238
RAU H15	346	0385	39	0146	0242
FMP BETA8	347	0246	21	0400	0253
STU BH16	348	0253	60	0248	0238
FMP BETA9	349	0435	39	0294	0344
RAU H16	350	0344	21	0398	0151
FMP H7	351	0151	60	0282	0288
STU BH17	352	0288	39	0194	0394
FMP H8	353	0394	33	0254	0121
STU BH18	354	1291	33	0002	0129
FMP H9	355	0129	33	0055	0179
STU BH19	356	0179	33	0550	0227
FMP H10	357	0227	33	0400	0277
STU BH20	358	0277	33	0128	0228
RAU H11	359	0225	21	0680	0233
FAD H13	360	0233	60	0250	0705
FAD H14	361	1705	32	0148	0275
FAD H15	362	0275	32	0248	0325
FAD H16	363	0325	32	0146	0573
FAD H17	364	0573	32	0146	0573
FAD H18	365	0573	32	0146	0573
FAD H19	366	0021	39	0121	1091
STU H1	367	1091	32	0680	0233
RAU H10	368	0257	21	0680	0233
FMP H10	369	0225	60	0121	1095
STU H0	370	0925	39	0304	0307
LDD 01	371	0544	21	0548	0201
STU 1977	372	0201	69	0304	0307
LDD 1	373	0307	24	1977	0730
LDD 2	374	0730	69	0148	0578
STU 1978	375	0578	24	1978	0281
LDD 3	376	1281	69	0094	0197
STU 1979	377	0197	24	1979	0332
LDD 4	378	0332	69	0143	0286
STU 1980	379	0296	24	1980	0333
LDD 5	380	0333	69	0092	0145
STU 1981	381	1145	24	1981	0184
LDD 6	382	0184	69	0091	0594
STU 1982	383	0594	24	1982	0525
LDD 7	384	0525	69	0040	0543
STU 1983	385	0543	24	1983	0136
LDD 8	386	0136	69	0239	0338
STU 1984	387	0338	24	1984	0338
LDD 9	388	0338	71	1984	0337
STU 1977	389	0337	69	0780	0363
LDD BETA1	390	0363	24	1977	1130
STU 1978	391	1030	69	0285	0388
LDD BETA2	392	0388	24	1978	0438
STU 1979	393	1431	69	0178	0438
LDD BETA3	394	0438	24	1979	0382
STU 1980	395	0382	69	0133	0166
LDD BETA4	396	0166	24	1980	0433
STU 1981	397	0433	69	0182	0525
LDD BETA5	398	0585	24	1981	0234
STU 1982	399	0234	69	0081	0284
LDD BETA6	400	0284	24	1982	0635
STU 1983	401	0635	69	0080	0533
LDD BETA7	402	0533	24	1983	0166
STU 1984	403	0336	69	0148	0231
LDD BETA8	404	0231	24	1984	0538
PCH 1977	405	0538	71	1977	0427
STU 1977	406	0377	69	1080	0000
LDD 03	407	1080	69	0583	0286
STU 1977	408	0286	24	1977	1130
LDD RHO	409	0444	69	0191	0644
STU 1978	410	0644	24	1978	1481
LDD L	411	1481	69	0078	1331
STU 1979	412	1331	24	1979	0432
LDD L	413	0432	71	1977	0427
STU 1977	414	0427	69	0882	0685
LDD H6	415	0685	24	1977	1130
STU 1978	416	1180	69	0162	0769
LDD H5	417	0769	24	1978	1361
STU 1979	418	1361	69	0016	0665
LDD H4	419	0665	24	1979	0929
STU 1980	420	0929	69	0066	0719
LDD H3	421	0719	24	1980	0633
STU 1981	422	0633	69	0310	0133
LDD H2	423	0133	24	1981	0334
STU 1982	424	0334	69	0710	1063
LDD H1	425	1063	24	1982	0735
STU 1983	426	0735	69	0680	0683
LDD H0	427	0683	24	1983	0336
STU 1984	428	0336	69	0548	0301
LDD H0	429	0301	24	1984	0588

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	PCH	1977		430	0568	71	1977	0527
	LOO	H7		431	0527	69	0282	0785
	STO	0493		432	0785	84	0493	0347
	LDO	H6		433	0346	69	0162	1015
	STO	0494		434	1015	84	0494	0247
	LDO	H5		435	0247	69	0016	0769
	STO	0495		436	0769	84	0495	0648
	LDO	H4		437	0648	69	0066	1019
	STO	0496		438	1019	84	0496	1049
	LDO	H3		439	1049	69	0510	1113
	STO	0497		440	1113	84	0497	1440
	LDO	H2		441	0550	69	0710	1163
	STO	0498		442	1163	84	0498	1444
	LDO	H1		443	0351	69	0680	0733
	STO	0499		444	0733	84	0499	0108
	LDO	H0		445	0102	69	0548	0401
	STO	0500		446	0401	84	0500	3003
	RESA	0007		447	1010	69	0007	0020
	RBB	0006		448	0209	83	0006	1065
	STU	N SEVEN		449	1065	60	0118	0623
	FMP	A	CON30	450	0623	81	1271	1611
	STU	N		451	1631	39	2300	0600
	AXA	0001		452	0600	81	450	0593
	NZA	0005		453	0593	50	0001	1099
	AXB	0001	CON31	454	1099	40	1522	0353
	RAU	N		455	0152	52	0001	0358
	FSB	ONE		456	0358	60	1278	0783
	BTU	N	CON30	457	0783	33	0288	0375
	RFC	0006	CON32	458	0375	21	1278	1631
	RBL	LAMOA	C	459	0353	89	0006	0259
	LDO	X	E00NR	460	0259	66	7350	1195
	STL	X		461	1195	20	1000	0403
	RAM	LAMOA	C	462	0403	69	0282	0606
	RAU	8002		463	0606	80	7000	0503
	FSM	X		464	0503	87	7350	1844
	BMI	CON33	CON34	465	1245	60	8002	0553
	AXC	0001		466	0553	38	7000	0577
	RAM	LAMOA	C	467	0577	46	1310	1681
	RAU	8002		468	1681	58	0001	0386
	BXC	0001		469	0386	67	7350	1295
	FSM	X		470	1295	60	8002	0603
	BMI	CON34	CON35	471	0603	59	0001	0309
	NZA	0005		472	0309	38	7000	0502
	RBL	LAMOA	C	473	0627	46	1681	1731
	BTU	X	E00NR	474	1731	46	1681	0633
	STL	X		475	0638	66	7350	1445
	RAM	LAMOA	C	476	1445	20	1000	0653
	RAU	8002		477	0653	62	0306	0208
	FSM	X		478	0306	59	0001	0412
	BMI	CON36	CON37	479	0412	60	7000	0502
	AXC	0001		480	0703	67	7350	1495
	RAM	LAMOA	C	481	1495	80	8002	0677
	RAU	8002		482	0753	39	7000	0727
	BXC	0001		483	0677	46	1280	1781
	FSM	X		484	1781	58	0001	0433
	BMI	CON37	CON38	485	0436	67	7350	1645
	AXC	0001		486	1645	60	8002	1008
	RAM	LAMOA	C	487	1008	59	0001	0359
	RAU	8002		488	0359	38	7000	0727
	BXC	0001		489	0727	46	1761	1731
	FSM	X	CON35	490	1781	60	7350	1595
	BMI	CON38	CON39	491	1595	58	0001	0500
	AXC	0001		492	0501	32	7350	0667
	RAM	LAMOA	C	493	0667	34	0158	0408
	RAU	8002		494	0408	81	0328	1114
	BXC	0001		495	1115	66	0512	0717
	FSM	X		496	0717	20	1000	1053
	BMI	CON40	E00NR	497	1003	69	0356	0050
	AXC	0001		498	0356	59	0001	0562
	RAM	LAMOA	C	499	0562	69	0001	0562
	RAU	8002		500	1103	67	7350	1645
	BXC	0001		501	1645	60	8002	1155
	FSM	X		502	1153	38	7000	0777
	BMI	CON39	CON39	503	0777	46	1430	1831
	AXC	0001		504	1831	38	0001	0528
	RAM	LAMOA	C	505	0526	67	7350	1695
	RAU	8002		506	1695	60	8002	1020
	BXC	0001		507	1203	59	0001	0409
	FSM	X		508	0409	38	7000	1027
	BMI	CON40	CON35	509	1027	46	1480	1731
	AXC	0002	CON39	510	1480	34	1000	1831
	RAM	LAMOA	CON42	511	1831	58	0002	0688
	RAU	8002	CON42	512	0688	49	1141	0348
	BXC	0001		513	1141	59	0001	0559
	FSM	X		514	0348	80	0191	1745
	BMI	CON43	CON44	515	1745	46	0698	1149
	AXC	0001		516	1149	60	698	0580
	RAM	LAMOA	C	517	1213	20	1000	1853
	RAU	8002		518	1853	69	0406	0050
	BXC	0001		519	0406	20	0061	0064
	FSM	X	E00NR	520	0064	60	0040	1795
	BMI	CON41	CON46	521	1795	32	0061	0733
	AXC	0001		522	0738	46	1151	0392
	RAM	LAMOA	C	523	1151	66	0040	1843
	RAU	8002		524	1843	20	1000	1402
	BXC	0001		525	1403	69	0306	0050
	FSM	X	E00NR	526	0050	20	0061	0111
	BMI	CON42	CON46	527	0114	60	0040	1895
	AXC	0001		528	1895	32	0061	0798
	RAM	LAMOA	C	529	0798	46	1241	0798
	RAU	8002		530	1241	61	0040	1945
	BXC	0001		531	1945	34	1000	1755
	FSM	X		532	0508	21	1000	1453
	BMI	CON43	CON46	533	1453	60	0040	1795
	AXC	0001		534	1755	80	0000	1503
	RAM	LAMOA	C	535	1503	69	0556	0050
	RAU	8002		536	0050	60	0061	0166
	BXC	0001		537	0164	60	0040	1995
	FSM	X		538	1995	32	0061	1038
	BMI	CON39	CON46	539	1038	46	1831	0392
	AXC	0001		540	1149	69	0098	0551
	RAM	LAMOA	C	541	0551	34	1354	0351
	RAU	8002		542	0357	24	0760	1263
	BXC	0001		543	1263	60	0239	0643
	FSM	X		544	0643	39	0191	1291
	BMI	CON44	CON46	545	1291	21	1350	1533
	AXC	0001		546	1533	60	0148	1246
	RAM	LAMOA	C	547	1603	33	0191	0767
	RAU	8002		548	0767	81	0422	0245
	BXC	0001		549	0428	60	1350	1805
	FSM	X		550	1805	34	0061	0226
	BMI	CON45	CON48	551	0522	21	0126	0229
	AXC	0001		552	0229	65	0126	1881
	RAM	LAMOA	C	553	1881	30	1000	1755
	RAU	8002		554	1653	69	0606	0050
	BXC	0001		555	0050	60	0061	0214
	FSM	X	E00NR	556	0214	60	0061	1165
	BMI	CON49	CON46	557	1165	46	0168	0398
	AXC	0001		558	0398	60	0354	0303
	RAM	LAMOA	C	559	0509	32	0058	0525
	RAU	8002		560	0525	21	0354	0407
	BXC	0001		561	0407	39	0126	0176

	FMP N7N			862	0176	39	0354	0404
	STU S	C044B		863	0404	31	0000	0239
CON46	RSC 0001			864	0392	69	0001	0748
	RSV LAMDA	C		865	0748	61	7350	0396
	BTU X		C045D	866	0396	21	1000	703
CON50	LDD		EODNR	867	1703	69	0656	0050
	STL X	C		868	0656	20	7000	753
	RAM LAMDA	C		869	1753	67	7350	0546
	RAU B002	C		870	0546	60	8008	1855
	FBM X	C		871	1855	38	7000	0777
CON52	BW1 CON51	C	C045B	872	1077	46	1580	1031
	RAU ONE			873	1931	60	0058	1803
	FAD NN			874	1803	35	0760	1088
	STU MN			875	1088	21	0760	613
	FMP LAMDA	C		876	1413	39	7350	1640
	FMP NN			877	1640	39	0760	1010
	RBL B003			878	1010	66	8003	0000
CON51	STL X	C	C045D	879	1017	20	1000	1703
	LDD O4		C0454	880	1580	69	1033	0000
CON54	STO 1977			881	1033	69	0586	0889
	RSC 0006			882	0889	24	1977	1630
	LDD X	C		883	1630	69	0006	0600
	STO 1978			884	0636	69	7000	1853
	STO 1393			885	1853	24	1978	0586
	AXC 0001			886	0586	24	1393	0596
	LDD X	C		887	0596	58	0001	0202
	STO 1979			888	0202	24	1394	0607
	AXC 0001			889	1903	24	1979	0632
	STO 1980			890	0632	24	1394	0607
	STO 1395			891	0297	58	0001	0504
	AXC 0001			892	0504	69	7000	0554
	LDD X	C		893	0544	24	1980	1031
	STO 1981			894	1083	24	1395	0798
	AXC 0001			895	0798	58	0001	0607
	LDD X	C		896	0604	69	7000	0654
	STO 1982			897	0654	24	1981	1034
	STO 1396			898	0384	24	1396	1199
	AXC 0001			899	1199	58	0001	1905
	LDD X	C		900	1905	69	7000	0700
	STO 1982			901	0704	24	1982	1035
	STO 1399			902	1035	24	1397	0650
	AXC 0001			903	0650	38	0003	0700
	LDD X	C		904	0706	69	7000	0754
	STO 1983			905	0754	24	1983	0686
	STO 1398			906	0686	24	1398	0601
	LDD X P09			907	0601	69	0061	0264
	STO 1984			908	0408	24	1984	1138
	STO 1399			909	1138	24	1399	0852
	PCH 1977			910	0852	71	0558	0000
	RBB 0007	C	C0455	911	1127	83	0007	1133
	LDD ZERO		0	912	1133	69	0006	0071
CON55	STO SUM1			913	0071	24	0054	1177
	STO SUM2			914	1177	24	1680	1183
	RBA 0006			915	1183	31	0006	0333
	RAU OMEGA	B		916	0339	60	5400	0756
CON56	FAD LAMDA	A		917	0756	32	3390	1067
	STU TEMP1			918	1067	21	0572	0575
	RAU BETA	A		919	0575	60	3380	1085
	FVY TEMP2			920	1085	34	0572	0628
	FAD SUM1			921	0628	32	0024	0651
	STU SUM1			922	0651	21	0024	1287
	RAU BETA	A		923	1227	60	3380	1125
	FVY TEMP1			924	1125	34	0572	0672
	FVY TEMP2			925	0672	34	0572	0723
	FMP LAMDA	A		926	0722	39	3390	1690
	FAD SUM2			927	1690	32	1680	0800
	STU SUM2			928	0800	21	1680	1233
	AXA 0001			929	1233	30	0001	0389
	NZA C0456	C	C0457	930	0389	40	0339	0624
	RAU ONE			931	0693	60	0098	1004
	TSB B056			932	1004	33	0191	1117
	STU TEMP1			933	1117	21	0572	0625
	FMP L			934	0625	39	0078	1428
	FAO BUM2			935	1428	32	1680	0800
	STU TEMP2			936	0800	21	0612	1215
	RAU L			937	1215	60	0078	1887
	FAO SUM1			938	1883	38	0024	0701
	FMP TEMP2			939	0701	39	0572	0772
	FVY TEMP2			940	0772	34	0612	0668
	STU BPJ	B		941	0662	21	5360	1463
	ASB 0001			942	1463	32	0001	1069
CON58	NZB C0455	C	C0458	943	1069	42	1133	0673
	LDD ZERO			944	0673	69	0068	0121
	STO SUMBJ			945	0121	24	0074	1277
	RBA 0007			946	1277	21	0007	1433
CON59	RAU SUMBJ	A	C0459	947	1433	60	0074	0277
	FAO BPJ	A		948	0277	32	3360	1188
	STU SUMBJ			949	1188	21	0074	1429
	AXA 0001			950	1429	30	0001	1403
CON60	NZA C0459	C	C0460	951	1403	40	1433	1238
	RAU BETA	A		952	1238	31	0007	0669
CON61	RAU BPJ	A	C0461	953	0664	60	3360	1825
	FVY SUMBJ			954	1825	34	0074	0124
	STU BUJ	A		955	0124	21	370	0723
	AXA 0001			956	0723	50	0001	0329
	NZA C0461	C	C0462	957	0329	40	0664	1533
CON62	LDD		EODCL	958	1533	69	0736	0000
	LDD OS		C0463	959	0000	69	0439	0542
CON63	STO 1977			960	0542	24	1977	0600
	RBA 0007			961	1730	21	0007	0786
	LDD B001	A		962	0786	69	370	0773
	STO 1978			963	0773	24	1978	0682
	AXA 0001			964	0682	50	0001	1288
	LDD B001	A		965	1288	69	370	0833
	STO 1979			966	1023	24	1979	0732
	AXA 0001			967	0732	60	0001	1438
	LDD BU	A		968	1438	69	3370	1073
	STO 1980			969	1073	24	1980	1583
	AXA 0001			970	1583	50	0001	0509
	LDD BU	A		971	0509	69	3370	1123
	STO 1981			972	1123	24	1981	1163
	AXA 0001			973	0434	30	0001	1740
	LDD B001	A		974	1740	69	370	1177
	STO 1982			975	1173	24	1982	1185
	AXA 0001			976	1185	50	0001	1441
	LDD BU	A		977	1441	69	370	1233
	STO 1983			978	1223	24	1983	1036
	AXA 0001			979	1036	50	0001	0924
	LDD BU	A		980	0928	69	3370	1273
	STO 1984			981	1273	24	1984	1489
CON64	PCH 1977		C0464	982	1488	24	1977	1489
	RSC 0015	C	C0465	983	1477	89	0015	1633
CON65	RBA 0007			984	1633	31	0007	0584
	LDD ZERO			985	0589	69	0068	0171
	STO SUMFX		C0466	986	0171	24	0174	1527
	RAU OMEGA	A		987	1527	62	3400	1006
CON66	FMP T	C		988	1006	39	7350	0700
	STU W			989	0700	21	1054	0600
	RAM 0003			990	0607	67	8003	1415
	RAU 0003			991	1415	60	8008	1428
	FBM 135			992	1428	38	0206	1104
	BW1 C0470	C	C0471	993	1104	46	0657	0558

CON70	RAL	#			694	0657	65	1054	0559
	LOD			EODEA	695	0559	69	0712	0200
	RAU	8002			696	0712	60	0002	0201
	FMS	8J	A		697	0221	39	3370	0200
	FAD	SUMFX			698	0020	32	0174	0731
CON71	STU	SUMFX		CON71	699	0751	25	0174	0358
	AXA	00D1			700	0588	30	0001	0314
CON67	NZA	CON66		CON67	701	0314	40	1527	0218
	LOD			EODDL	702	0218	69	0271	0000
	LDD	T	C		703	0271	69	7350	1154
	STO	1977			704	1154	24	1977	1790
	LOD	SUMFX			705	1780	59	0174	1570
	STO	1978			706	1577	24	1978	0782
	PCM	1977			707	0782	71	1977	1607
	AXC	00D1			708	1627	58	0001	1603
	NZC	CON65		CON39	709	1603	48	1633	3831
CON39	HLL				710	1831	01	1038	1032
ZER0	00	0000	0051		711	0068	00	0000	0031
ONE	10	0000	0051		712	0088	10	0000	0041
TWO	20	0000	0051		713	0158	20	0000	0051
THREE	30	0000	0051		714	0288	30	0000	0051
FOUR	40	0000	0051		715	0690	40	0000	0051
FIVE	50	0000	0051		716	1440	50	0000	0051
SIX	60	0000	0051		717	0750	60	0000	0051
SEVEN	70	0000	0051		718	0118	70	0000	0051
01	01	0000	0000		719	0304	01	0000	0000
02	02	0000	0000		720	0780	02	0000	0000
03	03	0000	0000		721	0583	03	0000	0000
04	04	0000	0000		722	0386	04	0000	0000
05	05	0000	0000		723	0439	05	0000	0000
06	06	0000	0051		724	0008	10	0000	0053
07	07	0000	0000		725	0009	00	0000	0000
08	08	0000	0051		726	0372	10	0000	0051
09	09	0000	0051		727	0322	24	9938	6850
10	10	0000	0051		728	0272	31	2575	8349
11	11	0000	0047		729	0222	35	9137	1248
12	12	0000	0047		730	0172	17	1566	0047
13	13	0000	0044		731	0028	54	3028	0045
14	14	0000	0053		732	0286	69	0600	0044
15	15	0000	0053		733	0226	13	5000	0053
16	16	0000	0053			1335	20	0000	0053
17	17	0000	0053			1336	18	0000	0053
18	18	0000	0052			1337	10	0000	0053
19	19	0000	0052			1338	75	0000	0052
20	20	0000	0052			1339	50	0000	0052
21	21	0000	0052			1340	40	0000	0052
22	22	0000	0052			1341	30	0000	0052
23	23	0000	0052			1342	20	0000	0052
24	24	0000	0052			1343	15	0000	0052
25	25	0000	0051			1344	10	0000	0052
26	26	0000	0051			1345	80	0000	0051
27	27	0000	0051			1346	60	0000	0051
28	28	0000	0051			1347	50	0000	0051
29	29	0000	0051			1348	40	0000	0051
30	30	0000	0051			1349	20	0000	0051

Part - D

Three Delayed Neutron Group Program

General. The basic principles of the computer program used to solve the three group kinetic equations were the same as those used in solving the equations for six delayed neutron groups. The program was simplified considerably since the coefficients of the characteristic polynomial were greatly simplified. Since the roots of the characteristic equation were farther apart, in magnitude, additional trials for the roots between the poles of the equation were required in order that the Newton-Raphson method would converge to the proper root.

Input Data. The required input data is:

Position	Parameter
1387	λ_3 (smallest decay constant)
1388	λ_2
1389	λ_1
1377	β_3 (Associated with the smallest decay constant)
1378	β_2
1379	β_1
0141	Reactivity, ρ
0028	Prompt neutron lifetime, ℓ
0236	Average decay constant, λ_{av}
0020	Total fraction of delayed neutrons, β

Output DATA. The output data appears in the same order as in the six group program.

THREE
 DELAYED
 NEUTRON
 GROUP
 REACTOR
 KINETIC
 EQUATIONS

BLR 0440	0499		1	0000	00 0000	0000
BLR 0900	0999	CORRECT	2	0000	00 0000	0000
BLR 0800	0899	TRACE	3	0000	00 0000	0000
BLR 1300	1400		4	0000	00 0000	0000
BLR 1951	1960	READ	5	0000	00 0000	0000
BLR 1977	1986	PRINT	6	0000	00 0000	0000
BYN A	0500		7	0000	00 0000	0000
BYN AP	0490		8	0000	00 0000	0000
BYN I	1000		9	0000	00 0000	0000
BYN T	1350		10	0000	00 0000	0000
BYN BPJ	1360		11	0000	00 0000	0000
BYN B	1370		12	0000	00 0000	0000
BYN BETA	1380		13	0000	00 0000	0000
BYN LAMDA	1390		14	0000	00 0000	0000
BYN OMEGA	1400		15	0000	00 0000	0000
BYN START	1999		16	0000	00 0000	0000
BYN HLT	1111		17	0000	00 0000	0000
BYN ZZZ1		READ	18	0000	24 0003	0006
EOOCL		OUT	19	0006	69 0009	0018
CONT1	LDD ZZZ10	DATA	20	0018	24 1977	0030
	BYN 1977	ZERO	21	0030	24 1978	0031
	BYN 1978	SUBROUTINE	22	0031	24 1979	0032
	BYN 1979		23	0032	24 1980	0033
	BYN 1980		24	0033	24 1981	0034
	BYN 1981		24	0034	24 1982	0035
	BYN 1982		25	0035	24 1983	0036
	BYN 1983		26	0035	24 1984	0037
EOOEA	BYN 1984	ZZZ1	27	0036	24 1984	0037
	BYN AAA1	E TO THE X	28	0036	24 0033	0036
	STL AAA2	SUBROUTINE	29	0036	20 0011	0014
	RAM AAA2		30	0014	57 0011	0018
	STL AAA3		31	0015	20 0019	0022
	RAM AAA3		32	0022	60 0019	0023
	FMP AAA3		33	0023	39 0023	0023
	FAD AAA15		34	0076	32 0029	0005
	FMP AAA2		35	0005	39 0019	0046
	FAD AAA14		36	0069	32 0072	0049
	FMP AAA1		37	0049	39 0019	0119
	FAD AAA13		38	0119	32 0119	0265
	FMP AAA3		39	0099	39 0019	0169
	FAD AAA12		40	0169	32 0172	0145
	FMP AAA3		41	0149	39 0019	0219
	FAD AAA11		42	0219	32 0222	0199
	FMP AAA3		43	0199	39 0019	0265
	FAD AAA10		44	0265	32 0272	0249
	STU AAA1		45	0249	21 0004	0000
	FMP AAA4		46	0007	39 0004	0054
	STU AAA4		47	0054	21 0004	0057
	FMP AAA4		48	0057	39 0004	0104
	STU AAA4		49	0104	21 0004	0107
	FMP AAA2		50	0107	60 0011	0065
AAA5	RAM AAA5	AAA6	51	0065	46 0008	0319
	STU AAA10		52	0018	60 0272	0227
	FOV AAA4		53	0027	34 0004	0154
	STU AAA4	AAA6	54	0154	21 0004	0319
AAA6	RAM AAA4	AAA1	55	0319	65 0004	0053
EOONR	STO NRM	CONT2	56	0100	84 0103	0106
CONT6	RSA 0004	A CONT3	57	0106	81 0004	0062
	BYN I	A CONT3	58	0062	60 2300	0053
CONT3	FMP X	A CONT3	59	0055	39 1000	0150
	BYN 0001	A CONT4	60	0150	50 0000	0156
	NZA LOOP2	A CONT4	61	0156	40 0059	0010
LOOPS	FAD A	A CONT3	62	0059	32 2800	0055
CONT4	BYN X	A	63	0010	32 2800	0077
	STU Y		64	0077	21 0082	0085
	BYN 0003	B CONT6	65	0085	83 0003	0041
	RAM AP	B CONT6	66	0041	60 4490	0045
	FMP X		67	0045	39 1000	0200
CONT6	BYN 0001		68	0200	62 0001	0206
	NZR LOOP3	B CONT5	69	0206	42 0109	0060
LOOPS	FAD AP	B CONT6	70	0109	32 4490	0045
CONT5	RAM AP		71	0060	24 0000	0017
	STU YP		72	0017	21 0322	0025
	BYN Y		73	0025	31 0028	0037
	FOV YP		74	0037	34 0322	0028
	FAD X		75	0037	32 1000	0127
	STU XP		76	0127	21 0132	0135
	RAM X		77	0135	60 1000	0106
	FOV PREC		78	0106	44 0008	0056
	STU OIFF		79	0058	21 0112	0115
	RAM X		80	0115	60 1000	0156
	FOV PREC		81	0156	44 0008	0108
	STU OIFF		82	0108	21 0112	0165
	RAM X		83	0165	60 1000	0156
	FBB XP		84	0205	33 0132	0159
	BYN 8003		85	0159	67 8003	0047
	RAM 8002		86	0067	60 8002	0075
	FSM OIFF		87	0075	38 0112	0039
	BYN FINIB	ITER	88	0039	46 0042	0043
ITER	LOO XP	CONT2	89	0043	69 0132	0185
	BYN X	NRM	90	0185	24 1000	0106
	RAM XP		91	0042	65 0102	0103
FINIS	LOO 1387	LAMBDA3	92	1999	69 1387	0040
START	STO 1388	LAMBDA2	93	0040	24 0093	0046
	LOO 1388	LAMBDA2	94	0046	69 1388	0091
	STO 1389	LAMBDA1	95	0091	24 0044	0047
	LDD 1389	LAMBDA1	96	0047	69 1389	0092
	BYN 1		97	0092	24 0092	0048
	LOO 1377		98	0048	69 1377	0080

STU BETA3		99	0080	24	0083	0086
LDD 1378		100	0086	69	1378	0081
STO BETA2		101	0081	24	0084	0087
LDD 1379		102	0087	69	1379	0116
STO BETA1		103	0182	24	0235	0038
RAU RHO		104	0035	60	0141	0114
FBB ONE		105	0145	33	0098	0128
FMP L		106	0128	39	0088	0078
STU H4	H4	107	0178	31	0092	0075
RAU 2		108	0285	60	0095	0299
FAD 2		109	0289	32	0044	0079
FAD 3		110	0021	32	0093	0369
STU AD		111	0369	31	0044	0177
RAU AD		112	0177	60	0024	0079
FMP H4		113	0079	39	0232	0282
FAD RHO		114	0282	32	0141	0177
F88 BETA2		115	0177	33	0020	0097
STU H3	H3	116	0097	21	0020	0235
RAU 1		117	0255	60	0095	0349
FMP 2		118	0349	39	0044	0094
STU 12		119	0094	21	0148	0094
RAU 1		120	0001	60	0095	0399
FMP 3		121	0399	39	0093	0144
STU 13		122	0143	21	0198	0051
FMP 3		123	0051	60	0044	0549
RAU 23		124	0549	29	0093	0101
STU 23		125	0193	21	0248	0101
FAD 13		126	0101	60	0148	0153
FAD 23		127	0153	32	0198	0375
STU A1		128	0175	32	0448	0285
RAU BETA1		129	0285	21	0130	0335
FMP 2		130	0133	60	0235	0089
STU B12		131	0089	39	0044	0144
RAU 1		132	0144	21	0298	0151
FMP 1		133	0151	60	0084	0139
STU B21		134	0139	39	0139	0189
RAU BETA3		135	0195	21	0250	0203
FMP 1		136	0203	60	0083	0133
STU B31		137	0137	39	0095	0245
RAU BETA1		138	0245	21	0300	0253
FMP 3		139	0253	60	0231	0189
STU B13		140	0189	39	0093	0243
RAU BETA2		141	0243	21	0348	0102
FMP 3		142	0201	60	0084	0239
STU B23		143	0239	39	0093	0293
RAU BETA3		144	0293	21	0398	0187
FMP 2		145	0251	60	0083	0187
STU B32		146	0187	39	0044	0151
RAU 4		147	0194	21	0548	0301
FMP H4		148	0301	60	0330	0335
STU H2		149	0335	39	0044	0349
RAU AD		150	0339	21	0136	0289
FMP RHO		151	0289	60	0089	0075
STU H2		152	0189	39	0141	0191
FBB B12		153	0191	38	0136	0075
FBB B21		154	0075	33	0298	0075
FBB B31		155	0275	33	0250	0227
FBB B11		156	0227	33	0300	0177
FBB B33		157	0277	33	0348	0325
FBB B23		158	0325	33	0398	0177
FBB B32		159	0375	33	0548	0425
STU H3	H2	160	0425	21	0136	0339
RAU B12		161	0339	60	0339	0343
FMP 3		162	0303	39	0093	0343
STU B23		163	0343	21	0598	0133
RAU B21		164	0351	60	0250	0305
FMP 3		165	0305	39	0093	0393
STU B213		166	0393	21	0648	0177
RAU B31		167	0401	60	0300	0355
FMP 2		168	0355	39	0044	0144
STU B312		169	0244	21	0698	0501
RAU 1		170	0501	60	0095	0299
FMP 3		171	0299	39	0044	0177
STU A2		172	0294	39	0093	0543
RAU A2		173	0543	21	0748	0335
FMP H4		174	0551	60	0748	0353
STU H1		175	0353	39	0238	0189
RAU A1		176	0389	60	0130	0385
FMP RHO		177	0389	39	0144	0177
FAD H1		178	0241	32	0186	0063
FBB B123		180	0063	33	0598	0225
FBB B213		181	0598	33	0648	0275
FBB B312		182	0575	33	0698	0625
STU H1	H1	183	0625	21	0186	0439
RAU A2		184	0439	60	0748	0403
FMP RHO		185	0403	39	0141	0229
STU H2	HO.	186	0291	21	0086	0649
LDD	EOOCL	187	0649	69	0558	0000
LDD 01	PUNCH	188	0000	69	0405	0158
STO 1977		189	0188	24	1977	0180
LDD 1		190	0180	69	0095	0798
LDD 2		191	0191	24	0798	0131
LDD 3		192	0131	69	0044	0147
STO 1980		193	0147	24	1979	0432
LDD LAMW		194	0432	69	0093	0146
STO 1981		195	0146	24	1980	0183
STO 1981		196	0183	29	0236	0335
PCB 1977		197	0539	24	1981	0134
LDD	EOOCL	198	0134	71	1977	0329
LDD 02	PUNCH	199	0327	69	0230	0000
STO 1977		200	0230	69	0233	0286
LDD BETA1		201	0286	69	0235	0088
STO 1978		202	0280	69	0235	0088
LDD BETA2		203	0088	24	1978	0088
STO 1979		204	0181	69	0084	0237
LDD 1979		205	0237	24	1979	0339
STO 1980		206	0339	69	0083	0299
LDD BETA1		207	0336	24	1980	0283
STO 1981		208	0283	29	0200	0779
PCB 1977		209	0073	24	1981	0184
LDD	EOOCL	210	0184	71	1979	0329
LDD 03	RHO AND L	211	0377	69	0330	0000
STO 1977		212	0330	69	0333	0386
LDD RHO		213	0386	24	1978	0231
LDD L 1978		214	0380	69	0141	0344
STO 1979		215	0344	24	1978	0231
PCB 1977	EOOCL	216	0231	69	0022	0589
LDD		217	0281	24	1979	0589
LDD 04		218	0589	71	1977	0329
STO 1977		219	0487	69	0430	0000
LDD		220	0000	69	0383	0436
STO 1977		221	0436	24	1977	0530
LDD H4		222	0530	69	0632	0435
STO 1978		223	0435	24	1978	0331
LDD H3		224	0331	69	0002	0505
STO 1979		225	0505	24	1979	0632
LDD H2		226	0632	69	0136	0589
STO 1980		227	0589	24	1980	0433
LDD H1		228	0433	69	0186	0288
STO 1981		229	0639	24	1981	0834
LDD HO		230	0834	69	0096	0699

STD 1982			231	0699	24 1982	0535
PEW 1977			232	0535	71 1977	0587
LDO H			233	0537	69 0032	0585
STD 0496			234	0585	24 0496	0749
LDO H			235	0749	69 0002	0555
STD 0497			236	0555	24 0497	0750
LDO H			237	0350	69 0136	0689
STD 0498			238	0689	24 0498	0600
LDO H1			239	0601	69 0186	0739
STD 0499			240	0739	24 0499	0709
LDO HD			241	0102	69 0092	0500
STD 0500			242	0795	24 0500	0700
RBA 0004			243	0503	24 0004	0543
RBB 0003			244	0209	83 0003	0515
RAU FOUR			245	0215	60 0006	0123
STD N	CON30		246	0123	81 0128	0381
FMA A	A	FUNCTION	247	0381	39 2500	0400
STU A	B	*	248	0448	81 4490	0500
AXA 0001			249	0593	50 0001	1049
XA L00P5	CON31		250	1049	50 1049	0555
AXB 0001			251	0152	52 0001	0808
RAU 0001			252	0208	60 0288	0533
FSS ONE			253	0533	33 0098	0675
STU N	CON30		254	0675	21 0128	0381
RSC 0003	CON32		255	0533	39 0003	0259
RBL LAMDA	C	TRY LAMRDA FOR OMEGA	256	0259	66 7390	0295
LDO	E00NR		257	0295	20 1000	0600
STL X	C		258	0600	67 7390	0345
RAM LAMDA	C		259	0256	20 7000	0653
RAU 8002			260	0653	60 8002	0703
FSM X	C		261	0345	38 7000	0577
BMI CON33	CON34		262	0703	46 0431	0531
AXC 0001			263	0580	58 0001	0536
RAM LAMDA	C		264	0665	67 7390	0329
RAU 8002			265	0395	60 8002	0753
SAC 0001			266	0665	59 0001	0509
FSM X	C		267	0753	38 7000	0627
BMI CON34	CON35		268	0627	46 0431	0531
AXC 0001			269	0431	58 0001	0533
RBL LAMDA	C		270	0307	66 7390	0545
STL X	E00NR		271	0307	20 1000	1000
LDO			272	1003	69 0006	0100
SXC			273	0807	59 0001	0162
STL X 0001			274	0306	67 7390	0259
RAM LAMDA	C		275	0146	60 8002	1103
RAU 8002			276	0776	67 7390	0595
FBM X	C		277	0595	60 8002	1103
BMI CON36	CON37		278	1103	46 0630	0581
AXC 0001			279	0630	58 0001	0581
RAM LAMDA	C		280	0581	67 7390	0645
RAU 8002			281	0586	60 8002	1153
SXC 0001			282	0645	67 7390	0645
FSM X	C		283	0645	60 8002	1153
BMI CON37	CON35		284	1153	59 0001	0359
AXC 0001			285	0359	38 7000	0727
RAM LAMDA	C		286	0727	46 0431	0533
RAU 8002			287	0581	60 7390	0695
SXC 0001			288	0695	58 0001	0653
FSM X	C		289	0653	71 7190	0567
BMI CON38	CON35		290	0567	34 0070	0220
AXC 0001			291	0220	81 0074	0777
RAM LAMDA	C		292	0120	66 0074	0179
RAU 8002			293	0777	20 1000	0600
SXC 0001	E00NR		294	0179	60 7390	0356
STL X			295	1503	59 0001	0918
RAM LAMDA	C		296	0356	67 7390	0745
RAU 8002			297	0745	65 8002	1403
FSM X	C		298	1403	58 7000	1027
BMI CON38	CON70		299	1027	46 0680	0631
AXC 0001			300	0680	58 0001	0631
RAM LAMDA	C		301	0631	67 7390	0795
RAU 8002			302	0795	60 8002	1450
SXC 0001			303	1453	59 0001	0409
FSM X	C		304	0409	38 7000	1077
BMI CON70	CON35		305	1077	46 0431	0533
AXC 0001			306	0631	58 0001	0337
RAM LAMDA	C		307	0337	60 7390	1045
RAU 8002			308	1045	52 0074	0701
FAD D			309	0701	34 0070	0170
FOV TWD			310	0170	60 0124	1127
STU M			311	1127	66 0124	0229
RSL K			312	0229	20 1000	1503
STL X	E00NR		313	1503	69 0406	0100
LDO			314	0406	59 0001	0262
SXC 0001			315	0262	60 7000	1503
STL X	C		316	1503	67 7390	1095
RAM LAMDA	C		317	1095	60 8002	1600
RAU 8002			318	1600	38 7000	1177
FBM X	C		319	1177	46 0730	0681
BMI CON71	CON72		320	0730	58 0001	0681
AXC 0001			321	0681	67 7390	1145
RAM LAMDA	C		322	0686	60 8002	1650
RAU 8002			323	1145	59 0001	0509
SXC 0001			324	0509	38 7000	1227
FSM X	C		325	1227	46 0681	0533
BMI CON72	CON35		326	0681	60 7390	1195
AXC 0001			327	1195	52 0074	0795
RAM LAMDA	C		328	0751	34 0070	0220
RAU 8002			329	0220	21 0174	1277
FAD D			330	1277	66 0174	0279
FOV TWD			331	0279	20 1000	1703
STU M			332	1703	60 1000	0100
RSL K	E00NR		333	0506	20 7000	1753
STL X	C		334	1753	67 7390	1245
RAM LAMDA	C		335	1245	60 8002	1803
RAU 8002			336	1803	38 7000	1427
FBM X	C		337	1427	46 0780	0733
BMI CON73	CON39		338	0780	58 0001	0736
AXC 0001			339	0736	67 7390	1295
RAM LAMDA	C		340	1295	60 8002	1850
RAU 8002			341	1850	59 0001	0559
SXC 0001			342	0559	38 7000	1477
FSM X	C		343	1477	46 0731	0531
BMI CON39	CON35		344	0531	69 0284	0000
AXC 0001	E00CL		345	0000	24 1977	1030
LDO			346	1903	71 1977	0535
STU 1977			347	0535	69 0002	0555
PCN 1977			348	0583	58 0002	0583
AXC 0002	CON42		349	0583	49 0786	0387
BMI CON41	CON32		350	0786	59 0001	0259
SXC 0001			351	0259	60 0141	1445
RAM RHO	CON44		352	1445	65 0009	0113
RAU 8002			353	0113	20 1000	0204
RAL ZZ10	CON44	CALC OMEGA 7	354	0387	69 0157	0100
STL X	E00NR		355	0204	20 0661	0064
LDO			356	0100	40 0193	0197
STL X PDS			357	0197	32 0061	0437
RAU 3			358	0437	66 0093	0247
FAO XPOS	CON46		359	0090	20 1000	0254
BMI CON45			360	0090	69 0003	0259
RSL 3			361	0247	20 1000	0254
STL X	E00NR		362	0254	69 0003	0259
LDO			363	0259		

	STL XP08		363	0207	20	0061	0114
	RAU 3		364	0114	60	0093	0297
	FAD XP08		365	0297	32	0061	0337
	SMI CON47	CON46	366	0537	46	0061	0341
CON 47	RSU 3		367	0140	61	0093	0347
	FOV TWD		368	0140	34	0070	0370
	STU X		369	0270	21	1000	0304
	RAL X		370	0304	65	1000	0305
	STL X		371	0205	60	0093	0305
	LOD		372	0354	69	0257	0100
	STL XP08	E00NR	373	0354	37	0061	0100
	RAU 3		374	0164	60	0093	0397
	FAD XP08		375	0397	32	0061	0397
	SMI WELL	CON46	376	0587	46	0190	0341
WELL	RAU POINT		377	0190	60	0643	0547
	STU M	TRY	378	0547	34	0128	0170
TRY	RSU 3		379	0781	61	0093	0397
	FOV M		380	0587	34	0128	0170
	STU X		381	0178	21	1000	0404
	RAL X		382	0404	65	1000	0655
	LOD	E00NR	383	0655	69	0258	0100
	STL XP08		384	0258	20	0061	0214
	RAU 3		385	0214	60	0093	0347
	FAD XP08		386	0647	32	0061	0377
	SMI SAD	CON46	387	0637	46	0240	0341
SAD	RAU N		388	0240	60	0258	0333
	FAD ONE		389	0633	32	0093	0725
	STU M	TRY	390	0725	61	0128	0781
CON 44	LOD ONE		391	1099	69	0098	1001
	STU NNN	CALC PD 8	392	1001	24	0500	0307
	STO NNN	OMEGA	393	0307	84	0110	0163
	RAU LAMAY		394	0163	60	0236	0391
	FMP RHO		395	0391	39	0141	0544
	STU T		396	0541	21	1350	0554
	RAU BETAT		397	0254	60	0020	0773
	F88 RHO		398	0773	33	0411	0307
	STU R		399	0517	61	0422	1025
	RAU T		400	1025	60	1350	0307
	FOV R		401	0705	34	0422	0522
CON 48	STU 8	CON48	402	0522	61	0126	0329
	RAL 8		403	0329	65	0126	1031
	STL X		404	1031	20	1000	0604
	LOD	E00NR	405	0604	69	0357	0100
	STL XP08		406	0357	20	0061	0264
	RAU XP08		407	0264	60	0061	0265
CON 49	SMI CON49	CON46	408	0265	46	0128	0341
	RAU NNN		409	0128	60	0504	0609
	FAD ONE		410	0609	32	0093	1073
	STU NNN		411	1073	61	0504	0407
	FMP NNN		412	0407	39	0126	0176
CON 46	STU 8	E00CL	413	0176	39	0126	0654
	LOD XP08		414	0654	81	0126	0329
	STO 1977		415	0341	69	0061	0314
	PCH 1977		416	0394	64	1977	1078
	RBC 0001		417	0314	71	1977	1078
	REU LAMOA	C	418	1080	71	1977	1078
	STU X		419	1577	89	0001	0683
CON 50	LOD X	CON50	420	0683	61	7390	1493
	STL X	E00NR	421	1495	21	1000	0704
	RAU LAMOA	C	422	0704	69	0098	0704
	RAU 8002	C	423	0507	60	7000	0754
	F8M X		424	0754	67	7390	1543
	SMI CON51	CON52	425	1545	60	8000	1004
CON 52	RAU ONE		426	1000	38	7000	1627
	FAD ONE		427	1627	46	1130	1084
	STU MN		428	1081	60	0098	1054
	FMP LAMOA	C	429	1054	32	0110	0687
	STU MN		430	0687	21	0110	0513
	FMP LAMOA	C	431	0213	39	7390	0290
	RBL 8003		432	0290	39	0110	0305
CON 51	STL X	CON50	433	0160	66	8003	0267
	LOD	E00CL	434	0267	20	1000	0704
CON 54	STO 04	CON54	435	1130	59	0383	1036
	LOD XP08		436	1036	37	0126	1180
	RBC X 0003		437	1036	69	0383	1036
	STO X 1981	C	438	1180	89	0003	1086
	STO X 1396		439	1086	69	7000	1104
	AXC 0001		440	1104	24	1104	0334
	LOB X	C	441	0334	24	1396	1149
	STO 1982		442	1149	58	0001	1493
	STO 1397		443	0755	69	7000	1154
	AXC 0001	C	444	1154	24	1982	0535
	LOD X 0001		445	0535	34	1397	0550
	STO 1983		446	0550	58	0001	0556
	LOD XP08	C	447	0556	69	7000	1204
	STO 1984		448	1204	24	1983	1136
	LOD XP08		449	1136	24	0398	1051
	STO 1984		450	1051	59	0198	0364
	RBC 1399		451	0364	24	1984	0737
	PCH 1977		452	0737	24	1399	0900
	R88 0004		453	0202	21	1977	1677
CON 55	LOD ZERO	CON55	454	1677	83	0004	0783
	STO SUM1		455	0783	69	1186	0783
	STO SUM2		456	0789	24	0142	1595
	R8A 0003		457	1595	24	1988	1103
CON 56	RAU 0003	CON56	458	1101	81	0003	0557
	FAD LAMDA A	B	459	0557	60	5400	1005
	STU TEMP1	A	460	1005	32	3390	0117
	FOV TEMP1		461	0317	21	0572	1225
	RAU SUM1		462	1225	60	3390	0683
	FOV TEMP1		463	0685	34	0572	0622
	STU SUM1		464	0622	41	0142	0411
	FOV TEMP1	A	465	0411	21	0142	0411
	RAU SUM1		466	0645	60	3380	0735
	FOV TEMP1		467	0735	34	0572	0679
	FMP TEMP1	A	468	0672	34	0572	0722
	STU SUM2		469	0722	39	3390	0544
	FOV TEMP1	A	470	0340	38	1098	1175
	RAU SUM2		471	1175	21	1098	1151
CON 57	AXA 0001	CON57	472	1151	20	0198	0607
	NZA CON56		473	0607	40	0557	0111
	RAU ONE		474	0111	20	0198	0607
	F88 RHO		475	1284	53	0141	0367
	STU TEMP1		476	0367	21	0572	1225
	FMP L		477	1225	39	0258	0397
	FAO SUM2		478	0288	32	1098	1275
	STU TEMP2		479	1275	21	1275	0799
	RAU L		480	1033	60	0028	1083
	FAO SUM1		481	1083	32	0142	0512
	FMP TEMP1		482	0512	39	0572	0622
	STU TEMP2		483	0622	34	1930	1880
	RAU L	B	484	1880	21	1275	0263
	AXA 0001		485	0263	52	0001	0569
CON 58	NZA CON55	CON58	486	0569	69	0198	0173
	LOD ZERO		487	0173	69	1186	0139
	STO SUM1		488	1039	84	0292	1695
	R8A 0004	CON59	489	1695	21	1695	1001
CON 59	RAU SUM8J		490	1001	60	0192	0697
	STU SUM8J	A	491	0697	29	3360	0781
	AXA 0001		492	0781	21	0398	1745
	NZA CON59	CON60	493	1745	30	0000	1255
			494	1251	40	1001	0755

CON60	RSU	D004			495	1055	81	0004	0141
CON61	RSU	B004	A	CON61	496	0161	80	2360	3135
	FDV	SUMBJ			497	0315	34	0192	0242
	STU	8J			498	0243	21	3570	0223
	AXA	D001	A		499	0229	50	0001	0379
CON62	NZA	CON61		CON62	500	0379	40	0161	1133
	LOO	05		CON63	501	1133	69	1236	0000
CON63	STO	1977			502	1236	69	1089	0292
	RSU	D004		PUNCH	503	0292	24	1977	1430
	LOO	8J	A	COEFF	504	1430	81	0004	1206
	STO	1981			505	1286	69	3370	0273
	AXA	C001			506	0273	50	0001	0384
	STO	B004	A		507	0384	30	0001	0390
	AXA	C001			508	0390	69	3370	0384
	STO	1982			509	0320	24	1982	0785
	AXA	C001			510	0785	50	0001	0591
	LOO	8J	A		511	0591	69	3370	0373
	STO	1983			512	0373	24	1983	1436
	AXA	C001			513	1436	30	0001	0342
	LOO	8J	A		514	0342	69	3370	0423
	STO	1984			515	0423	24	1984	1037
CON64	PCH	1977		CON64	516	1037	71	1977	1277
CON65	RBC	0015		CON65	517	1277	89	0015	1183
	RSU	0004			518	1183	81	0004	1139
	LOO	ZERO			519	1139	69	1186	1189
CON66	STO	BUMFX		CON66	520	1189	24	0392	1795
	RSU	OMEGA	A		521	1795	60	3400	1105
	FMP	T	C		522	1105	39	7350	0600
	STU	W		FUNCTION	523	0600	21	1404	0657
	RAM	8003		OF	524	0657	67	8003	0365
	RSU	8002		TIME	525	0365	60	8002	0523
	F5M	15			526	0523	38	0266	1454
CON75	BMI	CON75		CON76	527	1454	46	0707	0308
	RAL	W			528	0707	65	1404	0659
	LOO	8J		EOGFA	529	0659	69	0312	0050
	RSU	8002			530	0312	60	8002	0071
	FMP	8J	A		531	0071	39	3370	0320
	FAD	BUMFX			532	0320	32	0392	0619
CON76	STU	BUMFX		CON76	533	0619	21	0392	0308
	AXA	C001			534	0308	50	0001	0414
CON67	NZA	CON66		CON67	535	0414	40	1795	0168
	LOO	T	C	EOOCL	536	0168	69	0151	0000
	STO	1977			537	0121	69	7350	1504
	LOO	BUMFX			538	1504	69	0392	1845
	STO	1978			539	1480	69	0392	1845
	PCH	1978			540	1845	24	1978	1131
	AXC	C001			541	1131	71	1977	1777
CON99	NZC	CON65		CON99	542	1777	58	0001	1833
	NOP	8000			543	1833	48	1183	1087
CON99	HLT				544	1027	08	1137	8000
	01	0000	0000		545	0725	01	1181	1181
	02	0000	0000		546	0465	01	0000	0000
	03	0000	0000		547	0233	02	0000	0000
	04	0000	0000		548	0333	03	0000	0000
	05	0000	0000		549	0323	04	0000	0000
	PREC	10	0000	0054	550	1029	05	0000	0000
	ZERO	00	0000	0051	551	0008	10	0000	0054
	ONE	10	0000	0051	552	1186	00	0000	0051
	TWO	20	0000	0051	553	0098	10	0000	0051
	THREE	30	0000	0051	554	0070	20	0000	0051
	FOUR	40	0000	0051	555	0650	40	0000	0051
	FIVE	50	0000	0051	556	0068	30	0000	0051
	SIX	60	0000	0051	557	0700	60	0000	0051
	SEVEN	70	0000	0051	558	0750	60	0000	0051
	EIGHT	80	0000	0051	559	1050	70	0000	0051
	NINE	90	0000	0051	560	0226	13	5000	0253
	POINT	11	0000	0051	561	0643	11	0000	0051
	AA10	10	0000	0051	562	0279	10	0000	0051
	AA11	24	9998	8850	563	0280	34	9998	6850
	AA12	31	2575	8349	564	1172	31	2575	8349
	AA13	25	9137	1248	565	0122	25	9137	1248
	AA14	17	1562	0047	566	0078	17	1562	0047
	AA15	84	3080	0045	567	0029	84	3080	0045
	AA16	69	0600	0044	568	0026	69	0600	0044
ZZ210	00	0000	0000		569	0009	00	0000	0000

Part - E

Reactivity as a Function of Time

The equation that required a numerical integration was

$$n(t) = n(0) \left[\frac{a_2 + (\beta - 1)}{a_2 + (\beta - 1) e^{-bt}} \right] e^{-\int_0^t \frac{\lambda dt'}{a_2 + (\beta - 1) e^{-bt'}}} + \int_0^t \frac{\lambda dt'}{a_2 e^{-bt'} + (\beta - 1)} \quad (108)$$

The two integrals of Eq. (108) were solved numerically by use of Simpson's Rule,

$$\int_0^t \frac{\lambda dt'}{a_2 + (\beta - 1) e^{-bt'}} = \int_0^t y(t') dt' = \frac{h}{3} \left[(y_0 + y_{2m}) + 4(y_1 + y_3 + \dots + y_{2m-1}) + 2(y_2 + y_4 + \dots + y_{2m-2}) \right] \quad (109)$$

where the interval from 0 \rightarrow t was divided into 2m increments.

Thus

$$h = \frac{t - 0}{20} = \Delta t \quad (110)$$

The two integrals of Eq. (108) are almost alike, thus the same part of the program was used to evaluate the two integrals. The two integrals were identified by index register A. Since the integrals were evaluated at various times after the reactivity change, time was indexed by register C. The integrals were evaluated between time (C) and time (C - 1) and then added to the value of the integral from time (0) to time (C).

A continuous looping system was devised such that the computer kept track of the point being calculated to determine whether it was an odd or an even time point. If the point was odd, the value of the function at that point was multiplied by four and if the point was even, the function evaluated at that point was multiplied by two, as indicated by Eq. (109). After the function was evaluated at a particular time, the next time point was obtained by adding to the last time point. The two end points y_0 and y_{2m} were determined by testing the time to see if it corresponded to the end point times.

Since the exponential function grew very large or very small as time increased, a test was used to see whether the exponential was larger or smaller than 135.0. An overflow by the computer would have destroyed the results of the computation.

The following symbolic terms were used in the program:

$$\text{SUM} - \int_{T_c}^{T_{c+1}} y(t)_A dt$$

$$\text{PART} - y_0 + 4y_1 + 2y_2 + 4y_3 + 2y_4 + \dots + y_{2m}$$

$$\text{SUM}_A - \sum_{c=0}^{c=n} \int_{T_c}^{T_{c+1}} y(t)_A dt$$

NOFT - Neutron density as a function of time.

T - Time.

$$A2 - 1 - \beta + \frac{\rho}{A}$$

A - Magnitude of the final value of k_{ex} .

$$\text{DELTA} - \frac{T_{c+1} - T_c}{20} = H$$

TWENTY - 20.0.

BMONE - $(\beta - 1)$

$$\text{ETOFL} - e^{-\int_0^t y(t_1) dt} + \int_0^t y(t_2) dt$$

Input Data. The data that is required is listed below:

Position	Description
0086	β , total fraction of delayed neutrons.
0106	b, constant of the reactivity term.
0091	$(\beta - 1)$
0044	Absolute reactivity. Positive value for negative insertion of reactivity, A.
0136	λ
0601	The value of time at which the flux is to be calculated, T_c .
0621	

Output Data. Two output cards are punched for each time point.

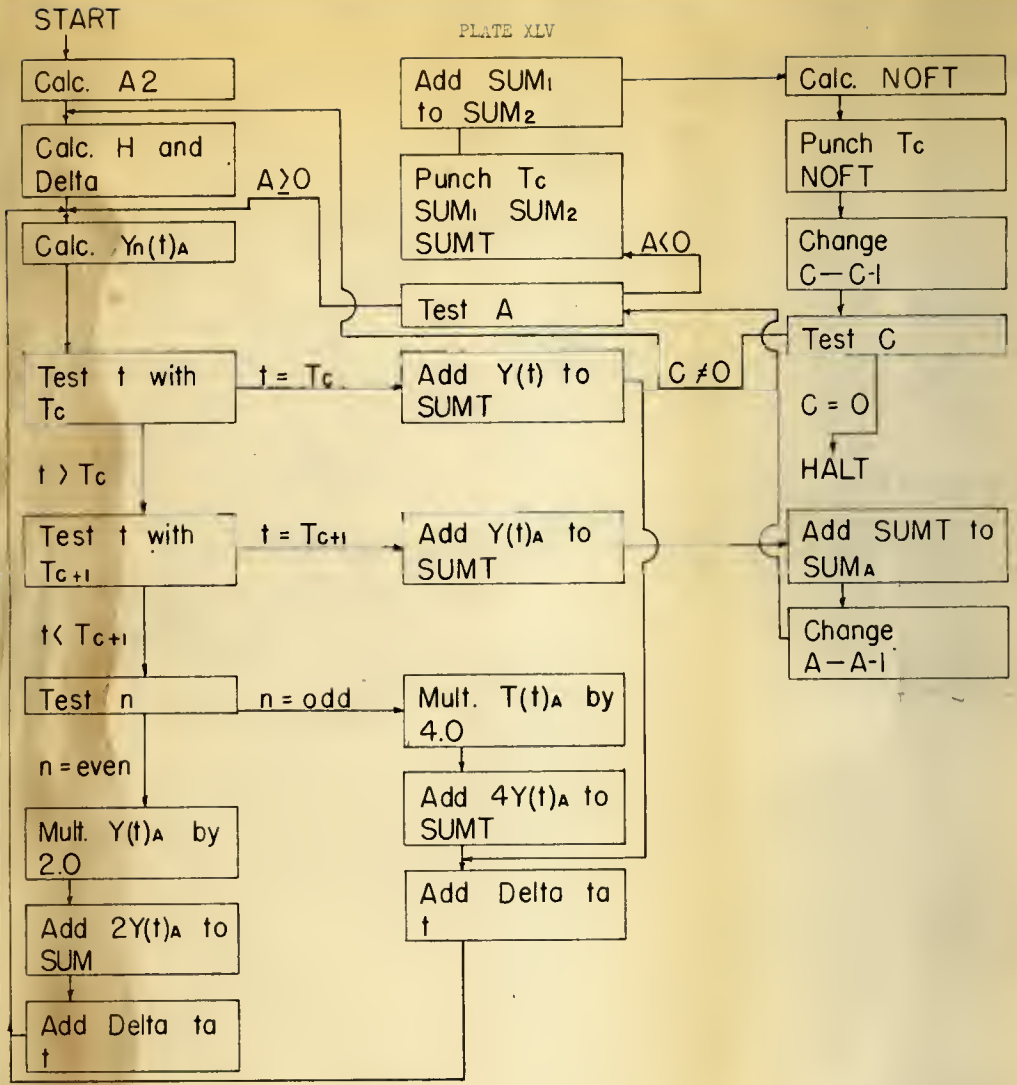
The data will appear in the following form.

Print position.	1977	1978	1979	1980	1981	1982
First line.	T(C)	SUM	SUM _{A=1}	SUM _{A=2}		
Second line.					T(C)	NOFT

EXPLANATION OF PLATE XLV

Flow diagram of the computer program for
the solution of the reactor kinetic equations
using one group of delayed neutrons and when
reactivity is a function of time.

PLATE XLV



REACTIVITY
AS A
FUNCTION
OF TIME

	BLR 0600	0700			1	0000	00	0000	0000
	BLR 0700	0800			2	0000	00	0000	0000
	BLR 0800	0900	CORRECTION		3	0000	00	0000	0000
	BLR 1951	1960	TRACE		4	0000	00	0000	0000
	BLR 1977	1984	READ		5	0000	00	0000	0000
	SYN T	0400	PRINT		6	0000	00	0000	0000
	SYN SUM	0630			7	0000	00	0000	0000
	SYN PART	0640			8	0000	00	0000	0000
	SYN START	1999			9	0000	00	0000	0000
E00CL	STO ZZZ1				10	0000	90	0000	0000
CONT1	LOD ZZZ1	CONT1	READ		11	0006	69	0009	0018
	STO 1977		OUT		12	0030	24	1978	0031
	STO 1978		DATA		13	0030	24	1978	0031
	STO 1979		ZERO		14	0031	24	1979	0032
	STO 1980		SUBROUTINE		15	0032	24	1980	0033
	STO 1981				16	0033	24	1981	0034
	STO 1982				17	0034	24	1982	0035
	STO 1983				18	0035	24	1983	0036
	STO 1984				19	0036	24	1984	0037
E00EA	STO AAA1	ZZZ1	E TO THE X		20	0050	24	0053	0056
	STL AAA2		SUBROUTINE		21	0056	20	0011	0014
	STW AAA2				22	0014	27	0011	0015
	STL AAA3				23	0015	20	0019	0022
	FMP AAA3				24	0022	60	0019	0023
	FMP AAA3.6				25	0023	39	0026	0027
	FAD AAA3.5				26	0076	32	0029	0035
	FMP AAA3				27	0005	39	0019	0029
	FAD AAA3.4				28	0069	32	0072	0049
	FMP AAA3				29	0149	39	0019	0159
	FAD AAA3.3				30	0149	38	0122	0099
	FMP AAA3				31	0099	39	0019	0169
	FAD AAA3.2				32	0169	32	0172	0149
	FMP AAA3				33	0169	39	0019	0219
	FAD AAA3.1				34	0219	32	0222	0199
	FMP AAA3				35	0199	39	0019	0269
	FAD AAA3.0				36	0269	32	0272	0249
	STU AAA4				37	0249	21	0004	0000
	FMP AAA4				38	0007	39	0004	0054
	STU AAA4				39	0054	21	0004	0057
	FMP AAA4				40	0057	39	0004	0104
	STU AAA4				41	0104	21	0004	0107
	FMP AAA4				42	0107	21	0004	0066
AAA5	STU AAA5	AAA6			43	0065	46	0018	0319
	FOV AAA5				44	0018	60	0272	0029
	STU AAA4	AAA6			45	0154	34	0004	0139
AAA6	STU AAA4	AAA1			46	0319	21	0004	0319
START	RAC 0020				47	1999	88	0020	0055
	RAA 0001				48	1999	88	0020	0055
	RAU ZERO				49	0035	80	0061	0063
	STU SUM				50	0061	60	0064	0369
	BXA 0001	A			51	0369	21	2630	0083
	STU SUM				52	0033	31	0001	0039
	RAU BETA				53	0039	21	2630	0133
	FOV A				54	0133	60	0086	0041
	FMP BETA				55	0041	34	0004	0041
	FAD ONE				56	0094	33	0086	0013
	STU X2	CON40			57	0013	33	0004	0041
CON40	RAU T	C	CON40		58	0043	21	0048	0001
	AXC				59	0001	60	6600	0105
	RAU T 0001	C			60	0105	58	0001	0111
	FOV TWENTY	C			61	0111	33	6600	0077
	STU DELTA				62	0077	34	0080	0120
	RAA 0001	CON13			63	0130	21	0084	0037
CON13	RSE 0001				64	0137	21	0042	0037
	RAU ZERO				65	0045	80	0001	0051
	STU SUM				66	0051	83	0001	0157
CON19	RAU T	C	CON19		67	0157	60	0064	0419
	FMP T				68	0419	21	0024	0127
	FMP T				69	0127	60	6600	0155
	FMP T				70	0155	21	0600	0103
	STU B				71	0103	39	0106	0156
	RAU B003				72	0106	39	0600	0150
	RAU B002				73	0100	39	0600	0150
	FBN ONE35				74	0150	39	0600	0200
	4MI CON10				75	0200	21	0204	0207
	NZA C CON14	CON31			76	0207	67	8003	0115
CON30	STU DELTA	CON15			77	0073	38	0126	0133
CON14	STU DELTA	CON16			78	0073	60	0202	0133
CON16	STL ARG	EO0EA			79	0206	40	0039	0010
	RAU B002				80	0109	20	0063	0109
	STU ETOS				81	0066	69	0469	0050
	NZA CON17	CON1R	BETAMINUB1		82	0469	60	8059	0177
CON17	FMP SWONE				83	0177	21	0082	0085
	STU BONE				84	0085	40	0032	0089
	FAD ONE				85	0085	39	0048	0085
	RAU L				86	0141	38	0048	0085
	STU PART	A	CON2D		87	0038	60	0136	0151
	RAU L	CON1			88	0183	21	0180	0230
CON31	STU PART				89	0025	21	0180	0230
CON30	FOV L				90	0183	21	2640	0093
	STU AG	A			91	0257	40	0050	0161
	STU ONE				92	0257	40	0048	0098
	RAU ZERO	CON2D			93	0257	40	0048	0098
					94	0250	21	0048	0193
					95	0241	34	0048	0098
					96	0241	34	0048	0098
					97	0143	60	0048	0203
					98	0203	21	0048	0203
CON51					99	0161	60	0064	0519

	STU PART A	CON20	100	0519	21	2640	0093
CON18	FHW		101	0089	39	0048	0148
	FAO SMDNE		102	0146	32	0091	0017
	STU ZERO		103	0017	01	032	0076
	RAU L		104	0075	60	0136	0291
	FOV SZERO		105	0291	34	0336	0376
	STU PART A	CON20	106	0376	60	2640	0093
CON15	RAL S	CON16	107	0010	65	0004	0109
CON20	RAU T	C	108	0198	60	6600	0203
	F88 T		109	0205	33	0600	0227
	NZU CONT3	CONHT4	110	0287	44	0081	0132
CON13	RAU T 0001	C	111	0081	60	6600	0227
	F88 T		112	0087	60	6600	0555
	NZU CONT5	CONHT6	113	0277	33	0600	0277
	AXC 0001		114	0114	44	0131	0182
CDNT5	AXC 0001		115	0131	58	0001	0137
	RAU T		116	0137	32	0011	0137
	NZB CONT7	CONHT8	117	0193	42	0046	0047
CDNT7	RAU PART A		118	0148	60	0046	0093
	FMP TWO		119	0095	59	0198	0248
	FAD SUMT		120	0248	39	0024	0101
	STU SUMT		121	0101	31	012	0182
	RAU T		122	0327	60	0600	0305
	RAU PART A	CDNT9	123	0305	32	0042	0158
CDNT8	FMP FOUR		124	0047	60	2640	0145
	FAO SUMT		125	0145	39	0398	0348
	STU SUMT		126	0348	32	0024	0151
	S8R 0002	CONH10	127	0151	31	0024	0377
	RAU T		128	0377	53	0002	0233
	RAU DELTA	CONHT9	129	0333	60	0600	0355
CON10	FAO SUMT		130	0355	32	0042	0151
CON14	RAU PART A		131	0138	60	2640	0079
	FAO SUMT		132	0195	32	0024	0201
	STU SUMT		133	0201	33	0004	0233
CON16	RAU PART A	CONH10	134	0182	60	2640	0245
	FAD SUMT		135	0245	39	0024	0251
	STU SUMT		136	0251	31	0024	0427
	FMP W		137	0427	39	0084	0134
	FOV THREE		138	0134	34	0134	0277
	STU SUMT		139	0237	21	0024	0477
	FAO SUMT	A	140	0477	32	2630	0307
	STU SUMT	A	141	0307	60	8002	0079
	SXA 0001		142	0283	31	0001	0139
	RAU PART A	CON63	143	0139	41	0092	0243
CON63	AXC 0001	CON13	144	0243	58	0001	0051
CON18	LDB	EODCL	145	0092	69	0295	0000
	LOO T	C	146	0295	69	6000	0491
	STD 1977		147	0253	24	1977	0280
	LOO SUMT		148	0280	60	6000	0577
	OTD 1978		149	0527	24	1978	0181
	RAA 0001		150	0181	80	0001	0589
	LOO SUMT	A	151	0287	69	2630	0333
	STO 1979		152	0333	24	1979	0238
	SXA 0001		153	0333	60	0013	0333
	LOO SUMT	A	154	0088	69	2630	0383
	STD 1980		155	0383	24	1980	0433
	PCW 1977		156	0433	71	1977	0577
	RAU SUMT	A	157	0577	60	2630	0333
	AXA 0001	A	158	0135	80	0001	0341
	F88 SUMT	A	159	0341	33	2630	0357
	STU EDNN		160	0357	60	0162	0165
	RAL EPDNN		161	0165	65	0020	0067
	LDB	EOD6A	162	0067	69	0620	0080
	RAU 8002		163	0200	60	0082	0079
	STU ETOTFL		164	0079	21	0184	0337
	RAU SMDNE		165	0337	60	0091	0345
	FAO A2		166	0345	32	0048	0125
	FOV 80NE		167	0125	34	0180	0330
	FMP ETOTFL		168	0330	39	0184	0334
	STU NOFT		169	0334	21	0138	0391
	LDB	EODCL	170	0391	69	0144	0000
	LOO T	C	171	0144	69	6600	0303
	STO 1981		172	0303	24	1981	0284
	LBO NOFT		173	0284	69	0138	0441
	STO 1982		174	0441	24	1982	0185
	PCW 1977		175	0185	71	1977	0927
	SXA 0001		176	0927	59	0001	0483
	NZC CON40	CON41	177	0483	48	0001	0387
CON41	00 0000	8000	178	0387	00	0000	0000
ZERO	00 0000	0000	179	0064	00	0000	0000
ONE	10 0000	0051	180	0064	10	0000	0051
TWO	20 0000	0051	181	0198	20	0000	0051
THREE	30 0000	0051	182	0187	30	0000	0051
FOUR	40 0000	0051	183	0398	40	0000	0051
FIVE	50 0000	0052	184	0080	20	0000	0052
SIX	60 0000	0053	185	0185	13	5000	0053
SEVEN	70 0000	0000	186	0009	00	0000	0000
EIGHT	80 0000	0000	187	0009	00	0000	0000
NINE	90 0000	0000	188	0072	10	0000	0051
AA110	00 0000	6850	189	0228	24	9998	6850
AA111	01 0000	8349	190	0172	31	3375	8349
AA112	02 0000	1248	191	0428	26	9137	1248
AA113	03 0000	0047	192	0072	17	1562	0047
AA114	04 0000	0045	193	0072	19	1562	0045
AA115	05 0000	0044	194	0072	17	1562	0044
AA116	06 0000	0051	195	0072	17	1562	0051
AA117	07 0000	0050	196	0619	10	0000	0050
AA118	08 0000	0050	197	0619	10	0000	0050
AA119	09 0000	0050	198	0617	40	0000	0050
AA120	10 0000	0050	199	0616	60	0000	0050
AA121	11 0000	0051	200	0615	60	0000	0051
AA122	12 0000	0051	201	0614	10	0000	0051
AA123	13 0000	0051	202	0614	10	0000	0051
AA124	14 0000	0051	203	0612	14	0000	0051
AA125	15 0000	0051	204	0611	16	0000	0051
AA126	16 0000	0051	205	0610	18	0000	0051
AA127	17 0000	0051	206	0609	20	0000	0051
AA128	18 0000	0051	207	0608	20	0000	0051
AA129	19 0000	0051	208	0607	30	0000	0051
AA130	20 0000	0051	209	0607	40	0000	0051
AA131	21 0000	0051	210	0605	40	0000	0051
AA132	22 0000	0051	211	0604	60	0000	0051
AA133	23 0000	0051	212	0601	70	0000	0051
AA134	24 0000	0051	213	0602	80	0000	0051
AA135	25 0000	0051	214	0601	10	0000	0052

Part - F

Tabulated Data, Analytical

Table 5. Flux ratio as a function of the value of the decay constant.

Group	: Value Used : For λ : (Sec.-1)	: React- : ** : : (\$)	: tivity : : :	Flux Ratios at Time (Sec.)			
				2	10	40	75
1	0.01216	-0.10	0.8698	0.7852	0.6194	0.4931	
2	0.02958	-0.10	0.8699	0.7856	0.6216	0.4968	
3	0.10747	-0.10	0.8670	0.7861	0.6209	0.4943	
4	0.2900	-0.10	0.8705	0.7867	0.6205	0.4940	
5	1.0002	-0.10	0.8705	0.7856	0.6195	0.4936	
*		-0.10	0.8698	0.7852	0.6192	0.4927	
1	0.01216	-0.50	0.5683	0.4019	0.2018	0.1105	
2	0.02958	-0.50	0.5684	0.4026	0.2043	0.1131	
3	0.10747	-0.50	0.5687	0.4034	0.2029	0.1108	
4	0.2900	-0.50	0.5699	0.4042	0.2023	0.1106	
5	1.002	-0.50	0.5698	0.4023	0.2018	0.1103	
*		-0.50	0.5683	0.4019	0.2016	0.1102	
1	0.01216	-1.50	0.3023	0.1712	0.0653	0.0298	
2	0.02958	-1.50	0.3024	0.1718	0.0665	0.0307	
3	0.10747	-1.50	0.3027	0.1723	0.0656	0.0297	
4	0.2900	-1.50	0.3039	0.1726	0.0654	0.0297	
5	1.0002	-1.50	0.3036	0.1713	0.0652	0.0296	
*		-1.50	0.3023	0.1712	0.0652	0.0296	
1	0.01216	0.01	1.0152	1.0273	1.0578	1.0890	
2	0.02958	0.01	1.0152	1.0272	1.072	1.0572	
3	0.10747	0.01	1.0152	1.0271	1.0573	1.0886	
4	0.2900	0.01	1.0151	1.0270	1.0574	1.0887	
5	1.0002	0.01	1.0151	1.0272	1.0576	1.0890	
*		0.01	1.0152	1.0273	1.0577	1.0891	
1	0.01216	0.10	1.1734	1.3407	1.9008	2.7335	
2	0.02958	0.10	1.1736	1.3401	1.8924	2.7074	
3	0.10747	0.10	1.1734	1.3391	1.8929	2.7169	
4	0.2900	0.10	1.1725	1.3374	1.8926	2.7174	
5	1.0002	0.10	1.1725	1.3398	1.8987	2.7305	
*		0.10	1.1737	1.3410	1.9014	2.7360	

Keepin's delayed neutron parameters.

Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

* Here, all of the accepted values of the decay constants were used.

** Here, the value for the particular decayed constant was the accepted value for the i-th group plus the experimental uncertainty.

Table 6. Flux ratio as a function of the total fraction of delayed neutrons.

β	: React- : tivity : : (\$)	Flux Ratios at Times (Sec.)				
		2	10	40	100	200
0.006	0.01	1.0160	1.0292	1.0618	1.1194	1.2112
0.007	0.01	1.0194	1.0338	1.0717	1.1391	1.2571
0.008	0.01	1.0161	1.0295	1.0624	1.1204	1.2210
0.006	0.10	1.1863	1.3736	2.0246	4.0070	12.306
0.007	0.10	1.1894	1.3844	2.0478	4.1021	12.852
0.008	0.10	1.1865	1.3780	2.0260	4.0128	12.338
0.006	0.30	1.8493	3.5377	25.615	1194.5	6.873 x 10 ⁻⁵
0.007	0.30	1.8558	3.5643	26.181	1256.8	7.577 x 10 ⁻⁵
0.008	0.30	1.8514	3.5442	25.738	1207.3	7.016 x 10 ⁻⁵
0.006	-0.10	0.8622	0.7701	0.5977	0.3969	0.2096
0.007	-0.10	0.8642	0.7722	0.6022	0.4029	0.2147
0.008	-0.10	0.8623	0.7703	0.5980	0.3981	0.2099
0.006	-0.50	0.5513	0.3795	0.1870	0.0672	0.0155
0.007	-0.50	0.5526	0.3808	0.1880	0.0676	0.0156
0.008	-0.50	0.5517	0.3799	0.1873	0.0673	0.0155

Hughes' delayed neutron parameters.

Prompt neutron lifetime, $\lambda = 8.0 \times 10^{-5}$ seconds.

Part - G

Tabulated Data, Experimental

Table 7. Experimental rod drop No. A-2

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.580	0.445	0.481
4	0.528	0.435	0.465
6	0.479	0.445	0.471
8	0.455	0.420	0.456
10	0.417	0.435	0.472
15	0.364	0.425	0.463
20	0.319	0.435	0.468
30	0.259	0.435	0.468
40	0.214	0.415	0.464
50	0.181	0.430	0.466
75	0.121	0.430	0.461
100	0.083	0.425	0.458

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 90

Table 8. Experimental rod drop No. A-3.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.521	0.560	0.603
4	0.458	0.575	0.589
6	0.424	0.550	0.585
8	0.394	0.535	0.574
10	0.363	0.535	0.580
15	0.315	0.525	0.564
20	0.274	0.530	0.568
30	0.219	0.525	0.563
40	0.179	0.525	0.560
50	0.149	0.520	0.562
75	0.099	0.515	0.544
100	0.069	0.495	0.527

Reactor: Argonaut
Rod dropped: Shim and Fine

Critical rod positions: Shim - 29.5
Fine - 100

Table 9. Experimental rod drop No. A-4.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.956	0.029	0.037
4	0.939	0.033	0.035
6	0.927	0.033	0.036
8	0.921	0.033	0.038
10	0.909	0.034	0.038
15	0.888	0.035	0.038
20	0.872	0.035	0.038
30	0.841	0.036	0.039
40	0.808	0.038	0.039
50	0.796	0.035	0.038
75	0.739	0.036	0.039
100	0.691	0.035	0.038

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 20

Table 10. Experimental rod drop No. A-5.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.594	0.420	0.454
4	0.541	0.410	0.444
6	0.500	0.410	0.432
8	0.470	0.400	0.433
10	0.438	0.400	0.432
15	0.486	0.390	0.427
20	0.341	0.395	0.428
30	0.282	0.395	0.426
40	0.234	0.390	0.424
50	0.198	0.390	0.424
75	0.134	0.395	0.422
100	0.093	0.390	0.422

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 80

Table 11. Experimental rod drop No. A-6.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.628	0.370	0.393
4	0.583	0.345	0.376
6	0.547	0.340	0.364
8	0.516	0.335	0.364
10	0.477	0.345	0.373
15	0.421	0.340	0.376
20	0.373	0.345	0.378
30	0.311	0.345	0.378
40	0.262	0.345	0.376
50	0.223	0.345	0.376
75	0.155	0.345	0.372
100	0.109	0.340	0.372

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 70

Table 12. Experimental rod drop No. A-7.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.683	0.285	0.306
4	0.618	0.299	0.325
6	0.583	0.292	0.318
8	0.556	0.283	0.312
10	0.529	0.281	0.319
15	0.471	0.283	0.312
20	0.422	0.290	0.312
30	0.354	0.294	0.316
40	0.306	0.292	0.313
50	0.267	0.287	0.307
75	0.188	0.288	0.313
100	0.140	0.282	0.304

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 60

Table 13. Experimental rod drop No. A-8.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.738	0.221	0.233
4	0.696	0.212	0.235
6	0.661	0.210	0.230
8	0.643	0.199	0.219
10	0.608	0.206	0.228
15	0.578	0.195	0.218
20	0.513	0.217	0.226
30	0.444	0.209	0.226
40	0.389	0.211	0.228
50	0.348	0.207	0.225
75	0.263	0.208	0.224
100	0.201	0.217	0.224

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 50

Table 14. Experimental rod drop No. A-9.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.823	0.133	0.141
4	0.790	0.130	0.141
6	0.768	0.125	0.136
8	0.737	0.131	0.141
10	0.714	0.131	0.145
15	0.666	0.132	0.145
20	0.632	0.131	0.144
30	0.562	0.135	0.145
40	0.514	0.133	0.146
50	0.469	0.135	0.146
75	0.382	0.135	0.145
100	0.306	0.138	0.147

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 40

Table 15. Experimental rod drop No. A-10.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.892	0.075	0.081
4	0.872	0.073	0.080
6	0.853	0.071	0.080
8	0.842	0.068	0.077
10	0.822	0.072	0.080
15	0.790	0.072	0.080
20	0.759	0.074	0.081
30	0.711	0.074	0.081
40	0.670	0.075	0.081
50	0.634	0.074	0.081
75	0.549	0.076	0.083
100	0.480	0.076	0.083

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 30

Table 16. Experimental rod drop No. A-11.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2	0.968	0.021	0.022
4	0.952	0.024	0.027
6	0.946	0.023	0.027
8	0.940	0.023	0.026
10	0.931	0.024	0.028
15	0.914	0.026	0.029
20	0.899	0.026	0.029
30	0.871	0.028	0.031
40	0.849	0.029	0.031
50	0.830	0.028	0.031
75	0.778	0.030	0.032
100	0.737	0.029	0.032

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 20

Table 17. Experimental rod drop No. A-12.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2			
4			
6			
8			
10	0.988	0.004	0.004
15			
20	0.983	0.006	0.005
30			
40			
50	0.974	0.004	0.004
75			
100	0.948	0.005	0.005

Reactor: Argonaut
Rod dropped: Shim

Critical rod positions: Shim - 10

Table 18. Experimental rod drop No. A-13.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.625	0.376	0.394
4	0.564	0.375	0.405
6	0.516	0.385	0.410
8	0.483	0.375	0.414
10	0.440	0.390	0.424
15	0.386	0.380	0.426
20	0.340	0.400	0.430
30	0.277	0.400	0.430
40	0.228	0.405	0.434
50	0.195	0.400	0.432
75	0.130	0.405	0.433
100	0.090	0.400	0.433

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100

Table 19. Experimental rod drop No. A-14.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2	0.604	0.400	0.435
4	0.553	0.390	0.422
6	0.515	0.385	0.410
8	0.475	0.390	0.426
10	0.439	0.400	0.430
15	0.384	0.395	0.428
20	0.341	0.395	0.429
30	0.275	0.405	0.434
40	0.229	0.405	0.432
50	0.194	0.400	0.433
75	0.131	0.405	0.431
100	0.091	0.405	0.426

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 90

Table 20. Experimental rod drop No. A-15.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.683	0.150	0.166
15			
20	0.598	0.150	0.164
30	0.536	0.149	0.163
40	0.482	0.150	0.165
50	0.438	0.150	0.163
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: 32.1

Table 21. Experimental rod drop No. A-16.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2	.	.	.
4	.	.	.
6	.	.	.
8	.	.	.
10	0.593	0.218	0.243
15	.	.	.
20	0.494	0.221	0.243
30	0.425	0.224	0.242
40	0.374	0.224	0.241
50	0.332	0.221	0.239
75	.	.	.
100	.	.	.

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 40

Table 22. Experimental rod drop No. A-17.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2	.	.	.
4	.	.	.
6	.	.	.
8	.	.	.
10	0.500	0.315	0.341
15	.	.	.
20	0.398	0.316	0.344
30	0.334	0.316	0.340
40	0.280	0.319	0.346
50	0.244	0.315	0.341
75	.	.	.
100	.	.	.

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 50

Table 23. Experimental rod drop No. A-18.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.436	0.415	0.434
15			
20	0.337	0.405	0.433
30	0.274	0.405	0.436
40	0.229	0.400	0.433
50	0.194	0.400	0.433
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 60

Table 24. Experimental rod drop No. A-19.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.387	0.490	0.525
15	0.297	0.480	0.513
20	0.237	0.490	0.513
30	0.196	0.480	0.510
40	0.165	0.475	0.508
50			
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 70

Table 25. Experimental rod drop No. A-20.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.372	0.520	0.560
15			
20	0.280	0.515	0.553
30	0.225	0.510	0.551
40	0.185	0.505	0.540
50	0.154	0.500	0.453
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 80

Table 26. Experimental rod drop No. A-21.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.365	0.540	0.575
15			
20	0.271	0.535	0.577
30	0.217	0.535	0.570
40	0.178	0.530	0.563
50	0.148	0.530	0.562
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 90

Table 27. Experimental rod drop No. A-22.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.362	0.540	0.583
15			
20	0.267	0.550	0.585
30	0.213	0.540	0.583
40	0.175	0.540	0.584
50	0.144	0.540	0.580
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100

Table 28. Experimental rod drop No. A-23.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.980	0.007	0.007
15			
20	0.972	0.007	0.008
30	0.965	0.007	0.008
40	0.959	0.007	0.008
50	0.945	0.007	0.008
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: -91.1

Table 29. Experimental rod drop No. A-24.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.985	0.014	0.016
15			
20	0.938	0.016	0.018
30	0.925	0.016	0.017
40	0.915	0.015	0.017
50	0.903	0.015	0.017
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 84.2

Table 30. Experimental rod drop No. A-25.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.908	0.034	0.038
15	0.875		
20	0.875	0.033	0.038
30	0.843	0.035	0.039
40	0.820	0.035	0.038
50	0.800	0.034	0.037
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 76.5

Table 31. Experimental rod drop No. A-26.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.870	0.050	0.055
15	0.828	0.049	0.054
20	0.828	0.049	0.054
30	0.792	0.049	0.054
40	0.760	0.050	0.054
50	0.732	0.048	0.053
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod positions: 73.0

Table 32. Experimental rod drop No. A-27.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.818	0.075	0.083
15			
20	0.754	0.075	0.083
30	0.702	0.076	0.085
40	0.662	0.076	0.084
50	0.622	0.076	0.085
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 67.1

Table 33. Experimental rod drop No. A-28.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.818	0.215	0.238
15			
20	0.502	0.215	0.235
30	0.436	0.215	0.234
40	0.385	0.214	0.233
50	0.341	0.214	0.231
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 50.0

Table 34. Experimental rod drop No. A-29.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.558	0.252	0.277
15			
20	0.456	0.255	0.277
30	0.390	0.255	0.277
40	0.338	0.255	0.276
50	0.296	0.254	0.274
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 45.3

Table 35. Experimental rod drop No. A-30.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.496	0.318	0.347
15			
20	0.400	0.314	0.342
30	0.335	0.316	0.340
40	0.286	0.311	0.340
50	0.245	0.314	0.340
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 39.3

Table 36. Experimental rod drop No. A-31.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.476	0.350	0.385
15			
20	0.375	0.345	0.375
30	0.310	0.345	0.375
40	0.263	0.340	0.376
50	0.225	0.345	0.374
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 35.2

Table 37. Experimental rod drop No. A-32.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (β), Hughes	Reactivity worth of rod drop (β) Keepin
2			
4			
6			
8			
10	0.440	0.400	0.429
15	0.340	0.400	
20	0.340	0.400	0.430
30	0.278	0.400	0.430
40	0.234	0.399	0.425
50	0.196	0.398	0.429
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 29.2

Table 38. Experimental rod drop No. A-33.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (β), Hughes	Reactivity worth of rod drop (β) Keepin
2			
4			
6			
8			
10	0.370	0.525	0.564
15			
20	0.274	0.530	0.567
30	0.219	0.530	0.563
40	0.179	0.520	0.558
50	0.151	0.520	0.556
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 8.6

Table 39. Experimental rod drop No. A-34.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.372	0.520	0.560
15			
20	0.279	0.515	0.554
30	0.224	0.515	0.550
40	0.185	0.505	0.540
50	0.154	0.505	0.543
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 10.8

Table 40. Experimental rod drop No. A-35.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$) Keepin
2			
4			
6			
8			
10	0.383	0.495	0.537
15			
20	0.289	0.495	0.532
30	0.232	0.500	0.530
40	0.191	0.490	0.524
50	0.161	0.485	0.521
75			
100			

Reactor: Argonaut
Rod dropped: Fine

Critical rod positions: Fine - 100
Final rod position: 17.0

Table 43. Experimental rod drop No. 9.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (β), Hughes	Reactivity worth of rod drop (β) Keepin
2	0.258	1.91	2.02
4	0.195	1.89	2.01
6			
8			
10			
15	0.107	1.80	1.96
20	0.089	1.79	1.95
30	0.066	1.77	1.92
40			
50	0.037	1.87	1.99
75	0.022	1.82	1.95
100	0.014	1.76	1.89

Reactor: TRIGA
Rod dropped: Shim

Critical rod positions: Shim - up
Reg. - 299
Safety - up

Table 44. Experimental rod drop No. 11.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (β), Hughes	Reactivity worth of rod drop (β) Keepin
2	0.281	1.56	1.65
4	0.208	1.75	1.88
6	0.169	1.84	1.99
8	0.145	1.85	2.00
10	0.135	1.80	1.95
15	0.108	1.78	1.94
20	0.090	1.76	1.91
30	0.062	1.83	1.99
40	0.047	1.83	1.99
50	0.040	1.77	1.90
75	0.023	1.77	1.90
100			

Reactor: TRIGA
Rod dropped: Shim

Critical rod positions: Shim - up
Reg. - 311
Safety - up

Table 49. Experimental rod drop No. 16.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.356	1.10	1.18
4	0.283	1.17	1.27
6	0.242	1.20	1.31
8	0.207	1.25	1.36
10			
15	0.150	1.25	1.35
20	0.126	1.25	1.34
30	0.097	1.21	1.36
40	0.074	1.24	1.34
50			
75	0.040	1.26	1.36
100			

Reactor: TRIGA

Rod dropped: Shim

Critical rod positions: Shim - 466
Reg. - 474
Safety - up

Table 50. Experimental rod drop No. 19.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2			
4			
6			
8			
10			
15	0.184	1.06	1.09
20	0.157	0.99	1.06
30	0.123	0.96	1.03
40	0.099	0.98	1.04
50	0.088	0.92	1.00
75	0.051	0.94	1.02
100	0.036	0.96	0.98

Reactor: TRIGA

Rod dropped: Regulating

Critical rod positions: Shim - up
Reg. - 519
Safety - up

Table 53. Experimental rod drop No. 26.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.339	1.23	1.26
4	.	.	.
6	.	.	.
8	.	.	.
10	0.168	1.40	1.51
15	0.133	1.42	1.54
20	.	.	.
30	0.082	1.43	1.52
40	0.063	1.43	1.54
50	0.050	1.43	1.54
75	0.029	1.42	1.53
100	.	.	.

Reactor: TRIGA
Rod dropped: Shim

Critical rod positions: Shim - 500
Reg. - down
Safety - 1/2

Table 54. Experimental rod drop No. 29-1.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.283	1.56	1.65
4	.	.	.
6	.	.	.
8	.	.	.
10	.	.	.
15	0.104	1.85	1.98
20	0.085	1.85	2.00
30	0.063	1.85	2.00
40	0.048	1.87	2.01
50	0.038	1.86	2.00
75	0.022	1.85	1.99
100	0.013	1.84	1.97

Reactor: TRIGA
Rod dropped: Regulating

Critical rod positions: Shim - 338
Reg. - 805
Safety - down

Table 59. Experimental rod drop No. 27.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.282	1.56	1.65
4	0.231	1.57	1.66
6	0.195	1.58	1.70
8	0.170	1.59	1.71
10	0.145	1.65	1.80
15	0.107	1.80	1.94
20	0.089	1.78	1.93
30	0.066	1.79	1.94
40	0.050	1.80	1.94
50	0.040	1.79	1.92
75	0.021	1.83	1.99
100	0.012	1.84	1.99

Reactor: TRIGA
Rod dropped: Shim

Critical rod positions: Shim - 600
Reg. - down

Table 60. Experimental rod drop No. 28.

Time after rod drop (Sec.)	Flux Ratio	Reactivity worth of rod drop (\$), Hughes	Reactivity worth of rod drop (\$), Keepin
2	0.292	1.49	1.57
4			
6			
8	0.150	1.81	1.98
10	0.136	1.82	1.95
15	0.107	1.83	1.96
20	0.085	1.83	1.97
30	0.064	1.86	1.98
40	0.049	1.88	1.96
50	0.037	1.87	1.99
75	0.022	1.86	1.98
100			

Reactor: TRIGA
Rod dropped: Shim

Critical rod positions: Shim - 693
Reg. - down
Safety - 1/4

THEORETICAL REACTOR KINETIC MODELS AND
EXPERIMENTAL VERIFICATION

by

MARVIN KEITH DRAKE

B. S., Kansas State University, 1959

AN ABSTRACT OF A THESIS

submitted in partial fulfillment of the

requirements for the degree

MASTER OF SCIENCE

Department of Nuclear Engineering

KANSAS STATE UNIVERSITY
OF AGRICULTURE AND APPLIED SCIENCE

1960

An analytical study was made of the reactor kinetic equations for a bare homogeneous thermal reactor assuming step changes of reactivity. The study considered two sets of delayed neutron parameters, Hughes' and Keepin's, and the results of the study are given in graphical form. Experimental measurements of the reactivity worth of the control rods in the Torrey Pines TRIGA reactor were made. The reactivity values obtained with positive period measurements were compared to those obtained by means of rod drop measurements. Agreement of the two methods indicated that the kinetics theory was valid when the reactor was operating at low power levels and reactivity changes were not extremely large ($\rho < \$2.00$).

A parameter analysis of the kinetic equations was made. The reactor parameters, prompt neutron lifetime, total fraction of delayed neutrons, and decay constants of the delayed neutrons, were evaluated over a relatively wide range of values. The importance of each parameter in the kinetic equations is given.

An approximate solution to the kinetic equations was obtained using three groups of delayed neutrons instead of the usual six groups. The delayed neutron parameters were empirically chosen to fit the computed data obtained with six delayed neutron groups.

Kinetics theory for a finite reactivity insertion rate was developed. The kinetics model included the assumption that the delayed neutrons could be approximated by one group of delayed

neutrons. The reactivity insertion rate was approximated by a function, of the form $A(1 - e^{-bt^n})$, to fit the reactivity insertion rate of the TRIGA reactor. The finite insertion rate equations predicted the neutron flux to a greater degree of accuracy for short times after the reactivity change, than the neutron flux as predicted by the usual six group kinetic equations with step input reactivities.

Three IBM-650 programs are given, solution of the kinetic equations for six groups of delayed neutrons, solution of the kinetic equations for three groups of delayed neutrons, and solution of the one delayed neutron group kinetic equations when reactivity is a function of time.