

1 **Characterization of Flow through random media via**
2 **Karhunen-Loève expansion: an Information Theory**
3 **perspective**

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7 **Abstract** We leverage on Information Theory to assess the fidelity of approx-
8 imated numerical stochastic groundwater flow simulations. We consider flow
9 in saturated heterogeneous porous media, where the Karhunen-Loève (KL)
10 expansion is used to express the hydraulic conductivity as a spatially corre-
11 lated random field. We quantify the impact of the KL expansion truncation on
12 the uncertainty associated with punctual values of hydraulic conductivity and
13 flow velocity. In particular, we compare the statistical dependence between
14 variables by considering (a) linear correlation metrics (Pearson coefficient of
15 correlation) and (b) metrics capable of accounting for nonlinear dependence
16 (coefficient of uncertainty based on mutual information). We test the selected
17 metrics by analyzing the relationship between hydraulic conductivity fields
18 generated via Monte Carlo sampling with different levels of truncation of the
19 KL expansion and the corresponding fluid velocity fields, obtained through
20 the numerical solution of Darcy’s flow. Our analysis shows that employing
21 linear correlation metrics leads to a general overestimation of the correlation
22 level and IT-based indicators are valuable tools to assess the impact of the
23 KL truncation on the output velocity values. We then analyze the impact of
24 the number of retained modes on the spatial organization of the velocity field.
25 Results indicates that (i) as the number of modes decrease the spatial correla-
26 tions of the velocity field increases; (ii) linear indicators of spatial correlation
27 are again larger than the nonlinear counterparts.

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29 Information Theory

30 1 Introduction

31 In the context of subsurface geology it is well known that the hydraulic proper-
32 ties of natural aquifers exhibit heterogeneity over diverse spatial length scales
33 [7, 25]. This ubiquitous heterogeneity and the typical poor level of characteriza-
34 tion of subsurface environments lead to a lack of knowledge about the hydraulic
35 properties of the hosting formations, which is a major source of uncertainty.
36 This motivates the interest towards applications of uncertainty quantification
37 and propagation through numerical modeling of subsurface flow and transport
38 processes.

39 In this work we consider the uncertainty stemming from incomplete knowl-
40 edge of the hydraulic conductivity of fully saturated heterogeneous porous for-
41 mations on the resulting steady state flow field. To address its parameterization
42 under uncertainty, permeability (or hydraulic conductivity) is often described
43 by random fields, where the heterogeneity structure can be characterized with
44 geostatistical methodologies [7]. These methods rely on an assumed statistical
45 distribution and spatial covariance function for the parameter under investiga-
46 tion. A classical assumption is to assign a lognormal distribution to hydraulic
47 conductivity, although different models have been proposed in the literature
48 [25]. When the log-conductivity is expressed by a multi-Gaussian random field,
49 the well-known Karhunen-Loève (KL) [21] expansion can be used to reduce
50 the stochastic dimension of the problem. This allows reducing the computa-
51 tional costs for uncertainty propagation, which can be achieved also through
52 the implementation of model reduction techniques, e.g. [22, 20, 19]. KL-based
53 approximations rely on a truncation of the exact expansion, which, in practice,
54 is equivalent to selecting a number of terms in the series, also called modes.
55 The number of modes retained in the KL expansion is typically determined
56 by ensuring that a given fraction of the variance (considered as a proxy of
57 the spatial variability) of the underlying field is retained. However, this ap-
58 proach does not ensure that a satisfactory level of variability of the output is
59 also retained (see e.g., [33]) particularly in the presence of nonlinear relations
60 between inputs and quantities of interest (QoI). Quantitative indicators are
61 therefore required to assess the foreseen effect of such approximations on QoI.
62 In this work our target is to characterize the propagation of uncertainty to
63 fluid velocity in fully saturated heterogeneous porous media. The simulation
64 of fluid velocity is fundamental to numerical studies aimed at assessing the im-
65 pact of heterogeneity in fluid flow and solute transport processes taking place
66 in heterogeneous aquifers and reservoirs. For instance, the velocity field and
67 its spatial structure are key information to simulate and predict transport and
68 dilution of solute mass in the subsurface [13, 32, 8]. To address these problems,
69 general indicators that can connect parameters describing the spatial hetero-

70 geneity of geomaterials properties to flow and transport emerging features are
71 sought in recent literature [2].

72 Developed after the work of Shannon [30], Information theory (IT) pro-
73 vides a suite of indicators that can be used to assess the similarity in informa-
74 tion content of two distributions. Arguably the most widespread quantitative
75 indicator derived from IT is Shannon entropy. In addition, IT provides quan-
76 titative indicators that can assess the mutual information between variables.
77 These indicators can also be used to describe correlation within a spatial and
78 temporal fields. For example, when nonlinear systems are of concern, the IT-
79 based mutual information can be used as an alternative to linear correlation,
80 as shown, e.g., by [26] in the context of the characterization of subsurface min-
81 eral distributions. IT-based indicators have been previously employed within
82 model-based assessment of groundwater flow and solute transport. For exam-
83 ple several studies have relied on the concept of entropy as an indicator of
84 uncertainty within risk assessment procedures [23,1] or to set up optimal ex-
85 perimental design for model discrimination [17]. IT mutual information has
86 also been shown to be an indicator of the degree of nonlinearity existing be-
87 tween output variables in flow and transport simulations, with a particular
88 focus on their spatial correlation [4]. As an alternative approach, the concept
89 of entrogram was introduced in [3]. This latter corresponds to the fraction of
90 the entropy of a variable sampled within a given spatial window and that of
91 its counterpart associated with the whole spatial domain, as a function of the
92 ratio between the spatial window and the whole domain sizes. The results in
93 [3] indicate that the entropic scale (i.e., a measure of the overall persistence of
94 a pattern of association) can be related to widely employed solute transport
95 descriptors.

96 In this work, we apply IT-based metrics within stochastic groundwater
97 flow simulations. We consider stochastic simulations based on Monte Carlo
98 sampling of the KL expansion of hydraulic conductivity. The selection of the
99 number of modes in the KL-expansion is typically guided by selecting a propor-
100 tion of the variance of the original hydraulic conductivity field to be retained.
101 Such proportion can be analytically determined and increases with the num-
102 ber of retained modes. However, for values of the variance of log-conductivity
103 above unity the degree of nonlinearity between input (hydraulic conductivity)
104 and output (flow velocity) increases [4]. A critical issue in this context is to
105 control and constrain a priori the output approximation accuracy. Our work
106 explores the possibility to employ IT-based metrics in the context of a quan-
107 titative assessment of the approximation resulting from KL series truncation.
108 While IT has been employed with a similar objective in other fields [16], to
109 the best of our knowledge this idea is explored here for the first time in the
110 context of stochastic groundwater flow simulation. In particular, for flow fields
111 associated with an increasing number retained modes in the KL expansion
112 we (*i*) compare the resulting entropy of the flow field, to investigate the level
113 of variability retained in the predicted quantity of interest; (*ii*) evaluate the
114 mutual information (MI) between flow fields obtained at different levels of res-

115 olution; *(iii)* compare the behavior of the spatial auto MI of the flow fields, to
 116 investigate how the spatial organization of the diverse flow fields behaves.

117 In the following, section 2 presents the considered problem setup, the IT
 118 indicators employed in the analysis and the considered quantities of interest.
 119 Section (3) present the results obtained by the application to groundwater
 120 flow with three different levels of heterogeneity, i.e. three selected variances of
 121 log-conductivity.

122 2 Methods

123 2.1 Problem setup

124 We consider steady flow in a two dimensional heterogeneous saturated porous
 125 medium, through Darcy's law and fluid mass conservation

$$\nabla \cdot \mathbf{u} = 0 \quad \mathbf{u} = -K(\mathbf{x})\nabla h \quad \mathbf{x} \in \Omega \quad (1)$$

126 where $\mathbf{u}[\text{LT}^{-1}]$ is the Darcy velocity, $K[\text{LT}^{-1}]$ is hydraulic conductivity, h is
 127 hydraulic head, Ω is a two-dimensional spatial domain. Note that, even though
 128 our results are limited to a two-dimensional set-up, they can be extended to
 129 three dimensional systems following the same line of reasoning. In the latter
 130 case, we expect a different quantitative behavior (the dimensionality of the
 131 problem is a crucial factor in the flow organization), but we expect a similar
 132 qualitative behavior. The Darcy velocity has two components, i.e., v and u
 133 which are transversal and longitudinal with respect to to the main flow di-
 134 rection, respectively. We consider in this study flow taking place in a squared
 135 domain of unit size. We impose impermeable boundaries along the right and
 136 left edges, while we set a uniform value of v and of h along the top and bot-
 137 tom boundary, respectively. Equation (1) is numerically solved upon employing
 138 a mixed two-field finite element approach (see e.g., [9]) implemented within
 139 the FreeFem++ environment [14]. We employ a structured triangular spatial
 140 discretization, considering 313600 triangles in order to ensure the accurate res-
 141 olution of the hydraulic conductivity spatial distribution (i.e., one correlation
 142 length of the conductivity field, see below, has been discretized by 16 elements
 143 [29]) and, in turn, of the ensuing velocity fields.

144 We assume here the log-conductivity $Y = \ln(K)$ as a spatially correlated
 145 Gaussian random field. Under such assumption, a reference field \bar{Y} can be
 146 defined through the Karhunen-Loève expansion [21]

$$\bar{Y} = \sum_{n=1}^{\infty} \xi_n \lambda_n f_n(\mathbf{x}) \quad (2)$$

147 where ξ_n are orthogonal Gaussian random variables with zero mean, λ_n and
 148 $f_n(\mathbf{x})$ are the eigenvalues and eigenfunctions that can be used to approximate
 149 the Y field with spatial covariance $C_Y(\mathbf{x}, \mathbf{y})$. In the following we assume an
 150 exponential separable covariance function $C_{Yref}(\mathbf{x}, \mathbf{y}) = \sigma_{Yref}^2 \exp(-|\mathbf{x}_1 -$

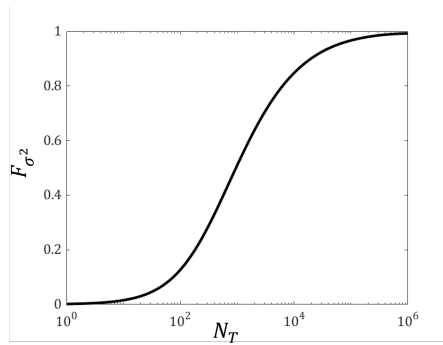


Fig. 1 Number of terms N_T to be included in (3) as a function of the target retained variance in the truncated field F_{σ^2} .

151 $\mathbf{y}_1/|\eta - |\mathbf{x}_2 - \mathbf{y}_2|/\eta)$, being $\sigma_{Y^{ref}}^2$ the variance of the conductivity field, η [L] the
 152 correlation length and with subscript 1 and 2 indicating the longitudinal and
 153 transverse spatial directions, respectively. The correlation length η is chosen
 154 as $1/35$ of the domain size. The analytical expressions provided in [34] are
 155 then used to compute the factors λ_n and the functions $f_n(\mathbf{x})$ in the expansion
 156 (2). The series (2) is employed in practice upon truncation to a finite number
 157 of terms, i.e.,

$$Y = \sum_{n=1}^{N_T} \xi_n \lambda_n f_n(\mathbf{x}) \approx \bar{Y} \quad (3)$$

158 The number of terms N_T included in the expansion can be selected upon
 159 requiring that the generated fields retain a given fraction F_{σ^2} of the reference
 160 field variance, i.e. $\sigma_Y^2 = F_{\sigma^2} \sigma_Y^2$, with $F_{\sigma^2} < 1$. In this context F_{σ^2} is often
 161 considered as an *a priori* indicator of the fidelity of the truncated expression
 162 in reproducing the original field. Figure 1 depicts F_{σ^2} versus N_T . Note that
 163 the value of F_{σ^2} increases monotonically with N_T , i.e. the number of terms
 164 increases with the required level of retained variance. Note that the Karhunen-
 165 Love expansion is valid for all variances of Y and can be applied considering any
 166 dimensionality of the system.

167 In this work we analyze Monte Carlo samples of log-conductivity fields
 168 characterized by three distinct reference variances, i.e. $\sigma_{Y^{ref}}^2 = 0.5, 1.5, 2.5$.
 169 For each of these three cases we consider a reference sample where the fields
 170 are generated upon truncating the expansion when the variance of the corre-
 171 sponding field attains a value of 0.97 of $\sigma_{Y^{ref}}^2$. Then, for each of the considered
 172 $\sigma_{Y^{ref}}^2$ we generate MC samples with F_{σ^2} comprised between 0.5 and 0.97 upon
 173 selecting a number of 1000 realizations per sample. In the following the results
 174 associated with $F_{\sigma^2} = 0.97$ are labeled as reference fields $Y^{ref}, u^{ref}, v^{ref}$. Vari-
 175 ables obtained for $F_{\sigma^2} < 0.97$ are considered as truncated fields and indicated
 176 with Y^{tr}, u^{tr}, v^{tr} . Note that, even though low level of F_{σ^2} (e.g., smaller than
 177 0.8) may be too restrictive in the context of practical applications, we here
 178 cover them in order to highlight the emergence of trends/patterns in the our
 179 results

180 2.2 Information Theory

181 Uncertainty in a discrete random variable, X , can be quantified via Shannon
182 Entropy [30]

$$H(X) = \sum_{i=1}^{N_X} p_i \ln(p_i^{-1}) \quad (4)$$

183 where N_X is the number of bins used to discretize the sample probability
184 distribution, and p_i is the probability associated with the i -th bin. Shannon
185 Entropy can be interpreted as a measure of the uncertainty associated with X ,
186 i.e., $H(X)$ is maximum (assigned the binning) in case p_i is uniform across all
187 the N_X bins while it equals zero when all the values of X fall into one single
188 bin. Note that employing the natural logarithm leads to have *nats* as unit of
189 measure for $H(X)$, other choices for the base of the logarithm are possible.

190 Statistical dependence between two random variables, i.e., X and Z , can
191 be characterized by the reduction in the uncertainty that the knowledge of one
192 variable entails for the other. This is formalized by the mutual information as

$$I(X; Z) = \sum_{i=1}^{N_X} \sum_{j=1}^{N_Z} p_{i,j} \ln(p_{i,j}/(p_i p_j)) \quad (5)$$

193 where N_Z is the number of bins associated with Z , p_j is the probability dis-
194 tribution for Z and $p_{i,j}$ is the joint probability distribution between X and
195 Z . Mutual Information turns out to be null in case of independent random
196 variables, while the equality $H(X) = H(Z) = I(X; Z)$ holds in case that the
197 knowledge of one variable is sufficient to predicted the other one exactly. Mu-
198 tual Information is again measured in *nats* as in (4). It is important to recall
199 here that mutual information is a nonlinear dependence metric, i.e., it is capa-
200 ble of detecting dependence between random variables which are not induced
201 by a linear relationship. On top of mutual information it is then convenient to
202 define the dimensionless uncertainty coefficient (UC) [31]

$$UC(X; Z) = \frac{2I(X; Z)}{H(X) + H(Z)} \quad (6)$$

203 where UC is bounded between zero, for independent variables, and one, when
204 there is an exact dependence between the two variables. In the following the
205 UC will be compared against the well known linear correlation coefficient (or
206 Pearson coefficient), i.e., $\rho(X; Z)$, which captures only the intensity of the lin-
207 ear dependence between two variables. Note that in principle $\rho(X; Z)$ can be
208 both larger or smaller than $UC(X; Z)$, for a given pair of variables $(X; Z)$. The
209 IT metrics employed in our work can be defined for continuous probability dis-
210 tributions where summations and probability mass functions are replaced by
211 integrals and probability density functions, respectively. This approach would
212 be characterized by a less intuitive and immediate interpretation, for example
213 Entropy could be negative, infinite or impossible to evaluate, see, e.g.[6,15].
214 Furthermore, the pdfs of the fluid velocity components are not associated with

215 a known analytical formulation but are here obtained as the result of numer-
 216 ical simulations. Employing a continuous approach would require subjecting
 217 these variables to quantization [6]. In general, the quality of IT metrics esti-
 218 mates increases (in a way which depends on the specific metric) with the level
 219 of quantization of the continuous variables [15]. These consideration lead us
 220 to treat the analyzed variables as discrete ones in line with several previous
 221 studies [28, 12, 24].

222 2.3 Quantities of interest

223 Our analysis considers three quantities of interest, i.e., (i) the log-conductivity
 224 Y and (ii) the two components of the velocity field v (longitudinal to the pres-
 225 sure gradient) and u (transverse). First, we focus on the analysis of punctual
 226 values of these variables, upon considering their distributions associated with
 227 diverse levels of truncation in the KL-expansion, as quantified by F_{σ^2} . This is
 228 obtained through two different types of analyses:

- 229 – we investigate the relationship between (i) the variance of the values of Y
 230 and that of v and u , as well as, between (ii) the Shannon Entropy, for the
 231 same variables, considering the diverse levels of F_{σ^2} .
- 232 – we analyse the relationship between the Linear Correlation Coefficient for
 233 (i) Y^{ref} and Y^{tr} and (ii) v^{ref}, u^{ref} and v^{tr}, u^{tr} , as a function of F_{σ^2} . We
 234 repeat the analysis by focusing on the corresponding uncertainty coeffi-
 235 cients..

236 These statistics are evaluated at each location of the domain, we then exclude
 237 from the analysis locations which have a distance from the boundary smaller
 238 than eight correlation lengths, to avoid the influence of domain boundaries.
 239 Furthermore, in the following we report the spatial averages of the statistics
 240 of interest.

241 The Information Theory metrics detailed in Section 2.2 refer to discrete
 242 random variables. Hereafter, we treat Y and v, u as discrete quantities in (4)-
 243 (6), i.e., we consider their probability mass distributions resulting from empir-
 244 ical frequencies. For each of the three considered $\sigma_{Y^{ref}}^2$, we employed a regular
 245 binning, made of 15 bins, ranging from the minimum to maximum values of
 246 the targeted quantities $Y^{ref}, v^{ref}, u^{ref}$, obtained for the reference field (cor-
 247 responding to $F_{\sigma^2} = 0.97$). The same bins are then employed to discretize the
 248 probability distributions associated with variables obtained for the same $\sigma_{Y^{ref}}^2$
 249 but smaller F_{σ^2} . Note that the number of bins employed in the discretization
 250 of the inspected variables can be thought as the achievable characterization
 251 accuracy (e.g., as dictated by hypothetical instrumental apparatuses). For ex-
 252 ample, our choice of binning of the log-conductivity field could be thought as
 253 be equivalent to consider a system with 15 distinct hydrofacies, each identified
 254 by a bin.

3 Results

Our analysis is subdivided into two main parts. In the first part of our analysis, we aim at quantifying variations in the uncertainty in the punctual values as a function of the number of retained modes in the KL-expansion (2) and at assessing the degree of correlation and mutual correlation between corresponding values associated with diverse levels of truncation. Then, we assess the impact of the retained number of modes on the spatial organization of the two velocity components. To this end, we evaluate the uncertainty coefficient and the linear correlation coefficient between pairs of v and u separated by increasing spatial lags (along the mean flow direction), for a selected F_{σ^2} .

3.1 Analysis of punctual values of velocity and log-conductivity

We analyze here the relationship between the uncertainty of the input field, Y , and of the two velocity components, u and v . We firstly focus on the quantification of uncertainty as rendered by the variance (see e.g., [11]) and by Shannon entropy. Figure 2 displays σ_Y^2 against σ_v^2 and σ_u^2 , for diverse values of F_{σ^2} (indicated by diverse markers) and variance of the reference field $\sigma_{Y^{ref}}^2$ (points indicated by diverse colors). As expected σ_v^2 and σ_u^2 are both directly proportional to σ_Y^2 [11]. In particular, we find that the results associated with a single value of $\sigma_{Y^{ref}}^2$ follow a specific trend, where σ_v^2 increases nonlinearly with σ_Y^2 while σ_u^2 display a trend close to linear (compare trends identified by symbols displayed in the same color in Figure 2a and b). This result is consistent with the idea that variations in the input log-conductivity are more deeply reflected by the longitudinal velocity rather than the transverse one. For both the velocity components the dependence of σ_v^2 and σ_u^2 on σ_Y^2 is specific to the selected value of $\sigma_{Y^{ref}}^2$.

To demonstrate the possible limitations in using the variance to characterize heterogeneity of velocity distributions (see also [5, 18]), Figure 3 reports the frequencies of the reference v for two values of $\sigma_{Y^{ref}}^2$ (note that values are normalized by the average longitudinal velocity). As the level of heterogeneity in Y increases the shape the pdf of velocity tends to display a increasing skewness towards high velocity values. For illustrative purposes, and to emphasize this transition, we compare the probability mass functions of v with that of Gaussian variables with equal mean and variance values. The change in the probability density shape is a result of the nonlinear relationship between log-conductivity and velocity, emerging for increasing $\sigma_{Y^{ref}}^2$, that has been reported in previous literature (e.g., [4]). The results shown in Figures 2-3 therefore suggest that the variance may not be an optimal indicator to characterize the velocity probability distributions for increasing heterogeneity.

For this reason, we resort to the Shannon entropy of v and u , which are shown as a function of σ_Y^2 in Figure 4. This analysis reveals a lack of generality in the relationship between $H(v)$ and σ_Y^2 (as well as, in those between $H(u)$ and σ_Y^2) which are strongly sensitive to the value of $\sigma_{Y^{ref}}^2$. In particular, the

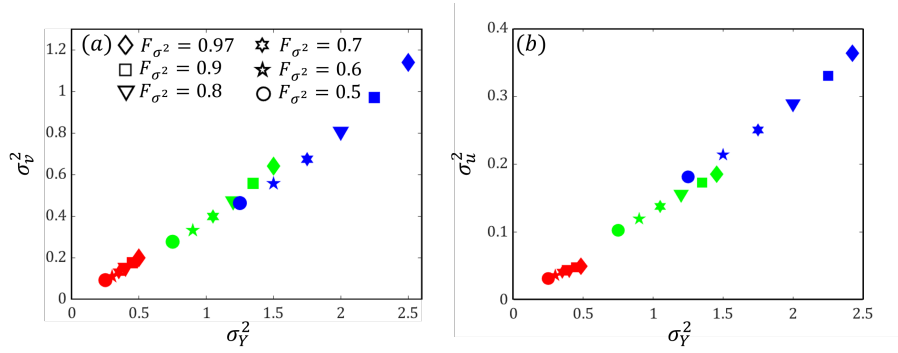


Fig. 2 Relationship between the variance of the log-conductivity, σ_Y^2 , and of the longitudinal (a) σ_v^2 and (b) transverse velocity components σ_u^2 , for $\sigma_{Y^{ref}}^2 = (0.5 \text{ (red)}, 1.5 \text{ (green)}, 2.5 \text{ (blue)})$ and diverse values of F_{σ^2} , indicated by different markers.

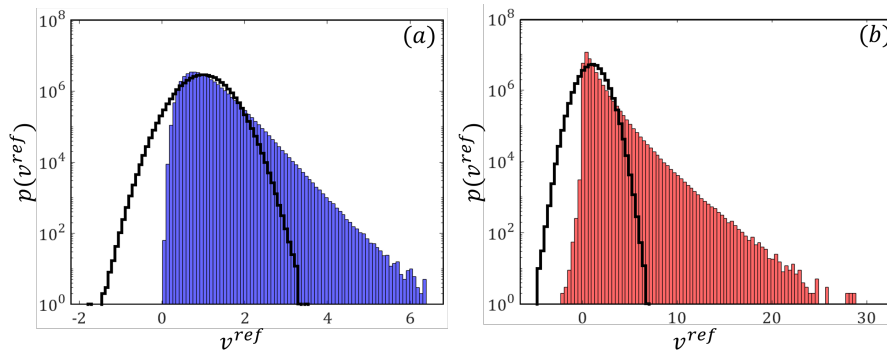


Fig. 3 Probability function for the longitudinal velocity component v^{ref} , i.e., v , for the reference field given (a) $\sigma_Y^2 = 0.5$ and (b) $\sigma_Y^2 = 2.5$. The probability function associated with a Gaussian variable with the same mean and variance is depicted (black curves) for each case.

297 relative change in $H(u)$ and $H(v)$ between the same levels of truncation F_{σ^2}
 298 largely differs as a function of $\sigma_{Y^{ref}}^2$. This result suggests that same relative
 299 change in the input variance is reflected in a different way on the velocity
 300 components pdfs, depending on the heterogeneity of the system.

301 To overcome this issue, we consider Shannon entropy H to describe the
 302 relation between Y and v, u . Figure 5a displays the relationship between
 303 the Shannon Entropy of the log-conductivity, $H(Y)$, and of the longitudinal
 304 velocity, $H(v)$, as a function of $\sigma_{Y^{ref}}^2$ and F_{σ^2} . Assigned a given value of $\sigma_{Y^{ref}}^2$,
 305 we observe that $H(v)$ increases linearly with $H(Y)$. This result shows that the
 306 uncertainty associated with v increases linearly with F_{σ^2} , or equivalently that
 307 $H(v)$ approaches its reference value as a linear function of $H(Y)$. However,
 308 Figure 5 also indicates that the slope of the $H(Y)$ - $H(v)$ linear relationship is
 309 specific to the investigated $\sigma_{Y^{ref}}^2$ values. In particular, we observe that $H(v)$
 310 is inversely proportional to $\sigma_{Y^{ref}}^2$ (assigned a value of F_{σ^2}). This result shows
 311 that the probability mass functions of v are more uniformly distributed across

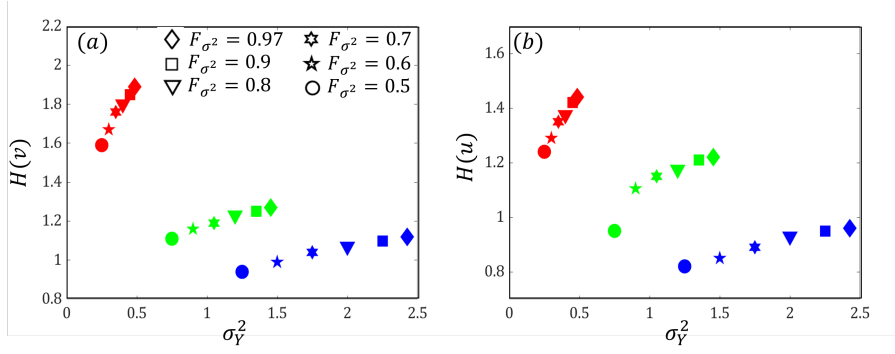


Fig. 4 Relationship between the variance of the log-conductivity, σ_Y^2 , and the Shannon Entropy of (a) longitudinal $H(v)$ and (b) transverse $H(u)$ velocity components, for $\sigma_{Yref}^2 = 0.5$ (red), 1.5 (green), 2.5 (blue) and diverse values of F_{σ^2} , indicated by different markers.

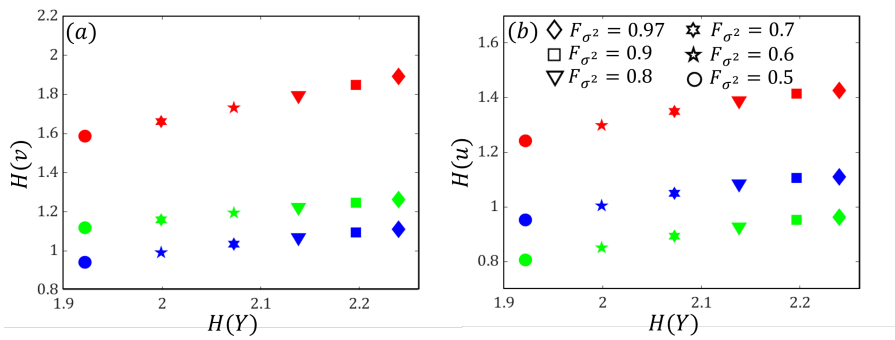


Fig. 5 Relationship between the Shannon Entropy of the log-conductivity, $H(Y)$, and of the (a) longitudinal $H(v)$ and (b) transverse $H(u)$ velocity components, for $\sigma_{Yref}^2 = 0.5$ (red), 1.5 (green), 2.5 (blue) and diverse values of F_{σ^2} , indicated by different markers.

312 their supports as σ_{Yref}^2 decreases, thus rendering increasing $H(v)$ values. Note
 313 that this result is related to our choice of employing a constant number of
 314 bins for all σ_{Yref}^2 , which implies that the bin size is highly variable across
 315 the different samples. In other words, this reflects the higher degree of spatial
 316 homogeneity of the values of v when the system heterogeneity is low, where
 317 the probability density displays relatively uniform values as a function of v .
 318 Qualitatively similar results are obtained for the transverse velocity, as shown
 319 in Figure 5b. Note, however, that $H(u)$ appears to be less sensitive than $H(v)$
 320 to σ_{Yref}^2 , i.e. markers identified with different colors are more close together in
 321 Figure 5b than in Figure 5a (note the different vertical axis scale). Moreover,
 322 the values of $H(u)$ tend to saturate to an asymptotic value for $F_{\sigma^2} > 0.9$. The
 323 analysis suggests that an increase in F_{σ^2} can be expected to have important
 324 effects on the resulting pdfs of v even for $F_{\sigma^2} > 0.9$, while the pdf of the
 325 transverse velocity u only displays a minor shape transition in the same range
 326 of values.

327 We analyze now the relationship between (i) the correlation between the
 328 truncated Y^{tr} and the reference Y^{ref} field and (ii) the correlation between
 329 the ensuing v^{tr}, u^{tr} (associated with the truncated Y) and the reference fields
 330 v^{ref}, u^{ref} . Our objective is to assess the significance of the linear correlation
 331 and of the uncertainty coefficients in quantifying the representativeness (or
 332 fidelity) of the velocity field obtained upon through a truncated field with re-
 333 spect to the reference one. Figure 6 shows the results obtained by associating
 334 each value of linear and nonlinear correlation coefficients computed for the
 335 log-conductivity with the those obtained for the corresponding velocity compo-
 336 nents. Figure 6a suggests that the relationship between $\rho(Y^{ref}; Y^{tr})$ and
 337 $\rho(v^{ref}; v^{tr})$ depends on $\sigma_{Y^{ref}}^2$, with increasing correlation detected for decreas-
 338 ing $\sigma_{Y^{ref}}^2$. On the other hand, Figure 6b shows that the relationship between
 339 $UC(Y^{ref}; Y^{tr})$ and $UC(v^{ref}; v^{tr})$ is insensitive to the degree of heterogeneity
 340 in the reference permeability field and that the two indicators practically co-
 341 incide, i.e. $UC(Y^{ref}; Y^{tr}) \approx UC(v^{ref}; v^{tr})$. The linear correlation coefficients
 342 are considerably larger than the corresponding uncertainty coefficients, for
 343 any considered value of F_{σ^2} and $\sigma_{Y^{ref}}^2$. In particular for $F_{\sigma^2} = 0.9$ we obtain a
 344 $\rho(v^{ref}; v^{tr}) \approx 0.93$ and $UC(v^{ref}; v^{tr}) \approx 0.5$. A similar behavior is observed for
 345 the transverse velocity u , where the maximum $\rho(u^{ref}; u^{tr})$ approaches unity
 346 while the corresponding value of $UC(u^{ref}; u^{tr})$ is below 0.7. Such striking
 347 quantitative differences between the two indicators can be explained upon
 348 observing the bivariate distributions obtained for the pairs $(v^{ref}; v^{tr})$ and
 349 $(u^{ref}; u^{tr})$, reported in Figure 7. In particular, we observe that the largest
 350 velocity values detected in the reference velocity distributions are consistently
 351 not found within the truncated fields v^{tr} . This discrepancy becomes more evi-
 352 dent for decreasing values of F_{σ^2} (see Figure 7b). This result is of potential
 353 relevance to a variety of applications in which the presence of large values of v
 354 (jointly with their spatial organization) plays a crucial role, e.g., early arrival
 355 times of dissolved chemicals at control planes (e.g., [10]), evaluation of the
 356 level of system connectivity (see e.g., [27]) and solute dispersion mechanisms
 357 (see e.g., [18]). Comparison of Figure 6c-d also shows that the truncated fields
 358 tend to underestimate the magnitude of the transverse velocity component.
 359 The implication of this result is that the tortuosity of the resulting stream-
 360 lines would be likely underestimated for the truncated fields with respect to
 361 the reference ones. This behavior would propagate to solute transport dynamics
 362 (see e.g., [5]), possibly resulting in incorrect predictions of solute time arrivals
 363 at control planes and inaccurate quantification of uncertainty.

364 Overall results in Figure 6 suggest that a linear correlation metric over-
 365 estimates the representativeness of the truncated fields with respect to the
 366 reference fields. Therefore, the use of linear metrics would be misleading if
 367 employed to assess the convergence of truncated fields to a higher fidelity ap-
 368 proximations. The low values of $UC(v^{ref}; v^{tr})$ suggests that surrogate models
 369 based on a KL-based expansion of log-conductivity may be practically unable
 370 to faithfully reproduce the full velocity pdfs, particularly as for what concern
 371 the large velocity tail of the distribution.

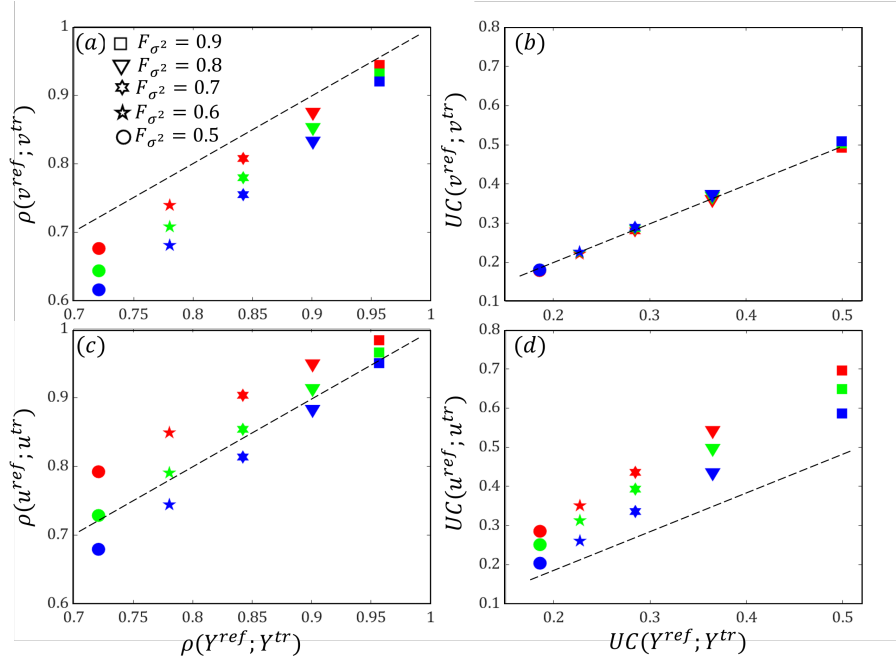


Fig. 6 Correlation analysis between truncated and reference fields quantified by (a) linear correlation coefficients $\rho(Y^{ref}; Y^{tr})$, $\rho(v^{ref}; v^{tr})$, (b) uncertainty coefficients $UC(Y^{ref}; Y^{tr})$, $UC(v^{ref}; v^{tr})$, (c) linear correlation coefficients $\rho(Y^{ref}; Y^{tr})$, $\rho(u^{ref}; u^{tr})$, (d) uncertainty coefficients $UC(Y^{ref}; Y^{tr})$, $UC(u^{ref}; u^{tr})$. Results are for $\sigma_{Y^{ref}}^2 = 0.5$ (red), 1.5 (green), 2.5 (blue) and diverse values of F_{σ^2} , indicated by different markers. Dashed black lines indicate the axes bisectors.

372 3.2 Analysis of spatial correlation

373 We extend here the analysis to the characterization of the spatial auto-correlation
 374 of the velocity components v and u , as rendered by different values of F_{σ^2} , i.e.
 375 with increasing number of modes retained in the KL-expansion. To this end we
 376 evaluate distinct two-points correlation metrics, i.e., (i) the linear correlation
 377 coefficient $\rho(v; v(s))$, $\rho(u; u(t))$ and (ii) the uncertainty coefficient $UC(v; v(s))$,
 378 $UC(u; u(t))$ (see Section 2.2) where s and t are spatial lags in the longitudinal
 379 and transverse directions, respectively. These lags are here normalized by η ,
 380 the correlation length assumed for Y^{ref} .

381 Figure 8 shows the trends of $\rho(v; v(s))$ and $UC(v; v(s))$ against s , for diverse
 382 values of F_{σ^2} , assigned $\sigma_{Y^{ref}}^2 = 0.5$ and $\sigma_{Y^{ref}}^2 = 2.5$. For both indicators we
 383 observe that spatial correlation tends to decrease with increasing F_{σ^2} . This
 384 result suggests that the use of truncated fields enforce larger spatial correlation
 385 than the reference one. We also observe that the linear correlation observed
 386 in the velocity components displays a slight inverse proportionality to $\sigma_{Y^{ref}}^2$,
 387 i.e. increasing heterogeneity reduces the linear correlation of the longitudinal
 388 velocity component, for a fixed lag. On the other hand, $UC(v; v(s))$ appears

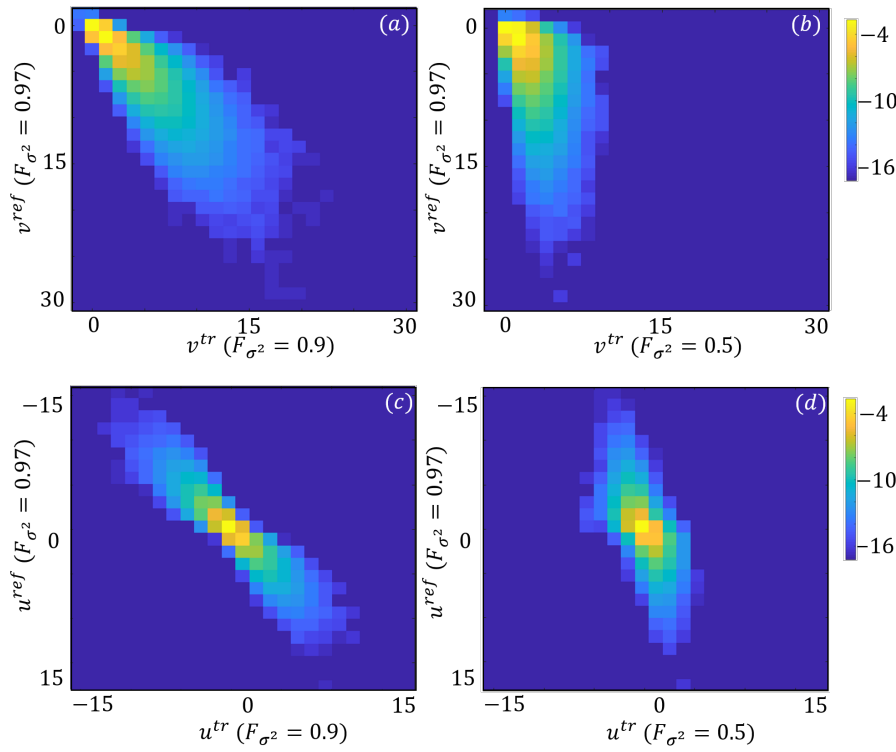


Fig. 7 Bivariate log-probability distributions between reference and truncated fields of the longitudinal v (a)-(b) and transverse u (c)-(d) velocity component. The truncated fields are obtained with (a) and (c) $F_{\sigma^2} = 0.9$, (b) and (d) for $F_{\sigma^2} = 0.5$, for $\sigma_{Y^{ref}}^2 = 2.5$.

389 to be less sensitive than $\rho(v; v(s))$ to the value assumed by $\sigma_{Y^{ref}}^2$. Figure 8
 390 also shows that $\rho(v; v(s))$ tends to be larger than $UC(v; v(s))$ assigned a given
 391 lag s . This result indicates that the linear correlation metric always scores
 392 higher than the nonlinear one, in line with results in Figure 6. To explain this
 393 result, Figure 9 displays the joint pdfs of $v; v(s)$ evaluated at spatial locations
 394 separated by lag of (a) one, (b) three and (c) six for $\sigma_{Y^{ref}}^2 = 0.5$ and (d-f) for
 395 $\sigma_{Y^{ref}}^2 = 2.5$. The joint probability distributions depicted in Figure 9 show two
 396 consistent features: (i) a high density in the upper left corner, aligned with
 397 the upper left - bottom right diagonal, and (ii) a fairly symmetrical dispersion
 398 of the probability density around this diagonal, whose intensity increases with
 399 spatial lag. The definition of $\rho(v; v(s))$ is such that feature (i) turns out to
 400 be the dominant one, i.e., $\rho(v; v(s))$ reflects the marked tendency of the joint
 401 probability density $p_{i,j}$ of displaying dominant terms when $i = j$. The linear
 402 correlation metric is instead poorly sensitive to the actual dispersion of the
 403 joint pdf. On the other hand, $UC(v; v(s))$ detects the dispersion of the joint
 404 pdf, which reflects non negligible probability of having uncorrelated pairs of v .
 405 This feature is emphasized as the spatial lag increases (e.g., compare panels a

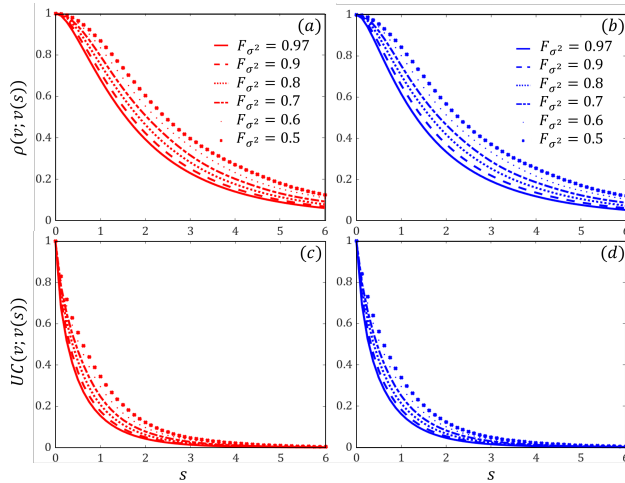


Fig. 8 (a)-(b) Linear Correlation Coefficients and (c-d) Uncertainty coefficients for pairwise values of the longitudinal component of velocity, i.e., $\rho(v; v(s))$, as a function of the spatial lag s , along the mean flow direction for (a,c) $\sigma^2(Y^{ref}) = 0.5$ and (b,d) $\sigma^2(Y^{ref}) = 2.5$. Lags are aligned with v and normalized by the Y^{ref} correlation length η . Diverse level of truncation (i.e., F_{σ^2}) of the KL-expansion are depicted, with different line types.

406 and c in Figure 9). These observations highlight the relevance of considering
 407 suitable summary metrics (such as the uncertainty coefficient) in order to
 408 characterize the whole behavior of the v and $v(s)$ joint pdf.

409 Figure 10 displays spatial correlation of the transverse flow velocity component
 410 u . As previously noted for v (see Figure 8), we observe that the spatial
 411 correlation of u measured by means of $\rho(u; u(t))$ or $UC(u; u(t))$ tends to
 412 decrease with F_{σ^2} , even though in a less marked fashion with respect to v .
 413 Figures 10a-b reveal that $\rho(u; u(t))$ assumes negative values, a feature related
 414 with the conservation of mass imposed by the flow equation (1). Note that,
 415 for $F_{\sigma^2} = 0.5, 0.6$ and t larger than two (approximately) the value of $\rho(u; u(t))$
 416 displays some mild oscillations. This feature is associated with the fact that as
 417 fewer modes are retained in the KL-expansion the ensuing Y^{tr} fields exhibit
 418 a periodic behaviour in space determined by the shape of the retained modes.
 419 Inspection of Figure 10c-d suggests that $UC(u; u(t))$ is smaller than $\rho(u; u(t))$
 420 for assigned $\sigma^2(Y^{ref})$ and F_{σ^2} , in line with the results obtained for v (see
 421 Figure 8). Furthermore, as fewer modes are retained in the KL-expansion (i.e.,
 422 small F_{σ^2}) the values of $UC(u; u(t))$ increases, given t . Inspection of the inserts
 423 in 10c-d allows observing that also $UC(u; u(t))$ has a fluctuating behaviour as
 424 a function of t for $F_{\sigma^2} = 0.5, 0.6$ (see previous discussion of Figure 10 a-b). We
 425 note that the correlation detected by $UC(u; u(t))$ is very similar for the two
 426 considered $\sigma_{Y^{ref}}^2$. This result reflects the fact the joint log-probability of pair-
 427 wise u values display a fairly similar shape in the two cases, as shown in Figure
 428 11. A close comparison of the curves in Figure 10c-d reveals that $UC(u; u(t))$
 429 tends to zero faster for $\sigma_{Y^{ref}}^2 = 0.5$ than for $\sigma_{Y^{ref}}^2 = 2.5$ (see results in log-

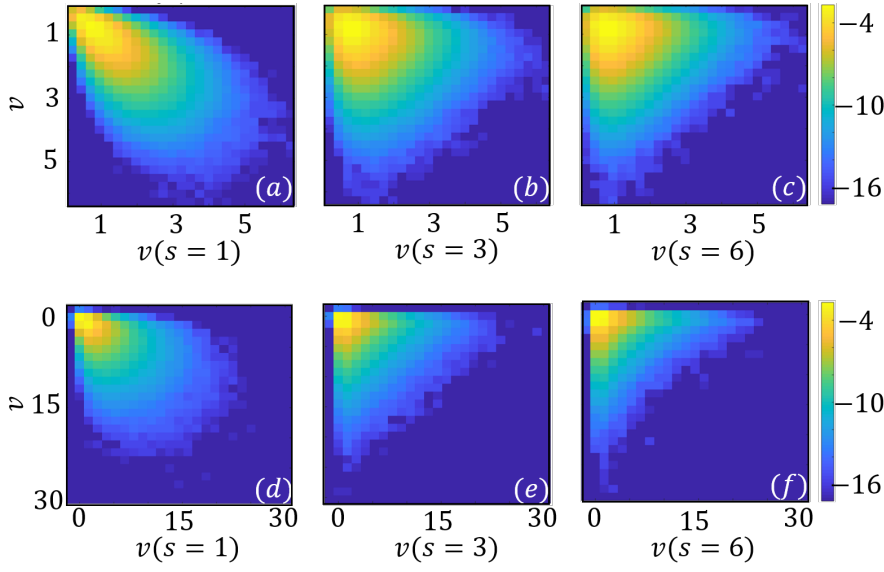


Fig. 9 Joint log-probability distribution of pairwise values of the longitudinal velocity, i.e., v , separated by different spatial lags s for $\sigma^2(Y^{ref}) = 0.5$ and (d-f) $\sigma^2(Y^{ref}) = 2.5$. Lags are aligned with v and normalized by the Y^{ref} correlation length η , results are for $F_{\sigma^2} = 0.97$.

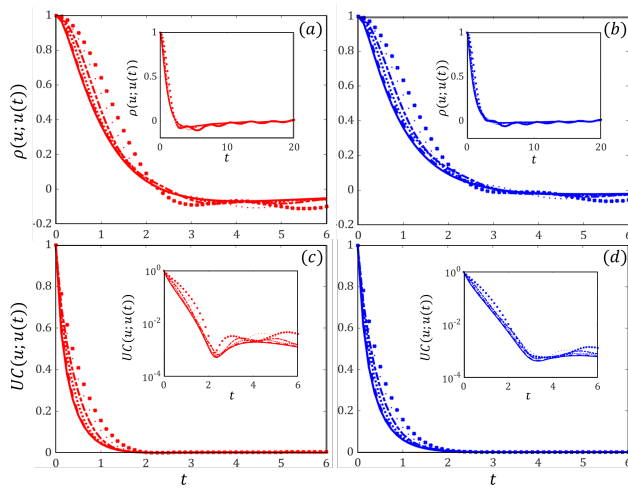


Fig. 10 (a)-(b) Linear Correlation Coefficients and (c-d) Uncertainty coefficients for pairwise values of the longitudinal component of velocity, i.e., $\rho(u; u(t))$, as a function of the spatial lag t , along the transverse direction for (a,c) $\sigma^2(Y^{ref}) = 0.5$ and (b,d) $\sigma^2(Y^{ref}) = 2.5$. Lags are aligned with u and normalized by the Y^{ref} correlation length η . Diverse level of truncation (i.e., F_{σ^2}) of the KL-expansion are depicted, with different line types (see legend in Figure 8). Inserts in (a)-(b) extend results to the interval $t \in [0, 20]$, in (c)-(d) display UC values with a log-axis scale.

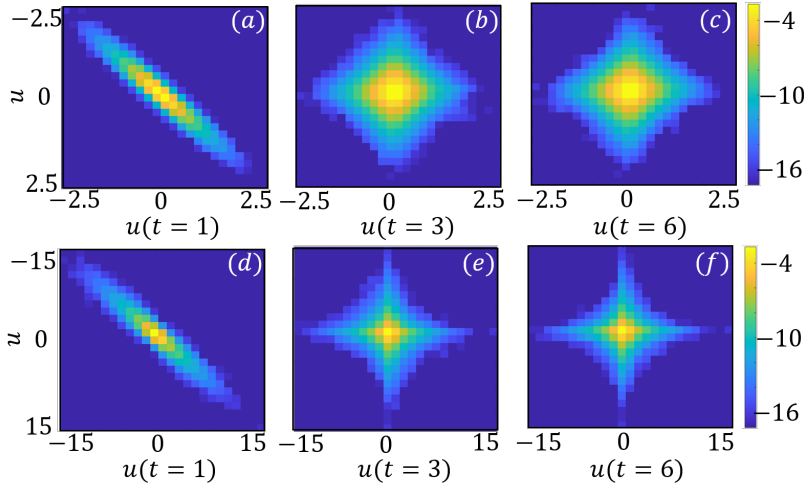


Fig. 11 Joint log-probability distribution of pairwise values of the transverse velocity, i.e., u , separated by different spatial lags t for $\sigma^2(Y^{ref}) = 0.5$ and (d-f) $\sigma^2(Y^{ref}) = 2.5$. Lags are aligned with u and normalized by the Y^{ref} correlation length η , results are for $F_{\sigma^2} = 0.97$.

430 scale reported in the inserts). This result is in line with the distribution $p_{i,j}$
 431 being more evenly distributed across the bins in Figure 11b-c than in Figure
 432 11d-e. These differences appear subtle but may play a role in determining
 433 the spatial correlation of transverse displacement in solute transport models,
 434 which play a relevant role in the upscaling of solute mixing, transport dynam-
 435 ics and associated risk assessment (e.g., [13, ?]). The spatial organization of the
 436 transversal flow component plays a crucial role in the transport dynamics of
 437 dissolved chemical (either passive or reactive ones). As such, it appears crucial
 438 to rely on summary metrics that can quantify the characteristics of joint pdfs
 439 of interest.

440 4 Conclusions

441 In the context of fluid flow within two-dimensional heterogeneous porous media
 442 we assess the impact of diverse degree of truncation of the Karhunen-Loève
 443 (KL) expansion of the log-conductivity (Y) field in terms of the ensuing (*i*)
 444 variability of the velocity field components and (*ii*) of their correlation in space.
 445 Analysis of point wise velocity components indicates that the variance can be
 446 an incomplete descriptor of the uncertainty associated with the velocity field
 447 and suggests to focus on the Shannon entropy as a metric to quantify the degree
 448 of uncertainty associated with both velocity components. In this context, one
 449 of the objectives of our study is to quantify the similarity of truncated input
 450 and output fields to a reference probability distribution. To this end, we tested

451 linear and nonlinear correlation metrics. We find that the linear correlation
452 coefficients, for both Y and the velocity components, are always larger than
453 the corresponding Uncertainty coefficients. This result shows that the linear
454 metric overestimates the degree of similarity with respect to the nonlinear
455 one. A direct implication of this comparison is that linear correlation metrics
456 may lead to misleading results when comparing velocity fields associated with
457 different levels of truncation of the KL-expansion. Our results also suggest that
458 truncated velocity and log-conductivity have the same uncertainty coefficients
459 if compared against their reference fields, a result which can be valuable to
460 assess the relation between input and output variables in the approximation
461 of stochastic groundwater flow outputs.

462 In line with these results, our analysis also shows that the Linear Correla-
463 tion Coefficient typically exceeds the Uncertainty Coefficient when evaluating
464 the spatial correlation between pairwise velocity components. Furthermore, we
465 found that as the number of modes retained in the KL-expansions decreases
466 both velocity components become more correlated in space (according to both
467 the linear and nonlinear metrics, even though with different degrees). These
468 results could be relevant in assessing the correlation structure in solute mass
469 velocity distributions, which are employed in upscaled or effective transport
470 models.

471 Overall our results suggest that it is beneficial to assess the impact of di-
472 verse degree of truncation of the input KL-expansion by means of Information
473 Theory metrics. The latter can provide a more comprehensive assessment of
474 the degree of uncertainty and of dependence (either among point-wise, across
475 levels of truncation, and pairwise velocity components, across space lags) than
476 standard metrics (e.g., variance and Linear Correlation coefficient).

477 IT metrics are capable of capturing salient features of point-wise velocity
478 components and their spatial organization and therefore they may help in
479 quantifying the impact of approximations (such as those entailed by the KL
480 truncation) on analyses of practical relevance. These include, for example,
481 uncertainty quantification and risk assessment for flow and solute transport
482 scenarios. Moreover, our results indicate that IT metrics could be employed
483 upscaling methodologies aimed at characterizing solute transport and mixing
484 in heterogeneous media, which are routinely constrained to observed point-
485 wise and two-points correlation indicators of the flow components.

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