



J Porous Mater  
DOI 10.1007/s10934-017-0492-0

# Advances in the study of nano-structured Co/MCM-41 materials: surface and magnetic characterization

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**Abstract** Co-modified mesoporous supports with MCM-41 structure and several metal loadings were successfully synthesized by a fast wet impregnation method. The nature and location of different Co species formed on the solids were inferred by TEM/SEM, XPS, XANES/EXAFS, adsorption of pyridine coupled to FT-IR spectroscopy and temperature dependence of magnetization. The presence of Co oxide species (clusters and Co<sub>3</sub>O<sub>4</sub> nanoparticles) inside the channels of all the samples could be evidenced by TEM, XPS and XANES/EXAFS. In this sense, the surface Co/Si atomic ratios obtained by XPS were notably lower than the corresponding bulk Co/Si ratios obtained by ICP, indicating that the Co atoms are mostly incorporated inside the mesopores channels of the silica matrix. Nevertheless, by TEM, Co<sub>3</sub>O<sub>4</sub> nanoparticles of small size segregated on the external surface of the silicate were also observed for the higher metal loadings. The temperature dependence of the magnetization performed for the Co/M(2.5) sample allowed to assign its

superparamagnetic behavior to the presence of clusters and Co<sub>3</sub>O<sub>4</sub> nanoparticles of very small size that grow inside the MCM-41 mesopores. Therefore, the analyses presented in this work indicate that a Co theoretical loading of 2.5 wt% leads to the formation of Co oxide nanospecies in the MCM-41 support with a particular superparamagnetic behavior. This sample with improved structural and magnetic properties result an attractive porous solid for drug hosting, to be applied in the field of the controlled release of medicaments.

**Keywords** EXAFS/XANES · Magnetic properties · XPS · Surface properties

## 1 Introduction

Mesoporous molecular sieves modified with metal oxides have different applications such as supports for gas sensors, electrode materials for batteries, supercapacitors, fuels cells, sorbents for separation, media for gas storage and drug delivery [1–6]. In this sense, Co is one of the transition metals extensively studied for their interesting properties. Regarding the novel applications, Co oxide has showed a good electrochemical performance in terms of specific capacity and cyclability [7] and nanosized Co<sub>3</sub>O<sub>4</sub> has exhibited high capacitance as electrode in supercapacitors. Moreover, supported Co oxides on mesoporous materials are of wide interest for their known applications in several areas. Thus, they have been used as catalysts for the Fischer Tropsch reaction, for the combustion of volatiles organic compounds [8, 9] as well as for other kinds of reactions because of their large specific areas and shape selective properties [10]. For its part, in view of their technological importance, the synthesis of magnetic systems with characteristic nanoscale dimensions has attracted much attention in the research field [11,

**Electronic supplementary material** The online version of this article (doi:10.1007/s10934-017-0492-0) contains supplementary material, which is available to authorized users.

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12]. Particularly, magnetic nanoparticles offer some attractive possibilities in biomedicine. One of the most important features of them is the possibility of controlling their size from a few nanometers up to tens of nanometers. Then, they reach dimensions comparable to those of biological entities allowing interacting with them, providing a controllable mean of tagging. In addition, the fact that the nanoparticles have magnetic properties opens the possibility to manipulate them with an external magnetic field gradient in order to make: delivering package such as anticancer drugs to a targeted region of the body such as a tumor, hyperthermia treatments, and magnetic resonance imaging contrast enhancements [13]. Meanwhile, the MCM-41 structure of mesoporous silicates, with a highly ordered structure of channels of well-defined size, is particularly suitable for the synthesis of nanospecies with magnetic properties. Bearing in mind the previous discussion, in this work we extend the study about Co nanospecies with magnetic properties loaded on MCM-41 silicates. Thus, the focus of this work is on its thorough identification, location and acidic properties in relation to their magnetic behavior through the application of specific spectroscopic techniques.

## 2 Experimental

### 2.1 Synthesis

The bare MCM-41 support was synthesized following the method already described in our previous report [14] using cetyltrimethylammonium bromide (CTA) as template and tetraethoxysilane (TEOS) as silicon source. Then, the silicate matrix was modified with Co by the wet impregnation method using several concentrations of  $\text{Co}(\text{NO}_3)_2 \cdot 6\text{H}_2\text{O}$  in aqueous solution in order to reach metal loadings of 1.0, 2.5, 5.0, 10.0 and 15.0 wt%. Then, water was evaporated and the solids calcined at 773 K for 9 h [15]. The calcined materials were designated as Co/M(x) where M indicates the presence of the MCM-41 structure and x indicates the nominal percentage of the metal loading.

### 2.2 Characterization

The Co content was determined by inductively coupled plasma optical emission spectroscopy (ICP) using an equipment VISTA-MPX CCD Simultaneous ICP-OES-VARIAN.

SEM micrographs were obtained in a JEOL model JSM 6380 LV. Gold coverage was applied to make samples conductive. The acceleration voltage was 20 kV. TEM images were obtained in a JEOL Model JEM-1200 EXII System, working voltage: 120 kV. A small drop of the dispersion

(sample in water–ethanol 50% solution) was deposited on a copper grid and then evaporated in air at room temperature.

Co K-XANES/EXAFS spectra were acquired in transmission mode by using a Double Crystal Si(111) monochromator at the D08B-XAFS2 beamline of the LNLS (Campinas, Brazil) [16–20]. Powder samples were mixed with an appropriate amount of boron nitride and pressed into pellets. The incident photon energy was in the range of 7630–8700 eV with an energy resolution about 1.2 eV. The ionization chambers were filled with  $\text{N}_2/\text{Ar}_2$  mixture under appropriate pressure for optimal absorption. For energy calibration purposes a Co metallic foil was placed between the second and the third ionization chamber. The photon energy was calibrated setting the first inflection point to the energy of the K absorption edge of metallic  $\text{Co}^0$  (7709 eV). Data reduction was carried out with Demeter package programs [21].

X-ray photoelectron spectra (XPS) were carried out in a multi-technical system (SPECS) equipped with a dual X-ray source Mg/Al model XR50, which has a hemispherical analyzer PHOIBOS 150 operating in fix analyzer transmission mode (FAT). Spectra were acquired with pass energy of 30 eV, using Al anode operated at 100 W. Work pressure in analysis chamber was lower than  $2.10^{-8}$  mbar. Binding energy values were corrected using as reference the adventitious C 1s line at 284.6 eV. The original XPS spectra were fitted by Gaussian curves using the conventional least squares method. Curve-fitting calculations were useful in determining each band location with confidence levels given by  $R^2 \geq 0.99$ .

FT-IR spectral measurements of pyridine adsorption on the samples were performed on a JASCO FT-IR 5300 spectrometer equipped with a DTGS detector. The range and resolution of acquisition data were 4600–400 and  $4 \text{ cm}^{-1}$ , respectively. A self-supporting wafer for each sample (~20 mg and 13 mm of diameter) was prepared, placed in a thermostatted cell with  $\text{CaF}_2$  windows connected to a vacuum line, and evacuated for 8 h at 673 K. The background spectrum was recorded first after cooling the sample to room temperature. Afterward, the solid wafer was exposed to pyridine vapors (Sintorgan, 99% purity) until the system was saturated to ~46 mmHg at room temperature; the contact time at this pressure was 12 h. The FT-IR spectrum for each sample was obtained after pyridine desorption by evacuation for 1 h at 473 or 573 K. All the spectra were recorded at room temperature before and after pyridine adsorption and desorption at each temperature.

The temperature dependence of the magnetization was measured following the conventional zero field cooling (ZFC) and the field cooling (FC) protocols, using a Quantum Design SQUID magnetometer MPMS XL7 with static field up to  $\mu_0 H = 10 \text{ mT}$ .

### 3 Results and discussion

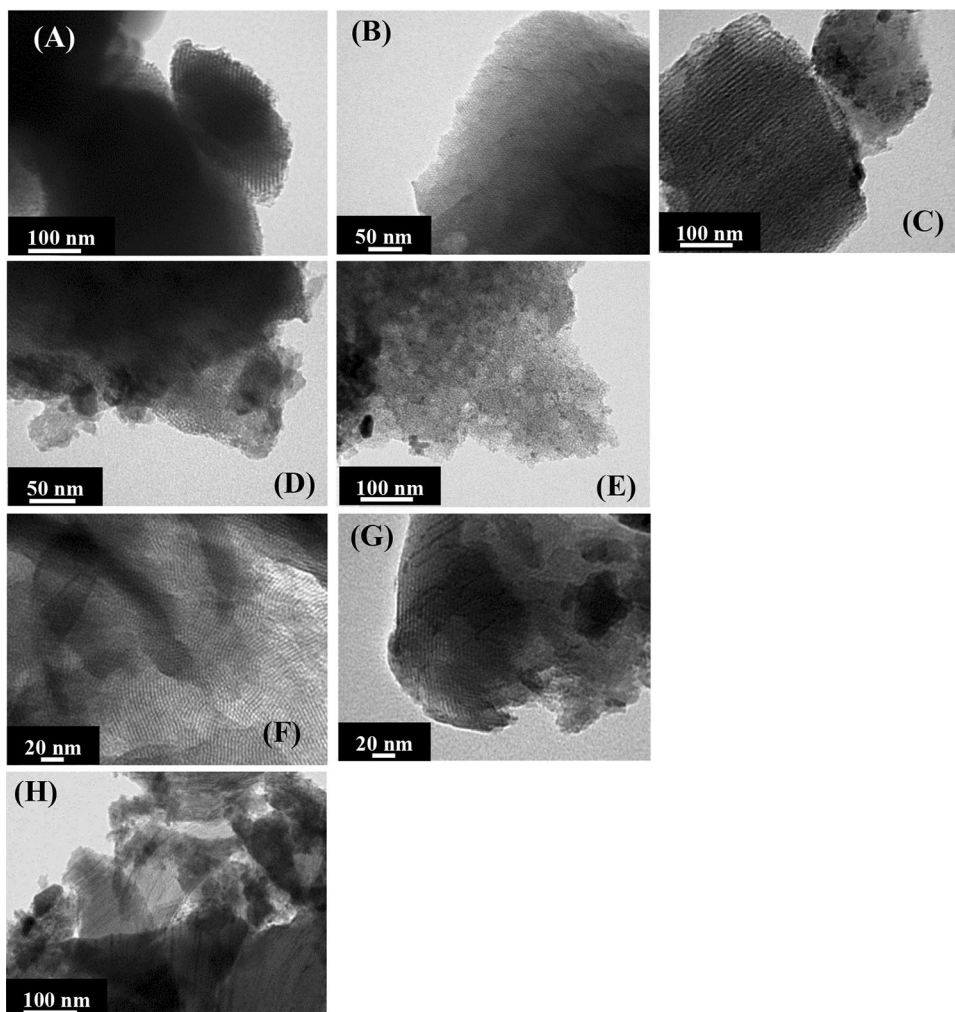
The chemical composition and physical properties of all of the synthesized materials are shown in Table S1. A discussion of these data have been reported previously, and they are also added here in order to support the discussion presented in this report (see also Fig. S1) [15].

Then, TEM images of some Co/M(x) samples are shown in Fig. 1. As it can be observed, it could be confirmed the presence of the well-ordered parallel straight mesochannels, characteristic of the hexagonal pore arrangement, present in the samples. In this sense, the images of Fig. 1c, h are views perpendicular to the direction of the hexagonal pore arrangement, where the straight mesochannels arraying along the long axis are shown [22–24]. For its part, a frontal view of the pores can be observed in Fig. 1a, d, f, g, where the regular order of the channels is preserved without being affected by the presence of Co species. In agreement with these observations, the already reported XRD patterns for Co/M(x) samples (Fig. S2) [15] show that the mesoporous order

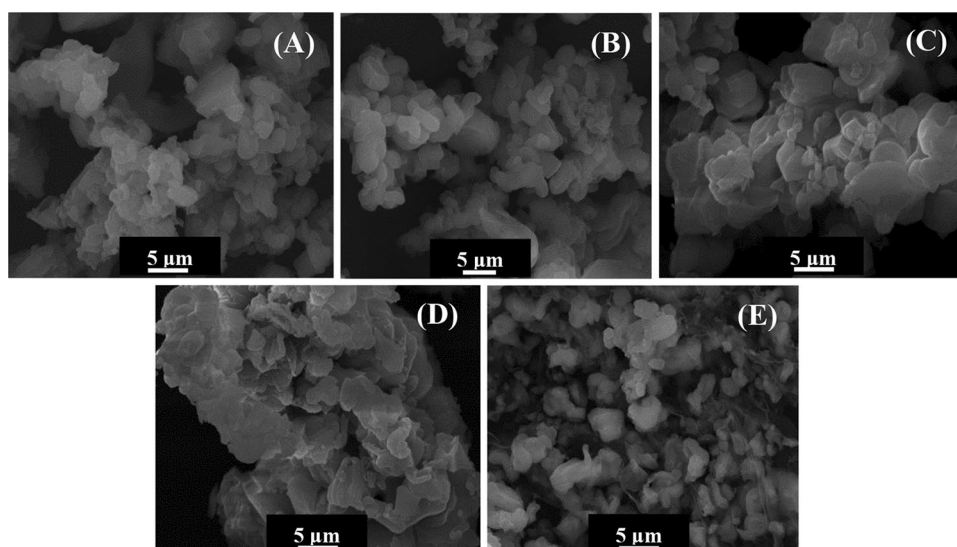
is sustained after the Co loading for all the metal contents. Moreover, no peaks corresponding to Co oxides appear in the high-angle XRD patterns (Fig. S3) for the lower Co loaded samples indicating that these species are amorphous or are clusters/particles too small to be detected by XRD. For its part, since the darker areas in TEM images represent the electronically more dense phases, metal oxide is considered to be present when irregular contrasts in the images are observed [25]. Thus, in Figs. 1g and 1h corresponding to samples with high Co loading, the incorporation of the cobalt oxide species inside the mesoporous channels could be confirmed. In fact, the impregnation with a high amount of Co nitrate followed by a calcination at 773 K, results in the formation of clusters or very small particles of oxide within the mesoporous silica channels evidenced by the darker regions along the mesopores [26].

The particle size and morphology of Co/M(x) samples were investigated by scanning electron microscopy. Figure 2 shows the SEM images of all of the samples. Here, particles with spherical-like morphology aggregated into larger

**Fig. 1** Transmission electron microscopy images: (a–b–c) Co/M(1), (d–e) Co/M(2.5), (f–g) Co/M(5) and (h) Co/M(15)



**Fig. 2** Scanning electron microscopy images: **a** Co/M(1), **b** Co/M(2.5), **c** Co/M(5), **d** Co/M(10) and **e** Co/M(15)

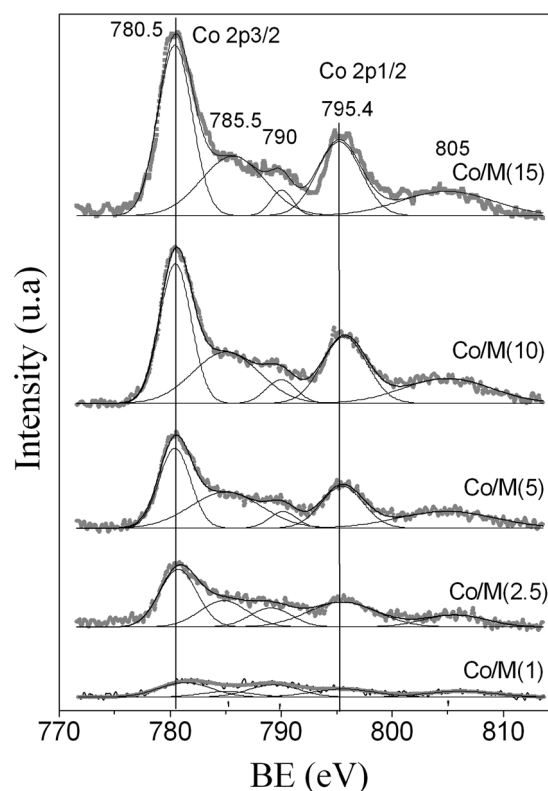


secondary particles that exist in various morphologies and sizes could be observed [28–30].

This irregular morphology could be consequence of the metal oxides deposited on the silica matrix at several locations which provoke the agglomeration. This kind of morphology was already observed for MCM-41 silicates modified with Cr and Ti [27, 31].

In order to study the chemical state and relative abundance of the elements on the solid surface, photoelectron spectroscopy was used. Figure 3 illustrates the Co 2p region of the X-ray photoelectron spectra of Co/M(x) samples with nominal Co loadings from 1.0 to 15.0 wt%. As it could be observed, the Co2p transition splits into two peaks, 2p3/2 (780.5 eV) and 2p1/2 (795.4 eV) with a spin–orbit splitting of around 15.0 eV, which is in accordance with data reported in the literature [32–34].

Nevertheless, it is known that it is difficult to determine the Co cations oxidation states only from the binding energies (BE) of the Co2p main lines, because similar values can be obtained for most of the Co oxides and hydroxides (CoO, Co<sub>2</sub>O<sub>3</sub>, Co<sub>3</sub>O<sub>4</sub>, CoOOH). Then, satellite peak information should be considered. In this sense, it is generally accepted that energy gaps, between the Co2p main lines and the satellites peaks, are highly related to the oxidation states of Co cations. Thus, as it can be observed in Fig. 3, the relatively sharp peak width, the 2p1/2 to 2p3/2 separation of around 15 eV and the energy gaps between the main lines and their satellite peaks indicate the co-existence of Co<sup>2+</sup> and Co<sup>3+</sup> on the surface of the solids. For its part according to the literature, the two satellite peaks detected at about 6 eV (785.5 eV) and 9–10 eV (790 eV) above the Co2p3/2 main line could be assigned to Co cation valence of 2+ and 3+ present in the Co<sub>3</sub>O<sub>4</sub> structure [35]. In this sense, the spectra of the samples synthesized in this work are similar to those corresponding to Co<sub>3</sub>O<sub>4</sub> oxide, which has two types



**Fig. 3** Co 2p core level photoelectron profile of Co/M(x) samples

of Co ions, tetrahedral Co<sup>2+</sup> and octahedral Co<sup>3+</sup> [36]. This behavior was already observed for materials calcined at high temperatures (higher than 673 K) where the weak satellites in the spectra were associated to the presence of the Co<sup>3+</sup> ions in octahedral sites that exist in the spinel cobaltite structure formed onto the mesoporous silicate [36]. Therefore, the XPS technique allowed confirming the presence of Co oxide

in the form of  $\text{Co}_3\text{O}_4$  already evidenced by UV–Vis DR [15]. It is important to note that, due to the low quantity of Co species in the surface (see Table 1), the characteristic Co 2p signals are hardly differentiable in the Co/M(1) sample.

On the other hand, for all of the synthesized solids, Table 1 lists the surface Co/Si atomic ratios determined by XPS and the bulk atomic Co/Si ratios determined by ICP. Generally, the surface Co/Si atomic ratio can be considered as the relative dispersion of transition metal ions on the structure of porous supports. In this sense, the values of this ratio were compared with the bulk Co/Si ratios [37]. For the Co/M(x) samples, the surface Co/Si atomic ratio increases with an increase in the bulk Co/Si ratio in the solids. Nevertheless, for all of the Co contents the surface Co/Si ratio is notably lower than the corresponding bulk Co/Si ratio, indicating that the Co atoms would be mostly incorporated inside the mesopores [26]. These results agree with the already observed by TEM analysis, where it could be observed the mesopores channels of the silica matrix filled with the Co oxide species.

Then, Fig. 4 shows the XPS spectra for the Si2p and O1s signals corresponding to the Co/M(15) sample taken as representative.

The Si2p BE of around 103.8 eV is characteristic of the MCM-41 silica. The XPS O1s spectra consist of two components with BE at around 530 and 533 eV corresponding to oxygen atoms from Co oxide species and from the silica support, respectively. As it can be observed in Table 1, the photoemission O1s for all of the Co/M(x) samples is mainly due to the contribution of oxygen from silica (higher than 90%). This fact would be other evidence of the presence of Co oxide species inside the mesoporous structure of the sieves. Nevertheless, a slight increase in the O1s percentage from the Co oxide species is observed with the metallic loading increasing which is according to the increased presence of the Co oxide species in the surface.

Co K-edge XAFS is a powerful tool for investigating the speciation and coordination structures of Co species formed in the silica mesoporous supports modified by the loading of this metal. Particularly, XANES spectra are frequently used as a “fingerprinting” to recognize a tetrahedral coordination

**Table 1** XPS binding energy (BE) values and surface atomic concentration of Co/MCM-41 samples

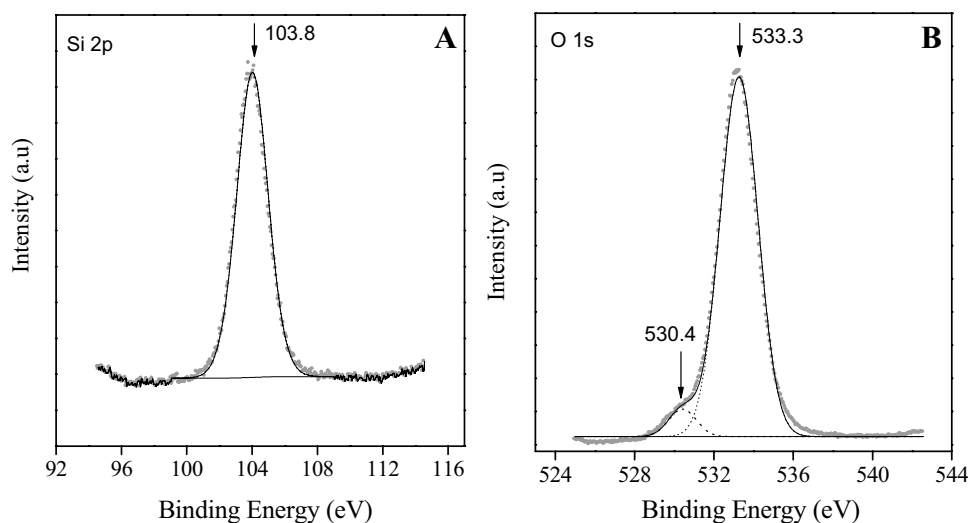
Sample	Co <sup>a</sup> (wt%)	Co bulk (at.%)	Co/Si bulk	Co2p BE (eV)	Co <sup>b</sup> surface (at.%)	Co/Si surface	O 1s BE (eV)	
							From Co <sub>x</sub> O <sub>y</sub>	From silica
Co/M(1)	0.80	0.27	0.008	780.5	0.14	0.004	530.2 (0.42) <sup>c</sup>	533.4 (99.58) <sup>c</sup>
Co/M(2.5)	2.20	0.76	0.023	779.6	0.29	0.009	530.2 (1.11) <sup>c</sup>	533.2 (98.89) <sup>c</sup>
Co/M(5)	4.35	1.52	0.046	779.9	0.46	0.015	530.5 (2.95) <sup>c</sup>	533.3 (97.05) <sup>c</sup>
Co/M(10)	8.81	3.17	0.098	780.2	1.29	0.039	530.4 (4.05) <sup>c</sup>	533.2 (95.94) <sup>c</sup>
Co/M(15)	11.90	4.38	0.137	780.2	1.92	0.060	530.4 (5.24) <sup>c</sup>	533.3 (94.75) <sup>c</sup>

<sup>a</sup>Determined by ICP

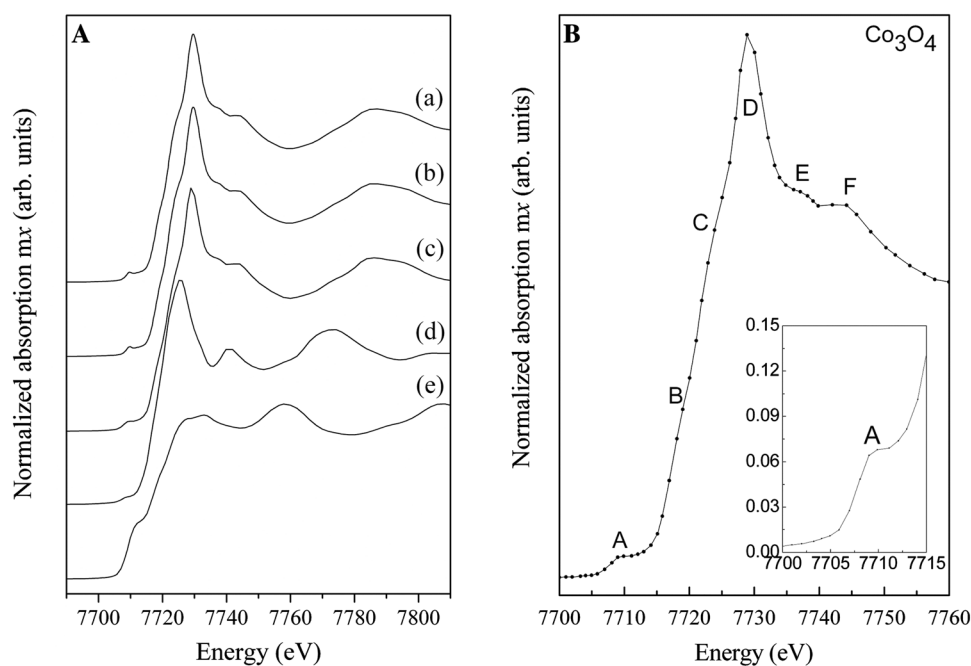
<sup>b</sup>Determined by XPS

<sup>c</sup>Between parenthesis the percentage composition of each species is shown

**Fig. 4** XPS spectra of (a) Si 2p and (b) O 1s for the Co/M(15) sample (taken as representative)



**Fig. 5** **A** XANES spectra at the Co K edge for samples: (a) Co/M(5), (b) Co/M(10) and the following compounds used as references (c) Co<sub>3</sub>O<sub>4</sub>, (d) CoO and (e) CoO. **B** XANES spectra at the Co K edge for the Co<sub>3</sub>O<sub>4</sub> oxide and a detail of their most important characteristics



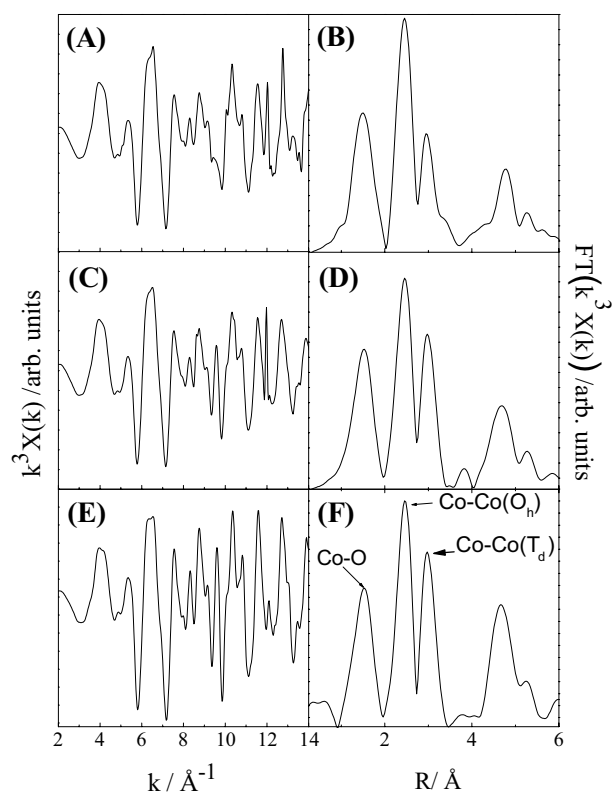
from an octahedral one [38] and changes in the average oxidation states [39, 40]. Figure 5A shows the XANES spectra corresponding to the samples: Co/M(5) (a), Co/M(10) (b) and those from standards compounds used as references (c) Co<sub>3</sub>O<sub>4</sub>, (d) CoO and (e) metallic Co. Figure 5B shows the Co K XANES spectra corresponding to the Co<sub>3</sub>O<sub>4</sub> with spinel structure. In this last figure, the A region (pre-edge region) is closely associated to the 3d electrons and to the symmetry. As it was mentioned before, it is known [41] that, at room temperature, the Co<sub>3</sub>O<sub>4</sub> spinel has the presence of Co ions in tetrahedral and octahedral environments in a stoichiometric ratio of Co<sup>+2</sup> [Co<sup>+3</sup>]<sub>2</sub>O<sub>4</sub>, and that Co<sup>+2</sup> from Co<sub>3</sub>O<sub>4</sub> oxide, tetrahedrally coordinated by O<sup>2-</sup> ions, is different from the octahedrally coordinated Co<sup>+2</sup> present in CoO oxide [42].

For the standard spinel structure analyzed in this work, the A pre-edge peak is attributed to a dipolar forbidden transition 1s–3d characteristic of the tetrahedrally coordinated Co<sup>+2</sup> [43], because this transition is forbidden when there is a reversal center in the metal cation, as it is the case of the non-distorted octahedrally coordination. Moreover, the pre-edge in the Co<sub>3</sub>O<sub>4</sub> spectra is assigned to the presence of tetrahedrally coordinated Co<sup>+2</sup> [44] due also to the peaks of this ion are more intense than those of octahedral Co<sup>+3</sup>. This fact is due to the tetrahedral ligand field allows a dipolar transition, while in the octahedral symmetry, only a quadrupolar transition is allowed for 1s–3d [45]. Another characteristic is that the pre-edge peak of the octahedral site is usually wider than the one corresponding to a tetrahedral site [44, 45]. According to the reported by T. Jiang and D.E. Ellis [46], the A pre-edge region corresponds to the transition 1s–3d already described, the shoulders B and C correspond

to contributions from multiple scattering in conjunction with the transition 1s–4s [44], the D peak corresponds to the allowed dipolar transition 1s–4p, and the resonances E and F correspond to contributions of multiple scattering. These are the main characteristics of the Co K XANES spectra of the spinel Co<sub>3</sub>O<sub>4</sub> and, as it can be observed in Fig. 5A, the spectra of the two analyzed Co/M(x) samples are similar to that corresponding to the Co<sub>3</sub>O<sub>4</sub> oxide. Then, as it was evidenced by the other techniques used here, the presence of Co<sub>3</sub>O<sub>4</sub> spinel on the MCM-41 silicates could be confirmed.

Figure 6 shows the normalized EXAFS signals,  $k^3\chi(k)$ , and their corresponding Fourier transforms for Co/M(5), Co/M(10) and the standard Co<sub>3</sub>O<sub>4</sub>.

Particularly, Fig. 6F shows the features assigned to each maximum of the Fourier transform [47]. The first one corresponds to the Co–O coordination sphere, the second one to the Co(II)–Co(II) coordination sphere (tetrahedrally coordinated), and the third one to the Co(III)–Co(III) coordination sphere (octahedrally coordinated) [48]. By comparison with the signals corresponding to the standard Co<sub>3</sub>O<sub>4</sub> oxide, the Co/M(x) samples analyzed in this work show a decrease in the intensity corresponding to Co–Co(T<sub>d</sub>), which is more notable for the sample Co/M(5), where there is also, a small variation in the intensity related to Co–Co(O<sub>h</sub>). This difference in the peak intensities (mainly for the lower loaded sample) with respect to the standard bulk oxide could be assigned to a certain disorder degree in the Co<sub>3</sub>O<sub>4</sub> nanospecies present as clusters or nanoparticles of very small size, more finely dispersed and strongly interacting with the mesoporous matrix surface [48].

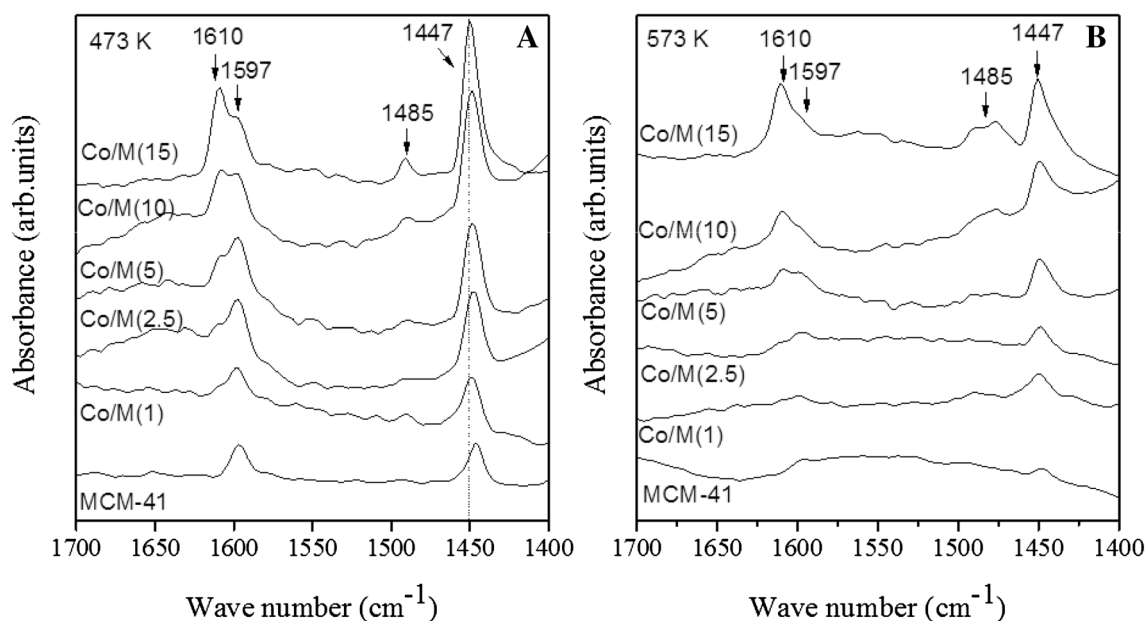


**Fig. 6** EXAFS signal and their Fourier transform for sample (a–b) Co/M(5), (c–d) Co/M(10) and the standard (e–f) Co<sub>3</sub>O<sub>4</sub>

On the other hand, at higher metal loading, the segregation of the oxide nanospecies of bigger size with more defined crystalline structure is enhanced. This could be

related to the fact that the presence of Co<sub>3</sub>O<sub>4</sub> was evidenced by XRD for theoretical Co loadings of 10 and 15 wt% and not for the smaller metal loadings.

Then, the chemisorption of pyridine followed by FT-IR studies was used to detect the presence and nature of acid sites on the synthesized solids [49–52]. Figure 7 shows the FT-IR spectra of the Co/M(x) samples, recorded after the adsorption of pyridine and subsequent evacuation at 473 and 573 K. It is important to note that in order to make a semi quantitative and comparative analysis, all the curves were affected by the pellet mass. It is known that pyridine is a basic molecule that can interact with acid sites via the nitrogen lone-pair electrons giving rise to different characteristic bands. Thus, these bands give information about the strength of Lewis and Brønsted acid sites. It can be observed that all samples show two bands at 1597 and 1447 cm<sup>-1</sup> corresponding to pyridine bonded to silanol groups whose hydroxyls are not capable of protonating pyridine [53–57]. These are the only bands found for the pure MCM-41 which disappear upon evacuation at 573 K indicating the weak interaction between pyridine and the Si–OH groups. For the Co containing solids, a band at 1610 cm<sup>-1</sup> corresponding to pyridine coordinately bonded to Lewis acid sites [50, 53, 56–60] could be observed, which is increasing proportionally with the Co loading. This band would be arising from the formation of a strong “electron–donor–acceptor” adduct between pyridine and Lewis sites, which result enough strong to retain pyridine even at 573 K. In addition, the systematic increase of the band at 1447 cm<sup>-1</sup> and its spectral shift towards higher wavenumber with the growing Co loading, could be interpreted in terms of the overlapping of

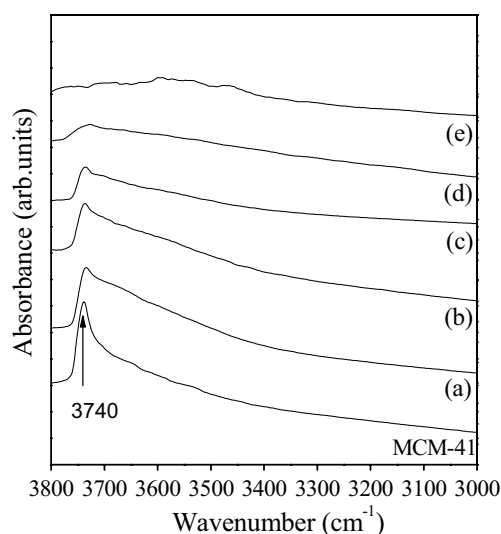


**Fig. 7** FT-IR pyridine adsorbed on the samples Co/M(x) and the bare MCM-41 material after evacuation at: **a** 473 K and **b** 573 K

both the hydrogen-bonded pyridine band ( $1447\text{ cm}^{-1}$ ) and a band attributed to the Lewis-type adduct which frequently appears at  $1450\text{ cm}^{-1}$  [58]. In fact, the retention of this band until  $573\text{ K}$  only for the Co/M(x) samples provides clear evidence that such band has a contribution from Lewis acid sites. Therefore, it is possible to suggest that the incorporation of Co in the siliceous matrix generates Lewis acid sites on the surface, which increase with the Co content. The origin of such Lewis acidity could arise from Co unoccupied molecular orbitals present in both the isolated Co species, coordinated to framework oxygen atoms in the channels, and the Co oxide species. Meanwhile, not all of the samples show bands at  $1540$  and  $1636\text{ cm}^{-1}$  characteristic of Brönsted acid sites [60]. Therefore, the band observed at  $1485\text{ cm}^{-1}$  in the Co/M(x) samples, frequently assigned to pyridine associated with both, Lewis and Brönsted acid sites, corresponds only to Lewis acid sites.

Considering the spectra of the samples in the hydroxyl range (Fig. 8), for samples with Co loading up to  $5\text{ wt}\%$ , it can be observed the presence of a band at  $3740\text{ cm}^{-1}$  corresponding to the stretching vibrations of isolated terminal hydroxyl groups [26, 61–65].

Nevertheless, its decreased intensity with respect to the bare MCM-41 could be attributed to the condensation of some Si–OH groups with Co species leading to Si–O–Co bonds as well as to their partial blocking by the presence of some Co oxide clusters [66]. In addition, a broad contribution between  $3700$  and  $3400\text{ cm}^{-1}$  that appears for the Co/M(x) samples with lower metal content could be assigned to the interaction between vicinal OH groups, as Si–OH and Co–OH groups [67]. Finally, the fact that the absorbance corresponding to the OH groups is strongly diminished



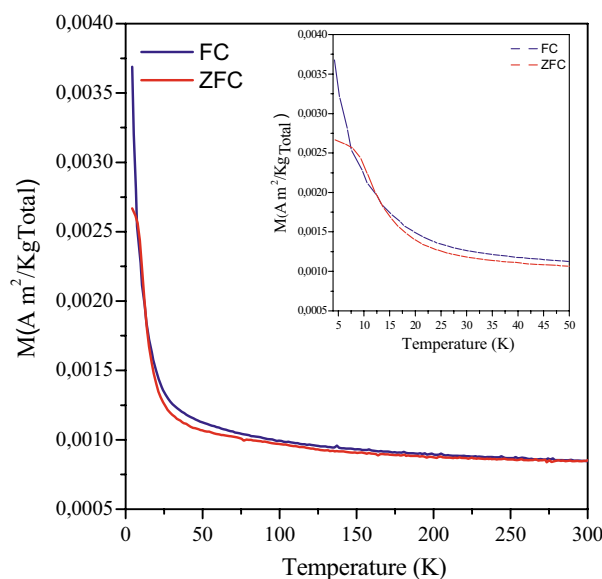
**Fig. 8** FT-IR spectra in the hydroxyl stretching region after degassing at  $673\text{ K}$  for samples: (a) Co/M(1), (b) Co/M(2.5), (c) Co/M(5), (d) Co/M(10), (e) Co/M(15)

for the samples with higher Co content [Co/M(10) and Co/M(15)] is likely due to enhanced amount of oxide nano-species which could be blocking the silanols, besides contributing to the Lewis acidity [68, 69].

In a previous report [15], the analysis of the curves of magnetization vs applied field, corresponding to the Co/M(x) samples, showed that the superparamagnetic contribution exhibit a maximum for a theoretical Co loading of  $2.5\text{ wt}\%$ . This behavior could be attributed to the presence of small Co nanospecies finely dispersed on the matrix. These species become larger for the higher Co loadings resulting in the decrease of the superparamagnetic contribution. In order to continue with this analysis, measurements of the temperature dependence of the magnetization were performed. Figure 9 shows the *zero-field-cooling* (ZFC) and *field-cooling* (FC) curves for the Co/M(2.5) sample, measured with an applied field of  $0.01\text{ T}$ . For this sample, a reversal behavior was observed for temperatures higher than  $\sim 150\text{ K}$ . Under this temperature, the magnetic system shows irreversibility with a maximum in the ZFC curve ( $T_B = 4\text{ K}$ ), which is associated with a distribution in the blocking temperatures of nanoparticles of different size. Therefore, considering an Arrhenius type relaxation for non-interacting nanoparticles, the  $\langle T_B \rangle$  corresponds to the blocking of an average volume:

$$\ln(\tau_m/\tau_0)T_B k_B = KV$$

where  $K$  is the effective anisotropy constant,  $k_B$  is the Boltzmann constant,  $\tau_m$  is the measure time ( $\tau_m \sim 100\text{ s}$  for SQUID measurements) and  $\tau_0$  is the measurement time characteristic of the nanoparticles ( $\tau_0 \sim 10^{-11} - 10^{-9}\text{ s}$ ). Takada et al. [70] reported that superparamagnetic nanoparticles of  $\text{Co}_3\text{O}_4$



**Fig. 9** Temperature dependence of the magnetization in the ZFC and FC states for an applied field for Co/M(2.5)



dispersed on amorphous SiO<sub>2</sub> matrix, with a size of around 3 nm, showed a T<sub>B</sub> of 3.4 K and K of 9 × 10<sup>5</sup> erg cm<sup>-3</sup>. In this sense, the T<sub>B</sub> of around 4 K observed for the Co/M(2.5) sample (Fig. 9) allowed to infer a nanoparticle size consistent with the presence of Co<sub>3</sub>O<sub>4</sub> inside the mesoporous of the MCM-41 matrix. This fact agrees with the statements arising from the results obtained by the other techniques. Considering the expected magnetic behavior in solids in order to be used in the controlled drug delivery systems a Co loading of 2.5 wt% lead to an adequate material for hosting drugs in their structure. Thus, an improved performance in the field of the controlled release of medicaments could be achieved.

#### 4 Conclusions

Several specific spectroscopic techniques were carefully applied in order to infer on the nature and location of different Co species formed on MCM-41 mesoporous silicates by a fast post synthesis method. All the analyses evidenced the presence of Co oxide nanospecies inside the pores of the mesostructure. Thus, TEM images showed that the pores are filled with the Co species from the darker regions along the channels. Nevertheless, for the higher Co loadings, Co<sub>3</sub>O<sub>4</sub> nanoparticles segregated on the external surface of the silicate were also observed. In accordance, XPS evidenced the co-existence of Co<sup>2+</sup> and Co<sup>3+</sup> on the surface of the solids which could be assigned to the spinel cobaltite structure (Co<sub>3</sub>O<sub>4</sub>) formed onto the mesoporous silicate. For all of the Co loadings, the surface Co/Si atomic ratios were notably lower than the corresponding bulk Co/Si ratios, indicating that the Co atoms are mostly incorporated inside the mesoporous channels of the silica matrix. In addition, comparing the XANES spectra of the Co/M(x) samples with some standard substances of Co, the presence of Co<sub>3</sub>O<sub>4</sub> spinel on the MCM-41 silicates could be confirmed. With respect to the acidic properties of the synthesized materials, all of them showed a Lewis acidity increasing with the metal content. The presence of Lewis acid sites was assigned to both, the isolated Co species coordinated to framework oxygen atoms in the channels and Co oxide species. Finally, the results of the spectroscopic characterization and the presence of a T<sub>B</sub> of around 4 K in the ZFC curve for the Co/M(2.5) sample, suggest that the Co oxide nanospecies of small size (clusters and Co<sub>3</sub>O<sub>4</sub> nanoparticles) formed inside the MCM-41 mesopores are the responsible of notable superparamagnetic behavior. Such feature gives rise to solids with interesting properties for novel applications such as “drug hosting” for subsequent delivery through an external magnetic field. Therefore, these porous solids with improved performance result very promising in the field of the controlled release of medicaments.

**Acknowledgements** The authors are grateful to UTN-FRC and CONICET for the financial support. The authors also thank Ph.D Eliana Vaschetto and Ph.D Silvia Mendieta for assistance with the TEM/SEM micrographs. Thanks are given to ANPCyT for the purchase of the SPECS multitechnique analysis instrument (PME8-2003).

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