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Preparation of atomic layer deposited vanadium dioxide thin films using tetrakis(ethylmethylamino) vanadium as precursor

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Abstract

Vanadium dioxide (VO₂) thin films were deposited by atomic layer deposition (ALD) using tetrakis(ethylmethylamino) vanadium precursor and H₂O oxidant at a temperature of 150 °C. Optimization of post-deposition annealing results in smooth, continuous VO₂ films (thickness, t ~ 30 nm) with small grains, exhibiting a transition from semiconducting to metal phase, typically known as the metal-insulator transition (MIT), at ~ 72 °C with a switching ratio ~ 10^2 . Such films were produced with high repeatability on a wafer scale and have been successfully utilized in resistively-coupled oscillators and self-selected resistive devices. Under smaller process window, thin films (t ~ 30 nm) with very large grains have also been produced, exhibiting the MIT ratio ~ 10^3 , which is the highest achieved for the ALD VO₂ films deposited on SiO₂ substrates. Both types of films were characterized again after 120 days to access their stability in air, a property that was rarely investigated.

Keywords: Vanadium Dioxide; Atomic Layer Deposition; Metal Semiconductor Transition; Stability

I. INTRODUCTION

Vanadium dioxide, VO₂, has attracted increasing attention recently due to its remarkable metal-insulator transition (MIT) at a transition temperature ~ 68 °C.^{1, 2} Recently, this material has been utilized in a wide variety of devices such as the smart windows,³ sensors,⁴ steep-sloped devices,^{5, 6} tunable capacitors,⁷ optical limiting modulators,⁸ memory adaptive filters,⁹ and neuromorphic computing.¹⁰ VO₂ has been grown by various methods, including pulsed laser deposition (PLD)¹¹ and rf magnetron sputtering.¹² While PLD can produce excellent films exhibiting a MIT of more than 3 orders of magnitude, this method is limited to small substrates. On the other hand, sputtering is scalable to large area, but control of film stoichiometry is not trivial.¹³ Atomic layer deposition (ALD) is therefore of interest due to the control of stoichiometry arising from its self-limiting growth, its repeatability and conformality over large substrate areas.¹⁴ Even though the deposition rate is much lower in ALD than in sputtering or PLD, the thickness requirement for devices these days are increasing smaller, making the ALD method more attractive both in research and industry.

Amongst various ALD precursors that have been developed VO_2 . tetrakis(ethylmethylamino) vanadium (TEMAV) has been increasingly used since it is a liquid precursor and the synthesized films mainly consist of +4 valency, leading to the growth of pure phase VO₂. ¹⁵⁻²⁰ Using this precursor, the window of the deposition temperatures for VO₂ has been reported to be between 125 and 175 °C. 15 The VO2 thin films were successfully deposited by ALD on various substrates, such as glass, SiO₂, Al₂O₃ and GaN.^{17, 21-23} The asdeposited ALD films are amorphous due to the low deposition temperature used, without a MIT, thus a post-deposition annealing is required. Many groups have reported the optimizations of the annealing conditions (temperature, time, ambient) and have shown that the annealing conditions are very important to achieve the desired VO₂ phase. 16, 17, 21 The

amount of MIT modulation, which is typically determined electrically or optically, depends very much on the substrates; for example, a resistivity modulation $\sim 10^2$ has been reported for the VO₂ films deposited on SiO₂ substrates, ^{17, 21} but a higher modulation $\sim 10^3$ has been reported for the VO₂ films deposited on GaN substrates. ²²

One aspect that previous reports have not paid much attention to is the production of the VO₂ films on a wafer scale, which is highly desirable for industry. However, it is also desirable in research community for easier incorporation of the VO₂ thin films in devices, as the typical fabrication steps, such as photolithography, are easier to process using full wafers than small pieces of substrates. However, it is difficult to achieve a good uniformity using a cross-flow reactor (for example, Savannah ALD100 system from Cambridge Nanotech), especially when using precursor like TEMAV which has a very low vapour pressure (~ 1 Torr at 107 °C).¹⁴

Therefore, we developed VO₂ thin films by ALD and post-deposition annealing with the emphasis on using (a) larger substrates up to a 4- or 8-inch wafer, and (b) the Si/SiO₂ substrates that will allow possible integration onto a CMOS platform. We have recently reported that employing TEMAV multiple pulsing mode (compare to a typical single pulsing) results in thin films with a higher deposition rate and a smaller chemical consumption without compromising the thickness uniformity.²⁴ The films we deposited have been successfully incorporated in resistively-coupled oscillators for neuromorphic application^{25, 26} and selector devices.²⁷

Here, we report the optimization of the post-deposition annealing conditions, employing a quartz tube furnace for annealing of small substrates and a Fiji ALD chamber for annealing of large substrates. The structural, optical and electrical properties of the VO₂ films that can be typically produced by the ALD are presented. Moreover, we also present the VO₂ thin films

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is the author's peer reviewed, accepted manuscript. However, the online version of record will be diffe PLEASE CITE THIS ARTICLE AS DOI: 10.1116/6.0000353 which have much larger grains and higher MIT modulation than the typical films and compare the two types of films in terms of their stability in air and potential use in devices.

II. EXPERIMENTAL DETAILS

VO₂ thin films were deposited using a cross flow reactor system (Savannah S100 ALD from Cambridge Nanotech) on various substrates, such as Si, SiO₂ and quartz, with various substrate sizes up to 4-inch diameter wafer. The multiple pulsing of TEMAV is done by the following sequence: (TEMAV pulse 0.02 s -delay 2s) repeat 8 times /purge 3 s /H₂O pulse 0.03 s /purge 5 s. The TEMAV cylinder was heated to 105 °C, yielding a vapor pressure ~ 1 Torr. During depositions, the chamber temperature was set to 150 °C, and the pressure is ~ 0.1 Torr which was achieved by N₂ flow at 20 standard cubic centimeters per minute (sccm).

Post-deposition annealing was performed either in a 1-inch diameter quartz tube furnace (Lindberg/Blue M, Thermo Scientific) or in a Fiji F200 ALD System from Cambridge Nanotech. The following annealing parameters were varied: temperatures (450 and 550°C), pressures (between 8 and 150 Pa), ambient conditions (O₂ or Ar/O₂) and annealing times (between 0.5 and 2 hours).

The structural and optical properties of the VO₂ films were analyzed by room temperature Raman spectroscopy (Renishaw InVia), X-ray diffraction (Bruker D8 Advance), scanning electron microscopy (FEI Magellan 400), atomic force microscopy (Agilent 5400) and UV-Vis-NIR optical Spectrometry (Cary 7000). Wafer maps of the film thicknesses were obtained using spectroscopic ellipsometry (J.A. Woollam M-2000), integrated with an automated xy stage and a 75 W high intensity Xenon arc lamp for the spectrum of light between 0.7 to 5 eV. The film thicknesses were measured at 49 locations over a 4-inch diameter wafer. The thickness uniformity is determined by the ratio of the standard deviation over the mean value of all points taken over the wafer, multiplied by 100% and subtracted from 100%.

To determine the MIT characteristics, the van der Pauw resistivity measurements were performed using a Keithley 4200 parameter analyzer and a prober with a home-built heater/substrate holder which temperatures were varied between 25 and 120 °C.

III. RESULTS

A. Physical properties

Raman spectroscopy is a fast, non-destructive method employed here as a preliminary test to probe the lattice vibrations in the VO₂. Figure 1(a) shows the Raman scattering spectra of the VO₂ films under various annealing conditions in a quartz tube furnace. Thin films annealed at 450 °C under vacuum (chamber pressure at 0.3 Pa) exhibit no discernible peaks, indicating that the films are still amorphous. On the other hand, thin films annealed in air at 550 °C result in pure V₂O₅ phase, in agreement with the previous studies.²⁸ Another source of reference for V⁵⁺ peak was obtained by the analyzing V₂O₃ powder from Alfa Aesar, before and after annealing in air at 550 °C which are also shown in Fig. 1(a). Upon annealing at 450 °C at 5 Pa in O₂, various peaks are observed; with the main peaks at 195, 224, 309, 399 and 616 cm⁻¹, which can be identified as those related to VO₂.^{29, 30} The large peak at ~ 510 and the broad peak at ~ 950 cm⁻¹ are related to the silicon substrate. When the O₂ pressure is increased to 100 Pa, films with mixed phase (VO₂ and V₂O₅) are formed.

It appears that if a vacuum or lower O₂ pressure was used, a higher temperature (up to ~500 °C) is required to produce VO₂, as also shown by Zhang *et al.*¹⁸ The evolution of film seems to follow from amorphous to polycrystalline phases of V⁴⁺, and then V⁵⁺. It is also noted that this evolution and the annealing parameter space are very similar among films synthesized using TEMAV.¹⁶⁻¹⁸

As shown in Fig. 1(b), due to the temperature gradient along the length of the quartz tube, the actual area that is suitable for annealing of VO_2 is rather small (the rectangular boxes measuring $\sim 2 \text{ cm} \times 4 \text{ cm}$). Therefore, for higher throughput and annealing of whole wafers, a

Fiji ALD system, with a capacity for substrate size up to an 8-inch diameter wafer, was used. As shown in Fig. 2(a), VO₂ samples (2 cm × 2 cm in size, total 9 pieces) were spread over the substrate holder of the ALD system, and the films were annealed at 480 °C for 2 hours under an Ar/O₂ flow at 200/200 sccm at a chamber pressure of 40 Pa. Figure 2(b) shows the Raman scattering spectra of the VO₂ thin films at various positions, and the peak intensity at 616 cm⁻¹ are shown in Fig. 2(a). Direct correlation between the peak intensity to the location of the sample is not done due to error within measurement, for example, the peak intensity will change slightly during laser warm up. However, Fig. 2 shows that the annealing was successfully performed in a Fiji ALD system for area of an 8-inch diameter wafer. The Raman peaks of VO₂ and V₂O₅ from both annealing systems and those from literature are compared in Table 1, generally the Raman data between our work and others agree well.²⁸⁻³⁰

The crystallinity of the synthesized films was analyzed by XRD. In Fig. 3, both the tube furance and Fiji appealed samples show a dominant peak at 20 = 27.0° which is the (011)

The crystallimity of the synthesized films was analyzed by XRD. In Fig. 3, both the tube furance and Fiji annealed samples show a dominant peak at $2\theta = 27.9^{\circ}$ which is the (011) reflection. The thicker film also shows a smaller peak at $2\theta = 39.8$ and 57.7° which are the (002) and (022) reflection (JCPDS card: 00-044-0252). These results confirm the formation of monolinic VO₂(M). There are other VO₂ polymorphs, such as VO₂(A), VO₂(B) and VO₂(R). The tetragonal VO₂(R) is achieved upon thermal heating of VO₂(M) films above 68 °C. The semi-metallic VO₂(B) is an intermediate phase that is formed under narrow annealing conditions. The insulating VO₂(A) has rarely been reported in the ALD deposited films. However, the growth of all three polymorphs, VO₂(A, B, and M) has been demonstrated by PLD on the SrTiO₃ substrates. The synthesized films are substrates.

B. Electrical properties

The resistivity values, ρ , as a function of temperature and the corresponding derivative, $d(\log \rho)/dT$, of the VO₂ thin films after annealing in quartz tube furnace or Fiji ALD system are shown in Fig. 4(a) and (b), respectively. A clear MIT modulation by temperature, as

shown by the low resistivity (metal) at high temperatures and high resistivity (insulator) at low temperatures, was observed. For the VO₂ thin films that were subjected to tube furnace annealing (Fig. 4(a)), the transition temperature during a heating cycle, T_h , is 72.0 °C (transition from insulator to metal), and the transition temperature during a cooling cycle, T_c , is 60.3 °C (transition from metal to insulator). For the VO₂ thin films that were subjected to annealing in Fiji system (Fig. 4(b)), the phase transitions occur at lower temperatures, with T_h and T_c at 60.5 °C and 47.5 °C, respectively. From the FWHM of the derivative profiles, annealing in Fiji system shows a narrower phase transition. The resistivity modulation is $\sim 10^2$ for both types of annealing, which is in the same order as that reported by Rampelberg *et al.*, ¹⁷ who also used the SiO₂ substrates. It should be noted that the resistivity modulation is dependent on the type of substrates used. An order of magnitude improvement on the resistivity modulation has been reported when the VO₂ is deposited on the latticed matched GaN substrates. ²²

C. Optical properties

As shown in the inset of Fig. 5(a), the optical transmission spectrum of VO₂ with thickness ~ 30 nm exhibits a transmittance ~ 50% at wavelength, λ , ~ 500 nm (visible region) and up to ~ 90% at λ ~ 2000 nm (near infrared region), which are consistent with previous reports.³¹ The physical appearance of the films grown on quartz substrates are greyish. The optical band gap, E_g , is extracted using Tauc's method using the relation:³²

$$\alpha h v = B \left(h v - E_{\sigma} \right)^{n}, \tag{1}$$

where α is the absorption coefficient, $h\nu$ is the photon energy, E_g is the band gap, B is a frequency dependent constant and n is either 0.5 or 2 (direct or indirect allowed transitions). The band gap can be deduced from the intercept of the extrapolated linear part of the plot of $(\alpha h\nu)^{0.5}$ vs. $h\nu$. A smaller gap ~ 0.55 eV and a larger gap ~1.35eV can be extracted (Fig. 5a),

which are consistent with the previously reported values.³³ The larger gap can be assigned to the onset of transitions from O 2p to V 3d states, and the smaller gap ~ 0.55 eV to the d-d transitions in the conduction band.^{2, 34, 35} Another way to determine band gap is the E04 method. As shown in Fig. 5(b), at the absorption coefficient of 10^4 cm⁻¹, the gap is 0.65 eV, broadly in agreement with the value determined by Tauc method in Fig. 5(a).

D. Thickness uniformity and film morphology

Figure 6 shows the thickness distribution of the VO₂ films over a 4-inch diameter wafer. These films were deposited using multiple pulsing of TEMAV as described in a previous paper, with a GPC of 0.06 ± 0.02 nm.²⁴ To improve the thickness uniformity, we rotated the wafer 180° halfway through the deposition. A deposition of 450 ALD cycles produces a film with average thickness of 25.1 ± 0.8 nm and thickness uniformity of 88.5% as shown in Fig. 6(a). Likewise, a deposition of 890 ALD cycles produces a film with average thickness of 47.7 ± 1.0 nm and a slightly better uniformity of 90.3% as shown in Fig. 6(b). The uniformity achieved within-wafer uniformity up to 98% has been previously reported for VO₂ deposited on a 12-inch diameter wafer.¹⁶ Such a good uniformity is possible due to efficient precursor gas distribution over the substrate through a shower head in the commercial ALD tool, the EX reactor from Tokyo Electron Limited. We showed here that a modest thickness uniformity could still be achieved using a cross flow type ALD system.

Thin films deposited by ALD are typically very smooth which is very useful for large area electronics. Amorphous ALD VO₂ thin films show a typical root mean square (rms) roughness of only 0.27 nm for a scan area of 2 cm \times 2 cm by the AFM (Fig. 7(a)). However, the required post-deposition annealing roughens the surface, with the emergence of polycrystallinity. Figure 7(b) shows a VO₂ thin film (t \sim 30 nm) with a roughness rms of 2.67 nm. The SEM micrograph in Fig. 7(c) shows a smooth, continuous film with a mixture of small, round grains (\sim 20 – 50 nm in diameter) and larger elongated grains (\sim 100 – 200 nm

in length). The density of the film (t \sim 25 nm on Si substrate) measured by X-ray reflectivity was shown to be 4.47 gcm⁻³, which is close to the density of bulk VO₂ which is 4.57 gcm⁻³.

IV. DISCUSSION

The smooth, continuous film with small and medium polycrystalline grains shown in Fig. 7 is of important requirement in producing VO₂ phase. During the annealing optimization, we found that when the annealing temperature is too high and/or too much oxygen, the grains coalesce, and the films are no longer continuous. It also coincides with the appearance of V₂O₅ phase shown by Raman spectroscopy. On the other hand, when the temperature is too low, the film is continuous but without granular structures seen under SEM, and there is very weak or no modulation of MIT. The annealing conditions optimized for both quartz tube furnace and the Fiji ALD system are reproducible, and they consistently produce smooth VO₂ films with MIT modulation ~ 10^2 (for t ~ 30 nm). As shown in Fig. 4, the typical resistivity at room temperature (below MIT) is ~ 1 Ω cm, and it reduces to ~ 0.01 Ω cm at higher temperature (above MIT).

For the annealing in quartz tube furnace, when the samples are slightly shifted towards higher temperatures from the optimized process window shown in Fig. 1(b), the annealed VO₂ films consist of round grains of various sizes (~ 50 to 250 nm in diameter) and very large, elongated grains (~ 400 to 1200 nm in length), as shown in Fig. 8(a). The MIT characteristics of such films with large grains (LG), along with the typical films with small grains (SG), are shown in Fig. 8(b). The resistivity below MIT is similar for both LG and SG films ($\sim 1~\Omega$ cm). However, above MIT, the resistivity of LG film is much lower ($\sim 0.001~\Omega$ cm) than that of SG film ($\sim 0.01~\Omega$ cm). The LG film thus exhibits MIT modulation $\sim 10^3$ (for t $\sim 30~\text{nm}$). This is the highest that has been reported for the VO₂ films deposited on SiO₂ substrates, and on a par with films deposited on GaN substrates.²²

To access the stability of VO2 films, the resistivity modulation was measured again after 60 days in air. As shown in Fig. 8(b) the whole resistivity profile is shifted up, and the room temperature resistivity is now $\sim 100~\Omega cm$. The resistivity of films with V_2O_5 phase is generally higher than that of films with VO_2 phase, and given that V^{5+} states are normally observed at the surface of VO_2 film, therefore, it seems that oxidation to higher valence state is happening in the LG film. However, it is still predominantly VO_2 , as evidenced by the MIT transition and the Raman shift spectra shown in Fig. 8(b) and (c).

On the other hand, the SG film shows stability even after 120 days in air. The difference of the air stability between the LG and SG films could be attributed to the morphology of these films. It clearly shows that a porous, rugged surface of the LG film is more susceptible to the ambient (moisture, oxidation, etc.) than a dense, continuous surface of the SG film, therefore an encapsulation layer would be required. Despite its smaller MIT modulation than LG films, the SG films are more attractive for incorporation in oscillator circuits due to their scalability and compatibility to CMOS platform. ^{25, 26}

V. CONCLUSIONS

Optimizations of the ALD VO₂ thin films deposited from TEMAV and water precursors were carried out, using different annealing systems for substrates of various sizes up to 8-inch diameter wafer. The physical, optical and electrical characteristics of the films were reported. Moreover, the VO₂ thin films with very large, elongated grains showed the highest MIT modulation reported for films deposited on SiO₂ substrates. However, the films with smaller grains and modest MIT modulation have a smooth continuous surface, and they are promising for integration in devices and circuits. Moreover, these films are found to be very stable in air.

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REFERENCES

- 1. F. J. Morin, Phys. Rev. Lett. 3, 34 (1959).
- 2. H. W. Verleur, A. S. Barker and C. N. Berglund, Phys. Rev. 172, 788 (1968).
- 3. I. Hamberg and C. G. Granqvist, J. Appl. Phys. **60**, 123 (1986).
- 4. B. Hu, Y. Ding, W. Chen, D. Kulkarni, Y. Shen, V. V. Tsukruk and Z. L. Wang, Adv. Mater. 22, 5 (2010).
- 5. G. Stefanovich, A. Pergament and D. Stefanovich, J. Phys.: Condens. Matter 12, 8837 (2000).
- 6. W. A. Vitale, C. F. Moldovan, M. Tamagnone, A. Paone, A. Schuler and A. M. Ionescu, IEEE Electron Device Lett. **36**, 972 (2015).
- 7. W. A. Vitale, L. Petit, C. F. Moldovan, M. Fernández-Bolaños, A. Paone, A. Schüler and A. M. Ionescu, Sens. Actuator A: Physcal **241**, 9 (2016).
- 8. M. Maaza, D. Hamidi, A. Simo, T. Kerdja, A. K. Chaudhary and J. B. K. Kana, Opt. Commun. 285, 4 (2012).
- 9. T. Driscoll, H.-T. Kim, B.-G. Chae, M. Di Ventra and D. N. Basov, Appl. Phys. Lett. 95, 043503 (2009).
- 10. Y. Zhou and S. Ramanathan, Proceedings of the IEEE 103, 1289 (2015).
- 11. C.-Y. Kim, S. H. Kim, S. J. Kim and K.-S. An, Appl. Surf. Sci. **313**, 368 (2014).
- 12. Y. Shigesato, M. Enomoto and H. Odaka, Jpn. J. Appl. Phys. 39, 6016 (2000).
- 13. H.-N. Cui, V. Teixeira, L.-J. Meng, R. Wang, J.-Y. Gao and E. Fortunato, Thin Solid Films 516, 1484 (2008).
- 14. V. P. Prasadam, N. Bahlawane, F. Mattelaer, G. Rampelberg, C. Detavernier, L. Fang, Y. Jiang, K. Martens, I. P. Parkin and I. Papakonstantinou, Mater. Today Chem. **12**, 396 (2019).
- 15. T. Blanquart, J. Niinistö, M. Gavagnin, V. Longo, M. Heikkilä, E. Puukilainen, V. R. Pallem, C. Dussarrat, M. Ritala and M. Leskelä, RSC Adv. 3, 1179 (2013).
- 16. P. A. Premkumar, M. Toeller, I. P. Radu, C. Adelmann, M. Schaekers, J. Meersschaut, T. Conard and S. V. Elshocht, ECS J. Solid State Sci. Technol. 1, P169 (2012).
- 17. G. Rampelberg, D. Deduytsche, B. De Schutter, P. A. Premkumar, M. Toeller, M. Schaekers, K. Martens, I. Radu and C. Detavernier, Thin Solid Films **550**, 59 (2014).
- 18. K. Zhang, M. Tangirala, D. Nminibapiel, W. Cao, V. Pallem, C. Dussarrat and H. Baumgart, ECS Transactions 50, 175 (2013).
- 19. A. P. Peter, K. Martens, G. Rampelberg, M. Toeller, J. M. Ablett, J. Meersschaut, D. Cuypers, A. Franquet, C. Detavernier, J.-P. Rueff, M. Schaekers, S. Van Elshocht, M. Jurczak, C. Adelmann and I. P. Radu, Adv. Funct. Mater. 25, 679 (2015).
- 20. A. C. Kozen, H. Joress, M. Currie, V. R. Anderson, C. R. Eddy and V. D. Wheeler, J. Phys. Chem. C **121**, 19341 (2017).
- 21. X. Lv, Y. Cao, L. Yan, Y. Li and L. Song, Appl. Surf. Sci. 396, 214 (2017).
- 22. B. P. Downey, V. D. Wheeler and D. J. Meyer, Appl. Phys. Express 10, 061101 (2017).
- 23. P. Dagur, A. U. Mane and S. A. Shivashankar, J. Crystal Growth 275, e1223 (2005).
- 24. K. M. Niang, G. Bai and J. Robertson, J. Vac. Sci. Technol. A 38, 042401 (2020).
- 25. E. Corti, B. Gotsmann, K. E. Moselund, A. M. Ionescu, J. Robertson and S. Karg, Solid State Electron. 168, 107729 (2020).
- 26. E. Corti, A. Khanna, K. Niang, J. Robertson, K. E. Moselund, B. Gotsmann, S. Datta and S. Karg, IEEE Electron Device Lett. **41**, 629 (2020).
- 27. Z. Wang, J. Kang, G. Bai, G. Zhong, B. Wang, Z. Yu, Y. Ling, Q. Chen, L. Bao, L. Wu, Y. Cai, J. Robertson and R. Huang, IEEE Electron Device Lett. 41, 1009 (2020).
- 28. J. Janicskó-Csáthy, A. Nagy, L. Nánai, W. Marine, S. P. Souto, M. Balkanski and T. F. George, J. Mater. Res. 17, 1096 (2011).
- 29. M. Pan, J. Liu, H. Zhong, S. Wang, Z.-f. Li, X. Chen and W. Lu, J. Crystal Growth 268, 178 (2004).
- 30. P. Schilbe, Phys. B **316**, 600 (2002).
- 31. Z. Chen, Y. Gao, L. Kang, J. Du, Z. Zhang, H. Luo, H. Miao and G. Tan, Sol. Energy Mater Sol. Cells 95, 8 (2011).
- 32. J. Tauc, R. Grigorovici and A. Vancu, Phys. Status Solidi 15, 627 (1966).
- 33. T. J. Huffman, C. Hendriks, E. J. Walter, J. Yoon, H. Ju, R. Smith, G. L. Carr, H. Krakauer and M. M. Qazilbash, Phys. Rev. B **95**, 075125 (2017).
- 34. H. Lu and J. Robertson, Phys. Status Solidi B **256**, 1900210 (2019).
- 35. H. Lu, Y. Guo and J. Robertson, Phys. Rev. Mater. 3, 094603 (2019).
- 36. A. Srivastava, H. Rotella, S. Saha, B. Pal, G. Kalon, S. Mathew, M. Motapothula, M. Dykas, P. Yang, E. Okunishi, D.D. Sarma, and T. Venkatesan, APL Materials 3, 026101 (2015).

Table 1. Comparison of Raman peaks of the polycrystalline VO₂ and V₂O₅ films between this work and other works.

		VO2		V	V ₂ O ₅
This work		Pan [29]	Schilbe [30]	This work	Janicsko- Csathy [28]
Tube	Fiji				
furnace	system				
141	138		149	103	104
195	194	194	199	145	149
224	222	225	225	197	197
261	257	258	259	284	285
		265	265	304	305
309	304	308	313	404	405
339	338	339	339	481	482
388	391	392	392	527	528
399		395	395		628
439	438	444	444	651	652
		453	453	701	702
		489	489		804
497		503	503	995	995
583		585	595	1020	1020
616	616	618	618		
667		650	670		
824	823	825	830		

FIG. 1. (a) Raman spectra of VO₂ thin films annealed under various annealing conditions in a quartz tube furnace, and (b) the temperature distribution measured across furnace for different setpoints. The red squares are where the VO₂ samples were placed during annealing.

FIG. 2. (a) The positions of the VO₂ samples placed across the 8-inch diameter substrate holder of a Fiji S200 ALD system for annealing at 450 °C for 2 hours in Ar/O₂ (100/100 sccm), and (b) Raman spectra of the VO₂ thin films located at various positions shown in (a). The peak intensity for the Raman shift at 616 cm⁻¹ are also shown in (a).

FIG. 3. Comparison of the XRD profiles of the VO₂ thin films annealed in either quartz tube furnce or Fiji ALD chamber. Peaks marked with * seem to be due to measurement artefact or substrate related.

FIG. 4. Resistivity, ρ , as a function of temperature and its derivative (d(log ρ)/dT) for the annealed VO₂ (t ~ 30 nm) thin films annealed in (a) quartz tube furnce and (b) Fiji ALD chamber. The transition temperatures during heating cycle (T_h) and cooling cycle (T_c) are also shown.

FIG. 5. (a) Transmittance (inset) and Tauc plot of absorption for a VO₂ film (t \sim 30 nm). The smaller gap (0.55 eV) can be assigned due to the *d-d* transitions, and the larger gap (1.3 eV) is due to O 2*p* to V 3*d* transitions. (b) Absorption coefficient as a function of energy for the same film.

FIG. 6. Thickness distribution for VO₂ films after (a) 450 ALD cycles, and (b) 890 ALD cycles. To achieve good uniformity, the wafers were rotated by 180° halfway through the deposition. The gas inlet and outlet of the cross flow reactor are shown in the map.

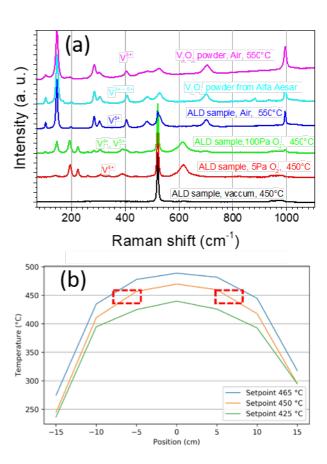
FIG. 7. AFM micrographs of the VO₂ film (t \sim 30 nm) (a) before and (b) after annealing, the rms of roughness are also shown. (c) SEM micrograph of the polycrystalline VO₂ film of same thickness. The scale bar in each image is 500 nm.

FIG. 8. (a) SEM micrograph of the VO₂ thin film (t \sim 30nm) with large grains (LG). The scale bar in SEM image is 500 nm. (b) Resistivity vs. temperature plots of the VO₂ films with LG and small grains (SG), measured as fresh sample and after 60 or 120 days. (c) Raman shift spectra of the LG film measured after 60 days.

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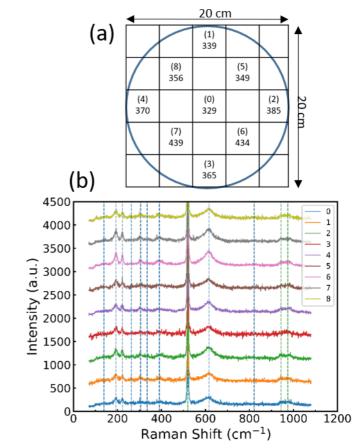


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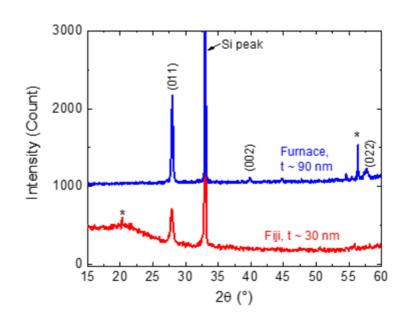




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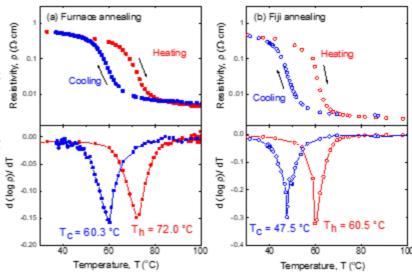


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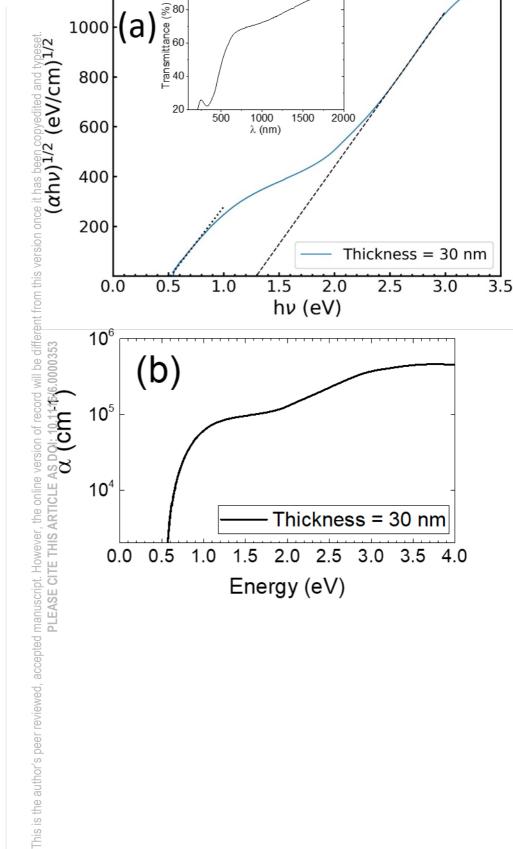






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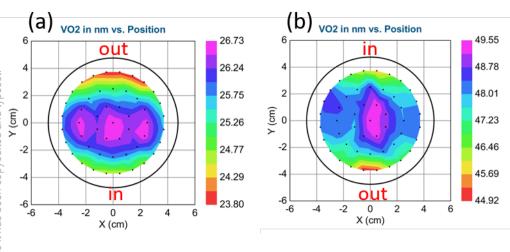
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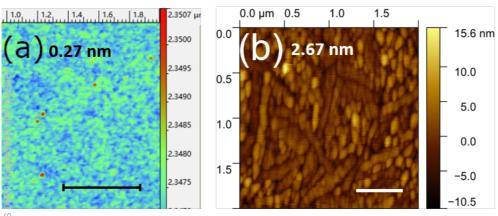


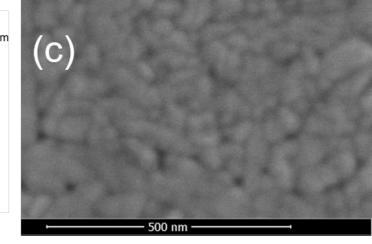


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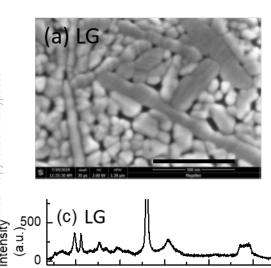
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