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## Spectroscopic and Time-Dependent DFT Study of the Photophysical Properties of Substituted 1,4-Distyrylbenzenes

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7 **Spectroscopic and time-dependent DFT study of**  
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11 **the photophysical properties of substituted 1,4-**  
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**ABSTRACT:** In this contribution we examine the photophysical properties of fifteen totally *trans-trans* 1,4-distyrylbenzene derivatives (DSBs) functionalized with different electron-donating (ED) and electron-withdrawing (EW) groups by experimental and computational methodologies. We use UV-Vis and fluorescence spectroscopy to determine the experimental optical properties such as the maximum absorption ( $\lambda_{abs}^{exp}$ ) and emission ( $\lambda_{em}^{exp}$ ) wavelengths, the HOMO-LUMO energy gaps ( $\Delta E_{abs}^{exp}$ ), the molar extinction coefficients ( $\epsilon$ ), the fluorescence quantum yields ( $\Phi_f$ ) and the fluorescence lifetimes ( $\tau$ ). We also calculate the experimental oscillator strengths ( $F_{osc}^{exp}$ ) and spontaneous emission decay rate ( $k_r^{exp}$ ) and correlate all these magnitudes to the corresponding calculated properties – maximum absorption ( $\lambda_{abs}^{cal}$ ) and emission ( $\lambda_{em}^{cal}$ ) wavelengths, vertical transition energies ( $\Delta E_{abs}^{cal}$ ), oscillator strength ( $F_{osc}^{cal}$ ) and spontaneous emission decay rate ( $k_r^{cal}$ ) – obtained by the TDDFT method. We analyze the effect of the electronic nature of the substituents on the properties of the DSBs, finding that the ED and EW groups lead to bathochromic shifts. This is consistent with the decrease of  $\Delta E$  values as the strength of ED and EW substituents increases. We find excellent correlations between calculated and experimental values for  $\lambda_{abs}$ ,  $\lambda_{em}$  and  $\Delta E_{abs}$  ( $r \sim 0.99-0.95$ ). Additionally,  $F_{osc}$  and  $k_r$  values are in good agreement ( $r \sim 0.87-0.72$ ) with the experimental properties. Overall, we find that for substituted 1,4-DSBs, computational chemistry is an excellent tool to predict structure-property relationships, which can be useful to forecast the properties of their polymeric analogs, which are usually difficult to determine experimentally.

## 1. Introduction

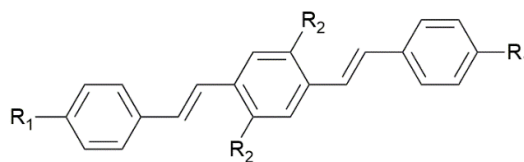
Poly(phenylenevinylene) (PPV) and its derivatives have attracted considerable attention because of their remarkable optical and electronic properties. It is well known that after the discovery of

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3 the optoelectronic properties of PPVs, especially the electroluminescence<sup>1-2</sup>, these materials have  
4 been exploited as an active layer in organic light-emitting diodes (OLEDs)<sup>3-13</sup> and other  
5 applications such as laser dyes<sup>14-16</sup>, photovoltaic cells<sup>17-23</sup>, chemosensors<sup>24-29</sup>, linkers for metal-  
6 organic frameworks (MOFs)<sup>30-32</sup>, dendrimers<sup>33-36</sup>, photoredox catalysts and photosensitizers<sup>37-39</sup>.  
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8 The synthesis of PPVs frequently presents challenges as it leads to a mixed *cis/trans* configuration  
9 over the vinyl segments of the polymeric structure; they show large dimensions and poor solubility  
10 in a wide range of solvents.<sup>40</sup>  
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20 Oligo(phenylenevinylene)s (OPVs) have drawn interest as structural models for PPVs, as their  
21 synthesis yields well-defined chemical structures that allow for easier optoelectronic  
22 characterizations. OPVs share many properties with their structural analog PPVs, especially in the  
23 case of segmented polymers. This is useful to determine which structural features of a PPV, or a  
24 segmented polymer, are the most appropriate to incorporate for a given application. This is known  
25 as “the oligomer approach”<sup>41</sup>. The OPVs smaller size allows computational studies of their  
26 optoelectronic properties with state-of-the-art density functional theory (DFT) methods that  
27 complement the experimental information.  
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39 Perhaps the most studied OPVs are distyrylbenzene (DSBs) derivatives. The optoelectronic  
40 properties of numerous substituted DSBs have been investigated by both experimental<sup>42-49</sup> and  
41 computational methods<sup>50-55</sup>. With the goal of understanding the structure-property relationships of  
42 these technological attractive materials, we reported the synthesis of a full *trans-trans* DSBs series  
43 with different electron-donor (ED) and electron-withdrawing (EW) substituents by Mizoroki-Heck  
44 cross-coupling reactions<sup>56</sup>. The interest in the *trans*-configuration of the vinyl bonds reflects that  
45 a more-planar conformation is obtained, which facilitates conjugation and improves their optical  
46 and electronic properties<sup>57-58</sup>.  
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The study of the structural effects on the photophysical properties of DSBs provides insights into the structural features that a phenylenevinylene derivative (oligomer or polymer) must have to be useful. Thus, in order to understand better the structure-property relationships of this type of materials, here we correlate the experimental optical characterization of DSBs – substituted with different ED or EW groups (Figure 1) – with their optical properties calculated by computational methods based on the time-dependent DFT (TDDFT).



Series 1			Series 2		
DSB	R <sub>1</sub>	R <sub>2</sub>	DSB	R <sub>1</sub>	R <sub>2</sub>
<b>1-H-DSB</b>	-H	-H	<b>2-H-DSB</b>	-H	-OCH <sub>3</sub>
<b>1-OH-DSB</b>	-OH	-H	<b>2-OH-DSB</b>	-OH	-OCH <sub>3</sub>
<b>1-OMe-DSB</b>	-OCH <sub>3</sub>	-H	<b>2-OMe-DSB</b>	-OCH <sub>3</sub>	-OCH <sub>3</sub>
<b>1-OCOMe-DSB</b>	-OCOCH <sub>3</sub>	-H	<b>2-OCOMe-DSB</b>	-OCOCH <sub>3</sub>	-OCH <sub>3</sub>
<b>1-CN-DSB</b>	-CN	-H	<b>2-CN-DSB</b>	-CN	-OCH <sub>3</sub>
<b>1-NO<sub>2</sub>-DSB</b>	-NO <sub>2</sub>	-H	<b>2-NO<sub>2</sub>-DSB</b>	-NO <sub>2</sub>	-OCH <sub>3</sub>
<b>1-COOMe-DSB</b>	-COOCH <sub>3</sub>	-H	<b>2-Me-DSB</b>	-CH <sub>3</sub>	-OCH <sub>3</sub>
<b>1-COMe-DSB</b>	-COCH <sub>3</sub>	-H			

**Figure 1.** Chemical structure of the studied DSBs.

## 2. Experimental section

### 2.1 Materials

The *trans-trans* 1,4-distyrylbenzene derivatives (DSBs) were previously synthesized by the Mizoroki-Heck cross-coupling reaction and the synthetic protocols are described in the literature<sup>56</sup>. The solvent used for the optical characterization (CHCl<sub>3</sub>) was purchased from Fisher Scientific and used without further purification.

### 2.2 Experimental procedures

DSBs compounds were dissolved in CHCl<sub>3</sub> to prepare 10 μM solutions. In order to obtain the

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3 desired concentrations for compounds with poor solubility (e.g., Series 1) ultrasonic and heating  
4 techniques were applied. UV-vis spectra were measured on an Agilent Cary 60 spectrophotometer.  
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6 The experimental HOMO-LUMO energy gaps ( $\Delta E_{abs}^{exp}$ ) were determined from the analysis of the  
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8 absorption edge according to equation 1:  
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$$\alpha \cdot h\nu = A \cdot (h\nu - \Delta E)^{1/2} \quad (1)$$

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13 where  $\alpha$  is the absorption coefficient,  $h$  is the Planck constant,  $\nu$  is the frequency of light and  $A$  is  
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15 the absorbance. By plotting  $(h\nu A)^2$  vs.  $h\nu$ , and extending the linearity edge of absorbance to the  
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17 intersect with the energy axis the  $\Delta E$  value can be obtained<sup>59-60</sup>. The molar extinction coefficients  
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19 ( $\epsilon$ ) were obtained from the slopes in the absorbance vs. concentration plots. To calculate the  
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21 experimental oscillator strengths ( $F_{osc}^{exp}$ ) we performed least squares fittings of the absorption  
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23 spectra expressed in wavenumbers to linear combinations of Gaussian-type functions. The  $F_{osc}^{exp}$   
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25 are related to the area under the absorption spectra and for Gaussian shaped bands are calculated  
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27 as<sup>61</sup>:  
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$$F_{osc}^{exp} = 4.67 \cdot 10^{-9} \frac{\text{mol cm}^2}{L} \epsilon_{max} \Delta\tilde{\nu}_{1/2}, \quad (2)$$

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38 where  $\Delta\tilde{\nu}_{1/2}$  is the band full width at half maximum (FWHM) value in  $\text{cm}^{-1}$ .  
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44 Fluorescence (FL) spectra were acquired on a PTI QuantaMaster<sup>TM</sup> 40 spectrofluorometer in  
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46  $\text{CHCl}_3$  solutions. The fluorescence quantum yields ( $\Phi_f$ ) were determined at  $\lambda_{exc}$  corresponding to  
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48 the absorption maxima. Quinine sulfate in 0.1 M  $\text{H}_2\text{SO}_4$  ( $\Phi_f = 0.54$ ) was used as the standard for  
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50 determination of fluorescence quantum yields ( $\Phi_f$ ) at  $\lambda_{exc}$  corresponding to the absorption  
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52 maximum of series 1. Perylene in ethanol ( $\Phi_f = 0.92$ ) was used as the standard for determination  
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54 of the  $\Phi_f$  at  $\lambda_{exc}$  corresponding to the second absorption maximum of series 2 ( $\lambda_2$ ). The  $\Phi_f$  of 1-  
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**NO<sub>2</sub>-DSB** was calculated also using perylene as standard. The  $\Phi_f$  values were calculated according to a method described in the literature<sup>62</sup>, using equation 3:

$$\Phi_{f(\text{DSB})} = \Phi_{f(\text{ST})} \frac{f_{(\text{ST})}}{f_{(\text{DSB})}} \frac{F_{(\text{DSB})}}{F_{(\text{ST})}} \left( \frac{\eta_{(\text{DSB})}}{\eta_{(\text{ST})}} \right)^2 \quad (3)$$

where  $\Phi_{f(\text{DSB})}$  and  $\Phi_{f(\text{ST})}$  are the fluorescence quantum yield of the sample and the standard, respectively.  $f_{(\text{DSB})}$  and  $f_{(\text{ST})}$  are the absorption factors, that is the fraction of the light impinging on the sample that is absorbed:  $f_{(\text{DSB})} = 1 - 10^{-A(\text{DSB})}$  and  $f_{(\text{ST})} = 1 - 10^{-A(\text{ST})}$ .  $F_{(\text{DSB})}$  y  $F_{(\text{ST})}$  are the integrated intensities (areas under the emission curve) of the sample and standard spectra, respectively.  $\eta_{(\text{DSB})}$  y  $\eta_{(\text{ST})}$  are the refractive indices of the solvent used in the sample and reference solutions, respectively. The fluorescence lifetimes ( $\tau$ ) were measured in an Easy-Life (PTI) system with samples with a maximum absorbance of 0.1 (at the excitation wavelength) in CHCl<sub>3</sub> under nitrogen atmosphere. Since it is possible to relate the radiative and non-radiative constants ( $k_r$  and  $k_{nr}$ ) with the  $\Phi_f$  and  $\tau$  values<sup>63</sup>,

$$\Phi_f = \frac{k_r}{k_r + k_{nr}} \quad (4)$$

$$\tau = \frac{1}{k_r + k_{nr}} \quad (5)$$

the experimental radiative constants ( $k_r^{exp}$ ) were calculated from equation 6:

$$k_r^{exp} = \frac{\Phi_f}{\tau} \quad (6)$$

### 3. Computational calculations

The ground and first excited state geometries of the 15 DSBs were optimized the CAM-B3LYP functional<sup>64</sup> and the def2-TZVPP<sup>65-66</sup> basis set considering CHCl<sub>3</sub> as the solvent with the conductor-like polarizable continuum model<sup>67</sup>. For all molecules, the initial geometry was built

with the  $C_{2h}$  symmetry point group. All calculations were performed with the ORCA computational package<sup>68</sup> with the resolution of identity and chain of spheres approximations<sup>69</sup>. All TDDFT calculations employed the Tamm-Dancoff approximation<sup>70</sup>.

Eight vertical transition energies ( $\Delta E_{abs}^{cal}$  and  $\Delta E_{em}^{cal}$ ) and oscillator strengths ( $\Delta F_{osc}^{cal}$ ) were obtained from single-point TDDFT calculations at the ground and excited state equilibrium geometries. The maximum absorption ( $\lambda_{abs}^{cal}$ ) and emission ( $\lambda_{em}^{cal}$ ) wavelengths were calculated from the  $\Delta E_{abs}^{cal}$  and  $\Delta E_{em}^{cal}$  values as  $\lambda = hc/\Delta E$ , where  $c$  is the speed of light.

The calculated radiative constants ( $k_r^{cal}$ ) from excited to ground state were calculated according to:

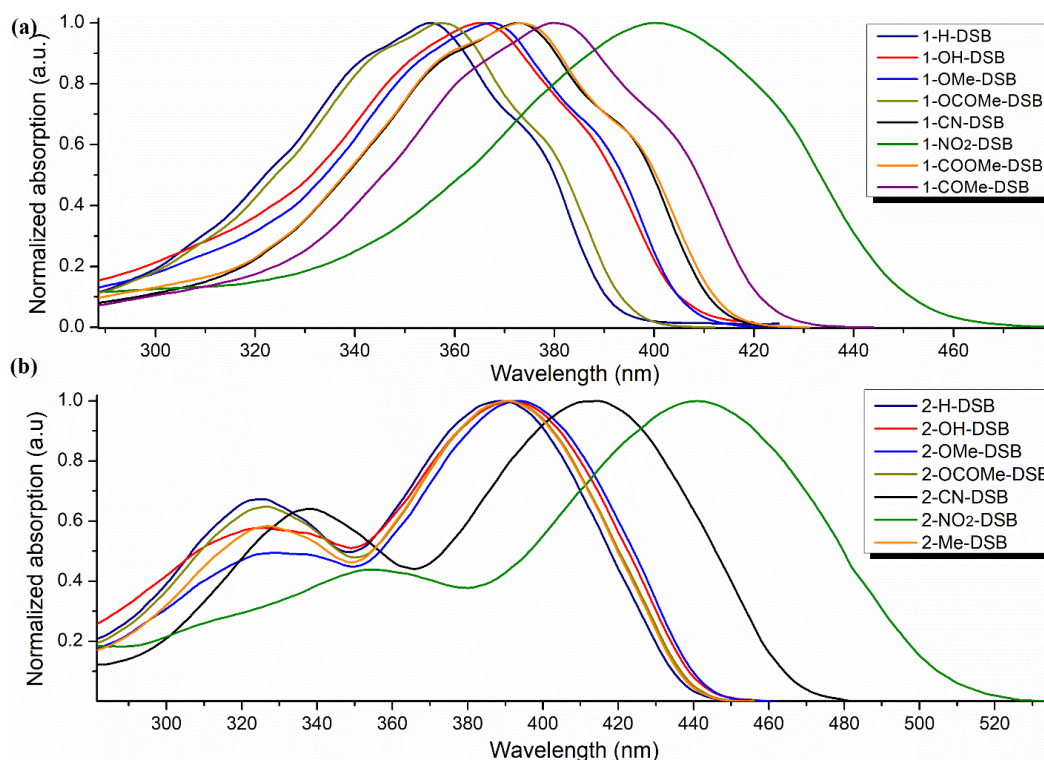
$$k_r^{cal} = \frac{4}{3} \frac{(\Delta E_{em}^{cal})^3}{c^3} \mu_{10}^2 \quad (7)$$

where  $\mu_{10}$  is the transition dipole strength evaluated at the excited state geometry<sup>71</sup>.

#### 4. Results and discussion

The 1,4-distyrylbenzenes synthesized by Mizoroki-Heck cross-coupling reaction were obtained as total *trans*-configuration systems<sup>56</sup>. The UV-vis spectroscopic characterization of these compounds is shown in Figure 2. The absorption spectra of the DSBs of series 1 (Figure 2a) present a maximum absorption band due to  $\pi \rightarrow \pi^*$  electronic transitions, while the absorption spectra of the DSBs of series 2 (Figure 2b), which have methoxy groups in the central ring, exhibit two maximum absorption bands, one of higher energy and lower intensity assigned to the  $n \rightarrow \pi^*$  transition, and one of lower energy and higher intensity due to the  $\pi \rightarrow \pi^*$  transitions of the DSB systems.





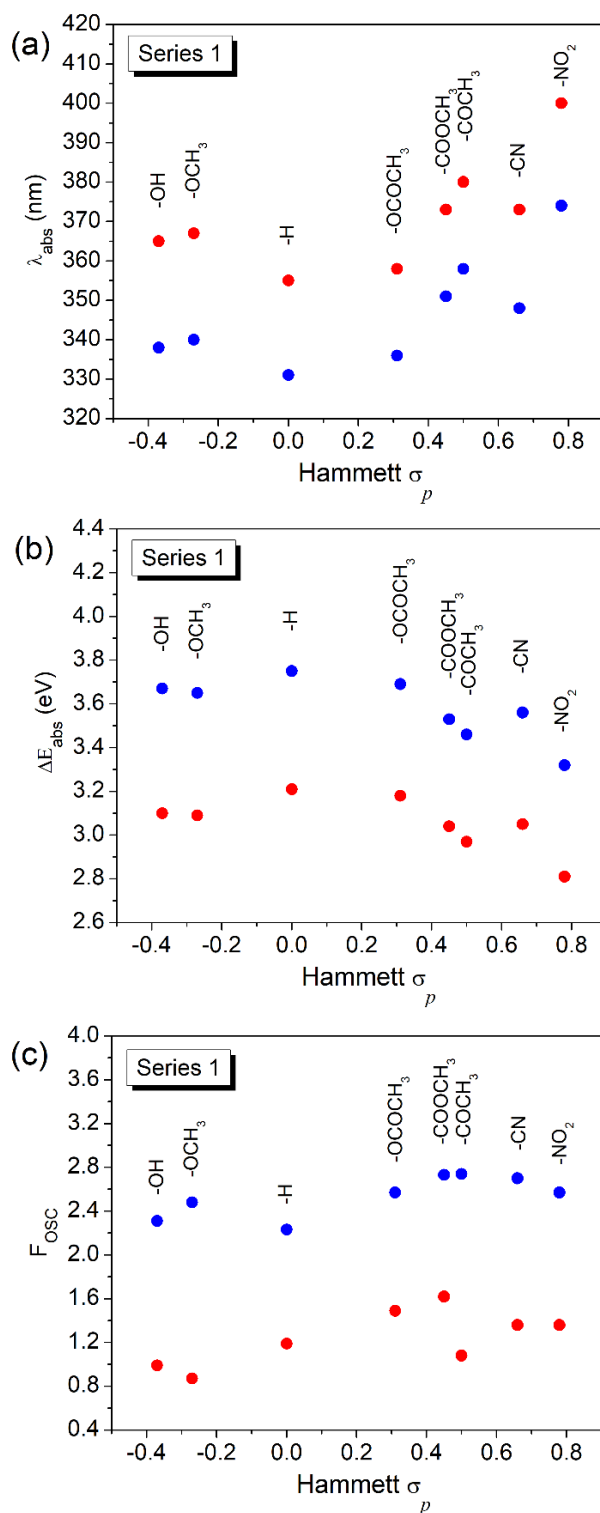
**Figure 2.** Normalized absorption spectra of DSBs of (a) series 1 and (b) series 2 measured in  $\text{CHCl}_3$ .

The calculated electronic transitions are presented in Tables S6 and S7 in the supporting information. Series 1 molecules only show strong  $S_{0 \rightarrow 1}$  transitions ( $\lambda_{abs}^{cal} = 331$  to  $374$  nm), whereas Series 1 molecules show strong  $S_{0 \rightarrow 1}$  transitions, ( $\lambda_{abs}^{cal} = 364$  to  $415$  nm) and less intense  $S_{0 \rightarrow 2}$  transitions ( $\lambda_{abs}^{cal} = 293$  to  $322$  nm), in good agreement with the experimental observations. The orbital pair contributions presented in Table S8 show that the most important contribution to  $S_{0 \rightarrow 1}$  transition comes from the HOMO-LUMO excitations, as the TDDFT combination coefficient of these orbitals is larger than 0.9 in all cases.

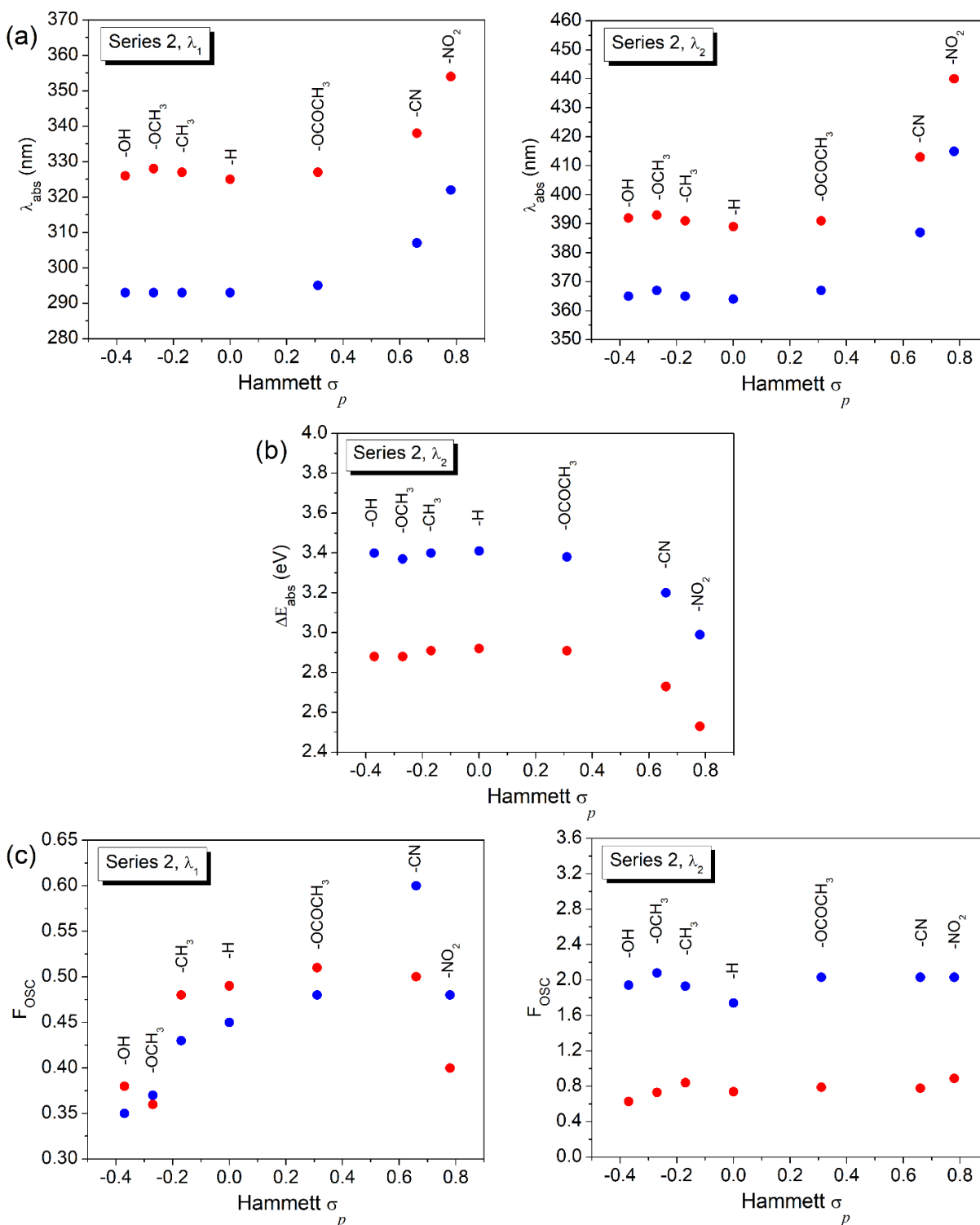
Figure 3 shows a comparison of the experimental absorption properties for DSBs of series 1 in  $\text{CHCl}_3$  obtained by UV-vis spectroscopy with the calculated by TDDFT. To facilitate understanding, the order of substituents in the figures was established according to the reported Hammett substituent constants ( $\sigma_p$ )<sup>72</sup>, which show how electron-donating or electron-withdrawing

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3 a group is relative to a neutral group (H) in the *para* position of the aromatic terminal rings of the  
4 DSBs. The experimental and calculated data are shown in the SI file (Table S1).  
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8 Figures 3a, 3b and 3c show the comparison between  $\lambda_{abs}^{exp}$  and  $\lambda_{abs}^{cal}$ ,  $\Delta E_{abs}^{exp}$  (calculated by eq. 1)  
9 and  $\Delta E_{abs}^{cal}$ , and  $F_{osc}^{exp}$  (calculated by eq. 2) and  $F_{osc}^{calc}$ , respectively, all as a function of the  
10 corresponding  $\sigma_p$  values. This figure shows that the computational studies describe very well the  
11 optical properties of these molecules. Likewise, Figure 4 shows the results obtained for series 2  
12 DSBs. Figure 4a shows the comparison between  $\lambda_{abs}^{exp}$  and  $\lambda_{abs}^{calc}$  determined at the two maximum  
13 absorption bands ( $\lambda_1$  and  $\lambda_2$ ), figure 4b compares the  $\Delta E_{abs}^{exp}$  and  $\Delta E_{abs}^{calc}$  values for the second  
14 absorption maximum ( $\lambda_2$ ), and figure 4c compares  $F_{osc}^{exp}$  (calculated by eq. 2) and  $F_{osc}^{calc}$  determined  
15 at  $\lambda_1$  and  $\lambda_2$ .  
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**Figure 3.** Experimental data (red) and computational calculations (blue) for the optical properties of DSBs-series 1 in  $\text{CHCl}_3$  as a function of  $\sigma_p$ : (a)  $\lambda_{\text{abs}}$  (b)  $\Delta E_{\text{abs}}$  (c)  $F_{\text{osc}}$ .



**Figure 4.** Experimental data (red) and computational calculations (blue) for the absorption properties of DSBs-series 2 in CHCl<sub>3</sub> as a function of  $\sigma_p$ : (a)  $\lambda_{\text{abs}}$  (b)  $\Delta E_{\text{abs}}$  (c)  $F_{\text{osc}}$ .

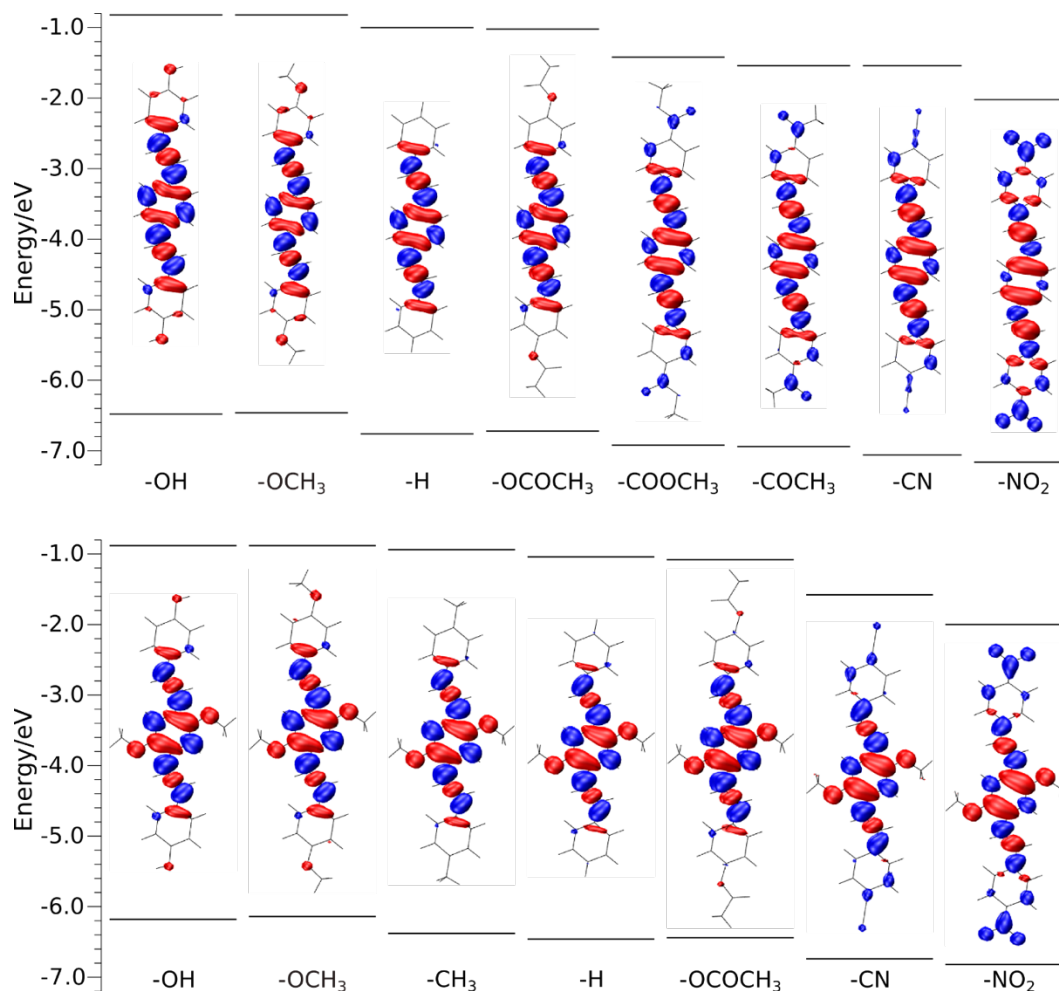
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3 Results in Figures 3 and 4, along with the spectra shown in Figure 2, show that in  $\text{CHCl}_3$  both  
4 ED and EW groups at the two ends of the DSB structure lead to bathochromic shifts of the  
5 absorption band compared to the unsubstituted DSBs (**1-H-DSB** and **2-H-DSB**). In particular, the  
6 ED groups studied have a moderate effect on the bathochromic shift ( $\Delta\lambda \sim 10$  nm), whereas EW  
7 groups have a greater effect ( $\Delta\lambda > 20$  nm); that is, greater shifts are found when the electron  
8 donating or electron withdrawing strength of the substituent increases. A similar trend has been  
9 reported in other studies of DSBs with ED and/or EW substituents on the terminal rings<sup>44-46</sup>.  
10 Although the absorption band shift caused by the ED and EW substituents seems to be smaller in  
11 DSBs of series 2, compared to series 1, it is important to note that the entire absorption spectrum  
12 for each of the DSBs of series 2 is red shifted, indicating that the presence of two ED substituents  
13 (methoxy groups) over the central ring strongly affect the optical properties of DSBs. Comparing  
14 the reference compounds in each series, the methoxy groups in the central ring of the **2-H-DSB**  
15 produce a large bathochromic shift ( $\sim 34$  nm) with respect to its unsubstituted **1-H-DSB** counterpart  
16 similar to computational predictions ( $\sim 32$  nm). A similar study reported by Chaieb et al., also  
17 showed that the introduction of ED groups on the central ring of the DSB system leads to a red  
18 shift<sup>45</sup>.  
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40 Results in Figures 3 and 4 also show that the EW groups in the DSB system lead to smaller  
41 HOMO-LUMO energy gaps ( $\Delta E$ ), being lowest for **1-NO<sub>2</sub>-DSB** and **2-NO<sub>2</sub>-DSB**. The ED groups  
42 also contribute to decreasing the  $\Delta E$  with respect to unsubstituted DSBs. The trend shows that  $\Delta E$   
43 decreases with the strength of ED and EW substituents, in agreement with the increased  
44 bathochromic shift. A similar trend in the decrease of  $\Delta E$  values with ED groups was found by Wu  
45 and coworkers<sup>44</sup>, who reported optical properties of DSBs substituted with other ED such as  $-\text{CH}_3$ ,  
46  $-\text{C}(\text{CH}_3)_3$ ,  $-\text{O}(\text{CH}_2)_4\text{CH}_3$ ,  $-\text{N}(\text{CH}_3)_2$ . Additionally, we probed that EW groups such as  $-\text{COOCH}_3$ ,  
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3 -COCH<sub>3</sub>, -CN and -NO<sub>2</sub> follow the same trend with an intensified effect on the decrease of  $\Delta E$   
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5 values. On the other hand, the  $\Delta E$  values obtained for DSBs of series 2 are lower than those  
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7 obtained for DSBs of series 1, which indicate that the introduction of methoxy groups in the central  
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9 ring favors a decrease in the energy gap.  
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13 In order to understand how the different substituents affect the optical properties of the DSB, we  
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15 calculate the orbital eigenvalues for the DSB series 1 and series 2 (Figure 5). The results reveal  
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17 that, relative to unsubstituted DSBs, EW groups decrease the energies of both HOMO and LUMO,  
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19 but the effect on the latter is larger. On the other hand, ED groups increase the HOMO and LUMO  
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21 energies, and they have a larger impact on the HOMO. For both EW and ED groups, these changes  
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23 in orbital energies lead to a HOMO-LUMO energy gap shrinkage and therefore to the  
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25 bathochromic shifts observed in the UV-vis spectra.  
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30 Figure 5 also displays the electron density changes that occur upon excitation in the DSBs of  
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32 series 1 and 2. For all molecules, the changes in the central aromatic ring and the connecting vinyl  
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34 units are similar, with alternating density increases and decreases between bonds. In DSBs with  
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36 ED groups, the electron density around the oxygen atom decreases upon excitation, whereas in  
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38 DSBs with EW groups, the density around the -NO<sub>2</sub>, -CN and carbonyl moieties increase upon  
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40 excitation.  
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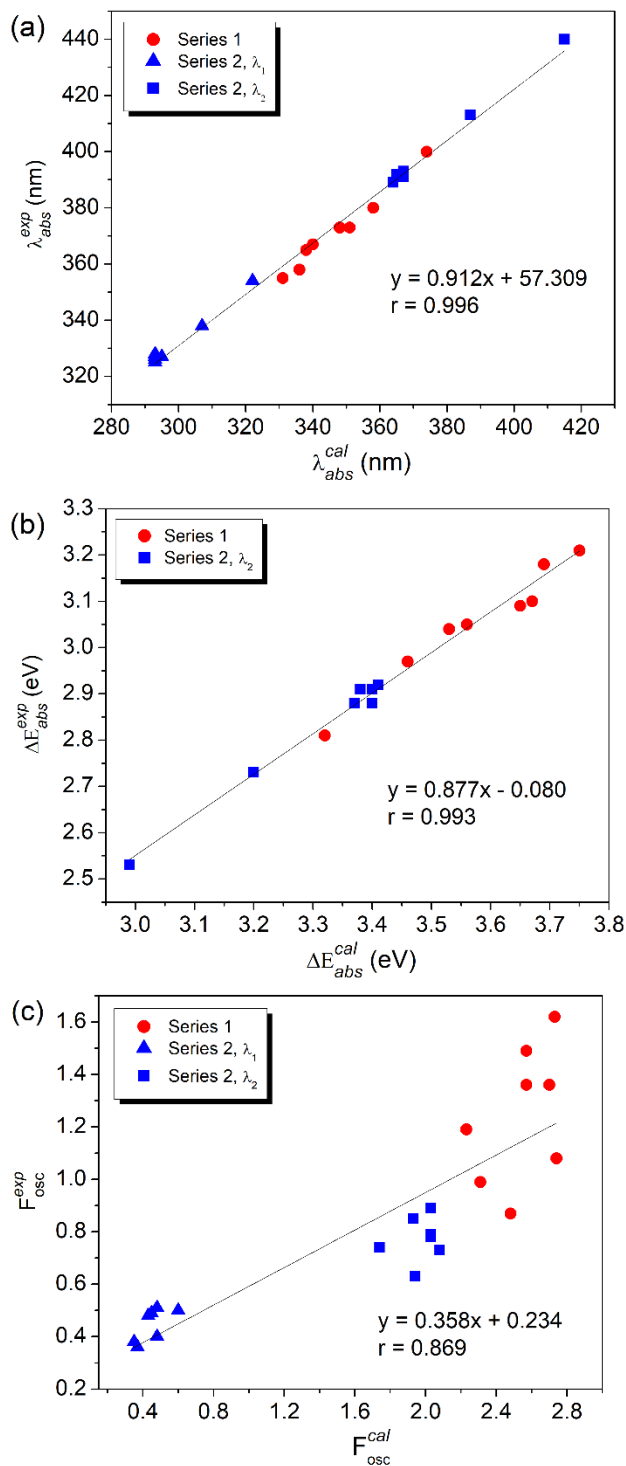
**Figure 5:** Calculated HOMO-LUMO eigenvalues and electron density changes upon excitation ( $S_0 \rightarrow S_1$ ) for the DSBs of series 1 (top) and series 2 (bottom). Blue electron density contours are drawn at 0.001 a.u. and red contours at -0.001 a.u.

Regarding the molar extinction coefficient ( $\epsilon$ ), no clear trends were observed (see Tables S1-S2). DSBs of series 1 showed values of  $\epsilon$  higher than those obtained for DSBs of series 2. Consequently, the experimental oscillator strength ( $F_{osc}^{exp}$ ), obtained from equation 2, shows higher values for DSBs of series 1 – ranging from 0.87 to 1.62 – and lower values for DSBs of series 2 – from 0.36 to 0.51 at  $\lambda_1$  and from 0.63 to 0.89 at  $\lambda_2$ . The same behavior was predicted by computational calculations, where DSBs of series 1 showed values of  $F_{osc}^{cal}$  higher than those obtained for DSBs of series 2. For DSBs of series 1 the  $F_{osc}^{cal}$  ranges from 2.23 to 2.74. In addition,

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3 the  $F_{\text{osc}}^{\text{cal}}$  values for series 2 are higher for the first transition ( $\lambda_2$ ), which is the most intense, ranging  
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5 from 1.74 to 2.08, whereas for the second transition ( $\lambda_1$ ) this quantity ranges from 0.35 to 0.60.  
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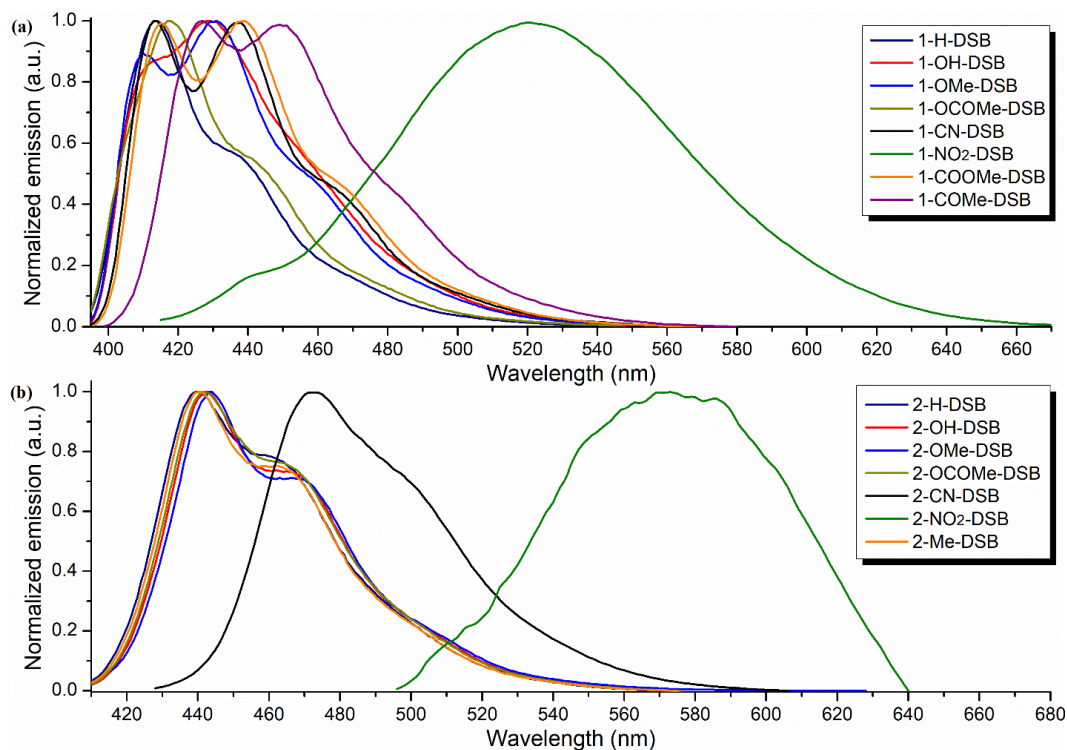
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9 The results shown above demonstrate that the computational calculations can be used to predict  
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11 the optical properties of DSBs. In order to support this, we correlate the results obtained from the  
12  
13 computational and experimental absorption properties studied for the DSBs of series 1 and 2  
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15 (Figure 6). In Figure 6a, which shows the correlation between  $\lambda_{\text{abs}}^{\text{cal}}$  and  $\lambda_{\text{abs}}^{\text{exp}}$  values, the DFT  
16  
17 results predict maximum absorption wavelengths that are about 20 nm smaller than the  
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19 experimental values, but with a high correlation ( $r = 0.996$ ). In figure 6b one can see the correlation  
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21 between  $\Delta E_{\text{abs}}^{\text{cal}}$  and  $\Delta E_{\text{abs}}^{\text{exp}}$  values. Again, DFT predicts values that, on average, are 0.50 eV higher  
22  
23 than the experimental values, with excellent correlation ( $r = 0.993$ ). It was observed that the  
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25 computational method correctly predicts that the energy gaps for DSBs of series 1 are larger than  
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27 those for DSBs of series 2. Finally, figure 6c shows the correlation of the experimental  $F_{\text{osc}}^{\text{exp}}$  values  
28  
29 (calculated by eq. 2) as a function of the calculated  $F_{\text{osc}}^{\text{cal}}$  values. In this case, although the  
30  
31 correlation is good ( $r = 0.87$ ), it is not ideal for quantitative predictions. Nevertheless, the  
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33 computational method correctly clusters the transitions of DSBs of series 1 as the most intense and  
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35 the second transitions of DSBs of series 2 as the least.  
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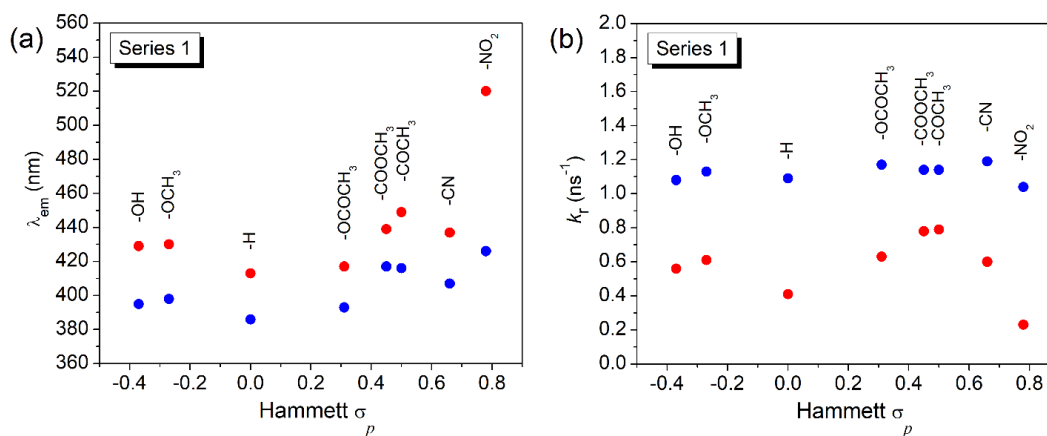
**Figure 6:** Correlation between the experimental and calculated absorption properties (a)  $\lambda_{abs}^{exp}$  and  $\lambda_{abs}^{cal}$ , (b)  $\Delta E_{abs}^{exp}$  and  $\Delta E_{abs}^{cal}$ , and (c)  $F_{osc}^{exp}$  and  $F_{osc}^{cal}$ . Red circles: DSBs of series 1. Blue triangles: DSBs of series 2 ( $\lambda_1$ ). Blue squares: DSBs of series 2 ( $\lambda_2$ ).

Next, we analyze the emission properties of the DSBs using FL spectroscopy and TDDFT calculations. The normalized emission spectra of the DSBs of series 1 and 2 are shown in Figure 7a and 7b, respectively.

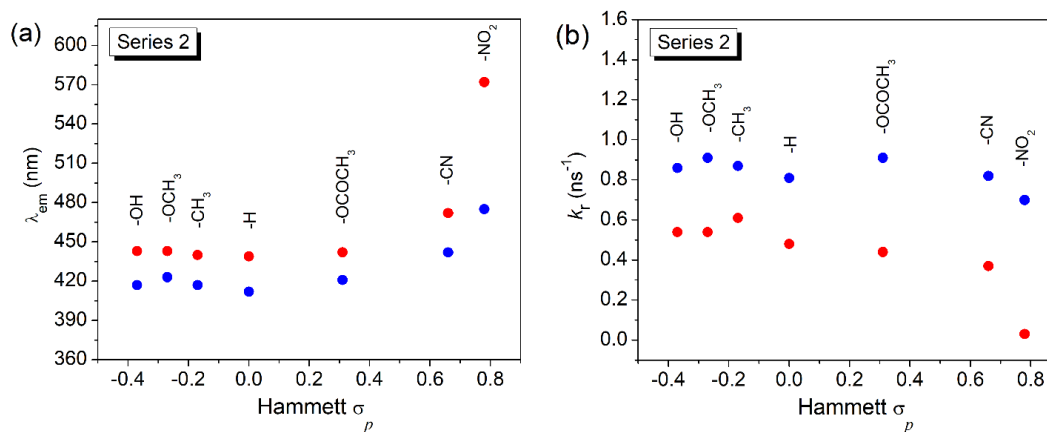


**Figure 7.** Normalized emission spectra of the DSBs of (a) series 1, and (b) series 2, measured in CHCl<sub>3</sub>.

Figure 8 shows a comparison of the experimental emission properties for DBSs of series 1 in CHCl<sub>3</sub> obtained by FL spectroscopy with the calculated by TDDFT. Figure 8a compares the experimental and calculated maximum emission wavelengths ( $\lambda_{em}$ ), and Figure 8b the experimental (determined by eq. 6) and calculated spontaneous emission decay rates ( $k_r$ ). As it is shown in Figures 3 and 4, the order of substituents in the figures was established according to the corresponding  $\sigma_p$  values. Furthermore, Figure 9 shows the comparison between  $\lambda_{em}^{exp}$  and  $\lambda_{em}^{cal}$ , and between  $k_r^{exp}$  and  $k_r^{cal}$  values for DBSs of series 2 in CHCl<sub>3</sub>.



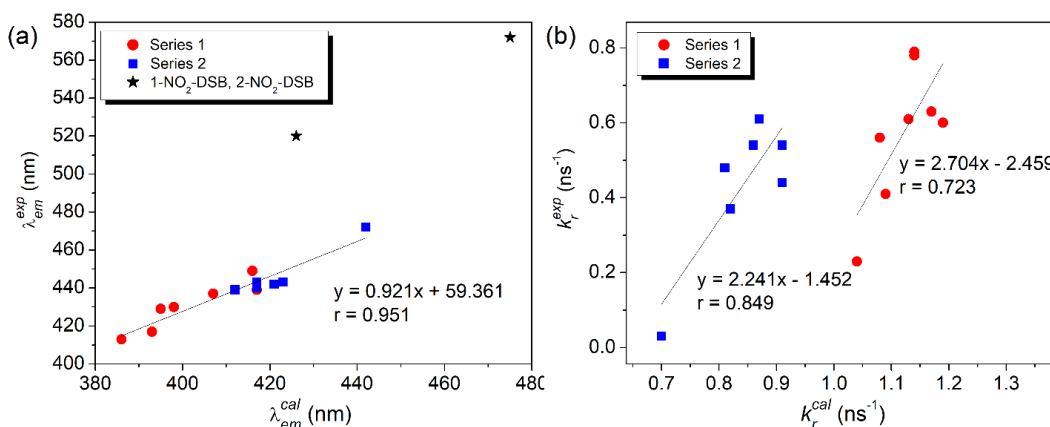
**Figure 8.** Experimental data (red) and computational calculations (blue) for the emission properties of DSBs-series 1 in CHCl<sub>3</sub> as a function of the corresponding  $\sigma_p$ : (a)  $\lambda_{em}$  (b)  $k_r$ .



**Figure 9.** Experimental data (red) and computational calculations (blue) for the emission properties of DSBs-series 2 in CHCl<sub>3</sub> as a function of the corresponding  $\sigma_p$ : (a)  $\lambda_{em}$  (b)  $k_r$ .

Results in Figures 8 and 9, along with the emission spectra in Figures 7, show that in CHCl<sub>3</sub> the ED and EW groups of the substituted DSB structures lead to a bathochromic shift of the emission band, which is greater in the case of EW groups. This agrees with the results from absorption spectroscopy. Additionally, we measure the fluorescence lifetimes ( $\tau$ ) for all DSBs under N<sub>2</sub>. However, in this case there seems to be no correlation between the values obtained and the electronic nature of the substituents (Tables S3-S4).

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3 As shown before, we correlate the experimental and the calculated values to determine how well  
4 this computational model can predict the emission properties of the DSBs of series 1 and 2 (Figure  
5 10). From this figure we can see an excellent correlation among the  $\lambda_{em}^{cal}$  and  $\lambda_{em}^{exp}$  values ( $r = 0.95$ ),  
6 although the absolute emission wavelength values are about 35 nm shorter than the experimental  
7 values. The computational method correctly predicts that the emission wavelengths of DSBs of  
8 series 1 are shorter than those for DSB of series 2. Notice that the **1-NO<sub>2</sub>-DSB** and the **2-NO<sub>2</sub>-**  
9 **DSB** were excluded from the linear regression presented, due to their particular emission behavior  
10 and photophysics that may result in unusual behavior for this moiety. This explains the difference  
11 in the Stokes shifts, while for the other DSBs the difference between the calculated and  
12 experimental Stokes shifts is not greater than 11 nm, for the nitro derivatives the experimental  
13 values are ~70 nm above the calculated ones (see Stoke shifts in Table S5). Thus, the calculations  
14 correctly predict that including the nitro group leads to the largest red shift but underestimate the  
15 effect of the geometry relaxation in the excited state, and consequently, the calculated emission  
16 wavelength is too short compared to the experimental value. The Pearson's r value for the  
17 correlation in Figure 10a including the **NO<sub>2</sub>-DSBs** is 0.88. On the other hand, Figure 10b compares  
18 the  $k_r^{exp}$  and the  $k_r^{cal}$  values. This figure reveals a good correlation between these quantities  
19 (Pearson's  $r = 0.72$  for DSBs-series 1 and 0.85 for DSBs-series 2), although it is not high enough  
20 for quantitative predictions. It was observed that the computational method correctly clusters the  
21 high values of  $k_r$  for DSBs of series 1, which experimentally are associated to low values of  $\tau$  (see  
22 equation 6), and the low values of  $k_r$  for DSBs of series 2 experimentally associated to high values  
23 of  $\tau$ .

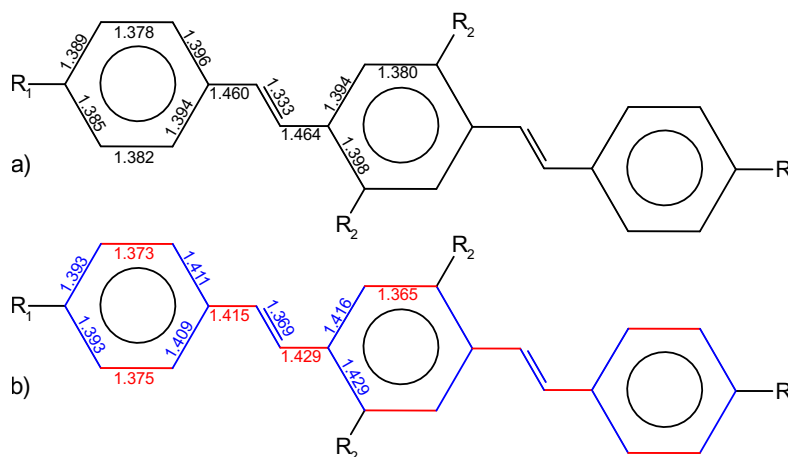


**Figure 10:** Comparison of experimental and calculated emission properties (a)  $\lambda_{em}^{exp}$  and  $\lambda_{em}^{cal}$ , (b)  $k_r^{exp}$  and  $k_r^{cal}$ . Red circles: DSBs of series 1. Blue squares: DSBs of series 2. Black stars: **1-NO<sub>2</sub>-DSB** and **2-NO<sub>2</sub>-DSB**, which are considered as outliers.

In general, the correlation found between experimental and calculated properties shows that TDDFT calculations are an excellent tool to predict structure-property relationships, and they can be useful to approximate the properties not only for DSBs but also for their polymeric analogs, which are usually difficult to determine by experimental methods. Thus, this work also makes a valuable contribution by presenting computational chemistry as an alternative to study conjugated organic materials.

Finally, we analyzed the calculated geometries of the DSBs and found that their geometries significantly change upon excitation and that these changes are similar for all molecules. The most significant changes occur in the C–C bond distances which are presented in Figure 11. We observe that the largest change occurs in the vinyl moiety, where the single bonds contract and the double bonds elongate, and that the changes in the central aromatic ring are larger than in the terminal ones. These geometry changes are coupled to the electron density changes presented in Figure 5, as a density increase upon excitation results in a shorter C-C bond and a density decrease in a

longer C-C bond. The fact that the geometric changes are very similar between the different DSB is consistent with the observation that the Stokes shifts are constant across all the molecules.



**Figure 11:** Calculated average C-C bond lengths in the ground state (a) and in the excited state (b). Bonds lengths that decrease upon excitation are shown in red and those that increase in blue.

The standard deviation of each C-C bond length is  $\leq 0.005$  Å.

## 5. Conclusions

The structure-property relationships of fifteen model DSBs were studied by spectroscopic (UV-vis and FL) techniques and TDDFT computational method. In both cases, it was observed that the ED groups studied have a moderate effect on the bathochromic shift of the absorption and emission bands, whereas the presence of EW groups produces larger shifts. Additionally, the presence of two methoxy groups on the central ring enhances the effect. This was attributed to the substituents changing the orbital energies leading to the reduction of the band gap energies. Thus, ED groups increase the HOMO and LUMO energies, affecting more the HOMO, and the EW groups decrease the energies of both HOMO and LUMO, but their effect on the LUMO is larger.

Good correlations were found between the  $\lambda_{abs}^{cal}$  and  $\lambda_{abs}^{exp}$  values (Pearson's  $r = 0.99$ ) and the  $\lambda_{em}^{cal}$  and  $\lambda_{em}^{exp}$  values (Pearson's  $r = 0.95$ ). The computational method correctly predicts the peak

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3 position for the single transition of DSBs of series 1 and for both transitions of DSBs of series 2  
4 and that the emission wavelengths of DSBs of series 1 are shorter than those for DSBs of series 2.  
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8 Excellent correlations were found also between the  $\Delta E_{abs}^{cal}$  and  $\Delta E_{abs}^{exp}$  (Pearson's  $r = 0.99$ ). The  
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10 computational method correctly predicts that the energy gaps for DSBs of series 1 are larger than  
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12 those for DSBs of series 2. Good agreements were found between the  $F_{osc}^{cal}$  and  $F_{osc}^{exp}$  values  
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14 (Pearson's  $r = 0.87$ ). Although not high enough for quantitative predictions, the theoretical method  
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16 correctly clusters the transitions of DSBs of series 1 as the most intense and the second transitions  
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18 of DSBs of series 2 as the least.  
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23 Distyrylbenzene motifs are important in many polymeric systems with applications in molecular  
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25 electronics, electroluminescent devices, OLEDs and optocouplers. However, the intimate  
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27 understanding of these chromophores in the polymers used for advanced optoelectronic  
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29 applications is hindered by the intrinsic properties of the polymer, including aggregation  
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31 phenomena, energy migration and limited solubility. The molecules reported here offer the  
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33 opportunity for a detailed understanding of these key distyrylbenzene moieties from both an  
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35 experimental and computational point of view. Thus, this study shows that the DSB systems  
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37 possess valuable optoelectronic properties, which can be adjusted with the inclusion of substituents  
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39 with strongly ED and/or EW nature. The study of the structure-property relationships by  
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41 computational techniques helps us to choose properly the structures that will be most appropriate  
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43 and efficient for the desired applications.  
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