

THESIS FOR THE DEGREE OF LICENTIATE OF ENGINEERING IN SOLID
AND STRUCTURAL MECHANICS

Experimental Validation of a Phased Array Ultrasonic Testing
Probe Model and Sound Field Optimization

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Cover:

Beam forming in steering and focusing case using a 16-elements linear phased
array probe with non-linear delay law

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ABSTRACT

New manufacturing technologies are developed to facilitate flexible product designs and production processes. However, the quality of the final products should not be compromised, especially for safety prioritized industries, e.g. aerospace industry. The assessment of product quality and integrity lies on various nondestructive inspection methods and the ultrasonic testing method, among others, is widely used as an effective approach. The phased array technique in the ultrasonic testing area shows more advantages comparing to conventional ones and is revealing more benefits to industrial applications. To incorporate new technique into practical operations, it needs to be qualified with practical experiments. Due to the extensive costs and considerable challenges with experimental works, the necessity of researching on numerical simulation models arises and several models had therefore been developed. The numerical simulation model implemented in the software, simSUNDT, developed at the Scientific Center of NDT (SCeNDT) at Chalmer University of Technology is one of these models for ultrasonic inspection. However, the validity of the models should be proved before supporting or replacing the experiments, and this validation work should be accomplished by experiments ultimately.

In the current work, the main purpose is to further validate the phased array probe model in simSUNDT by comparing simulation results with corresponding experiments. An experimental platform is built with the intention to fully control the operation conditions and the set of testing results. Well-defined artificial defects in test specimens are considered in both simulations and experiments. Comparisons in the end validate the current phased array probe model and could be treated as an alternative to experiments.

With the aid of this validated probe model, optimization of the generated sound field from a phased array probe is then conducted. The optimization aims at searching for a proper combination of main beam angle and focus distance of the probe at this stage, so that the echo amplitude from a certain defect reaches its potential maximum.

Keywords: Experiments, Phased array, Validation, Sound field optimization

PREFACE

This thesis includes the work conducted during the year 2018 - 2020, which is part of a research project *Adaptive Nondestructive Testing of Additive Manufacturing* (*Adaptiv oförstörande provning vid additiv tillverkning på Svenska*). This research project is a collaboration between GKN Aerospace Engine Systems and Chalmers University of Technology, funded by the Swedish innovation agency VINNOVA (reference number 2017-04856) within the national aeronautical research program (NFFP7). The actual work was carried out at the research group of Advanced Nondestructive Testing, department of Industrial and Materials Science, Chalmers.

During my studies at Chalmers, I got lots of help from my surroundings. Special thank is given to my supervisor Prof. Håkan Wirdelius for his great help along the way, and his ideas that always pull me out from my dead ends. He taught me that there is no unique correct answer in researches, but one should try the best to find the way to solve a problem. I also appreciate the help from Anders Rosell from GKN in all forms of discussion about my work. The technical helps from Farham Farhangi and Daniel Snögren from the Swedish Qualification Centre (SQC) are also highly appreciated when I first touched the experimental work and software. Furthermore, I also want to say thanks to other members in our research group, it is good to share the Fika time with them!

Finally, I am really grateful to my beloved family that always support and encourage me in China, and to my best friend Hao Wang staying with me when I was in hard time.

Gothenburg, 2020
Xiangyu Lei

THESIS

This thesis consists of an extended summary and the following appended papers:

Paper A Xiangyu Lei, Håkan Wirdelius, Anders Rosell. "Experimental Validation of a Phased Array Probe Model in Ultrasonic Inspection". *Submitted for publication*

Paper B Xiangyu Lei, Håkan Wirdelius, Anders Rosell. "Experimental Validation and Application Of A Phased Array Ultrasonic Testing Model On Sound Field Optimization". *Submitted for publication*

The appended papers were collaboratively prepared with co-authors. The author of this thesis was responsible for the major progress of the work, i.e. designing and performing of experiments, simulations, post-processing and writing of the papers, all with the assistance of the co-authors.

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Part I

Extended Summary

1 Introduction

Ultrasound as a method of nondestructive evaluation (NDE) and nondestructive testing (NDT) has been used for centuries in both medical and industrial fields thanks to its rapid and safe characters.

Ultrasound or the ultrasonic wave, as indicated by its nomination, is the sound wave that vibrates at frequencies over 20 kHz. Typical frequencies used in ultrasonic NDE applications are ranged from 50 kHz to some MHz. The fundamental theory of using ultrasound to perform evaluation and characterization is that the sound waves can propagate in solids, liquids and gas. The features, such as propagation velocity and attenuation of the sound waves, can be used to characterize the material properties in terms of its structure, elastic properties, etc. When propagating in an object, the waves interact differently to the variations in the object, for example in a fabricated structure, when encountering a pore (air) that has a different density compared to the surrounding structures, the waves scatter as echoes, so that one can detect these echoes and have an understanding of the different composition inside the object.

Today, many industries are using ultrasound as one of the inspection media that ensure the structural integrity and quality of the manufactured components. New techniques in this area are continuously developing along with the arising and application of the new manufacturing technologies.

1.1 Background

In aerospace industries, in order to fulfill the increasing demands of aeroengines manufacturing with low environmental impact and reduced fuel consumption, advanced production and lightweight technologies are developed, e.g. additive manufacturing (AM). This new technology enables effective manufacturing, materials utilization and energy optimization, which is suitable for lightweight components production. However, new manufacturing methods come along with problems, e.g. heat caused deformation and the appearance of unknown vital defects inside the components specifically related to the production process. Thus, the application of these new production technologies demand higher and more reliable inspection methods for quality insurance and to maintain certain safety margins. In common sense, destructive testing (cut-up) is one way of accomplishing it, but it is

expensive, time-consuming and obviously, the destroyed components should not be used any longer and is not suitable for small quantity of objects. To overcome these drawbacks and to ensure the quality at an early stage as well as reducing production costs, NDT methods should be considered.

Ultrasound, or refer to Ultrasonic Testing (UT), among others, has been proven to be an effective and powerful inspection method within this context. The phased array technique in this area, referred to Phased Array Ultrasonic Testing (PAUT), is even more advanced that enables more flexible and faster operations, which compared to the conventional single-element UT has the possibility of enhanced signal-processing.

Before applying the new inspection techniques and evaluation procedures, comprehensive qualifications must be performed for both NDT techniques and related personnel [1]. Traditionally, the qualification processes are based on extensive experiments towards actual test specimens, which means that many variables should be characterized and limited to situations related to the specific application scenario. In addition to the massive cost of producing representative test specimens, the challenges of manufacturing characterized defects in critical locations are considerable. Such experiments could however, be assisted or even partly be replaced by mathematical models that had been developed in recent decades [2, 3]. These models include such as CIVA [4, 5], Thompson-Gray Measurement Model used in UTSim [6], Finite Element Method (FEM) model [7, 8], Elastodynamic Finite Integration Technique (EFIT) [9, 10], etc. With the help of these models, the corresponding experiments can be emulated numerically and the involved physical principles can be studied for further development. For instance, Azar et al. [11] took advantage of a numerical model to simulate the pressure field of a PA probe in beam steering and focusing cases, to show e.g. the impact of focusing to detection resolution in the near field. Puel et al. [12] utilized numerical simulation with a evolutionary algorithm based optimization method to reach an optimal design and setting of phased array probe. Other similar studies can be seen in e.g. [13, 14].

However, the numerical models should be thoroughly validated before practical use. This is performed either by comparing the model with other models that had been validated, or by comparing the model with experimental results, which should ultimately be done in order to show that the model truly reflects the reality [15, 16].

The UT simulation model developed by Chalmers University of Technology is implemented into software, simSUNDT [17]. The conventional probe model involved has been validated [18–20] by comparing the simulation with the experimental benchmark study [21], which was started by the World Federation of NDE Centres and the experimental results were provided by Commissariat a l'énergie

atomique (CEA, France). The PA probe model [22] however, was validated only to a limit extent and thus requires further validation work. After the PA probe validation, the model is capable of facilitating the sound field optimization in order to get the maximized echo amplitude towards a defect, performed in the present work.

1.2 Aims and limitations

This thesis is part of a research project *Adaptive Nondestructive Testing of Additive Manufacturing*, which has the overall objective to incorporate the developed tools into inspection technologies during manufacturing processes, with the intention of quality control and assessment. The tool should also be optimized towards well-defined manufacturing defects, which potentially may occur in manufacturing processes. To accomplish this goal, the ultrasonic wave propagation within complex geometries should be studied through corresponding mathematical models. The overall research questions are hereby lie on e.g., how well the mathematical models can reflect on the physical inspections, what flexibility can be provided by the models, how can the models help with production and process optimizations, etc. The present work is thus focused on further experimental validation of the newly developed PA probe model in the software, simSUNDT. The experiments are performed at the NDT lab of the Scientific Center of NDT (SCeNDT) at Chalmers University of Technology using a newly built mechanized operation platform. After the validation of the model, it is then used to investigate the approach of sound field optimizations towards a backwall surface breaking defect, with the intention to retrieve a maximized corner echo. This is mainly conducted by adjusting the combination of beam angle and focus distance of a PA probe, to optimize the sound field for specific defect characteristics at this stage.

However, some limitations are involved. In the validation work, the considered defects only address well-defined artificial ones, i.e. SDH and surface breaking defect. Source of attenuation including material damping properties of the test specimen and contact conditions that could influence beam divergence characters are not included in the corresponding simulations, because the validation focus is on the PA probe model instead of material properties. These limitations could provide error source to the validation results.

1.3 Thesis structure

This thesis is structured in following sections according to the aims. Section 1 provides background information as well as the objectives and limitations of the

current work. Section 2 introduces basic knowledge related to the ultrasonic testing method. The software simSUNDT is also briefly introduced with theoretical base and capabilities. Section 3 presents a simple and general approach to phased array delay law derivation based on geometrical wave path compensation. Section 4 summarizes the practical information of experiments used in probe model validation work. Section 5 gives an overview of the optimization work and the applied algorithm. Section 6 lists the summary of appended papers. The extended summary ends with Section 7, where some concluding remarks are provided with the future direction of work.

2 Technical Background

2.1 Ultrasonic Testing (UT)

2.1.1 Basic principles of ultrasound generation and receiving

The ultrasonic testing method is all based on the ultrasonic waves, which is generated from a specific transducer or probe. The essential part inside a conventional probe is the piezoelectric element that converts the imposed electrical pulse into mechanical movement of the element, and vice versa. The vibration frequency of this element is determined by its thickness as indicated by equation (3.98) in [23], i.e. the thinner the element, the higher the frequency. When the element is vibrating, a pulse of sound will be generated in the adjacent medium (solid, liquid and gas) in the form of acoustic waves. The waves can propagate within the medium and scattered if strikes an object, e.g. pore or inclusion, which has a different *acoustic impedance* comparing to the surrounding medium. The scattered waves then travel in all directions depending on the property of the object and some of them return back to the probe. The piezoelectric element receives the vibration and converts it back to electrical pulse, which is then amplified as the output signal. This signal is represented as pulse amplitude versus elapsed time. By knowing the wave traveling speed in the medium, the line distance between the probe and the object can thus be determined. Therefore, the ultrasonic testing method, like all other NDT methods, is indirect by nature, i.e. the measured outputs are in the form of pulse vs. time and that other quantitative information are not directly obvious, e.g. the size of the object, which need further interpretations. Nevertheless, the quantitative information could be determined to some extent in post-processing if moving the probe to collect signals from different positions and orientations.

2.1.2 Near field and far field

As the piezoelectric element converts an electrical pulse into mechanical movement, the generated sound waves travel along their paths. Due to the reason that the actual element is not a point but has an area, the waves generated from each point of the area must have different path lengths along the traveling direction, which cause interference both constructively and destructively that end up with intensity fluctuations in a short distance. The region within this distance is called *near field* and beyond that is called *far field*, where the wave path differences are smaller, leading to minor intensity fluctuations and the intensity decreases slowly with growing distance, as expected. The near field length N for a certain probe can be expressed by:

$$N = \frac{D^2}{4\lambda} \quad (2.1)$$

where λ is the wavelength and D is the diameter of a circular probe.

As a simple example, the normalized sound pressure level along the axis of a piston oscillator, on which the displacement is the same for all points, can be seen in Figure 2.1 [24], where p_0 is the average pressure value in the near field region. It is seen that the sound pressure level fluctuates in the near field and reaches its last maximum at near field length. After that, the pressure decreases slowly as the distance increases and returns to the average pressure level at the third times of the near field length. Practically, most experiments are performed in the far field due to the nature of unstable wave pressure in near field.

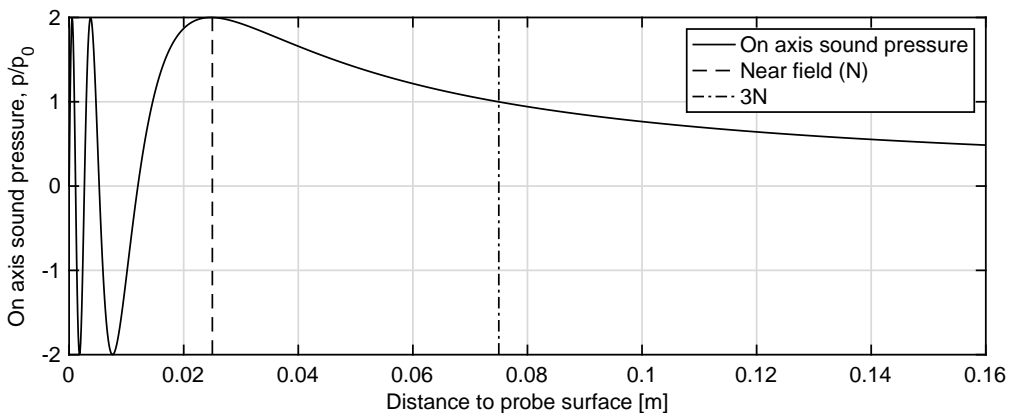


Figure 2.1: *Sound pressure variation along the axis of a piston oscillator as a function of distance to the probe surface*

2.1.3 Wave types

For sound waves, there are two types of traveling in bulk materials, namely *longitudinal waves* and *transverse waves*. They are differed by the movement directions of the medium particles. There are however other types of waves, such as *Rayleigh waves* and *Lamb waves*, also available for certain inspection purposes [23], but the longitudinal and transverse waves are mostly used in ultrasonic testing.

Different wave modes can be converted. This phenomenon is called *mode conversion*, which happens when the waves strike on the material interface at an oblique angle. The waves are reflected in the same medium, called *reflection*, but it can also travel into the second medium, called *transmission*. It should however be noted that the mode converted waves will propagate in a different direction, which can be determined through *Snell's law*. This law is used to determine the direction (angle) of reflected and transmitted mode-converted waves when the waves strike on the material interface at an oblique angle. Refer to Figure 2.2 it can be expressed as:

$$\frac{\sin \theta_i}{v_i} = \frac{\sin \theta_{rs}}{v_{rs}} = \frac{\sin \theta_{rl}}{v_{rl}} = \frac{\sin \theta_{ts}}{v_{ts}} = \frac{\sin \theta_{tl}}{v_{tl}} \quad (2.2)$$

where

θ_i - incident wave angle

θ_{rs} - reflected transverse wave angle

θ_{rl} - reflected longitudinal wave angle

θ_{ts} - transmitted transverse wave angle

θ_{tl} - transmitted longitudinal wave angle

v is the wave speed in corresponding medium for each case

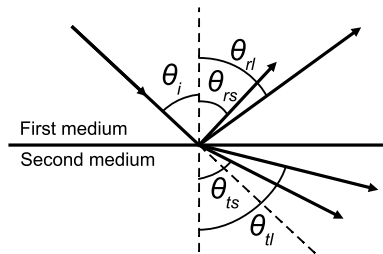


Figure 2.2: *Wave mode conversion and Snell's law demonstration*

It is worth mentioning that the Snell's law only shows the propagation directions of the waves without any information on their amplitudes. Looking at Snell's law, it can be noticed that when the transmitted wave angle reaches 90° in the second

medium, it would propagate along the material interface. This corresponding incident wave angle is thus called the *critical angle*. Knowing the fact that the longitudinal waves generally travels faster than the transverse waves in mediums, the longitudinal wave angle should be larger than the transverse wave angle.

2.1.4 Data presentation

As mentioned in Section 2.1.1, the ultrasonic testing is an indirect method since the result presentations need to be post-processed before interpretation. The direct variables used in conventional inspections are *time-of-flight* and *amplitude*. All other information and result presentation formats are derived based on these variables. Conventionally, there are three formats used to visualize the inspection results, known as A-scan, B-scan and C-scan, each of which presents the data in a different way for evaluation, see an inspection example in Figure 2.3. Previously mentioned pulse amplitude vs. time-of-flight direct presentation is denoted as A-scan, directly obtained through inspection equipment and provides a straightforward echo indication at which distance an object exists. B-scan shows a cross-section view of the test specimen. It is generated by re-organizing and color-scaling an A-scan amplitude along the vertical axis (time-of-flight axis) at a scan position, and repeating for all scan positions of the probe along the horizontal axis. In this way, given the wave speed in the test specimen, the depth of the defect and its approximate linear dimension in the scan direction can be determined. C-scan shows a color-scaled top view of the defects in test specimen for a raster scanning sequence, parallel to the scanning surface. It is obtained through color-scaling the maximum echo amplitudes at each scanning position. C-scan can also be an echo dynamic curve for an one-line scanning sequence, showing the maximum echo amplitudes at each scan position (the case in Figure 2.3). It is useful when presenting the defect distribution in the scan plane. Obviously, the scanning position information is needed to generate the B-scan and C-scan. With a modern ultrasonic inspection equipment, together with a mechanized system equipped with encoders providing probe position information, it is possible to visualize these scan presentations simultaneously during real-time inspection.

2.1.5 Probe configuration

Basically, there are two ways to configure the ultrasonic probes during inspections depending on the number of probe/probes used. If only one probe is used for both generating and receiving the waves, it is called *pulse-echo* mode. This mode simplifies the overall experimental setup, but the *dead zone* occurs in the received signals, shown in Figure 2.3. The dead zone refers to a region where the initial

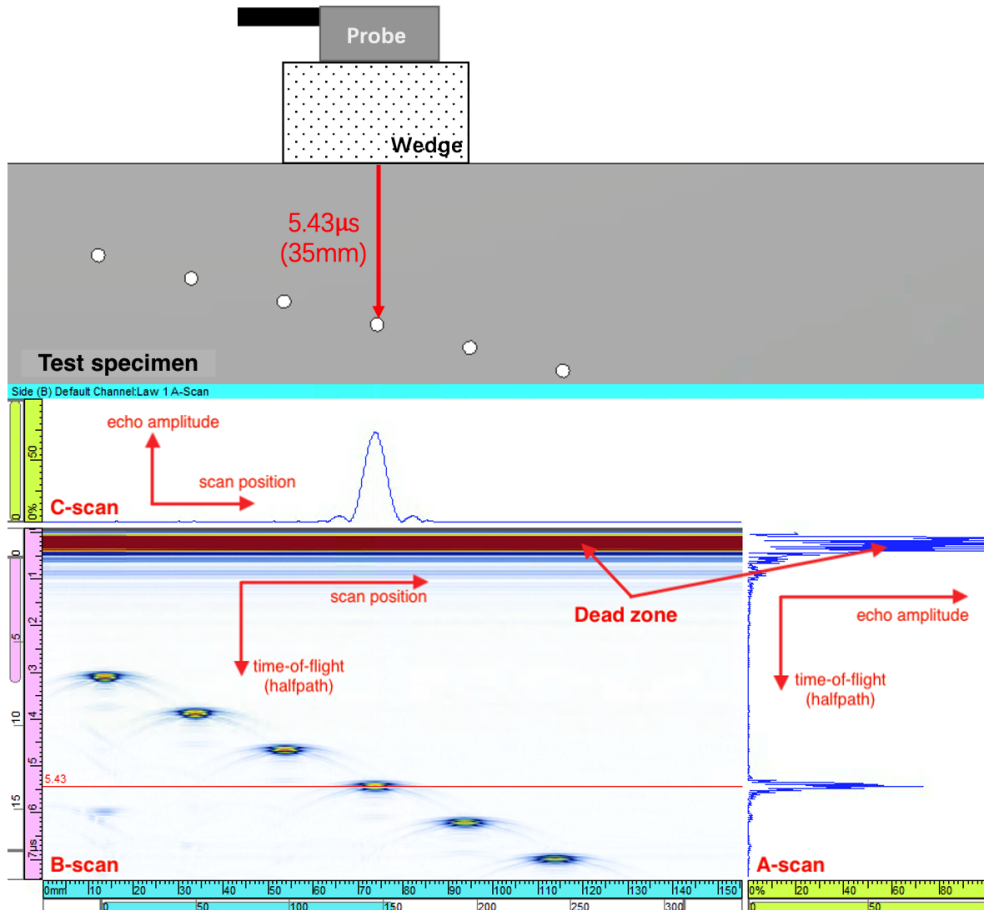


Figure 2.3: An example of inspection situation, where the current probe center sits above a defect (35 mm depth). The results are presented in A- B- and C-scan (echo dynamic curve in this case of a one-line scan) and the dead zone is also visible in A- and B-scan

large electrical pulse generating the waves is also treated as a received signal, which has a large initial echo amplitude in the result presentations. It can totally mask an echo from a real defect within this depth region of the specimen. This drawback can be overcome if one probe is used for generating the waves while another probe is for receiving at a different position along the wave path, called *pitch-catch* mode, because the initial large electrical pulse will not be detected by the receiving probe. This mode is therefore more appropriated for near-surface defects detection and detection on thin materials. Time of Flight Diffraction Technique (TOFD) is one of the examples using pitch-catch configuration.

2.2 Phased Array UT (PAUT)

The *phased array* (PA) configuration generally refers to a series of transducers in an array setup, ordered linearly or in matrix form. *Phased array ultrasonic testing* (PAUT) therefore refers to an ultrasonic testing method that utilizes the specific probes, which consist of an array of small piezoelectric elements. The element shape is in most cases rectangular since they are cost-effective to produce, which is used in two types of arrays, namely *linear* and *2D arrays* based on how the individual rectangular elements are arranged. Figure 2.4 for example, shows a sketch of a linear PA probe surface with corresponding terminologies used in this field. There are also other types of PA probes with overall angular shapes [25] in PAUT, but the linear and 2D arrays are most generally used in NDE.

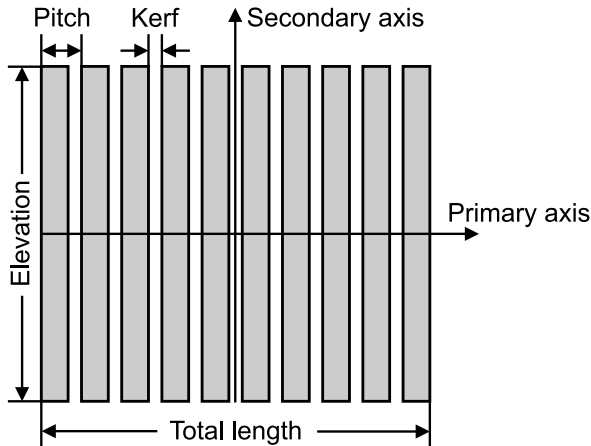


Figure 2.4: *The terminology of a linear phased array probe*

Each of these small elements can be triggered individually by electronic pulse, and the signal response can also be received independently. The advantages of this configuration comparing to the conventional single-element UT are revealed by the flexibility of sound beam *steering* and *focusing* through constructive interference of waves, together with the possibility of fast visualization generation [25]. As it is known for conventional single-element UT, the beam steering and focusing is accomplished by using corresponding wedge that has an angle or curved, which needs to be applied based on a certain application scenario. Whereas for the PA probe, each element can be driven separately, thus it is of nature that there can be linear time shifts between pulses. The individual small element now approximates as a point source that emits spherical waves, while the generated synthetic wave front after constructive interference is hereby a plane wave and can travel in

different directions based on different combinations of time shifts. This time shifts combination is called *delay law*, and its general derivation is seen in Section 3. Furthermore, if the applied delay law is nonlinear, the generated wave front can also emulate the focusing effect, just as the conventional single-element ultrasonic probe with a curved wedge, that the sound beam is focused at a certain point where the sound pressure reaches maximum. The beam steering and focusing can also be combined in more complex operations if needed.

When each of these elements receive signals, the corresponding time delay laws generally can also be used to shift the individual signals so that all of them appear at the same time and can be summed up to obtain a single and large signal in the end. Thus, in this transmitting and receiving process with proper delay laws, the effect of using a PA probe is the same as using a conventional single-element probe associated with the corresponding angled or curved wedge. The advantage of using a PA probe is that the beam steering and focusing effect are obtained simply by adjusting the proper time delay laws, with no need of moving or changing the physical probe as in the conventional inspections. This flexibility enables many ultrasonic measurements in a rapid and simple way, for example a rapid sectorial scan in a region using a single PA probe and real-time imaging based on acquired data. Figure 2.5 shows an inspection case as an example, using a single PA probe on an AM component (Titanium Alloy 6AL4V) with side-drilled holes (SDHs) at depth of 14 mm. The SDHs under inspection have diameters of 1.2 mm, 0.8 mm and 0.4 mm from left to right. The probe contacts directly on the surface and its position is unchanged as shown in the figure. The delay laws together with the pulse sequences applied are: (a) 16-elements aperture travels linearly along the array without focusing effect; (b) 16-elements aperture travels linearly along the array with focusing at SDHs; (c) 64-elements aperture sectorial sweeping from -45° to 45° with focusing at SDHs and (d) Full Matrix Capture (FMC) that will be introduced in Section 2.3. The corresponding data visualization in B-scan can be seen in Figure 2.6, where the FMC data is visualized using an advanced imaging algorithm, Total Focusing Method (TFM), also introduced in Section 2.3. It can be noticed in Figure 2.6(a)-2.6(c) that weaker ghost indications are also visible below the SDH indications, which probably come from bottom sphere reflections, whereas the ghost indications are eliminated in Figure 2.6(d) by TFM.

The near field length of a general rectangular PA probe can also be calculated using equation (2.1), where the probe diameter D is now the total length of the PA probe (see Figure 2.4) if the aspect ratios between the total length and the elevation is larger than 0.6 [13]. By focusing effect, the resolution within PA probe's near field can be improved [11], as compared by Figure 2.6(a) and Figure 2.6(b). It is also noted that the generated sound beam in the far field is similar to the one generated by the conventional single-element probe that has the same

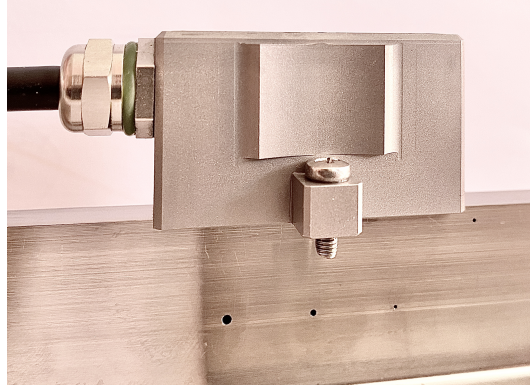
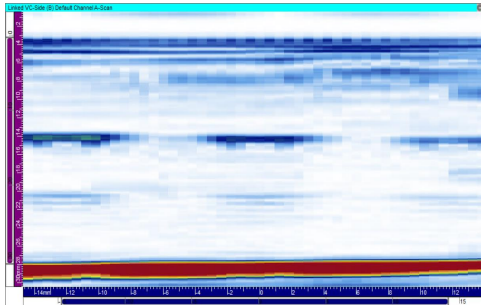
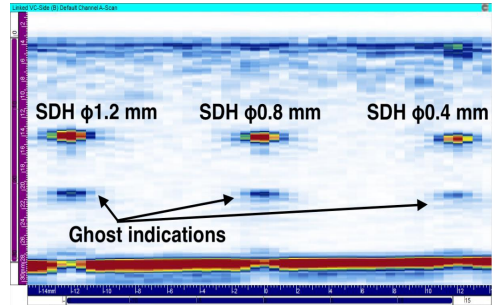


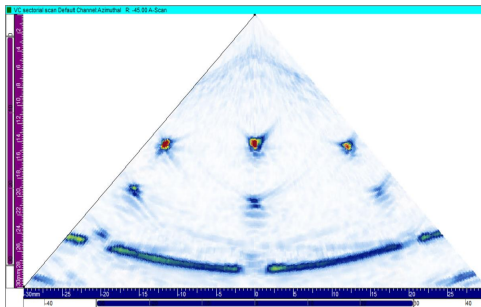
Figure 2.5: PA probe inspects an AM component with SDHs (diameter of 1.2 mm, 0.8 mm and 0.4 mm from left to right) at depth of 14 mm from scanning surface



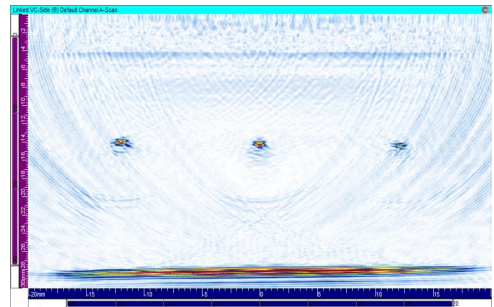
(a) Linear scan without focus



(b) Linear scan with focus at SDHs



(c) Sectorial scan



(d) TFM algorithm

Figure 2.6: Four B-scans of flexible scanning and imaging possibilities using PA probe on an AM component with SDHs in the middle, as the case shown in Figure 2.5

overall size as the PA probe. Because of these advantages, PAUT had been applied in medical applications since decades [26, 27] and is now attracting more attentions from the industrial perspectives. For example, Shan et al. [28] built up an PAUT inspection system for defect detection in steel structures. A SDH in the test block can be quantitatively evaluated from the system. Lopez et al. [29] explored the possibility of using PAUT on AM component inspections and it was concluded to be suitable also for quantitative evaluations. Qin et al. [30] introduced an improved delay law calculation for PAUT technique based on the relation between the observed peak offset and propagation distance in high-density polyethylene used in the nuclear power plant pipes, in order to improve the ultrasound field intensity at the focal point of interest and to increase the imaging sensitivity. Gros et al. [31] discussed about more applications and the advantages of using PAUT in industries, as well as the necessity of numerical modeling and some future development possibilities of the technique.

2.3 Full Matrix Capture (FMC)

Using an ultrasonic PA probe, *Full Matrix Capture* (FMC) is a data acquisition method capturing complete set of time-domain signals from each possible pair of transmitter-receiver combination, e.g. time-domain signal from i -th element (transmitter) to j -th element (receiver), S_{ij} , illustrated in Figure 2.7. For an array with N elements, FMC can generate N^2 possible time-domain data. With this complete set of data, any beam-forming scheme of conventional PAUT can be emulated in the post-processing with no need of further experiments.

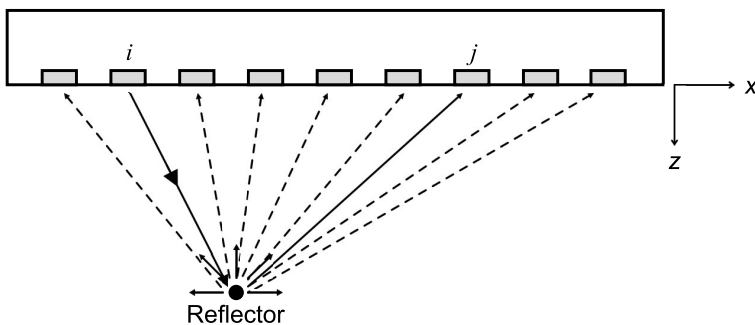


Figure 2.7: *Full Matrix Capture (FMC) illustration*

To fully use the FMC data, an imaging algorithm called Total Focusing Method (TFM) is preferred, which has better imaging performance comparing to conventional algorithms [32], see also comparisons in Figure 2.6. TFM is

performed by firstly discretizing the imaging region of interest into frame grids, then advancing all pairs of FMC signals using the corresponding time-of-flight of this pair to the imaging grid point. The image is finally formed by summing these signals at time $t = 0$ and repeating for all image grid points. Due to the massive computation effort, TFM imaging is practically performed during post-processing.

2.4 simSUNDT

The *simSUNDT* software consists of a Windows[®]-based pre- and post-processor, as well as a mathematical kernel *UTDefect* [19, 33] that conducts the actual mathematical modeling and computation. *UTDefect* was developed at Chalmers University of Technology. The 3D elastodynamic wave equation, which defines the wave propagation in a homogeneous half space, is solved using vector wave functions [33].

The contact probe can be modeled in elliptic and rectangular shape. The modeling assumes that the probe is placed on the surface of an elastic half-space, which has no traction on the surface except beneath the probe. This enables the possibilities of simulating any types of the probe available on the market, by specifying related parameters such as wave types, element size and shape, angles, frequency ranges, contact conditions, etc. The traction is derived so that a plane wave is generated in the far field, shown in equation (2.3) for longitudinal, vertical transverse (SV) and horizontal transverse (SH) wave types, respectively.

$$\mathbf{t} = \begin{cases} Agi\mu k_p \left[\left(\frac{k_s^2}{k_p^2} - 2 \sin^2 \gamma \right) \hat{z} + \delta \sin 2\gamma \hat{x} \right] e^{-ik_p x \sin \gamma}, \text{Longitudinal probe} \\ Agi\mu k_s [\sin 2\gamma \hat{z} - \delta \cos 2\gamma \hat{x}] e^{-ik_s x \sin \gamma}, \text{SV probe} \\ Agi\mu k_s \delta \cos \gamma \hat{y} e^{-ik_s x \sin \gamma}, \text{SH probe} \\ 0, \text{elsewhere} \end{cases} \quad (2.3)$$

where \hat{x} , \hat{y} and \hat{z} are the unit vectors in corresponding directions. A is the displacement amplitude and the function g enables reduction of edge effects. μ is the Lamé constant of the elastic half space. k_p and k_s are longitudinal and transverse wave numbers, respectively. γ is the wave angle. $\delta \in [0,1]$ is a constant that represents the coupling effect, where $\delta=0$ is fluid coupling and $\delta=1$ is glued condition.

The wave propagation is governed by the elastodynamic equation of motion, in which the displacement field, \mathbf{u} , is involved:

$$k_p^{-2} \nabla \nabla \cdot \mathbf{u} - k_s^{-2} \nabla \times \nabla \times \mathbf{u} + \mathbf{u} = \mathbf{0} \quad (2.4)$$

The total displacement field is a summation of the incident field (\mathbf{u}^i) and the scattered field (\mathbf{u}^s), as:

$$\mathbf{u} = \mathbf{u}^i + \mathbf{u}^s \quad (2.5)$$

The incident fields can be expanded in terms of regular spherical partial vector waves ($\text{Re}\Psi_n$), and the scattered field caused by various of defects can be expanded in its outgoing spherical partial vector waves (Ψ_n):

$$\begin{cases} \mathbf{u}^i = \sum_n a_n \text{Re} \Psi_n \\ \mathbf{u}^s = \sum_n f_n \Psi_n \end{cases} \quad (2.6)$$

Volumetric and crack-like defects are available types of defect to be modeled. Specifically, volumetric defects include a spherical/spheroid cavity (pore), a spherical inclusion (isotropic material differing from the surrounding material, i.e. slag) and a cylindrical cavity (SDH). Crack-like defects include rectangular/circular crack (lack of fusion) and strip-like crack (fatigue crack). There is option to model the surface roughness for the rectangular and strip-like crack, and the degree of closure can be modeled for the circular crack. Tilting planar back surface could also be modeled for the strip-like crack, but otherwise it is assumed parallel to the scanning surface. The surface-breaking strip-like crack and rectangular crack close to the back surface can be used to model the corresponding defects in the test piece.

The modeling of these defects are important and the methods of solution can be for example, various types of surface integral equations, null field approach (T-matrix method) and FEM. UTDefect incorporates the T-matrix method [34] and all information regarding the defects is included in the transition matrix, as well as providing the linear relation between the expansion coefficients of the incoming (a_n) and scattered (f_n) wave fields in equation (2.6):

$$f_n = \sum_{n'} T_{nn'} a_{n'} \quad (2.7)$$

To incorporate the probe model into the T-matrix formulation, the displacement field needs to be transformed from the plane vector waves centered at the contact area into spherical vector wave functions oriented and centered at the defect [33].

To model the receiver, a reciprocity argument [35] is applied. In the end, the electrical signal response is expressed by:

$$\delta\Gamma \sim \sum_{nn'} a_n^b T_{nn'} a_{n'}^a \quad (2.8)$$

With the transmitting probe is characterized by $a_{n'}^a$, the defect by $T_{nn'}$ and the receiving probe by a_n^b .

To simulate the entire testing procedure of an actual NDT situation, a calibration option is available towards a reference reflector including for example, the *side-drilled hole* (SDH) represented by the cylindrical cavity [36] and the *flat-bottomed hole* (FBH) approximated by an open circular crack.

The model geometry can be limited by a reflecting backwall and be described as a plate with finite or infinite thickness bounded by the scanning surface, on which the scanning sequence are defined by rectangular mesh.

In addition, it is also possible to suppress the unexpected wave component in the simulation to eventually facilitate the analysis of the received signal. The configurations of the probe can be chosen among pulse-echo, separate with fixed transmitter and tandem configuration (TOFD).

These principles are the same for the phased array probe model, that element is represented by the boundary conditions (traction), from which the plane wave is generated in the far field with a certain angle. The individual boundary conditions are translated into the main coordinate system and a phased array wave front with certain nominal angle is formulated by constructive phase interference. The formulated nominal angle can also be altered by specific delay law, but it should be noted that this is only possible for small angles if no wedge is specified.

3 Delay law derivation

As mentioned in Section 2.2 that the delay laws are used to steer and focus the generated beams for the PA probes. It is easy to imagine that a focusing effect by nature requires also beam steering of some probe elements. In other words, beam steering can be treated as a special case to beam focusing with infinite focusing length. In this section, the general delay laws for linear PA probe are derived for both beam steering and focusing. It is basically done by compensating for the geometrical wave path lengths between each element, as each element of the array emits a wave that travels in a straight line at high wave frequencies [37].

Before the derivation, some nomenclatures of the normal parameters used for the PA probe, the wedge, material properties and inspection case are firstly introduced in Section 3.1. The wedge is generally involved in practical applications that facilitates the mechanized operation and coupling possibilities, and the first element of a PA probe is set to the lower edge side of the wedge, to be consistent with the actual industrial application. For more general cases, an angled wedge is used in these derivations.

3.1 Nomenclatures

Below lists the nomenclatures specifically used for these derivations. It should be noted that all these parameters are known before specifying a delay law, except for Φ_1 inside the wedge.

Table 3.1: Nomenclatures used specifically for delay law derivations

Notation	Nomenclatures
N	Total number of elements in a PA probe
p	Pitch between elements
ϵ	Wedge angle
h_1	Vertical distance of the center of the first element to the wedge bottom
c_1	Wave speed in the first medium (wedge)
c_2	Wave speed in the second medium (test specimen)
D_f	Vertical focusing depth from the surface of test specimen
Φ_1	Primary beam impinging angle in the first medium (wedge)
Φ_2	Refracted (expected) beam angle in the second medium (test specimen)

3.2 Beam steering

The geometrical parameters used in beam steering case are indicated in Figure 3.1. Assume that all probe elements on the wedge surface are activated in the beam forming process. For this pure steering case, the generated wave should be a plane wave, which means that each wave path should be in parallel with each other and propagates in a specified refracted angle Φ_2 in the second medium. Take the i -th and the last element for example, the generated wave front in the second medium should thus be represented by line BE in Figure 3.1, which is perpendicular to all wave paths. To accomplish that, when the wave from the i -th element reaches point E , the wave from the last element should just reach point B . The time difference between these two waves should therefore be the time delay of the wave that has the shortest time of flight, so that all waves reach the wave front at the same time. Based on this principle, the time of flight of the i -th wave can be derived and the delay law is obtained.

As indicated in Figure 3.1 that the distance on the wedge surface between the i -th element and the first element is $(i - 1)p$, and to the last element is $(N - i)p$, thus the height of the center of the i -th element (point F) to the wedge bottom (point C), h_i , can be expressed by:

$$h_i = h_1 + (i - 1)p \sin \epsilon \quad (3.1)$$

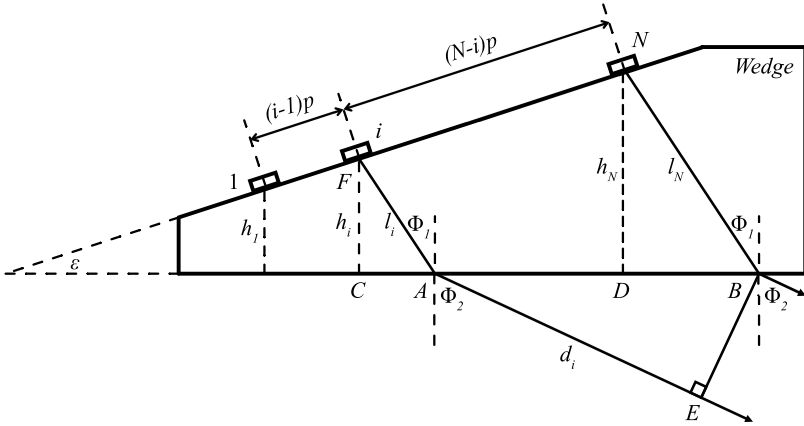


Figure 3.1: Geometrical parameters for delay law derivation in pure steering case

The wave path length for the i -th element in the wedge, l_i , from the center of this element (point F) to point A (intersecting point between the wave path and the wedge bottom) is expressed in term of h_i by:

$$l_i = \frac{h_i}{\cos \Phi_1} = \frac{h_1}{\cos \Phi_1} + p(i-1) \frac{\sin \epsilon}{\cos \Phi_1} \quad (3.2)$$

Using equation (3.1) for the i -th element and the last element, the horizontal distance between the center of the element and its intersecting point, CA_i and BD , respectively, can be expressed by:

$$CA_i = h_i \tan \Phi_1 = [h_1 + (i-1)p \sin \epsilon] \tan \Phi_1 \quad (3.3)$$

$$BD = h_N \tan \Phi_1 = [h_1 + (N-1)p \sin \epsilon] \tan \Phi_1 \quad (3.4)$$

The horizontal distance between the i -th element and the last element, DC_i , can be written as:

$$DC_i = (N-i)p \cos \epsilon \quad (3.5)$$

Combining equations (3.3), (3.4) and (3.5), the horizontal distance between the intersection points A and B for the i -th element and the last element, BA_i , can finally be written as:

$$BA_i = BD + DC_i - CA_i = p(N-i)(\sin \epsilon \tan \Phi_1 + \cos \epsilon) \quad (3.6)$$

Then, using equation (3.6) the wave path length of the i -th wave in the second medium, d_i , can be expressed by:

$$d_i = BA_i \sin \Phi_2 = p(N - i)(\sin \epsilon \tan \Phi_1 + \cos \epsilon) \sin \Phi_2 \quad (3.7)$$

Combining equation (3.2) and (3.7), the time of flight of the i -th wave from the element center (point F) to its corresponding point E , indicated in Figure 3.1, is expressed by:

$$\begin{aligned} t_i &= \frac{l_i}{c_1} + \frac{d_i}{c_2} \\ &= \frac{1}{c_1} \left[\frac{h_1}{\cos \Phi_1} + p(i - 1) \frac{\sin \epsilon}{\cos \Phi_1} \right] \\ &\quad + \frac{1}{c_2} [p(N - i)(\sin \epsilon \tan \Phi_1 + \cos \epsilon) \sin \Phi_2] \end{aligned} \quad (3.8)$$

Finally, the time delays Δt_i are obtained simply through subtracting all the individual time of flight, t_i ($i=1, \dots, N$), by the largest one:

$$\Delta t_i = t_i - \max(t_i) \quad (3.9)$$

For example in a simple case, a PA probe with 5 elements:

time of flight, t_i : $5\mu s, 4\mu s, 3\mu s, 2\mu s, 1\mu s$;

time delay, $|\Delta t_i|$: $0\mu s, 1\mu s, 2\mu s, 3\mu s, 4\mu s$

It is worth noting from the equation (3.8) and (3.9) that, the time delays between elements are a linear relation for pure steering case.

3.3 Beam focusing

When the beam focusing effect is considered in the delay law generation, the goal is to have all refracted waves reach a common point at certain depth, D_f , in the second medium, which is the *focusing depth* one expected and specified in advance. Unlike the pure steering case, since the refracted waves are not in parallel in the focusing case, the corresponding primary impinging waves inside the wedge should not be in parallel either. In other words, each wave path has a different unique impinging angle at the wedge bottom. This is very important to notice because all the impinging angles are not the same as the one calculated by Snell's law using the nominal refracted angle Φ_2 . This specified nominal refracted angle Φ_2 in practice however, refers to the refracted angle of the overall synthetic wave beam axis in the second medium, see Figure 3.2.

The corresponding nominal impinging angle, Φ_1 , calculated by Snell's law in the wedge, again, refers to the impinging angle of the overall synthetic beam axis

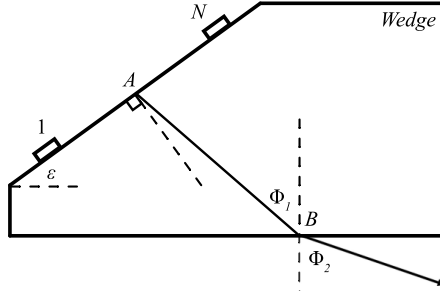


Figure 3.2: *Illustration of the nominal refracted angle in beam focusing case*

in the wedge, represented by a line from the center of active aperture on the wedge surface (point A) to the synthetic beam exit point (point B) at wedge bottom. This exit point B is hereby determined provided that the center of active aperture and Φ_1 are known. Note that the line AB does not have to be a wave path of any element but it is a representative of the overall synthetic sound beam in beam focusing case.

Understanding the above principles, the time-of-flight of each wave and the delay law derivation for focusing case are followed, refer to Figure 3.3, where the geometrical parameters are indicated. Assume again that all probe elements are active on the wedge surface in this focusing case.

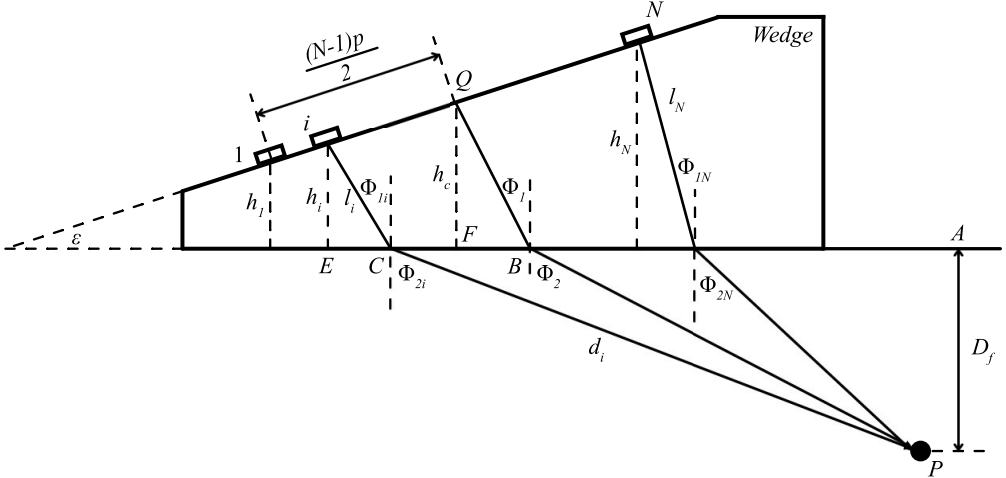


Figure 3.3: *Geometrical parameters for delay law derivation in focusing case*

Since the vertical focusing depth D_f and the expected refracted angle Φ_2 are specified in advance, the horizontal distance between the focusing point (point P)

and the synthetic main beam exit point (point B), AB , is a known value:

$$AB = D_f \tan \Phi_2 \quad (3.10)$$

For the main beam path, the horizontal distance between the center of active aperture on the wedge surface (point Q) and the main beam exit point (point B), BF , can be calculated based on the height of point Q to the wedge bottom. This vertical distance, h_c , is expressed in a similar way as equation (3.1):

$$h_c = h_1 + \frac{(N-1)p}{2} \sin \epsilon \quad (3.11)$$

$$BF = h_c \tan \Phi_1 = \left[h_1 + \frac{(N-1)p}{2} \sin \epsilon \right] \tan \Phi_1 \quad (3.12)$$

Then, the horizontal distance between the center of the i -th element and the center point of the active aperture (point Q), FE_i , is expressed by:

$$FE_i = \left[\frac{(N-1)p}{2} - (i-1)p \right] \cos \epsilon \quad (3.13)$$

It can be seen from equation (3.13) that the value of FE_i for the elements locate after the aperture center point are negative. Combining equation (3.10), (3.12) and (3.13), the horizontal distance between the focusing point P and the center of the i -th element can generally be written by:

$$AE_i = AB + BF + FE_i \quad (3.14)$$

As mentioned earlier and refer to Figure 3.3 that each wave impinging the wedge bottom at an unique angle, i.e. $\Phi_{1i} \neq \Phi_1$, and refracted at a corresponding unique angle, i.e. $\Phi_{2i} \neq \Phi_2$, for the i -th wave, it is thus no ease to determine the impinging point for an individual wave. However, there should be an impinging point for the i -th element (point C) that satisfies the Snell's law for its corresponding angles Φ_{1i} and Φ_{2i} , namely [37]:

$$\frac{\sin \Phi_{1i}}{c_1} = \frac{\sin \Phi_{2i}}{c_2} \quad (3.15)$$

Using equation (3.1) for element height, h_i , to rewrite the equation (3.15) and

write it as a function of the unknown variable EC_i for the i -th element, $f(EC_i)$:

$$\begin{aligned}
 f(EC_i) &= \frac{\sin \Phi_{1i}}{c_1} - \frac{\sin \Phi_{2i}}{c_2} \\
 &= \frac{EC_i}{l_i} \frac{1}{c_1} - \frac{AC_i}{d_i} \frac{1}{c_2} \\
 &= \frac{EC_i}{\sqrt{EC_i^2 + h_i^2}} \frac{1}{c_1} - \frac{AE_i - EC_i}{\sqrt{(AE_i - EC_i)^2 + D_f^2}} \frac{1}{c_2}
 \end{aligned} \tag{3.16}$$

Finding a value of EC_i that makes the equation (3.16) equals to zero means that the Snell's law in equation (3.15) is fulfilled, thus the impinging point C_i for the i -th element is found. Furthermore, its corresponding angles are found by:

$$\begin{aligned}
 \Phi_{1i} &= \arctan\left(\frac{EC_i}{h_i}\right) \\
 \Phi_{2i} &= \arctan\left(\frac{AE_i - EC_i}{D_f}\right)
 \end{aligned} \tag{3.17}$$

Using equation (3.17), the wave path length of the i -th element in both mediums can be expressed by:

$$\begin{aligned}
 l_i &= \frac{h_i}{\cos \Phi_{1i}} \\
 d_i &= \frac{D_f}{\cos \Phi_{2i}}
 \end{aligned} \tag{3.18}$$

Thus, the time-of-flight of the i -th wave from its element to the focusing point can be calculated using the equation (3.18):

$$t_i = \frac{l_i}{c_1} + \frac{d_i}{c_2} \tag{3.19}$$

Finally, the same as in the pure steering case, the corresponding delay law for focusing case can be obtained using equation (3.9).

4 Experimental setup

This section summarizes the experimental setup, including the UT equipment, the mechanized gantry system, the test specimens used in the experiments, and the design of experiments.

4.1 UT equipment

The UT equipment mainly includes the data acquisition hardware (TOPAZ64) and a commercial PA probe with corresponding plastic wedges.

TOPAZ64 (Figure 4.1) is a portable 64-channel phased array ultrasonic testing equipment with FMC and TFM capabilities incorporated. Ultrasonic inspection data is communicated in real time between TOPAZ64 and computer by Gigabyte Ethernet cable connection, and is processed by corresponding software UltraVision on the computer.



Figure 4.1: *Data acquisition hardware unit, TOPAZ64*

The commercial ultrasonic probe from Zetec is a 64-elements linear phased array longitudinal-wave probe with the notation of LM-5MHz. The nominal center frequency is 5 MHz and bandwidth is 74%. Refer to Figure 2.4, each element has a size of 0.5 mm along primary axis and 10 mm in secondary axis (elevation). The kerf between elements are 0.1 mm and the total aperture of the probe is thus 38.3 mm in primary axis and 10 mm in elevation.

Two specific plastic wedges (Figure 4.2), with and without wedge angle denoted as LM-55SW and LM-0LW, respectively, are used in all experiments to protect the surface of the PA probe and to generate angled beam properly. The skew holes on two sides of the wedge facilitate fixation of the probe on mechanized system, and the skew holes on the wedge surfaces are used to fix the PA probe. The angled wedge (LM-55SW) helps the probe generate 55° transverse waves into carbon steel (wave speed = 3230 m/s) without any delay law, or 40° to 70° transverse waves by certain delay laws.



Figure 4.2: *Plastic wedges, LM-55SW (left) and LM-0LW (right)*

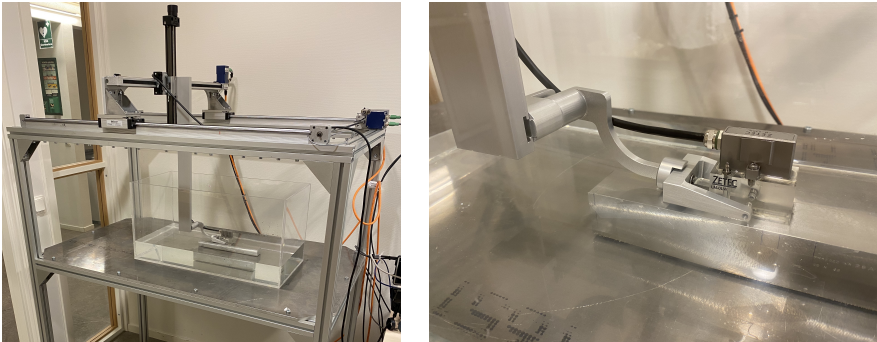


Figure 4.3: *Mechanized gantry system on the experimental platform (left) and probe fixation (right)*

4.2 Mechanized gantry system

The mechanized gantry system is built on the experimental platform (Figure 4.3 (left)) in order to provide a stable inspection condition with high repeatability between experiments. Currently, the system is motor controlled only in the horizontal plane (x-y plane), while the z-axis position is manually adjusted by a guide skew on top. The embedded encoders in the motors can provide z-axis position information within x-y plane, which enable different data presentation probabilities, e.g. B- and C-scan. The water tank containing test specimens is placed on the platform and the wedge with probe is clamped on the z-axis by a spring-loaded fork, as shown in Figure 4.3 (right).

4.3 Test specimens

Three flat surface test specimens (#1, #2, #3) with well-defined artificial defects are used in the validation work. One (#1, aluminum) with 15 SDHs (3 mm in diameter) at depth range from 20 mm to 90 mm in step of 5 mm, drilled through the width of specimen and one of the SDHs can be used as calibration defect. As mentioned in Section 2.4 that FBH could also be treated as the calibration defect in simSUNDT and in experiments, however, due to its manufacturing uncertainties [22], it is not considered in the current work. The second test specimen (#2, stainless steel) has *Electric Discharge Machining (EDM)* notches on the specimen surface, denoted as surface breaking cracks with height of 15 mm, 0.5 mm, 2 mm, 5 mm and 10 mm, ordered as in the specimen. Another test specimen (#3, stainless steel) has 6 SDHs (2 mm in diameter) at depth range from 10 mm to 60 mm in step of 10 mm, drilled through the width of specimen.

The dimension and acoustic properties of these three test specimens are summarized in Table 4.1. The profile sketches of these specimens in length-height plane are shown in Figure 4.4.

Table 4.1: Dimension and acoustic properties of two test specimens

No.	Length (mm)	Height (mm)	Width (mm)	L-wave speed (m/s)	T-wave speed (m/s)
#1	320	100	30	6320	3130
#2	500	35	50	5573	3150
#3	250	65	39	5640	3110

4.4 Design of validation experiments

The configuration and the design of experiments are summarized as followed. Pulse-echo mode is used for the PA probe in the experiments and simulations. The validation of the probe model addresses the comparison of the maximum echo amplitudes on test specimen #1 (**Paper A**), data presentations towards defects on test specimen #2 and #3 (**Paper B**) between the experiments and corresponding simulations. Only one SDH at 50 mm depth in specimen #3 is used. The PA probe performs a continuous one-line scan in a single run over all defects on the scanning surface from one end to the other as in the simulation, and the maximum echo amplitude of each defect, expressed in percentage of screen height, could be retrieved in post-processing. All maximum echo amplitudes of the defects (A_{def}) are normalized to the respective calibration SDH (A_{cal}) and

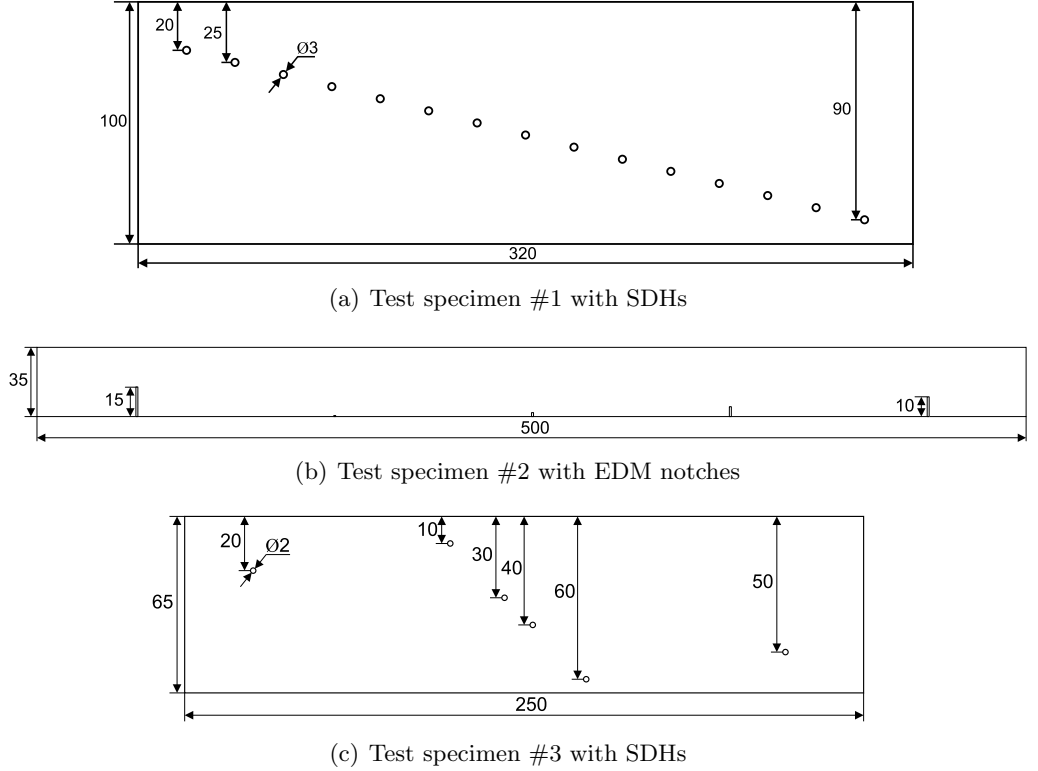


Figure 4.4: *Sketch profiles of the test specimens*

expressed in decibel (dB) using equation (4.1), where G_{def} and G_{cal} is the applied gain to the defect and calibration SDH amplitude, respectively.

$$dB = 20 \log\left(\frac{A_{def}}{A_{cal}}\right) - (G_{def} - G_{cal}) \quad (4.1)$$

Inspections under non-angled and 45° angled conditions with and without focusing effect are considered overall. Only the central 16 elements are active in the unfocused inspection to avoid ghost images, whereas all 64 elements are active to generate proper focusing effect at certain depths of interest. Each set of the experiments is repeated five times by moving the probe continuously in a one-line scan from the same starting point on the scanning surface. All other experimental conditions are unchanged between repetitions to ensure consistency. The experimental results are therefore expressed by the mean value with error bars indicating the variations.

5 Sound field optimization

The sound field optimization aims at optimizing the generated sound field in order to get a maximized echo amplitude (optimization objective) from a defect with certain depth, angle and size. The previous validated PA probe simulation model in simSUNDT is used as a computation kernel and the optimization process is realized using software modeFrontier (**Paper B**). Initially, a well-defined surface breaking crack (as in test specimen #2 in validation work, see Figure 4.4(b)) is studied, which has an angle of 90-degrees towards bottom surface and a size of 10 mm located at the bottom. The decision variables to the optimization problem in the current case are the beam angle and focus distance of the wave beam, which in the end provide a synthetic wave front to get a maximized echo amplitude from the defect of interest.

The optimization algorithm used is heuristic Simplex based on Nelder-Mead method [38], which is used for non-linear single-objective optimization problems. It compares the objective values at simplex vertices and moves the simplex towards the optimal solutions based on different operations, i.e. reflection, expansion, contraction and shrink, see Figure 5.1, according to the evaluation of current objective values. Simplex does not calculate the derivatives and is more robust than gradient-based algorithms, which motivates the choice of this method.

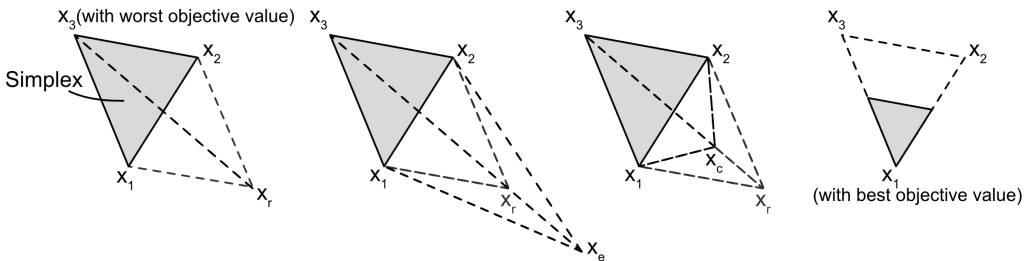


Figure 5.1: *Four operations in Nelder-Mead based Simplex optimization algorithm (left to right): reflection, expansion, contraction and shrink*

6 Summary of appended papers

6.1 Paper A

The newly developed and implemented phased array probe model in simSUNDT for advanced ultrasonic testing is further validated by comparing the simulation results with corresponding experiments, in terms of the maximum echo amplitudes. To master the experimental process and data, a mechanized gantry system and platform was built at the NDT lab of the Scientific Center of NDT (SCeNDT) at Chalmers University of Technology, on which all experiments were performed. Two test specimens with side-drilled holes (SDHs), considered as predefined artificial defects and different materials are involved for validation and practical purposes. Good correlations can be seen from the comparisons and this model is concluded as an acceptable option to the corresponding experimental work. In addition, the relation between depth and the true beam angle is investigated, which is essential to guarantee an accurate inspection. It is also to show the flexibility of parametric studies using a simulation model.

6.2 Paper B

The model validation work further continues in this paper by data presentation comparisons with corresponding experimental results, i.e. A-, B- and C-scans. The defect types considered are side-drilled holes (SDHs) and Electric Discharge Machining (EDM) notches, where the direct echo and corner echo are the respective received signals. The comparisons show generally satisfactory correlations. Upon validation of the model, it is then used in sound field optimization. The optimization addresses searching for a proper combination of main beam angle and focus distance at this stage, so that a maximized echo amplitude towards a certain defect, which has specific characters (size and tilt angle), can be retrieved. The initial methodology of the optimization process is investigated in this paper.

7 Conclusion and future plans

With the increasing demands of quality and integrity assurance of the manufactured components in industries utilizing advanced production technologies, ultrasonic inspection as one of the nondestructive evaluation methods plays important role. To better understand the method and facilitate qualification procedures, corresponding numerical simulation models were developed, which need to be validated thoroughly.

In this thesis, a newly developed and implemented phased array ultrasonic testing (PAUT) model in simSUNDT, an ultrasonic testing simulation software developed at Scientific Center of NDT (SCeNDT) at Chalmers University of Technology, is further validated by comparing the simulated results with corresponding experiments, in terms of maximum echo amplitudes and different types of data presentations. To enable stabilized inspection operations and recording the probe position for data presentation, a mechanized inspection system is built at NDT lab in SCeNDT.

Good correlations are observed from these comparisons in general and the phased array probe model is hereby concluded as validated. It is also noticed from the comparisons when beam focusing effect is involved that the focusing accuracy of a phased array beam decreases with increasing specified depth. In addition, the relation between depth and true beam angle is discovered when analyzing the results, which emphasizes the necessity of calibration before actual inspections.

After the model validation, it can then be used in the sound field optimization work, which at this stage investigates an appropriate combination of main beam angle and focusing distance towards a certain defect, so that the echo amplitude from this defect is maximized.

As the current validation works focus mainly on well-defined artificial defects, i.e. SDHs and cracks, future possibilities thorough the validation could involve more complex AM related defect and material types. Besides, the sound field optimization could proceed on exploring optimized delay law of the phased array technique for better signal processing in the final stage.

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