Article D: 1000-7032 (2007) 01-0001-06

# Influence of Intercalated L i on Electronic Structures and Optical Properties of V<sub>2</sub> O<sub>5</sub>

LIZhi-yang, ZHOU Chang-jie, LNWei, WUQi-hui<sup>\*</sup>, KANG Jun-yong

(Department of Physics and Son iconductor Photonics Research Center, Xiam en University, Xiam en 361005, China)

**Abstract:** The density of states and band structures of  $-V_2O_5$  and Li-intercalated  $V_2O_5$  (Li<sub>x</sub> $V_2O_5$ , x = 0.5 and 1.0) have been studied using a first-principles calculation based on density function theory with the boal density approximation. The results indicate that  $V_2O_5$  is an indirect-gap semiconductor, the intercalation of Li will not change its way of electron transition. While, the intercalation of Li bowers the energy of conduction band, and then narrows the band gap. At the same time, due to the intercalation of Li, the split-off in the conduction band of  $V_2O_5$  disappears because of the split of conduction band. The Fermi level of Li<sub>x</sub> $V_2O_5$  increases dramatically due to the electron transfer from Li 2 s to the  $V_2O_5$  host, which is probably the main reason why the optical band-gap augments with the Li intercalation

Key words:first-principles calculation; $V_2O_5$ ;Li intercalation;electronic structure;optical propertiesCLC number:O471.5PACC:6170Y;7125C;8140Document code:A

## 1 Introduction

Recently, a lot of studies have been focused on the vanadium pentoxide  $(V_2O_5)$ , because it is a kind of very important inorganic material Due to its special crystal and electronic structure, V<sub>2</sub>O<sub>5</sub> has been used widely as cathode material for intercalation batteries<sup>[1]</sup>, oxidation catalyst in industry<sup>[2]</sup>, as</sup></sup> well as gas sensors and electrode for electrochromic devices  $[3^{5}]$ . As using in electrochromic devices,  $V_2O_5$  can be intercalated and de-intercalated by Li and thus could be defined as both a cathodically and an anodically coloring materials  $[6^{-9}]$ . Generally,  $V_2O_5$  is yellow, but the color will change to green and then blue with the intercalation of Li During the de-intercalation of Li, its color will reversibly change back from blue to green and then yellow. Experimental results have showed that the transmission in the near-ultraviolet region changes strongly and the transmission in the near-infrared region has a slight variation with the increase of the amount of intercalated lithium; while the optical absorption edge moves to the shorter wavelength<sup>[6~10]</sup>. It was also found that the absorption tail in the low photonic energy region increases with the intercalation of Li It has been considered that the intercalated Li will donate its 2s electron to the host, and leading to the reduction of partial V<sup>5+</sup> to V<sup>4+[11]</sup>. The reaction can be showed as:

$$V_{2}^{5+}O_{5} + xe + xLi^{+} Li_{x}^{+}V_{x}^{4+}V_{1-x}^{5+}O_{5}$$
(1)

The reduced  $V^{4+}$  may play an important role in increasing the absorption tail in the low photonic energy region<sup>[12]</sup>. So far, the mechanism of change of optical properties of  $V_2O_5$  due to the intercalation/ de-intercalation of L i is still not clear Little work has been done in the study of the influence on the

Received date: 2006-08-25; Revised date: 2006-11-15

Foundation item: Project supported by the Major State Basic Research Development Program of China (001CB610505); the National Nature Science Foundation of China (60376015, 60336020); the Project of Young Talents Innovation of Fujian Province (2005J005)

**B iography:** L I Zhi-yang, male, was bom in 1981, Fujian Province, Master student His work focuses on growth and characterization of oxide materials thin films

E-mail: zhy\_li@xmu edu cn

<sup>\*:</sup> Corresponding Author, E-mail: qihui\_wu@xmu edu cn, Tel: (0592)2184302, Fax: (0592)2187737

optical properties of  $V_2O_5$  with Li intercalation/deintercalation by theoretical calculation It is well known that the band-to-band transition of electron is the most important part of optical absorption in the semiconductors There are two transition ways: direct-gap and indirect-gap transition Close to the absorption edge, the absorption coefficient (*h*) and the energy of incident photon (*h*) can be expressed as,

$$(h) = B(h - E_g)^m / h$$
 (2)

where *B* is transition probability,  $E_g$  is width of band gap. For the direct-gap transition, m = 1/2, and m = 2 for the indirect-gap transition. Therefore, the absorption coefficient has close relationship with the band structures of semiconductors. In present work, we use a first-principles calculation based on density function theory (DFT) with the local density approximation (LDA) to calculate the density of states (DOSs) and band structures of  $-V_2O_5$  and Li-intercalated  $V_2O_5$ . The purpose of this study is to understand, from theoretical view, the influence of intercalation of Li on the optical properties of  $V_2O_5$ .

## 2 Calculation and Crystal Structure

V ienna *ab-initio* Simulation Package (VASP), which has been well tested in solid-state  $physics^{[13^{-15}]}$ , is used in our study. VASP is a complex package for performing *ab-initio* quantum mechanical molecular dynamic simulations with plane wave basis sets and ultra-soft V and erbilt p seudopotentials

 $-V_2O_5$  forms an orthorhombic  $(P_{mmn})$  layer crystal (see Fig 1) with the lattice constants of a = 1.1510, b = 0.3563 and c = 0.4369 nm<sup>[16]</sup>. The layers with normal direction along [001] are









composed of distorter VO<sub>6</sub> octahedral (see Fig 2) sharing edges or corners There are three different oxygen atoms in  $V_2O_5$ : one vanadyl oxygen O (1), one bridging oxygen O (2), and three trip le bonding oxygens O (3). The unit cell in this study contains 14 atoms, four vanadium and ten oxygen atoms

## 3 Results and Discussion

### 3.1 The Electron ic Structures of $-V_2O_5$

The total energy of cell has been calculated with varying lattice constants, and it exhibits a minimum at lattice constants of a = 1. 149 7, b = 0.357 5 and c = 0.435 5 nm. It is only about 0.3% relative error comparing with those lattice constants measured from experimental study<sup>[16]</sup>. The good agreement suggests that this theoretical calculation method is suitable in this study.

The total DOSs, the partial DOSs of V and O, and the energy band structure of  $V_2O_5$  are calculated by using the optimized lattice constants

The total DOSs (Fig 3) show that the electron states mainly distribute in the two energy regions, one is in - 5. 12 ~ 0 eV corresponding to the valence band (VB) and the other is in 1. 61 ~ 4. 03 eV where the conduction band (CB) appears A narrow split-off with width of 0. 605 eV in CB is found, which has also been reported in experimental study<sup>[10]</sup>. Fig 4 shows the partial DOSs of V and O. It is obvious that the O 2p state mostly locates in the VB and the CB is almost constructed by V 3d orbital<sup>[17]</sup>. It is interesting to see that O (3) contributes very little in the CB, and the contribution of



Fig 3 The total DOSs of  $V_2O_5$ .



Fig 4 The DOSs of V 3d and O 2p in V<sub>2</sub>O<sub>5</sub> O (1) vanadyl oxygen, O (2) bridging oxygen, and O (3) trip le bonding oxygen



Fig 5 The band structure of  $V_2O_5$ .

O(1) in the CB locates in higher energy band than that of O(2). These results suggest that different O have different bonding properties with V.

The band structures (Fig 5) show that there is a band gap of 1. 61 eV between the bottom of the CB at (0, 0, 0) point and the top of the VB at U (0,

, ) point, which indicates that the  $V_2O_5$  is an indirect-gap semiconductor. The band gap is smaller than the experimental results, which are usually between 1. 90 and 2. 46 eV. This is reasonable for using LDA<sup>[13]</sup>.

### 3. 2 The Electron ic Structures of L i $_5$ V<sub>2</sub>O<sub>5</sub>

We introduce one L i atom on the line of vanadyl oxygen and vanadium without changing the lattice constants of  $V_2O_5^{[18]}$  as shown in Fig 6 The chemical formulation thus can be written as L  $i_{1.5}V_2O_5$ .



Fig 6 The location of L i in the cell of  $V_2O_5$ .

During the calculation Li is relaxed along *c* axis W hen the total energy of the system reaches a minimum, the location of Li shifts a little bit to vanadyl O ion Then, the DOSs and the band structures of  $L_{i_15}V_2O_5$  have been calculated Fig 7 shows the comparison of the total DOSs of  $L_{i_15}V_2O_5$  and  $V_2O_5$ . It is obvious to notice that the narrow split-off in the CB of  $V_2O_5$  disappears in that of  $L_{i_15}V_2O_5$ . The width of VB region of  $L_{i_15}V_2O_5$  is - 6 17 ~ 0 eV, which is about 1. 02 eV wider than that of  $V_2O_5$ . CB moves to the higher energy compared with that of V 3d in  $V_2O_5$ ; while the DOSs of  $V_B$  3d in CB.



Fig 7 The comparison of the total DOSs of  $V_2O_5$  and  $L\,i_{\!1\,5}V_2O_5$ .

Fig 8 displays the DOSs of V 3d states The locations of  $V_A$  and  $V_B$  are indicated in Fig 6. There are strong differences between the DOSs of  $V_A$  3d and  $V_B$  3d states due to the different influences from intercalated Li The DOSs minimum of  $V_A$  3d in the



Fig 8 The DOSs of V: (a) the V of  $V_2O_5$ ; (b)  $V_A$  and (c)  $V_B$  of  $L_{i_0,5}V_2O_5$  as shown in Fig 6.

CB moves to the higher energy compared with that of V 3d in  $V_2O_5$ ; while the DOSs of  $V_B$  3d in CB shifts to lower energy. The DOSs of  $V_A$  3d in the VB becomes broader and moves to lower energy, at the same time, it decreases in the higher energy region of VB. The CB of  $V_B$  3d does not show obvious change, except for a little bit broadening Fig 9 shows that, compared with these of O in  $V_2O_5$ , the DOSs of the nearest vanadyl oxygen OA 2p also shifts to lower energy and increases in higher energy region in the  $V_B$  due to the effect of intercalated Li Correspondingly, little influence is found for O<sub>B</sub> 2p state, which is farer away from the intercalated Li atom. Experimental data have reported that the intercalated Li will transfer its 2s electron to V ion and then leads to the reduction of  $V^{^{5\,+}}\,$  to  $V^{^{4\,+\,\,[\,11\,]}}.$  The strong effects on the DOSs of  $V_{\scriptscriptstyle A}\,$  3d and  $O_{\scriptscriptstyle A}\,$  2p found in our calculation results somewhat support the experimental data



Fig 9 The DOSs of O (1) 2p: (a) the O (1) of  $V_2O_5$ , (b)  $O_A$  and (c)  $O_B$  of L  $i_{0.5}V_2O_5$  as shown in Fig 6.

Fig 10 shows the band structures and Fermi levels  $(E_{\rm F})$  of V<sub>2</sub>O<sub>5</sub> and Li<sub>0.5</sub>V<sub>2</sub>O<sub>5</sub>. The band-gap

of  $L_{i_1,5} V_2 O_5$  is about 1. 33 eV, which is about 0. 231 eV smaller than that of  $V_2O_5$ . Since the bottom of CB and the top of VB still locate on (0, (0, 0) and U(0, ...) respectively, the properties of electron transition of  $L_{i_1,5} V_2 O_5$  do not change. The band structures also suggest that there is not any inpurity level caused by intercalated Li, but the symmetrical characteristic of system is reduced The energy bands that are degenerated split off move to the low energy region The intercalation of Li lowers the energy of conduction band and narrows the band gap. But why dose the optical absorption edge move to the shorter wave with Li intercalation as reported by experimental measurements<sup>[10]</sup>? During the calculation, we found that the Fermi level of  $L_{i_1,5}V_2O_5$ increases to 1. 85 eV owing to the electron transfer from Li2s to the  $V_2O_5$  host, which already locates in the CB region and is about 0. 52 eV higher than the bottom of CB. As it is known that the energy bands below the Fermi level would be filled up by electrons, so the electron from VB only can transfer to CB above the Fermi level, which practically increases the optical band-gap because of the intercalation of Li This might be the main reason why the absorption edge moves to shorter wave after the intercalation of Liproved by the experimental measurements



Fig 10 The band structures of  $V_2O_5$  and  $Li_{0.5}V_2O_5$ .

#### The Electron ic Structures of L iV<sub>2</sub>O<sub>5</sub> 3.3

Using the same model, we continue to introduce another L i into the L  $i_{0.5}V_2O_5$  cell The location of L i is close to  $O_B$  atom in Fig 6. The chemical formulation is thus  $LN_2O_5$ , its band structure compared with  $V_2O_5$  and  $Li_{0.5}V_2O_5$ , and the  $E_F$  is shown in Fig 11.

4



Fig 11 The band structures of  $Li_{0.5}V_2O_5$  and  $LiV_2O_5$ .

The band-gap of  $LN_2O_5$  is about 0. 24 eV narrower than that of  $V_2O_5$ . The broadening of VB is more evident, but the bands on the bottom of CB degenerate again, and make the band-gap of  $LN_2O_5$  a little larger than that of  $Li_{h,5}V_2O_5$ . More amount of Li intercalation still does not change the positions of the bottom of CB and the top of VB. In addition, the Fermi level increases to 1. 91 eV, which is about 0. 06 eV higher than that of  $L_{i_{0.5}}V_2O_5$  and 0. 52 eV higher than the bottom of CB. Therefore, it makes the optical band-gap ever bigger This result agrees well with that the optical absorption edge move to the shorter wave with the more amount of Li intercalation

### 4 Conclusion

By using DFT calculation with LDA, the equilibrium lattice constants, DOSs and energy band structures of  $V_2O_5$ ,  $L\dot{h}_5V_2O_5$  and  $LN_2O_5$  have been investigated We found that the electron transfer property of  $V_2O_5$ , which is indirect band gap, does not change with the increase intercalation amount of Li, but the Fermi level of host increases with the increase amount of intercalated Li That might be the main reason for the optical absorption edge of  $L\dot{h}_5V_2O_5$  and  $LN_2O_5$  moving to the shorter wave region, as the electron can only transfer from VB to CB above the Fermi level

### **References:**

- [1] Ivanova T, Harizanova A, Surtchev M. Formation and investigation of sol-gel TD<sub>2</sub>-V<sub>2</sub>O<sub>5</sub> system [J]. *Mater Lett*, 2002, 55 (5): 327-333.
- [2] Grzybowska-Swierkosz B. Vanadia-titania catalysts for oxidation of o-xylene and other hydrocarbons [J]. Appl Catal A, 1997, **157**(1): 263-310.
- [3] Bondarenka V, Grebinskij S Thin films of poly-vanadium-molybdenum acid as starting materials for humidity sensors [J]. Sensors and Actuators B, 1995, **28**(3): 227-231.
- [4] TongM S, Dai G R, Gao D S, et al The electronic properties and uses of V<sub>2</sub>O<sub>5</sub> thin films [J]. Mater Rev (材料导报), 2000, 14(10): 36-38 (in Chinese).
- [5] Lee S H, Hyeonsik M C, Maeng J S, et al Raman spectroscopic studies of amorphous vanadium oxide thin films [J]. Solid State Ionics, 2003, 165(1-4):111-116
- [6] Benmoussa M, Outzourhit A, Bennouna A, et al Electrochromism in sputtered V<sub>2</sub>O<sub>5</sub> thin films: structural and optical studies [J]. Thin Solid Films, 2002, 405 (1-2): 11-16.
- [7] Ottaviano L, Pennisi A, Simone F, et al RF sputtered electrochrom ic V<sub>2</sub>O<sub>5</sub> films [J]. Opt Mater, 2004, 27 (2): 307-313.
- [8] Yoo S J, L in J W, Sung Y E Inproved electrochromic devices with an inorganic solid electrolyte protective layer [J]. Solar Energy Materials & Solar Cells, 2006, 90 (4): 477-484.
- [9] Cesar O A, Luis O S B. Optical and electrochemical properties of V<sub>2</sub>O<sub>5</sub> Ta sol-gel thin films [J]. Solar Energy Materials & Solar Cells, 2006, 90 (4): 444-451.
- [10] Wu GM, Du K F, Xia C S, et al Optical absorption edge evolution of vanadium pentoxide films during lithium intercalation [J]. Thin Solid Films, 2005, 485 (1-2): 284-289.
- [11] WuQH, Thissen A, Jaegermann W. Photoelectron spectroscopic study of Li intercalation into V<sub>2</sub>O<sub>5</sub> thin films [J]. Surf Sci, 2005, 578 (1-3): 203-212

- [12] Sayede A D, Mathieu C, Khelifa B. Understanding the atomic force microscopy image of the V<sub>2</sub>O<sub>5</sub> and Li<sub>0.03</sub>V<sub>2</sub>O<sub>5</sub> (001) surface using *ab-initio* calculations [J]. *Catalysis Today*, 2006, **113** (3-4): 263-269.
- [13] Kresse G, Furthmuller J. Efficiency of *ab-initio* total energy calculations for metals and semiconductors using a plane-wave basis set [J]. *Comput Mater Sci*, 1996, 6(1):15-50.
- [14] Xu Qunhe, Kang Junyong Structure properties of Li-related defects in ZnO [J]. Chin J. Luon in (发光学报), 2006, 27 (4): 509-513 (in English).
- [15] Zhang Hua, Feng Xia, Kang Junyong Electronic structures of Zn<sub>2</sub> SiO<sub>4</sub> [J]. Chin J. Lum in (发光学报), 2006, 27 (5):750-754 (in English).
- [16] Wu Q H. Photoelectron spectroscopy of intercalation phases: Na and Li in  $V_2O_5$  thin films and  $LMn_2O_4$  [D]. *Ph D. D issertation*, Damstadt, Germany, 2003.
- [17] Wu Q H, Thissen A, Jaegermann W, et al Resonant photoemission spectroscopy study of electronic structure of V<sub>2</sub>O<sub>5</sub>
  [J]. Chan. Phys Lett, 2006, 430 (4-6): 309-313.
- [18] Zhou J, Chen W, Xu Q, et al The study on the position of Li<sup>+</sup> and electronic structure in PBO /V<sub>2</sub>O<sub>5</sub> composite film [J].
  J. Molecular Sci (分子科学学报), 2001, 17(2):71-76 (in Chinese).

# $Li嵌入对 V_2 O_5$ 电子结构及光学性质的影响

# 李志阳,周昌杰,林 伟,吴启辉\*,康俊勇

(厦门大学 物理系 /半导体光子学研究中心 , 福建 厦门 361005)

**摘要**:采用第一性原理局域密度近似法计算了  $V_2O_5$ 的电子态密度和能带结构以及 Li嵌入后对其电子结构 和光学性质的影响。计算结果表明, $V_2O_5$ 是间接带隙半导体,Li的嵌入并没有改变其电子的跃迁方式。但 Li 的嵌入使得  $V_2O_5$ 导带能量下移,禁带宽度减小,导带中原有的劈裂被分裂的能级填满;同时致使价带出现展 宽。电子态密度计算结果表明 Li的嵌入对临近的 O和 V的电子结构有较大的影响。Li 2s电子的注入提高 了  $V_2O_5$ 的费米能级并导致其进入导带。由于价带中的电子只能跃迁到费米能级以上的导带空能级,这致使 体系实际的光学带隙增大。同时随着 Li注入量的进一步增加,价带的展宽更为明显,费米能级亦呈升高的趋 势,使得光学带隙随着 Li注入量的增加而增大。

关 键 词:第一性原理; V<sub>2</sub>O<sub>5</sub>; Li嵌入; 电子结构; 光学性质
 中图分类号: O471.5
 PACC: 6170Y; 7125C; 8140
 文献标识码: A
 文章 编 号: 1000-7032(2007)01-0001-06

收稿日期: 2006-08-25; 修订日期: 2006-11-15

基金项目:国家重点基础研究发展规划(001CB610505);国家自然科学基金(60376015,60336020);福建省青年人才创新基金 (2005J005)资助项目

作者简介: 李志阳 (1981 - ), 男, 福建人, 硕士研究生, 主要从事氧化物薄膜材料的生长和表征的研究。 E-mail: \_zhy\_li@ xmu edu cn

<sup>\*:</sup>通讯联系人; E-mail: qihui\_wu@xmu edu cn, Tel: (0592)2184302, Fax: (0592)2187737