

# Radio observations of evaporating objects in the Cygnus OB2 region

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## ABSTRACT

We present observations of the Cygnus OB2 region obtained with the Giant Metrewave Radio Telescope (GMRT) at the frequencies of 325 MHz and 610 MHz. In this contribution we focus on the study of proplyd-like objects (also known as free-floating Evaporating Gas Globules or frEGGs) that typically show an extended cometary morphology. We identify eight objects previously studied at other wavelengths and derive their physical properties by obtaining their optical depth at radio-wavelengths. Using their geometry and the photoionization rate needed to produce their radio-continuum emission, we find that these sources are possibly ionized by a contribution of the stars Cyg OB2 #9 and Cyg OB2 #22. Spectral index maps of the eight frEGGs were constructed, showing a flat spectrum in radio frequencies in general. We interpret these as produced by optically-thin ionized gas, although it is possible that a combination of thermal emission, not necessarily optically thin, produced by a diffuse gas component and the instrument response (which detects more diffuse emission at low frequencies) can artificially generate negative spectral indices. In particular, for the case of the Tadpole we suggest that the observed emission is not of non-thermal origin despite the presence of regions with negative spectral indices in our maps.

**Key words.** ISM: individual objects: Cygnus OB2 – Radio continuum: stars – star: formation – stars: protostar

## 1. Introduction

The term protoplanetary disk (or, in short, proplyd) was introduced by O'dell et al. (1993) to describe circumstellar disks around young low-mass stars, with the size of our own planetary system, revealed by Hubble Space Telescope images of the Orion Nebula (M42). These proplyds consist of a rotating neutral accreting disk with its surface ionized by an external source, usually a hot massive star. Hence, proplyds are usually found inside HII regions.

Photoevaporation of a circumstellar disk due to nearby young massive stars is an efficient mechanism for molecule and atom destruction, resulting in an important mass loss and producing ultra compact HII regions (Hollenbach et al. 1994). Photoevaporation is produced either by EUV photons ( $h\nu > 13.6$  eV) that ionize the gas and rise its temperature, or by UVL photons ( $6$  eV  $< h\nu < 13.6$  eV) that dissociate molecules and also heat the gas. In 1998, Johnstone et al. presented a model to explain the photoevaporation of circumstellar disks by an external source of ultraviolet radiation and applied it to the Orion proplyds. Their model regards the sound speed in the ionized gas and the escape velocity of the gas from the protostellar gravitational field to find out the shape of the heated ionized material from the disk surface, which finally forms a thermal disk wind with a cometary structure.

Proplyds have been observed in Orion at different wavelengths (O'dell et al. 1993; Henney & Arthur 1998; Graham et al. 2002; Bally et al. 2005; Ricci et al. 2008; Miotello et al. 2012; Mann et al. 2014; Kim et al. 2016; Eisner et al. 2016), as well as in other molecular clouds, such as the Carina Neb-

ula (Smith et al. 2003), NGC 3603 (Brandner et al. 1999; Mücke et al. 2002), the Lagoon Nebula (Stecklum et al. 1998; Masqué et al. 2014), NGC 2024 (Stapelfeldt 2002), Sgr A\* (Yusef-Zadeh et al. 2015), Cygnus OB2 (Wright et al. 2012; Sahai et al. 2012a; Guarcello et al. 2014; Schneider et al. 2016) for instance. Their characteristic sizes range from 40 AU to 1000 AU (Bally et al. 1998; Henney & O'Dell 1999; Sahai et al. 2012a). The masses of the ionized disks are similar to those measured in the non-ionized protostellar disks in Taurus and Ophiuchus molecular clouds (typical masses ranging between 0.003 – 0.07  $M_{\odot}$ , Andrews & Williams 2007; Mann & Williams 2010), with the only exception of the disks within 0.3 pc to  $\theta^1$  Ori C that show a slightly lighter mass distribution.

In the Cygnus OB2 region, Sahai et al. (2012a) defined a new kind of objects called frEGGs (free-floating Evaporating Gas Globules). Like proplyds, these objects are also externally ionized and present extended cometary morphologies. The main difference is that proplyds harbor protostars in more evolved stages with protoplanetary disks already formed, whereas frEGGs are found to have larger molecular masses, indicating an earlier evolutionary stage in which the protostar is still undergoing strong accretion. The frEGGs are like the Evaporating Gas Globules (EGGs) defined by Hester et al. (1996), but found in isolation, detached from their parental molecular clouds. In addition, frEGGs are about 10 times larger than proplyds (frEGGs are typically greater than 20000 AU), so they can contain not just one but several young protostars with their circumstellar disks (Sahai et al. 2012b). FrEGGs are also more massive than Orion proplyds with total molecular gas masses exceeding 1–2  $M_{\odot}$  and contain molecular material inside the photoionized region. They

are usually at further distances from their ionizing source than those in Orion. Previous radio observations on frEGGs toward NGC 3603 and G5.97–1.17 (Mücke et al. 2002; Masqué et al. 2014) found that some of these objects present negative spectral indices, associated with non-thermal emission. The non-thermal radiation from these ionized sources implies the presence of a population of relativistic electrons. Such particles can either be locally accelerated in the source or come from nearby cosmic-ray sources. In the case of G5.97–1.17, Masqué et al. (2014) suggest that electrons are accelerated at the shock produced by the photoevaporating flow. In this contribution, we present 325 and 610 MHz GMRT continuum observations toward the Cygnus OB2 region and report on the characteristics of the population of externally ionized proplyd-like objects also known as frEGGs. The paper is structured as follows. We review previous studies on these objects located in Cygnus in Section 2. In Section 3 we describe how the interferometric observations were taken and calibrated and the process for image production. We present the results in Section 4. In Section 5 we interpret the results and finally we summarize our main findings in Section 6.

## 2. Proplyds and frEGGs in Cygnus

Cygnus is a large, northern-sky region of active star formation at an average distance of 2.5 kpc (Reipurth & Schneider 2008), which hosts nine stellar associations and a dozen of bright clusters. Nearly at the center of this complex region lies the Cygnus OB2 association, probably one of the most massive associations in the Galaxy. The distance to the OB2 association is 1.4 kpc (Rygl et al. 2012). Cygnus OB2 contains pulsars (Bednarek 2003; Camilo et al. 2009), supernova remnants (Kimura et al. 2013; Boubert et al. 2017) and binary stellar systems (De Becker et al. 2004; Lyne et al. 2015), among other non-thermal radio-emitting sources. Recently, using IR and H $\alpha$  observations, a sample of ten externally-ionized objects has also been identified toward this region (Wright et al. 2012, hereafter W12). Most of them are tadpole-shaped with a bright ionization front head roughly pointing to the central cluster of massive stars, and a tail extending in the opposite direction.

A more recent study with Herschel FIR observations at 70  $\mu$ m in the Cygnus region (Schneider et al. 2016) classified several objects ionized in their peripheries by external UV radiation as pillars, globules, evaporation gas globules and proplyd-like objects, measuring densities and temperatures for all of them. The EGGs and proplyd-like objects are 0.1–0.2 pc in size, with masses of 10–30  $M_{\odot}$ , densities of  $2.2\text{--}15.0 \times 10^4 \text{ cm}^{-3}$  and average dust temperatures of 17 K. They also show a clear orientation toward the center of the association Cygnus OB2.

Perhaps the most studied of these proplyd-like objects is IRAS 20324+4057, dubbed the Tadpole (Sahai et al. 2012a), toward which there is a battery of observations at different wavelengths (W12; Sahai et al. 2012a; Guarcello et al. 2014; Schneider et al. 2016). IRAS 20324+4057 is located at  $RA = 20:34:13.27$ ,  $Dec = 41:08:13.8$  [J2000], about 15' (6 pc at a distance of 1.4 kpc) from the center of the Cygnus OB2 association. It consists of an extended cometary nebula ( $7.7 \times 10^4$  AU) oriented East-West, brighter and well defined in the northern edge and with no apparent axial symmetry, as seen in optical wavelengths (Sahai et al. 2012a). Its morphology and orientation suggest a structure formed by the dynamic pressure of wind passing from a distant source or sources located west of the object, presumably a stellar wind from Cyg OB2 #8. Sahai et al. (2012b) observed the Tadpole using the VLA at 22 GHz and 8.5 GHz. They derived a negative spectral index suggesting the presence

of non-thermal emission. The confirmation of non-thermal radiation associated with frEGGs would suggest that these objects are possible accelerators of cosmic-rays, and perhaps even putative gamma-ray sources. Hence, frEGGs could become new laboratories for the investigation of particle acceleration processes in slow ( $< 10^3 \text{ km s}^{-1}$ ) flows, and possibly also as counterparts of unidentified gamma-ray sources. Because of their non-symmetries, large sizes and distances between the head and tail, these objects are different from those detected in the Orion Nebula.

## 3. GMRT observations

The field of IRAS 20324+4057 was observed with the Giant Metrewave Radio Telescope (GMRT)<sup>1</sup> in Nov 2013 at the frequency bands of 325 MHz and 610 MHz, during 6 h and 3 h, respectively. The phase center was set at  $RA = 20:32:50$ ,  $Dec = 41:16:50$  (J2000), near the center of the Cyg OB2 association. The source 3C48 was used as the flux calibrator, and 2038+513 and 2052+365 were monitored as phase calibrators at 325 and 610 MHz respectively.

The data reduction and imaging were performed with the Astronomical Imaging Processing System (AIPS, Greisen 2003) following standard procedures; see for instance Marcote et al. (2015). Also as in Marcote et al. (2015), in addition, we applied a factor for the system temperatures at each frequency bands, to correct for the emission of the Galaxy. The factors,  $2.2 \pm 0.22$  for the 325 MHz band and  $1.3 \pm 0.09$  for the 610 MHz band, were derived from data taken in Jan 2018, exclusively to this purpose (see Marcote et al. 2015, Appendix A, for the characteristics of the corrective process).

We obtained a 325 MHz image of the field with a synthesized beam of  $7.81'' \times 6.60''$ , an rms of  $0.2 \text{ mJy beam}^{-1}$ , and an intensity peak of  $0.270 \text{ Jy beam}^{-1}$ . The corresponding image at 610 MHz was built with a synthesized beam of  $7.60'' \times 5.97''$ , resulting in an rms of  $0.2 \text{ mJy beam}^{-1}$ , and an intensity peak of  $0.201 \text{ Jy beam}^{-1}$ .

The imaging step produced a phase shift -specifically during the self-calibration stage- between the intensity peak at both frequencies. This was corrected by selecting several random point sources (radio galaxies preferentially) across all the field of view, determining the coordinates of the intensity peak at both frequencies, computing the average offset between them, and applying this positional shift to the 610 MHz image ( $RA_{\text{offset}} = 3.19 \pm 1.39''$ ,  $Dec_{\text{offset}} = 0.34 \pm 0.99''$ ). This step was crucial before producing spectral index images, because combining the offsetted images produced a gradual spectral index gradient with the same orientation as the positional shift in all the sources detected in the field of view.

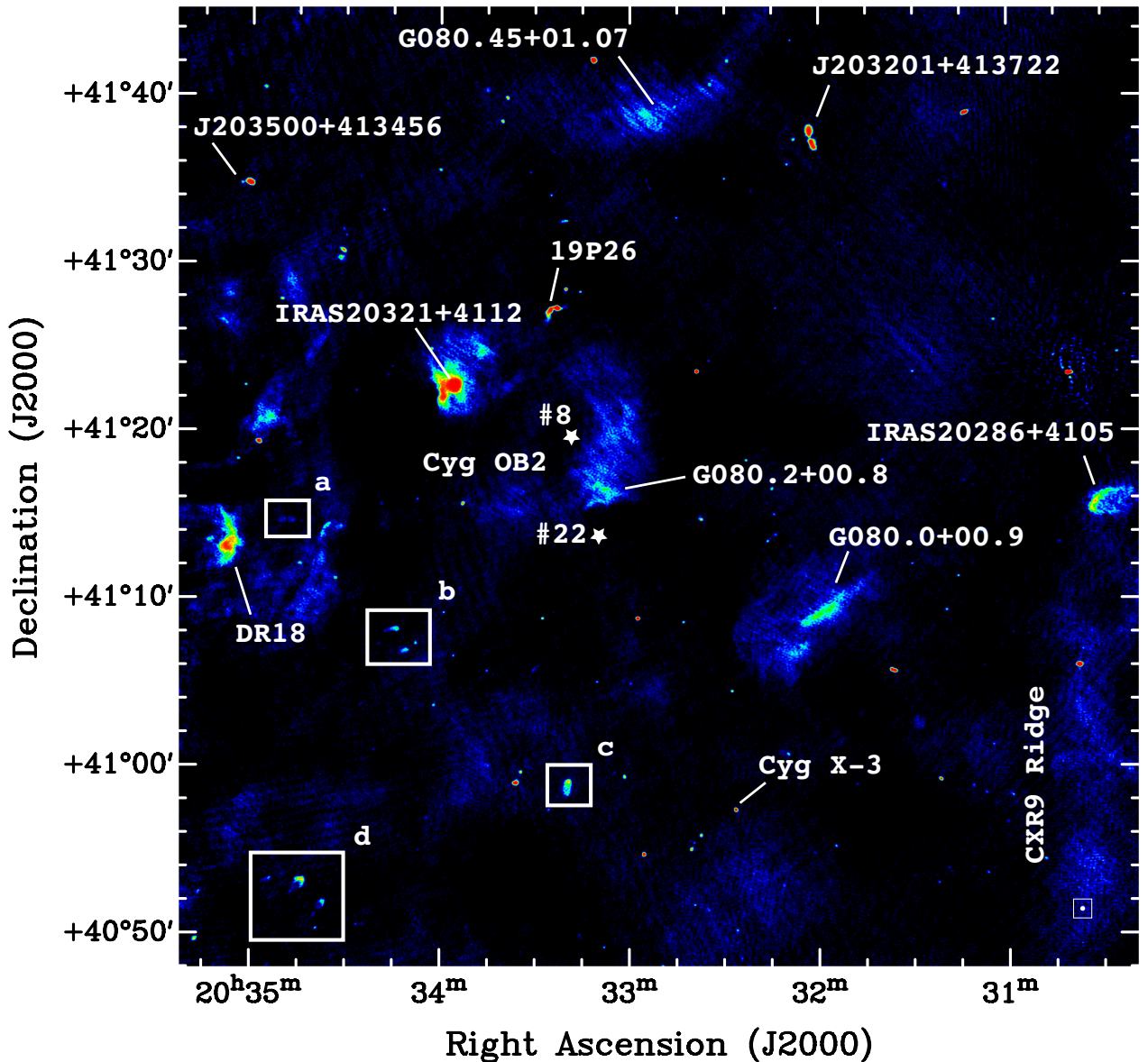
To generate the maps of spectral indices in MIRIAD, we convolved and regridded the 610 MHz image to exactly the same grid as in the 325 MHz one (thus the pixel size in both maps is the same,  $1.5'' \times 1.5''$ ).

## 4. Results

### 4.1. Observed parameters

The Cygnus OB2 image at 325 MHz is shown in Figure 1. The primary beam of the GMRT observations includes the Cygnus OB2 association itself, prominent HII regions spreading several arcminutes, such as DR18 (Comerón & Torra 1999) and

<sup>1</sup> www.gmrt.ncra.tifr.res.in

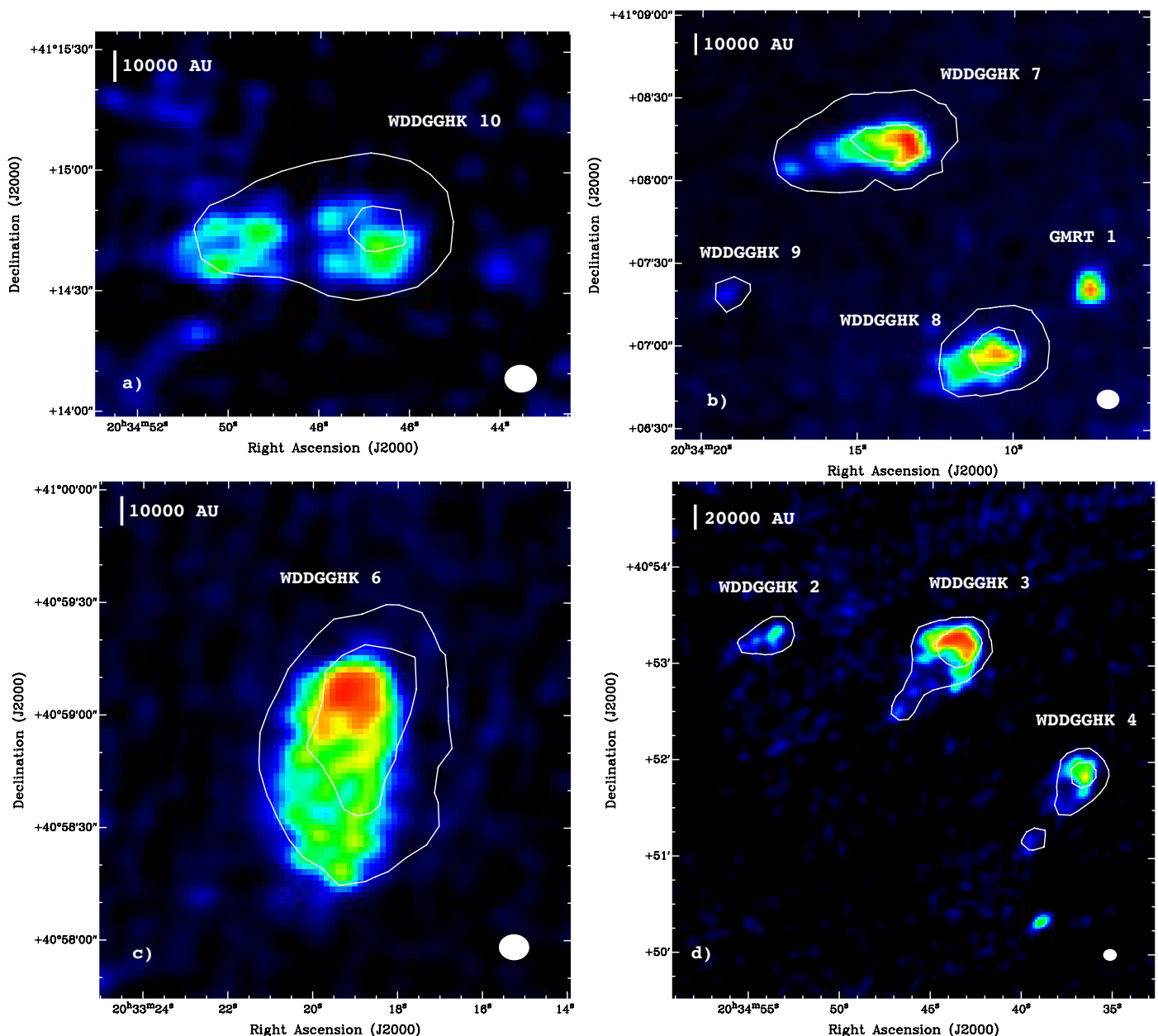


**Fig. 1.** Cygnus OB2 image at 325 MHz; the synthesized beam is  $7.81'' \times 6.60''$ ,  $P.A. = 88.6$  degrees, the rms is  $0.2 \text{ mJy beam}^{-1}$ , and the intensity peak is  $0.270 \text{ Jy beam}^{-1}$ . The boxes in the figures show the regions where the proplyd sources studied here lie: source WDDGGHK 10 in box a, WDDGGHK 7, 8, 9 in box b, WDDGGHK 6 in box c, WDDGGHK 2, 3, 4 in box d. The white stars mark the position of putative ionizing agents.

080.45+01.07 (Marti et al. 2007), radio galaxies like the double-sided NVSS J203201+413722 (Butt et al. 2006; Marti et al. 2007) and many others unresolved within the  $\sim 7''$  angular resolution, and outflows such as IRAS 20286+4105. In addition, this radio image contains most of the photoionized globules or frEGGs previously found by W12. The analysis of these frEGG sources is the focus of the present work. Only objects WDDGGHK 1 and WDDGGHK 5 from [W12]’s list lie outside of the GMRT field of view. Objects WDDGGHK 7, WDDGGHK 8 and WDDGGHK 9 were studied by Sahai et al. (2012a) as well, who called them objects A (or Tadpole), B (or Goldfish) and C. Close to these three sources there is another radio source with negative spectral index that we named GMRT 1 and we could not classify nor discard as another frEGG. Following W12, the O type stars in the region are distributed over a large area and hence, although the center of the association is probably between Cyg OB2 #8 and Cyg OB2 #22 (indicated by white stars in Fig. 1), other luminous stars are probably responsible for the ioniza-

tion of the frEGGs. Figure 2 presents four zoom maps corresponding to these locations, where the color scale indicates the flux at 325 MHz and the white contours indicate the Herschel (Pilbratt et al. 2010) infrared emission, very similar with the radio emission distribution, which shows the cometary morphology of the frEGGs.

Table 1 summarizes the positions, fluxes at 325 MHz and 610 MHz, spectral indices between both frequencies and sizes of each frEGG in the GMRT field of view. Note that sources WDDGGHK 2, WDDGGHK 3 and WDDGGHK 4 lie outside of the primary beam at 610 GHz and the corresponding fluxes are thus missing. We use a  $3\sigma$  threshold to measure the integrated fluxes and estimate linear sizes, which range from about  $10''$  to  $60''$  (15 000 AU to 90 000 AU at the assumed distance of 1.4 kpc), similar to those obtained from optical images (W12). The average spectral indices ( $\alpha = \log(S_2/S_1) / \log(\nu_2/\nu_1)$ ) are mostly around zero.



**Fig. 2.** Zoom in to the four white boxes marked in Fig. 1, 325 MHz emission in color scale, and IR contours of 470 and 1070 mJy beam<sup>-1</sup>.

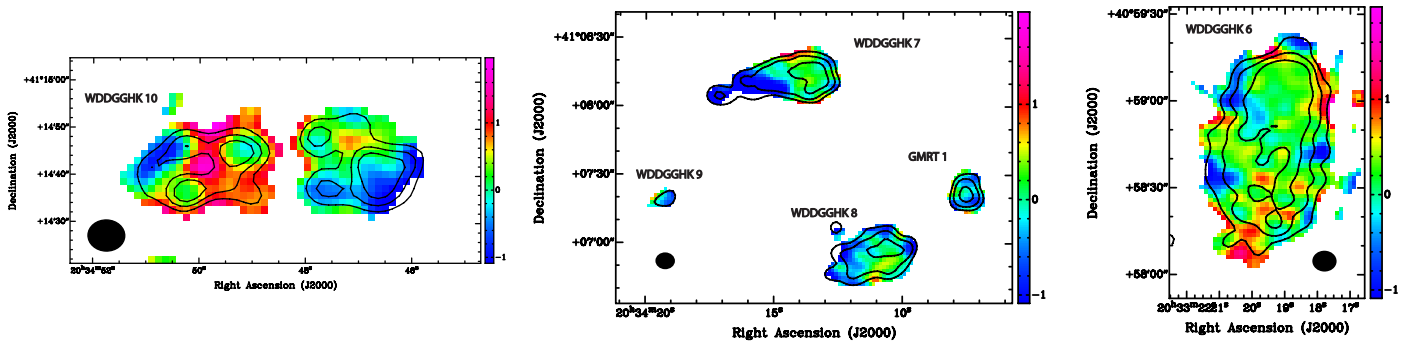
We build spectral index distribution images (Fig. 3) implementing in MIRIAD its mathematical expression. The smaller beam image was convolved with a beam corresponding to the one with the larger beam, forcing them to have the same synthesized beam. Some of these maps show differences on the spectral index across the sources, but we comment more on these in Section 4.2. We determine the average optical depth ( $\tau$ ) of the frEGGs using the following equation, considering spherical symmetry and homogeneity and the Rayleigh-Jeans approach (Altenhoff et al. 1960),

$$\tau = 0.08235(T_e/K)^{-1.35}(\nu/\text{GHz})^{-2.1}(E.M./\text{pc cm}^{-6}) \quad (1)$$

where  $T_e$  is the electron temperature,  $\nu$  the observing frequency, and  $E.M.$  the emission measure. In general, the emission is optically thin.

Table 2 presents, besides  $\tau$  values, the physical properties that we obtain analyzing the morphology of the sources (length and orientation), derived from Mezger & Henderson (1967) and

Ainsworth et al. (2016) expressions for HII regions, evaluated at the frequency of 325 MHz. In doing this, as a first order approximation we consider a model of free-free emission in an optically thin homogeneous sphere. We adopted the values  $T_e = 10^4$  K and  $d = 1.4$  kpc (Rygl et al. 2012) to estimate the electron density, mass of ionized gas, emission measure and rate of Lyman photons needed to ionize each source, the latter with  $(N_i/s^{-1}) = 7.6 \times 10^{43}(S_\nu/\text{mJy})(\nu/\text{GHz})^{0.1}(T_e/10^4\text{K})^{-0.45}(D/\text{kpc})^2$  (Garay et al. 1987). The densities are typical of photoionized regions ( $10^2 \text{ cm}^{-3}$ ) with ionized gas masses of  $0.002 M_\odot$ – $0.2 M_\odot$ . The photon rates range between  $10^{43} \text{ s}^{-1}$ – $10^{45} \text{ s}^{-1}$ , which can be produced by early B-type stars. We give more insight on the physical properties in Section 4.2, where we review the main characteristics of each frEGG in our sample.



**Fig. 3.** Spectral index distribution. *Left:* source WDDGGHK 10 with 325 MHz contours of 0.5, 0.7 and 0.9 mJy beam<sup>-1</sup>. *Center:* sources WDDGGHK 7, 8, 9, and GMRT 1, with 325 MHz contours of 0.5, 1 and 2 mJy beam<sup>-1</sup>. *Right:* source WDDGGHK 6 with 325 MHz contours of 0.5, 1 and 2 mJy beam<sup>-1</sup>. Since the error of the spectral index is largest at the edges of the sources because the emission is weaker, extreme values of the spectral index can be generated in those areas.

**Table 1.** Observed parameters for the frEGGs on Cygnus OB2.

ID	RA(J2000) (h,m,s)	Dec(J2000) (°,′,″)	$S_{325}$ (mJy)	$S_{610}$ (mJy)	$\alpha$	$\theta_{325}$ (″)	$P.A._{325}$ (°)	$\theta_{610}$ (″)	$P.A._{610}$ (°)
WDDGGHK2	20:34:53.55	+40:53:20.34	5.5±0.4	—	—	27	132	—	—
WDDGGHK3	20:34:43.49	+40:53:17.98	54±1	—	—	57	138	—	—
WDDGGHK4	20:34:36.577	+40:51:49.90	15.2±0.6	—	—	36	141	—	—
WDDGGHK6	20:33:19.276	+40:59:06.26	68±1	79±1	0.22±0.07	35	177	33	176
WDDGGHK7	20:34:13.385	+41:08:15.11	19.8±0.6	17.8±0.8	-0.2±0.2	28	108	25	98
WDDGGHK8	20:34:10.438	+41:06:57.24	15.2±0.4	13.7±0.6	-0.2±0.2	21	112	22	118
WDDGGHK9	20:34:19.207	+41:07:19.34	0.4±0.1	0.4±0.1	0±1	11	127	8	119
WDDGGHK10	20:34:46.639	+41:14:41.82	9.9±0.2	12.2±0.3	0.3±0.1	32	90	28	89

**Notes.** (1) Geometric sizes,  $\theta$ , obtained from  $\sqrt{ab}$ , where  $a$  and  $b$  are major and minor axis respectively. (2) Typical uncertainties in the measured sizes are 1.5″, and in  $P.A.$  are 8°.

**Table 2.** Physical properties for the frEGGs of Cygnus OB2.

ID	$n_e$ (cm <sup>-3</sup> )	$M_{\text{ion}}$ ( $M_{\odot}$ )	$E.M.$ (pc cm <sup>-6</sup> )	$\tau$	$N_{\text{Lyman}}$ (s <sup>-1</sup> )
WDDGGHK2	100±10	0.025±0.003	2500±700	0.009±0.003	$7.32 \times 10^{44}$
WDDGGHK3	98±5	0.24±0.01	55500±700	0.019±0.002	$7.18 \times 10^{45}$
WDDGGHK4	103±8	0.063±0.005	3800±800	0.013±0.003	$1.99 \times 10^{45}$
WDDGGHK6	230±20	0.13±0.09	18000±3000	0.06±0.01	$9.05 \times 10^{45}$
WDDGGHK7	170±20	0.050±0.005	8000±2000	0.029±0.007	$2.66 \times 10^{45}$
WDDGGHK8	230±30	0.03±0.03	11000±3000	0.04±0.01	$1.99 \times 10^{45}$
WDDGGHK9	100±30	0.002±0.001	1100±800	0.004±0.003	$5.22 \times 10^{43}$
WDDGGHK10	100±8	0.042±0.003	3200±600	0.011±0.002	$1.33 \times 10^{45}$

**Notes.** The values of  $n_e$ ,  $M_{\text{ion}}$ ,  $E.M.$  were determined with equations (5), (6) and (7) from Ainsworth et al. 2016.

#### 4.2. Individual sources

The observed frEGGs are distributed south of the OB2 association and are grouped in four distinct regions (boxes  $a$ ,  $b$ ,  $c$  and  $d$  in Fig. 1). There are two isolated objects (WDDGGHK 6 and WDDGGHK 10) and two groups of three objects each (boxes  $b$  and  $d$ , respectively in Figs. 1 and 2). All the sources show an elongated morphology and spectral indices around zero. The morphology in radio wavelengths matches perfectly the emission detected by Herschel at 70  $\mu\text{m}$ . The electron density and the orientation of the sources is similar within the members of the two mentioned groups. We summarize the main findings of each object below.

**WDDGGHK 2:** This elongated (37000 AU) source, also known as IPHASX J2034533+405321 is located at the south-east edge of the 325 MHz primary beam (box  $d$  in Figs. 1 and 2),

lying outside of the 610 MHz primary beam. It is grouped with objects WDDGGHK 3 and WDDGGHK 4, with which shares a northwest-southeast orientation and a similar electron density around 90-100 cm<sup>-3</sup>. As with these other two objects, the tail of WDDGGHK 2 seems to comprise two lateral threads. It has been also detected in infrared, submillimeter and optical wavelengths (Wright et al. 2012; Schneider et al. 2016).

**WDDGGHK 3:** It is located about 0.5 pc west of WDDGGHK 2 and 0.5 pc east of WDDGGHK 4, and is one of the strongest sources of our sample at 325 MHz. This object presents a clear cometary shape with a prominent head, showing a bow-shock structure (also detected in optical wavelengths, see Comerón et al. 2002) and a faint tail due southeast. Its 80000 AU length makes it the largest frEGG in our sample. Additionally, it has the largest content of ionized gas (0.236  $M_{\odot}$ ), needing

an ionizing photon rate of  $10^{45.8} \text{ s}^{-1}$ . WDDGGHK 3 contains a protostar probably associated with IRAS 20328+4042, a  $100 L_{\odot}$  young stellar object embedded in a massive gaseous envelope or disk (Comerón et al. 2002; Roy et al. 2011).

**WDDGGHK 4:** This elongated source is spatially resolved at 325 MHz (50 000 AU in length). It splits into two clear parts: a head with a bowshock shape pointing northwest ( $P.A. = 141^{\circ}$ ) and a frayed tail, comprised of faint shreds, one of them disconnected from the rest (Fig. 2). Previous FIR observations between  $500 \mu\text{m}$  and  $60 \mu\text{m}$  toward the head of this object have identified a dusty core, BLAST C62, possibly harboring a stellar nursery (Roy et al. 2011). About  $1.8'$  south of WDDGGHK 4's head there is an unresolved source with a flux density of  $2.7 \text{ mJy}$  (at 325 MHz) that we do not include in our analysis of frEGGs, although it could be related with this group of photoionized sources because of its proximity and its similar orientation, maybe a tail emanating behind the bow shock.

**WDDGGHK 6:** Box *c* in Figure 1 contains this isolated source (see also, Figs. 2 and 3) with a north-south orientation ( $P.A. = 177^{\circ}$ ). As many of these objects, it has a very strong arc-shaped head, followed by a frayed tail with two or three shreds, also observed in  $H\alpha$  and  $8 \mu\text{m}$  images (W12). It is one of the densest sources as well ( $230 \text{ cm}^{-3}$ )<sup>2</sup> It needs an ionization photon rate of  $10^{45.9} \text{ s}^{-1}$ . Roy et al. (2011) associated this photoionized region with IRAS20315+4046, whose emission may be originated by a  $400 L_{\odot}$  protostar(s).

**WDDGGHK 7:** Dubbed the *Tadpole* because of its morphological resemblance with the aquatic animal, it has been the focus of several works (Wright et al. 2012; Sahai et al. 2012a; Guarcello et al. 2014; Schneider et al. 2016). It is  $\sim 12'$  south-west of DR18 (box *b* in Fig. 1) and it consists of the characteristic bowshock head, followed by a fainter tail (Figs. 2 and 3). The object is oriented almost east-west ( $P.A. = 108^{\circ}$ ), similarly to its frEGGs neighbors WDDGGHK 8 and possibly WDDGGHK 9. From the radio emission detected with the GMRT we derive an average spectral index of  $-0.2 \pm 0.2$ . However, the spatial distribution of the spectral index is not homogeneous, varying from positive at the head to negative at the tail. In section 4.3 we analyze its spectral energy distribution (SED) in detail.

**WDDGGHK 8:** This source, also called the *Goldfish*, is about  $1'$  south of the Tadpole. It is as dense as the Tadpole ( $233 \text{ cm}^{-3}$ ) but with a slightly smaller ionized mass. We find an average spectral index of  $-0.16 \pm 0.19$  which, in this case, is homogeneously distributed throughout the source (i.e., no apparent variations), that within the errors it is consistent with emission from an optically thin plasma.

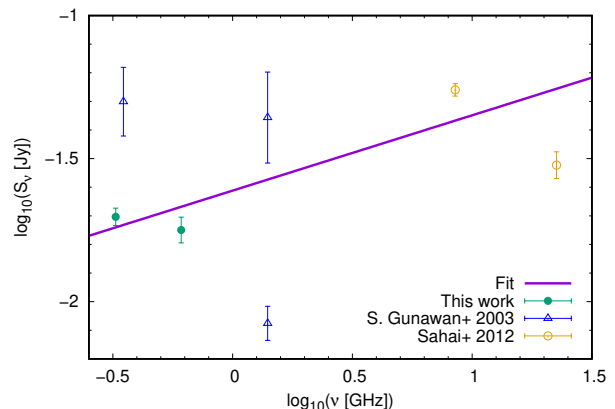
**WDDGGHK 9:** Another neighbor of the Tadpole but barely detected by GMRT observations. The weakness of the emission makes the estimate of its physical properties very uncertain. We estimate its length to be  $\sim 15\,000 \text{ AU}$  and its orientation to be  $P.A. = 127^{\circ}$ , similar to those of the Tadpole and the Goldfish.

**WDDGGHK 10:** Isolated source located inside the box *a* in Figure 1. It is  $3/5$  west of the prominent  $4'$  long bowshock of the DR18 ionized region. Figure 2 shows that WDDGGHK 10 comprises two main radio components, oriented east-west, interestingly the same orientation aimed by the DR18 bow shock. NIR and optical images (W12) show WDDGGHK 10 has a head-tail bow-shock morphology, in good agreement with the GMRT image, which displays stronger emission toward the head of the shock (west) and somewhat fainter toward the tail (east). At radio wavelengths, the contrast between head and tail is not as large as in the other frEGGs studied here. Maybe a reason for this feature

**Table 3.** Flux density values of WDDGGHK7 (the Tadpole) from radio to the IR range.

$\nu$ (GHz)	$S_{\nu}$ (Jy)	Resolution ( $''$ )	Reference
0.325	$0.019 \pm 0.0006$	7.8	This work
0.350	$0.05 \pm 0.006$	55	S. Gunawan et al. 2003
0.610	$0.017 \pm 0.0008$	7.6	This work
1.4	$0.0084 \pm 0.0005$	13	S. Gunawan et al. 2003
1.4C	$0.044 \pm 0.007$	55	S. Gunawan et al. 2003
8.5	$0.055 \pm 0.0012$	3.2	Sahai et al. 2012
22.5	$0.030 \pm 0.0014$	3.2	Sahai et al. 2012

**Notes.** 1.4C: Convolved to the 350 MHz beam.



**Fig. 4.** SED of the Tadpole (WDDGGHK7). We have fitted a power-law to the available radio interferometer data (dashed green line); the purple line shows the sum of both components.

could be that the shreds of the tail of this object are apparently twisted and cross each other at the position of the eastern radio peak (noticeable in the optical and NIR images; Wright et al. 2012). This would explain the emission enhancement from the tail. On average, this source has the steepest spectral index in our sample ( $\alpha = 0.3$ ). It is worth to note that the two isolated frEGGs have average positive spectral indices, while the other sources have flat or slightly negative spectral indices. The spatial distribution of the spectral index toward WDDGGHK 10 (Fig. 3) is also interesting since it varies from the head at  $\alpha \sim -0.1$ , growing in the center to  $\alpha \sim 0.5$  and decreasing toward the easternmost part of the tail to  $\alpha \sim 0.3$ . These variations in the spectral index could be due to differences in the optical depth along the source.

Regarding the protostellar content of WDDGGHK 10, the only IR source inside the photoionized region found by Wright et al. (2012) is IRAS 20329+4104, which is about  $10''$  ( $14\,000 \text{ AU}$ ) away from the peak emission in radio from this frEGG.

#### 4.3. Tadpole spectral energy distribution

In this section we analyze the SED of the object WDDGGHK7, the Tadpole, built up with data from previous radio and infrared observations (Table 3). We selected the Tadpole as an example of a frEGG with a flat/negative spectral index (regarding to our data), and the one with a suspicious negative spectral index at higher frequencies (Sahai et al. 2012a).

At radio wavelengths we use a power-law obtaining an average slope of 0.26 (see Fig. 4) which is consistent with ther-

<sup>2</sup> This value represents a lower limit.

mal emission from a partially optically thin plasma. The thermal emission is expected to be stationary unless the thermodynamical quantities of the gas (e.g.,  $n_e$ ) suffer a quick evolution in a year-scale, which is highly unlikely. Therefore, the mismatch between our observations and the values reported by Setia Gunawan et al. (2003) should be explained by differences in beam-size and calibration errors.

Although partial power-law fits in specific ranges of the radio spectrum show negative spectral indices, the overall index from frequencies below 1 GHz to 22.5 GHz is consistent with thermal emission. If the data at 8 and 22 GHz by Sahai et al. (2012a) is interpreted as non-thermal emission, one would expect that the flux detected at lower frequencies of 0.325-1.4 GHz to be much higher. The only way to reconcile this would be to summon the action of absorption processes which reduce the observed low-frequency emission (Melrose 1980). However, this would lead to a positive spectral index at low frequencies which is in contradiction to the flat spectrum reported we observe. Therefore, we favour an explanation involving calibration uncertainties and/or the presence of diffuse emission which is lost (not detected) at high frequencies (as the LAS is smaller at higher frequencies), which could artificially generate the observed slightly negative spectral index. More observations at intermediate frequencies will help to settle this issue.

## 5. Discussion

In this section, we present a short discussion about the sources causing the ionization of the proplyd-like objects found in Cygnus OB2. We also discuss the possible origin of the flat radio-spectral indices of some of them, focusing particularly on the Tadpole, which has been observed at several wavelengths and its radio emission has been interpreted as non-thermal in the past (Sahai et al. 2012a).

### 5.1. Sources causing the ionization

In Table 4 we report the photon rate necessary to ionize the Cygnus OB2 frEGGs. It ranges between  $5 \times 10^{43}$  and almost  $10^{46}\text{s}^{-1}$ , corresponding to main sequence stars with spectral types B3-B1 (Panagia 1973). According to Wright et al. (2012), the ionizing stars responsible for photoionizing all the Cygnus OB2 frEGGs could be OB2 #8 ( $RA=20:33:15.07$ ,  $Dec=41:18:50.47$ ) or OB2 #22 ( $RA=20:33:08.79$ ,  $Dec=41:13:18.21$ ), although they do not provide quantitative evidence. These systems (see e.g., Wright et al. 2015) are a binary comprised by an O6 I plus an O4.5 III stars (OB2 #8A) along with an O6.5 III (OB2 #8B), and another binary with an O3 If plus an O6 V stars (OB2 #22). The aggregated ionization rates of these systems are  $7.52 \times 10^{49}$  and  $7.00 \times 10^{49}$  photons per second, respectively. We use the Lyman photon rates tabulated by Martins et al. (2005), which have into account line-blanketing and wind effects. In addition to these stellar systems we have tested another two massive systems favorably located in the area so that they could also be responsible for ionizing all the frEGGs (Wright et al. 2015): OB2 #9 (comprised by an O5 I plus an O3.5 III), and MT 516 (an O5.5 V star, with the number 516 in the system by Massey & Thompson 1991). These spectral types lead to total ionizing photon fluxes of  $9.30 \times 10^{49}$  and  $1.26 \times 10^{49}$  per second, respectively. From all four massive stellar systems, the only with accurate Gaia DR2 parallax reported is OB2 #9 ( $\pi = 0.60$ ), at a distance of 1663 pc. Now, to test if these stars are ionizing the

Cygnus OB2 frEGGs we use as a first criterion the prescription in Garay et al. (1987). In particular, their equation 4,

$$N_i = N_{i0} \frac{\Omega}{4\pi} e^{-\tau_{\text{Lyman}}},$$

provides the ionizing photon rate  $N_i$  incident on a source that subtends a solid angle  $\Omega$  as seen from a star with a  $N_{i0}$  ionizing photon rate. Here  $\Omega = \pi * (r/d)^2$ , where  $d$  is the distance from the ionized object to the ionizing source and  $r$  the projected radius of the object as seen from the ionizing stars. For the calculations we use the projected distance as  $d$  (therefore the obtained  $N_i$  would be upper limits), and assume that the Lyman continuum optical depth  $\tau_{\text{Lyman}}$  is negligible. We compare these rates with the ionization photon flux needed to ionize the frEGGs derived from their associated radio continuum emission (Table 4). A second criterion to evaluate if a star is effectively ionizing the OB2 frEGGs is the relative agreement between the orientation of the projected line linking the star with the head of a frEGG and the position angle of that frEGG (Table 5). The positional alignment is quite good for stars #22, #9 and MT 516, with deviations less than 1.5 times the P.A. uncertainty, in general (although #9 and WDDGGHK 7-8 are not that well aligned). On the contrary, #8 shows large deviations for some frEGGs, whether in the group of WDDGGHK 2, WDDGGHK 7 or with WDDGGHK 10. Thus, system OB2 #8 seems not adequately located to be a good candidate. Regarding the Lyman photon fluxes, none of the stellar systems provide a high enough rate to ionize the southern WDDGGHK 2-4 frEGGs, specially WDDGGHK 3, which needs  $7.2 \times 10^{45}$  photons  $\text{s}^{-1}$ . The single star MT 516, although is closer to all the frEGGs is of a later type and cannot ionize the sources on its own. Only systems #9 and #22 are close to the required values. They provide higher than needed fluxes for frEGGs WDDGGHK 7-10, slightly lower for WDDGGHK 6 and half of the needed rate for WDDGGHK 3. The reason for a higher than needed rate may be simply a projection effect, while lower than needed rates may indicate that a combination of fluxes from different stellar systems are responsible for the final ionization. When adding the contribution from both stellar systems (#9 and #22), the resulting photon flux is higher than needed, which may be lowered if unprojected distances are used in the calculations. Therefore, the most favorable ionizing sources for the Cygnus OB2 frEGGs are OB2 #9 and OB2 #22, and small amounts of photons would also be provided by the other stellar systems. It is noticeable, that all the stellar candidates could provide Lyman photon fluxes hundreds of times the rate needed to ionize WDDGGHK 9.

### 5.2. Spectral indices

In Table 1 we report the measured spectral indices extracted from the fluxes obtained at the two observing frequencies: 325 MHz and 610 MHz. The determination of the spectral indices produces small values with relatively large uncertainties. All spectral indices are below the canonical 0.6 index typical of HII regions (e.g., Churchwell & Walmsley 1975). Moreover, two of the frEGGs (WDDGGHK 7 and WDDGGHK 8) show negative spectral indices of  $-0.2 \pm 0.2$ , consistent with flat indices if the uncertainties are taken into account. As already pointed out in previous sections, a negative spectral index may indicate the presence of non-thermal emission. However, in this case, the reason behind the low spectral indices measured in the Cygnus OB2 frEGGs may be the existence of significant amounts of diffuse gas (extended and tenuous emission), preferentially detected at lower frequencies by the GMRT (see also Ramachandran et al.

**Table 4.** Differences between needed ionizing photon rates and those expected from nearby massive Cygnus OB2 stars.

Source	$N_i^{\text{needed}}$ ( $\text{s}^{-1}$ )	$N_i^{\#9}$ ( $\text{s}^{-1}$ )	$\Delta N_i \#9$ (%)	$N_i^{\#22}$ ( $\text{s}^{-1}$ )	$\Delta N_i \#22$ (%)	$N_i^{\#8}$ ( $\text{s}^{-1}$ )	$\Delta N_i \#8$ (%)	$N_i^{\text{MT 516}}$ ( $\text{s}^{-1}$ )	$\Delta N_i \text{MT 516}$ (%)
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
WDDGGHK2	$7.32 \times 10^{44}$	$6.19 \times 10^{44}$	-15	$5.02 \times 10^{44}$	-31	$4.30 \times 10^{44}$	-41	$1.31 \times 10^{44}$	-82
WDDGGHK3	$7.19 \times 10^{45}$	$4.25 \times 10^{45}$	-41	$3.47 \times 10^{45}$	-52	$2.90 \times 10^{45}$	-60	$9.35 \times 10^{44}$	-87
WDDGGHK4	$2.02 \times 10^{45}$	$2.21 \times 10^{45}$	9	$1.81 \times 10^{45}$	-10	$1.49 \times 10^{45}$	-26	$4.88 \times 10^{44}$	-76
WDDGGHK6	$9.05 \times 10^{45}$	$8.35 \times 10^{45}$	-8	$7.93 \times 10^{45}$	-12	$4.48 \times 10^{45}$	-50	$2.85 \times 10^{45}$	-68
WDDGGHK7	$2.63 \times 10^{45}$	$4.16 \times 10^{45}$	58	$3.37 \times 10^{45}$	28	$2.69 \times 10^{45}$	2	$1.18 \times 10^{45}$	-55
WDDGGHK8	$2.02 \times 10^{45}$	$3.71 \times 10^{45}$	84	$3.07 \times 10^{45}$	52	$2.33 \times 10^{45}$	15	$1.16 \times 10^{45}$	-42
WDDGGHK9	$5.32 \times 10^{43}$	$9.43 \times 10^{44}$	1671	$7.66 \times 10^{44}$	1338	$6.26 \times 10^{44}$	1075	$2.54 \times 10^{44}$	377
WDDGGHK10	$1.31 \times 10^{45}$	$3.37 \times 10^{45}$	157	$2.44 \times 10^{45}$	86	$2.85 \times 10^{45}$	118	$5.39 \times 10^{44}$	-59

**Notes.** The distances between the ionizing candidates and the WDDGGHK 3 group range 1300''-1900'' ( $\sim 13$  pc); to WDDGGHK 6, distances range 600''-1200'' ( $\sim 7.5$  pc); to the WDDGGHK 7 group, distances range 550''-900'' ( $\sim 6$  pc); to WDDGGHK 10, distances range 1000''-1100'' ( $\sim 8.5$  pc). (1): Name of the source; (2): Rate of ionizing Lyman photons needed to produce the measured radio emission; (3), (5), (7) and (9): Rates of Lyman photons incident at each of the frEGGs (see text) and produced by Cygnus OB2 stars #9 (O5 I and O3.5 III), #22 (O5 I and O3.5 III), #8 (considering #8A and #8B, which comprises an O6 I plus an O5.5 III, and an O6.5 III), and the star RPL841 (O5.5 V), respectively. The spectral type of the stars were adopted from Wright et al. (2015). The ionizing photon rates from the corresponding spectral type, using the models based on observational  $T_{\text{eff}}$  scales by Martins et al. (2005); (4), (6), (8) and (10): Difference in % of the needed and the provided ionizing photon rates,  $100 \times (N_i' - N_i^{\text{needed}})/N_i^{\text{needed}}$ . Negative percentages indicate a deficit of ionizing photons from the star, while positive percentages indicate the opposite, an excess of ionizing photons (or a larger than projected distance to the object).

**Table 5.** Differences between the *P.A.* of frEGGs and its orientation with respect to nearby massive Cygnus OB2 stars.

Source	<i>P.A.</i> ( $^\circ$ )	<i>P.A.</i> from #9 ( $^\circ$ )	$\Delta \#9$ ( $\sigma$ )	<i>P.A.</i> from #22 ( $^\circ$ )	$\Delta \#22$ ( $\sigma$ )	<i>P.A.</i> from #8 ( $^\circ$ )	$\Delta \#8$ ( $\sigma$ )	<i>P.A.</i> from MT 516 ( $^\circ$ )	$\Delta \text{MT 516}$ ( $\sigma$ )
(1)	(2)	(3)	(4)	(5)	(6)	(7)	(8)	(9)	(10)
WDDGGHK2	132	138	0.8	136	0.5	145	1.7	139	0.9
WDDGGHK3	138	141	0.4	139	0.1	147	1.2	143	0.7
WDDGGHK4	141	144	0.4	142	0.1	151	1.3	147	0.8
WDDGGHK6	177	180	0.4	172	0.7	178	0.1	184	0.9
WDDGGHK7	108	120	1.6	116	1.1	136	3.7	118	1.3
WDDGGHK8	112	125	1.7	121	1.2	141	3.9	125	1.7
WDDGGHK9	127	121	0.8	116	1.5	135	1.1	119	1.1
WDDGGHK10	90	91	0.1	88	0.3	105	2.0	85	0.7

**Notes.** (1): Name of the source; (2): Position angle (*P.A.*) of each frEGG (we estimate the typical uncertainty of these measurements to be  $8^\circ$ ); (3), (5), (7) and (9): *P.A.* measured from the considered Cygnus OB2 ionizing candidates to each frEGG; (4), (6), (8) and (10): Difference (absolute value) between the frEGG *P.A.s* and the orientation with respect to the ionizing star candidates, in terms of the uncertainty of the frEGG's *P.A.* measurement.

2017). Supporting this argument, we check for any well-known HII region lying within the field of view of our Cygnus OB2 GMRT observations. We found for instance DR 18. Using our observations, we obtain a radio-spectral index of  $0.3 \pm 0.01$  for this extended object, which is also lower than the expected 0.6 value. For DR 18, it is evident also that for extended objects, the GMRT is more sensitive to the extended weak emission at 325 MHz than at 610 MHz. This effect can be quantitatively accounted for by reporting the number of pixels with significant emission at each frequency. Again, the continuum emission of DR 18 at 325 MHz spreads on 4% more pixels than at 610 MHz. Similar results are obtained for most of the frEGGs: for WDDGGHK 8 the diffuse emission extends over a 10% more pixels at 325 MHz, while for WDDGGHK 10, this percentage rises to 30%. This effect is more pronounced in the tail of WDDGGHK 7 and WDDGGHK 8, which is the region with more diffuse emission and negative values of the spectral index (Fig. 3). As a consequence the radio-spectral indices are smaller than they are in reality and therefore we are no further interpreting them as originated by non-thermal emission. We just make the caveat for future usage of our results that diffuse emission have to be ac-

counted for when estimating the spectral index in these sources. Then the values of  $\tau$ ,  $n_e$ ,  $M_{\text{ion}}$  and  $E.M.$  from Table 2 were determined at the frequency of 325 MHz since they are not affected by the loss of emission at the frequency of 610 MHz.

A special case among the studied frEGGs is that WDDGGHK 7, the Tadpole, since Sahai et al. (2012a) have reported the detection of a negative spectral index between 8.5 GHz and 22.0 GHz. Our low-frequency spectral index is consistent with a slightly negative or flat value, which cannot be reconciled with Sahai et al. (2012a)'s data (Figure 4). As already mentioned, we only detect negative spectral index values between  $-1.0$  and  $-0.5$  along the tail of this object (Fig. 3). Hence, in case there is any non-thermal emission it should belong to this part of the object. However, 15% of Tadpole's pixels at 325 MHz are not seen at 610 MHz and most of them are lost in the tail. This suggests that Tadpole's tail is probably affected by diffuse emission, which artificially produces the negative index. When studying together the GMRT measured fluxes with those of Sahai et al. (2012a) for the Tadpole (Figure 4) a power-law fit produces a slope of  $0.3 \pm 0.1$ . This is probably a more reliable value for the radio-spectral index of this source. It shows that thermal emis-



sion dominates Tadpole's radio emission (a black-body is a good fit for its infrared emission in addition), although more data at intermediate and larger frequencies would improve this fit to finally constrain the nature of the radio emission.

## 6. Summary and conclusions

We study the externally ionized objects of the Cygnus OB2 region previously identified in Wright et al. (2012) through GMRT 325 MHz and 610 MHz observations. After describing them observationally we determine their physical properties, such as their electronic densities, ionized mass, optical depth and the amount of photons needed to be ionized. We compare the ionizing photon rate with that produced by different massive stars in the neighborhood searching for good ionizing candidates. We also compare the orientation of the frEGGs with the position angle between the ionizing stars and the frEGGs. We conclude that the stellar systems Cyg OB2 #9 and Cyg OB2 #22 are probably responsible for their ionization.

We also obtain radio spectral indices between 325 MHz and 610 MHz for five frEGGs of our sample. Some sources show an average flat spectral index consistent with thermal free-free emission. However, some of the spectral index maps show regions of negative values, that we interpret as a possible effect of the presence of diffuse gas, since the interferometer would be more sensitive to the extended diffuse gas at low frequencies, resulting in spectral index values close or below to zero. Still our observations are not conclusive about the existence of regions with non-thermal emission in the frEGGs, although they suggest that the main contribution to the continuum radio-emission is thermal.

Regarding the object WDDGGHK 7 (the Tadpole), thought to be non-thermal at radio-wavelengths, we collect data in the literature to build its SED. We fit a power-law function to the radio-continuum measurements obtaining a 0.3 spectral index consistent with thermal emission. Hence, according to our data, most of the radio-emission from frEGGs would not be non-thermal.

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