

J-PAS: forecasts on dark energy and modified gravity theories

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ABSTRACT

The next generation of galaxy surveys will allow us to test one of the most fundamental assumptions of the standard cosmology, i.e. that gravity is governed by the general theory of relativity (GR). In this paper, we investigate the ability of the Javalambre Physics of the Accelerating Universe Astrophysical Survey (J-PAS) to constrain GR and its extensions. Based on the J-PAS information on clustering and gravitational lensing, we perform a Fisher matrix forecast on the effective Newton constant, μ , and the gravitational slip parameter, η , whose deviations from unity would indicate a breakdown of GR. Similar analysis is also performed for the DESI and Euclid surveys and compared to J-PAS with two configurations providing different areas, namely an initial expectation with 4000 deg² and the future best case scenario with 8500 deg². We show that J-PAS will be able to measure the parameters μ and η at a sensitivity of 2–7 per cent, and will provide the best constraints in the interval $z = 0.3$ – 0.6 , thanks to the large number of ELGs detectable in that redshift range. We also discuss the constraining power of J-PAS for dark energy models with a time-dependent equation-of-state parameter of the type $w(a) = w_0 + w_a(1 - a)$, obtaining $\Delta w_0 = 0.058$ and $\Delta w_a = 0.24$ for the absolute errors of the dark energy parameters.

Key words: dark energy – large-scale structure of Universe.

1 INTRODUCTION

The success of the general theory of relativity (GR) is unquestionable. For about a hundred years now, GR has remained unchanged and capable of explaining observations and experiments in a number of regimes, such as the dynamics of the Solar system, gravitational wave emission, the energetics of supermassive black holes, and quasars [see e.g. Will (2014) for the status of experimental tests of

GR]. When extrapolated to cosmological scales, Einstein’s theory has also provided a very good description of the evolution of the Universe, which is obtained at the cost of postulating the existence of both dark matter as well as a dark energy component, i.e. an additional field with fine-tuned properties responsible for the current cosmic acceleration (Sahni & Starobinsky 2000; Padmanabhan 2003; Peebles & Ratra 2003; Copeland, Sami & Tsujikawa 2006).

Given the unnatural properties of dark energy (Weinberg 1989), a promising alternative to the standard scenario (GR plus dark energy) is based on infra-red modifications to GR, leading to a weakening

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of gravity on cosmological scales and thus to late-time acceleration. In the past few decades, a number of modified or extended theories of gravity (MG) have been proposed (Dvali, Gabadadze & Porrati 2000; Sahni & Shtanov 2003; Capozziello 2002; Carroll et al. 2004; Santos et al. 2007; see also Sotiriou & Faraoni 2010; Capozziello & De Laurentis 2011; Clifton et al. 2012a; Ferreira 2019 for recent reviews). In general, these ideas explore as much as they can the loopholes of Lovelock’s theorem, while preserving GR on astrophysical scales. Recently, the number of allowed MG theories was significantly restricted (Baker et al. 2017; Creminelli & Vernizzi 2017; Ezquiaga & Zumalacárregui 2017), given the tight bound on the speed of propagation of gravitational waves, $|c_{\text{gw}}/c - 1| \lesssim 10^{-15}$, obtained from the binary neutron star merger GW170817 (Abbott et al. 2017). In the near future, other constraints are also expected from black hole imaging, as recently reported by the Event Horizon Telescope.¹

Cosmological observations are also able constrain MG theories at the largest scales, as has been shown by e.g. the Planck experiment (Aghanim et al. 2018). In this context, the large-scale structure surveys that will become available in the coming years will play the major role (Ferreira 2019). Those surveys can be categorized in two main types: (i) spectroscopic surveys, obtaining high-quality spectra (and corresponding high-quality redshift measurements thereof), typically targeting a pre-selected subsample of extragalactic objects BOSS (e.g. Dawson et al. 2013), eBOSS (Dawson et al. 2016), DESI (Flaugher & Bebek 2014; Aghamousa et al. 2016), *Euclid* (Laureijs et al. 2011; Amendola et al. 2018), etc., and (ii) photometric surveys, probing the sky at deeper magnitudes in a reduced number of filters, providing significantly larger catalogues of sources, but at the expense of a poorer spectral characterization (e.g. DES, Abbott et al. 2005; LSST, Abell et al. 2009).

An intermediate regime is represented by the so-called spectrophotometric surveys (COMBO-17, Wolf et al. 2003; ALHAMBRA, Moles et al. 2008; COSMOS, Ilbert et al. 2009; MUSYC, Cardamone et al. 2010; CLASH, Postman et al. 2012; SHARDS, Pérez-González et al. 2013; PAU, Martí et al. 2014; J-PLUS, Cenarro et al. 2019a; J-PAS, Benítez et al. 2014; SPHEREx, Korngut et al. 2018) that combine deep imaging with multicolour information obtained through combination of broad-, medium- and narrow-band filters. In this way, a low-resolution spectrum (also known as ‘pseudo-spectrum’) is obtained for every pixel in the survey’s footprint, and in particular for each and all sources present in the joint catalogue extracted from the combination of all bands. This allows providing high-quality photometric redshift estimations for a much larger number of objects compared with spectroscopic surveys, on top of 2D information for those sources that are spatially resolved.

This paper discusses the expected cosmological implications of J-PAS Benítez et al. (2014) on dark energy and modified gravity theories. As is well known, the main body of observations currently available comes from distance measurements that map the expansion history of the Universe at the background level. However, these measurements alone are not enough to discriminate between a dark energy fluid and modifications to GR, as different models can predict the same expansion history (Kunz 2012). Additional observational information is thus required in order to break the model degeneracy and, in particular, the growth of structures and gravitational lensing, which is directly sensitive to the growth of dark matter perturbations – in contrast with measurements based on

galaxies, neutral hydrogen or any other baryonic tracer – are among the most promising avenues in this respect.

Here, we consider the J-PAS information on clustering and gravitational lensing and perform a Fisher matrix forecast on the effective Newton constant, μ , and the gravitational slip parameter, η (defined in Section 3), assuming two configurations of area for J-PAS, i.e. 4000 and 8500 deg². For completeness, we also discuss the constraining power of J-PAS for dark energy models with a time-dependent equation-of-state parameter $w(a)$, and compare all J-PAS forecasts with those expected by the DESI (Flaugher & Bebek 2014; Aghamousa et al. 2016) and *Euclid* surveys (Laureijs et al. 2011; Amendola et al. 2018). In this sense, this work updates some of the results contained in Benítez et al. (2014) and also makes new forecasts, including several new scenarios. Further analysis on interactions in the dark sector can be found in Costa et al. (2019).

2 THE J-PAS SURVEY

The Javalambre Physics of the Accelerating Universe Astrophysical Survey (J-PAS; Benítez et al. 2014) is a spectrophotometric survey to be conducted at the Observatorio Astrofísico de Javalambre (hereafter OAJ), a site on top of *Pico del Buitre*, a summit about ~2000 m high above sea level at the Sierra of Javalambre, in the Eastern region of the Iberian peninsula. The Javalambre Survey Telescope (JST/T250), a 2.5 m diameter, altazimuthal telescope, will be on charge of J-PAS. JST will be equipped with the Javalambre Panoramic Camera (JPCam), a 14-CCD mosaic camera using a new large format e2v 9.2 k-by-9.2 k 10 μm pixel detectors, and will incorporate a 54 narrow- and 4 broad-band filter set covering the optical range (Marín-Franch et al. 2017). The field of view covered by JPCam is close to 5 deg², and thus the JST/JPCam system constitutes a system specifically defined to optimally conduct spectrophotometric surveys. J-PAS is not the first survey being carried out at the OAJ, since the Javalambre Local Universe Photometric Survey (J-PLUS), conducted by the Javalambre Auxiliary Survey Telescope (JAST/T80), has already covered about 1600 deg² with 12 broad- and narrow-band filters (some of them in common to J-PAS). We refer the reader to Benítez et al. (2014) and Cenarro et al. (2019b) for more details on J-PAS and J-PLUS, respectively.

3 DARK ENERGY AND MODIFIED GRAVITY PARAMETRIZATIONS

In recent years, many different models of dark energy or MG have been proposed as alternatives to the standard Λ cold dark matter (Λ CDM) cosmology. The possibility of confronting such alternatives with observations in a largely model-independent way has motivated the development of theoretical frameworks in which general modifications can be captured in a few effective parameters that can be directly tested by observations (Clifton et al. 2012b; Silvestri, Pogosian & Buniy 2013).

In this section, we introduce the phenomenological parametrizations of dark energy and MG that will be considered throughout the paper.

3.1 Dark energy

In the context of GR, dark energy is understood as a smooth (non-clustering) energy component with a sufficient negative pressure, p , to violate the strong energy condition ($\rho + 3p \geq 0$, where ρ is the energy density) and accelerate the Universe. Many different

¹<https://eventhorizontelescope.org>

models of dark energy have been proposed in recent years (see e.g. Peebles & Ratra 2003; Copeland et al. 2006; Barboza & Alcaniz 2008; and references therein), based on fluid descriptions with different equations of state or the inclusion of an additional scalar field, as in the quintessence models.

Rather than focusing on particular models, we will consider a phenomenological description of dark energy as a perfect fluid with an equation of state given by the parametrization (Chevallier & Polarski 2001; Linder 2003)

$$w(a) = w_0 + w_a(1 - a), \quad (1)$$

which reduces to the standard Λ CDM model for values of $w_0 = -1$ and $w_a = 0$. Note also that this effective modification with respect to the standard cosmology mainly affects the background evolution. Notice that the dark energy component could acquire cosmological perturbations that are already taken into account in the CAMB code (Lewis, Challinor & Lasenby 2000).

3.2 Modified gravity

We will consider for simplicity the case of MG theories that include additional scalar degrees of freedom. Extensions of the model-independent approach for modified theories including additional vector fields can be found in Resco & Maroto (2018b).

Let us then consider the scalar-perturbed flat Friedmann–Lemaître–Robertson–Walker (FLRW) metric, written in the longitudinal gauge (Amendola & Tsujikawa 2010; Tsujikawa, De Felice & Alcaniz 2013):

$$ds^2 = -(1 + 2\Psi)dt^2 + a^2(t)(1 + 2\Phi)d\mathbf{x}^2. \quad (2)$$

The modified Einstein equation to first order in perturbations can be written as

$$\delta\bar{G}^\mu_\nu = 8\pi G \delta T^\mu_\nu, \quad (3)$$

where the perturbed modified Einstein tensor $\delta\bar{G}^\mu_\nu$ can in principle depend on both the metric potentials Φ and Ψ , and the perturbed scalar field $\delta\phi$. On the other hand, at late times the only relevant energy component is non-relativistic matter so that

$$\delta T^0_0 = -\rho_m \delta_m, \quad (4a)$$

$$\delta T^0_i = -\rho_m v_i, \quad (4b)$$

$$\delta T^i_j = 0, \quad (4c)$$

where v_i is the three velocity of matter, ρ_m is the total matter density, and $\delta_m = \delta\rho_m/\rho_m$ is the corresponding matter density contrast, which is related to the galaxy density contrast δ_g via the bias factor b , as $\delta_g = b \delta_m$.

Using the Bianchi identities in the modified Einstein tensor, we find that in the sub-Hubble regime ($k \gg aH$, $H = \dot{a}(t)/a(t)$ is the Hubble parameter) there are only two independent Einstein equations, which together with the scalar field equation of motion lead to the following set of equations to first order in perturbations in Fourier space:

$$k^2 \Phi = 4\pi G a^2 \mu \eta \rho_m \delta_m, \quad (5)$$

$$k^2 \Psi = -4\pi G a^2 \mu \rho_m \delta_m. \quad (6)$$

Here, for simplicity, we have restricted ourselves to the case of second-order operators and used the so-called quasi-static approximation, in which time derivatives can be neglected with respect

to the spatial ones. Notice that the quasi-static approximation is a good one for models with large speed of sound of dark energy perturbations and can be safely employed for current galaxy surveys. For future large surveys, it could be inappropriate on scales close to the Hubble horizon. Also as shown in Sawicki & Bellini (2015), it should never be used for the integrated Sachs–Wolfe effect analysis.

Note that on the sub-Hubble scales, δ_m agrees with the density perturbation Δ used in Silvestri et al. (2013) since $\Delta = \delta_m + \frac{3aHv}{k}$. Therefore, in the quasi-static approximation, a general modification of Einstein's equations can be written in terms of two arbitrary functions of time and scale $\mu(a, k)$ and $\eta(a, k)$ (Pogosian et al. 2010; Silvestri et al. 2013). These parameters can be understood as an effective Newton constant, $G_{\text{eff}}(a, k)$, given by

$$\mu(a, k) = \frac{G_{\text{eff}}}{G}, \quad (7)$$

and the gravitational slip parameter

$$\eta(a, k) = -\frac{\Phi}{\Psi}, \quad (8)$$

which modifies the equation for the lensing potential that depends upon the combination $(\Psi - \Phi)/2$. Thus, deviations from $\mu = \eta = 1$ indicate a breakdown of standard GR. Notice that alternative parametrizations have been considered in the literature such as (μ, Σ) with $\Sigma = \mu(1 + \eta)/2$. We have preferred to use (μ, η) , since they have a more direct physical interpretation.

The modified equations can be rewritten as

$$k^2 \Psi \simeq -4\pi G_{\text{eff}} a^2 \rho_m \delta_m, \quad (9)$$

and

$$\frac{\Psi - \Phi}{2} \simeq -\frac{3G_{\text{eff}}}{2G} \frac{1 + \eta}{2} \left(\frac{aH}{k}\right)^2 \Omega_m(a) \delta_m. \quad (10)$$

where $\Omega_m(a) = \Omega_m a^{-3} E(a)^{-2}$ is the matter density parameter and $E(a) = H(a)/H_0$, with the Hubble constant written as $H_0 = 100h \text{ km s}^{-1} \text{ Mpc}^{-1}$.

Using the standard conservation equation, $T^{\mu\nu}_{; \nu} = 0$, we obtain the continuity and Euler equations, which in the sub-Hubble regime and for non-relativistic matter, reduce to

$$a \dot{\delta}_m = -\theta, \quad (11)$$

$$a \dot{\theta} = -aH\theta + k^2 \Psi, \quad (12)$$

where $\theta = i(\mathbf{k} \cdot \mathbf{v})$.

Taking the time derivative of equation (11) and using (12), we obtain the modified growth equation that reads

$$\delta_m'' + \left(2 + \frac{H'}{H}\right) \delta_m' - \frac{3}{2} \mu(a, k) \Omega_m(a) \delta_m \simeq 0, \quad (13)$$

where the prime denotes derivative with respect to $\ln a$.

Notice that in general, in typical modified gravity theories such as $f(R)$ or scalar-tensor models, the effective Newton constant and slip parameter generically depends on both scale k and time a . For simplicity, in our analysis we will limit ourselves to two particular classes of effective parameters, namely scale-independent parametrizations with $\mu = \mu(a)$ and $\eta = \eta(a)$ and time-independent parametrizations, i.e. $\mu = \mu(k)$ and $\eta = \eta(k)$, in order to independently assess the sensitivity of the surveys to scale and time dependences. Notice that, in any case, this type of parametrizations will only be valid in a given range of scales and times. At high redshifts or sub-galactic scales, we expect the standard $\mu = \eta = 1$ values to be recovered.

In the scale-independent case, two particularly relevant examples will be analysed. On one hand, the constant in time case and, on the other, the parametrization proposed in Simpson et al. (2013), which is usually employed in the literature (Ade et al. 2016),

$$\mu(a) = 1 + (\mu_0 - 1) \frac{1 - \Omega_m(a)}{1 - \Omega_m}, \quad (14)$$

$$\eta(a) = 1 + (\eta_0 - 1) \frac{1 - \Omega_m(a)}{1 - \Omega_m}. \quad (15)$$

This parametrization ensures that at high redshift the standard GR values are recovered.

4 FISHER MATRICES FOR GALAXY AND LENSING POWER SPECTRA

The Fisher matrix formalism provides a simple way to estimate the precision with which certain cosmological parameters could be measured from a set of observables once the survey specifications and the fiducial cosmology are fixed. Thus, given a set of parameters $\{p_\alpha\}$, the Fisher matrix \mathbf{F}^p is just the inverse of the covariance matrix in the parameters space. It provides the marginalized error for the p_α parameter as $\sqrt{F_{\alpha\alpha}^{-1}}$. The corresponding 1σ region is just an ellipsoid in the parameter space since the probability distribution function (PDF) are assumed to be Gaussian in the Fisher formalism. If we are interested in obtaining errors for a different set of parameters $\{q_\alpha\}$, the Fisher matrix of the new parameters simply reads

$$\mathbf{F}^q = \mathbf{P}^t \mathbf{F}^p \mathbf{P}, \quad (16)$$

where $\mathbf{P} = \mathbf{Q}^{-1}$ and $Q_{\alpha\beta} = \partial q_\alpha / \partial p_\beta$, evaluated on the fiducial model.

In the following, we provide general expressions for the Fisher matrices for the galaxy power spectrum in redshift space and for the lensing convergence power spectrum, both in different redshift and in k (or ℓ) bins. We will apply them separately to J-PAS (Benitez et al. 2014), DESI (Aghamousa et al. 2016), and *Euclid* (Laureijs et al. 2011) galaxy surveys and for J-PAS and *Euclid* lensing surveys.

4.1 Fisher matrix for galaxy clustering

Following Amendola et al. (2013, 2014), let us introduce the following dimensionless parameters A and R ,

$$A = D b \sigma_8, \quad (17)$$

$$R = D f \sigma_8, \quad (18)$$

where $D(z) = \delta_m(z)/\delta_m(0)$ is the growth factor, b is the bias, and f is the growth function defined by

$$D(z) = \exp \left[\int_0^{N(z)} f(N') dN' \right], \quad (19)$$

being $N(z) = -\log(1+z)$. The σ_8 constant corresponds to $\sigma_8 = \sigma(0.8 \text{ Mpc } h^{-1})$, where

$$\sigma^2(z, R) = D^2(z) \int \frac{k'^2 dk'}{2\pi^2} P(k') |\hat{W}(R, k')|^2, \quad (20)$$

being $P(k)$ the matter power spectrum. We use a top-hat filter $\hat{W}(R, k)$, defined by

$$\hat{W}(R, k) = \frac{3}{k^3 R^3} [\sin(kR) - kR \cos(kR)]. \quad (21)$$

Then, the galaxy power spectrum in redshift space is (Seo & Eisenstein 2003)

$$P(k_r, \hat{\mu}_r, z) = \frac{D_{Ar}^2 E}{D_A^2 E_r} (A + R \hat{\mu}^2)^2 \hat{P}(k) e^{-k^2 \hat{\mu}_r^2 \sigma_r^2}, \quad (22)$$

where sub-index r denotes that the corresponding quantity is evaluated on the fiducial model, $\hat{P}(k) \equiv P(k)/\sigma_8^2$, $\sigma_r = (\delta z(1+z))/H(z)$ with $\delta z(1+z)$ the photometric redshift error, and D_A is the angular distance which, in a flat universe, reads $D_A = (1+z)^{-1} \chi(z)$, with

$$\chi(z) = H_0^{-1} \int_0^z \frac{dz'}{E(z')}. \quad (23)$$

The dependences $k = k(k_r)$, $\hat{\mu} = \hat{\mu}(\hat{\mu}_r)$, and the factor $\frac{D_{Ar}^2 E}{D_A^2 E_r}$ are due to the Alcock–Paczynski effect (Alcock & Paczynski 1979; see also Amendola & Tsujikawa 2010):

$$k = Q k_r, \quad (24)$$

$$\hat{\mu} = \frac{E \hat{\mu}_r}{E_r Q}, \quad (25)$$

$$Q = \frac{\sqrt{E^2 \chi^2 \hat{\mu}_r^2 - E_r^2 \chi_r^2 (\hat{\mu}_r^2 - 1)}}{E_r \chi}. \quad (26)$$

If we consider different galaxies as dark matter tracers with bias b_i , the galaxy power spectrum is (White, Song & Percival 2008; McDonald & Seljak 2009)

$$P_{ij}(k_r, \hat{\mu}_r, z) = \frac{D_{Ar}^2 E}{D_A^2 E_r} (A_i + R \hat{\mu}^2) \times (A_j + R \hat{\mu}^2) \hat{P}(k) e^{-k^2 \hat{\mu}_r^2 \sigma_r^2}, \quad (27)$$

where $A_i = D b_i \sigma_8$. Then, considering a set of cosmological parameters $\{p_\alpha\}$, the corresponding Fisher matrix for clustering of different tracers and for a given redshift bin centred at z_a is (Abramo 2012; Abramo, Secco & Loureiro 2016)

$$F_{\alpha\beta}^C(z_a) = \frac{V_a}{8\pi^2} \int_{-1}^1 d\hat{\mu} \int_{k_{\min}}^\infty dk k^2 \frac{\partial P_{ij}(k, \hat{\mu}, z_a)}{\partial p_\alpha} \Big|_r \times C_{jl}^{-1} \frac{\partial P_{lm}(k, \hat{\mu}, z_a)}{\partial p_\beta} \Big|_r C_{mi}^{-1} e^{-k^2 \Sigma_\perp^2 - k^2 \hat{\mu}^2 (\Sigma_\parallel^2 - \Sigma_\perp^2)}, \quad (28)$$

where

$$\Sigma_\perp(z) = 0.785 D(z) \Sigma_0, \quad (29)$$

$$\Sigma_\parallel(z) = 0.785 D(z) (1 + f(z)) \Sigma_0, \quad (30)$$

with $\Sigma_0 = 11 h^{-1} \text{ Mpc}$ for our fiducial value of $\sigma_8 = 0.82$ in the modified gravity case, and $\Sigma_0 = 6.5 h^{-1} \text{ Mpc}$ for the dark energy case due to the reconstruction procedure (Seo & Eisenstein 2007). Finally, k_{\min} is fixed to $0.007 h \text{ Mpc}^{-1}$ (Amendola et al. 2014). Thus, the exponential cut-off (Seo & Eisenstein 2007) removes the contribution from non-linear scales across and along the line of sight. The factor 0.785 takes into account the different normalization of $(1+z)D(z)$ at high redshifts compared to Seo & Eisenstein

(2007).² The data covariance matrix is

$$C_{ij} = P_{ij} + \frac{\delta_{ij}}{\bar{n}_i}, \quad (31)$$

where $\bar{n}_i = \bar{n}_i(z_a)$ is the mean galaxy density of tracer i in the z bin a . Finally, V_a is the total volume of the a th bin. For a flat Λ CDM model, $V_a = \frac{4\pi f_{\text{sky}}}{3} (\chi(\bar{z}_a)^3 - \chi(\bar{z}_{a-1})^3)$ where f_{sky} is the sky fraction of the survey and \bar{z}_a the upper limit of the a th bin. For the particular case in which we have only one tracer we recover from (28) the standard Fisher matrix of clustering for the power spectrum (22) at z_a (Seo & Eisenstein 2003),

$$F_{\alpha\beta}^C(z_a) = \frac{V_a}{8\pi^2} \int_{-1}^1 d\hat{\mu} \int_{k_{\text{min}}}^{\infty} k^2 V_{\text{eff}} \left. \frac{\partial \ln(P(k, \hat{\mu}, z_a))}{\partial p_\alpha} \right|_r \times \left. \frac{\partial \ln(P(k, \hat{\mu}, z_a))}{\partial p_\beta} \right|_r e^{-k^2 \Sigma_\perp^2 - k^2 \hat{\mu}^2 (\Sigma_\parallel^2 - \Sigma_\perp^2)} dk. \quad (32)$$

Here, V_a is the volume of the redshift slice z_a , and the effective volume is given by

$$V_{\text{eff}} = \left(\frac{\bar{n}(z_a) P(k, \hat{\mu}, z)}{1 + \bar{n}(z_a) P(k, \hat{\mu}, z)} \right)^2. \quad (33)$$

Finally, if we are interested in estimating errors in different k -bins, we sum the information for all z bins in each k_q bin of width Δk_q , so that

$$F_{\alpha\beta}^C(k_q) = \sum_a \frac{V_a}{8\pi^2} \int_{-1}^1 d\hat{\mu} \int_{\Delta k_q} dk k^2 \left. \frac{\partial P_{ij}(k, \hat{\mu}, z_a)}{\partial p_\alpha} \right|_r \times C_{jl}^{-1} \left. \frac{\partial P_{lm}(k, \hat{\mu}, z_a)}{\partial p_\beta} \right|_r C_{mi}^{-1} e^{-k^2 \Sigma_\perp^2 - k^2 \hat{\mu}^2 (\Sigma_\parallel^2 - \Sigma_\perp^2)}. \quad (34)$$

4.2 Fisher matrix for weak lensing

The main observable for the weak lensing measurements is the convergence power spectrum. Using the Limber and flat-sky approximations, we obtain (Lemos, Challinor & Efstathiou 2017)

$$P(\ell) = \int_0^\infty dz \frac{H_0^2 \Omega_m^2}{H(z)} K^2(z) \frac{\mu^2 (1 + \eta)^2}{4} D^2(z) P\left(\frac{\ell}{\chi(z)}\right), \quad (35)$$

where $K(z)$ is defined as

$$K(z) = \frac{3H_0}{2} (1+z) \int_z^\infty \left(1 - \frac{\chi(z')}{\chi(z)}\right) n(z') dz', \quad (36)$$

being $n(z)$ the source galaxy density function as a function of the redshift. For a redshift tomography analysis, we can generalize the convergence power spectrum as (Hu 1999)

$$P_{ij}(\ell) \simeq H_0 \sum_a \frac{\Delta z_a}{E_a} K_i(z_a) K_j(z_a) L_a^2 \hat{P}\left(\frac{\ell}{\chi(z_a)}\right), \quad (37)$$

where we have discretized the integral (35) and defined the dimensionless parameter L as (Amendola et al. 2013)

$$L = \Omega_m D \frac{\mu(1+\eta)}{2} \sigma_8, \quad (38)$$

²Note that there is a typo in the normalization factor of 0.785 on Seo & Eisenstein (2007). We thank Cássio Pigozzo for pointing this out.

where $L_a = L(z_a)$. The function K_i is related to the weak lensing window function for the i -bin by

$$K_i(z) = \frac{3H_0}{2} (1+z) \int_z^\infty \left(1 - \frac{\chi(z')}{\chi(z)}\right) n_i(z') dz', \quad (39)$$

where $n_i(z)$ is the density function for the i -bin, which is obtained as follows: let us first consider the source galaxy density function for the survey (Ma, Hu & Huterer 2005),

$$n(z) = \frac{3}{2z_p^3} z^2 e^{-(z/z_p)^{3/2}}, \quad (40)$$

where $z_p = z_{\text{mean}}/\sqrt{2}$, being z_{mean} the survey mean redshift. Then, within the i -bin we have a new distribution function which is defined to be equal to $n(z)$ inside the bin and zero outside. Now, taking into account the photometric redshift error, $\sigma_i = \delta z (1 + z_i)$, we obtain

$$n_i(z) \propto \int_{\bar{z}_{i-1}}^{\bar{z}_i} z'^2 e^{-(z'/z_p)^{3/2}} e^{-\frac{(z'-z)^2}{2\sigma_i^2}} dz', \quad (41)$$

where \bar{z}_i is the upper limit of the i -bin. Then, the Fisher matrix for weak lensing is given by (Eisenstein, Hu & Tegmark 1999)

$$F_{\alpha\beta}^L = f_{\text{sky}} \sum_\ell \Delta \ln \ell \frac{(2\ell+1)\ell}{2} \text{Tr} \left[\frac{\partial \mathbf{P}}{\partial p_\alpha} \mathbf{C}^{-1} \frac{\partial \mathbf{P}}{\partial p_\beta} \mathbf{C}^{-1} \right], \quad (42)$$

where \mathbf{P} and \mathbf{C} are the matrix of size $n_b \times n_b$ with,

$$C_{ij} = P_{ij} + \gamma_{\text{int}}^2 \hat{n}_i^{-1} \delta_{ij}, \quad (43)$$

$\gamma_{\text{int}} = 0.22$ being the intrinsic ellipticity (see for instance Hilbert et al. 2017). Notice that we are not considering the effect of possible systematic errors in the shear measurements (Huterer et al. 2006). Finally, \hat{n}_i denotes the number of galaxies per steradian in the i th bin,

$$\hat{n}_i = n_\theta \frac{\int_{\bar{z}_{i-1}}^{\bar{z}_i} n(z) dz}{\int_0^\infty n(z) dz}, \quad (44)$$

where n_θ is the areal galaxy density. We sum in ℓ with $\Delta \ln \ell = 0.1$ from $\ell_{\text{min}} = 5$ (Amendola et al. 2014) to ℓ_{max} with $\ell_{\text{max}} = \chi(z_{a'}) k_{\text{max}}$ where $\alpha' = \min(\alpha, \beta)$ and $k_{\text{max}}(z_a)$ is defined so that $\sigma(z_a, \pi/2k_{\text{max}}(z_a)) = 0.35$ using (20), i.e. we only consider modes in the linear regime.

Finally, if we are interested in estimating errors in different ℓ -bins, we introduce a window function in the Fisher matrix (42) in order to take into account only the information of a bin ℓ_a of width $\Delta \ell_a$,

$$F_{\alpha\beta}^L(\ell_a) = f_{\text{sky}} \sum_\ell \Delta \ell \frac{(2\ell+1)}{2} W_a(\ell) \text{Tr} \left[\frac{\partial \mathbf{P}}{\partial p_\alpha} \mathbf{C}^{-1} \frac{\partial \mathbf{P}}{\partial p_\beta} \mathbf{C}^{-1} \right], \quad (45)$$

where $W_a(\ell)$ is defined as

$$W_a(\ell) = \theta \left(\ell - \left[\ell_a - \frac{\Delta \ell_a}{2} \right] \right) \theta \left(\left[\ell_a + \frac{\Delta \ell_a}{2} \right] - \ell \right), \quad (46)$$

being $\theta(x)$ is the Heaviside function.

4.3 Fiducial model and surveys specifications

The fiducial J-PAS cosmology Costa et al. (2019) assumed in our analysis is the flat Λ CDM model with the parameters $\Omega_m = 0.31$, $\Omega_b = 0.049$, $n_s = 0.96$, $h = 0.68$, $H_0^{-1} = 2997.9 \text{ Mpc } h^{-1}$, and $\sigma_8 = 0.82$ which are compatible with Planck 2018 (Aghanim et al. 2018).

For this cosmology, the $E(z)$ function defined previously is given by

$$E(z) = \sqrt{\Omega_m(1+z)^3 + (1-\Omega_m)}, \quad (47)$$

whereas the growth function can be written as

$$f_\Lambda(z) = \left(\Omega_m(1+z)^3 \frac{1}{E^2(z)} \right)^\gamma, \quad (48)$$

with the growth index $\gamma = 0.545$ (Linder & Cahn 2007). For the fiducial cosmology, the linear matter power spectrum $\hat{P}(k)$ takes the form

$$\hat{P}(k) \propto k^{n_s} T^2(k), \quad (49)$$

where the transfer function has been obtained from CAMB (Lewis et al. 2000). Then, we impose the normalization

$$\int \frac{k'^2 dk'}{2\pi^2} \hat{P}(k') |\hat{W}(8 \text{ Mpc } h^{-1}, k')|^2 = 1, \quad (50)$$

since we have taken out σ_8^2 from the power spectrum and have inserted it in the definitions (17) and (18). In the dark energy case, we will consider derivatives of the transfer function with respect to w_0 and w_a parameters when calculating the corresponding Fisher matrices. However in the modified gravity case this is no longer as the dependence of the transfer functions on the modified gravity parameters is not explicitly known. For the bias, we consider four different types of galaxies: luminous red galaxies (LRGs), emission-line galaxies (ELGs), bright galaxies (BGS), and quasars (QSO; Mostek et al. 2013; Ross et al. 2009). Each type has different fiducial bias given by

$$b(z) = \frac{b(0)}{D(z)}, \quad (51)$$

being $b_0 = 0.84$ for ELGs, $b_0 = 1.7$ for LRGs, and $b_0 = 1.34$ for BGS. For *Euclid* survey we use a fiducial bias for ELGs of the form $b(z) = \sqrt{1+z}$ (Laureijs et al. 2011), while the bias for quasars is $b(z) = 0.53 + 0.289(1+z)^2$.

Finally, we summarize the surveys specifications necessary to compute the different Fisher matrices. For clustering, we have considered: redshift bins and galaxy densities for each bin which can be found in the left-hand panel of Table A1 for J-PAS, in the centre panel of Table A1 for DESI and in the right-hand panel of Table A1 for *Euclid*. We consider two configurations of total area for J-PAS, namely 8500 and 4000 deg² which correspond to fractions of the sky of $f_{\text{sky}} = 0.206$ and $f_{\text{sky}} = 0.097$, respectively. $f_{\text{sky}} = 0.339$ for DESI with 14 000 deg² and $f_{\text{sky}} = 0.364$ for *Euclid* with 15 000 deg². The redshift error is $\delta z = 0.003$ for galaxies and QSO in J-PAS, $\delta z = 0.0005$ for galaxies in DESI and $\delta z = 0.001$ for QSO in DESI and galaxies in *Euclid*.

For the weak lensing analysis we have used: redshift bins and the fraction of the sky f_{sky} , which are the same as in the clustering analysis; mean redshifts for the galaxy density which are $z_{\text{mean}} = 0.5$ for J-PAS and $z_{\text{mean}} = 0.9$ for *Euclid*; the angular number density n_θ (in galaxies per square arcminute) that can be found in Table A8 for J-PAS with three different photometric errors. For *Euclid*, $n_\theta = 35$ galaxies per square arcminute with $\delta z = 0.05$.

5 RESULTS

5.1 Galaxy clustering

5.1.1 Dark energy

The dark energy equation of state is one of the main drivers of modern galaxy surveys. Low-redshift measurements of the scale of baryonic acoustic oscillations (BAOs) in galaxy clustering constitute a straightforward, nearly systematic-free way of measuring distances using the ‘cosmic standard ruler’ provided by the acoustic horizon at the epoch of baryon drag (Seo & Eisenstein 2003). These distances are measured both along the line of sight [since $d\chi = cdz/H(z)$] as well as across the line of sight (using the angular-diameter distance, which for an object of size dL subtending an angle $d\theta$ reads $d\theta = dL/D_A$). The different dependencies of $H(z)$ and $D_A(z)$ on cosmological parameters help break degeneracies, improving the constraints.

In order to derive these constraints, the BAOs derived from galaxy clustering must be compared against the high-redshift measurement of the acoustic horizon from observations of the cosmic microwave background (Ade et al. 2016). In terms of the Fisher matrix analysis, this means that one should include priors that codify the CMB constraints on the acoustic horizon, so we have considered from Aghanim et al. (2018) the acoustic horizon $r_{\text{drag}} = 147.18 \pm 0.29$ Mpc. Here, we chose the standard procedure of including those priors as additional Fisher matrices that are added to the full Fisher matrix (for all parameters and all slices), before slicing and eventually inverting those matrices to find the constraints.

It is important to note that one may break degeneracies and improve measurements by measuring not only the BAO features but also the shape of the power spectrum. However, since the shape measurements are much more sensitive to systematic errors than the pure BAO measurements (Seo & Eisenstein 2003; White et al. 2008), by isolating the former from the latter one obtains more robust constraints. For that reason, it has become standard practice to first derive constraints from each redshift slice on $H(z)$ and $D_A(z)$, and then project those constraints into the cosmological parameters.

It has been pointed out that the smearing of the BAO scale caused by mode-coupling in the nonlinear regime can be partially undone (at least on large scales) by the procedure known as *reconstruction* (Seo & Eisenstein 2007). For our dark energy constraints, we assume that a simple, conservative reconstruction procedure has been applied to all data sets, which would lower the non-linear scale Σ_0 from $11 h^{-1}$ Mpc to $6.5 h^{-1}$ Mpc.

The procedure for extracting constraints from BAOs while isolating as much as possible the systematics from the unknown broad-band shape of the power spectrum and non-linear effects has been well established (Seo & Eisenstein 2003). We have followed this standard procedure, which in our case means that our basic (parent) Fisher matrices include not only the ‘global’ degrees of freedom $\theta^{\text{glob}} = \{\Omega_k, \Omega_b, \Omega_c, h, n_s\}$ but also ‘local’ parameters, which are unknown on each redshift slice: $\theta^{\text{loc}} = \{H(z), D_A(z), f\sigma_8(z), b\sigma_8(z), P_{\text{shot}}(z)\}$. If there are more than one tracer available on a given slice, there are as many bias factors in that slice.

After marginalizing against every other parameter in the parent Fisher matrix, we obtain constraints for the radial and angular-diameter distances on each redshift slice (for dark energy constraints we employed slices of $\Delta z = 0.2$, and rescaled DESI and *Euclid* parameters to match that choice). Finally, the Fisher matrices in

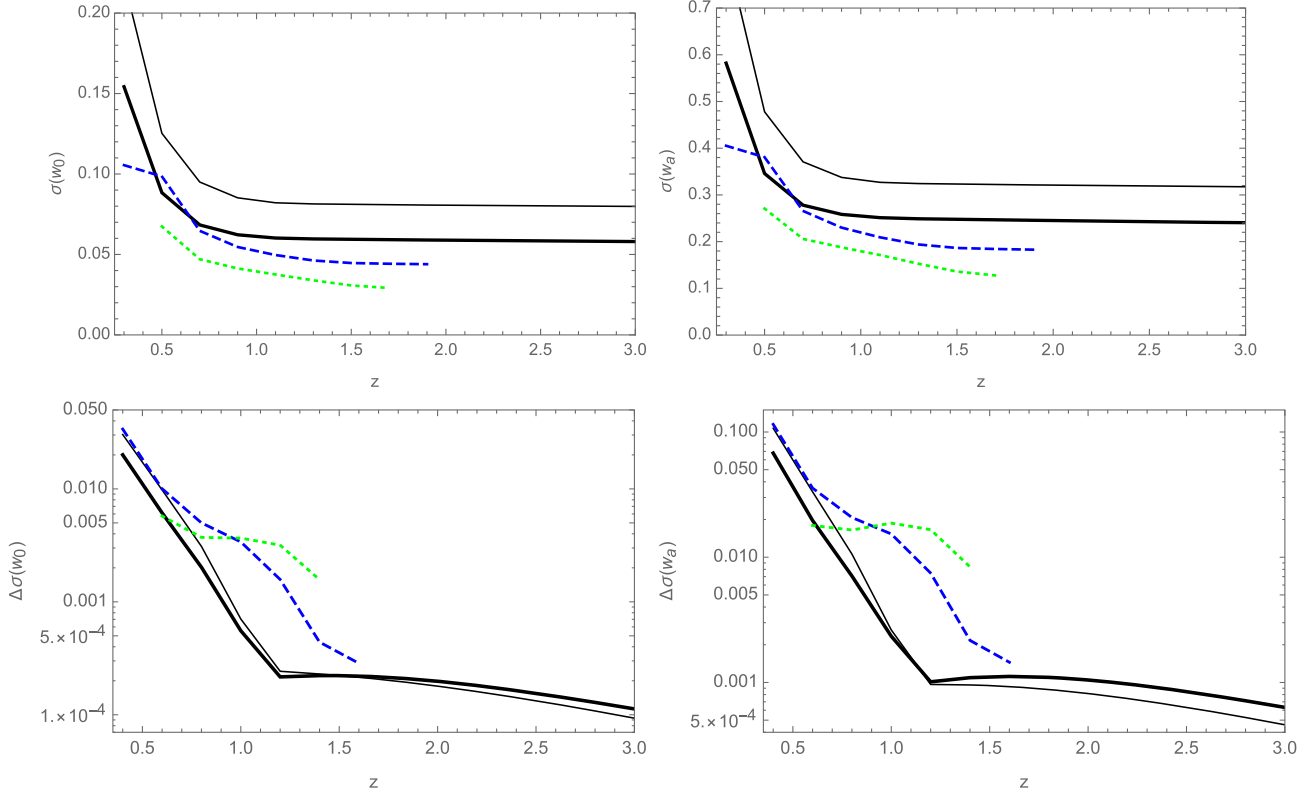


Figure 1. *Top panel:* Constraints on w_0 (left) and w_a (right) as we increase the depth of the surveys. Here, we consider only the clustering information. The errors for J-PAS (8500 deg², black solid lines; 4000 deg², black thin lines) combine ELGs, LRGs, and QSOs; those for DESI (14 000 deg², blue dashed lines) combine the BGS sample, ELGs, LRGs, and QSOs; and those for *Euclid* (15 000 deg², green dotted lines) include only ELGs. *Bottom panel:* Added value of each successive redshift slice (assuming here bins of $\Delta z = 0.2$) on w_0 (left) and w_a (right).

terms of these parameters are used to derive constraints on the desired cosmological parameters – in our case, $\{\Omega_m, w_0, w_a\}$. This last step requires that we use the BAO scale, which is imposed in terms of a suitable prior derived from Planck data.

As mentioned earlier, our model for dark energy parametrizes the equation of state using two parameters, such that $w(a) = w_0 + w_a(1 - a)$ (Chevallier & Polarski 2001; Linder 2003). The joint measurement of w_0 and w_a has been the standard metric for comparing surveys in terms of their power to constrain dark energy (Albrecht et al. 2006). In Fig. 1, we compare the constraints on w_0 and w_a for two areas of J-PAS, together with those for DESI and *Euclid*. In the top panel, we show how the constraints improve as we include successive redshift slices, and in the bottom panel, we show the added value of each successive slice for those constraints. In Fig. 2, we plot 1σ contour error for w_0 and w_a using the information of all redshift bins. We summarize the marginalized errors for w_0 and w_a in Table A2.

5.1.2 Modified gravity

For MG scenarios, we have the following independent parameters: A_i , R , and E with i denoting the different tracers. Because we have checked that marginalizing with respect to a non-Poissonian shot noise component has a minimal effect, for simplicity, we do not consider the shot noise term as a free parameter in this case. However, we are interested in obtaining errors for the effective Newton constant parameter μ and the growth function f . Thus, we first consider as parameters the dimensionless quantities A_i , R ,

and E for each redshift bin. Using the definitions of the A_i and R parameters, we obtain for $\partial P_{ij}(k_r, \hat{\mu}_r, z_a)/\partial p_\alpha$,

$$\frac{\partial P_{ij}(k_r, \hat{\mu}_r, z_a)}{\partial A_i} = \left[\frac{\delta_{ii}}{A_i + R \hat{\mu}^2} + \frac{\delta_{ij}}{A_j + R \hat{\mu}^2} \right] P_{ij}, \quad (52a)$$

$$\frac{\partial P_{ij}(k_r, \hat{\mu}_r, z_a)}{\partial R} = \left[\frac{\hat{\mu}^2}{A_i + R \hat{\mu}^2} + \frac{\hat{\mu}^2}{A_j + R \hat{\mu}^2} \right] P_{ij}, \quad (52b)$$

$$\frac{\partial P_{ij}(k_r, \hat{\mu}_r, z_a)}{\partial E} = \left[\frac{1}{E} + 2R \hat{\mu}^2 (1 - \hat{\mu}^2) \Xi + \frac{2\Delta z_a}{E^2 H_0 \chi(z_a)} \right] P_{ij} \quad (52c)$$

where

$$\Xi = \left(\frac{1}{A_i + R \hat{\mu}^2} + \frac{1}{A_j + R \hat{\mu}^2} \right) \left(\frac{1}{E} - \frac{\Delta z_a}{E^2 H_0 \chi(z_a)} \right),$$

and the length of the bin Δz_a appears since we have discretized the integration in equation (23) in order to calculate the derivative with respect to E . Following Amendola et al. (2013), in the calculation of $\partial P_{ij}(k_r, \hat{\mu}_r, z_a)/\partial E$ we do not consider the dependence of $P_{ij}(k_r, \hat{\mu}_r, z)$ on E through k since we do not know its explicit k dependence in a model-independent way.

Once we have obtained the Fisher matrix for $[A_i, R, E]$, we project first into $[A_i, f, E]$, and then to $[A_i, \mu, E]$ using equations (16) and (48) and the approximate analytic expression for $f = f(\mu, z)$ (Resco & Maroto 2018a),

$$f(\mu, z) = \frac{1}{4} \left(\sqrt{1 + 24\mu} - 1 \right) f_\Lambda(z), \quad (53)$$

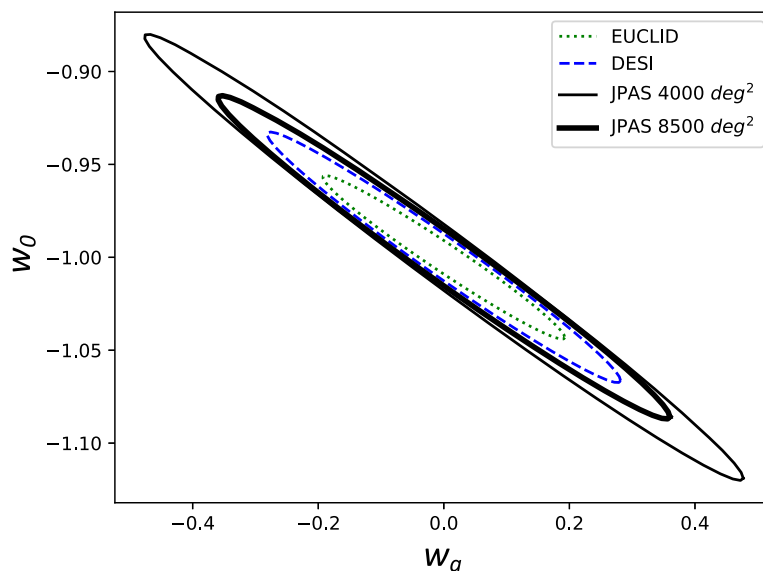


Figure 2. 1σ contour error for w_0 and w_a for J-PAS (8500 deg^2 , black solid lines; 4000 deg^2 , black thin lines) combine ELGs, LRGs, and QSOs; those for DESI (14 000 deg^2 , blue dashed lines) combine the BGS sample, ELGs, LRGs, and QSOs; and those for *Euclid* (15 000 deg^2 , green dotted lines) include only ELGs.

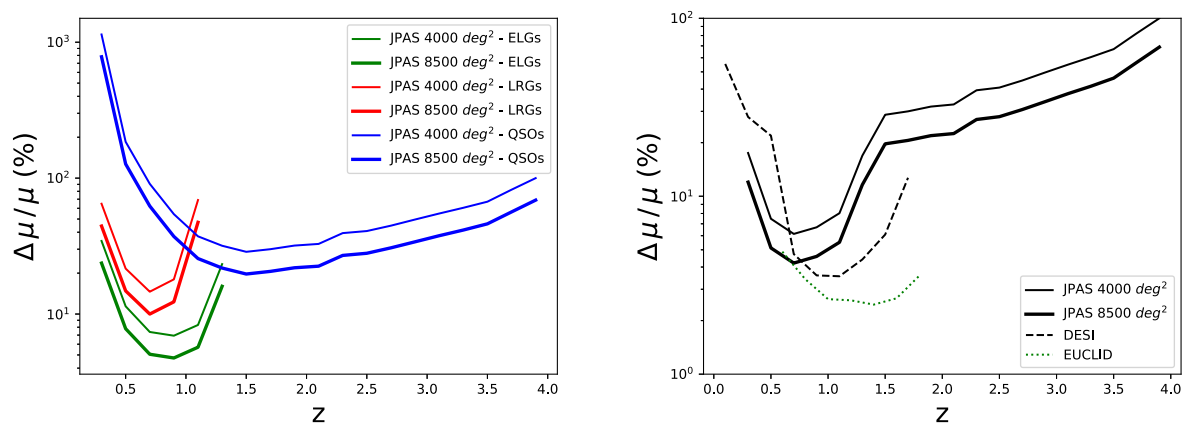


Figure 3. *Left:* Tomographic relative errors on scale independent and constant μ for J-PAS with ELGs, LRGs, and quasars. *Right:* Same as in the left plot but for J-PAS (ELGs+LRGs+QSOs), DESI (BGS+ELGs+LRGs+QSOs), and *Euclid* (ELGs) using clustering information (right).

which is valid for time-independent μ . Thus, using equation (32) we obtain the errors for f and then those for μ . Forecasts for the relative errors in μ and $f(z)$ in the different redshift bins can be found in Table A4 and in Table A5 for J-PAS, in Table A3 for DESI and in Table A6 for *Euclid*. In Fig. 3, we plot these results for the three surveys. As we can see, ELGs provide the tightest constraints for J-PAS. Compared to *Euclid* or DESI, we find that J-PAS provides the best precision in the redshift range of $z = 0.3$ – 0.6 . Notice this is also the case in the 4000 deg^2 configuration. This is mainly thanks to the large number of expected ELG detection in that redshift range that compensates the smaller fraction of sky of J-PAS as compared to other surveys.

In Fig. 4, we show $f(z)$ and $f\sigma_8(z)$ with the expected error bars. Errors for μ in different k -bins are obtained using (equation 34) and can be found in Table A7 and in Fig. 5 (left). We find that the best precision is obtained for scales around $k = 0.1 h \text{ Mpc}^{-1}$, which are slightly below *Euclid* and DESI best scales. Finally, in Fig. 7 (left) we show errors for the Hubble dimensionless parameter $E(z)$ in the different redshift bins. Once more, J-PAS provides better

precision below $z = 0.6$, but also thanks to QSOs observation at higher redshifts, J-PAS will be able to measure the expansion rate in the practically unexplored region up to redshift $z = 3.5$ with precision below 30 per cent.

5.2 Weak lensing

In this section, we obtain the errors on the η parameter using weak lensing information. First, we compute the Fisher matrix for $[E, L]$ in each bin which has the following form:

$$\begin{pmatrix} E_1 E_1 & E_1 L_1 & E_1 E_2 & E_1 L_2 & \dots \\ L_1 E_1 & L_1 L_1 & L_1 E_2 & L_1 L_2 & \dots \\ E_2 E_1 & E_2 L_1 & E_2 E_2 & E_2 L_2 & \dots \\ L_2 E_1 & L_2 L_1 & L_2 E_2 & L_2 L_2 & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}. \quad (54)$$

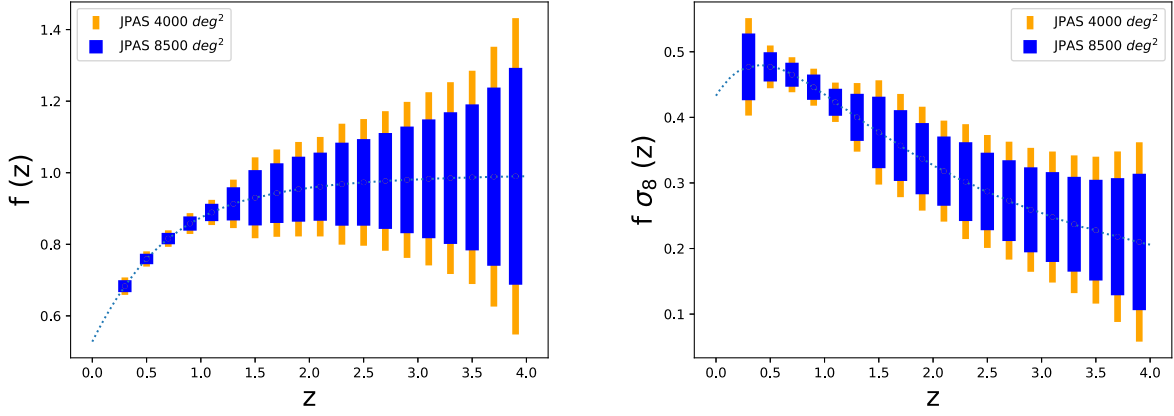


Figure 4. Growth function and $f\sigma_8$ function for the fiducial cosmology with error bars for J-PAS 8500 and 4000 deg^2 , using ELGs+LRGs+QSOs.

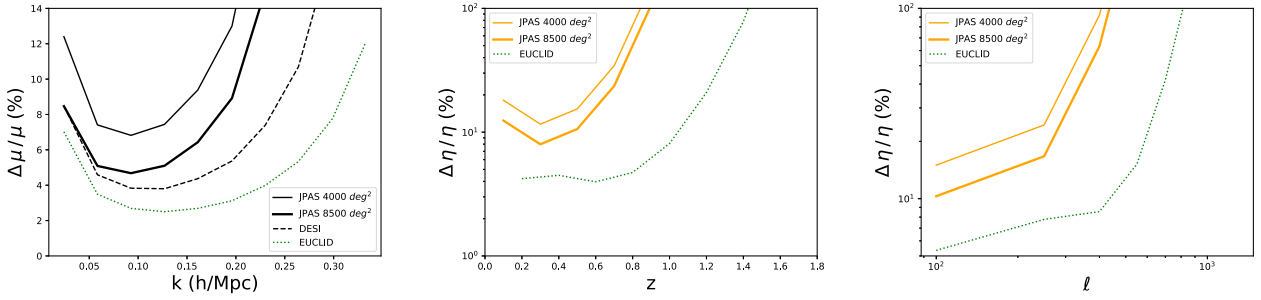


Figure 5. *Left:* Relative errors on constant $\mu(k)$ for J-PAS (ELGs+LRGs+QSOs), DESI (BGS+ELGs+LRGs+QSOs), and *Euclid* (ELGs) using clustering information. *Centre:* Tomographic relative errors of scale-independent constant η for J-PAS (ELGs+LRGs) and *Euclid* (ELGs) using lensing information. *Right:* Relative errors on constant $\eta(\ell)$ for J-PAS (ELGs+LRGs) and *Euclid* (ELGs) using lensing information.

Then, we obtain the expressions for the derivatives of the convergence power spectrum. The simplest case corresponds to the derivative with respect to L ,

$$\frac{\partial P_{ij}}{\partial L_a} = 2H_0 \frac{\Delta z_a}{E_a} K_i(z_a) K_j(z_a) L_a \hat{P} \left(\frac{\ell}{\chi(z_a)} \right). \quad (55)$$

For the derivative with respect to E , we need the expression,

$$\begin{aligned} \frac{\partial K_i(z_b)}{\partial E_a} &= \frac{3(1+z_b)\Delta z_a}{2E_a^2} \left[-\hat{\theta}(z_a - z_b)\chi(z_b) \int_{z_a}^{\infty} \frac{n_i(z')}{\chi(z')^2} dz' \right. \\ &\quad \left. + \theta(z_b - z_a) \int_{z_b}^{\infty} \left(1 - \frac{\chi(z_b)}{\chi(z')} \right) \frac{n_i(z')}{\chi(z')} dz' \right], \quad (56) \end{aligned}$$

where we have discretized the integration in equation (23) in the different bins and we have introduced Heaviside functions such that $\hat{\theta}(0) = 0$ and $\theta(0) = 1$. Then the derivative with respect to E reads

$$\begin{aligned} \frac{\partial P_{ij}}{\partial E_a} &= -H_0 \frac{\Delta z_a}{E_a^2} K_i(z_a) K_j(z_a) L_a^2 \hat{P} \left(\frac{\ell}{\chi(z_a)} \right) \\ &\quad + H_0 \sum_b \frac{\Delta z_b}{E_b} \frac{\partial K_i(z_b)}{\partial E_a} K_j(z_b) L_b^2 \hat{P} \left(\frac{\ell}{\chi(z_b)} \right) \\ &\quad + H_0 \sum_b \frac{\Delta z_b}{E_b} \frac{\partial K_j(z_b)}{\partial E_a} K_i(z_b) L_b^2 \hat{P} \left(\frac{\ell}{\chi(z_b)} \right). \quad (57) \end{aligned}$$

As in the clustering case, we have not considered derivatives of $\hat{P}(k)$.

Now, it is necessary to change the initial parameters $[E, L]$ to the new ones $[E, \eta]$. Using equation (38), we obtain $\frac{\partial \eta}{\partial L} = \frac{2}{L}$ and $\frac{\partial \eta}{\partial E} = 0$. For time-independent parameters, we show in Table A9 and in Fig. 5 (middle) the relative errors in η for the different redshift bins for J-PAS and *Euclid*. Again, J-PAS provides the best errors in the range of $z = 0.3$ – 0.6 . In order to obtain the errors of η in different ℓ -bins, we compute the Fisher matrix (equation 45). We first change from $[E, L]$ to $[E, \eta]$ in each redshift bin and then sum the information of η for the different redshift bins. The corresponding errors can be found in Table A10 for J-PAS and *Euclid* as well as in Fig. 5 (right).

5.3 Clustering+weak lensing

Finally, in this section, we analyse the case in which information from clustering and lensing is combined. We first take the Fisher matrix of parameters $[A_i, \mu, E]$ for clustering and $[E, \eta]$ for weak lensing and build the full matrix with parameters $[A_i, \mu, E, \eta]$. This matrix has the form,

$$\begin{pmatrix} A_1 A_1 & A_1 \mu_1 & 0 & A_1 E_1 & \dots \\ \mu_1 A & \mu_1 \mu_1 & 0 & \mu_1 E_1 & \dots \\ 0 & 0 & \eta_1 \eta_1 & \eta_1 E_1 & \dots \\ E_1 A_1 & E_1 \mu_1 & E_1 \eta_1 & E_1 E_1 & \dots \\ \dots & \dots & \dots & \dots & \dots \end{pmatrix}, \quad (58)$$

where EE is the sum of terms EE for clustering and lensing. By inverting this Fisher matrix, we obtain the errors for μ and η . These

results are shown in Table A11 for J-PAS and in Table A12 for *Euclid*. Finally, Fig. 6 compares the sensitivity of both surveys for time-independent μ and η in the different redshift bins. For completeness, we also show the same comparison for the function $E(z)$ in Fig. 7. As we can see, the combination of clustering and lensing information improves the sensitivity in around a 10 per cent for all the parameters. We sum all the information in the whole redshift range for μ and η and plot their error ellipses in the right-hand panel of Fig. 8. These results are summarized in Table A13.

So far we have limited ourselves to time-independent μ and η parameters. For scale-independent parameters, we consider the case in equations (14) and (15). Using the analytical fitting function for this particular expressions obtained in Resco & Maroto (2018a), we obtain errors for μ_0 and η_0 with fiducial values $\mu_0 = \eta_0 = 1$. We plot on the left-hand panel of Fig. 8 error ellipses for μ_0 and η_0 , and we summarize these errors in Table A13.

6 DISCUSSION AND CONCLUSION

Over the past years, cosmological observations have been used not only to constrain models within the context of GR but also the theory of gravity itself (see e.g. Okumura et al. 2016). In general, MG theories introduce changes in the Poisson equation which relate the density perturbations δ with the gravitational potential Ψ , thus modifying the amplitude and evolution of the growth of cosmological perturbations. Furthermore, gravitational lensing is directly sensitive to the growth of dark matter perturbations – in contrast with measurements based on galaxies, neutral hydrogen, or any other baryonic tracer. These theories, therefore, also introduce modifications in the equation that determines the lensing potential and controls the motion of photons. Thus, observations of the distribution of matter on large scales at different redshifts, and of the weak lensing generated by those structures, provide a new suite of tests of GR on cosmological scales (TsujiKawa et al. 2013; Huterer et al. 2015; Joyce et al. 2015).

In this work, we have investigated the ability of the J-PAS survey to constrain dark energy and MG cosmologies using both the J-PAS information on the galaxy power spectra for different dark matter tracers, with baryon acoustic oscillations and redshift-space distortions, as well as the weak lensing information by considering the convergence power spectrum. Our analysis considers phenomenological parametrization of dark energy and modified gravity models, as discussed in Section 3.

Following Amendola et al. (2013), we have adopted a model-independent parametrization of the power spectra of clustering and weak lensing. This parametrization considers all the free and independent parameters that are needed to describe such power spectra in the linear regime. In this analysis, we have fixed the initial dark matter power spectrum $\hat{P}(k)$ to the fiducial model, corresponding to a flat Λ CDM cosmology. As mentioned above, rather than focusing on specific dark energy or MG theories, we have considered a phenomenological approach described in terms of a set of parameters that can be contrasted with observations. Thus, in the dark energy case, the widely used (w_0, w_a) CPL parametrization has been assumed. For MG theories, two cases have been considered. First, for time-independent μ and η , we have performed both a tomographic redshift bin analysis and an analysis in k -bins. By summing over all the redshift range, we have obtained the best errors for the modified gravity parameters. Secondly, for scale-independent parameters, we have considered the particular

parametrization in terms of μ_0 and η_0 (equations 14 and 15) usually employed in the literature.

J-PAS will be able to measure different tracers, e.g. LRG, ELG, and QS. In order to contextualize the J-PAS measurements, we have performed the same Fisher analysis for DESI and *Euclid* surveys. In the case of DESI, in addition to LRGs, ELGs, and QSOs, a bright galaxy sample (BGS) will be also measured at low redshifts, while *Euclid* will measure only ELGs. In the dark energy analysis, we have found that J-PAS will measure w_0 with precision below 6 per cent that can be compared with the 4.5 per cent for DESI and 3 per cent for *Euclid*. The absolute error in w_a is found to be below 0.24 for J-PAS, 0.19 for DESI, and 0.13 for *Euclid*. From the tomographic analysis, we find that using the clustering information alone, J-PAS will allow to measure the expansion rate $H(z)$ with precision 3 per cent in the best redshift bin ($z = 0.7$) and the μ parameter with a precision around 5 per cent in the best redshift bin. From lensing alone, J-PAS will be able to measure η with a precision around 8 per cent in the best redshift bin. The combination of clustering and lensing will allow to improve the precision in μ down to 4 per cent in the best bin. Considering the information in the whole redshift range, we have found that J-PAS will be able to measure time-independent μ and η with precision better than 3 per cent for both parameters. For μ_0 and η_0 , we have obtained errors of 10 per cent and 5 per cent, respectively.

When compared to future spectroscopic surveys such as DESI or spectroscopic and photometric ones such as *Euclid*, we have shown that from clustering and lensing information, J-PAS will have the best errors for redshifts between $z = 0.3$ and 0.6, thanks to the large number of ELGs detectable in that redshift range. Note also that thanks to QSOs observation at higher redshifts, J-PAS will be able to measure the expansion rate and MG parameters in the practically unexplored region up to redshift $z = 3.5$ with precision below 30 per cent.

In the whole redshift range, the J-PAS precision in both μ and η will be a factor of 1.5–2 below *Euclid* in their respective best bins. For the (time-dependent) μ_0 – η_0 parametrization (equations 14 and 15), we have shown that J-PAS is closer to *Euclid* than in the constant case. This is due to the fact that low-redshift measurements are more sensitive to μ_0 and η_0 than high-redshift ones, such that at low-redshift J-PAS precision surpasses that of *Euclid*.

Finally, it is worth mentioning that by increasing the precision in the determination of the dimensionless Hubble parameter using e.g. the J-PAS sample of type Ia supernovae, and taking into account information from the non-linear power spectra, it can be expected that the sensitivity to the μ and η parameters will increase. Additionally, considering the cross-correlation between galaxy distribution and galaxy shapes will also allow to improve the precision of J-PAS in the determination of dark energy and MG parameters.

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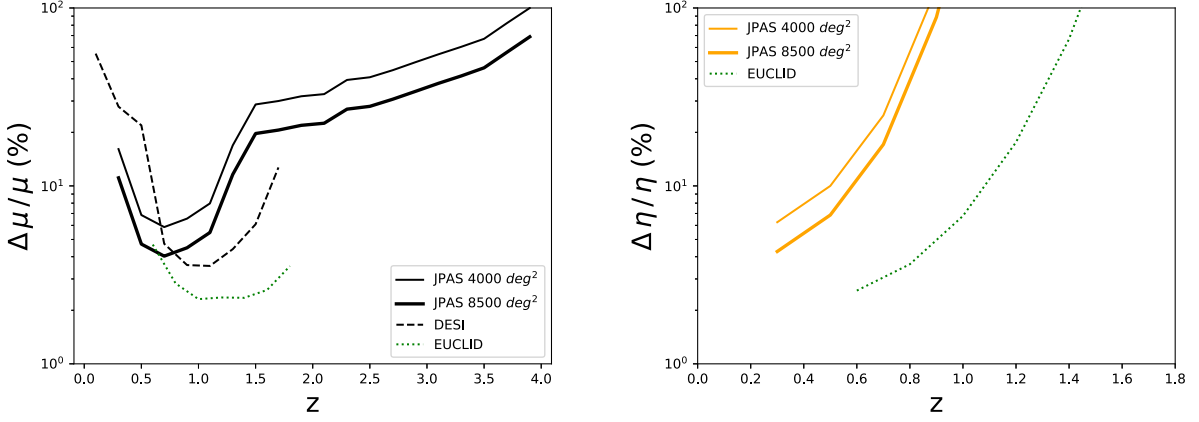


Figure 6. *Left:* Tomographic relative errors for scale-independent constant μ for J-PAS (ELGs+LRGs+QSOs), DESI (BGS+ELGs+LRGs+QSOs), and *Euclid* (ELGs) using clustering and lensing information. In the case of DESI and J-PAS quasars only clustering information is taken into account. *Right:* Tomographic relative errors for scale-independent constant η for J-PAS (ELGs+LRGs+QSOs) with $\delta z = 3$ per cent and *Euclid* (ELGs) using clustering and lensing information.

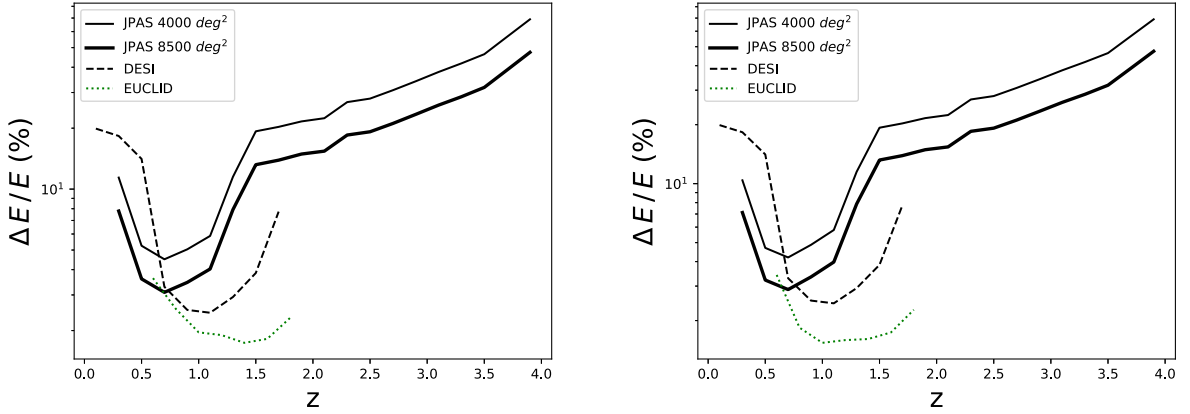


Figure 7. Relative errors for $E(z)$ for J-PAS (ELGs+LRGs+QSOs), DESI (BGS+ELGs+LRGs+QSOs), and *Euclid* (ELGs) using clustering information (left-hand panel), and using clustering and lensing information (right-hand panel). In the case of DESI and J-PAS quasars, only clustering information is taken into account. For lensing in J-PAS, the redshift error is $\delta z = 3$ per cent.

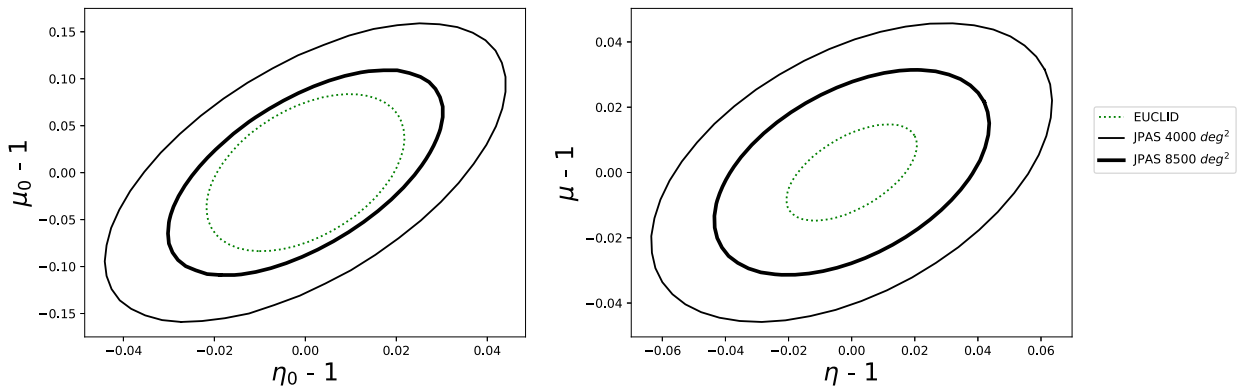


Figure 8. *Left:* 1σ contour error for μ_0 and η_0 defined in equations (14) and (15). *Right:* 1σ contour error for scale-independent constant μ and η . All in J-PAS (ELGs+LRGs+QSOs) and *Euclid* (ELGs) surveys combining clustering and lensing information for 8500 and 4000 deg^2 .

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APPENDIX: DATA TABLES

See Tables [A1–A13](#).

Table A1. In left-hand panel: Redshift bins and densities of luminous red galaxies, emission-line galaxies, and quasars for J-PAS. In centre panel: Redshift bins and densities of bright galaxies, luminous red galaxies, emission-line galaxies and quasars for DESI. In right-hand panel: Redshift bins and densities of emission-line galaxies for *Euclid*. Galaxy densities in units of $10^{-5} \text{ h}^3 \text{ Mpc}^{-3}$.

z	J-PAS		
	LRG	ELG	QSO
0.3	226.6	2958.6	0.45
0.5	156.3	1181.1	1.14
0.7	68.8	502.1	1.61
0.9	12.0	138.0	2.27
1.1	0.9	41.2	2.86
1.3	0	6.7	3.60
1.5	0	0	3.60
1.7	0	0	3.21
1.9	0	0	2.86
2.1	0	0	2.55
2.3	0	0	2.27
2.5	0	0	2.03
2.7	0	0	1.81
2.9	0	0	1.61
3.1	0	0	1.43
3.3	0	0	1.28
3.5	0	0	1.14
3.7	0	0	0.91
3.9	0	0	0.72

z	DESI			
	BGS	LRG	ELG	QSO
0.1	2240	0	0	0
0.3	240	0	0	0
0.5	6.3	0	0	0
0.7	0	48.7	69.1	2.75
0.9	0	19.1	81.9	2.60
1.1	0	1.18	47.7	2.55
1.3	0	0	28.2	2.50
1.5	0	0	11.2	2.40
1.7	0	0	1.68	2.30

<i>Euclid</i>	
z	ELG
0.6	356
0.8	242
1.0	181
1.2	144
1.4	99
1.6	66
1.8	33

Table A2. Absolute errors for w_0 and w_a for *Euclid*, DESI, and JPAS (with 8500 and 4000 deg^2), considering clustering information.

Survey	Δw_0	Δw_a
<i>Euclid</i>	0.029	0.128
DESI	0.045	0.186
J-PAS 8500	0.058	0.238
J-PAS 4000	0.079	0.316

Table A3. Redshift bins, fiducial values for μ and f , and their errors for DESI forecast with clustering information, using BGS+ELGs+LRGs+QSOs. Relative errors are per cent errors.

z	DESI clustering				
	μ	$\Delta\mu/\mu$	f	Δf	$\Delta f/f$
0.1	1	55.4	0.585	0.085	14.5
0.3	1	27.9	0.683	0.037	5.47
0.5	1	21.9	0.759	0.048	6.32
0.7	1	4.73	0.816	0.016	1.96
0.9	1	3.59	0.858	0.014	1.62
1.1	1	3.55	0.889	0.014	1.58
1.3	1	4.41	0.913	0.017	1.87
1.5	1	6.09	0.930	0.022	2.40
1.7	1	12.7	0.943	0.044	4.66

Table A4. Redshift bins, fiducial values for μ and f , and their errors for J-PAS forecast with clustering information, 4000 deg² and using ELGs+LRGs+QSOs. Relative errors are per cent errors.

z	μ	J-PAS clustering 4000 deg ²						
		$\Delta\mu/\mu$	f	Δf	$\Delta f/f$	$f\sigma_8$	$\Delta f\sigma_8$	$\Delta f\sigma_8/f\sigma_8$
0.30	1	17.5	0.683	0.024	3.57	0.477	0.074	15.6
0.50	1	7.47	0.759	0.021	2.81	0.477	0.033	6.83
0.70	1	6.14	0.816	0.023	2.84	0.465	0.027	5.75
0.90	1	6.69	0.858	0.029	3.39	0.446	0.028	6.33
1.10	1	8.03	0.889	0.035	3.96	0.423	0.030	7.10
1.30	1	16.9	0.913	0.068	7.42	0.400	0.052	13.1
1.50	1	28.7	0.930	0.113	12.1	0.377	0.080	21.1
1.70	1	30.0	0.943	0.122	12.9	0.357	0.079	22.1
1.90	1	31.9	0.954	0.132	13.9	0.337	0.079	23.5
2.10	1	32.8	0.961	0.139	14.4	0.318	0.077	24.2
2.30	1	39.4	0.968	0.169	17.4	0.302	0.088	29.0
2.50	1	40.8	0.973	0.177	18.2	0.287	0.086	30.0
2.70	1	44.7	0.977	0.195	20.0	0.273	0.090	33.0
2.90	1	49.6	0.980	0.218	22.2	0.259	0.094	36.5
3.10	1	54.9	0.983	0.242	24.7	0.248	0.100	40.4
3.30	1	60.5	0.985	0.268	27.2	0.237	0.105	44.4
3.50	1	67.1	0.987	0.298	30.2	0.228	0.112	49.2
3.70	1	82.2	0.989	0.363	36.7	0.218	0.130	59.7
3.90	1	100	0.990	0.442	44.6	0.210	0.152	72.5

Table A5. Redshift bins, fiducial values for μ and f and their errors for J-PAS forecast with clustering information, 8500 deg² and using ELGs+LRGs+QSOs. Relative errors are per cent errors.

z	μ	J-PAS clustering 8500 deg ²						
		$\Delta\mu/\mu$	f	Δf	$\Delta f/f$	$f\sigma_8$	$\Delta f\sigma_8$	$\Delta f\sigma_8/f\sigma_8$
0.30	1	12.0	0.683	0.017	2.45	0.477	0.051	10.7
0.50	1	5.12	0.759	0.015	1.93	0.477	0.022	4.68
0.70	1	4.21	0.816	0.016	1.95	0.465	0.018	3.95
0.90	1	4.59	0.858	0.020	2.32	0.446	0.019	4.34
1.10	1	5.51	0.889	0.024	2.72	0.423	0.021	4.87
1.30	1	11.6	0.913	0.046	5.09	0.400	0.036	8.97
1.50	1	19.7	0.930	0.077	8.32	0.377	0.055	14.5
1.70	1	20.6	0.943	0.083	8.84	0.357	0.054	15.1
1.90	1	21.9	0.954	0.091	9.52	0.337	0.054	16.1
2.10	1	22.5	0.961	0.095	9.90	0.318	0.053	16.6
2.30	1	27.0	0.968	0.116	12.0	0.302	0.060	19.9
2.50	1	28.0	0.973	0.121	12.5	0.287	0.059	20.6
2.70	1	30.7	0.977	0.134	13.7	0.273	0.062	22.6
2.90	1	34.0	0.980	0.149	15.2	0.259	0.065	25.0
3.10	1	37.7	0.983	0.166	16.9	0.248	0.068	27.7
3.30	1	41.5	0.985	0.184	18.6	0.237	0.072	30.4
3.50	1	46.1	0.987	0.204	20.7	0.228	0.077	33.7
3.70	1	56.4	0.989	0.249	25.2	0.218	0.089	41.0
3.90	1	68.9	0.990	0.303	30.6	0.210	0.104	49.8

Table A6. Redshift bins, fiducial values for μ and f and their errors for *Euclid* forecast with clustering information, using ELGs. Relative errors are per cent errors.

z	μ	<i>Euclid</i> clustering			
		$\Delta\mu/\mu$	f	Δf	$\Delta f/f$
0.6	1	4.88	0.789	0.017	2.12
0.8	1	3.42	0.838	0.014	1.65
1.0	1	2.64	0.875	0.012	1.32
1.2	1	2.60	0.902	0.012	1.31
1.4	1	2.46	0.922	0.011	1.19
1.6	1	2.67	0.937	0.012	1.23
1.8	1	3.58	0.949	0.014	1.50

Table A7. Centres of bins k_a in units of h^{-1} Mpc, fiducial values for μ and their relative errors (per cent) for *Euclid* forecast using ELGs, DESI forecast using BGS+ELGs+LRGs+QSOs, and J-PAS forecast using ELGs+LRGs+QSOs with 8500 and 4000 deg^2 . All for clustering information.

k	μ	<i>Euclid</i>	DESI	$\Delta\mu/\mu$	
				JPAS 8500 deg^2	JPAS 4000 deg^2
0.024	1	7.02	8.48	8.47	12.4
0.058	1	3.49	4.59	5.09	7.41
0.093	1	2.69	3.83	4.68	6.82
0.127	1	2.50	3.80	5.10	7.44
0.161	1	2.69	4.37	6.43	9.38
0.196	1	3.12	5.37	8.92	13.0
0.230	1	3.99	7.39	15.0	21.8
0.264	1	5.34	10.7	29.6	43.2
0.299	1	7.78	17.6	67.6	98.6
0.333	1	1.21	32.6	153	223

Table A8. n_θ values for J-PAS with different galaxies and redshift errors in galaxies per square arcminute.

δz	n_θ values for J-PAS		
	LRG	ELG	LRG+ELG
0.003	0.52	2.48	3.00
0.01	2.02	6.21	8.23
0.03	3.25	9.07	12.32

Table A9. Redshift bins, ℓ_{max} values, fiducial values for η , and relative errors (per cent). In left table, errors for J-PAS, using LRG+ELG galaxies with $\delta z = 0.03$. We show only errors using ELG+LRG and lensing information. In right table, errors for *Euclid* using lensing information.

z	ℓ_{max}	η	J-PAS lensing	
			8500 deg^2	4000 deg^2
0.1	40	1	12.4	18.1
0.3	130	1	7.98	11.6
0.5	238	1	10.6	15.4
0.7	366	1	23.6	34.4
0.9	514	1	106	154
1.1	686	1	–	–
1.3	884	1	–	–

z	<i>Euclid</i> lensing		
	ℓ_{max}	η	$\Delta\eta/\eta$
0.2	83	1	4.21
0.4	182	1	4.48
0.6	300	1	3.97
0.8	437	1	4.72
1.0	597	1	8.10
1.2	782	1	20.9
1.4	994	1	78.3
1.6	1240	1	490
1.8	1510	1	–

Table A10. Centres of bins ℓ_a , fiducial values for η and relative errors (per cent) for J-PAS, using LRG+ELG galaxies with $\delta z = 0.03$ and for *Euclid* using lensing information.

ℓ	η	$\Delta\eta/\eta$		
		<i>Euclid</i>	8500 deg^2	4000 deg^2
100	1	5.35	10.3	15.0
250	1	7.78	16.7	24.4
400	1	8.55	63.3	92.3
550	1	15.2	360	524
700	1	42.1	–	–
850	1	130	–	–
1000	1	176	–	–

Table A11. Redshift bins, relative errors (per cent) for η , μ , and $E(z)$ for J-PAS considering clustering and lensing information (with $\delta z = 3$ per cent and ELGs+LRGs+QSOs).

z	J-PAS clustering +lensing					
	$\Delta\eta/\eta_{8500}$	$\Delta\eta/\eta_{4000}$	$\Delta\mu/\mu_{8500}$	$\Delta\mu/\mu_{4000}$	$\Delta E/E_{8500}$	$\Delta E/E_{4000}$
0.3	4.28	6.25	11.1	16.1	7.12	10.4
0.5	6.86	10.0	4.71	6.86	3.22	4.70
0.7	17.1	24.9	4.03	5.87	2.88	4.20
0.9	88.8	129	4.49	6.55	3.34	4.87
1.1	–	–	5.47	7.97	3.98	5.80
1.3	–	–	11.6	16.9	7.88	11.5

Table A12. Redshift bins, relative errors (per cent) for η , and μ for *Euclid*, considering clustering and lensing information.

z	<i>Euclid</i> clustering +lensing		
	$\Delta\eta/\eta$	$\Delta\mu/\mu$	$\Delta E/E$
0.6	2.58	4.68	3.42
0.8	3.63	2.83	1.84
1.0	6.78	2.31	1.54
1.2	17.6	2.36	1.59
1.4	66.9	2.35	1.61
1.6	415	2.60	1.74
1.8	–	3.54	2.27

Table A13. Relative errors (per cent) for constant μ and η , and μ_0 and η_0 for *Euclid* and JPAS (with 8500 and 4000 square degrees), considering clustering and lensing information.

Survey	$\Delta\mu/\mu$	$\Delta\eta/\eta$	$\Delta\mu_0/\mu_0$	$\Delta\eta_0/\eta_0$
<i>Euclid</i>	0.98	1.37	7.13	3.38
J-PAS 8500	2.08	2.89	9.66	4.58
J-PAS 4000	3.03	4.21	14.1	6.68

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