# Scattering Parameters of VHF Semiconductor Devices 

Young Dae Kim

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# SCATTERING PARANETERS OF VHF <br> SEMICONDUCTOR DEVICES 

BY YOUNG DAE KIM

A thesis submitted
in partial fulfillment of the requirements for the degree Master of Science, Department of Electrical Engineering, South Dakota

State University

January, 1970

# SCATTERING PARAMETERS OF VHF SEMICONDUCTOR DEVICES 

This thesis is approved as a creditable and independent investigation by a candidate for the degree, Master of Science, and is acceptable as meeting the thesis requirements for this degree, but without implying that the conclusions reached by the candidate are necessarily the conclusions of the major department.
' Thesis Adviser Date
Head, Electrical
Engineering Department Date

## ACKNOWLEDGMENTS

The author wishes to express his appreciation and gratitude to Dr. Virgil G. Ellerbruch and Dr. Franklin C. Fitchen, whose guidance and advice made this investigation possible, and to the National Science Foundation for partial financial support of this investigation.
Y.D.K.

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## CHAPTER I

## IN'TRODUC'TION

Since Campbell and Foster ${ }^{1}$ first used scattering parameters in studying the properties of ideal transformer networks, much work has been done with scattering parameters in the analysis of microwave circuits and general lumped parameter networks. An excellent summary for microwave circuits appears in Montgomery, Dicke, and Purcell. Application of scattering parameters to network synthesis was attempted by Oono and Yasuura. ${ }^{3}$ Youla ${ }^{4}$ also extended scattering parameter theory by complex normalization. He and Penfield ${ }^{5}$ later applied scattering parameters in analyzing negative resistance amplifiers in conjunction with the development of the tunnel diode.

The rapid development of high-frequency technologies in the past decade requires improved high-frequency measurement techniques because of the difficulty in measuring commonly-accepted immittance parameters at frequencies above 100 MHz . This difficulty stemmed from the fact that in measuring the $\mathrm{h}-, \mathrm{y}$-, or z -parameters the circuit is either in a short or open condition. The scattering parameter measurement technique is one way of overcoming this problem since the measurement circuit employs finite terminations and therefore provides more stable wide-band
measurements.
Along with the introduction of the scattering parameter measurement technique in high-frequency transistor measurements, attempts were made to characterize the transistors with directly measurable scattering parameters. Lange, ${ }^{6}$ Weinert, ${ }^{7}$ Anderson, ${ }^{8}$ Froehner, ${ }^{9}$ and Bodway ${ }^{10}$ are among those engineers who introduced the use of scattering parameters in transistor circuit design. An excellent analysis of transistor circuit design with generalized 2-port scattering parameters is given by Bodway. At the present time, the scattering parameter technique is one of the standard methods used for high-frequency transistor characterization and design.

The major advantages of the scattering parameter method for high-frequency transistor measurements and characterization are:

1. A very stable wide-band measurement, which allows wide-band swept-frequency measurement, can be made.
2. Matching networks are also measured in terms of scattering parameters for reasons of simplicity at low frequencies, and because of necessity at high frequencies. Thus, unified circuit design is possible with easily measurable scattering parameters.
3. Scattering parameter circuit design is more closely
related to power relations than any other design method. This gives a clearer understanding of amplifier design procedures.
4. The signal flow graph ${ }^{11}$ can be easily employed to visualize the scattering parameter design procedures.

In addition to these advantages, the scattering parameter method has a few disadvantages. Since they are parameters derived for traveling waves, voltage and current relationships are not easily seen. Engineers are not as familiar with scattering parameters as they are with immittance parameters. These disadvantages can be overcome when one becomes familiar with scattering parameters and can relate them to existing immittance parameters of the equivalent circuits of the device.

As an effort to fill the gap between conventional methods and the scattering parameter method in high-frequency semiconductor device measurement and characterization, this paper aims at the following goals:

1. To assemble a measuring system to accurately measure 2-, 3-, and 4-port scattering parameters of semiconductor devices at VHF frequencies.
2. To measure scattering parameters of bipolar transistors, single- and dual-gate MOSFETs, and an IC differential amplifier.
3. To explain the behavior of measured scatterinc parameters of high-frequency semiconductor devices with their equivalent circuits.
4. To apply the scattering parameter method in obtaining equivalent circuit parameters.
5. To analyze mathematically and experimentally the scattering parameter relations between a dual-gate MOSFET and a single-gate MOSFET, and between an IC differential amplifier and its constituent transistors.

Information of this kind has not appeared in the literature. In addition to these, scattering parameters are introduced from the classical transmission line theory and expanded utilizing convenient signal flow graphs. Mixer scattering parameters are also first suggested in this paper as an extension of linear 2-port parameters.

The measurements were made with a modified standard set upl2 at the Microwave Laboratory of the South Dakota State University.

## CHAPTER II

SCATTERING PARAMETERS
A. Introduction to S-Parameters ${ }^{13}$

The voltages and currents on a transmission line are generally analyzed as two traveling waves, reflected and incident. These two waves propagate in opposite directions along the line.

When this analysis method is used for a linear n-port network where the access terminals of each port are connected to a transmission line that is equipped with a device designed to sample these two waves, the terminal voltages and currents at each port can be represented by the amplitudes and phase angles.

Voltage and current equations for the ith port can be written as

$$
\begin{align*}
& v_{i}=v_{i}^{+}+v_{i}^{-} \\
& i_{i}=i_{i}^{+}-i_{i}^{-} \tag{2-1}
\end{align*}
$$

The superscripts (+) and (-) stand for inward (toward the network) and outward (away from the network) traveling waves at the ith port, respectively. $V_{i}$ and $i_{i}$ are the terminal voltage and current at the ith port. (Refer to Fig. 2-1.)

The characteristic impedance of a transmission line


Fig. 2-1. An n-port network.


Fig. 2-2. A model for $E q$. (2-7).
is defined as

$$
\begin{equation*}
z_{o i}=\frac{v_{i}^{+}}{i_{i}}=\frac{v_{i}^{-}}{i_{i}^{-}} \tag{2-2}
\end{equation*}
$$

Impedance, $Z_{\text {oi }}$, relates the voltage and current waves on the it port transmission line.

Eq. (2-1) and (2-2) can be simultaneously solved in terms of $v_{i}$ and $i_{i}$ with the following results:

$$
\begin{align*}
& v_{i}^{+}=\frac{1}{2}\left(v_{i}+Z_{o i^{i}}\right) \\
& v_{i}^{-}=\frac{1}{2}\left(v_{i}-Z_{o i_{i}}\right) \tag{2-3}
\end{align*}
$$

Examination of Eq. (2-3) shows that they are evidently a linear transformation from the terminal variables $v_{i}$ and $i_{i}$ to a new set of variables, $\mathbf{v}_{\mathbf{i}}{ }^{+}$and $\mathbf{v}_{\mathbf{i}}{ }^{-}, Z_{o i}$ is an rbitracy constant with the dimension of impedance and can be defined as the normalization impedance.

For a linear n-port network, the new port variables, $\mathbf{v}_{\mathbf{i}}{ }^{+}$and $\mathbf{v}_{\mathbf{i}}{ }^{-}$, are linearly related and the relationship can be expressed in a compact form as follows:

$$
\left(\begin{array}{c}
v_{1}{ }^{-} \\
v_{2}{ }^{-} \\
\ldots \\
v_{n}
\end{array}\right)=\left(\begin{array}{lllll}
s_{11} & s_{12} & s_{13} & \ldots & s_{1 n} \\
s_{21} & s_{22} & s_{23} & \ldots & s_{2 n} \\
\ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots \ldots . . \\
s_{n 1} & s_{n 2} & s_{n 3} & \ldots & s_{n n}
\end{array}\right) \quad\left(\begin{array}{c}
v_{1}^{+} \\
v_{2}^{+} \\
\ldots . \\
v_{n}
\end{array}\right)
$$

or $\quad \overline{\mathrm{V}}^{-}=\overline{\mathrm{S}} \overline{\mathrm{V}}^{+}$
$\overline{\mathrm{V}}^{-}$and $\overline{\mathrm{V}}^{+}$are column matrices with components $\mathrm{v}_{\mathbf{i}}{ }^{-}$and $\mathbf{v}_{\mathbf{i}}{ }^{+}$, respectively, and $\overline{\mathrm{S}}$ is a square matrix with components $S_{i j}$. The $S_{i j}$ components will be defined in the following discussion.

The constant matrix $\overline{\mathrm{S}}$ is defined as the scattering matrix and each component in the array is a scattering aramter or simply an S-parameter. Eq. (2-3) show that the scattering parameters are a function of the network and the impedance $Z_{01}$. For practical applications, it is convenient to work with real positive impedances for all ports such that

$$
\begin{equation*}
z_{o 1}=z_{o 2}=\ldots=z_{\text {on }}=z_{o} \tag{2-5}
\end{equation*}
$$

Furthermore, it turns out to be more natural to normalive $\overline{\mathrm{V}}^{+}$and $\overline{\mathrm{V}}^{-}$with respect to $\sqrt{L_{0}}$. Later, it is justified that this choice for a normalizing term gives clearer network power relationships without affecting the S-parameters.

These normalized voltage waves, which will be expressed as power waves from now on, have the same form as Eqs. (2-4). i.e.,

$$
\begin{equation*}
\bar{B}=\overline{\mathrm{S}} \overline{\mathrm{~A}} \tag{2-6}
\end{equation*}
$$

In this equation, $\bar{B}$ and $\bar{A}$ are scattered and incident power wave matrices whose components $b_{i}$ and $a_{i}$ are defined by

$$
\begin{align*}
& a_{i}=\frac{1}{2 \sqrt{2_{o}}}\left(v_{i}+z_{o} i_{i}\right) \\
& b_{i}=\frac{1}{2 \sqrt{z_{o}}}\left(v_{i}-z_{\rho_{i}}\right) \tag{2-7}
\end{align*}
$$

Although these relationships are derived from a transmission line voltage wave concept, they can also be verified by lumped-element circuit theory. 14

By referring to Fig. 2-2 and using Kirchhoff's voltage law, the following equations are written

$$
\begin{aligned}
& \mathbf{v}_{\mathbf{g}}=\mathbf{v}_{i}+z_{o_{i}} i_{i} \\
& \mathbf{v}_{r}=v_{i}-z_{o} i_{i}
\end{aligned}
$$

These values are substituted into Eq. (2-7) with the result

$$
\begin{align*}
& \mathrm{a}_{\mathrm{i}}=\frac{\mathbf{v}_{\mathbf{s}}}{2 \sqrt{2_{0}}} \\
& \mathrm{~b}_{\mathrm{i}}=\frac{\mathrm{v}_{\mathrm{r}}}{2 \sqrt{2_{0}}} \tag{2-8}
\end{align*}
$$

From this, the following definitions are made:
$\left|a_{i}\right|^{2,}$ available power from source which has an impedance of $2_{0}$.
$\left|b_{i}\right|^{2}$, power that flows from the eth port of a network
when the port is terminated in $\mathrm{Z}_{0}$.
$\left|a_{i}\right|^{2}-\left|b_{i}\right|^{2}=\operatorname{Re}\left(v_{i} i_{i}{ }^{*}\right)$
1 power delivered to network through fth port, if positive; power received through ith port, if negative.

Accordingly, the S-parameters of the matrix are defined as follows:

$$
\begin{align*}
& s_{i i}=\left.\frac{b_{i}}{a_{i}}\right|_{a_{j}}=0, j=1,2, \ldots, n, \text { but } \neq i . \\
& S_{i j}=\left.\frac{b_{i}}{a_{j}}\right|_{a_{k}}=0, k=1,2, \ldots, n, \text { but } \neq j . \tag{2-9}
\end{align*}
$$

$S_{i i}$ and $S_{i j}$ are voltage reflection and transmission coefficients. The squares of the magnitude of $S_{i i}$ and $S_{i j}$ are more useful quantities. They are defined as follows:

$$
\begin{aligned}
& \left|S_{i i}\right|^{2}=\frac{\left|b_{i}\right|^{2}}{\left|a_{i}\right|^{2}}: \begin{array}{l}
\text { power reflection coefficient } \\
\text { at } i t h \text { port }
\end{array} \\
& \left|S_{i j}\right|^{2}=\frac{\left|b_{i}\right|^{2}}{\left|a_{j}\right|^{2}} \quad \begin{array}{l}
\text { transducer power gain from the } \\
j t h \text { port to the ith port. }
\end{array}
\end{aligned}
$$

It should be noted that all of the preceding discussions are based on the assumption that the normalization impedance is $Z_{o}$ in every instance.

For the case of arbitrary source and termination impedances, new equations must be derived. These functions are derived through the use of signal flow graphs.

From the 1-port network of Fig. 2-3, calculation of the input reflection coefficient yields

$$
\begin{equation*}
S_{11}=\frac{Z-Z_{0}}{Z+Z_{0}}=\frac{Y_{0}-Y}{Y_{0}+Y} \tag{2-10}
\end{equation*}
$$

where $Y_{0}$ and $Y$ are reciprocals of $Z_{0}$ and $Z_{\text {. }}$
For the definition of Eq. (2-10) the source impedance is assumed to be $Z_{0}$, so there are no mismatch reflections, and $\mathrm{a}_{\mathrm{i}}$ is solely determined by $\mathbf{v}_{\mathrm{s}}$ without considering $\mathrm{b}_{\mathrm{i}}$ • However, for the case of an arbitrary source impedance $Z_{s}$, there are two kinds of incident waves: one generated by the source voltage and the other is the wave reflected at the source impedance. In Fig. 2-4, this relationship is shown for a 1-port network:
where

$$
\begin{align*}
& a_{1}=a_{s}+s_{s} b_{1}  \tag{2-11}\\
& a_{s}=\frac{v_{s} / \longdiv { z _ { o } }}{z_{s}+z_{o}}
\end{align*}
$$

and $S_{s}$ is the source reflection coefficient defined by Eq. (2-10) when $Z$ is replaced by $Z{ }^{2}$. The signal flow diagram for the 1-port network is shown in Fig. 2-5 (a).

(a)

(b)

Fig. 2-3. (a) Impedance and (b) Admittance representations for 1 -port network.


Fig. 2-4. 1-port network with arbitrary source impedance.


- (a)


Fig. 2-5. Signal flow graphs: (a) 1-port:


Fig. 2-6. 2-port signal flow graph without
port 2 generator.

From the 1-port signal flow graph, the 2-port signal flow graph can be developed as illustrated in Fig. 2-5 (b). From the flow graphs, the desired S-parameter relationship can be derived using a node absorption formula or Mason's rule. These derivations are given in later sections. B. Two-Port Formulation 8,10

In this section, S-parameters for 2-port networks are studied employing signal flow graphs. Since well analyzed 2-port formulations are documented in the literature, only those relationships essential to this thesis will be presented.

In the previous section, a 2 -port signal flow graph was derived from the 1 -port scattering variable relationships. Referring to Fig. 2-5 (b) it is seen that the scattering relations of 2 -port networks are clear.

In Fig. 2-5 (b), port 1 is assumed as the input port and port 2 as the output port. No signal generator is connected to port 2 and $a_{s 2}$ node disappears from the graph. $S_{s}$ and $S_{1}$ denote the reflection coefficients at source and load impedances, respectively.

With the relationships between power wave variables, $a^{\prime} s$ and b's, and the S-parameters which are shown in the flow graph of Fig. 2-6, the following set of equations is written

$$
\begin{align*}
& b_{1}=s_{11} a_{1}+s_{12} a_{2} \\
& b_{2}=s_{21} a_{1}+s_{22} a_{2}  \tag{2-12}\\
& a_{1}=a_{s}+s_{s} b_{1} \\
& a_{2}=s_{1} b_{2}  \tag{2-13}\\
& a_{s}=\frac{v_{s} \sqrt{L_{o}}}{z_{s}+z_{o}} \\
& s_{s}=\frac{z_{s}-z_{0}}{z_{s}+z_{0}} \\
& s_{1}=\frac{z_{1}-z_{0}}{z_{1}+z_{0}} \tag{2-14}
\end{align*}
$$

and
where
measurement procedures are discussed in Chapter IV.
Now the 2 -port transducer power gain which is one of the most important design parameters is to be considered. For arbitrary source and load impedances the transducer power gain is deified as follows:

$$
\begin{align*}
G_{T} & =\frac{\text { Power delivered to the load }}{\text { Power available from the source }} \\
& =\frac{P_{L}}{P_{a v S}} \tag{2-15}
\end{align*}
$$

From the definition of $\left|a_{1}\right|^{2}$ and $\left|b_{1}\right|^{2}$, and from Eq. (2-13),

$$
\begin{equation*}
P_{L}=\left|b_{2}\right|^{2}-\left|a_{2}\right|^{2}=\left|b_{2}\right|^{2}\left(1-\left|S_{1}\right|^{2}\right) \tag{2-16}
\end{equation*}
$$

Power available from the source $P_{a v S}$ can be expressed compactly using $a_{s}$. Since $a_{s}$ can be defined as the power wave applied to input port from an arbitrary source with voltage and internal impedance $v_{s}$, and $Z_{s}$, respectively. Then $\left|a_{s}\right|^{2}$ is the power available to the $Z_{o}$ line and has a definite relationship with $P_{\text {avs }}$ which is defined as maximum power available at a matched load for a given source. That is,

$$
\begin{equation*}
P_{\mathrm{avS}}=\frac{\left|\mathrm{a}_{\mathrm{s}}\right|^{2}}{1-\left|\mathrm{s}_{\mathrm{s}}\right|^{2}} \tag{2-17}
\end{equation*}
$$

The factor ( $1-\left|S_{s}\right|^{2}$ ) is the power transmission coefficient at the discontinuity between $Z_{s}$ and $Z_{0}$.

Transducer power gain is

$$
\begin{equation*}
G_{T}=\frac{\left|b_{2}\right|^{2}}{\left|a_{s}\right|^{2}}\left(1-\left|s_{s}\right|^{2}\right)\left(1-\left|s_{1}\right|^{2}\right) \tag{2-18}
\end{equation*}
$$

Applying Mason's rule to Fig. 2-6,

$$
\frac{b_{2}}{a_{s}}=\frac{s_{21}}{\left(1-s_{11} s_{s}\right)\left(1-s_{22} s_{1}\right)-s_{21} S_{12} s_{s} s_{1}} \quad(2-19)
$$

Combining Eqs. (2-18) and (2-19),

$$
G_{T}=\frac{\left|s_{21}\right|^{2}\left(1-\left|s_{s}\right|^{2}\right)\left(1-\left|s_{1}\right|^{2}\right)}{\left|\left(1-s_{11} s_{s}\right)\left(1-s_{22} s_{1}\right)-s_{21} S_{12} s_{s} s_{1}\right|^{2}}(2-20)
$$

Eq. (2-20) is an expression of the transducer power gain for arbitrary source and load impedance in terms of 2-port S-parameters.

The input and output reflection coefficients for arbtray source and load impedances are also important for imperdance matching. In general the input or output impedance varies with the load or source impedance. Also, the input and output reflection coefficients deviate from $S_{11}$ and $S_{22}$ for arbitrary terminations other than $Z_{0}$.

First, for an arbitrary load reflection coefficient $S_{1}$, the input reflection coefficient $S_{11}$. can be derived using the flow graph in Fig. 2-5 (b) and the node absorption formulas

$$
\begin{equation*}
\mathrm{S}_{11} \cdot=\mathrm{S}_{11}+\frac{\mathrm{S}_{12} \mathrm{~S}_{21} \mathrm{~S}_{1}}{1-\mathrm{S}_{22} \mathrm{~S}_{1}} \tag{2-21}
\end{equation*}
$$

In the same manner, the output reflection coefficient is derived for the case of an arbitrary source reflection coefficient $S_{s}$. The resulting output reflection coefficient is denoted as $\mathrm{S}_{22}{ }^{\circ}$.

$$
\begin{equation*}
s_{22}{ }^{\prime}=s_{22}+\frac{s_{12} S_{21} s_{s}}{1-s_{11} s_{s}} \tag{2-22}
\end{equation*}
$$

In circuit theory, maximum power transfer is obtained by matching the complex impedance source with the complex conjugate of the load impedance. In using the $S$-parameter method, this matching is also represented as complex conjugate matching. For instance, for an output reflection coefficient $S_{22}{ }^{\prime}$, the load reflection coefficient of $S_{1}=S_{22}{ }^{\circ *}$ will match with $\mathrm{S}_{22}$ ' resulting in maximum power transfer. Simultaneous matching conditions for $S_{s}$ and $S_{l}$ can be obtained setting $S_{s}$ and $S_{1}$ equal to $S_{11}{ }^{* *}$ and $S_{22}{ }^{* *}$, respectively, in Eqs. (2-21) and (2-22).

The stability of a 2-port network can also be considered by using Eqs. (2-21) and (2-22). The magnitude of the reflection coefficient is always less than unity when a transmission line is terminated witha passive element. For stability criterion the critical condition exists when
$S_{11}{ }^{\prime}$ and $S_{22}$ equal unity. If the magnitude of either reflection coefficient is larger than unity, the network may oscillate. This oscillation is possible even for the case of a passive source impedance because there is sufficient reflected power to maintain oscillation.

If $S_{11}{ }^{\circ}$ and $S_{22^{\prime}}$ are equated to unity in Eqs. (2-21) and $(2-22)$, then $S_{1}$ and $S_{s}$ can be found in terms of 2-port S-parameters. The loci of $S_{l}$ and $S_{s}$ obtained in this manner trace out circles on the reflection coefficient plane, the Smith Chart. These circles, which represent the boundary between the stable and unstable regions of $S_{s}$ and $S_{1}$, are defined as stability circles.
C. Unilateral Case

When the reverse gain of a 2-port network is zero, ide. $S_{12}=0$, power flows only in the forward direction independent of the impedance associated with the load and source. This is the unilateral case of a 2-port network. In an actual device, however, reverse gain always exists but is usually much smaller in magnitude than the forward gain, and can be ignored.

If the reverse gain, $S_{12}$, is assumed as zero, the total number of 2 -port $S$-parameters is reduced from 4 to 3. For this condition, the path connecting $a_{2}$ and $b_{1}$ on the flow graph of Fig. 2-6 vanishes. Eqs. (2-21) and (2-22), for the unilateral case $S_{12}=0$, reduce to

$$
\begin{equation*}
s_{11}=s_{11} \tag{2-21}
\end{equation*}
$$

and

$$
\begin{equation*}
s_{12}=s_{22} \tag{2-22}
\end{equation*}
$$

It should be noted that these parameters are independent of $\mathrm{Z}_{\mathrm{s}}$ and $\mathrm{Z}_{1}$. Thus a simultaneous conjugate match is always possible with passive $Z_{1}$ and $Z_{s}$, if the reflection coefficients $S_{11}$ and $S_{22}$ do not have a magnitude greater than unity. For the unilateral case maximum transducer power gain can be attained when the source and load reflection coefficients are
and

$$
\begin{equation*}
S_{s}=S_{11} \tag{2-23}
\end{equation*}
$$

$$
S_{1}=s_{22}^{*}
$$

Once $S_{S}$ and $S_{l}$ are found from Eq. (2-23), the corresponding source and load impedances can be read directly from the Smith Chart or can be calculated using Eq. (2-14). These impedances are the complex conjugates of input and output impedances of the 2 -port network.

Unilateral transducer gain can be derived from Eq. (2-20) by letting $S_{12}=0$.

$$
\begin{equation*}
G_{T U}=\frac{\left|s_{21}\right|^{2}\left(1-\left|s_{s}\right|^{2}\right)\left(1-\left|s_{1}\right|^{2}\right)}{\left|\left(1-s_{11} S_{s}\right)\left(1-s_{22_{1}}\right)\right|^{2}} \tag{2-20}
\end{equation*}
$$

When the input and output impedances are conjugately matched to the source and load impedances, maximum unilateral gain $G_{U}$ is obtained. From Eq. (2-20).

$$
\begin{equation*}
G_{U}=\frac{\left|S_{21}\right|^{2}}{\left(1-\left|S_{11}\right|^{2}\right)\left(1-\left|s_{22}\right|^{2}\right)} \tag{2-24}
\end{equation*}
$$

It should be noted that $G_{U}$ is invariant regardless of source and load impedances.

Let us define the unilateral figure of merit $U$ as

$$
\begin{equation*}
U=\frac{\left|s_{11} S_{22} S_{12} S_{21}\right|}{\left(1-\left|s_{11}\right|^{2}\right)\left(1-\left|s_{22}\right|^{2}\right)} \tag{2-25}
\end{equation*}
$$

It can be shown from Eqs. (2-20), (2-24), and (2-25) that

$$
\begin{equation*}
\frac{1}{(1+U)^{2}}<\frac{G_{T}}{G_{T U}}<\frac{1}{(1-U)^{2}} \tag{2-26}
\end{equation*}
$$

In order to justify the unilateral design, the figure of merit may be evaluated to check the error in power gain calculation (error from using Eqs. (2-20)' and (2-24), rather than the exact Eq. (2-20)). Depending upon the purpose of a circuit, maximum permissible $U$ can be determined from the design accuracy required.

It should be noted that the magnitudes of $S_{11}$ and $S_{22}$ are assumed to be smaller than unity for all the preceding equations.
D. Three-Port Formulation ${ }^{17}$

In this section, the relationships between 3-port and

2-port parameters will be discussed. The most common active device network is a 2-port network. However, many active devices have more than 3 terminals. In order to represent a device that has more than 3 terminals, 2-port parameters are not sufficient for characterizing the device for all possible arbitrary configurations.

Even in 3-terminal devices such as transistors, a 3-port representation is sometimes convenient even though there is redundancy. A 3-port model of a 3-terminal device is exactly the case in which the indefinite admittance matrix representation is more general and versatile than the general definite admittance matrix representation. This is especially true for the cases of port interconnections and arbitrary terminations at the ports. The application to 3-terminal transistors is fully appreciated in Bodway's paper.

A 3-port signal flow graph is shown in Fig. 2-7 (a). Port 1 and port 2 are the input and output ports, respectively, and port 3 is terminated with an arbitrary impedance whose equivalent reflection coefficient is $S_{3}$. A 2-port model can be obtained using the node absorption formula to eliminate the nodes of port 3. The resulting 2 -port $S$-parameters obtained in this manner are as follows:

$$
s_{11(3)}=s_{11}+\frac{S_{31} S_{13} S_{3}}{1-S_{33} s_{3}}
$$



Fig. 2-7. (a) 3-port signal flow graph;
(b) Derived 2-port absorbing port 3 nodes. $\left(a_{s 2}=a_{s 3}=0\right)$

$$
\begin{align*}
& S_{21(3)}=S_{21}+\frac{S_{31} S_{23} S_{3}}{1-S_{33} S_{3}} \\
& S_{22(3)}=S_{22}+\frac{S_{32} S_{23} S_{3}}{1-S_{33} S_{3}} \\
& S_{12(3)}=S_{12}+\frac{S_{13} S_{32} S_{3}}{1-S_{33} S_{3}} \tag{2-27}
\end{align*}
$$

Subscript (3) designates the absorbed port. The resulting flow graph is also shown in Fig. 2-7 (a). Note that the choice of the port to be eliminated is arbitrary.

For a 3-terminal device such as illustrated in Fig. 2-8, the scattering variables are not independent from Kirchhoff's Lav: and the resulting 3-port scattering arameters have a characteristic that is different from that of general 3-port network. Theoretically 4 out of the 9 components are sufficient in order to characterize the ( $3 \times 3$ ) 3-port scattering matrix for a 3-terminal device.

Refer to Fig. 2-8, from Kirchhoff's current law at the reference or ground node

$$
i_{1}+i_{2}+i_{3}=0
$$

Also, no current will flow if the potentials applied to every terminal pair are made equal, ie.

$$
i_{1}=i_{2}=i_{3}=0 \text {, if } v_{1}=v_{2}=v_{3}
$$

The above two equations can be expressed in scattering variables

$$
a_{1}+a_{2}+a_{3}=b_{1}+b_{2}+b_{3}
$$

and

$$
\begin{equation*}
a_{1}-b_{1}=a_{2}-b_{2}=a_{3}-b_{3}=0 \text {, if } a_{1}=a_{2}=a_{3} \tag{2-28}
\end{equation*}
$$

The conditions as set forth in the preceding equations are summarized as follows:


Eq. (2-29) dictates the dependency of the individual S-parameters. It can be seen that 4 independent S-parameters determine the remaining $S$-parameters according to Eq. (2-29). Although in practice all of the 9 parameters are measured for precision, the number of measurements can be reduced to 4 and the rest can be obtained from Eq. (2-29).

If the upper limits of Eq. (2-29) are extended from 3 to $n$ the equation will then be good for all n-terminal n-port networks.

Before concluding this chapter it is worthwhile to discuss another possible 3-port representation for a 3-terminal


Fig. 2-8. 3-port representation of 3-terminal device with common reference.


Fig. 2-9. 3-port representation of 3-tèrminal device without common reference.
device. Fig. 2-9 shows the connection. This connection is very common in actual devices; however, it is not so practical when making parameter measurements. Also, the inherent feedback of the circuit may cause instability when measurements are made.

This circuit can be thought of as the dual of the circuit discussed in Fig. 2-8, but there are no obvious relationships between the two sets of parameters. There is no impedance matrix defined for the connection of Fig. 2-8 and for the network of Fig. 2-9 no admittance matrix is defined.

## CHAPTER III

## APPLICATION TO HIGH-FREQUENCY

## SEMICONDUCTOR DEVICES

In the preceding chapter, the concept of scattering relations was introduced along with definitions of 2-port and 3-port parameters. In this chapter high-frequency semiconductor devices are discussed from the viewpoint of relating the $S$-parameters to existing equivalent circuits. A. General Review of High-Frequency Semiconductor Devices 1. Bipolar Transistors. At high frequencies the performance of a bipolar transistor is limited by diffusion capacitance, junction transition capacitance, and basespreading resistance. The hybrid $P_{i}$ equivalent ${ }^{18,19}$ circuit of Fig. 3-1 is one equivalent circuit that models the approximate high-frequency performance of a bipolar transistor. However, the equivalent circuit is not valid at frequencies where the transition time of minority carriers across the base region cannot be neglected. For this case, some of the parameters in Fig. 3-1 are frequency dependent. At a single frequency, the parameters of a transistor vary with the operating point. In general, diffusion capacitance is proportional to the forward junction current and junction transition capacitance is inversely proportional to the absolute value of square or cube root of the junction


Fig. 3-1. Junction transistor equivalent circuit for common-emitter configuration.
potential difference. Whether the square or cube root applies depends on the impurity profile, i.e., for a homogenously-doped $p-n$ junction, the power is $1 / 2$, and for a linearly-graded junction, the power is $1 / 3 .{ }^{18}$

Since, in a forward-biased junction, diffusion capacitance is predominant, $C_{b}{ }^{\prime}$ is mainly diffusion capacitance. At a reverse-biased junction, such as between the collector and base, the predominant capacitance is transition capacitance. Therefore, $C_{b}{ }_{c}$ is mainly transition capacitance. In the saturation region, i.e., under a very low collector voltage condition, the collector-base junction becomes forward-biased; in this case, diffusion capacitance and a low dynamic resistance would replace $C_{b}{ }^{\prime} c$ and $r_{b}{ }^{\prime} c$ of the active region model.

At frequencies within the VHF band, S-parameters of a device may exhibit the characteristics of the equivalent circuit in Fig. 3-1, since the time delay due to minority carrier transit time in the frequency range is not significant for a graded-bade junction transistor. For a typical silicon npn transistor, $2 N 3478$, it is reported that transit time is about $112 \mathrm{ps} .{ }^{8}$ Using this transit time, a calculation for phase shift at 200 NWHz yields about 10 degrees. Ihis will appear directly in $S_{21}$, but not explicitly in other parameters since only $\mathrm{S}_{21}$ is directly related with the minority carrier transit time.

At high frequencies, the input impedance of a bipolar device consists of $r_{b b}, r_{b} e^{\prime}$ and $C_{b} e^{\prime}$, if we neglect the feedback elements, $r_{b}{ }^{\prime}$ and $C_{b}{ }^{\prime} \cdot S_{11}$ is then determined by this impedance, and will follow a constant resistant circle on the Smith Chart, if frequencies are high enough to reduce the capacitive impedance of $C_{b} e^{\text {much }}$ below that of $r_{b} e^{\cdot}$ A plot of $S_{22}$ follows a constant conductance circle and the conductance is the parallel conductance of $r_{b}{ }_{c}$ and $\mathbf{r}_{\text {ce }}$. For this plot it is assumed that impedance between $b^{\prime}$ and $e$ is negligible.
$S_{21}$ and $S_{12}$ will be discussed in Chapter $V$, along with discussions of parameter variations due to bias voltages. 2. MOSFET. ${ }^{20}$ One of the main limitations of highfrequency bipolar devices depends upon the relatively low minority carrier velocity in base region. From this point of view, the radically different structure of the field effect transistor, in which current movement depends on majority carriers controlled by electric field, can be expected to show better high-frequency performance. However, this was not true for early junction FETs in which the large ineffective areas of the gate provided large internal feedback capacitance and large channel capacitance, resulting in very low cut-off frequency. Furthermore, highly-developed bipolar transistor technologies began to produce GHz range devices of good performance, and in the light of this fact,

(a)


Fig. 3-2. (a) Structure of depletion mode MOSFET; (b) Equivalent circuit.
it is natural that the early Fs'rs failed to receive attention as high-frequency devices.

Recent improvements in NOSFE'T technology present a good prospect for these devices in high-frequency applications. The reverse feedback capacitance is reduced by geometrical gate offset, and dual-gate units give effective electro-static shielding between drain and source. These devices also have many properties that cannot be found in bipolar devices: excellent cross modulation characteristics, small AGC power consumption due to high DC input resistance, high power amplification, and good thermal stability. The noise figure is also comparable to that of good bipolar devices.

These characteristics are favorable for many applications, such as $R F$ amplifiers, gain-controlled oscillators, etc. Thus, the study of MOSFE'Ps with S-parameters is particularly interesting for their future applications at higher frequencies.

Fig. 3-2 shows the physical structure and a complete equivalent circuit of an $n$-channel single-gate MOSFET. The diodes and substrate terminals are features resulting from the construction of the devices. The diodes are $p-n$ junctions formed between the heavily-doped source and drain, and the lightly-doped substrate. To some degree, the substrate may also form a p-n junction with the channel, and can
be used as a second gate, but due to difficulty in controlling the transconductance the terminal is usually grounded for operation in the common-source connection.

The principal difference between the equivalent circuit of bipolar transistors and MOSFETs is the addition of series combination, $r_{c}$ and $C_{c}$, that connect the gate and source. This combination represents the distributed nature of the channel and the gate. $g_{m}$ is controlled by the voltage across $C_{c}$ so the reduction in gain at high frequencies is related to the time constant $r_{c} C_{c}$.
$r_{g s}$ and $r_{g d}$ represent the leakage paths associated with the oxide layer of the gate and source, and the gate and drain, respectively. $r_{d s}$ is given by the slope of the common source output characteristics and varies widely with $\mathrm{v}_{\mathrm{GS}}$.

The capacitances $C_{g d}, C_{g s}$, and $C_{d s}$ include intrinsic and extrinsic capacitances between corresponding electrodes. The intrinsic portion of $C_{g d}$ decreases with increasing drain voltage due to the effective widening of separation between drain and gate by carrier depletion region.

One more consideration has to be given to the equivalent circuit in Fig. 3-2 (b). For an offset gate MOSFET, a resistance appears in series with the drain terminal because of the bulk resistance of the unmodulated region. The output impedance at high frequencies then exhibit a second
order effect.
At frequencies above 400 MHz , deviations from the simple low frequency theory are reported by Kolk and Johnson. ${ }^{20}$ According to their report, significant differences were observed in input and forward transfer admittance above 400 MHz , and at low frequencies in the case of reverse transfer admittance. They also mention that some of the deviations may be accounted for by small but significant inductances in the transistor assembly and measurement equipment; however, the experimental data indicate that further refinement of the equivalent circuit is necessary at these frequencies. Further discussion of this will be given in Chapter $V$; the remainder of this section will be devoted to dual-gate MOS FETs.

As was discussed early in this section, one of the effective ways to reduce feedback capacitance is by adding another gate between drain and gate of a single-gate MOSFET. Actually this becomes a cascode amplifier if gate 2 is RF grounded, by creating a common-source and common-gate pair. Fig. 3-3 shows a schematic diagram for the cascode connection. The input and output impedance values in this connection are about eciual to those of common-source and commongate connections, respectively. This amplifier gain is also about the same as that of single-common-source stage. But reverse gain or feedback becomes much smaller than that for a single common-source or common-gate stage.


Fig. 3-3. $\begin{aligned} & \text { Equivalent circuit of dual-gate } \\ & \text { MOSFET. }\end{aligned}$


Fig. 3-4. Cascaded two-port for dual-gate
MOSFET study.

To understand the behavior of a dual-gate MOSFEY using S-parameters, it is convenient to derive a 2-port expression for a cascaded amplifier. Referring to Fig. 3-4, it is seen that S -parameters for a cascaded amplifier are obtained from 2 sets of single-stage S-parameters. Then, substituting common-source and common-gate $S$-parameters for a cascode amplifier, which can represent a dual-gate MOSFeT, S-parameters of a dual-gate MOSFET will be obtained. Numerical justifications will be given in Chapter $V$.
3. Differential Amplifier-Integrated Circuit. A configuration that can provide good thermal stability without sacrificing low frequency gain is the differential amplifier. As in Fig. 3-5, identical transistors are made on an IC chip, with $Q_{1}, Q_{2}$, and $Q_{3}$ interconnected to give minimum wiring distance. However, this type of IC is not confined to differential amplifier applications; by utilizing the proper access terminals they are usable as cascode amplifiers, commoncollector and common-emitter pair amplifiers, balanced mixers, oscillators, etc.

The S-parameters of a differential amplifier with balanced input and output are exactly the same as those of a single transistor if the normalization impedance is doubled. For the circuit in Fig. 3-6, the impedance due to $Q_{3}$ does not appear in the amplifier action so that $Q_{1}$ and $Q_{2}$ are in a series connection, and their impedance parame-


Fig. 3-5. IC differential amplifier


Fig. 3-6. Balanced input and output for ideal differential operation.
ter add. For example, the S-parameters for a particular transistor are measured with respect to 50 ohms then the S-parameters of a differential amplifier circuit consisting of 2 identical transistors are the addition of those of the single transistor and the reference impedance is now 100 ohms.

When $Q_{3}$ is used as the current source as in Fig. 3-5, this type of differential amplifier circuit can be a 5-terminail network and may be characterized by 4 -port parameters. Assuming symmetry of the parameters due to identical characteristics of $Q_{1}$ and $Q_{2}$, the number of parameters required to characterize this IC circuit is 8 . If it is assumed that the output impedance of $Q_{3}$ is very large, the number of independent parameters is reduced to 4.

By using the flow graphs of Fig. 3-5, the 4-port relations can be written as follows:

$$
\left(\begin{array}{l}
b_{1}  \tag{3-1}\\
b_{1} \\
b_{2} \\
b_{2 \cdot}
\end{array}\right)=\left(\begin{array}{llll}
s_{11} & s_{11} \cdot & s_{12} & s_{12} \cdot \\
s_{1 \cdot 1} & s_{1 \cdot 1}, & s_{1 \cdot 2} & s_{1 \cdot 2} \\
s_{21} & s_{21 \cdot} & s_{22} & s_{22 \cdot} \\
s_{2 \cdot 1} & s_{2 \cdot 1 \cdot} & s_{2 \cdot 2} & s_{2 \cdot 2 \cdot}
\end{array}\right) \quad\left(\begin{array}{l}
a_{1} \\
a_{1} \\
a_{2} \\
a_{2}
\end{array}\right]
$$

If $Q_{1}$ and $Q_{2}$ are identical then

$$
S_{i j}=S_{i^{\prime} j^{\prime}}
$$

and

$$
\begin{equation*}
S_{i j} \prime=S_{i^{\prime} j} \tag{3-2}
\end{equation*}
$$

Also, the sum of each column or row is unity for very high output impedance of $Q_{3}$. This was discussed in Chapter II where an n-port network with a common reference was consicered.

In a practical case it is convenient to measure 8 parameters for a complete set from Eqs. (3-1) and (3-2). Four parameters can be measured for a good approximation in the case where the output impedance of $Q_{3}$ is very high.

Once the scattering matrix of Eq. (3-1) is known then S-parameters for every possible configuration can be derived. For instance, S-parameters for a differential mode can be obtained where new scattering variables are defined.

From Fig. 3-5 and the relationships of impedance and scattering variables in Chapter II, differential and common mode variables are represented as

$$
\begin{align*}
& v_{1} \pm v_{1} \cdot=\sqrt{L_{0}}\left[\left(a_{1} \pm a_{1}\right)+\left(b_{1} \pm b_{1},\right)\right] \\
& i_{1} \pm i_{1} \cdot=\frac{1}{\sqrt{L_{0}}}\left[\left(a_{1} \pm a_{1},\right)-\left(b_{1} \pm b_{1},\right)\right] \tag{3-3}
\end{align*}
$$

For port 2 and port $2^{\prime}$ the above equations can be extended if the subscripts 1 and $1^{\circ}$ are changed to 2 and $2^{\prime}$, respecttively. A simple linear transformation of Eq. (3-1) finally
yields

$$
\begin{equation*}
\overline{\mathrm{B}}_{\mathrm{dc}}=\overline{\mathrm{C}} \overline{\mathrm{~S}} \overline{\mathrm{C}}^{-1} \overline{\mathrm{~A}}_{\mathrm{dc}} \tag{3-4}
\end{equation*}
$$

where

$$
\bar{A}_{d c}=\left(\begin{array}{l}
a_{1}-a_{1} \cdot \\
a_{1}+a_{1} \cdot \\
a_{2}-a_{2} \cdot \\
a_{2}+a_{2} \cdot
\end{array}\right) \quad \bar{B}_{d c}=\left(\begin{array}{l}
b_{1}-b_{1} \\
b_{1}+b_{1} \\
b_{2}-b_{2} \\
b_{2}+b_{2}
\end{array}\right) \quad \bar{c}=\left(\begin{array}{cccc}
1 & -1 & 0 & 0 \\
1 & 1 & 0 & 0 \\
0 & 0 & 1 & -1 \\
0 & 0 & 1 & 1
\end{array}\right)
$$

and $\overline{\mathrm{S}}$ is a scattering matrix defined by Eq. (3-1).
If $Q_{1}$ and $Q_{2}$ are assumed to be identical then no common mode variables will appear for the differential mode of operation with balanced input and output. In other words, for identical $Q_{1}$ and $Q_{2}, a_{1}+a_{1}, a_{2}+a_{2}, b_{1}+b_{1}$, , and $b_{2}+b_{2}$. terms will vanish substituting the conditions of $v_{1}=-v_{1}$., $v_{2}=-v_{2}, i_{1}=-i_{1}$, , and $i_{2}=-i_{2}$, into Eq. (3-3). Then 2 sets of differential mode terms are left in $\bar{A}_{d c}$ and $\bar{B}_{d c}$, resulting in a 2 -port relation. In this case the scattering matrix $\overline{\mathrm{C}} \overline{\mathrm{S}}^{\mathrm{C}}{ }^{-1}$ has 4 significant components relating the differential mode terms of $\bar{A}_{d c}$ and $\bar{B}$ d ce Referring to Eq. (3-3), it is seen that 4 parameters represent the differendial mode $S$-parameters that are normalized with $2 Z_{0}$. (Refer to Eq. $(2-7)$ and note that port current represented by Eq. (3-3) is 2 times the actual port current.) The resulting 4 differential mode $S$-parameters obtained from the
scattering matrix of Eq. (3-4) are given below.

$$
\begin{align*}
& S_{11 d}=\frac{1}{2}\left(S_{11}+S_{1 \cdot 1},-S_{11},-S_{1 \prime 1}\right) \\
& S_{22 d}=\frac{1}{2}\left(S_{22}+S_{2 \prime 2},-S_{22 \prime}-S_{2 \prime 2}\right) \\
& S_{21 d}=\frac{1}{2}\left(S_{21}+S_{2 \prime 1},-S_{21},-S_{2 \prime 1}\right) \\
& S_{12 d}=\frac{1}{2}\left(S_{12}+S_{1 \cdot 2},-S_{12},-S_{1 \prime 2}\right) \tag{3-5}
\end{align*}
$$

The normalization impedance of the differential mode $S$-parameters in Eq. (3-5) is $2 Z_{0}$ while 4 -port parameters are measured with respect to a normalization impedance of $Z_{0}$. Then Eq. (3-5) is combined with Eq. (3-2) to give the S-parameters of the differential amplifier with identical $Q_{1}$ and $Q_{2}$ of Fig. 3-5.

Eq. (3-5) provides the means whereby the S-parameters of the constituent transistors in a differential amplifier can be measured. For instance, the common-emitter input reflection coefficient of $Q_{1}$ and $Q_{2}$ in Fig. 3-5 is $S_{11}-S_{11}$, and can be obtained by 4-port measurement.
B. S-Parameters as Indirect Means of Measurement for Other Parameters
At frequencies above 100 NHz , direct measurement ${ }^{22}$ of $\mathrm{h}-\mathrm{y}, \mathrm{y}$, or z -parameters, ordinarily used in active circuit design at low frequencies, is very inconvenient because of the difficulty in establishing the required short or open
circuit condition. Also, a short or open circuit frequently causes the device to oscillate.

Short circuit admittance or y-parameters are the most popular high-frequency parameters. It is possible to measure them up to microwave frequencies with the slotted line technique. But there are two major inconveniences when compared with the $S$-parameter measurement technique. First, for wide-band measurement, the slotted line should be tuned every time the frequency is varied. This tuning requires too much time and labor. Second, as discussed before, the short or open condition of a device easily can result in an oscillatory state. This causes unstable measurement.

Because of these difficulties, it is expedient to derive y-parameters or any other immittance parameters from the more precise and easily obtainable S-parameters which are measured with a more stable wide-band measurement system. Conversion tables are provided in Appendix I.

Other equivalent circuit parameters, such as hybrid Pi, normal Pi, ${ }^{19}$ or others, can also be derived from the measured $S$-parameters as long as the equivalent circuits are valid at the frequency of interest. However, it should be noted that there is no reason to insist on equivalent circuits at much higher frequencies where the lumped element representation fails.

At relatively low frequencies, i.e. 200 MHz , the equivalent circuits are still useful for understanding a transistor. As was shown earlier in this chapter, the equivalent circuit of the MOSFET can be thought of as a derived form of the normal Pi circuit of Fig. 3-7. The normal Pi elements can be derived directly from the $y$-parameters without introducing conversion errors. This provides a convenient way to approach the equivalent circuit from measured parameters. Also, referring to Fig. 3-7 and the scattering admittance conversion table in Appendix $I$, it is possible to use approximate formulas for VHF band transistors. The approximate equations are given for the elements of the normal Pi circuit.

$$
\begin{aligned}
A & =y_{1}+\frac{S_{12}\left(S_{21}-2\right)}{M} \\
B & =\frac{2 S_{12}}{M} \\
C & =\frac{-2 S_{21}}{M} \\
\text { where } \quad y_{1} & =\frac{1-S_{11}}{1+S_{11}}
\end{aligned}
$$



$$
\begin{aligned}
& A=y_{11}+y_{12} \\
& B=-y_{12} \\
& C=y_{21}-y_{12} \\
& D=y_{22}+y_{12}
\end{aligned}
$$

Fig. 3-7. Normal Pi model.

$$
y_{2}=\frac{1-S_{22}}{1+S_{22}}
$$

and

$$
M=\left(1+S_{11}\right)\left(1+S_{22}\right)
$$

A, B, C, and D are admittances normalized with $1 / Z_{0}$ Eqs. (3-6) are derived with the assumption that the magnitude of $M$ is much larger than that of $S_{12} S_{21}$ which is always true in VHF transistors for normal operating conditions. $y_{1}$ and $y_{2}$ are normalized input and output admittances that can be obtained directly from a Smith Chart plot of $S_{11}$ and $S_{22}$. From the frequency response of $A, B, C$, and $D$, the MOSFET equivalent circuit of Fig. 3-2 can be approximated by trial and error.

The hybrid Pi circuit of Fig. 3-1 can also be derived mathematically for a bipolar transistor through a double conversion, i.e., first to y-parameters and then to hybrid Pi parameters. However, since the equivalent circuit is not an exact representation of the device performance, the reliability is inferior to the direct measurement method. An approximate method is more useful and is given in Appendix II.
C. Mixer S-Parameters

In Chapter II, the possibility of S-parameters as amplifier design parameters was fully described in the discussion of 2-port formulation. The practical design
procedures appear in Weinert, ${ }^{7}$ Anderson, ${ }^{8}$ Froehner, ${ }^{9}$ and Bodway. ${ }^{10}$ However, no attempt has been made in the literature to characterize a mixer circuit with S-parameters. It will be shown that $S$-parameters can also be used in designing a mixer circuit.

Conventionally, short circuit admittance parameters have been employed in mixer circuit design. However, at frequencies where y-parameters cannot be easily obtained it is desirable to define mixer parameters in terms of S -parameters. Once the $S$-parameters are known, designing a mixer is similar to amplifier design.

It can be noted in Fig. 3-8 that there are three major frequencies associated with a mixer stage. Only 2 of these are pertinents the RF input signal and IF output signal. When an ideal IF filter is provided all frequencies are attenuated except the IF signal. For the balanced modulator circuit as in Fig. 3-9, elimination of local oscillator signal from an $R F$ input end can easily be achieved. Then a linear relationship between the $R F$ and IF signal amplitudes can be measured as in the normal 2-port S-parameter case.

From Fig. 3-8, 3 important $S$-parameters for z mixer design can be defined as follows:

$$
\left.S_{11}(R F)=\frac{b_{1}(R F)}{a_{1}(R F)} \right\rvert\, \quad a_{2}(R F)=0
$$



Fig. 3-8. Mixer model.


Fig. 3-9. An example of a balanced mixer.

$$
\begin{align*}
& S_{22}(I F)=\left.\frac{b_{2}(I F)}{a_{2}(I F)}\right|_{a_{1}(I F)=0} \\
& S_{21}(M I X)=\left.\frac{b_{2}(I F)}{a_{1}(R F)}\right|_{a_{2}(I F), a_{2}(R F)=0} \tag{3-7}
\end{align*}
$$

These parameters are a function of the local oscillator input level as well as the DC operating point. If the IF input impedance is $Z_{0}$, ie. 50 ohms, $S_{11}(R F)$ is the input reflecttion coefficient for the RF signal. $S_{22}(I F)$ is the output reflection coefficient for the IF signal. $S_{22}(I F)$ can be measured approximately after removing the $R F$ and local oscillator signals and terminating the RF input terminals with $Z_{0} \cdot S_{21}(M I X)$ denotes conversion gain. Once the prameters are obtained, unilateral (mixer) transducer gain can be calculated and proper matching networks designed to give maximum gain.

## CHAPTER IV

## MEASUREMENT OF S-PARAMETTERS

## AT VHF FREQUENCIES

In Chapter II it was suggested that $S$-parameters can be obtained from the ratio of b's to a's according to Eq. (2-9). This is for the case where a signal source of internal impedance $Z_{0}$ is connected to one port and the remaining ports are terminated in $Z_{0}$. Under this condition the a's are all zero except at the port where the signal source is connected. Actual measurement procedures for obtaining S-parameters will be treated in this chapter. A. Measurement Set-Up for 2-Port S-Parameters

A block diagram for making 2 -port $S$-parameter measurements is shown in Fig. 4-1. Referring to the block diagram it can be seen that the over-all accuracy of a given measurement is determined by many factors such as VSWR (due to bias networks, directional couplers, transistor jig or any discontinuity in the signal path), loss in defective connectors, errors in vector voltmeter (VVM), etc. The VVM is the most important component of the measurement system. In order to get reliable results it is necessary to use a high-quality instrument.

The specifications for the equipment and the necessary accessories corresponding to the blocks in Fig. 4-1 are as


Fig. 4-1. Block diagram for 2-port

1. Signal Generator: GR 1201-A

Variable frequency signal generator with a 50 ohm coaxial output. The signal output is attenuable continuously from several volts down to zero volts. Frequency ranges of $40-250 \mathrm{MHz}$ and $250-920 \mathrm{MHz}$ are available in this type.
2. Bias Tee: MICRO LAB/FXR HW-O2N

To mix or separate RF signals and DC bias currents, bias tees must be used. The monitor tees, MICRO LAB/FXR HW-02N, are employed for this purpose. Characteristics for HW-O2N are as follows:
a. Maximum VSWR $=1.3$
b. Impedance $=50$ ohms
c. Maximum insertion loss $=0.2 \mathrm{~dB}$

## 3. Dual Directional Coupler: HP 774D

For sampling the incident, reflected, or transmitted wave, dual directional couplers are used. HP 774D Couplers, which are of strip line structure, are especially good for VHF applications. The close tracking of the auxiliary arms makes these couplers useful for wide-band measurement. S-parameters are determined by the ratios of the outputs of two arms so coupling variations with frequency are eliminated. The characteristics of the HP 774D by the
manufacturer are as follows:
a. Frequency range $=215-450 \mathrm{MHz}$
b. Minimum directivity $=40 \mathrm{~dB}$
c. Coupling attenuation $=20 \mathrm{~dB}$
d. Maximum primary line $S W R=1.15$
4. Vector Voltmeter: HP 8405A

The HP 8405A Vector Voltmeter measures the magnitude and the phase difference between two voltage vectors from 1 to 1000 MHz . It uses a sampling method. The a's and b's are measured with the VVM. and the reflection coefficients are the ratios of $b_{i}$ to $a_{i}$.

Bridge tees and isolators, which do not appear in the block diagram of Fig. 4-1, are employed when readings are taken. Specifications of the HP 8405A VVM are as follows:
a. Frequency range $=1-1000 \mathrm{MHz}$
b. Voltage ratio accuracy $= \pm 2 \%$
c. Phase accuracy $=$ at single frequency $\pm 1.5^{\circ}$
d. Voltage range (rms) $=300 \mathrm{uV}-1.0 \mathrm{~V}$ for A channel and 100 uV - 1.0 V for B channel at frequency of $10-500 \mathrm{MHz}$.
5. Test Jig or Transistor Mounting

The test jig is a device which accomodates the transistor during test. There are two types

(a)

(b)

Fig. 4-2. Test jigs for (a) 2-port and (b) 3-port measurements.
of test jigs: strip-line and coaxial.
For this experiment, several test jigs were attempted so that 2-port, 3-port, and 4-port parameters could be measured. Fig. 4-2 shows the test jigs used. Coaxial jigs where several pieces of flexible dielectric line, RG A/J with socket pins built inside the lines that are an integral part of the inner conductors, have been used.
The lines are cut to identical lengths so the paths for the reflected and transmitted waves are equal. By doing this the ports are interchangeable thereby reducing set-up time. The path differences observed for each branch of the jigs can be minimized by trial and error. The residual path differences caused a phase shift of less than $3^{0}$ at frequencies between 150 and 250 MiHz .
6. Adjustable Line: GR 874-LA

Since the phase differences between incident and scattered center are to be measured, it is necessary for the waves to have equal electrical lengths. An adjustable coaxial line with an impedance of 50 ohms was used for this purpose.
7. 50 Ohm Coaxial Terminations: HP IModel 908 A

For VHF frequencies precise terminations over
a wide range of measurement frequencies are necessary so there are no undesired reflections at the termination points during measurement. At relatively low frequencies the terminations can be constructed with carbon film or solid resistors as shown in Fig. 4-3. Biasing terminals are also shown in the figure.
8. DC Supplies It is desirable to use very stable and finely adjustable DC supplies. A high resistance may be connected in series with the bias supply to give smoother control of bias current during bipolar device testing.

## B. Reference Plane

It has been shown that when a traveling wave encounters a discontinuity on a transmission line, twp scattered waves traveling in opposite directions are generated. S-parameters can be obtained by comparing the amplitudes and phase angles of the scattered waves with those of the original wave.

However, measurements cannot be made directly at the scattering center (port terminals). This is because of the physical size of the directional couplers and connecting leads (50-ohm line between measurement system and device to be tested). As a result, the scattered waves travel finite


Fig. 4-3. Termination with DC block.
distances along lossless lines before they are measured. To obtain the relationships between the scattered and original waves at the scattering center, phase compensation is made by adjusting the path length of the original (incident)
wave
Traveling distances for the waves are shown in Fig. 4-4 for a 2-port measurement. $L_{i}, L_{r}$, and $L_{t}$ are the electrical path lengths (excluding that of the device to be tested) of incident, reflected, and transmitted waves measured from an arbitrary input point. $L_{r}$ and $L_{t}$ are equal if the test jig is symmetrical and both directional couplers are identical. $L_{i}$ can be adjusted by changing the length of the adjustable line. When the length $L_{i}$ is made equal to $L_{r}$ and $L_{t}$, the amplitude and phase relationships for the incident and reflected waves at the measurement point and the original scattering point are identical. Further, the relationships are independent of frequency.

The scattering plane $P$ in Fig. $4-4$ is often called the "reference plane." 6 The reference plane can be located by employing short or open termination in place of the device to be tested. The phase difference between incident and reflected waves for short and open ends are $180^{\circ}$ and $0^{\circ}$, respectively; however, the fringing effect at an open end may not allow precise positioning of the reference plane. The short termination method is more commonly employed for


Fig. 4-4. Block diagram showing reference plane calibration.
that practical reason. A "through-section" ${ }^{12}$ can also be used to adjust the path lengths but is not necessary in a test set-up where $L_{r}$ and $L_{t}$ are equal.
C. Measurement Procedures

1. Preliminary check of equipment before mounting test jig (channel A probe connected to $V$ in Fig. 4-1).

After every connector is firmly locked, the amplitude difference between channel A and channel $B$ is read without test jig in position. The amplitude ratios of channel A and channel B outputs in this manner should be very close to unity with open or short termination in place of the test jig, and approach zero with a 50 -ohm termination. Typical observed values are 0.99 for open or short termination and 0.01 for 50 -ohm termination at 150 MHz . These amplitude errors can be attributed to the directional couplers and VVM.
2. Set-up of reference plane.

The test jig is mounted in position and the shorting conductor is plugged into the socket. The adjustable line should be changed until the phase meter of the VVM reads 180 degrees. The test jig is then reversed by connecting the input to the
opposite end. The phase differences observed in this manner are due to the difference between the input and output path lengths. I'his difference has to be minimized by careful construction of the test jig. The adjustable line should be readjusted so that the same reference plane is attained when the input and output ports are interchanged during calibration.
3. DC supplies consideration.

After the reference plane is set up, the transistor to be tested is plugged in. Then proper DC voltages are provided to the transistor through the bias tees.
4. Signal generator.

The signal generator is first set at the desired frequency with a minimum output level. Then the output is slowly increased until the reading of channel A voltage reaches full scale on the -50 dB or 1 mV range. Assuming that the coupling factor of the directional coupler, $H P$ 774D, is 25 dB maximum, the magnitude of the input signal is less than 20 mV .
5. S-parameters measurement (2-port parameters 0 .

Refer to Fig. 4-1. Initially channel B is connected to the $V$ arm of directional coupler 1 in
order to measure reflection coefficients $S_{11}$ and
$S_{22}$. Port number 1 or 2 must be defined in terms of the terminals of the device being tested. For example, in bipolar transistor measurements it may be convenient to define the base and the collector terminals as port 1 and port 2, respectively. If port 1 of the device being tested is connected to directional coupler $1, S_{11}$ can now be measured. If the channel A voltage is set at full scale, then the voltage of channel $B$ and the phase angle between channel A and channel B represents the magnitude and phase of the input reflection coefficient, $S_{11}$. Measurement of the output reflection coefficient, $S_{22}$, is exactly the same as that of $S_{11}$ except that the input and output ends of test jig are interchanged. The DC supply must also be changed.

The forward gain $S_{21}$ is obtained by plugging the channel $B$ probe into arm $W$ of directional coupler 2. The ratio of channel $A$ and channel $B$ voltage amplitudes and phase angles between them are the polar expression of $S_{21} \cdot S_{12}$, the reverse gain, is measured in the same way as $\mathrm{S}_{21}$, but input and output ports are interchanged.

If the base and the collector of a bipolar
device (or the gate and drain of a MOSFET) are properly assigned port 1 and port 2, respectively, then $S_{12}$ is usually very small for those devices in common-emitter (or common-source) configuration. These small values can be read accurately using the $\log$ step scale of VVM. Also, in order to read a small $\mathrm{S}_{12}$, injection of large signal is necessary. A minimum of 100 mV applied at the output port was required for satisfactory reading.
6. Three-port parameters and differential-amplifier type IC measurement.
The same set-up for 2-port parameter measurement was used, only replacing test jigs. In 3-port measurements, auxiliary termination with proper DC bias is connected to the remaining port.
7. Error approximation.

By varying the frequency of the signal generator over the frequency range where measurements are to be made, the maximum deviation of amplitude and phase errors can be approximated. The maximum deviation in the frequency range from 150 MHz to 250 MHz was $\pm 1 \%$ in amplitude and $\pm 2$ degrees in phase. It should also be noted, however, that amplitude errors are measured at full scale where reading accuracy is a maximum.

The fringing effect error is also significant particularly for the reflection coefficients phase angles when ports are terminated in a high impedance. When the shorting conductor is removed from the socket of the test jig, the phase deviation from $0^{\circ}$ is measured without changing the adjustable line. This measures the fringing effect error. (Refer to Appendix III-B.) A maximum of 5 degrees positive phase error was observed at frequencies between 150 and 250 MHz .

## EXPERIVENTAL DATA AND DISCUSSION OF RESUD'i's

S-parameters of several VHF semiconductor devices are obtained with the measurement set-up discussed in Chapter IV. This chapter will be devoted to the presentation of the experimental results based on the discussions of Chapter III. A. Bipolar Transistor - Common-Emitter 2-Port Parameters The VHF range is a relatively low frequency category for these devices. At VHF frequencies, the time delay due to minority carrier transit time is negligible and the hybrid Pi equivalent circuit of Fig. 3-1 can be employed in explaining experimental data.

Two sets of curves, $S-I_{C}$ and $S-V_{C E}$ in common-emitter configurations, are measured in order to study the nonlinear behavior. The resulting curves are shown in Figs. 5-1 and 5-2. The frequency was fixed at 150 MHz for these data.

1. S-Parameters versus Collector Current Curves.

$$
\text { 1-1. } S_{11}-I_{C} \text { (Fig. 5-1 (a)). The results are plot- }
$$ ted on a Smith Chart (Fig. 5-3) which gives the input impedance with a $50-0 h m$ terminated output. The reciprocal of this impedance is approximately $y_{11}$, since 50 -ohm output termination impedance is relatively small when compared with the output impedance of the device. Internal feedback is also very small at 150 MHz .



Fig. 5-1(a). Common-emitter $\mathrm{S}_{11}-\mathrm{I}_{\mathrm{C}}$ curve.


Fig. 5-1(b). Comnon-emitter $\mathrm{S}_{22}-\mathrm{I}_{\mathrm{C}}$ curve.



## The Smith Chart shows that the locus of $S_{11}$ with

 increasing $I_{C}$ lies on a constant resistant circle. This implies the effect of the base spreading resistance, $r_{b b}$. of Fig. 3-1. As was discussed in Chapter III, $r_{b b}$, is constant regardless of the DC operating point. The value of $r_{b b}$, from the Smith Chart of Fig. 5-1 is constant at about 50 ohms.$$
\text { 1-2. } S_{22}-I_{C} \text { (Fig. 5-1 (b)). } S_{22} \text { is generally not }
$$ as sensitive to $I_{C}$ variations as $S_{11}$; however, the magnitude of $S_{22}$ decreases gradually with increasing $I_{C}$, while the angle, $/_{22}$ is almost constant.

These variations can be explained by a decreasing dynamic resistance, $r_{c e}$. The variation of $/_{22}$ in the region of $I_{C}$ less than 1 mA is mainly due to the variations in the capacitance $\mathrm{C}_{\mathrm{b}} \mathrm{cc}^{\text {. }}$ This capacitance is a function of the collector-base junction voltage.

$$
\text { 1-3. } \mathrm{S}_{21}-\mathrm{I}_{\mathrm{C}} \text { (Fig. 5-1 (c)). This curve is analo- }
$$ gous to $g_{m}-I_{C}$ curves (or more closely $y_{21}-I_{C}$ curves). The result can be checked in Appendix $I . S_{21}$ can be approximated assuming $S_{12}=0$ by

$$
\begin{equation*}
s_{21}=-\frac{1}{2} y_{21}\left(1+s_{11}\right)\left(1+s_{22}\right) \tag{5-1}
\end{equation*}
$$

where $y_{21}$ is normalized forward transfer admittance.
From Eq. (5-1) it can be seen that $S_{21}$ is in proportion to $I_{C}$ (or $g_{m}$ where $g_{m}=\lambda I_{C}$ ) for small collector currents
(i.e. less than 1 mA ). Note that the impedance between the $e$ and $b^{\prime}$ nodes of the hybrid $P i$ circuit is much larger than $r_{b b}$, and $y_{21}$ is approximately proportional to $g_{m}$ for this collector current. The variations due to the product, $\left(1+S_{11}\right)\left(1+S_{22}\right)$ of Eq. (5-1), do not appear explicitly in the magnitude data of $S_{21}$.
$\left|S_{21}\right|$ deviates from a straight line with increasing $I_{C}$ because the impedance between $b^{\prime}$ and e decreases inversely with $I_{C}$. The terms $y_{21}$ and $\left(1+S_{11}\right)$ decrease as the results in Fig. 5-1 (c) show.

If the measurement frequency is sufficiently high the impedance between $b^{\prime}$ and $e$ is mainly due to $C_{b} e^{\prime}$ i.e., $\omega r_{b} e^{C_{b}}{ }^{\prime}{ }^{\prime} \gg 1$. Therefore, at high frequencies $y_{21}$ can be approximated in terms of $g_{m}, C_{b} e^{\prime}$, and $r_{b b}$, where the feedback elements $C_{b}{ }^{\prime} c$ and $r_{b}{ }^{\prime}$ are neglected.

$$
\begin{equation*}
y_{21} / Z_{o}=\frac{g_{m} / j \omega C_{b^{\prime} e}}{r_{b b^{\prime}}+1 / j \omega C_{b^{\prime} e}} \tag{5-2}
\end{equation*}
$$

The denominator ( $r_{b b}{ }^{\prime}+1 / j \omega C_{b^{\prime}}$ ) is approximately equal to the input impedance of the device at this frequency. The input impedance can then be expressed in terms of the S-parameter relationships

$$
\begin{equation*}
\left(r_{b b}+1 / j \omega c_{b^{\prime} e^{\prime}}\right) / L_{o}=\left(1+S_{11}\right) /\left(1-S_{11}\right) \tag{5-3}
\end{equation*}
$$

where $S_{11}$ is the common-emitter input reflection coefficient. When these results are substituted into Eq. (5-1), the phase angle of $S_{21}$ takes a simple form:

$$
\begin{equation*}
\angle S_{21}=90^{\circ}+\angle 1-S_{11}+\angle 1+S_{22} \tag{5-4}
\end{equation*}
$$

For collector currents greater than 1 mA , Eq. (5-4) agrees with the data given in Fig. 5-1. For example, the phase angle from Eq. (5-4) using measured values for $S_{11}$ and $S_{22}$ at $I_{C}=2 \mathrm{~mA}$ is about $122^{\circ}$ and the measured phase angle of $S_{21}$ is $118^{\circ}$. This result can be used to check the validity of the bipolar hybrid Pi circuit, i.e., another application of Eq. $(5-4)$ is the estimation of minority carrier transit time. Since Eq. $(5-4)$ is derived with the assumption of real $g_{m}$, the phase shift due to minority carrier transit time is not considered. The phase difference between medsure $S_{21}$ and calculated $S_{21}$ then implies this phase shift, which is, for the above data, $4^{\circ}$. In the light of the assumptions for Eq. (5-4), this angle, $4^{\circ}$, is negligible. At higher frequencies these angles will be more significant.

$$
\text { 1-4. } S_{22}-I_{C} \text { (Fig. 5-1 (d)). The same argument for }
$$ Eq. (5-1) can be applied for $S_{12}$ by replacing $y_{21}$ and $S_{21}$ in the equation with $\mathrm{y}_{12}$ and $\mathrm{S}_{12}$, respectively. For negligible $\left|\mathrm{S}_{12} \mathrm{~S}_{21}\right|$

$$
\begin{equation*}
S_{12}=-\frac{1}{2} y_{12}\left(1+S_{11}\right)\left(1+S_{22}\right) \tag{5-5}
\end{equation*}
$$

where $y_{12}$ is normalized reverse transfer admittance.
$Y_{12}$ is now given in terms of hybrid Pi parameters.

$$
\begin{equation*}
y_{12} Z_{0}=-\frac{Y_{b^{\prime} c}}{Y_{b^{\prime} e}} \cdot \frac{1}{r_{b b^{\prime}}+1 / Y_{b^{\prime} e}} \tag{5-6}
\end{equation*}
$$

where

$$
\begin{aligned}
& Y_{b^{\prime} c}=1 / r_{b^{\prime} c}+j \omega c_{b^{\prime} c} \\
& Y_{b^{\prime} e}=1 / r_{b^{\prime} e}+j \omega c_{b^{\prime} e}
\end{aligned}
$$

but ( $r_{b_{b}}+1 / Y_{b^{\prime} e^{\prime}}$ ) is approximately the input impedance, for $Y_{b}{ }^{\prime}$ is assumed to be negligible. Then the scattering parameter relationship is as follows

$$
\begin{equation*}
\left(r_{b b^{\prime}}+1 / Y_{b^{\prime}}\right) / R_{o}=\left(1+S_{11}\right) /\left(1-S_{11}\right) \tag{5-7}
\end{equation*}
$$

By substituting Eq. (5-6) and Eq. (5-7) into Eq. (5-5), $\mathrm{S}_{12}$ is finally expressed in the following forms

$$
\begin{equation*}
S_{12}=\frac{Y_{b^{\prime} c}}{Y_{b^{\prime} e}}\left(1-S_{11}\right)\left(1+S_{22}\right) \tag{5-8}
\end{equation*}
$$

Since the $Y_{b}{ }_{c}$ and $\left(1+S_{22}\right)$ terms in Eq. (5-8) are not sensitive to collector current, ( $1-S_{11}$ ) and $Y_{b \cdot e}$ terms are the main source of $S_{12}$ variation. However, the magnitude of $\mathrm{S}_{12}$ in Fig. 5-1 (d) is approximately constant over wide variations of $I_{C}$, implying that the magnitudes of $\left(1-S_{11}\right)$ and $Y_{b} e$ increase with increasing $I_{C}$.

Where the measurement frequency is sufficiently high,
i.e. $\omega r_{b} e^{C_{b} e^{\prime} \gg 1, Y_{b} e^{\prime}}$ is approximately $j \omega C_{b^{\prime} e^{\prime}}$. Then the phase relations for Eq. (5-8) can be written as

$$
\begin{equation*}
S_{12}={\underline{Y_{b}^{\prime} c}}-90^{\circ}+/ 1-S_{11}+/ 1+S_{22} \tag{5-9}
\end{equation*}
$$

but $\mathcal{Y}_{b^{\prime} \underline{c}}$ does not vary with $I_{C}$ and can be assumed constant. Eqs. (5-9) and (5-4) show that the curves for $/ \mathrm{S}_{12}$ and $/ \mathrm{S}_{21}$ are parallel with a constant angular difference of $180^{\circ}$ M $\underline{b}^{\circ}{ }^{\circ}$. This constant angle in Fig. $5-1$ is about $65^{\circ}$ indicoating $Y_{b^{\prime} c}$ is about $115^{\circ}$.
2. S-Parameters versus Collector Voltage Curves.

Severe parameter changes can be observed in the saturaLion region. This is where the collector voltage fails to provide sufficient reverse bias at the collector junction for encouraging the drift movement of minority carriers. To understand parameter behavior at low collector voltage, a transistor can be thought of as consisting of two diodes. At low collector voltages the base bias voltage may be higher than the collector voltage and the two diodes are effectively forward-biased. Input impedance then becomes much smaller in this case as compared with the input impedance when sufficient collector voltage is applied.

$$
\text { 2-1. } S_{11}-V_{C E} \text { (Fig. 5-2 (a)). } S_{11}-V_{C E} \text { curve }
$$

shows the behavior just discussed. Where $V_{C E}$ is smaller than 1 V , the magnitude and phase curves drop down very quickly according to the forward-biased diode characterise-


Fig. 5-2(a). Common-emitter $S_{11}-V_{C E}$ curve.


Fig. 5-2(b). Common-emitter $\mathrm{S}_{22} \mathrm{~V}_{\mathrm{CE}}$ curve.



Fig. 5-2(c). Common-emitter $\mathrm{S}_{21}-\mathrm{V}_{\mathrm{C}} \mathrm{E}^{\text {curve. }}$



But the curves become much flatter once sufficient collector voltages are supplied; input impedance now mainly depends on the base-to-emitter diode. (See Fig. 5-3.)

$$
\text { 2-2. } \mathrm{S}_{22}-\mathrm{V}_{\mathrm{CE}} \text { (Fig. 5-2 (b)). In the saturation }
$$ region the same explanation is valid for the $S_{22}-V_{C E}$ curve; the output impedance corresponding to $S_{22}$ shows the forwardbiased diode characteristic until sufficient reverse bias is provided between the base-collector junction. $S_{22}$ varies with increasing $V_{C E}$ mainly due to the decrease of $C_{b}{ }^{\circ} c$ and the increase of dynamic resistance $r_{c e}$.

2-3. $S_{21}-V_{C E}$ (Fig. 5-2 (c)). By recalling Eq. (5-1) the behavior of $S_{21}$ is more clearly understood. Since $I_{C}$ is held constant, $y_{21}$ can also be assumed a constant. The phase variation then is approximately the sum of the varialions of $/ 1+S_{11}$ and $/ 1+S_{22}$. The data shown in Fig. 5-2 (c) agree very well with this type variation.

## It can also be noted that $S_{21}$ is almost constant once

 $\mathrm{V}_{\mathrm{CE}}$ passes from the saturation region. This point can be easily understood by referring to the hybrid equivalent circuit. Under normal operating conditions most of the hybrid $P_{i}$ parameters are unaffected by changes in collector voltage. $C_{b}{ }^{\prime}$ c is collector voltage dependent, but is not significant in the $S_{21}$ calculations.$$
\text { 2-4. } \mathrm{S}_{12}-\mathrm{V}_{\mathrm{CE}} \text { (Fig. 5-2 (d)). Lastly, } \mathrm{S}_{12}-\mathrm{V}_{\mathrm{CE}}
$$

curves are shown. The decrease of $\left|S_{12}\right|$ with increase of $V_{C E}$
implies an effective decrease in the feedback admittance, $Y_{b}{ }^{\prime}$. See Eq. (5-8).

By now, interpretation of the experimental results for a bipolar device has been made. At higher frequencies, such as in microwave frequency range, the transit time of minority carriers would have to be taken into account. Investigatior: of this point of view is beyond the scope of this thesis and is not discussed.
B. MOSFETs (Single-Gate and Dual-Gate) - Common-Source 2-Port Parameters

Measurements on MOSFET transistors are made in the same manner as for bipolar transistor measurements. During dualgate MOSFET measurements, additional bias for gate 2 is provided through an RF-grounded terminal. Grounding of this terminal eliminates parasitic effects. The case and substrate are also in common with the source terminal for common-source operation. The test jig is shown in Fig. 4-2(a).

Reading $\left|S_{11}\right|$ and $\left|S_{22}\right|$ requires great care since these values are generally close to unity for low-frequency (i.e. 150 MHz ) common-source measurements. Sophisticated calibration is required for these measurements.

A heat sink may have to be prepared to prevent thermal runaway or excessive excursions of the operating point whenever the transistor is operated at large DC input levels. Severe changes in the forward transmission coefficient are
observed for large DC inputs. Readjustment of the bias voltages to compensate for the decrease of drain current can result in a wide variation of parameter values.

1. Single-gate MOSFET - Common-source $S-V_{D S}$ Curves. 1-1. $S_{11}-V_{D S}$ (Fig. 5-4 (a)). A family of curves is obtained for various drain voltages. From typical values of the hybrid Pi parameters, for a high-frequency MOSFET (Fig. 3-2 (b)), it can be seen that $S_{11}$ will largely depend on $C_{c}$ and $r_{c}$, the gate-channel capacitance and the channelsource bulk resistance. But the behavior of $r_{c}$ is different from that of $r_{b b}$, of bipolar devices; $r_{c}$ as well as $C_{c}$ vary with the gate bias voltage $\mathrm{V}_{\mathrm{GS}}$ •

The input impedance of a MOSFET can be approximated by $r_{c}$ and $C_{c}$ as a series impedance and shunt capacitance $C_{g s}, r_{g s}, C_{g d}$, and $r_{g d}$ are neglected. Now the input impedance $Z_{i}$ can be represented in terms of $r_{c}, C_{c}$, and $C_{g s}$.

$$
\begin{equation*}
z_{i}=\frac{r_{c}+1 / j \omega c_{c}}{\left(c_{c}+c_{g s}\right) / c_{c}+j \omega r_{c} C_{g s}} \tag{5-10}
\end{equation*}
$$

For typical MOSFET data, ${ }^{20} \mathrm{C}_{\mathrm{gs}}$ is much smaller than $\mathrm{C}_{\mathrm{c}}$ and $\omega r_{c} c_{g s}$ term is negligible at 200 NHz . Then

$$
\begin{equation*}
z_{i}=\frac{c_{c}}{c_{c}+c_{g s}} r_{c}+\frac{1}{j \omega\left(c_{c}+c_{g s}\right)} \tag{5-11}
\end{equation*}
$$

Eq. (5-11) indicates that the contribution of $C_{g s}$ to the $r_{c}$ and $C_{c}$ series impedance causes only a small amount of



$+140^{\circ}$
Fig. 5-4(c). Common-source $\mathrm{S}_{21}-\mathrm{V}_{\mathrm{DS}}$ curve (Single-gate MOSFET).



Fig. 5-5. Smith Chart plot of common-source $S$-parameters:

$$
\text { (a) } S_{11}-V_{D S} i^{\prime}(\mathrm{b}) \mathrm{S}_{22}-V_{D S} \text { (Single-gate MOSFET). }
$$

decrease in the magnitude. In other words input impedance still can be approximated by a series impedance consisting of $r_{c}$ ' and $C_{c}$ ' where $r_{c}$ ' and $C_{c}$ ' are defined by the equivalent resistance and capacitance of $Z_{i}$ in Eq. (5-11). This impedance can be obtained directly from a Smith Chart plot of $S_{11}$. (Refer to Fig. 5-5.) The variations of $r_{c}{ }^{\prime}$ and $1 / \omega c_{c}$ ' observed in this manner are in the range of 10 to 30 ohms and 200 to 250 ohms, respectively, at 200 NHz . Assuming $C_{c}=5 . C_{g s}$, the variations of $r_{c}$ and $1 / \omega C_{c}$ are approximately in the range of 12 to 36 ohms and 240 to 300 ohms, respectively. The exact value depends on the operating point.

When $V_{D S}$ is increased the gate-channel voltage will also slightly increase because of the DC potential drop across the channel-source DC resistance. The increase of $/ \mathrm{S}_{11}$ in Fig. 5-4 (a) implies this effect; once $\mathrm{V}_{\mathrm{DS}}$ reaches the pinch-off region, the $D C$ drop will not change, resulting in constant $/ \mathrm{S}_{11}$. Anomalous inversion of $/ \mathrm{S}_{11}$ curves (bend downward) at low drain voltages seems to be the result of another nonlinearity of the MOS capacitor.

$$
\text { 1-2. } \left.\mathrm{S}_{22}-\mathrm{V}_{\mathrm{DS}} \text { (Fig. } 5-4(\mathrm{~b})\right) \text {. A set of drain }
$$

characteristics can be used to explain the behavior of $S_{22}$. In the triode region the dynamic output resistance is relatively small, resulting in low $\left|S_{22}\right|$ But this resistance increases and reaches saturation in the pinch-off region.
resulting in an $\left|S_{22}\right|$ that is close to unity.
In conjunction with the dynamic resistance variation, the drain-source capacitance decreases. The capacitance changes because of an effective increase in the width of the drain-source depletion region. ${ }^{21}$ The positive phase angles of $\mathrm{S}_{22}$ in the figure are partly due to phase error of the measurement system (fringing effect error) and transistor lead lengths.

$$
\text { 1-3. } S_{21}-V_{D S} \text { (Fig. 5-4 (c)). } S_{21} \text { data are cons- }
$$ dered under two conditions: $S_{21}$ versus $V_{D S}$ and $S_{21}$ versus $\mathbf{V}_{\text {GS }} \cdot S_{21}-V_{G S}$ relationship can be read directly from the figure for an arbitrary $V_{D S}$. On the other hand, $\left|S_{21}\right|-V_{D S}$ curves look like that of $g_{m}-V_{D S}$ relationship except at very small $\mathrm{V}_{\mathrm{DS}}$ •

This can be explained using Eq. (5-1), in which $\mathrm{S}_{12}$ is assumed negligible, ie. $\left|S_{12} S_{21}\right| \ll\left|\left(1+S_{11}\right)\left(1+S_{22}\right)\right|$. Eq. (5-1) can be applied to this MOSFET.
$y_{21}$ can now be derived from the hybrid $P_{i}$ equivalent circuit in Fig. 3-2 (b).

Assuming $r_{c} 《 1 / \omega C_{c}$ and $r_{g d} \gg 1 / \omega C_{g d}$

$$
\begin{equation*}
y_{21} / Z_{o}=g_{m}+j \omega c_{g d} \tag{5-12}
\end{equation*}
$$

When Eq. $(5-12)$ is substituted into Eq. $(5-1), S_{21}$ can be expressed as

$$
\begin{equation*}
s_{21}=-\frac{1}{2}\left(g_{m}+j \omega c_{g d}\right)\left(1+S_{11}\right)\left(1+s_{22}\right) Z_{o} \tag{5-13}
\end{equation*}
$$

$g_{m}$ and $C_{g d}$ in Eq. (5-13) are obviously functions of $V_{\text {Dis }}$. For an arbitrary $V_{G S}$, $g_{m}$ increases with increasing $V_{\text {BS }}$ (triode region) and it becomes constant when sufficient $V_{D S}$ is applied (pinch-off region). The $C_{g d}-V_{D S}$ curve is the opposite of $g_{m}-V_{D S}$ relationship; $C_{g d}$ decreases rapidly and approaches a minimum value, the extrinsic drain-gate capacitance. ${ }^{20} \mathrm{~g}_{\mathrm{m}}$ and $\mathrm{C}_{\mathrm{gd}}$ also vary with $\mathrm{V}_{\mathrm{GS}}$. With increasing $V_{G S}$, both $g_{m}$ and $C_{g d}$ increase.

First, the $S_{21}-V_{D S}$ relationship can clearly be understood using Eq. (5-13). At $\mathrm{V}_{\mathrm{DS}}=0, \mathrm{~g}_{\mathrm{m}}$ is also 0 and $\mathrm{S}_{21}$ is determined by $\mathrm{C}_{\mathrm{gd}}$. For this case, it can be shown that $\mathrm{y}_{21}$ is equal to $y_{12}$ and $S_{21}$ is equal to $S_{12}$.

With increasing $V_{D S}, g_{m}$ dominates $C_{g d}$. For sufficiently high $V_{D S}, C_{g d}$ is negligible and $S_{21}$ can be approximated by

$$
\begin{equation*}
s_{21}=-\frac{1}{2} g_{m}\left(1+s_{11}\right)\left(1+s_{22}\right) \tag{5-14}
\end{equation*}
$$

Since $g_{m}$ is a constant in the pinch-off region, $\left|j_{21}\right|$ will approximate a constant in view of the variations of the $\left(1+S_{11}\right)\left(1+S_{22}\right)$ product. Phases can also be checked by Eq. $(5-14)$. At $V_{D S}=10 \mathrm{~V}$, an angular difference of about $10^{\circ}$ is found between the measured $/ \mathrm{S}_{21}$ and the calculated $\underline{S_{21}}$. The assumptions for $\mathrm{Eq} .(5-14), \mathrm{g}_{\mathrm{m}} \geqslant \omega \mathrm{C}_{\mathrm{gd}}$ are justified.

For explaining the $|\stackrel{s}{21}|$ absolute maximum of the $\left|\mathrm{S}_{21}\right|-\mathrm{V}_{\mathrm{GS}}$ relationship, Eq. (5-14) is useful. In the
equation, $S_{21}$ is a function of $g_{m}$ and $\left(1+S_{22}\right)$, if the $\left(1+S_{11}\right)$ term is assumed constant. $g_{m}$ and $\left(1+S_{22}\right)$ are distinctly functions of $V_{G S}$ ' one increases rapidly while the other decreases rapidly with increasing $V_{G S}$. Thus, the product will have an absolute maximum and so will $S_{21}$ if the variation of ( $1+S_{11}$ ) term in Eq. (5-14) is not so significant with respect to $V_{G S}$ •

$$
\text { 1-4. } \mathrm{S}_{12}-\mathrm{V}_{\mathrm{DS}} \text { (Fig. 5-4 (d)). Eq. }(5-5) \text { again is }
$$

employed to explain these curves. If $r_{g d}$ is ignored, $y_{12}$ can be approximated in terms of $C_{g d}$ such that

$$
\begin{equation*}
-y_{12} L_{o}=j \omega c_{g d} \tag{5-15}
\end{equation*}
$$

and combining Eq. (5-5) with Eq. (5-15),

$$
\begin{equation*}
S_{12}=\frac{1}{2} j \omega C_{g d}\left(1+S_{11}\right)\left(1+S_{22}\right) \tag{5-16}
\end{equation*}
$$

For a very low bias voltage such as $\mathrm{V}_{\mathrm{GS}}=-2.5 \mathrm{~V}$, the variations of $\left(1+S_{11}\right)\left(1+S_{22}\right)$ are not significant and $\left|S_{21}\right|$ will be directly proportional to the feedback capacitance $C_{g d}$. $S_{12}-V_{D S}$ curve for $V_{G S}=-2.5 \mathrm{~V}$ is very close to the $C_{g d}-V_{D S}$ curve. ${ }^{20}$

With increasing $V_{G S}$, the $S_{12}$ curve deviates due to $\left(1+S_{22}\right)$ term and also due to the nonlinearity of $C_{g d}$. Effectively, $C_{g d}$ can be increased by reducing the depletion layer at the drain-gate region with an increased $\mathrm{V}_{\mathrm{GS}}$. Note that the $\mathrm{S}_{12}-\mathrm{V}_{\mathrm{DS}}$ curves converge at $\mathrm{V}_{\mathrm{DS}}>10$ to a minimum
value indicating the effect of the extrinsic feedback capacitance.
2. Frequency Characteristics of Single-Gate MOSFET. Common-source S-parameters are measured and presented in Fig. 5-6. A normal Pi calculation is performed according to Eq. (3-6) and the results are plotted in Fig. 5-7.

It is seen that the rapid increase of the real part of A can be regarded as a second order effect of the $R-C$ series circuit ${ }^{21}$ combination of $r_{c}{ }^{\prime}$ and $C_{c}{ }^{\prime}$. A 25 ohm resistor and 5 pF capacitor for the R-C series circuit gives excellent agreement over the frequency range.

The apparent inductive reactance of $S_{22}$ is due to fringing effects. If an angular error of $5^{\circ}$ (see Appendix Fig. A-III-4) is subtracted from $/ S_{22}$ then the true value of $S_{22}$ is resistive and capacitive.

The frequency response of $D$ does not show the exact behavior indicated by the simplified equivalent circuit of Fig. 3-2 (b). Some modifications, such as the addition of a drain-channel resistance, may be made for a more precise representation. 20

The imaginary part of $B$ represents the capacitive susceptance of $C_{g d}$, whose value is 0.15 pF . This value agrees with the 0.2 pF as given in the manufacturer's data (RCA MOS-160). ${ }^{23}$

The real part of $B$ shows a negative conductance which cannot be explained by the equivalent circuit. It may be


Fig. 5-6(a). Frequency response of common-source

$$
\begin{aligned}
& \mathbf{v}_{\mathrm{DS}}=+15 \mathrm{v} \\
& \mathbf{v}_{\mathrm{GS}}=-2 \mathrm{v} \\
& 3 \mathrm{~N}_{142}
\end{aligned}
$$


$\mathrm{S}_{12}$

Fig. 5-6(b) Frequency response of common-source



Fig. 5-7(a). Normal $P_{i}$ versus frequency curves (Single-gate MOSFET 3N142).


Fig. 5-7(b). Normal Pi versus frequency curves (Single-gate NOSFET 3N142).
regarded as the result of improper reference plane (i.e. longer $L_{i}$ than required) or phase shift due to the distributed nature of the impedance in the channel region. The data show that the phenomenon is too distinct to be explained by measurement error, and a phase lead is also unlikely from a distributed R-C circuit viewpoint. Similar results are reported for $y$-parameters. ${ }^{20}$

The phase variation of $C$ is more gradual than expected, but it is not so distinct as in $\mathrm{S}_{12}$. It should be noted that the magnitude of $C$ is practically constant over the measurement frequency; the value of $C$ in Fig. 5-7 is approximately 8,000 umhos, which agrees with the $g_{m}$ given by the manufacturer. ${ }^{23}$
3. Dual-Gate MOSFET - Common-Source S-V $\mathrm{V}_{\mathrm{DS}}$ Curves. As was discussed earlier in Chapter III, a dual-gate MOSFET may be analyzed as a cascode amplifier consisting of 2 single-gate MOSFETs. (Refer to Fig. 3-4.) Dual-gate MOSFET S-parameters, with gate 2 grounded, can be studied from common-source and common-gate $S$-parameters of a singlegate MOSFET.

Typical values of the S-parameters of a single-gate device are obtained from Fig. 5-4 and Appendix IV. The DC operating point of each single-gate MOSFET is made to agree with the drain current and drain voltage of the dual-gate MOSFET that is studied.

Common Source $\left(200 \mathrm{MHz}, \mathrm{V}_{\mathrm{GS}}=+2.0 \mathrm{~V}\right.$, and $\left.\mathrm{V}_{\mathrm{DS}}=6 \mathrm{~V}\right)$

$$
\begin{aligned}
& \mathrm{S}_{11}=0.94 \angle-24^{\circ} \\
& \mathrm{S}_{22}=0.92 \angle-0.4^{\circ} \\
& \mathrm{S}_{21}=0.65 \angle+154^{\circ} \\
& \mathrm{S}_{12}=0.025 \angle+82^{\circ}
\end{aligned}
$$

Common_Gate $\left(200 \mathrm{MHz}, \mathrm{V}_{\mathrm{SG}}=+2.0 \mathrm{~V}\right.$, and $\left.\mathrm{V}_{\mathrm{DG}}=6 \mathrm{~V}\right)$

$$
\begin{aligned}
& \mathrm{S}_{11}{ }^{\prime}=0.45 \angle-30^{\circ} \\
& \mathrm{S}_{22^{\prime}}=0.96 \angle-1^{\circ} \\
& \mathrm{S}_{21},=0.30 \angle-15^{\circ} \\
& \mathrm{S}_{12}{ }^{\prime}=0.05 . \angle+85^{\circ}
\end{aligned}
$$

By substituting above measured values into the reflectlion coefficient formulas it can be shown that the input and output reflection coefficients of a dual-gate device can be approximated by $S_{11}$ and $S_{22}$. The forward and reverse transmission coefficients are calculated as $0.30 \angle+120^{\circ}$ and $0.0013 \angle+148^{\circ}$, respectively.

From the calculated values, one can expect that a dual-gate device will show almost the same characteristics as a common-source stage except for the large decrease in internal drain-gate feedback.

By referring to Fig. $5-8$, the $S-V_{D S}$ curves, the above approximation method can be fully justified. The only exception is that the magnitude of the reverse transmission


Fig. 5-8(a). Common-source $\mathrm{S}_{11}-\mathrm{V}_{\mathrm{DS}}$ curve
(Dual-Eate MOSFET).


Fig. 5-8(b). Common-source $\mathrm{S}_{22}-\mathrm{V}_{\text {DS }}$ curve (Dual-gate MOSFET).



Fig. 5-8(c). Common-source $\mathrm{S}_{21}-\mathrm{V}_{\mathrm{DS}}$ curve (Dual-gate MOSFET).


Fig. 5-8(d). Common-source $\mathrm{S}_{12}-\mathrm{V}_{\text {DS }}$ curve (Dual-gate MOSFET).
coefficient $S_{12}$ is a little larger than the values calculated from single-gate $S$-parameters. If substantial increase of internal feedback due to the geometry of a dual-gate MOSFET is considered, every behavior of a dualgate device can be understood in terms of single-gate S-parameters.
4. Frequency Characteristics of Dual-Gate MOSFET. Significant differences can be found in $S_{12}$ and $S_{21}$ when the curves of Fig. 5-9 are compared with those of the single-gate device in Fig. 5-6. The phase angle of $S_{12}$ is rather irregular and the slope of $/_{21}$ is steeper than that of a single-gate device. The irregular variations of $/ S_{12}$ are also confirmed at the higher frequencies. (Refer to Appendix V.) It is, however, impressive that the magnitude of $S_{12}$ increases very uniformly in spite of a phase drift.
5. Interchangeability of Gate 1 and Gate 2. From the original purpose of gate 2 in a dual-gate MOSFET, it makes sense to use gate 2 as the AGC gate with RF ground, and gate 1 as the control gate for signal input. However, the roles of the gates may be interchanged sucl. that gate 1 is used as the control gate. S-parameters for each connection are then certainly different. For example, $\left|S_{12}\right|$ is much larger when gate 2 is used as the control gate, and the phase angle is around $+90^{\circ}$ for that case.

Usually, for stable amplifier purposes, it is better


Fig. 5-9(a). Frequency response of common-source


Fig. 5-9(b). Frequency response of common-source
to use gate 2 rather than gate 1 for AGC, but for oscillator applications specific phase angles of S-parameters may be required and the use of gate 1 may be more preferable for this purpose. In Appendix VII S-parameters for the two different connections are given.
6. Consideration of MOSFET as an Amplifier. Since $\left|S_{11}\right|$ and $\left|S_{22}\right|$ for the common-source configuration are close to unity at frequencies below 200 MHz (Fig. 5-6), a unilateral figure of merit $U$ cannot be neglected, particularly for single-gate MOSFETs. Typical values of $U$ and $G_{U}$ at 200 MHz are 0.645 and 54.5 , respectively. The lower the frequency, the worse the figure $U$ becomes. In this case mismatching methods ${ }^{10,22}$ may be employed using Eq. (2-20)'.

For dual-gate MOSFEts, $G_{U}$ and $U$ are obtained and plotted in Appendix VI. The unilateral figure of merit in this case is between 0.10 and 0.13 for the frequency range.

It should be noted that $U$ and $G_{U}$ are also functions of the DC operating point so proper selection of the operating point is necessary for optimization.

So far, single- and dual-gate MOSFET S-parameters have been discussed. It has been shown that other 2 -port parameters can also be calculated from measured S-parameters and that these values agree with the data presented in the literature. Dual-gate devices can be analyzed on the basis of measured single-gate S -parameters.
C. Three-Port Parameters

A set of 3-port parameters has been attained for a bipolar transistor. In order to determine 9 components of ( $3 \times 3$ ) matrix, 9 individual measurements (18 numbers, i.e. magnitudes and phase angles of 9 complex numbers) have been made. At $I_{C}=1 \mathrm{~mA}$ and $\mathrm{V}_{\mathrm{CE}}=6 \mathrm{~V}$, the scattering matrix obtained at 200 MHz is as follows

## Table I

$$
\left[\begin{array}{lll}
S_{\text {ee }} & S_{e b} & S_{e c} \\
S_{b e} & S_{b b} & S_{b c} \\
S_{c e} & S_{c b} & S_{c c}
\end{array}\right]=\left[\begin{array}{lll}
\underline{0.39 \angle+94^{\circ}} & \underline{1.05 L-27^{\circ}} & \underline{0.107 \angle+59^{\circ}} \\
\underline{0.58 \angle+51^{\circ}} & \underline{0.72 L-52^{\circ}} & \underline{0.075 \angle+70^{\circ}} \\
\underline{1.00 /-50^{\circ}} & \underline{1.00 \angle+127^{\circ}} & \underline{0.86 /-19^{0}}
\end{array}\right]
$$

3-Port S-Parameters, GM 0290
where e, b, and c denote the emitter, base, and collector terminals.

According to Eq. (2-29), each row and column can be summed in order to check whether each sum is unity. However, in actual devices, some RF current may flow directly to ground through stray capacitances producing an additional current path and the sums of each row and column may deviate from unity. In other words, the device is no more a 3terminal device due to stray capacitance between the device and ground. The data given above also show some deviations from 3-terminal device theory. Certainly, some measurement
errors were present.
From the data, 2-port S-parameters of any configuration (or, more generally, any termination) can be derived using the formulas of Eq. (2-27). For instance, common-emitter 2 -port parameters can be obtained by substituting $S_{3}=-1$ (short termination) into the formulas. Here, the port indices 1, 2, and 3 are replaced with b, c, and e, respectively, for convenience.

Two sets of common-emitter S-parameters, one calculated from 3-port parameters and the other from direct 2-port measurements, are given as follows:

## Table II

|  | Calculated from <br> 3-port $S$-parameters | Results of direct <br> 2-port measurement |
| :--- | :---: | :--- |
| $\mathrm{S}_{11}$ | $0.56 \angle-103^{\circ}$ | $0.65 \angle-85^{\circ}$ |
| $\mathrm{S}_{22}$ | $0.76 \angle-20^{0}$ | $0.81 \angle-10^{\circ}$ |
| $\mathrm{S}_{21}$ | $2.00 \angle+128^{\circ}$ | $1.87 \angle+115^{\circ}$ |
| $\mathrm{S}_{12}$ | $\underline{0.04 \angle+35^{\circ}}$ | $\underline{0.08 \angle+57^{\circ}}$ |

Common-Emitter S-Parameters
$\mathrm{GM} 0290,200 \mathrm{MHz}, \mathrm{V}_{\mathrm{CE}}=6 \mathrm{~V}, \mathrm{I}_{\mathrm{C}}=1 \mathrm{~mA}$
The two sets of parameters are not in close agreement, although the proximity of their values can be appreciated.

The use of different test jigs and the parasitic effects due to stray capacitances are considered as the main sources of error. For further application of 3-port parameters, refer to Bodway. ${ }^{17}$
D. Four-Port Parameters

For the IC of Fig. 3-5, complete 4-port S-parameters are obtained; however, only 8 parameters are measured in order to determine a full set because of the symmetry of $Q_{1}$ and $Q_{2}$. The 4-port parameters obtained for the RCA CA 3049 integrated circuit are as follows:

Table III

| S-parameters | 2 mA | 4 mA | 6 mA |
| :---: | :---: | :---: | :---: |
| $\mathrm{S}_{11}\left(11^{\prime}\right)$ | $0.185 /-25.3^{0}$ | $0.77, \angle-26.5^{\circ}$ | $0.71 \angle-28.5^{\circ}$ |
| S 1 $^{\prime \prime}$ (11') | $0.184 /+31^{\circ}$ | $0.237 . L+21.4^{0}$ | $0.26 / \pm 16.4^{\circ}$ |
| S22(2.2.) | $0.90 \angle-19^{\circ}$ | $0.88 /-12^{0}$ | $0.86 /-12.5^{\circ}$ |
| S2.2(22.) | $0.035 \angle+21^{\circ}$ | 0.041. $\angle+51.5^{\circ}$ | $0.043 /+32^{0}$ |
| S21(211) | 1.10/+123.5 ${ }^{\circ}$ | 1.56/+115 ${ }^{\circ}$ | $1.95 /+110.2^{0}$ |
| S21(21.) | 1.03 $/-66^{\circ}$ | 1.46/-72.5 ${ }^{\circ}$ | 1.70/-77.2 ${ }^{\circ}$ |
| S12(1'2') | $0.040 /+85.5^{\circ}$ | $0.039 /+86^{\circ}$ | $\underline{0.038 /+83.5^{\circ}}$ |
| S 12(12') $^{\prime}$ | $0.014 /+130^{\circ}$ | $0.015 /+123^{\circ}$ | $0.014 /+118^{\circ}$ |
| 4-Port S-Parameters <br> RCA CA 3049, $200 \mathrm{NMHz}, \mathrm{V}_{\mathrm{C}}=+3 \mathrm{~V}, \mathrm{~V}_{\mathrm{B}}=+1.5 \mathrm{~V}$ |  |  |  |

The 2-port common-emitter S-parameters of the coilstituent transistors $Q_{1}$ and $Q_{2}$ can be obtained using Eq. (3-5). The results are given in the following table.

Table IV

|  | $I_{C}=1 \mathrm{~mA}$ | 2 mA | 3 mA |
| :---: | :---: | :---: | :---: |
| $\mathrm{~S}_{11}$ | $\underline{0.72 /-38^{\circ}}$ | $\underline{0.60 /-40^{\circ}}$ | $\underline{0.55 /-45^{\circ}}$ |
| $\mathrm{S}_{22}$ | $\underline{0.89 /-20.2^{\circ}}$ | $\underline{0.86 /-20.1^{\circ}}$ | $\underline{0.84 \angle-20.1^{\circ}}$ |
| $\mathrm{S}_{21}$ | $\underline{2.10 /+118^{\circ}}$ | $\underline{2.98 \angle+110^{\circ}}$ | $\underline{3.60 \angle+102^{\circ}}$ |
| $\mathrm{S}_{12}$ | $\underline{0.034 \angle+62^{\circ}}$ | $\underline{0.032 \angle+61^{\circ}}$ | $0.030 \angle+61^{\circ}$ |

CA 3049, $Q_{1}$ and $Q_{2}$ C-E S-parameters $200 \mathrm{MHz}, \mathrm{V}_{\mathrm{CE}}=2.2 \mathrm{~V}$ (approximated by assuming $\mathrm{V}_{\mathrm{BE}}=0.7 \mathrm{~V}$ for silicon transistors)

These results are insufficient data to discuss the general characteristics of IC transistors; however, a little bit of information can be obtained by making a comparison with the data on discrete transistors. Several discrete VHF bipolar transistors and IC transistors are compared in Table V. Two isolated transistors on an IC chip (CA 3018) are also separately measured and the data are added in the table.

## Table V



Comparison of Several Bipolar Devices
(C-E S-parameters)
In Table V, the difference between IC transistors and discrete transistors can be found in $/ \mathrm{S}_{22}$; calculation (see Appendix II) shows that in IC transistors output capacitance (obtained from $S_{22}$ ) values are substantially larger than those of the feedback capacitance. The difference can be explained by collector-to-substrate capacitance in IC which ranged up to 2 pF .24

## CHAPTER VI

## CONCLUSIONS

Scattering parameters of several VHF semiconductor devices have been obtained from a special measurement setup. High-frequency characteristics of the device have been analyzed by means of their measured scattering parameters. Hybrid $\mathrm{Pi}_{\mathrm{i}}$ equivalent circuits of bipolar transistors and single-gate MOSFETs can be extended to VHF frequency range in interpreting most of $S$-parameter data. However, modifications are necessary in view of the fact that some parameters such as $S_{21}$ in bipolar transistors (due to minority carrier transit time) and $S_{12}$ of a MOSFET which started showing an anomalous frequency response.

The S-parameters of bipolar devices generally follow expected behavior as derived from the basic theory. The bipolar transistor hybrid Pi equivalent circuit can be approximated from common-emitter S-parameters (see Appendix II). The minority carrier transit time; which appears in $\mathrm{S}_{21}$ of common-emitter configuration, is not significant at 150 MHz . When sufficient reverse bias is applied across the base-collector region, the differences given by Eq. (5-4) will be a good approximat: on of the phase delay due to minority carrier transit time. This linear phase shift should be considered in $g_{m}$ for the bipolar hybrid Pi circuits at
higher frequencies.
There are formulas that given conversions from one set of parameters to another. Conversions between S-parameters and $y$-parameters are most convenient in the light of parameter values for the ordinary configurations for VHF devices. The normal Pi equivalent circuit obtained from y-parameters is useful for a single-gate MOSFsir representation. From the circuit, most desired design parameters such as $g_{m}$, $C_{g d}, C_{c}, r_{c}$, etc. can be obtained.

The common-source S-parameters of a dual-gate MOSFET can be closely approximated by two sets of single-gate MOSFET S-parameters: one of a common-source, and the other of a common-gate. The significant difference between the actual data and estimated values for single-gate MOSFET S-parameters can be found in the magnitudes of reverse transmission coefficient: the measured $S_{12}$ is much larger than values estimated from single-gate MOSFET S-parameters.

Differential amplifier type ICs are fully characterized by 4-port S-parameters. Constituent transistor $S$-parameters can be obtained using a simple formula. From calculated values of constituent transistor S-parameters, no significant differences between an IC transistor and a discrete transistor can be found. Only the phase angle of $S_{22}$ of the IC transistor shows a relatively higher shunt capacitance than that of a discrete transistor of the same kind.

As amplifier design parameters, a set of MOSFET S-parameters in VHF frequency range suggests that unilateral design for maximum gain may not be used, because of the relatively large unilateral figure of merit. This is because the magnitudes of $S_{11}$ and $S_{22}$ of the common-source MOSFET S-parameters are usually very close to unity at frequencies below 200 MHz . It also suggests that these devices should be assumed as voltage amplifiers for design purposes at these frequencies as was true for vacuum tube designs.

Since the common-emitter or common-source output impedances are usually very large for a normal DC operating point, significant fringing effect error (short and open position difference of a reference plane) can be observed. At frequencies between 150 MHz and 250 MHz , a maximum of +10 degrees could be found for the 2 -port test jig of Fig. 4-2 (a). So some of these angles should be subtracted from the corresponding $/ \mathrm{S}_{22}$ in order to eliminate the error. In conclusion, the following statement can be made regarding $S$-parameters. S-parameters for VHF semiconductor devices can be obtained with sufficient accuracy and the results can be justified using simple equivalent circuits. The next step is to apply S-parameters to high-frequency active circuit design.

The superiority of the $S$-parameter method is distincts
first in the convenience of making measurements; second in the simple design methods which are based on power relationships. The ample and accurate data obtained from the S-parameter method then can be combined with computers, resulting in fast and precise high-frequency circuit designs.

## REFERENCES

1. Campbell, G.A. and R.M. Foster, "Maximum Output Networks for Telephone Substation and Repeater Circuits," Transe of AIEE, Vol. 39, 1920.
2. Montgomery, C.. R.H. Dicke, and E.M. Purcell, Principles of Microwave Circuits. McGraw-Hill, Inc., 1948.
3. Oono, Y. and K. Yasuura, "Synthesis of Finite Passive 2n-Terminal Networks with Prescribed Scattering Matrices," Annales des Telecommunication, Vol. 9 , 1954.
4. Youla, D.C., "On Scattering Matrices Normalized to Complex Port Numbers," Proc. I.R.E.. Vol. 49, July, 1961 .
5. Penfield, P. Jr., "Noise in Negative Resistance Amplifier," IRE Trans, - CT, Vol. CT-7, June, 1960.
6. Lange, J. "Microwave Transistor Characterization Including S-Parameters," S-Parameters - Circuit Analysis and Design, Hewlett-Packard, Application Note 95, September, 1968.
7. Weinert, "Scattering Parameters Speed Design of HighFrequency Transistor Circuits," Electronics. September, 1966.
8. Anderson, D., "S-Parameter Techniques for Faster, More Accurate Network Design," Hewlett-Packard Journale February, 1967.
9. Froehner, W., "Quick Amplifier Design with Scattering Parameters," Electronics. October, 1967.
10. Bodway, G., "Two-Port Power Analysis Using Generalized Scattering Parameters," Microwave Journal. May, 1967.
11. Robichaud, Boisvert, and Robert, Signal Flow Graphs and Applications. Prentice-Hall, 1962.
12. Hewlett-Packard, Inc., "Transistor Parameter Measurement," Application Note 77-1, February, 1967.
13. Collin, R., Foundation for Microwave Engineering, McGraw-Hill, 1966.
14. Carlin and Giordano, Network Theory, Prentice-Hall, 1964.
15. Mason, S.J., "Feedback Theory: Some Properties of Signal Flow Graphs," Proc. I.R.E.. September, 1953.
16. , "Feedback Theory: Further Properties of Signal Flow Graphs," Proc. I.R.E., July, 1956.
17. Bodway, G., "Circuit Design and Characterization of Transistors by Means of Three-Port Scattering Parameters," Microwave Journal, Vol. II, No. 5, May, 1968.

18. Cochrun, B., Transistor Circuit Engineering, The MacMillan Co.., 1967.
19. Wallmark and Johnson, Field Effect Transistors. Prentice-Hall, 1966.
20. Walston and Miller, Transistor Circuit Design, McGrawHill, 1963.
21. Sevin, Field Effect Transistors, McGraw-Hill, 1965.
22. RCA Product Guide, "MOS Field Effect Transistor," MOS-160, December, 1967.
23. RCA Product Guide, "Linear Integrated Circuit," File No. 378, January, 1969.

## APPENDICES

## APPENDIXI

CONVERSION TABLE BETWEEN SCATTERING AND IMMITTANCE PARAMETERS

$$
\begin{array}{ll}
S_{11}=\frac{\left(z_{11}-1\right)\left(z_{22}+1\right)-z_{12} z_{21}}{\left(z_{11}+1\right)\left(z_{22}+1\right)-z_{12} z_{21}} & z_{11}=\frac{\left(1+S_{11}\right)\left(1-S_{22}\right)+S_{12} S_{21}}{\left(1-S_{11}\right)\left(1-S_{22}\right)-S_{12} S_{21}} \\
S_{12}=\frac{2 z_{12}}{\left(z_{11}+1\right)\left(z_{22}+1\right)-z_{12} z_{21}} & z_{12}=\frac{2 S_{12}}{\left(1-S_{11}\right)\left(1-S_{22}\right)-S_{12} S_{21}} \\
S_{21}=\frac{2 z_{21}}{\left(z_{11}+1\right)\left(z_{22}+1\right)-z_{12} z_{21}} & z_{21}=\frac{2 S_{21}}{\left(1-S_{11}\right)\left(1-S_{22}\right)-S_{12} S_{21}} \\
S_{22}=\frac{\left(z_{11}+1\right)\left(z_{22}-1\right)-z_{12} z_{21}}{\left(z_{11}+1\right)\left(z_{22}+1\right)-z_{12} z_{21}} & z_{22}=\frac{\left(1+S_{22}\right)\left(1-S_{11}\right)+S_{12} S_{21}}{\left(1-S_{11}\right)\left(1-S_{22}\right)-S_{12} S_{21}}
\end{array}
$$

Noter $\mathrm{h}=, \mathrm{y}$-, and z -parameters listed above and in the succeeding pages are all normalized to $Z_{0}$. See end of this appendix.

$$
\begin{array}{ll}
s_{11}=\frac{\left(1-y_{11}\right)\left(1+y_{22}\right)+y_{12} y_{21}}{\left(1+y_{11}\right)\left(1+y_{22}\right)-y_{12} y_{21}} & y_{11}=\frac{\left(1+s_{22}\right)\left(1-s_{11}\right)+s_{12} s_{21}}{\left(1+s_{11}\right)\left(1+s_{22}\right)-s_{12} s_{21}} \\
s_{12}=\frac{-2 y_{12}}{\left(1+y_{11}\right)\left(1+y_{22}\right)-y_{12} y_{21}} & y_{12}=\frac{-2 s_{12}}{\left(1+s_{11}\right)\left(1+s_{22}\right)-s_{12} s_{21}} \\
s_{21}=\frac{-2 y_{21}}{\left(1+y_{11}\right)\left(1+y_{22}\right)-y_{12} y_{21}} & y_{21}=\frac{-2 s_{21}}{\left(1+s_{11}\right)\left(1+s_{22}\right)-s_{12} s_{21}} \\
s_{22}=\frac{\left(1+y_{11}\right)\left(1-y_{22}\right)+y_{12} y_{21}}{\left(1+y_{11}\right)\left(1+y_{22}\right)-y_{12} y_{21}} & y_{22}=\frac{\left(1+s_{11}\right)\left(1-s_{22}\right)+s_{12} s_{21}}{\left(1+s_{22}\right)\left(1+s_{11}\right)-s_{12} s_{21}} \\
s_{11}=\frac{\left(h_{11}-1\right)\left(h_{22}+1\right)-h_{12} h_{21}}{\left(h_{11}+1\right)\left(h_{22}+1\right)-h_{12} h_{21}} & h_{11}=\frac{\left(1+s_{11}\right)\left(1+s_{22}\right)-s_{12} s_{21}}{\left(1-s_{11}\right)\left(1+s_{22}\right)+s_{12} s_{21}} \\
s_{12}=\frac{2 h_{12}}{\left(h_{11}+1\right)\left(h_{22}+1\right)-h_{12} h_{21}} & h_{12}=\frac{2 s_{12}}{\left(1-s_{11}\right)\left(1+s_{22}\right)+s_{12} s_{21}}
\end{array}
$$

$$
\begin{array}{ll}
S_{21}=\frac{-2 h_{21}}{\left(h_{11}+1\right)\left(h_{22}+1\right)-h_{12} h_{21}} & h_{21}=\frac{-2 s_{21}}{\left(1-s_{11}\right)\left(1+s_{22}\right)+s_{12} s_{21}} \\
S_{22}=\frac{\left(1+h_{11}\right)\left(1-h_{22}\right)+h_{12} h_{21}}{\left(h_{11}+1\right)\left(h_{22}+1\right)-h_{12} h_{21}} & h_{22}=\frac{\left(1-s_{22}\right)\left(1-s_{11}\right)-s_{12} s_{21}}{\left(1-s_{11}\right)\left(1+s_{22}\right)+s_{12} s_{21}}
\end{array}
$$

If $\mathrm{H}, \mathrm{Y}$, and $Z$ are the actual parameters, conversion is obtained as follows:
$z_{11}=z_{11} Z_{0}$
$Y_{11}=y_{11} \mu_{0}$

$$
H_{11}=h_{11} Z_{0}
$$

$$
z_{12}=z_{12} z_{0}
$$

$$
Y_{12}=y_{12} Z_{0}
$$

$$
\mathrm{H}_{12}=\mathrm{h}_{12}
$$

$$
z_{21}=z_{21} z_{0}
$$

$$
Y_{21}=y_{21} / Z_{0}
$$

$$
\mathrm{H}_{21}=\mathrm{h}_{21}
$$

$$
z_{22}=z_{22} z_{0}
$$

$$
Y_{22}=y_{22} / Z_{0}
$$

$$
H_{22}=h_{22} / Z_{0}
$$

## APPENDIX II

## APPROXIMATION OF HYBRID PI ELEMENTS

BY S -PARAMETER METHOD

1. $r_{b b}$.

From Smith Chart plot of $\mathrm{S}_{11} \mathrm{II}_{\mathrm{C}}$ curve at sufficiently high frequency (i.e. $\omega C_{b}{ }^{\prime} e^{r} b^{\prime} e^{\gg}$ ) the resistance value of the constant resistance circle is approximately $r_{b b}$. . (Refer to Fig. 5-3.)


Fig. A-II. Hybrid Pi equivalent circuit for a bipolar device.

Note: $S_{12^{S}} 21$ and $Y_{b}{ }^{\prime}$ c are assumed negligible in the derivations.
2. $Y_{b^{\prime} e}\left(=1 / r_{b}{ }^{\prime} e+j \omega C_{b^{\prime} e^{\prime}}\right)$

$$
\begin{equation*}
Y_{b^{\prime} e}=1 /\left(z_{i}-r_{b b^{\prime}}\right) \tag{II-1}
\end{equation*}
$$

where $Z_{i}$ is input impedance obtained from Smith Chart plot of $\mathrm{S}_{11}$.
3. $E_{m}$

$$
\begin{equation*}
\mathcal{E}_{\text {III }} / \theta=-\frac{1}{2} Y_{b} \cdot e^{S_{21}} /\left(1-S_{11}\right)\left(1+S_{22}\right) \tag{II-2}
\end{equation*}
$$

$\theta$ is the phase shift due to minority carrier transit time. 4. $Y_{b}{ }^{\prime} c\left(=1 / r_{b} c+j \omega C_{b}{ }^{\prime}\right)$

$$
Y_{b} \cdot c=\frac{1}{2} Y_{b} \cdot e_{12} /\left(1-S_{11}\right)\left(1+S_{22}\right)
$$

(I I-3)
5. $Y_{c e}\left(=1 / r_{c e}+j \omega C_{c e}\right)$

$$
\begin{equation*}
Y_{c e}=Y_{o}-Y_{b} c \tag{II-4}
\end{equation*}
$$

where $Y_{0}$ is output admittance obtained from Smith Chart plot of $\mathrm{S}_{22}$.

## APPENDIX III

## FURTHER DISCUSSIONS ON REFERENCE PLANE

A. The Effect of Improper Reference Plane


Channel A

$$
\begin{aligned}
& \text { Fig. A-III-1. } \text { Phase relation of incident and } \\
& \text { scattered waves at the measurement } \\
& \text { point. }
\end{aligned}
$$

If the ratio of the voltage waves $v_{B} / v_{A}$ gives a desired $S$-parameter, then the effect of remote measurements for the parameters will be as follows:

$$
s_{B A}=\frac{v_{B} e^{-j \beta l}}{v_{A} e^{-j \beta(l+\Delta l)}}=\left(\frac{v_{B}}{v_{A}}\right) e^{j \beta \Delta l}=S_{B A} e^{j \beta \Delta l}
$$

since

$$
\begin{equation*}
\beta=2 \pi / \lambda=2 \pi f / v \tag{III-1}
\end{equation*}
$$

where $\mathbf{v}$ is the velocity of a wave traveling along a transmission line and $f$ is the frequency.

$$
\begin{equation*}
S_{B A}=S_{B A} e^{j Ө f} \tag{III-2}
\end{equation*}
$$

where

$$
\theta=(2 \pi / v) \Delta l
$$

The measured $S$-parameter $S_{B A}{ }^{\prime}$ will experience a linear phase error which is a function of frequency.

For positive $\Delta \ell$ (longer reference signal path) the measured $S$-parameter $S_{B A}$ " will have a negative phase error corresponding to $\theta f$.

$\Delta l<0$
$\Delta l>0$

Fig. A-III-2. Phase error due to improper reference plane.
B. Fringing Effect Error - Short and Open Position Difference

This is the error due to the end effect of a transmission line. Theoretically if $\Delta l=0$

$$
\begin{aligned}
S_{B A}=S_{B A} & =1 \angle+180^{\circ} & & \text { for short end, } \\
& =1 \angle 0^{\circ} & & \text { for open end. (III-3) }
\end{aligned}
$$

But in actual case the open and shoit termination positions differ from each other by a small angle.

To study this relationship a measurement is performed as follows:


Fig. A-III-3. Measurement of fringing effect error.
The adjustable line is adjusted for a short termination such that VVM shows $180^{\circ}$ over a wide range of frequencies. Then the short is renoved and the phase deviation from $0^{\circ}$ is measured. The general curves are as follows


Fig. A-III-4. An example of fringing effect.

So if reference plane is set at a point using the short termination method, the high impedance (open condition) termination will show significant position errors depicted in the figure. For high impedance measurements the angular error should be subtracted from the measured values.

## APPENDIX IV

COMMON-GATE S-PARAMETERS OF SINGLE-GATE MOSFET (S-V ${ }_{\text {DG }}$ CURVES)


Fig. A-IV-1. Common-gate $S_{21}-V_{D G}$ curve (Single-gate MOSFET)


Fig. A-IV-2. Common-gate $\mathrm{S}_{22}-\mathrm{V}_{\mathrm{DG}}$ curve
(Single-gate MOSFiT)


Fig. A-IV-3. Common-gate $S_{21}-V_{D G}$ curve (Single-gate MOSFi®T).


Fig. A-IV-4. Common-gate $\mathrm{S}_{12}-\mathrm{V}_{\mathrm{DG}}$ curve (Single-gate MOSFET)
APPENDIX VFREQUENCY RESPONSE OF DUAL-GATE MOSFETCOMMON-SOURCE S-PARAMETERS BEYOND VHF


Fig. A-V-1. Frequency response of dual-gate MOSFET common-source S-parameters beyond VHF.


Fig. A-V-2. Frequency response of dual-gate MOSFET common-source S-parameters beyond VHF


Fig. A-V-3. Smith Chart plot of frequency response for dual-gate NOSFEI common-source S-parameters beyond VHF: (a) $\mathrm{S}_{11}$ and (b) $\mathrm{S}_{22}$ versus freçuency curves.

## APPENDIX VI

## MAXIMUM UNILATERAL GAIN AND UNILATERAL FIGURE OF MERIT VERSUS FREQUENCY CURVES



Fig. A-VI. Maximum unilateral gain $G_{U}$ and unilateral figure of merit $U$ versus frequency curves of dual-gate MOSFET common-source stage beyond VHF .
APPENDIX VII
COMMON-SOURCE S-PARAMETERS VERSUS GATE 2 VOLTAGECURVES OF DUAL-GATE MOSFET


Fig. A-VII-1. $\begin{gathered}\text { Common-source } \mathrm{S}_{11}-\mathrm{V}_{\mathrm{G} 2} \text { curves } . \\ \text { (Dual-gate MOSFET) }\end{gathered}$


Fig. A-VII-2. Common-source $S_{22}-V_{G 2}$ curves (Dual-gate MOSFE'T)


Fig. A-VII-3. Common-source $\mathrm{S}_{21}-\mathrm{V}_{\mathrm{G} 2}$ curves (Dual-gate MOSFET)



Fig. A-VII-5. Common-source $\mathrm{S}_{11}-\mathrm{V}_{\mathrm{G} 2}$ curves (Dual-gate MOSFE'T).


Fig. A-VII-6. Common-source $\mathrm{S}_{22}-\mathrm{V}_{\text {G2 }}$ curves (Dual-gate MOSFCT).


Fig. A-VII-7. Common-source $S_{21}-V_{G 2}$ curves (Dual-gate MOSFET)


Fig. A-VII-8. Common-source $\mathrm{S}_{12}-\mathrm{V}_{\mathrm{G} 2}$ curves (Dual-gate MOSFET).

