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Wafer-Scale Fabrication of Sub-10 nm TiO₂-Ga₂O₃ *n-p* Heterojunctions with Efficient Photocatalytic Activity by Atomic Layer Deposition

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Abstract

Wafer-scale, conformal, two-dimensional (2D) TiO₂-Ga₂O₃ *n-p* heterostructures with a thickness of less than 10 nm were fabricated on the Si/SiO₂ substrates by the atomic layer deposition (ALD) technique for the first time with subsequent post-deposition annealing at a temperature of 250 °C. The best deposition parameters were established. The structure and morphology of 2D TiO₂-Ga₂O₃ *n-p* heterostructures were characterized by the scanning electron microscopy (SEM), X-ray photoelectron spectroscopy (XPS), electrochemical impedance spectroscopy (EIS), etc. 2D TiO₂-Ga₂O₃ *n-p* heterostructures demonstrated efficient photocatalytic activity towards methyl orange (MO) degradation at the UV light ($\lambda = 254$ nm) irradiation. The improvement of TiO₂-Ga₂O₃ *n-p* heterostructure capabilities is due to the development of the defects on Ga₂O₃-TiO₂ interface, which were able to trap electrons faster.

Keywords: TiO₂-Ga₂O₃, *n-p* heterostructures, atomic layer deposition, 2D semiconductors

Background

Fabrication of 2D *p-n* heterojunctions of semiconductor oxides is one of the key directions of future development of nanostructures with unique distinguishable properties, as they are able to combine various outstanding features of both semiconductors at the nanoscale [1–4]. However, it is extremely challenging to fabricate them defects-free over the wafer area, particularly when the thickness of each oxide is only few nanometers [2]. In order to overcome numerous manufacturing challenges, ALD technology has already established clear and unprecedented advantages in the development of conformal nano- and monolayers of the semiconductor oxides and their 2D heterostructures with the thickness less than 10 nm on wafer-scale with high aspect ratio [3, 5–7]. In addition, various new approaches were also initiated

recently for the development of 2D heterostructures with enhanced functional capabilities [8–11]. They specifically targeted both oxygen evolution reaction (OER) and hydrogen evolution reaction (HER), as a core processes for various renewable energy systems [12]. However, in comparison to HER, OER with multistep, four-electron process evolved is severely constrained by its sluggish kinetics [13]. Thus, more efforts have therefore been devoted to improve the conductivity of heterostructures and control the electronic structures of their surface active sites through the modulation of their morphology, constituent compositions, and/or dopants [8, 12]. Moreover, regulating the surface-adsorbed species may also provide an alternative valuable approach to fine-tuning the interfacial properties, particular at nanostructured heterojunctions, and the electronic structures of active materials [14].

More importantly, it was demonstrated that the decreasing free energy of the OER intermediates at the nano-interface would remarkably enhance the inherent electrochemical performance of catalyst [13]. In this regard, surface engineering is well illustrated to improve

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the accessibility of the reactants and to alter the electrochemical activity of the catalysts [14]. To achieve such enhancements of electrochemical properties of nanostructured heterostructures, various technological approaches have been utilized. Among them, the ALD technique can be used to deposit wafer-scaled nanomaterials with controlling their deposition rate at the Angstrom scale. Additional vital advantage of ALD is its self-limited nature by depositing materials in an atomic layer-to-layer [5, 6].

The alternative approach represents the development of 2D C-MOFs via the combination of “through-space” and “through-bind” strategies [14]. In particular, hexahydroxytriphenylene ligand-based 2D C-MOFs possess M-O4 (M—transitional metals) as their secondary building units and provide discrete metal-replicable layers as promising reactive sites for OER [14]. Moreover, these C-MOFs can remain stable in high pH solution, which is quite important for OER. Thus, all these above-mentioned recent advancements indirectly confirmed that no other technologies of making 2D nanostructures, including *sol-gel*, chemical vapor deposition (CVD), RF sputtering, etc., are capable to deliver uniformed deposition at the Ångstrom level over the large areas of Si/SiO₂ wafers with precise control of the deposition rate and thickness. Therefore, most of the developed recipes for ALD of 2D nanostructures using specific precursors possess valuable *know-how* and represent a highly repeatable process on the semi-industrial scale [5, 15, 16].

One of the main 2D semiconductors successfully utilized in the different photovoltaic applications is titanium dioxide (TiO₂), which is a typical *n*-type semiconductor with wide bandgap $E_g = \sim 3.2$ eV [5, 15–18]. There are numerous scientific reports focused on the different approaches for improvement of its properties such as changing thickness of nanostructured 2D TiO₂ down to monolayer [15, 16], doping TiO₂ by other nanostructures semiconductors [5, 17], surface functionalization of 2D TiO₂ [18] and making *n-p* heterojunctions [19]. In addition, low electron/hole recombination is blamed for the low quantum yields, which is still a big obstacle for the improvement of photocatalytic activity. Therefore, fabrication of efficient *n-p* heterojunctions has been proposed and attempted with the different levels of success during the last few years [4, 17, 20–22]. Specifically, it was found that the fabricated *n-p* heterojunctions could sufficiently reduce the recombination rate of the photo-generated electron/hole pairs with the following enhancement of the overall photocatalytic activity [1, 23, 24]. Thus, the combination of *p*- and *n*-type semiconductor oxides has paved the way for the further development of *n-p* heterojunctions and optimization of their photocatalytic capabilities [25].

In this regard, 2D surface functionalization of 2D *n*-type TiO₂ by ALD of another *p*-type semiconductor

on the top of TiO₂ represents a unique strategy of making *n-p* heterojunctions and combining various outstanding properties of both semiconductors [5]. On the other hand, semiconductor oxides with a d¹⁰ electron configuration have recently attracted considerable attention for their superior activities as potential dopant. This is mainly owing to their conduction bands being formed by hybridized *sp* orbits with a large dispersion, which enabled them to generate electrons with the large mobility [26]. Gallium oxide (Ga₂O₃), as a typical representative of such d¹⁰ semiconductor oxides, belongs to the group of transparent semiconducting oxides with a wide band gap and electrical conductivity. It exhibits the largest band gap with $E_g = 4.8$ eV and thus a unique transparency from the visible into the UV region and good luminescence properties [27]. β -Ga₂O₃ is reported to be the most stable polymorph among five existing polymorphs of Ga₂O₃ within the high-temperature range [28]. Moreover, nontoxic β -Ga₂O₃ displayed significant potential for photocatalytic air purification, particularly for the elimination of toxic aromatic compounds [29]. Therefore, all these distinguishable properties of β -Ga₂O₃ [30] substantiated a lot of efforts for the best suitable technologies of Ga₂O₃ deposition at the nanoscale [31–33].

Notwithstanding the great attempts dedicated to the ALD of 2D semiconductor oxides during last few years, authors wish to stress that so far 2D TiO₂-Ga₂O₃ *n-p* heterostructures with the thickness less than 10 nm have not yet been reported. In this work, 2D TiO₂-Ga₂O₃ *n-p* heterostructures were ALD-fabricated on wafer-scale for the first time using Ti(N(CH₃)₂)₄ and C₃₃H₅₇GaO₆ as TiO₂ and Ga₂O₃ precursors, respectively. Their optimal deposition parameters were established and structural and photocatalytic properties were investigated.

Results and Discussion

Figure 1 illustrates the fabrication process of 2D TiO₂-Ga₂O₃ *n-p* heterostructures on the Si/SiO₂ substrate. Figure 2 schematically depicts the details of ALD depositions. After depositions, wafers were diced on 1.0 × 1.0 cm segments (Fig. 2a) for further testing. For ALD 2D TiO₂-Ga₂O₃ *n-p* heterostructures Ti(N(CH₃)₂)₄ and C₃₃H₅₇GaO₆ (Strem Chemicals Inc., USA) were used as TiO₂ and Ga₂O₃ precursors, respectively. Their graphical interpretations are given in Fig. 2b, c. The growth per cycle (GPC) yielded from the slopes of growth curves shown in Fig. 2d, e was calculated to be around 0.7 Å/cycle and 0.16 Å/cycle for TiO₂ and Ga₂O₃, respectively. The growth curves were linear without any nucleation delay for both TiO₂ and Ga₂O₃ samples, implying that the self-limited property of ALD growth process and the film thickness could be developed precisely by varying the number of ALD cycles. The lower growth rate of 2D Ga₂O₃ nano-films makes its

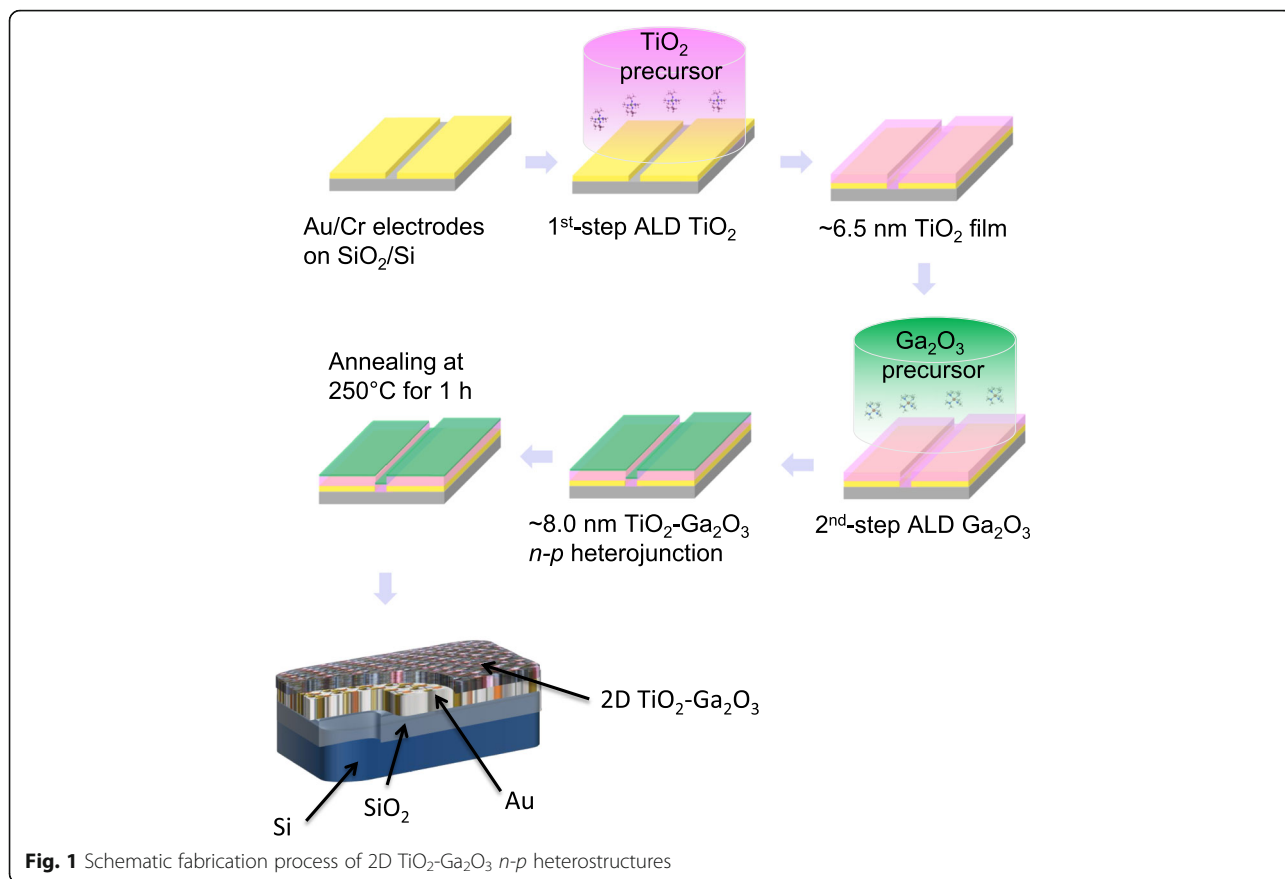


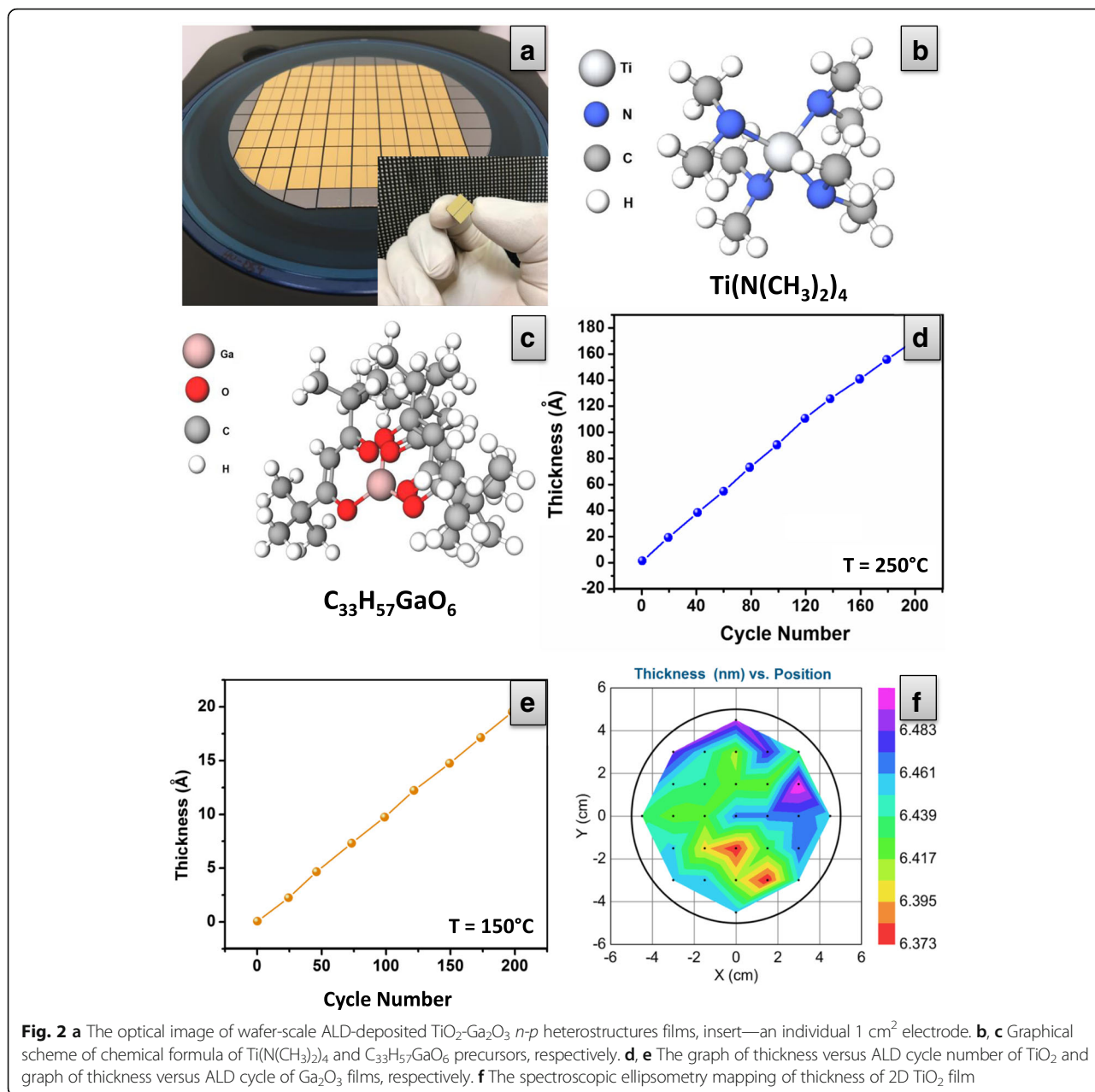
Fig. 1 Schematic fabrication process of 2D TiO₂-Ga₂O₃ n-p heterostructures

applications on the doping and modification possible [34]. Noteworthy, the optimal ALD deposition parameters for each precursor are usually established after several initial trials [6]. After each deposition cycle the variable angle in situ ellipsometry measurements (J.A. Woollam M2000 DI) were carried to monitor the uniformity and to measure the thickness of films. For example, Fig. 2f illustrates the in situ ellipsometry measurements for 2D TiO₂ with the average thickness of ~ 6.45 nm. Since the thickness measurements were found difficult on heterostructure, the Ga₂O₃ film growth was followed, using in situ ellipsometry measurement, on SiO₂/Si substrate that was placed on the heater block, together with the sample. After the deposition, the Ga₂O₃ film thickness on heterostructure was confirmed by comparing the amount of material deposited on it and the reference SiO₂/Si using X-ray fluorescence measurements [19]. 2D Ga₂O₃ films had an average thickness of ~ 1.5 nm, which resulted in the total thickness of 2D TiO₂-Ga₂O₃ heterostructures to be ~ 8.0 nm. All fabricated samples were subsequently annealed in the air for 1 h at 250 °C with a heating rate of 0.5 °C/min.

Figure 3 shows SEM surface morphology images for both ALD-fabricated 2D TiO₂ (thickness ~ 6.5 nm) and

Ga₂O₃ (thickness ~ 1.5 nm) nano-films. It is noteworthy that the TiO₂ nano-grains in the fabricated films were uniformly distributed over Si/SiO₂ wafer and varied in size from approximately ~ 30 to ~ 70 nm prior to Ga₂O₃ deposition. Figure 3a depicts surface morphology of TiO₂ nano-film consisting of the flat nano-particles. Then, the ALD-developed ~ 1.5-nm-thick Ga₂O₃ nano-films were fabricated on the top of ~ 6.5-nm-thick TiO₂ nano-films. The ALD-developed sub-10 nm Ga₂O₃-TiO₂ heterostructures were subsequently annealed at 250 °C. Thus, Fig. 3b depicts crystalline surface morphology of the Ga₂O₃ in heterostructure after annealing. The Ga₂O₃ nano-film consists of uniformly distributed Ga₂O₃ nano-grains with the average size from ~ 80 to ~ 110 nm. Owing to the extremely thin nature of the ALD-fabricated nano-films, employment of the X-ray diffraction technique for crystallinity investigation of these films was not possible.

Chemical composition and bonding states of 2D TiO₂-Ga₂O₃ heterostructures were studied by XPS with Fig. 4a representing the TiO₂-Ga₂O₃ heterostructure scan survey. The charge shift spectrum was calibrated for C1s peak at 284.8 eV. Three main elements of Ti, O, and Ga are clearly observed. In addition, C1s peak was also detected as it was originated from the reference to calibrate the binding energies of the peaks. Figure 4b



depicts high-resolution two quasi-symmetrical Ga 2p_{1/2} and Ga 2p_{3/2} peaks for Ga-O bonding at 1145.2 eV and 1118.4 eV with a separation distance of 26.8 eV, which is consistent with the binding energy of Ga 2p for doped β-Ga₂O₃ [35, 36]. The weak energy peak for Ga 3d is centered at 21.1 eV, which is caused by the presence of Ga-O bond reported for *p*-type β-Ga₂O₃ films [37], but not observed for the *n*-type β-Ga₂O₃ structures [38]. The Ga 3d peak is asymmetrical, which was ascribed to the hybridization of Ga 3d and O 2s states near the valence band [39]. Figure 4c displayed the high-resolution scan of Ti 2p. The doublet peaks demonstrated in Fig. 4c correspond to Ti 2p_{3/2} and Ti 2p_{1/2} with the spin-orbital

splitting of 6.2 eV, which were attributed to Ti⁺⁴ oxidation state. It should be noted that the obtained XPS results in this investigation are slightly different from our previous report on the development of TiO₂ monolayer [15] and bi-layer [3] grown by ALD. This difference is reasonable considering the amount of Ti in the samples.

The O 1s peak in the XPS spectrum (Fig. 4d) could be deconvoluted into two major peaks. The main binding energy component centered at 531.53 eV is attributed to oxygen vacancies or OH⁻¹ adsorbed species on the surface [38]. The second binding energy peak at 530.01 eV can be the characteristic of the lattice oxygen in the TiO₂-Ga₂O₃ heterostructure. Very relevant to this investigation was

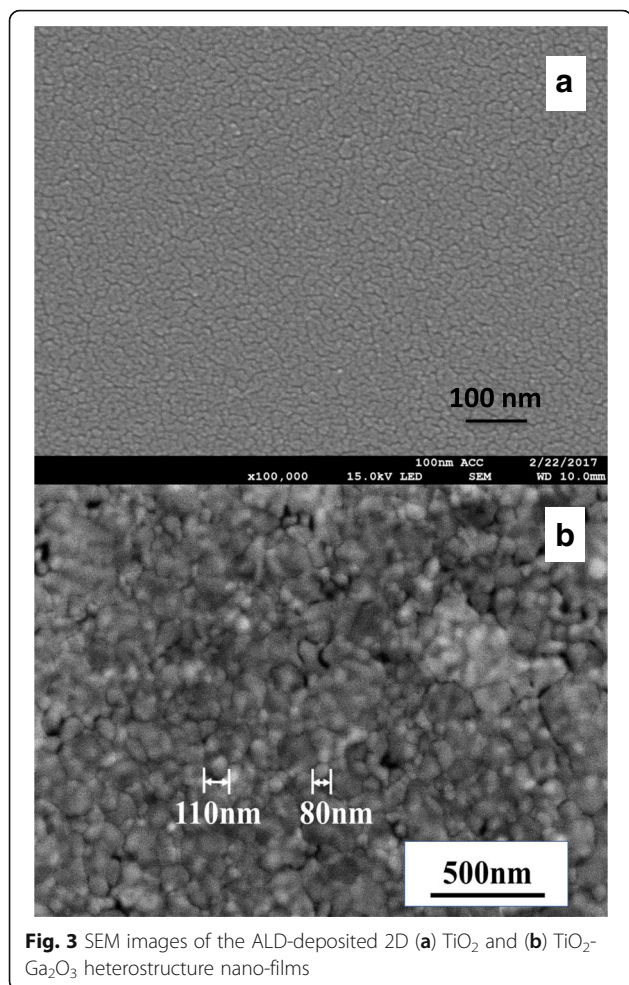


Fig. 3 SEM images of the ALD-deposited 2D (a) TiO₂ and (b) TiO₂-Ga₂O₃ heterostructure nano-films

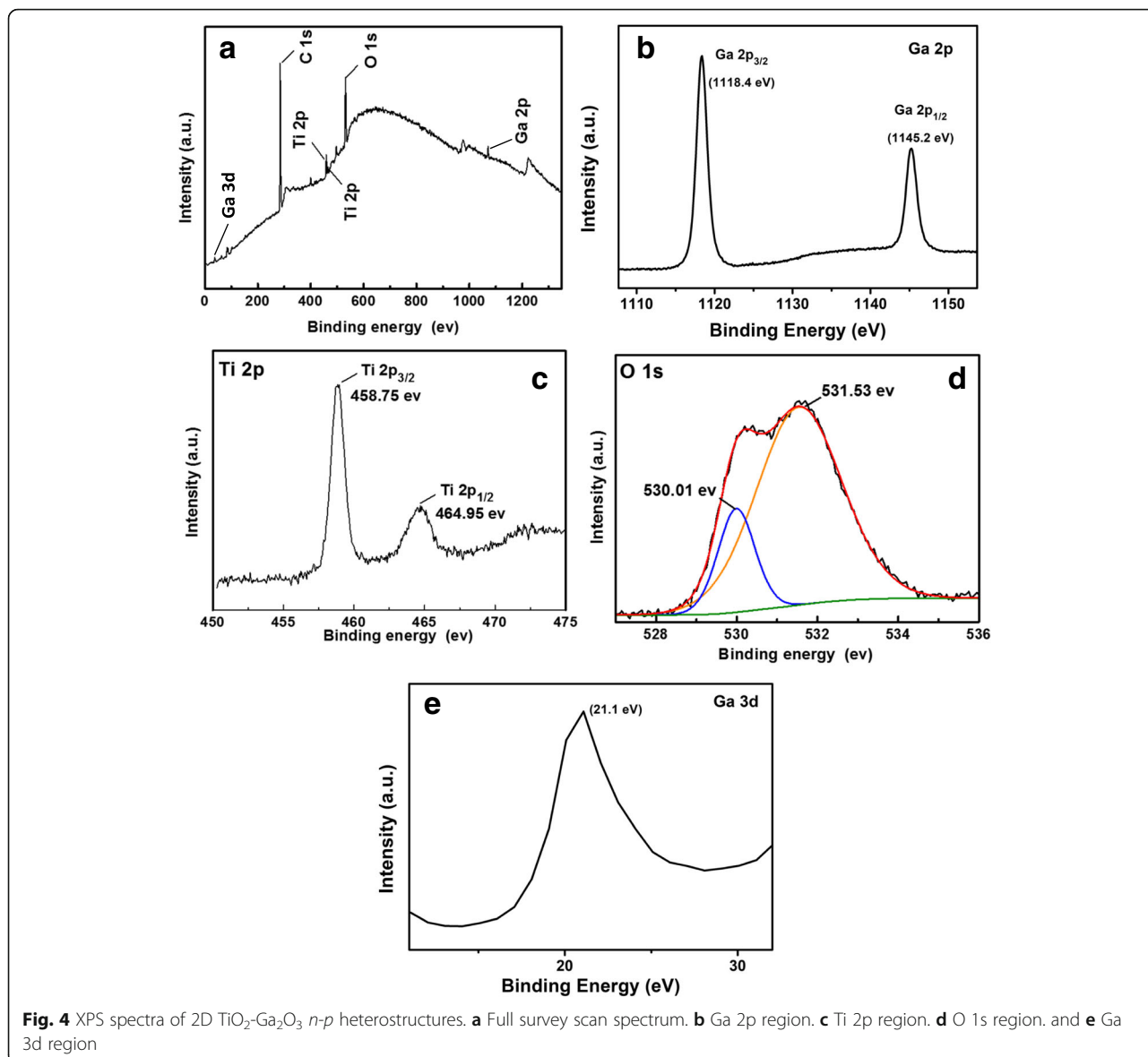
our previous study on ALD TiO₂ bi-layer confirming the influence of SiO₂ substrate, where the bottom oxygen of TiO₂ is shared with SiO₂ making 2D TiO₂ slightly non-stoichiometric [3]. Thus, this non-stoichiometry plays a critical role in 2D TiO₂-Ga₂O₃ heterostructure while the thickness of Ga₂O₃ ALD on the top of TiO₂ is only ~ 1.5 nm. The enlarged energy peak for Ga 3d is presented in Fig 4e. Presence of Ga 3d peak in the spectrum is confirmation of the *p*-type conductivity for Ga₂O₃ in the heterostructure, as being reported [37]. For further investigation of the conductivity type of 2D β-Ga₂O₃, additional 4.8-nm-thick Ga₂O₃ samples were subjected to the Hall coefficient measurements at *T* = 25 °C. The measured Hall coefficient value of 8.292×10^4 cm³/C independently confirmed the stable *p*-type performance of 2D Ga₂O₃.

Figure 5 expresses the plotted EIM measurements of the spectra for 2D TiO₂ (~ 3.5 nm), Ga₂O₃ (~ 3.5 nm), and 2D TiO₂-Ga₂O₃ heterostructures (~ 8.0 nm), respectively. EIS measurements were carried out in air at the temperature of 25 °C and the impedance results were obtained using the Randles equivalent circuit. It is noteworthy that the fitted Nyquist plots in Fig. 5 revealed the

charge-transfer resistance ($R_{ct} = 4.5$ kΩ) of 2D TiO₂-Ga₂O₃ heterostructures with a thickness of ~ 8.0 nm being about 2.7-fold lower than that of ALD-developed 2D TiO₂ ($R_{ct} = \sim 12.5$ kΩ) and even slightly lower than that of 2D Ga₂O₃ ($R_{ct} = \sim 6.0$ kΩ). This fact further designates that 2D TiO₂-Ga₂O₃ heterostructures possess a much faster charge-transfer characteristics than that of 2D TiO₂ and Ga₂O₃. Although the measured impedance value for of Ga₂O₃ was slightly higher than the reported value for 2D ALD-fabricated Ga₂O₃ [40], this was partially due to the sub-nanometer thickness of the Ga₂O₃ film [40] compare to the ~ 3.5-nm-thick Ga₂O₃ in our experiments and was also partially owing to the fact that the developed 2D Ga₂O₃ was not fully crystallized at the annealing temperature of 250 °C.

All FTIR spectra of 2D Ga₂O₃, TiO₂ and TiO₂-Ga₂O₃ *n-p* heterostructures are summarized in Fig. 6. As spectra for 2D TiO₂ and Ga₂O₃ are nearly overlapping each other, they therefore were presented separately in Fig 6a and Fig. 6b, respectively, in comparison with the spectrum of 2D TiO₂-Ga₂O₃ heterostructures. The peaks centered at about 1594 cm⁻¹ are attributed to the O-H stretching and bending modes of the hydrated oxide surface and the adsorbed water [41]. Moreover, the adsorption of atmospheric CO₂ on the surface of gallium oxide is characterized by the detection of bands at 1519 cm⁻¹ and 1646 cm⁻¹, which resulted from preparation and processing of the samples in ambient air [42]. More interesting results were observed in the perturbation area, presented as inserts in Fig. 7a and Fig. 7b, respectively. The IR band at 607.9 cm⁻¹ is due to vibration of the Ga-O bond of GaO₆ octahedra in Ga₂O₃ lattice [43]. Its intensity has the maximum in FTIR spectrum of 2D Ga₂O₃ and decreased in the FTIR spectrum of 2D TiO₂-Ga₂O₃ *n-p* heterostructure. Compared with the FTIR spectrum of 2D Ga₂O₃ nano-film, a new peak at 464 cm⁻¹ appeared in FTIR spectrum for ALD-fabricated 2D TiO₂-Ga₂O₃ *n-p* heterostructures. This peak is near overlapping typical characteristic peak at 470 cm⁻¹ for TiO₂ [15].

Photoluminescence (PL) technique is usually employed to investigate the migration, transfer and recombination rate of the photo-induced electrons-holes pairs in semiconductors. Figure 7 shows the room temperature (25 °C) PL spectra of ALD-fabricated 2D TiO₂-Ga₂O₃ *n-p* heterostructures annealed at 250 °C with the details of the measured bandgap for TiO₂ and Ga₂O₃, respectively. There are two peaks in the PL spectra for the 2D TiO₂-Ga₂O₃ *n-p* heterostructures (presented in insert in Fig. 7): one is called near band edge emission (NBE), which is in the UV region due to the recombination of free excitons through an exciton-exciton collision process; and the second one is called deep level emission (DPE), which is caused by the impurities and/or

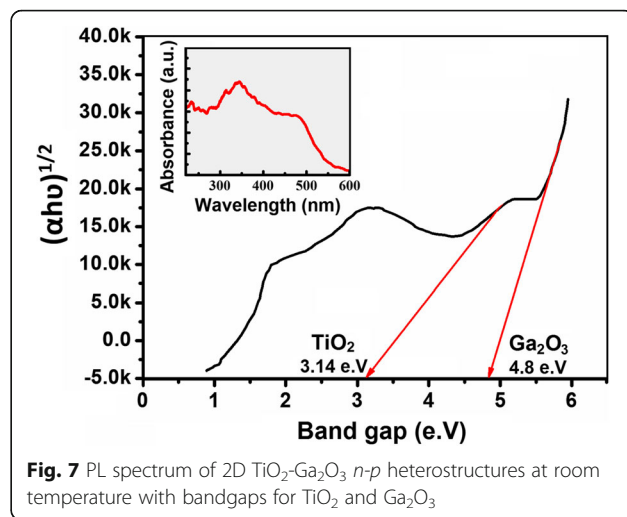
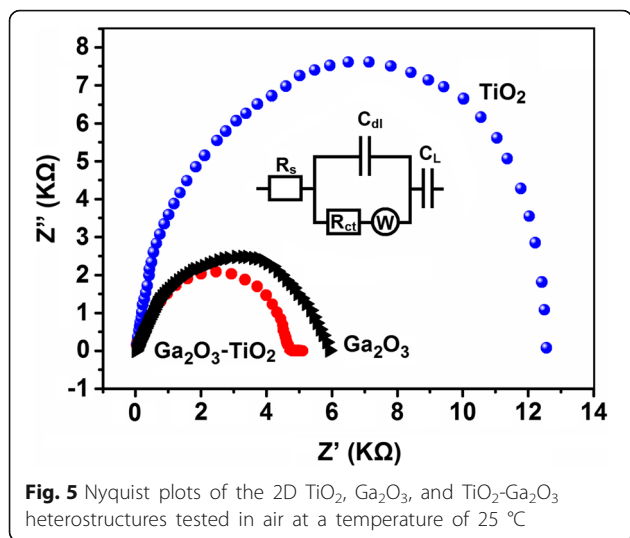


structural defects in the crystal [41]. The DPE intensity in 2D TiO₂-Ga₂O₃ *n-p* heterostructures is lower than that in Ga₂O₃ [44], which indicates more efficient transfer and separation of the charge carriers owing to the electron-hole transfer in the heterojunctions between TiO₂ and Ga₂O₃. Noteworthy, the DPE of 2D TiO₂-Ga₂O₃ *n-p* heterostructures is shifted towards the UV region whereas DPE of Ga₂O₃ is within the visible light region [44]. In addition, the selected annealing temperature of 250 °C did not allow full crystallization of Ga₂O₃ nano-film in the heterostructure, which was reflected by the unchanged value of its bandgap (4.8 eV). However, in our previous investigation, it was found that further increase of the annealing temperature (above 250 °C) of such extremely-thin films causes their disintegration with the following

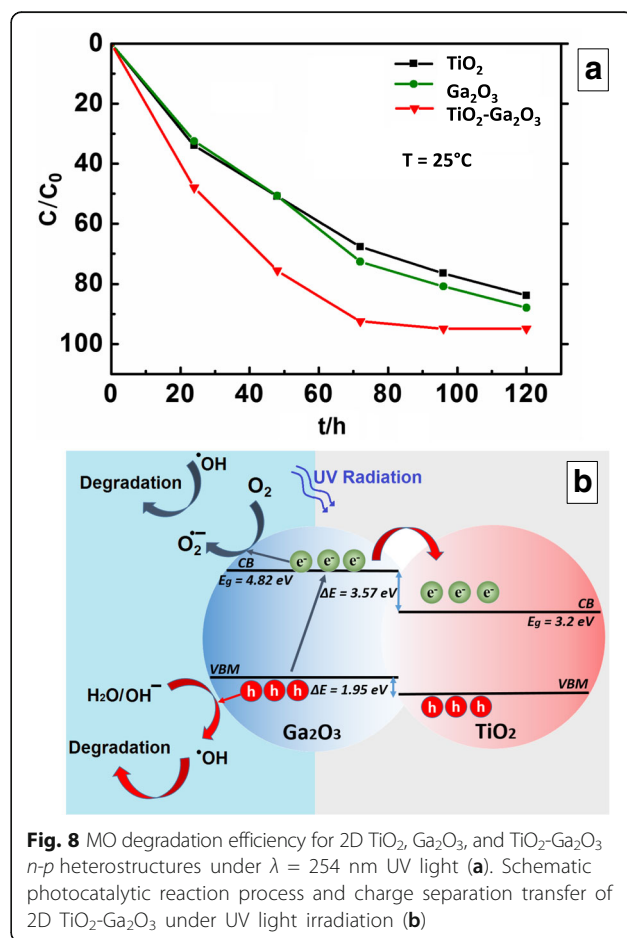
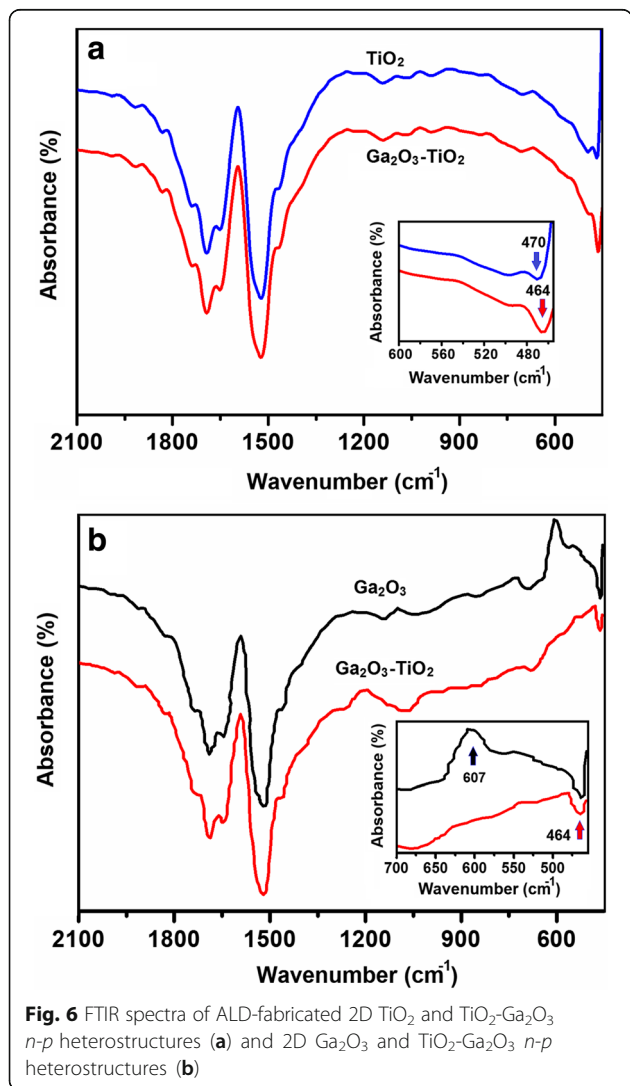
agglomeration of their nano-grains into island-like nanostructure [6]. On the contrary, the bandgap for TiO₂ slightly changed to ~3.14 eV compared to its microstructural counterpart.

Consequently, all the above material characterization experiments clearly confirmed the successful development of conformal and uniform sub-10 nm TiO₂-Ga₂O₃ *n-p* heterostructures. Thus, these 2D TiO₂-Ga₂O₃ *n-p* heterostructures were ALD-fabricated impurity-free on the wafer-scale and subsequently annealed at 250 °C for the establishment of developed *n-p* nano-interface.

The photocatalytic degradation of MO under the UV light irradiation ($\lambda = 254\text{nm}$) was carried out at the room temperature (25 °C) to evaluate the photocatalytic activity of ALD-fabricated 2D TiO₂, Ga₂O₃ and 2D TiO₂-Ga₂O₃ *n-p* heterostructures. As presented in Fig. 8a, 2D



TiO₂-Ga₂O₃ *n-p* heterostructure demonstrated higher photocatalytic activity compared to both 2D TiO₂ and Ga₂O₃ under the same UV irradiation. Specifically, using 2D TiO₂-Ga₂O₃ *n-p* heterostructure as the catalyst, MO degradation efficiency reached ~90% within 70 h, while the values for 2D Ga₂O₃ and TiO₂ were



approximately ~70% and ~65%, respectively, at the same time. Considering the fact that 2D Ga₂O₃ has not been fully crystallized under the annealing temperature of 250 °C, it is assumed that the weak chemical bond developed between 2D TiO₂ and Ga₂O₃ is good enough to ensure the successful role of *n-p* heterojunction for the photocatalytic activity.

The photocatalytic degradation mechanism by 2D TiO₂-Ga₂O₃ *n-p* heterostructure under $\lambda = 245$ nm UV light irradiation is proposed in Fig. 8b. It is a common knowledge that the photocatalytic degradation of dyes mainly involves several active radical species such as hydroxyl radicals ($\cdot\text{OH}$), holes (h^+) and electrons (e^-) [45]. The direct contact between 2D Ga₂O₃ and TiO₂ induced the development of heterojunction owing to the different energy levels. Under $\lambda = 254$ nm UV light irradiation, both Ga₂O₃ and TiO₂ were excited to generate electrons and holes simultaneously. Large numbers of defects consisting of robust acceptor state in the band-gap trap holes and prevent recombination. Various defect bands promote the electron-hole pair separation rate. The enhanced photo-catalytic performance is mainly derived from the large numbers of acceptor states accompany with Ga₂O₃ defects especially in its not fully crystallized phase. The acceptor states not only expand the light absorption edge of UV but also retard the rate of electron-hole pair recombination. In this regard, both large number of defects and acceptor states is responsible for enhancing the photocatalytic performance of 2D TiO₂-Ga₂O₃ *n-p* heterostructure. At the same time, holes in the VB of TiO₂ can migrate into the VB of Ga₂O₃. Thus, the concentration of photo-generated holes on the Ga₂O₃ surface increases. The photo-generated holes play a vital role in the photo-degradation process of 2D TiO₂-Ga₂O₃ *n-p* heterostructures. Therefore, the increasing concentration of the photo-generated holes in the VB of Ga₂O₃ could also lead to its high photocatalytic activity. Moreover, the higher-specific surface area fabricated after annealing may additionally improve the overall photocatalytic activity of 2D TiO₂-Ga₂O₃ *n-p* heterostructures. The absorption and desorption of molecules on the surface of the catalyst is the first step in the degradation process [46, 47]. Thus, higher surface-to-volume ratio in the surface morphology of the TiO₂-Ga₂O₃ *n-p* heterostructures provides more unsaturated surface coordination sites. Therefore, the annealed 2D TiO₂-Ga₂O₃ *n-p* heterostructures possess higher-specific surface area caused by numerous ultrathin nano-grains, as presented in SEM characterization. Consequently, high surface-to-volume ratio combined with the suitable nano-interfaces obtained for the 2D TiO₂-Ga₂O₃ *n-p* heterostructures resulted in its great photocatalytic activity towards the efficient MO degradation.

Conclusions

In this work, wafer-scale 2D TiO₂-Ga₂O₃ *n-p* heterostructures with the average thickness of ~8.0 nm were successfully fabricated for the first time via a two-step ALD process by using Ti(N(CH₃)₂)₄ and C₃₃H₅₇GaO₆ as TiO₂ and Ga₂O₃ precursors, respectively. Their optimal deposition parameters were established. The 2D TiO₂-Ga₂O₃ *n-p* heterostructures were annealed at 250 °C for the structural stabilization and development of the *n-p* nano-interface. Subsequently, 2D TiO₂-Ga₂O₃ *n-p* heterostructures were utilized for efficient MO degradation at the room temperature under the UV light ($\lambda = 254$ nm) irradiation. 2D TiO₂-Ga₂O₃ *n-p* heterostructures have clearly demonstrated unique capabilities and higher photocatalytic activity than that of pure 2D TiO₂ and Ga₂O₃ for MO degradation. Specifically, the effect of *n-p* heterojunction between *n*-type TiO₂ and *p*-type Ga₂O₃ enabled a higher concentration of the photo-generated holes and larger-specific surface area, which ultimately led to its higher photocatalytic activity. Therefore, sub-10 nm, 2D *n-p* heterostructures can be potentially exploited as promising nano-materials for the practical photocatalytic devices.

Methods

Synthesis 2D *n-p* Heterostructure

All reagents and precursors were purchased from the commercial sources and represented analytical grade. They were used as received without further purification. The 4-in. Si/SiO₂ wafers (12 Ω /cm) were utilized as substrates for ALD depositions, where the thickness of the native oxide was ~1.78–1.9 nm. 2D TiO₂-Ga₂O₃ *n-p* heterostructures were prepared by a two-step fabrication method. Prior to ALD depositions, in order to reduce the influence of Si wafer on electrical measurements, an additional ~100-nm-thick SiO₂ insulating layer was applied by CVD, (Oxford Instruments Plasmalab 100). After that 150-nm-thick Au/Cr films were deposited on SiO₂/Si by the Electron Beam Evaporator method (Nanochrome II (Intivac, USA)) to develop electrodes for subsequent investigations. All ALD fabrications were carried out on Savannah S100 (Ultratech/Cambridge Nanotech). A pulse time of 5 s was used for both the Ga(TMHD)₃ and O₂ plasma, at a pressure of 3×10^{-3} mbar.

Characterization

The surface morphology and elemental analysis of ALD-fabricated sub-10 nm TiO₂-Ga₂O₃ heterostructures were characterized by scanning electron microscopy (SEM, SU-500) and energy dispersive X-ray (EDX) spectroscopy (EDS, JEOL). Fourier transform infrared (FTIR) spectra were taken using a NEXUS Thermo Nicolet IR-spectrometer in the range 4000–400 cm⁻¹ with a

spectral resolution 2 cm^{-1} . In order to investigate the surface chemistries of the developed samples, X-ray photoelectron spectroscopy (XPS) was employed in the ESCALAB system with AlK X-ray radiation at 15 kV. All XPS spectra were accurately calibrated by the C1s peak at 284.6 eV for the compensation of the charge effect. Hall effect measurement system (HMS3000) was employed at the room temperature to measure the Hall coefficient of Ga_2O_3 thin films by using a 0.55T magnet. EIS and all electrical measurements for 2D TiO_2 , Ga_2O_3 , and TiO_2 - Ga_2O_3 heterostructures were carried out on AutoLab PGSTAT204 (Metrohm Autolab, B.V., Netherlands). Room temperature photoluminescence (PL) spectra of ALD 2D TiO_2 - Ga_2O_3 heterostructures were performed on an F-4600 fluorescent spectrophotometer (Hitachi Corp., Tokyo, Japan), and the maximal excitation wavelength was $\lambda = 200\text{ nm}$, and the filter was $\lambda = 300\text{ nm}$. The photocatalytic activity of 2D TiO_2 , Ga_2O_3 and 2D TiO_2 - Ga_2O_3 heterostructures for the MO ($\text{C}_{14}\text{H}_{14}\text{N}_3\text{NaO}_3\text{S}$) degradation in aqueous solution under the UV light was evaluated by measuring the absorbance of the irradiated solution. For this study, 2D TiO_2 - Ga_2O_3 heterostructures were placed into 100 mL of MO solutions with a concentration of 6 mg/L and a pH of 6.5. The solutions were continuously stirred in the dark for 2 h before illumination in order to reach the absorption-desorption equilibrium between MO and the 2D TiO_2 - Ga_2O_3 heterostructures. Then the solutions were irradiated by a 30 W low-pressure UV lamp ($\lambda = 254\text{ nm}$), which was located at the distance of 50 cm above the top of the dye solution. During the process, 5 mL solutions were pipetted every 12 h for the absorbance determination by a UNIC UV-2800A spectrophotometer using the maximum absorbance at 465 nm. All experiments were performed under the ambient condition and room temperature. The degradation efficiency of MO was defined as

$$D = [(A_0 - A_t) / A_0] \times 100\%, \quad (1)$$

where D is degradation efficiency, A_0 is the initial absorbance of MO solution, and A_t is the absorbance of MO solution after UV irradiation within the elapsed time t .

Abbreviations

EDS: Energy dispersive spectroscopy; FTIR: Fourier transform infrared; MO: Methyl orange; PL: Photoluminescence; SEM: Scanning electron microscopy; UV-vis: Ultraviolet-visible; XPS: X-ray photoelectron spectroscopy

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Availability of Data and Materials

The crystal structure and chemical bonding structure of the as-prepared samples were characterized by XPS (Fig. 4), electrochemical impedance spectroscopy (Fig. 5), FTIR (Fig. 6), and FL (Fig. 7) measurements. Surface morphology of the samples was investigated by SEM (Fig. 3). Photocatalytic tests were examined by UV light ($\lambda = 254\text{ nm}$) irradiation (Fig. 8).

Authors' Contributions

HX and SZ conceived the idea and designed the growth experiment and investigation process. HX, FH, and SW performed the growth experiments and photocatalytic tests. RMR and CD fabricated heterojunction samples by ALD. SW and HF performed FL and Raman tests. CX and SW performed FTIR, SEM, and XPS tests. HX, MW, LL, and SZ discussed all the results. HX, FH, CX, and SZ wrote the manuscript. All authors read, discussed, and corrected the manuscript, and approved the final manuscript.

Competing Interests

The authors declare that they have no competing interests.

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