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Electronic transport properties of (fluorinated) metal phthalocyanine

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Fadlallah, M M.; Eckern, U; Romero, A H.; and Schwingenschlogl, U, "Electronic transport properties of (fluorinated) metal phthalocyanine" (2015). *Faculty & Staff Scholarship*. 2098.
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To cite this article: M M Fadlallah *et al* 2016 *New J. Phys.* **18** 013003

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OPEN ACCESS

RECEIVED
25 August 2015REVISED
17 November 2015ACCEPTED FOR PUBLICATION
27 November 2015PUBLISHED
21 December 2015

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Keywords: phthalocyanine, charge transport, spin transport, magnetic moment, density functional theory

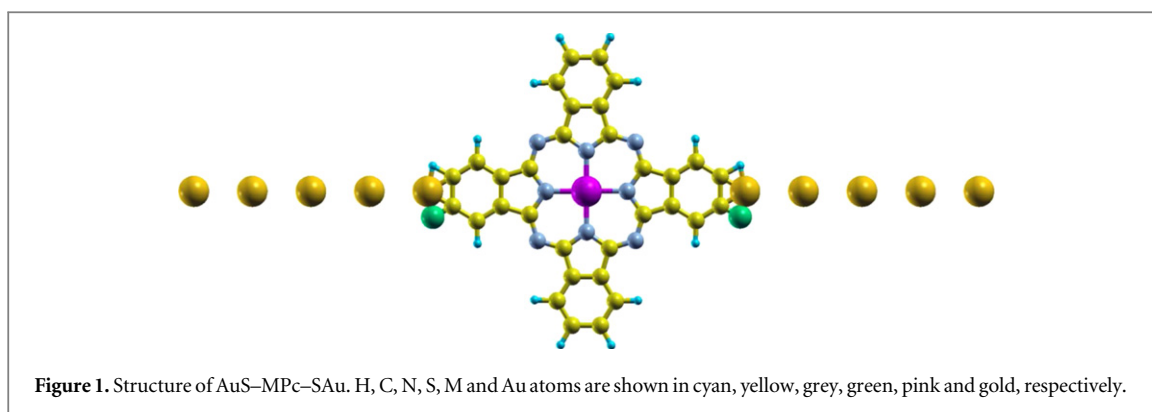
Abstract

The magnetic and transport properties of the metal phthalocyanine (MPc) and F₁₆MPc (M = Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn and Ag) families of molecules in contact with S–Au wires are investigated by density functional theory within the local density approximation, including local electronic correlations on the central metal atom. The magnetic moments are found to be considerably modified under fluorination. In addition, they do not depend exclusively on the configuration of the outer electronic shell of the central metal atom (as in isolated MPc and F₁₆MPc) but also on the interaction with the leads. Good agreement between the calculated conductance and experimental results is obtained. For M = Ag, a high spin filter efficiency and conductance is observed, giving rise to a potentially high sensitivity for chemical sensor applications.

1. Introduction

Metal phthalocyanines (MPcs) constitute a family of medium sized molecular semiconductors, which are of considerable interest for numerous applications such as chemical sensors [1], fuel cells [2], solar cells [3], and optoelectronic devices [4]. A large number of MPcs can be synthesized [5] and geometrically arranged in large areas at low cost [6], which is important from the technological point of view. In particular, the photoconductivity of MPcs has been studied intensively with the purpose of increasing the electrical conductivity of devices based on these molecules [7]. The improvement of the organic semiconductor devices relies on the quality of the metal–organic interfaces [8], in particular, on the efficiency of charge injection and on the mobility of the charge carriers [9, 10]. For example, there is a charge transfer shift of the electronic levels at the CuPc/Au interface in the early stages of CuPc deposition on Au [11], and new occupied molecular states are created [12, 13].

The replacement of the M atom as well as the substitution of H by F (fluorinated MPc) alters the gap between the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital [14], the latter resulting in n-type conduction, in contrast to the p-type behavior of the standard MPcs [15, 16]. The charge transport through the molecule can be measured directly [17, 18] or indirectly by photo-induced electron transfer [19], time resolved microwave conductivity [20], and scanning tunneling microscopy [21]. CuPc has been contacted with Au atomic chains on NiAl substrate by Nazin *et al* [22], who observed a shift and splitting of the molecular orbitals as well as modifications of the electrode orbitals by scanning tunneling microscopy. Field-effect transistors and metal–insulator–semiconductor diodes have been used to study the transport through CuPc for different leads such as Ca, Au, and F₄TCNQ/Au, demonstrating both electron or hole transport with a strong dependence on the geometry of the molecule–metal contact [23]. CuPc sandwiched between two semi-infinite Au electrodes has been investigated theoretically in [24, 25]. The transmission coefficient, $T(E)$, shows two peaks near the Fermi energy (E_F) which have been dissected in terms of molecular orbitals. The electronic states of CuPc hardly change when leads are attached.



On the other hand, the Landauer approach and Green's function formalism have been used to address the quantum transport in MPC structures. MPCs with $M = \text{Mn, Fe, Co, Ni, Cu}$ and Zn sandwiched between semi-infinite armchair single-wall carbon nanotubes (SWNT) have been considered in [26] using the SMEAGOL package [27, 28]. Within the generalized gradient approximation (GGA) it has been concluded that FePc and MnPc can be used as spin filters, and that the spin filter efficiency (SFE) increases when N-doped graphene nanoribbons are used as leads for FePc [29]. To overcome the weak interaction between Au contacts and MPCs, S atoms have been added, leading to a distinct molecular bonding. The transport properties of such AuS-MPC-SAu ($M = \text{Cu, Mn}$) systems have been studied by the WanT code, with the result that electronic correlations are likely to be irrelevant [30]. Additionally, the authors have concluded that CuPc is a molecular conductor, and MnPc a spin filter.

In this article we extend previous works on AuS-MPC-SAu and AuS- F_{16} MPC-SAu junctions by including spin polarization, and by considering a variety of different metals ($M = \text{Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn}$ and Ag). The effect of local electronic interactions on the central metal atom are also studied systematically. In section 2 we discuss the computational method, and present in section 3 the magnetic properties. Then section 4 the transmission properties and finally section 5 the SFE and electronic conductance are discussed. A summary is given in the concluding section 6.

2. Computational method

The transport properties are investigated using non-equilibrium Green's function and density functional theory as implemented in the SMEAGOL [27, 28] and SIESTA packages [31]. The wave functions are expanded in atomic orbitals with an energy cutoff of 300 Ry^6 . We use a double zeta plus polarization basis. Test calculations show no significant change when a single zeta basis is used for the M atom; however, when a single zeta basis set is used for the whole structure, the results are clearly of inferior quality. The nuclei and core electrons are represented by Troullier-Martins pseudopotentials [32]. We employ the local density approximation with Coulomb interaction ($\text{LDA} + U$). The on-site Coulomb interaction, U , is applied to the d orbitals on the central metal ions. Forces below 0.04 eV \AA^{-1} are achieved in the structural optimization of MPC, using the conjugate gradient method. Since for metallic chains a Peierls gap is expected around the Fermi energy, we do not relax the structure after attaching the leads to the molecule [33]. We use the first two of the five-Au-atom leads, to calculate the surface Green's function, while the other three, closer to the molecule, are assumed to be part of what is usually called 'scattering region'. The Au-Au distance is chosen to be 2.89 \AA [22, 23, 30], and the interface distance between MPC and Au is set to 1.37 \AA [23].

The structure of the devices is shown in figure 1. As a general remark, we note that the Mulliken analysis shows a considerable charge redistribution due to the leads (Au atoms), mainly involving the M atom, and the N and C atoms close to the central atom and at the contact between leads and molecule.

3. Magnetic moments (MMs)

The electronic and magnetic properties of *isolated* MPC and F_{16} MPC molecules have been studied before [34], where it was found that the MM is carried mainly by the metal atom. Below we will compare, in particular, isolated molecules with molecules connected to leads, with emphasis on the differences. In addition, since the central atom is a transition metal, we will encounter all possible d orbital occupations, which often give rise to

⁶ Surprisingly, there was a convergence problem with the 300 Ry cutoff for F_{16}TiPc . In this case, we managed with a cutoff of 250 Ry.

Table 1. Total and metal MMs (in μ_B) of AuS–MPc–SAu and AuS–F₁₆MPc–SAu at different U .

			Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ag
MPc	$U = 0$	total	0.8	2.9	2.1	2.3	5.3	5.9	1.6	1.5	2.2	0.4	1.6
		metal	0.1	0.8	2.9	4.0	4.6	4.1	1.3	0.1	0.5	0.0	0.2
	$U = 4$	total	0.8	3.9	3.2	3.7	5.3	5.9	4.6	3.2	1.8	0.4	1.9
		metal	0.0	1.1	3.0	4.2	4.8	4.2	2.7	1.6	0.5	0.0	0.2
	$U = 8$	total	0.8	3.8	3.2	4.1	5.3	5.8	4.4	2.9	1.6	0.4	2.1
		metal	0.0	1.1	3.0	4.2	4.9	4.6	2.8	1.7	0.5	0.0	0.2
F ₁₆ MPc	$U = 0$	total	0.8	1.6	1.2	4.0	4.7	4.1	1.0	0.0	1.1	0.0	0.7
		metal	0.0	1.5	2.4	4.0	4.7	4.1	1.1	0.0	0.5	0.0	0.2
	$U = 4$	total	0.8	1.5	3.2	4.0	4.9	3.4	2.9	1.9	1.1	0.0	0.8
		metal	0.0	1.5	2.6	4.1	5.0	4.2	2.6	1.5	0.5	0.0	0.3
	$U = 8$	total	0.8	1.4	3.3	4.0	5.0	3.4	3.0	1.9	1.1	0.0	0.9
		metal	0.0	1.3	2.3	4.1	5.0	4.3	2.8	1.6	0.5	0.0	0.3

Table 2. Spin filter efficiency for AuS–MPc–SAu and AuS–F₁₆MPc–SAu in %.

			Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ag
AuMPc	$U = 0$	–18.0	0	0	0	0	0	99.9	0	0	0	0	99.9
	$U = 4$	–3.0	0	0	0	0	0	51.7	0	0	8.3	68.4	99.9
	$U = 8$	–16.0	0	0	0	0	0	–2.5	0	0	0	72.9	97.9
AuF ₁₆ MPc	$U = 0$	48.8	99.9	14.2	0	0	0	0	0	0	0	0	99.5
	$U = 4$	70.4	98.0	64.3	0	0	0	0	0	0	–99.9	0	99.2
	$U = 8$	74.1	99.0	–90.6	0	0	0	0	0	0	–99.3	0	98.6

considerable correlation effects. Thus it will be important to also investigate, not only for a few selected cases but systematically, the effect of a local Coulomb interaction on the electronic structure as well as on transport properties. In this context, we note that U often is considered to be a parameter, to be determined by comparing specific physical quantities, such as the band gap or the MM, with experiment. However, the value obtained for U in such a way depends on the quantity considered. On the other hand, U can in principle be computed by constrained density functional theory. Both approaches appear not to be useful for the present systematic study where the system properties are strongly modified by the coupling to the leads. Thus our concept is to vary U systematically, within a reasonable range (from $U = 0$ to 8 eV, i.e., from weak to strong correlation) in order to elucidate the interaction-dependent trends in the changes of the system properties. For this goal, it appears sufficient to rely on the three values 0, 4, and 8 eV.

The calculated MMs are summarized in table 1. We first discuss the total MMs for the AuS–MPc–SAu system. For $M = \text{Sc}$ the MM differs slightly from that of the isolated molecule ($1 \mu_B$), and there is no effect of the U parameter, whereas the MM of the $M = \text{Ti}$ system is higher than for the isolated molecule ($2 \mu_B$), increasing with U . Apparently the d-d electron interaction of Ti enhances the charge transfer from the gold chain to the molecule when correlation increase. For V and Cr the MM for $U = 0$ is smaller than that of the isolated molecule ($3 \mu_B$ for VPc, $2 \mu_B$ for CrPc). As U increases from 4 to 8 eV, the MMs get closer to the isolated molecule value. For Mn junctions the MM is mainly located on Mn ($4.8 \mu_B$ for plain MnPc) and hardly depends on U .

For the Fe case, the MM of isolated FePc happens to be close to $2 \mu_B$, but is found to be strongly increased in the junction due to the overlap between Au orbitals and molecular orbitals, the MM on Fe being roughly doubled. Apparently this coupling induces a transition from a low-spin to a high-spin state, unlike the case of, e.g., Mn which is already in a high-spin state for isolated MnPc; see [34], tables 1 and 2. Thus Fe is an exceptional example where the coupling to the leads has a pronounced effect on the magnetic properties. Unfortunately, it seems not to be possible to explain this surprising fact in terms of the standard picture of atomic d orbitals in a square planar crystal field (see also below).

For the last five columns in the table, Co ... Ag, we note, first of all, that the total MM of the respective isolated molecules are close to 1, 0, 1, 0, $1 \mu_B$, respectively [34]. Thus it is apparent that the coupling to the leads has a significant effect in all cases for MPc, while the disturbance due to the leads is rather small for the fluorinated counterparts (considering first $U = 0$). For both, MPc and fluorinated MPc, we find a considerable correlation dependence for Co and Ni. However, due to the strong coupling to the leads for MPc, this dependence cannot straightforwardly explained by the electronic configurations of the respective isolated metal atoms. On the other hand, there is hardly any U -dependence for Cu, Zn and Ag (filled d shells).

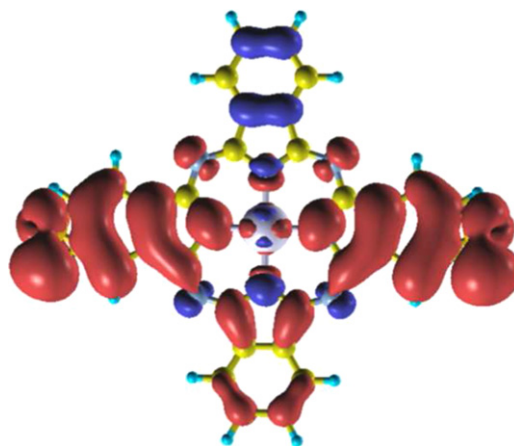


Figure 2. Spin isosurface for AuS–AgPc–SAu as obtained from SIESTA (i.e., for $U = 0$), taking into account the states between -2 and 0 eV; see figure 8. Red: majority spin, blue: minority spin. Shown is only the part containing the central molecule and the first (left and right) Au atom.

Table 1 also shows that the MM on the central metal atom in most cases differs considerably from the total MM. The U -dependence of the metal MM generally is small, except for Co and Ni.

Because of the high electronegativity of F, the MMs of the AuS–F₁₆MPC–SAu systems generally are less than or equal to those of AuS–MPC–SAu. For $M = \text{Cr, Cu, and Zn}$ the MMs of the AuS–(F₁₆MPC)–SAu junctions agree with those of the isolated molecules, independent of U . In most cases, the MMs do not change when increasing U from 4 to 8 eV, especially in the fluorinated systems. The MMs generally are slightly higher for the interacting case ($U = 4$ or 8 eV) than for zero Coulomb interaction, which is reasonable since increasing U will lead to a decrease of hybridization among the d electrons.

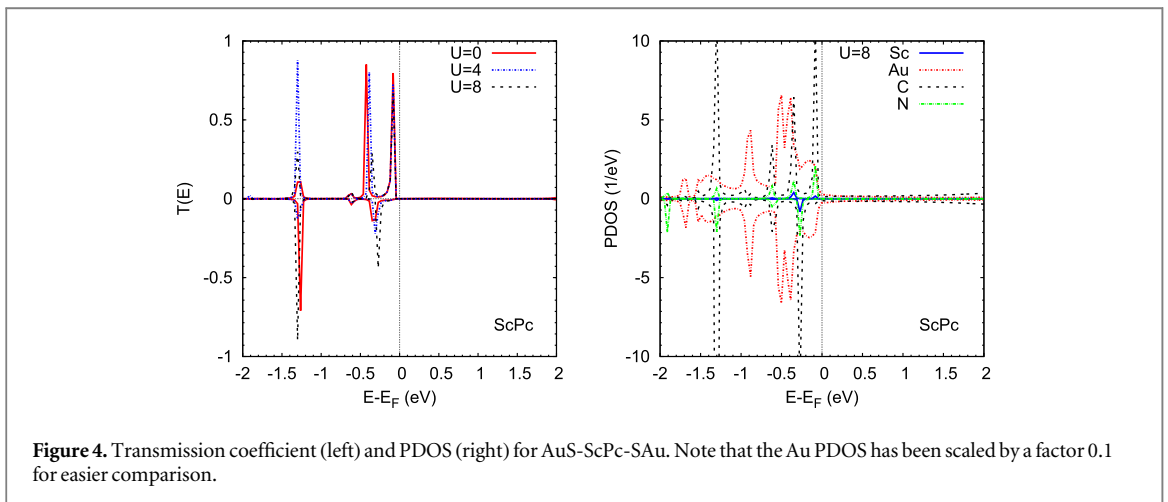
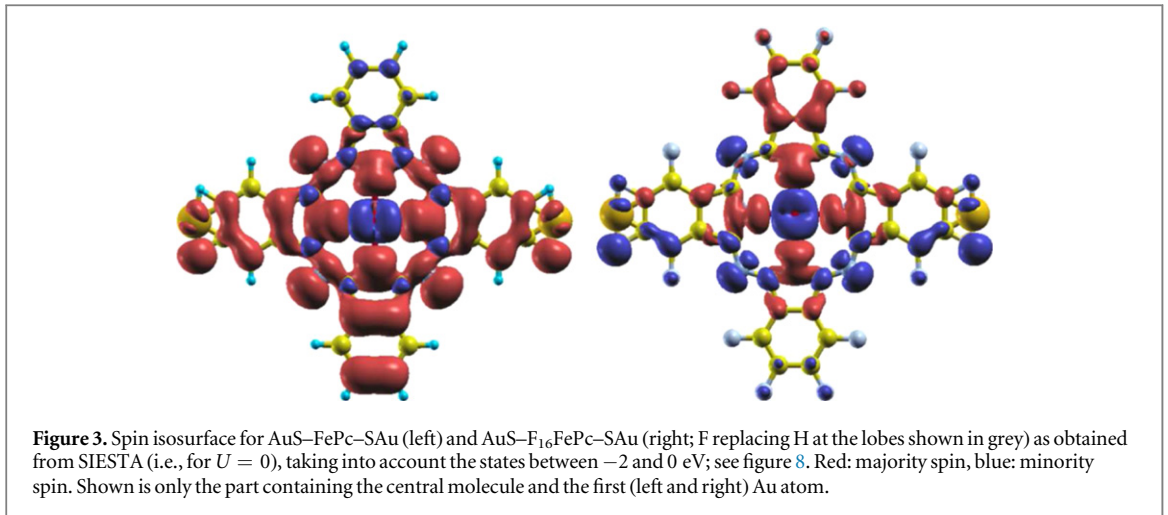
In order to obtain further insight into the above results, we also study the charge and spin density isosurfaces of the considered systems. For example, the charge isosurface of isolated ScPc was presented already in [34] (figure 2), with the result that there is hardly any charge on the metal. In the transport situation, this does not change; however, due to the good coupling to the Au leads, which is mediated by the above-mentioned addition of an S atom, some charge is being transferred to the contact region. The good coupling is also apparent in the spin isosurface. As a function of the interaction parameter U , it is found that the spin density is elongated along the transport direction, and that some spin density is shifted towards the leads, which is a general trend obtained from our data.

As another example, the spin density isosurface of AuS–AgPc–SAu is shown in figure 2. For clarity of presentation, we show here and in the next figure only the spin isosurface for the molecule plus the first Au atom on both sides, i.e. we cut out the remaining part of the leads. Of course, the isosurfaces were computed for the complete system. In figure 2 the above-mentioned good coupling is clearly visible as a relatively strong spin contributions at and near the contact. Also apparent is the small MM on the metal; note the tiny contribution from the minority spin on Ag. Similar to ScPc, the U -dependence is smooth, following the general trends mentioned above.

In figure 3 we present the spin density isosurfaces for AuS–FePc–SAu and AuS–F₁₆FePc–SAu. Most notable is the fact that the spin on Fe appears to be the same for the two cases, in agreement with the results for the MM, see table 1. For FePc, one also observes large contributions to the majority spin density from N atoms, and from C atoms which are not bonded with N. However, the C atoms which are bonded with N as well as the central metal atom contribute to the minority spin density. There is hardly any contribution from the H atoms. Concerning the fluorinated FePc system (right), the spin density is clearly smaller than for FePc; see table 1. This reduction can be attributed to the reduction in the spin density of the C and N atoms.

4. Transmission properties

In the following, we present the transmission coefficients for selected systems. Note that the d orbitals at the M atom have a D_{4h} symmetry with a crystal field that splits the orbitals as a_{1g} (d_{z^2}), b_{2g} (d_{xy}), e_g (d_{xz} , d_{yz}) and b_{1g} ($d_{x^2-y^2}$) states. For some systems, like F₁₆ScPc and TiPc, $T(E)$ is found to be negligible: the projected density of states (PDOS) of these junctions are mostly formed by sharp lines, which indicates a negligible coupling of molecular orbitals with the atomic orbitals in the leads. In the following, we concentrate on a few representative cases.



ScPc. The transmission coefficient obtained for the AuS–ScPc–SAu system is shown in figure 4. The spin components are distinguished by positive and negative values. A tiny shift of the $T(E)$ peaks is noticeable as a function of U , see left hand side of figure 4. The PDOS for $U = 8$ eV on the right hand side of figure 4 shows that the C and Au states dominate in the energy range -2 and 2 eV, whereas N and Sc states appear only below E_F . The a_{1g} states are part of the HOMO and the e_g states are part of HOMO–1 spin down and HOMO–2 spin up. The transmission peaks at -0.2 , -0.3 and -0.4 eV are mainly due to HOMO, HOMO–1 and HOMO–2, respectively, which are localized on Sc and Pc. The transmission at -1.3 eV is dominated by the molecular orbitals for both spins. The overlap between molecular orbitals and Sc orbitals is a signature that can be observed in $T(E)$ as the spatial overlap region becomes larger. The changes in $T(E)$ when increasing the interaction parameter are very small.

Fluorinated VPc. Figure 5 shows $T(E)$ and the PDOS for AuS–F₁₆VPc–SAu where all transmission peaks are below the Fermi energy. We find small variations in $T(E)$ for energies below -0.7 eV as the value of U is changed. This is likely due to the fact that when U increases, electronic localization is favored and the transmission becomes more localized in a smaller energy window, as shown in the left hand side of figure 5. Near the Fermi energy the direction of spin and the corresponding PDOS strongly depend on U . The PDOS for $U = 4$ eV is depicted on the right hand side of figure 5. The degenerate e_g states contribute to all molecular states along the junction and also to the transmission coefficient below the Fermi energy. At -0.8 and -0.9 eV we find contributions mainly from the b_{1g} orbital.

We note in passing that the transmission behavior for AuS–MnPc–SAu and AuS–F₁₆MnPc–SAu is similar to the case of ScPc and fluorinated VPc. The transmission coefficient is practically independent of U , and very quite small. For a different lead material (namely a SWNT), however, it has been suggested that MnPc can be used as spin filter [26].

FePc. We now consider Fe in the center of the Pc molecule. The effect of the U parameter is very clear in this case. Some transmission coefficient peaks happen to be independent of U , in particular, those at -1.3 (e_g and

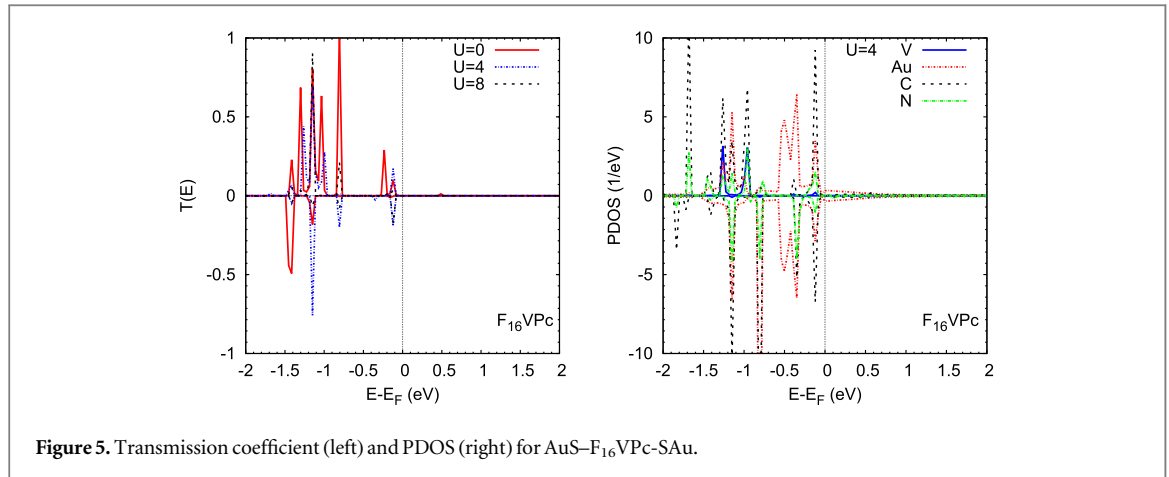


Figure 5. Transmission coefficient (left) and PDOS (right) for AuS–F₁₆VPc–SAu.

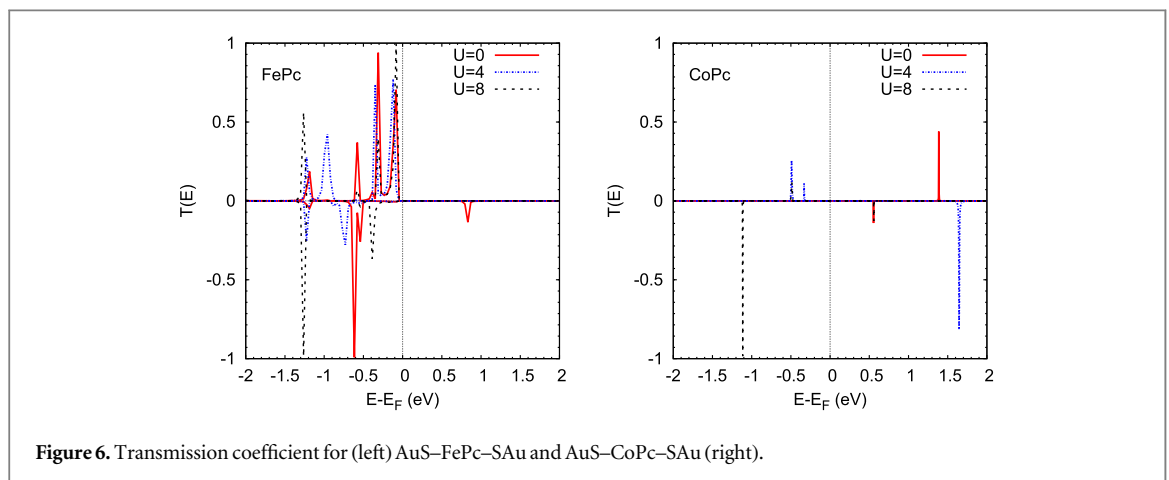


Figure 6. Transmission coefficient for (left) AuS–FePc–SAu and AuS–CoPc–SAu (right).

molecular states) for both spins, -0.4 eV ($a_{1g}b_{2g}$ and molecular states), and -0.1 eV (b_{1g} and all other states of Au and the molecule) for spin up. From -1 to -0.7 eV, some $T(E)$ peaks appear only for $U = 4$ eV, and at -0.4 eV for $U = 8$ eV only in case of spin down. Furthermore, at 0.9 eV ($e_g b_{1g}$) but only for spin down, there is a single $T(E)$ peak for $U = 0$. For $U = 0$ and 8 eV, $T(E)$ has peaks at -0.6 eV ($e_{1g}b_{1g}$ and molecular states). We find zero transmission at the Fermi energy, in contrast to the case where FePc is connected to a SWNT [26, 29].

CoPc. For the case of Co in the center of the Pc molecule, the $T(E)$ peaks are sharp lines, and their positions vary with U . There are two peaks below and one peak above E_F for every U value, see right hand side of figure 6. For $U = 0$, the peak in $T(E)$ at -0.5 eV is related to molecular as well as Au and b_{1g} states. There are two additional peaks at 0.5 eV (molecular orbitals with b_{1g}) and at 1.4 eV (molecular orbitals with $e_g b_{1g} a_{1g}$).

CuPc. For the AuS–CuPc–SAu junction, several transmission peaks are found below the Fermi energy, and, in particular, they are enhanced for $U = 4$ eV, see figure 7. Above the Fermi energy, there is a single peak, which changes position and intensity as function of U , the position shifting away from E_F as U increases (because of increasing repulsion between electrons). The energy difference between the first peak below and the first one above E_F for $U = 4$ is smaller than the corresponding one for $U = 0$ and 8 eV.

AgPc. Last but not least we discuss AuS–AgPc–SAu and AuS–F₁₆AgPc–SAu (figure 8). The $T(E)$ peaks demonstrate a good coupling between lead and molecule as shown in the left hand side of figure 8 (AuS–AgPc–SAu). We find peaks for spin up states at E_F , independent of U , which indicates that this configuration has a high potential as a spintronic device. The change of the interaction parameter does not affect the positions of the peaks below the Fermi energy. However, we observe a strong dependence of the peak position and intensity on U above the Fermi energy: as U increases, the peak moves further away from E_F , and the intensity increases almost by a factor of two. This indicates that as electron localization increases, the transmission also increases for that particular state. For the AuS–F₁₆AgPc–SAu junction, right hand side of figure 8, the transmission peak near -0.1 eV is finite for both spin polarizations, and independent of U . The transmission near the Fermi energy is dominated by molecular states with $a_{1g} b_{1g}$ orbitals.

For completeness we discuss the occupation of the d orbitals on the metals which were addressed in detail in this section. First, we find only minor differences between MPc and their fluorinated counterparts. Second, the

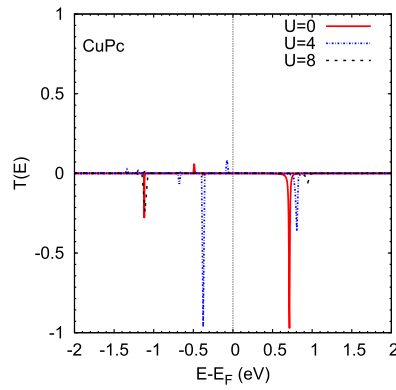


Figure 7. Transmission coefficient for AuS–CuPc–SAu.

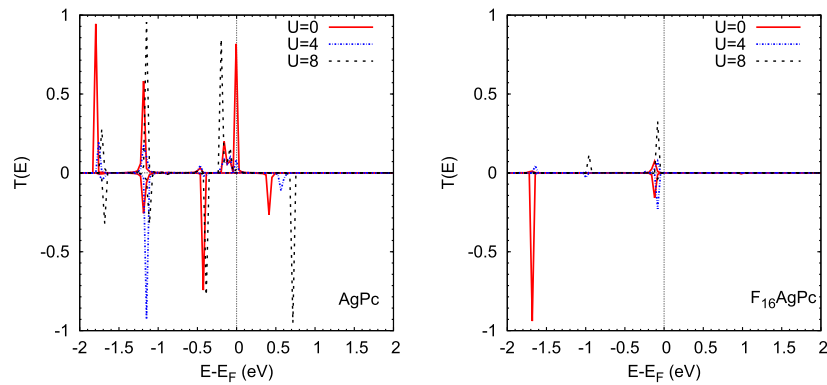


Figure 8. Transmission coefficient for AuS–AgPc–SAu (left) and AuS–F₁₆AgPc–SAu (right).

occupation numbers of the five orbitals add up almost to the values to be expected from the periodic table. Sc: ≈ 0.7 , decreasing by about 20 % with increasing U ; V: ≈ 3.1 , almost independent of U ; Fe: ≈ 5.5 , decreasing by about 10 % with increasing U ; Co: close to 7, independent of U ; Cu: ≈ 9.3 , almost independent of U ; Ag: ≈ 9.4 , increasing by about 1 % with increasing U . Considering the simple (from an atomic point of view) case of Sc, it is already apparent that an interpretation in terms of the standard level picture is incomplete. This conclusion is further supported by the case of Fe, which has a rather high MM as discussed above: We find that all five d orbitals are occupied by close to one up electron, while the spin down contribution is mainly due to an e_g orbital (≈ 0.5 for $U = 4$ eV, dropping to ≈ 0.1 when U is increased to 8 eV, consistent with the increase of the metal MM from 4.2 to 4.6 μ_B , cf table 1). For the V system, the d orbital occupations depend only weakly on the interaction: The e_g (up) and b_{1g} (up) occupations are slightly increased when U is increased, while the others decrease by a small amount. For the Co system, we find the up occupations to be close to one, while the occupations of the b_{2g} (down), a_{1g} (down), and one of the e_g (down) orbitals are considerably smaller. For Cu and Ag, all occupations are close to one, except for the a_{1g} (down) orbital (≈ 0.6 and ≈ 0.7 , respectively).

5. Spin filter efficiency and conductance

The SFE for an electronic device, defined as

$$\text{SFE} = \frac{T_{\uparrow}(E_F) - T_{\downarrow}(E_F)}{T_{\uparrow}(E_F) + T_{\downarrow}(E_F)}, \quad (1)$$

is the ability of a device to pass a particular spin component. Here, $T_{\uparrow}(E_F)$ and $T_{\downarrow}(E_F)$ denote the transmission coefficients of spin up and spin down electrons at the Fermi energy, respectively. Our findings are summarized in table 2, the positive values corresponding to the case where ‘up’ is the majority spin, while a negative sign means that the majority spin is ‘down’. We note that the magnitude and the sign of the SFE can vary with the interaction parameter. In some cases the majority spin in F₁₆MPc is opposite to the majority spin in MPc, e.g.,

for Sc and Cu. Our results suggest that junctions containing Sc and Ag, as well as FePc, ZnPc, F₁₆TiPc, F₁₆VPc and F₁₆CuPc could potentially be used in spintronic devices. Note, however, that ScPc is chemically unstable [35, 36]. In several of the considered molecules we observe zero SFE (independent of U), in particular, for CrPc, CoPc, NiPc, and their fluorinated counterparts.

However, a reasonable electronic conductance is also essential for the performance of a device, hence we determine in addition the conductance

$$G = G_0 \left[T_{\uparrow}(E_F) + T_{\downarrow}(E_F) \right], \quad G_0 = 2e^2/h. \quad (2)$$

For the uncorrelated case, $U = 0$, the junction which includes AgPc shows the highest conductance, close to $0.22 G_0$. The next higher one is CuPc with $1.8 \times 10^{-4} G_0$. This value is smaller than the experimental result found in aromatic molecules by two orders of magnitude [37, 38], but our result roughly agrees with another theoretical calculation [25]. Differences in theoretical results generally depends on the construction of the junctions, as well as the employed calculational method and the details of the (weak) interaction between molecule and leads. From a practical point of view, the results will also depend on the interaction between the junction and the substrate, which is not taken into account here and in most other studies.

The conductance of ZnPc is found to be $3 \times 10^{-5} G_0$, which is the same order as previous experimental and theoretical results for an oligo-porphyrin molecular wire by using Zn in the molecule center [43]. We mention in passing that for TiPc and VPc, which are also chemically unstable [36, 39–42], a vanishing conductance is obtained.

In F₁₆MPc, the conductances in most cases are less than for the corresponding MPc. The highest conductance is found for F₁₆AgPc, namely $0.002 G_0$, the next highest value being $1.5 \times 10^{-4} G_0$ for F₁₆FePc. In addition, the conductance of F₁₆NiPc and F₁₆CuPc is of the same order as the conductance of F₁₆FePc.

The electronic conductance is somewhat insensitive to an increase of U for all systems, in particular, it remains zero for those cases where we find zero conductance at $U = 0$. On the other hand, when the conductance is finite at $U = 0$, we find it either to remain constant or to decrease with increasing U . For example, for $U = 4$, the conductances for CuPc, ZnPc and AgPc are 6×10^{-8} , 3×10^{-6} and $1.1 \times 10^{-2} G_0$, respectively. When we increase U to 8 eV the conductances for the same systems drop to 0, 3×10^{-8} and $7.3 \times 10^{-3} G_0$, respectively.

6. Summary

In this paper we studied systematically the magnetic, electronic and transport properties of metal phtahlocyanines and fluorinated MPcs connected to Au leads, including some MPc molecules which are chemically unstable. We employed the LDA + U approach, with values of the Coulomb interaction parameter ranging from $U = 0$ to $U = 8$ eV. The MMs are largely determined by the hybridization between d metal and Au states near the Fermi energy. The MMs (table 1) to some extent vary with the electron-electron interaction on the central metal atom, U , which also for some systems considerably modifies the SFE (table 2). Considering equation (1), it is apparent that the SFE is a very sensitive quantity whenever the transmission coefficients for both channels are very small.

In particular, the MMs are found to increase with increasing correlation, or are roughly independent of U . For the transport properties, the situation is less clear: it appears that for some systems the SFE hardly depends on U , while others are quite sensitive to correlation effects. There can be a remarkable difference between MPc and its fluorinated counterpart. In detail, the results can be explained by the respective electronic structure near the Fermi surface, which, however, is not only determined by the orbitals of the central metal atom but also by the coupling to the leads. In order to finally clarify the role of electronic correlation in the series of junctions considered, additional experimental efforts will be needed. Overall our theoretical results compare well with other theoretical and experimental results for the electronic conductance where these are available.

The structures with Ag in the center of the junction are a notable exception: they show both a good SFE and a reasonable electronic conductance since the contribution of Ag states at the Fermi energy is large. Hence AgPc and F₁₆AgPc junctions can potentially be used in spintronic devices.

Acknowledgments

We thank A A Maarouf, I Rungger, and C Schuster for fruitful discussions, and acknowledge financial support by the Deutsche Forschungsgemeinschaft (through TRR 80). AHR acknowledges the support by the Marie Curie Actions from the European Union in the international incoming fellowships (grant PIIFR-GA-2011-911070), and the Donors of the American Chemical Society Petroleum Research Fund under contract 54075-ND10 for partial support of this research. Research reported in this publication was supported by the King Abdullah University of Science and Technology (KAUST).

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