

# THE UNIVERSITY of EDINBURGH

# Edinburgh Research Explorer

# Deciphering the Structural Evolution and Electronic Properties of Magnesium Clusters: An Aromatic Homonuclear Metal Mg17 Cluster

#### Citation for published version:

Xia, X, Kuang, X, Lu, C, Jin, Y, Xing, X, Merino, G & Hermann, A 2016, 'Deciphering the Structural Evolution and Electronic Properties of Magnesium Clusters: An Aromatic Homonuclear Metal Mg17 Cluster', *The Journal of Physical Chemistry A*, vol. 120, no. 40, pp. 7947-7954. https://doi.org/10.1021/acs.jpca.6b07322

#### **Digital Object Identifier (DOI):**

10.1021/acs.jpca.6b07322

#### Link:

Link to publication record in Edinburgh Research Explorer

**Document Version:** Peer reviewed version

Published In: The Journal of Physical Chemistry A

#### **General rights**

Copyright for the publications made accessible via the Edinburgh Research Explorer is retained by the author(s) and / or other copyright owners and it is a condition of accessing these publications that users recognise and abide by the legal requirements associated with these rights.

#### Take down policy

The University of Édinburgh has made every reasonable effort to ensure that Edinburgh Research Explorer content complies with UK legislation. If you believe that the public display of this file breaches copyright please contact openaccess@ed.ac.uk providing details, and we will remove access to the work immediately and investigate your claim.



## **Deciphering the Structural Evolution and Electronic Properties of**

# Magnesium Clusters: a New Aromatic Homonuclear Metal Mg<sub>17</sub>

## Cluster

Xin-Xin Xia,<sup>1,2</sup> Xiao-Yu Kuang,<sup>1,a)</sup> Cheng Lu,<sup>2,3,b)</sup> Yuan-Yuan Jin,<sup>1</sup> Xiao-Dong Xing,<sup>1</sup> Gabriel Merino,<sup>4</sup> and Andreas Hermann<sup>5,c)</sup>

<sup>1</sup>Institute of Atomic and Molecular Physics, Sichuan University, Chengdu 610065, China

<sup>2</sup>Department of Physics, Nanyang Normal University, Nanyang 473061, China

<sup>3</sup> Department of Physics and High Pressure Science and Engineering Center, University of

Nevada, Las Vegas, Nevada 89154, United States

<sup>4</sup> Departamento de Física Aplicada, Centro de Investigación y de Estudios Avanzados, Unidad

Mérida, Km 6 Antigua Carretera a Progreso, Apdo. Postal 73, Cordemex, 97310 Mérida,

Yucatán, Mexico

<sup>5</sup> Centre for Science at Extreme Conditions and SUPA, School of Physics and Astronomy, The

University of Edinburgh, Edinburgh EH9 3JZ, United Kingdom

# ABSTRACT

The structures and electronic properties of low energy neutral and anionic Mg<sub>n</sub> (n = 3-20) clusters have been studied by utilizing a widely adopted CALYPSO structure searching method coupled with density functional theory calculations. A large number of low energy isomers are optimized at the B3PW91 functional with the 6-311+G(d) basis set. The optimized geometries clearly indicate that a structural transition from

<sup>&</sup>lt;sup>a)</sup> Electronic mail: <u>scu\_kuang@163.com (Xiao-Yu</u> Kuang).

<sup>&</sup>lt;sup>b)</sup> Electronic mail: <u>cheng.lu@unlv.edu</u> (Cheng Lu).

<sup>&</sup>lt;sup>c)</sup> Electronic mail: <u>a.hermann@ed.ac.uk (</u> Andreas Hermann)

hollow three-dimensional configurations to filled-cage-like structures occurs at n = 16 for both neutral and anionic clusters. Based on the anionic ground state structures, photoelectron spectra are simulated using time-dependent density functional theory (TD-DFT) and compared with experimental results. The good agreement validates that the current ground state structures, obtained from the symmetry-unconstrained searches, are true global minima. A detailed chemical bonding analysis distinctly indicates that the Mg<sub>17</sub> cluster is the first neutral locally  $\pi$ -aromatic homonuclear all-metal cluster, which perfectly satisfies Hückel's well-known 4N+2 rule.

## I. Introduction

Magnesium is an interesting metal that exhibits strong chemical activity. Its unique features have initiated extensive research in diverse fields such as superconductivity,<sup>1</sup> hydrogen storage,<sup>2,3</sup> nanomaterials,<sup>4,5</sup> even biomedicine.<sup>6,7</sup> Apart from their direct influence in these fields, magnesium clusters exhibit a number of unique phenomena and features. One of the most striking phenomena is the transition from weak van der Waals bonding to metallic bonding as the clusters grow in size. The critical size for nonmetallic to metallic behavior was experimentally established to be n = 18 for anions,<sup>8</sup> and suggested to be n = 20 for neutral clusters.<sup>9</sup> From a theoretical perspective, the precise size of the insulator-to-metal transition is yet to be determined.<sup>10-16</sup>

There has been a plethora of theoretical studies on the geometric structures and electronic properties of magnesium clusters.<sup>10-22</sup> Jellinek *et al.*<sup>11</sup> investigated neutral

and anionic magnesium clusters up to n = 22 using gradient-corrected density functional theory, and found that the electron binding energies in the anionic clusters, derived from the gap between the two most external electrons, agree with electron photodetachment experiment data. Lyadin et al. investigated the structural evolution and electronic shells of neutral and cationic magnesium clusters using ab initio theoretical methods, where results showed that the metallic evolution is a slow and non-monotonous process.<sup>20</sup> Exhaustive structure optimization revealed the growth behaviors of larger neutral magnesium clusters and indicated that most of the ground state structures are non-symmetric in the size range of n > 20.<sup>22</sup> It is worth pointing out that the neutral ground state structures found in the different studies are not identical.<sup>11,20,22</sup> Thus, although structures of magnesium clusters have been extensively studied and specific structures have been described and analysed, the true ground state structures might still be debatable. Their electronic properties have been investigated by both experiment and theory, but these do not provide systematic information about their chemical bonding. It is therefore timely to (i) compare theoretical photoelectron spectra with the experimental PES data to verify the lowest-energy structures of Mg clusters; (ii) investigate the most probable fragmentation channels for magnesium clusters; and (iii) understand the chemical bonding of Mg clusters, which also serve to explain their stabilizing mechanisms. To that end, we invoke the concept of aromaticity in our analyses. Aromaticity is

traditionally confined to the realm of organic chemistry to describe cyclic, delocalized  $\pi$  bonding in planar and conjugate molecules possessing (4*N* + 2)  $\pi$  electrons.<sup>23-29</sup>

(Merino G. Sola M. Celebrating the 150th anniversary of the Kekulé benzene structure. Phys. Chem. Chem. Phys. 2016, 18, 11587-11588). Recently, this concept has been extended to inorganic molecules including organometallic compounds,<sup>30</sup> (Add the next reference Fernandez, I.; Frenking, G.; Merino, G. Aromaticity in Metallabenzenes and Related Compounds. Chem. Soc. Rev. 2015, 44, 6452-6463.) transition-metal systems<sup>31-34</sup> and, in particular, all-metal clusters.<sup>35-37</sup> The gallium-gallium bond is exemplary for our understanding of electronic structure and chemical bonding in organometallic chemistry, with the first example of metalloaromaticity.<sup>30</sup> Combined anion photoelectron spectroscopy and quantum chemical calculations on metal-boron clusters indicated that various planar M@B<sub>n</sub><sup>-</sup> clusters feature planar aromatic boron rings, with delocalized  $\pi$ -electrons and the metal atom strongly covalently bound to the surrounding boron atoms.<sup>31,32</sup> (Add the next reference: Islas, R.; Heine, T.; Ito, K.; Schleyer, P. v. R.; Merino, G. Boron Rings Enclosing Planar Hypercoordinate Group 14 Elements. J. Am. Chem. Soc. 2007, 129, 14767-14774) A series of bimetallic metal anionic clusters with chemical composition  $MAl_4^-$  (M = Li, Na, or Cu) were synthesized and studied with photoelectron spectroscopy and *ab initio* calculations. This study showed that Al4<sup>2-</sup> maintains a square planar structure and aromaticity due to two delocalized  $\pi$  electrons present in all three MAl<sub>4</sub><sup>-</sup> complexes.<sup>35</sup> Considering the concept of aromaticity, other small alkali metal and alkaline earth metal clusters were studied and for Li<sub>2</sub>Mg<sub>2</sub> it was shown that the cyclic  $\sigma$ -aromatic structures are more stable than the classical linear Li-Mg-Mg-Li structure.<sup>36</sup> These investigation extended the aromaticity concept to the

realm of all-metal species and highlighted its importance. Like aromaticity, which is then well-established in inorganic and all-metal systems, the concept of antiaromaticity (cyclic conjugated systems with  $4N \pi$  electrons)<sup>38</sup> has recently been extended beyond organic molecules as well. The Al4<sup>4-</sup> anion in the Li<sub>3</sub>Al4<sup>-</sup> mixed-metal cluster was shown to feature  $\pi$ -antiaromaticity (and  $\sigma$ -aromaticity).<sup>39</sup> (Islas, R.; Heine, T.; Merino, G. Structure and Electron Delocalization in Al42- and Al44-. J. Chem. Theory Comp. 2007, 3, 775-781.) In recent years, three-dimensional all-metal clusters, for instance the Sb<sub>4</sub> unit in  $[Ln(\eta^4-Sb_4)_3]^{3-}$  (Ln = La, Y, Ho, Er, Lu)<sup>40</sup>, have been studied. Based on the results of chemical bonding analyses and Breslow's 4*N* rule, the  $[Ln(\eta^4-Sb_4)_3]^{3-}$  (Ln = La, Y, Ho, Er, Lu) compound was found to be the first locally  $\pi$ -antiaromatic all-metal system. Despite aromaticity and antiaromaticity having been observed in inorganic and all-metal systems, these concepts have been applied relatively little to homonuclear metal clusters. However, it is mandatory to do an exhaustive and comparative study in order to establish if an all-metal aromatic cluster can be classified as aromatic (ADD the next reference: Hoffmann R. The Many Guises of Aromaticity. Am. Sci. 2015, 103, 18) With this in mind, the exploration of the geometric structures, electronic properties and the nature of the chemical bonding in magnesium clusters, is an intriguing proposition.

In order to study systematically and in-depth the structural evolution and electronic properties of magnesium clusters, we have carried out comprehensive structure searches on neutral and anionic magnesium clusters in the size range from n = 3 to 20, by combining a systematic exploration of the potential energy surface using

CALYPSO (Crystal structure AnaLYsis by Particle Swarm Optimization)<sup>41-44</sup> with density functional theory calculations. Subsequently, we reexamine the structure of particular neutral and anionic low-energy isomers of magnesium clusters with respect to those reported in previous experiments or theoretical calculations. We then study the stabilizing mechanism due to the electronic properties of specific neutral and anionic magnesium clusters and provide new insights for further theoretical and experimental explorations. The paper is organized as follows. The details of the computational method are presented in Sec. II. Then, our results and discussion are described in Sec. III. Finally, the main conclusions are summarized in Sec. IV.

# **II.** Computational details

The structures of low-lying isomers of neutral and anionic magnesium clusters were obtained using CALYPSO, which provides a local version of the particle swarm optimization (PSO) algorithm to explore the free-energy surfaces for any given (non-)periodic system. The algorithm can predict stable structures depending only on the chemical composition. It has successfully predicted structures for various systems ranging from clusters<sup>45</sup> to surface reconstructions<sup>43,46</sup> and crystal structures.<sup>42,44,47</sup> Here, structure predictions are performed for neutral and anionic magnesium clusters up to 20 atoms. Each generation contains 50 structures, 60% of which are generated by the PSO, while the others are new and will be generated randomly. We have followed 30 generations for each cluster to achieve convergence of the potential energy surface sampling. The searches generated 1000–1500 isomers for low energy

neutral and anionic magnesium clusters. Among those isomers, the fifteen energetically lowest-lying isomers are selected as candidates for the global minimum structure. Low energy structures within 3 eV of the global minimum structure are further optimized with subsequent frequency calculations. The calculations are performed using the all-electron density functional theory method with the B3PW91<sup>48-50</sup> generalized gradient approximation functional. The 6-311+G(d) basis set is selected for the confirmation of the lowest energy structures of magnesium clusters. The choice of B3PW91/6-311+G(d) level of theory set is based on a previous report.<sup>20</sup> The effect of spin multiplicity (up to septet and octet) is taken into account and no symmetry constraints are enforced in the geometric optimization procedure. All calculations are performed with the Gaussian09 program package.<sup>51</sup> The photoelectron spectra of the anionic magnesium clusters are simulated with the time-dependent density functional theory (TD-DFT) method. Chemical bonding analyses (B3PW91/6-311+G(d)) are conducted using the adaptive natural density partitioning (AdNDP)<sup>52</sup> method. The nucleus-independent chemical shift (NICS) and multicenter bond order are calculated by using the Multiwfn 3.3.8 program package.<sup>53</sup>

# **III. Results and discussions**

#### A. Geometric structure

We have performed a comprehensive structure search and all of the previously reported structures, including experimental and theoretical ones, are successfully reproduced in our search results. Based on those results, we optimize the candidates of low-lying isomers and display the global minima structures of neutral and anionic magnesium clusters with up to 20 atoms in Figure 1. The vibrational frequencies are also calculated and listed in Table S1 (in the Supporting Information, SI) to assure true global minima for the ground state structures. Moreover, other typical low-energy isomers of all clusters, together with their corresponding symmetry, are displayed in Figure S1 (see SI).

From Figure 1, we can see that the neutral and corresponding anionic magnesium clusters have similar ground-state structures and follow the same structural evolution. For both neutrals and anions, only the Mg<sub>3</sub><sup>0/-</sup> clusters are trigonal plane structures. For  $n \ge 4$ , the lowest energy structures of magnesium clusters form three-dimensional configurations. As the number of atoms increases, we find that eventually, for  $n \ge 16$ , one of the Mg atoms is fully encapsulated within the magnesium framework. So, n = 16 is a structural transition point from hollow three-dimensional configurations to filled-cage-like structures. All of the lowest energy structures of neutral and anionic Mg<sub>n</sub> (n = 3-20) clusters, except for Mg<sub>14,16</sub> and Mg<sub>12,14,20</sub><sup>-</sup>, which are of  $C_1$  symmetry, possess relatively high point symmetry. Additionally, for n = 3-5, 9, 10, 13, 15, 19, the neutral and anionic clusters have the same geometries and point symmetries.

The optimized lowest energy structures found in our searches are largely in agreement with those discussed in earlier theoretical and experimental findings. However, we find a Mg<sub>13</sub> cluster of  $C_s$  symmetry (Figure 1) to be the lowest energy structure in our study, while the Mg<sub>13</sub> cluster with  $C_1$  symmetry (Figure S1), previously studied theoretically by Lyalin *et al.*<sup>20</sup>, is shown to be a transition state. For

Mg<sub>19</sub>, Lyalin *et al.*<sup>20</sup> claimed that global minimum structure has  $C_{2\nu}$  symmetry whereas our result found it to be of  $C_s$  symmetry, which is in agreement with the result reported by Heidari *et al.*<sup>22</sup> The structure of Mg<sub>20</sub> is in agreement with the result of Jellinek *et al.*,<sup>11</sup> and different from that reported by Lyalin *et al.*<sup>20</sup> The global minimum structure of Mg<sub>8</sub><sup>-</sup> is found here to be of  $C_2$  symmetry, while the  $D_{4d}$ symmetry structure proposed by Jellinek *et al.*<sup>11</sup> is not the lowest energy structure, as shown by the results of our harmonic vibration analysis. Note that the lowest energy Mg<sub>n</sub> (n = 3-20) clusters are found to prefer low spin state, except in the case of Mg<sub>18</sub><sup>0/-</sup> where the ground state are triplet and quartet, respectively. The spin states of all these lowest energy clusters are in good agreement with the results of Acioli *et al.*<sup>12</sup>

#### B. Electronic properties of anionic Mg<sub>n</sub><sup>-</sup> clusters with ground state structures

In order to confirm the validity of the ground state structures (shown in Figure 1) of  $Mg_n^-$  (n = 3-20) clusters, their photoelectron spectra were simulated using TD-DFT. The simulated spectra of the ground state structures are displayed in Figure 2, along with the available experimental spectra<sup>9</sup> for comparison. Simulated spectra of other low-energy isomers are shown in Figure S2 in the SI. The vertical detachment energy (VDE) was taken from the first peak position of the spectra and the adiabatic detachment energy (ADE) for the neutral clusters was measured by the corresponding intersection between the baseline and the rising edge of the first peak. Comparing the calculated and experimental numbers in Table 1, we can see that both the VDE and

the ADE values extracted from the simulated PES are in satisfactory agreement with the experimental data, supporting the reliability of our theoretical approach. The spectra themselves are in overall good agreement with experiment throughout, though spectral weight seems to be off in some cases (see  $Mg_{11,18,19}$ ). The photoelectron spectra for  $Mg_n^-$  with n = 3-13 show only one major peak in the range of low binding energy ( $\leq 2.0$  eV), while several recognizable peaks are found in the binding energy range of 2.0–3.3 eV. Notice that for cluster sizes n = 16-19, the VDE's are larger than 2.0 eV, followed by a drop to 1.88 eV for n = 20. From the simulated PES, evidently the energy spacing between the first two main spectral features, which can be interpreted as the HOMO-LUMO gap for neutral magnesium clusters, has a local maximum (compared to their neighbors' gaps) for n = 4, 10 and 20. This phenomenon is consistent with the neutral HOMO-LUMO gap analysis discussed below and offers a first hint towards stability of magic number clusters. Overall, the comparison of the experimental and simulated PES results lends support to the proposed ground state structures of magnesium clusters.

The experimental and theoretical values of ADE and VDE, listed in Table 1, have been plotted as function of cluster size in Figure 3 to explore more deeply their electronic properties. The theoretical values of ADE increase from 0.80 to 1.89 eV for n = 3-9 (see Figure 3a). Afterward, the theoretical ADE reveals obvious odd–even oscillation from n = 13 to 18. The theoretical ADEs successfully reproduce the overall experimental trend, and yield two striking minima at n = 10 and 20, which are related to the shell model introduced by Thomas *et al.*<sup>9</sup> As shown in Figure 3b, the theoretical values of the VDE keep increasing overall with increased numbers of Mg atoms, and two obvious minima appear again at n = 10 and 20. Again, the calculated VDEs are in good agreement with the experimental results, with only n = 6, 12, 18, 19 deviating by more than 0.1eV. Note that for most cluster sizes, the deviation between theoretical values of ADE and VDE is less than 0.13 eV, with exceptions n = 6-8, 11, and 18. This is in agreement with the observation that all clusters, except for Mg<sub>6-8,11,18</sub>, basically retain the structural framework of the neutral ground state in the corresponding anionic cluster.

#### C. Relative stabilities and HOMO-LUMO gaps

The inherent stability of a given cluster might be determined by its binding energy  $(E_b)$  per atom that, for neutral and anionic magnesium clusters, can be defined as:

$$E_{b} = \left[ \left( n-1 \right) E \left( \mathrm{Mg} \right) + E \left( \mathrm{Mg}^{Q} \right) - E \left( \mathrm{Mg}^{Q}_{n} \right) \right] / n \quad , \qquad Q \quad = \quad 0, \qquad -1,$$
(1)

where *E* is the total energy of the corresponding atom or cluster. The calculated results of  $E_b$  are summarized in Table 1 and are plotted as function of cluster size *n* in Figure 4a. For both the neutral and anionic magnesium clusters, the binding energies mostly increase with cluster size, and both curves show similar size dependence (see Figure 4b). Three weak local peaks of the binding energy occur at n = 4, 10 and 15 for neutral clusters and n = 4, 9 and 15 for anionic clusters, respectively. This indicates that Mg<sub>4,10,15</sub> and Mg<sub>4,9,15</sub><sup>-</sup> are more stable than their adjacent sized clusters. It is also

worth noticing that the  $E_b$  values of anionic  $Mg_n^-$  cluster are always higher than their neutral counterparts. This implies that the binding energy of an additional electron is larger in the clusters than in a single  $Mg^-$  anion.

The second-order difference of the energy  $(\Delta^2 E)$  is another important parameter that can reflect the relative stabilities of neutral and anionic clusters, and can be defined here as

$$\Delta^2 E = E\left(\mathrm{Mg}_{n-1}^{\mathcal{Q}}\right) + E\left(\mathrm{Mg}_{n+1}^{\mathcal{Q}}\right) - 2E\left(\mathrm{Mg}_n^{\mathcal{Q}}\right), \ Q = 0, -1.$$
<sup>(2)</sup>

The  $\Delta^2 E$  values for neutral and anionic Mg<sub>n</sub> (n = 3-20) clusters are listed in Table S2 and shown in Figure 4b. Figure 4b reveals for both neutral and anionic clusters irregularly oscillating behavior in the region of n = 3-9 and odd-even oscillations for n = 12-18. Several pronounced peaks are found at n = 4, 7, 13, 15 and 17, signifying that the clusters Mg<sub>4</sub><sup>0/-</sup>, Mg<sub>7</sub><sup>0/-</sup>, Mg<sub>13</sub><sup>0/-</sup>, Mg<sub>15</sub><sup>0/-</sup>, and Mg<sub>17</sub><sup>0/-</sup> are more stable compared to their neighbors. Furthermore, the neutral Mg<sub>10</sub> and anionic Mg<sub>9</sub><sup>-</sup> clusters are relatively more stable due to their locally maximal  $\Delta^2 E$  values.

The HOMO–LUMO gaps are calculated for the ground state  $Mg_n^Q$  (n = 3-20, Q = 0, -1) structures. The energy gap between the highest occupied molecular orbital and the lowest unoccupied molecular orbital is an indicator of relative stability.<sup>54</sup> The results for the clusters in their ground state structures are summarized in Table 1 and plotted in Figure 4c. By comparing the  $E_{gap}$  values of neutral and anionic magnesium clusters, it can be seen that the  $E_{gap}$  curve for neutral clusters shows a decreasing trend

with size, with significant fluctuations, while a much less pronounced trend is perceived for anionic clusters. In addition, the values for neutral clusters are always higher than for anionic clusters, indicating that neutral Mg<sub>n</sub> clusters are relatively more stable than anionic clusters. Local maxima values of  $E_{gap}$  are found at n = 4, 7,11, 13, and 15 for neutral and n = 10, 16, and 18 for anionic clusters, respectively. This suggests that these cluster sizes have a stronger stability than their size-adjacent clusters.

#### **D.** Fragmentation channels

Potential fragmentation paths can be studied from a thermodynamic viewpoint (not considering kinetic barriers) and the fragmentation energies  $E_f$  of the ground state magnesium clusters can be expressed as:

$$E_f = E\left(\mathrm{Mg}_p\right) + E\left(\mathrm{Mg}_{n-p}^{\mathcal{Q}}\right) - E\left(\mathrm{Mg}_{n}^{\mathcal{Q}}\right), Q = 0, -1$$
(3)

Generally speaking, if the fragmentation energy is negative for a particular fragmentation channel, the initial cluster is unstable and may dissociate spontaneously by releasing the amount of energy  $E_{f}$ . In this work, we have calculated all possible fragmentation channels for  $Mg_n^Q$  (n = 3-20, Q = 0, -1) clusters and the respective values are positive in all cases, implying that the clusters are stable and must obtain energy to realize the fragmentations. The smaller the fragmentation energies of cluster are, the less stable the clusters will be. The easiest (least energy-costly) fragmentation channels and corresponding fragmentation energies ( $E_f$ ) for neutral and anionic

magnesium clusters are summarized in Table 2. From the table, it is evident that the  $Mg_n^{Q} \rightarrow Mg + Mg_{n-1}^{Q}$  channel is the most popular route for all neutral and anionic magnesium clusters. For clarity, the fragmentation energies of this most probable fragmentation channel are displayed in Figure 4d as functions of cluster size *n*. We notice that, except for n = 8, the two curves of fragmentation energies show quantitatively very similar behavior across the cluster sizes n = 3-18. Several energy maxima found at n = 4, 9, 15 and 17 suggest that initial clusters  $Mg_4^{0/-}$ ,  $Mg_9^{0/-}$ ,  $Mg_{15}^{0/-}$  and  $Mg_{17}^{0/-}$  are more stable than their neighbors and could be difficult to dissociate.

#### E. Chemical bonding analysis

From a comprehensive analysis of the electronic properties of the magnesium clusters, it is clear that the neutral Mg<sub>17</sub> cluster is relatively stable, which caught our interest. The neutral cluster Mg<sub>17</sub>, a filled-cage-like structure, has  $D_{4d}$  symmetry, four peripheral Mg<sub>4</sub> fragments and one central Mg atom. To gain insight into its bonding properties, we analysed the molecular orbitals (MO) of the neutral square Mg<sub>4</sub> unit (Figure S3a), which, as isolated Mg<sub>4</sub> cluster, is a metastable species. The highest occupied MO (HOMO) of Mg<sub>4</sub> (Figure S3b) is a completely delocalized  $\pi$  orbital. There are also four lone-pair MOs and two pairs of  $\sigma$  electrons can be shared by four Mg–Mg bonds, resulting in a  $\sigma$  bond order of 0.5. Both the  $\pi$  bonding patterns and the lone-pair MOs of Mg<sub>4</sub> are analogous to the previously discovered aromatic Al<sub>4</sub><sup>2-.35</sup> In principle the system satisfied the Hückel rule for a  $\pi$ -aromatic system. To provide

further evidence for a possible aromatic character, we calculated its nucleus-independent chemical shift (NICS)<sup>55,56</sup> and multicenter bond order<sup>57</sup> (see Table S3). NICS(0)<sub>zz</sub>, NICS(0.5)<sub>zz</sub> and NICS(1)<sub>zz</sub> are -37.1, -36.6, and -34.0 ppm, respectively. NICS<sub>zz</sub> is identical to the z-component of the induced magnetic field (Merino, G.; Heine, T.; Seifert, G. The Induced Magnetic Field in Cyclic Molecules. Chem. Eur. J. 2004, 10, 4367-4371 AND Heine, T. Islas, R.; Merino, G.  $\sigma$  and  $\pi$  Contributions to the Induced Magnetic: Indicators for the Mobility of Electrons in Molecules. J. Comput. Chem. 2007, 28, 302-309 AND 58. Islas, R.; Heine, T.; Merino, G. The Induced Magnetic Field. Acc. Chem. Res. 2012, 45, 215-228). All of the large negative NICS indices suggested that Mg<sub>4</sub> cluster can be classified as aromatic. In addition, the total value of multicenter bond order is 0.0322, thus reconfirming the aromatic character of the Mg<sub>4</sub> cluster.

The ground state structure of neutral Mg<sub>17</sub> found here contains two square Mg<sub>4</sub> frameworks. In order to improve our understanding of the bonding in the Mg<sub>17</sub> cluster, we performed chemical bonding analyses using the AdNDP method. The detailed results are displayed in Figure 5. The 34 valence electrons in neutral Mg<sub>17</sub> cluster can be divided into two sets. The first set consists of delocalized  $\sigma$ -bond elements, while the other set is composed of delocalized  $\pi$ -bond elements. In the first set, the AdNDP analysis revealed eight 4c–2e  $\sigma$ -bonds with ON = 1.81 |e|, which are responsible for the bonding within the eight peripheral cyclo-Mg<sub>4</sub> units, and eight 4c–2e  $\sigma$ -bonds with ON = 1.71 |e|, which are likely due to the bonding between the central Mg with peripheral Mg<sub>3</sub> units. The second set includes two 8c–2e  $\pi$ -bonds, which are best

described as two completely delocalized  $\pi$  bonds with 4 electrons in total and embody the strong bond among eight peripheral Mg atoms, and one 9c–2e  $\pi$ -bond with an occupation number (ON) of 1.77 |e|, which has an interesting sandwich shape and creates strains in the interior between the central Mg atom and the two apical Mg<sub>4</sub> units. The 9c–2e  $\pi$  bonds on two Mg<sub>4</sub> units are quite similar to the 4c–2e bonds in the neutral square planar Mg<sub>4</sub> cluster (Figure S3c), which in turn are similar to the Sb<sub>4</sub> fragments of 5c–2e  $\pi$ -bonds in [Ln( $\eta^4$ -Sb<sub>4</sub>)<sub>3</sub>]<sup>3–</sup>complex<sup>40</sup>, rendering it locally  $\pi$ -aromatic according to Hückel's 4N + 2 rule. Why should such peculiar Mg<sub>4</sub> fragments occur in the Mg<sub>17</sub> cluster? They are stabilised in Mg<sub>17</sub> due to the strong interactions between the Mg<sub>4</sub> units and the central Mg atom and more importantly due to the presence of local aromaticity. Therefore, the neutral Mg<sub>17</sub> cluster can be considered to be  $\pi$ -aromatic.

Moreover, Mg<sub>17</sub> cluster is a 34 electron system, which possesses a closed electronic shell according to the jellium superatom model.[cite: JPCC 113, 2664 (2009)] For a detailed illustration of the electronic structure of the Mg<sub>17</sub> cluster, we plotted its molecular orbitals (Figure S4). Combining the multicentre bonding and molecular orbitals analyses on the  $\pi$ -aromatic Mg<sub>17</sub> cluster indicates that the electronic configuration of Mg<sub>17</sub> can be labeled as  $1S^21P^61D^42S^21D^61F^{14}$ .

### **IV. Conclusions**

In summary, we have studied the structural evolution and electronic properties of neutral and anionic Mg clusters having up to 20 atoms using the unbiased CALYPSO

structure searching method combined with density functional theory calculations. The ground state structures of both neutral and corresponding anionic Mg clusters have similar geometric structures and show the same structural evolution with increasing atom number. The simulated photoelectron spectra are in excellent agreement with experimental measurements. On the basis of fragmentation energy calculations, we identified that the Mg<sub>n</sub><sup>Q</sup>  $\rightarrow$  Mg<sub>n-1</sub><sup>Q</sup> + Mg channel is the route favored by all the neutral and anion Mg<sub>n</sub> clusters (n = 3-20), but is endothermic in all cases. A detailed chemical bonding analysis reveals that the neutral Mg<sub>17</sub> cluster is the first locally  $\pi$ -aromatic cluster among medium-sized homonuclear metal clusters. To some extent, we expect that this result may provide a new twist on the concept of aromaticity and stimulate theoretical analyses of the chemical bonding in other known or as yet unknown homonuclear metal clusters.

## Acknowledgments

This work was supported by the National Natural Science Foundation of China (No. 11274235, 11304167 and 11574220), the 973 Program of China (2014CB660804), the Special Program for Applied Research on Super Computation of the NSFC-Guangdong Joint Fund (the second phase) and the Program for Science & Technology Innovation Talents in Universities of Henan Province (No. 15HASTIT020). The work in México was supported by Conacyt (Grant CB-2015-252356).

# References

- Monteverde, M.; Núňez-Regueiro, M.; Rogado, N.; Regan, K. A.; Hayward,M. A.; He, T.; Loureiro, S. M.; Cava, R. J. Pressure Dependence of the Superconducting Transition Temperature of Magnesium Diboride. *Science* 2001, 292, 75–77.
- Er, S.; Wijs, G. A. de; Brocks, G. Tuning the Hydrogen Storage in Magnesium Alloys. J. Phys. Chem. Lett. 2010, 1, 1982–1986.
- Nevshupa, R.; Ares, J. R.; Fernández, J. F.; del, Campo, A.; Roman, E. Tribochemical Decomposition of Light Ionic Hydrides at Room Temperature. J. Phys. Chem. Lett. 2015, 6, 2780–2785.
- Chen, L. Y.; Xu, J. Q.; Choi, H.; Pozuelo, M.; Ma, X. L.; Bhowmick, S.; Yang, J. M.; Mathaudhu, S.; Li, X. C. Processing and Properties of Magnesium Containing a Dense Uniform Dispersion of Nanoparticles. *Nature* 2015, 528, 539–545.
- Barcaro, G.; Ferrando, R.; Fortunelli, A.; Rossi, G. Exotic Supported CoPt Nanostructures: From Clusters to Wires. J. Phys. Chem. Lett. 2010, 1, 1111–1115.
- Stephenson, A. E.; DeYoreo, J. J.; Wu, L.; Wu, K. J.; Hoyer, J.; Dove, P. M. Peptides Enhance Magnesium Signature in Calcite: Insights into Origins of Vital Effects. *Science* 2008, 322, 724–727.
- Yoo, J.; Aksimentiev, A. Improved Parametrization of Li<sup>+</sup>, Na<sup>+</sup>, K<sup>+</sup>, and Mg<sup>2+</sup> Ions for All-Atom Molecular Dynamics Simulations of Nucleic Acid Systems. J.

Phys. Chem. Lett. 2012, 3, 45–50.

- Diederich, T.; Döppner, T.; Braune, J.; Tiggesbäumker, J.; Meiwes-Broer, K. Electron Delocalization in Magnesium Clusters Grown in Supercold Helium Droplets. *Phys. Rev. Lett.* 2001, 86, 4807–4810.
- Thomas, O. C.; Zheng, W. J.; Xu, S. J.; Bowen, K. H. Jr. Onset of Metallic Behavior in Magnesium Clusters. *Phys. Rev. Lett.* 2002, 89, 213403.
- Kumar, V.; Car, R. Structure, Growth, and Bonding Nature of Mg Clusters. *Phys. Rev. B* 1991, 44, 8243–8255.
- Jellinek, J.; Acioli, P. H. Magnesium Clusters: Structural and Electronic Properties and the Size-Induced Nonmetal-to-Metal Transition. *J. Phys. Chem. A* 2002, 106, 10919–10925.
- Acioli, P. H.; Jellinek, J. Electron Binding Energies of Anionic Magnesium Clusters and the Nonmetal-to-Metal Transition. *Phys. Rev. Lett.* 2002, 89, 213402.
- Köhn, A.; Weigend, F.; Ahlrichs, R. Theoretical Study on Clusters of Magnesium. *Phys. Chem. Chem. Phys.* 2001, 3,711–719.
- Reuse, F.; Khanna, S. N.; Coulon, V. de; Buttet, J. Behavior of Magnesium Clusters Under Electron Attachment and Detachment. *Phys. Rev. B* 1989, 39, 12911–12914.
- 15. Reuse, F.; Khanna, S. N.; Coulon, V. de; Buttet, J. Pseudopotential Local-Spin-Density Studies of Neutral and Charged Mg<sub>n</sub> ( $n \le 7$ ) Clusters. *Phys. Rev. B* **1990**, 41, 11743–11759.

- Delaly, P.; Ballone, P. Buttet, J. Metallic Bonding in Magnesium Microclusters. *Phys. Rev. B* 1992, 45, 3838–3841.
- Davidson, E. R.; Frey, R. F. Density Functional Calculations for Mg<sup>+</sup> Clusters. J. Chem. Phys. 1997, 106, 2331–2341.
- Eriksson, L. A. Accurate Density Functional Theory Study of Cationic Magnesium Clusters and Mg<sup>+</sup>-Rare Gas Interactions. J. Chem. Phys. 1995, 103, 1050–1056.
- Reimann, S. M.; Koskinen, M.; Häkkinen, H.; Lindelof, P. E.; Manninen, M. Magic Triangular and Tetrahedral Clusters. *Phys. Rev. B* 1997, 56, 12147–12150.
- Lyalin, A.; Solov'yov, I. A.; Solov'yov, A. V.; Greiner, W. Evolution of the Electronic and Ionic Structure of Mg Clusters with Increase in Cluster Size. *Phys. Rev. A* 2003, 67, 063203.
- De, S.; Ghasemi, S. A.; Willand, A.; Genovese, L.; Kanhere, D.; Goedecker, S. The Effect of Ionization on the Global Minima of Small and Medium Sized Silicon and Magnesium Clusters. *J. Chem. Phys.* 2011, 134, 124302.
- Heidari, I.; De, S.; Ghazi, S. M.; Goedecker, S.; Kanhere, D. G. Growth and Structural Properties of Mg<sub>N</sub> (N = 10–56) Clusters: Density Functional Theory Study. J. Phys. Chem. A 2011, 115, 12307–12314.
- Hückel, E. Quantentheoretische Beiträge zum Benzolproblem. Z. Phys. 1931, 70, 204–286.
- 24. Hückel, E. Quantentheoretische Beiträge zum Problem der aromatischen und ungesattigten Verbindungen. III. Z. Phys. **1932**, 76, 628–648.

- 25. Minkin, V. I.; Glukhovtsev, M. N.; Simkin, B. Ya. Aromaticity and antiaromaticity: electronic and structural aspects. Wiley, New York, 1994.
- Zhang, A.; Han, Y. H.; Yamato, K.; Zeng, X. C.; Gong, B. Aromatic Oligoureas: Enforced Folding and Assisted Cyclization. *Org. Lett.* 2006, 8, 803–806.
- 27. Jiang, J. Y.; Lima, O. V.; Pei, Y.; Zeng, X. C.; Tan, L.; Forsythe, E. Dipole-Induced, Thermally Stable Lamellar Structure by Polar Aromatic Silane. *J. Am. Chem. Soc.* 2009, 131, 900–901.
- Yang,Y. A.; Feng, W.; Hu, J. C.; Zou, S. L.; Gao, R. Z.; Yamato, K.; Kline, M.; Cai, Z. H.; Gao, Y.; Wang, Y. B.; Li, Y.; B.; Yang, Y. L.; Yuan, L. H.; Zeng, X. C.; Gong, B. Strong Aggregation and Directional Assembly of Aromatic Oligoamide Macrocycles. *J. Am. Chem. Soc.* **2011**, 133, 18590–18593.
- 29. Boldyrev, A. I.; Wang, L. S. All-Metal Aromaticity and Antiaromaticity. *Chem. Rev.* 2005, 105, 3716–3757.
- Robinson. G. H. Gallanes, Gallenes, Cyclogallenes, and Gallynes: Organometallic Chemistry about the Gallium-Gallium Bond. *Accounts Chem. Res.* 1999, 32, 773–782.
- Romanescu, C.; Galeev, T. R.; Li, W. L.; Boldyrev, A. I.; Wang, L. S. Transition-Metal-Centered Monocyclic Boron Wheel Clusters (M@B<sub>n</sub>): A New Class of Aromatic Borometallic Compounds. *Accounts Chem. Res.* 2013, 46, 350–358.
- 32. Li, W. L.; Jian, T.; Chen, X.; Chen, T. T.; Lopez, G. V.; Li, J.; Wang, L. S. The Planar CoB<sub>18</sub><sup>-</sup> Cluster as a Motif for Metallo-Borophenes. *Angew. Chem. Int. Edit.*

**2016,** 128, 1–6.

- Zubarev, D. Yu.; Averkiev, B. B.; Zhai, H. J.; Wang, L. S.; Boldyrev, A. I. Aromaticity and Antiaromaticity in Transition-Metal Systems. *Phys. Chem. Chem. Phys.* 2008, 10, 257–267.
- Galeev, T. R.; Boldyrev, A. I. Recent Advances in Aromaticity and Antiaromaticity in Transition-Metal Systems. *Annu. Rep. Prog. Chem. Sect. C* 2011, 107, 124–147.
- 35. Li, X.; Kuznetsov, A. E.; Zhang, H. F.; Boldyrev, A. I.; Wang, L. S. Observation of All-Metal Aromatic Molecules. *Science* **2001**, 291, 859–861.
- Alexandrova, A. N.; Boldyrev, A. I. σ-Aromaticity and σ-Antiaromaticity in Alkali Metal and Alkaline Earth Metal Small Clusters. J. Phys. Chem. A 2003, 107, 554–560.
- Mercero, J. M.; Boldyrev, A. I.; Merino, G.; Ugalde, J. M. Recent Developments and Future Prospects of All-Metal Aromatic Compounds. *Chem. Soc. Rev.* 2015, 44, 6519–6534.
- 38. Breslow, R. Antiaromaticity. Accounts Chem. Res. 1973, 6, 393–398.
- Kuznetsov, A. E.; Birch, K. A.; Boldyrev, A. I.; Li, X.; Zhai, H. J.; Wang, L. S. All-Metal Antiaromatic Molecule: Rectangular Al<sub>4</sub><sup>4-</sup> in the Li<sub>3</sub>Al<sub>4</sub><sup>-</sup> Anion. *Science* 2003, 300, 622–625.
- Min, X.; Popov, I. A.; Pan, F. X.; Li, L. J.; Matito, E.; Sun, Z. M.; Wang, L. S. Boldyrev, A. I. All-Metal Antiaromaticity in Sb<sub>4</sub>-Type Lanthanocene Anions. *Angew. Chem. Int. Ed.* 2016,55,5531–5535.

- 41. Wang, Y. C.; Lv, J.; Zhu, L.; Ma, Y. M. Crystal Structure Prediction via Particle-Swarm Optimization. *Phys. Rev. B* 2010, 82, 094116.
- 42. Wang, Y. C; Lv, J.; Zhu, L.; Ma, Y. M. CALYPSO: A Method for Crystal Structure Prediction. *Comput. Phys. Commun.* **2012**, 183, 2063–2070.
- 43. Wang, Y. C; Miao, M. S; Lv, J.; Zhu, L.; Yin, K.; Liu, H. Y.; Ma, Y. M. An Effective Structure Prediction Method for Layered Materials Based on 2D Particle Swarm Optimization Algorithm. J. Chem. Phys. 2012, 137, 224108.
- 44. Wang, Y. C; Ma, Y. M. Perspective: Crystal Structure Prediction at High Pressures. J. Chem. Phys. 2014, 140, 040901.
- 45. Lv, J.; Wang, Y. C.; Zhu, L.; Ma. Y. M. Particle-Swarm Structure Prediction on Clusters. J. Chem. Phys. 2012, 137, 084104.
- 46. Lu, S. H.; Wang, Y. C.; Liu, H. Y.; Miao, M. S.; Ma, Y. M. Self-Assembled Ultrathin Nanotubes on Diamond (100) Surface. *Nat. Commun.* **2014**, *5*, 3666.
- 47. Zhu, L.; Liu, H. Y; Pickard, C. J.; Zou , G. T.; Ma, Y. M. Reactions of Xenon With Iron and Nickel are Predicted in the Earth's Inner Core. *Nat. Chem.* 2014, 6, 644–648.
- 48. Becke, A. D. Density-Functional Thermochemistry. III. The Role of Exact Exchange. J. Chem. Phys. 1993, 98, 5648–5652.
- Perdew, J. P.; Wang, Y. Pair-distribution Function and its Coupling-Constant Average for the Spin-Polarized Electron Gas. *Phys. Rev. B* 1992, 46, 12947–12954.
- 50. Perdew, J. P.; Ziesche, P.; Eschrig, H. Electronic Structure of Solids' 91.

Akademie Verlag, Berlin, 1991.

- 51. Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Montgomery, J. A. Jr, Vreven, T.; Kudin, K. N.; Burant, J.C.; Millam, J. M.; Iyengar, S. S.; Tomasi, J.; Barone, V.; Mennucci, B.; Cossi, M.; et al. *Gaussian 09*, revision C.0; Gaussian, Inc.: Wallingford, CT, 2009.
- Zubarev, D. Y.; Boldyrev, A. I. Developing Paradigms of Chemical Bonding: Adaptive Natural Density Partitioning. *Phys. Chem. Chem. Phys.* 2008, 10, 5207–5217.
- Lu, T.; Chen, F. Multiwfn: A Multifunctional Wavefunction Analyzer. J. Comput. Chem. 2012, 33, 580–592.
- Pears, R. G. The Principle of Maximum Hardness. Accounts Chem. Res. 1993, 26, 250–255.
- Schleyer, P. v. R.; Maerker, C.; Dransfeld, A.; Jiao, H. J.; Hommes, N. J. R. v. E.; Nucleus-Independent Chemical Shifts: A Simple and Efficient Aromaticity Probe. *J. Am. Chem. Soc.* 1996, 118, 6317–6318.
- 56. West, R.; Buffy, J. J.; Haaf, M.; Müller, T.; Gehrhus, B.; Lappert, M. F.; Apeloig,
  Y. Chemical Shift Tensors and NICS Calculations for Stable Silylenes. *J. Am. Chem. Soc.* 1998, 120, 1639–1640.
- Giambiagi, M.; Giambiagi, M. S. de; Mundim, K. C. Definition of a Multicenter Bond Index. *Struct. Chem.* 1990, 1, 423–427.

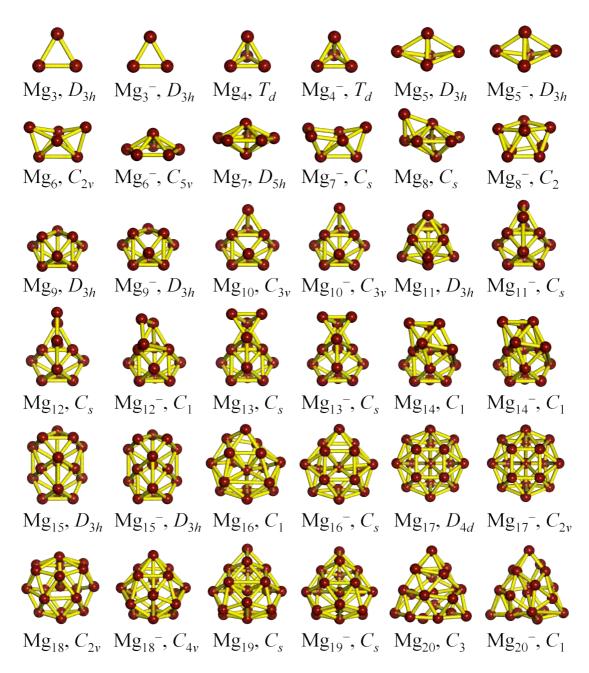


Figure 1. The lowest energy structures of  $Mg_n^Q$  (n = 3-20, Q = 0, -1) clusters.

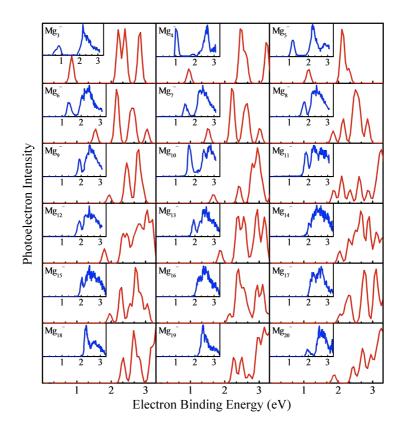


Figure 2. Comparison of the simulated photoelectron spectra (outer) with the experimental PES (inset) from reference 9 of lowest energy  $Mg_n^-$  (n = 3-20) clusters.

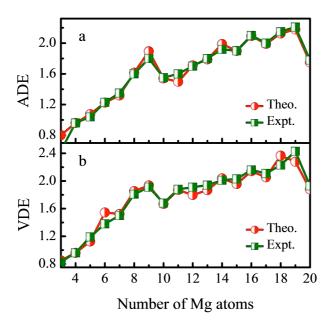


Figure 3. Adiabatic detachment energies (ADEs) and vertical detachment energies (VDEs) of  $Mg_n^-$  (n = 3-20) clusters: red circles, theory; green squares, experiment.

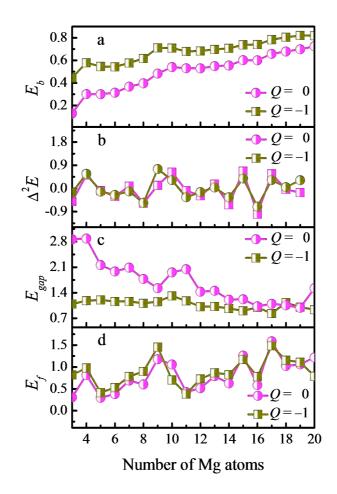


Figure 4. The averaged binding energies  $E_b$  (a), second-order energy differences  $\Delta^2 E$ (b), HOMO–LUMO energy gaps  $E_{gap}$  (c) and fragmentation energies of the most probable fragmentation channels (d) for lowest energy  $Mg_n^Q$  (n = 3-20, Q = 0, -1) clusters as functions of cluster size n.

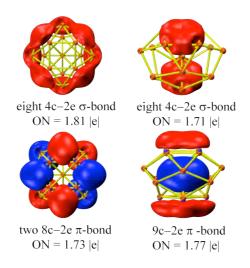


Figure 5. Chemical bonding analyses of neutral Mg<sub>17</sub> cluster using the AdNDP

method. ON stands for occupation number.

Table 1. The calculated binding energies ( $E_b$ ) and HOMO–LUMO energy gaps  $E_{gap}$  of the lowest-energy Mg<sub>n</sub><sup>Q</sup> (n = 3-20, Q = 0, -1) clusters. And the data of theoretical and experimental vertical detachment energies (VDEs) and adiabatic detachment energies (ADEs) for the lowest-energy Mg<sub>n</sub><sup>-</sup> (n = 3-20) clusters. All of energies are shown in the unit of eV.

n	Mg <sub>n</sub>		Mg <sub>n</sub>		VDE		ADE	
	$E_b$	$E_{gap}$	$E_b$	$E_{gap}$	Theo.	Expt. <sup>a</sup>	Theo.	Expt. <sup>a</sup>
3	0.13	2.86	0.45	1.09	0.85	0.83	0.80	0.60
4	0.30	2.89	0.58	1.18	0.96	0.96	0.96	0.96
5	0.30	2.16	0.55	1.20	1.13	1.19	1.07	1.04
6	0.31	1.98	0.54	1.16	1.54	1.38	1.23	1.23
7	0.37	2.09	0.58	1.16	1.52	1.50	1.32	1.35
8	0.40	1.78	0.62	1.11	1.85	1.81	1.61	1.60
9	0.48	1.53	0.71	1.15	1.94	1.91	1.89	1.80
10	0.54	1.96	0.71	1.31	1.67	1.68	1.54	1.55
11	0.53	2.04	0.68	1.17	1.87	1.88	1.50	1.60
12	0.53	1.43	0.68	1.02	1.80	1.91	1.71	1.70
13	0.55	1.45	0.70	1.02	1.87	1.94	1.79	1.80
14	0.55	1.21	0.71	0.97	2.04	2.01	1.99	1.92
15	0.60	1.22	0.74	0.90	1.96	2.03	1.90	1.90
16	0.60	1.03	0.74	1.00	2.14	2.16	2.09	2.10
17	0.66	1.10	0.78	0.83	2.05	2.11	1.99	2.00
18	0.68	1.06	0.81	1.13	2.37	2.23	2.12	2.15
19	0.70	0.99	0.82	1.00	2.28	2.43	2.18	2.21
20	0.72	1.52	0.82	0.94	1.88	1.93	1.75	1.78

<sup>*a*</sup>Refernce 9

Cluster	$Mg_p + Mg_{n-p}$	$E_f(eV)$	Cluster	$Mg_p + Mg_{n-p}$	$E_f(eV)$
Mg <sub>3</sub>	$Mg + Mg_2$	0.31	Mg <sub>3</sub> <sup>-</sup>	$Mg + Mg_2^-$	0.82
Mg <sub>4</sub>	$Mg + Mg_3$	0.82	Mg <sub>4</sub> <sup>-</sup>	$Mg + Mg_3^-$	0.98
Mg <sub>5</sub>	$Mg + Mg_4$	0.30	Mg <sub>5</sub> <sup>-</sup>	$Mg + Mg_4^-$	0.41
Mg <sub>6</sub>	$Mg + Mg_5$	0.38	Mg <sub>6</sub>	$Mg + Mg_5$	0.53
Mg <sub>7</sub>	$Mg + Mg_6$	0.69	Mg <sub>7</sub> <sup>-</sup>	$Mg + Mg_6^-$	0.78
Mg <sub>8</sub>	$Mg + Mg_7$	0.60	Mg8	$Mg + Mg_7$	0.90
Mg9	$Mg + Mg_8$	1.18	Mg9 <sup>-</sup>	$Mg + Mg8^{-}$	1.46
Mg <sub>10</sub>	$Mg + Mg_9$	1.06	$Mg_{10}$	$Mg + Mgg^{-}$	0.70
Mg <sub>11</sub>	$Mg + Mg_{10}$	0.43	Mg <sub>11</sub> <sup>-</sup>	$Mg + Mg_{10}$	0.38
Mg <sub>12</sub>	$Mg + Mg_{11}$	0.51	$Mg_{12}^{-}$	$Mg + Mg_{11}^{-}$	0.73
Mg <sub>13</sub>	$Mg + Mg_{12}$	0.79	Mg <sub>13</sub> <sup>-</sup>	$Mg + Mg_{12}^{-}$	0.87
Mg <sub>14</sub>	$Mg + Mg_{13}$	0.63	Mg <sub>14</sub>	$Mg + Mg_{13}$	0.83
Mg <sub>15</sub>	$Mg + Mg_{14}$	1.26	Mg <sub>15</sub> <sup>-</sup>	$Mg + Mg_{14}^{-}$	1.17
Mg <sub>16</sub>	$Mg + Mg_{15}$	0.59	Mg <sub>16</sub>	$Mg + Mg_{15}$	0.78
Mg <sub>17</sub>	$Mg + Mg_{16}$	1.59	$Mg_{17}^{-}$	$Mg + Mg_{16}$	1.48
Mg <sub>18</sub>	$Mg + Mg_{17}$	1.02	$Mg_{18}^{-}$	$Mg + Mg_{17}^{-}$	1.16
Mg <sub>19</sub>	$Mg + Mg_{18}$	1.06	Mg <sub>19</sub> <sup>-</sup>	$Mg + Mg_{18}$	1.11
Mg <sub>20</sub>	$Mg + Mg_{19}$	1.22	$Mg_{20}^{-}$	$Mg + Mg_{19}^{-}$	0.79

Table 2. The easiest fragmentation channels and corresponding fragmentation energies (*E<sub>f</sub>*) for Mg<sub>n</sub><sup>Q</sup> (n = 3-20, Q = 0, -1) clusters.