

Numerical Solution of Coupled System of Nonlinear Partial Differential Equations Using Laplace-Adomian Decomposition Method

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ABSTRACT

Aim of the paper is to investigate applications of Laplace Adomian Decomposition Method (LADM) on nonlinear physical problems. Some coupled system of non-linear partial differential equations (NLPDEs) are considered and solved numerically using LADM. The results obtained by LADM are compared with those obtained by standard and modified Adomian Decomposition Methods. The behavior of the numerical solution is shown through graphs. It is observed that LADM is an effective method with high accuracy with less number of components.

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1. INTRODUCTION

The partial differential equations (PDEs) have so many essential applications of science and engineering such as wave propagation, shallow water waves, fluid mechanics, thermodynamic, chemistry and micro electro mechanic system, etc. It is difficult to handle nonlinear part of these systems, although most of the scientists applied numerical methods to find the solutions of these systems that based on linearization, perturbed and discretizations. Debnath [1] applied the characteristics method and Logan [2] used the Riemann invariants method to handle systems of PDEs. Wazwaz [3] used the Adomian decomposition method (ADM) to handle the systems of PDEs. Laplace Decomposition Method (LDM) is free of any small or large parameters and has advantages over other approximation techniques like perturbation, LDM requires no discretization and linearization, therefore, results obtained by LDM are more efficient and realistic. This method has been used to obtain approximate solutions of a class of nonlinear ordinary and PDEs [4-7]. In this paper, we compute numerical solutions to nonlinear systems of PDEs by using LADM. The numerical solutions become easier and higher accuracy than the standard Adomian Decomposition Method (ADM).

2. LADM for NONLINEAR COUPLED of PDEs

Consider the general nonlinear coupled of PDEs written in an operators form (see [8])

$$\begin{cases} L_t(u) + R_1(u) + M_1(u) + N_1(u, v) = f_1(x, t) \\ L_t(v) + R_2(v) + M_2(v) + N_2(u, v) = f_2(x, t) \end{cases} \quad (1)$$

Subject to the initial conditions

$$\begin{cases} u(x, 0) = g_1(x), 0 \leq x \leq \ell_1, \\ v(x, 0) = g_2(x), 0 \leq x \leq \ell_2, \end{cases} \quad (2)$$

where ℓ_1 and ℓ_2 are real constants and the notations of $L_t = \frac{\partial}{\partial t}$, R_1 and R_2 symbolized the linear spatial differential operators, the notations M_1, M_2, N_1 and N_2 symbolized the nonlinear differential operators and $f_1(x, t), f_2(x, t)$ are given functions. The method consists of first applying the Laplace transform to both sides of equations in system (1) and then by using initial conditions, we have

$$\begin{cases} \mathcal{L}\{u\} = \frac{g_1(x)}{s} + \frac{1}{s} \mathcal{L}\{f_1(x, t)\} - \frac{1}{s} [\mathcal{L}\{R_1(u)\} + \mathcal{L}\{M_1(u)\} + \mathcal{L}\{N_1(u, v)\}] \\ \mathcal{L}\{v\} = \frac{g_2(x)}{s} + \frac{1}{s} \mathcal{L}\{f_2(x, t)\} - \frac{1}{s} [\mathcal{L}\{R_2(v)\} + \mathcal{L}\{M_2(v)\} + \mathcal{L}\{N_2(u, v)\}] \end{cases} \quad (3)$$

in the Laplace decomposition method we assume the solution is in an infinite series, given as follows



$$\begin{cases} u(x,t) = \sum_{k=0}^{\infty} u_k(x,t), \\ v(x,t) = \sum_{k=0}^{\infty} v_k(x,t), \end{cases} \quad (4)$$

where $u_k(x,t)$ and $v_k(x,t)$ are to be recursively computed. Also the nonlinear terms M_1, M_2, N_1 and N_2 are decomposed as an infinite series of Adomian polynomials (see [9,10]),

$$M_1(u) = \sum_{k=0}^{\infty} A_k, N_1(u,v) = \sum_{k=0}^{\infty} B_k, \quad (5)$$

For the nonlinear operators $M_1(u)$, and $N_1(u,v)$ the Adomian's polynomials can be generated for all forms of nonlinearity. They are defined by the following relations

$$\begin{aligned} A_n(u_0, u_0, \dots, u_n) &= \frac{1}{n!} \frac{d^n}{d\lambda^n} \left[M_1 \left(\sum_{k=0}^n \lambda^k u_k(x,t) \right) \right]_{\lambda=0}, \\ B_n(u_0, u_0, \dots, u_n; v_0, v_0, \dots, v_n) &= \frac{1}{n!} \frac{d^n}{d\lambda^n} \left[N_1 \left(\sum_{k=0}^n \lambda^k u_k(x,t), \sum_{k=0}^n \lambda^k v_k(x,t) \right) \right]_{\lambda=0}. \end{aligned} \quad (6)$$

Substituting (4) and (6) into (3) and applying the linearity of the Laplace transform, we get the following recursively formula

$$\begin{cases} \mathcal{L}\{u_0(x,t)\} = \frac{g_1(x)}{s} + \frac{1}{s} \mathcal{L}\{f_1(x,t)\} \\ \mathcal{L}\{v_0(x,t)\} = \frac{g_2(x)}{s} + \frac{1}{s} \mathcal{L}\{f_2(x,t)\} \end{cases} \quad (7)$$

and

$$\begin{cases} \mathcal{L}\{u_{k+1}(x,t)\} = -\frac{1}{s} \mathcal{L}\{R_1(u_k(x,t)) + A_k + B_k\}, \quad k \geq 0 \\ \mathcal{L}\{v_{k+1}(x,t)\} = -\frac{1}{s} \mathcal{L}\{R_2(v_k(x,t)) + C_k + D_k\}, \quad k \geq 0 \end{cases} \quad (8)$$

Applying the inverse Laplace transform, we can evaluate $u_k(x,t)$ and $v_k(x,t)$

3. APPLICATIONS

In order to verify numerically whether the proposed methodology leads to the accurate solutions, we evaluate LADM using the approximation for some examples of non-linear systems of PDEs. To show the efficiency of the present methods for our problems in comparison with the exact solution, we report the absolute error. The calculations in this paper have been done using the Maple 18 software. The results are listed in tables 1-2 and figures 1-4 below.

Example 1: Consider the coupled system of nonlinear PDE of the form [11]

$$\begin{cases} p_t + v_x w_y - v_y w_x = -p \\ v_t + w_x p_y + p_x w_y = v \\ w_t + p_x v_y + p_y v_x = w \end{cases} \quad (9)$$

with the following initial conditions



$$\begin{cases} p(x, y, 0) = e^{x+y}, \\ v(x, y, 0) = e^{x-y} \\ w(x, y, 0) = e^{y-x}. \end{cases} \quad (10)$$

The exact solution of the system (9) is [11].

$$\begin{cases} p(x, y, t) = e^{x+y-t}; \\ v(x, y, t) = e^{x-y+tt}; \\ w(x, y, t) = e^{y-x+tt} \end{cases}$$

Solution: To solve the system of equations (9)-(10) by means of LADM, we construct a correctional functional which reads

$$\begin{cases} \mathfrak{F}\{p\} = \frac{1}{s} p(x, y, 0) + \frac{1}{s} \mathfrak{F}\{-p - v_x w_y + v_y w_x\} \\ \mathfrak{F}\{v\} = \frac{1}{s} v(x, y, 0) + \frac{1}{s} \mathfrak{F}\{v - w_x p_y - p_x w_y\} \\ \mathfrak{F}\{w\} = \frac{1}{s} w(x, y, 0) + \frac{1}{s} \mathfrak{F}\{w - p_x v_y - p_y v_x\} \end{cases} \quad (11)$$

We can define the Adomian polynomials as follows

$$\begin{cases} A_n = \sum_{k=0}^n \left(\frac{\partial}{\partial x} v_{n-k}(x, y, t) \right) \left(\frac{\partial}{\partial y} w_{n-k}(x, y, t) \right), \\ B_n = \sum_{k=0}^n \left(\frac{\partial}{\partial y} v_{n-k}(x, y, t) \right) \left(\frac{\partial}{\partial x} w_{n-k}(x, y, t) \right), \\ C_n = \sum_{k=0}^n \left(\frac{\partial}{\partial x} w_{n-k}(x, y, t) \right) \left(\frac{\partial}{\partial y} p_{n-k}(x, y, t) \right), \\ E_n = \sum_{k=0}^n \left(\frac{\partial}{\partial x} p_{n-k}(x, y, t) \right) \left(\frac{\partial}{\partial y} w_{n-k}(x, y, t) \right), \\ F_n = \sum_{k=0}^n \left(\frac{\partial}{\partial x} p_{n-k}(x, y, t) \right) \left(\frac{\partial}{\partial y} v_{n-k}(x, y, t) \right), \\ G_n = \sum_{k=0}^n \left(\frac{\partial}{\partial y} p_{n-k}(x, y, t) \right) \left(\frac{\partial}{\partial x} v_{n-k}(x, y, t) \right). \end{cases} \quad (12)$$

We define an iterative scheme



$$\left\{ \begin{array}{l} \mathfrak{L}\{p_0(x, y, t)\} = \frac{1}{s} e^{x+y} \\ \mathfrak{L}\{v_0(x, y, t)\} = \frac{1}{s} e^{x-y} \\ \mathfrak{L}\{w_0(x, y, t)\} = \frac{1}{s} e^{y-x} \\ \mathfrak{L}\{p_{n+1}(x, y, t)\} = \frac{1}{s} \mathfrak{L}\{-p_n(x, y, t) - A_n + B_n\}, \quad n \geq 0 \\ \mathfrak{L}\{v_{n+1}(x, y, t)\} = \frac{1}{s} \mathfrak{L}\{v_n(x, y, t) - C_n - E_n\}, \quad n \geq 0 \\ \mathfrak{L}\{w_{n+1}(x, y, t)\} = \frac{1}{s} \mathfrak{L}\{w_n(x, y, t) - F_n - G_n\}, \quad n \geq 0 \end{array} \right. \quad (13)$$

the Adomian's polynomials A_n, B_n, C_n, E_n, F_n and G_n are generated according to (6), we can give the first few Adomian's polynomials as follows

$$A_0 = 1, A_1 = t^2 + 1, A_2 = \frac{1}{4}t^4 + t^2 + 1, A_3 = \frac{1}{36}t^6 + \frac{1}{4}t^4 + t^2 + 1$$

$$B_0 = 1, B_1 = t^2 + 1, B_2 = \frac{1}{4}t^4 + t^2 + 1, B_3 = \frac{1}{36}t^6 + \frac{1}{4}t^4 + t^2 + 1$$

$$C_0 = -e^{2y}, C_1 = e^{2y}(t^2 - 1), C_2 = -\frac{1}{4}e^{2y}(t^4 - 4t^2 + 4),$$

$$C_3 = \frac{1}{36}e^{2y}(t^6 - 9t^4 + 36t^2 - 36), E_0 = e^{2y}, E_1 = -e^{2y}(t^2 - 1),$$

$$E_2 = \frac{1}{4}e^{2y}(t^4 - 4t^2 + 4), E_3 = -\frac{1}{36}e^{2y}(t^6 - 9t^4 + 36t^2 - 36),$$

$$F_0 = -e^{2x}, F_1 = e^{2x}(t^2 - 1), F_2 = \frac{1}{4}e^{2y}(t^4 - 4t^2 + 4),$$

$$F_3 = \frac{1}{36}e^{2x}(t^6 - 9t^4 + 36t^2 - 36), G_0 = e^{2x}, G_1 = -e^{2x}(t^2 - 1),$$

$$G_2 = \frac{1}{4}e^{2x}(t^4 - 4t^2 + 4), G_3 = -\frac{1}{36}e^{2x}(t^6 - 9t^4 + 36t^2 - 36),$$

Applying the inverse Laplace transform, finally, According to (11) the 0^{th} components

$p_0(x, y, t) = e^{x+y}, v_0(x, y, t) = e^{x-y}$ and $w_0(x, y, t) = e^{y-x}$ written as follows:

$p_0(x, y, t) = e^{x+y}, v_0(x, y, t) = e^{x-y}, w_0(x, y, t) = e^{y-x}$, where, $k \geq 0$. So, we get the following components:

$$p_1(x, y, t) = e^{x+y}t, v_1(x, y, t) = e^{x-y}t, w_1(x, y, t) = e^{y-x}t,$$

$$p_2(x, y, t) = \frac{1}{2}e^{x+y}t^2, v_2(x, y, t) = \frac{1}{2}e^{x-y}t^2, w_2(x, y, t) = \frac{1}{2}e^{y-x}t^2,$$

$$p_3(x, y, t) = -\frac{1}{6}e^{x+y}t^3, v_3(x, y, t) = \frac{1}{6}e^{x-y}t^3, w_3(x, y, t) = \frac{1}{6}e^{y-x}t^3$$

$$p_{30}(x, y, t) = \frac{1}{265252859812191058636308480000000} e^{x+y} t^{30},$$

$$v_{30}(x, y, t) = \frac{1}{265252859812191058636308480000000} e^{x-y} t^{30},$$

$$w_0(x, y, t) = \frac{1}{265252859812191058636308480000000} e^{y-x} t^{30},$$

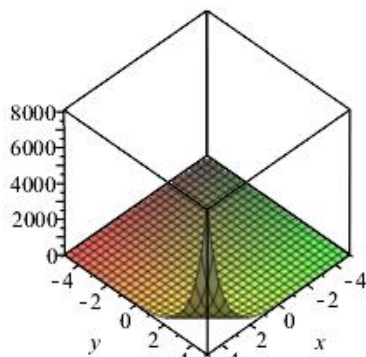
Similarly, we can also find other components, and the approximate solution for calculating 50^{th} . Using (4), the series solutions are therefore given by

$$p(x, y, t) = -\frac{1}{120} e^{x+y} (t^5 - 5t^4 + 20t^3 - 60t^2 + 120t - 120 + \dots),$$

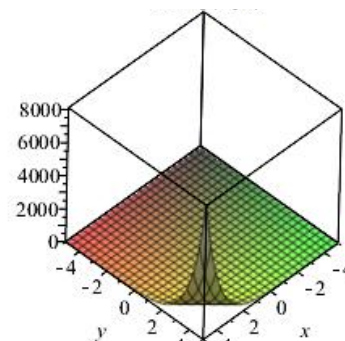
$$v(x, y, t) = \frac{1}{120} e^{x-y} (t^5 + 5t^4 + 20t^3 + 60t^2 + 120t + 120 + \dots),$$

$$w(x, y, t) = \frac{1}{120} e^{y-x} (t^5 + 5t^4 + 20t^3 + 60t^2 + 120t + 120 + \dots).$$

By using the Taylor expansion for e^t and e^{-t} we can find the exact solutions. The numerical behaviors of approximate solutions of LADM of $p(x, y, t)$, $v(x, y, t)$ and $w(x, y, t)$ with different values of time are compared with the exact solution are shown through figure1(a)-(c)

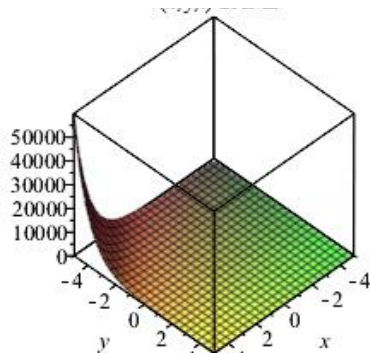


p(x,y,t) LADM

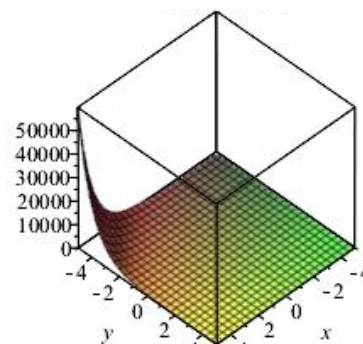


p(x,y,t) exact

(a)



v(x,y,t) LADM



v(x,y,t) exact

(b)

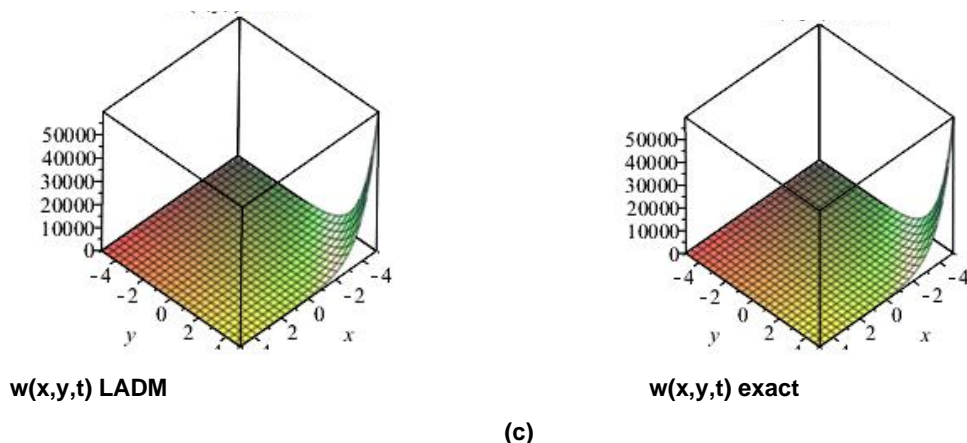


Figure 1. (a) Exact and numerical solution of $p(x,y, t)$, $-5 \leq x \leq 5$, $-5 \leq y \leq 5$ for different values of t ; (b) Exact and numerical solution of $v(x,y, t)$, $-5 \leq x \leq 5$, $-5 \leq y \leq 5$ for different values of t ; (c) Exact and numerical solution of $w(x,y, t)$, $-5 \leq x \leq 5$, $-5 \leq y \leq 5$ for different values of t

Example 2: The coupled Burgers system

$$\begin{aligned} u_t - u_{xx} - 2uu_x + (uv)_x &= 0 \\ v_t - v_{xx} - 2vv_x + (uv)_x &= 0 \end{aligned} \quad (14)$$

where $q_s = \frac{\partial}{\partial s} q(x, t)$, $q_{ss} = \frac{\partial^2}{\partial s^2} q(x, t)$, with initial conditions $u(x, 0) = \sin x$, $v(x, 0) = \sin x$. The

exact solution of the system (14) is $u(x, t) = v(x, t) = \sin x e^{-t}$ (see [11] and [12]).

Solution: To solve the system of equation (14) by means of LADM, we construct a correctional functional which reads

$$\begin{aligned} \mathfrak{L}\{u\} &= \frac{1}{s} u(x, 0) + \frac{1}{s} \mathfrak{L}\{u_{xx} + 2uu_x - (uv)_x\} \\ \mathfrak{L}\{v\} &= \frac{1}{s} v(x, 0) + \frac{1}{s} \mathfrak{L}\{v_{xx} + 2vv_x - (uv)_x\} \end{aligned} \quad (15)$$

We can define the Adomian polynomials as follows

$$\begin{cases} A_n = \sum_{k=0}^n \frac{\partial^2}{\partial x^2} u_k(x, t), & B_n = 2 \sum_{k=0}^n u_k(x, t) \frac{\partial}{\partial x} u_{n-k}(x, t), \\ C_n = \sum_{k=0}^n \frac{\partial}{\partial x} (u_{n-k}(x, t) v_{n-k}(x, t)), & F_n = \sum_{k=0}^n \frac{\partial^2}{\partial x^2} v_k(x, t) \\ E_n = 2 \sum_{k=0}^n v_k(x, t) \frac{\partial}{\partial x} v_{n-k}(x, t), \end{cases} \quad (16)$$

We define an iterative scheme

$$\begin{cases} \mathfrak{L}\{u_0(x, t)\} = \frac{1}{s} \sin x \\ \mathfrak{L}\{v_0(x, t)\} = \frac{1}{s} \sin x \\ \mathfrak{L}\{u_{n+1}(x, t)\} = \frac{1}{s} \mathfrak{L}\{A_n + B_n - C_n\}, & n \geq 0 \\ \mathfrak{L}\{v_{n+1}(x, t)\} = \frac{1}{s} \mathfrak{L}\{F_n + E_n - C_n\}, & n \geq 0. \end{cases} \quad (17)$$

Applying the inverse Laplace transform, finally, the Adomian's polynomials A_n, B_n, C_n, E_n and F_n are generated according to (6), we can give the first few components of the Adomian's polynomials respectively as follows



$$\begin{aligned}
 A_0 &= -\sin x, B_0 = 2 \sin x \cos x, C_0 = 2 \sin x \cos x, E_0 = 2 \sin x \cos x, F_0 = -\sin x, \\
 A_1 &= \sin x (t - 1), B_1 = -4t \sin x \cos x, C_1 = 2(t^2 + 1) \sin x \cos x, E_1 = -4t \sin x \cos x, \\
 F_1 &= \sin x (t - 1), A_2 = -\sin x + t \sin t - \frac{1}{6} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \\
 &+ 2t \cos x \sin x (t^2 + 3t + 3), B_2 = 2 \sin x \left(\frac{1}{6} \cos x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \right. \\
 &+ \left. \frac{2}{3} t \sin^2 x (t^2 + 3t + 3) \right) + 2t^2 \sin x \cos x + \frac{1}{3} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x \\
 &+ 3)) \cos x, C_2 = \frac{1}{3} \left(\frac{1}{6} \cos x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \right) + \frac{2}{3} t \sin^2 x (t^2 + 3t + 3) \\
 &\sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) + 2t^2 \sin x \cos x + 2 \sin x \cos x, \\
 E_2 &= 2 \sin x \left(\frac{1}{6} \cos x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \right) + \frac{2}{3} t \sin^2 x (t^2 + 3t + 3) \\
 &+ \frac{2}{3} t \sin^2 x (t^2 + 3t + 3) + 2t^2 \sin x \cos x + \frac{1}{3} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x \\
 &+ 3)) \cos x, F_2 = -\sin x + t \sin t - \frac{1}{6} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \\
 &+ 2t \cos x \sin x (t^2 + 3t + 3)
 \end{aligned}$$

According to (7), the 0^{th} components $u_0(x, t)$ and $v_0(x, t)$ written as follows: $u_0(x, t) := \sin(x)$
 $v_0(x, t) := \sin(x)$, where, $n \geq 0$. So, we get the following components:

$$\begin{aligned}
 u_1(x, t) &= -t \sin x, v_1(x, t) = -t \sin x \\
 u_2(x, t) &= \frac{1}{6} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \\
 v_2(x, t) &= \frac{1}{6} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3)) \\
 u_3(x, t) &= \frac{1}{630} t^2 \sin^3 x (1260 - 70t^4 - 84t^3 + 525t^2 + 40t((2t^4 + 14t^3 \\
 &+ 42t^2 + 63t + 42) \cos + 42)) + \frac{1}{630} (-630 - 21 \cos(3t^4 - 35t^3 - 80t^2 - 60t + 60) \\
 &- 5t^2(8 \cos^3 x (2t^4 + 14t^3 + 42t^2 + 63t + 42) + 21) + 14t(\cos^2 x (5t + 6) \\
 &(2t^3 - 15t - 30) + 45)) t \sin x \\
 v_3(x, t) &= \frac{1}{630} t^2 \sin^3 x (1260 - 70t^4 - 84t^3 + 525t^2 + 40t((2t^4 + 14t^3 \\
 &+ 42t^2 + 63t + 42) \cos + 42)) + \frac{1}{630} (-630 - 21 \cos(3t^4 - 35t^3 - 80t^2 - 60t + 60) \\
 &- 5t^2(8 \cos^3 x (2t^4 + 14t^3 + 42t^2 + 63t + 42) + 21) + 14t(\cos^2 x (5t + 6) \\
 &(2t^3 - 15t - 30) + 45)) t \sin x
 \end{aligned}$$

Similarly, we can also find other components, and the approximate solution for calculating 5^{th} . Using (4), the series solutions are therefore given by

$$\begin{aligned}
 u(x, t) &= \sin x - t \sin x + \frac{1}{6} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3 + \dots)) \\
 v(x, t) &= \sin x - t \sin x + \frac{1}{6} \sin x (3t^2 - 2t(2(t^2 + 3t + 3) \cos x + 3 + \dots))
 \end{aligned}$$

and

Figures 2. (a) and (b) show the exact and numerical solution of system (14) with 10^{th} terms by LADM.

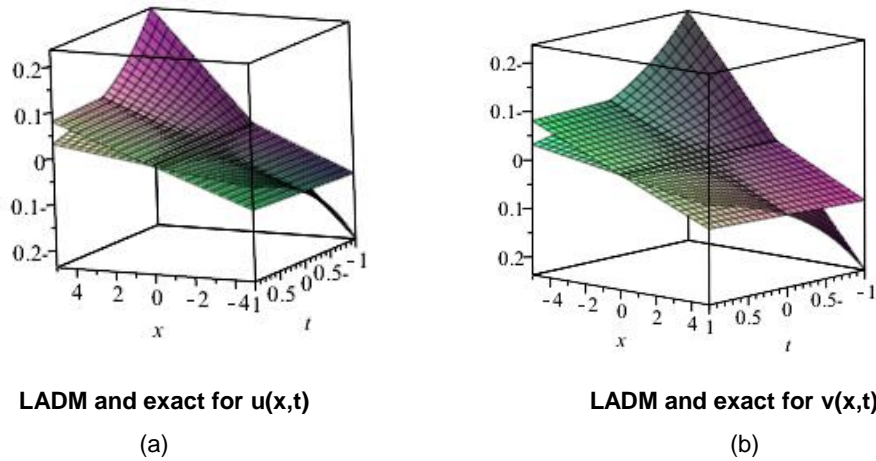


Figure 2. (a) show the exact and LADM numerical solution $u(x, t)$ of example 2, $-5 \leq x \leq 5$, $-1 \leq t \leq 1$; (b) show the exact and LADM numerical solution $v(x, t)$ of example 2, $-5 \leq x \leq 5$, $-1 \leq t \leq 1$.

Example 3: we consider the nonlinear system [13]

$$\begin{aligned} u_t &= uu_x + vv_y \\ v_t &= vv_y + uv_x, \end{aligned} \tag{18}$$

with the initial conditions $u(x, y, 0) = x + y$; $v(x, y, 0) = x + y$, $0 \leq x, y \leq 1$.

The exact solution given as $u(x, y, t) = \frac{x+y}{1-2t}$; $v(x, y, t) = \frac{x+y}{1-2t}$, $0 \leq t \leq T$.

Solution: Taking the Laplace transform on both sides of Eqs. (10) then, by using the differentiation property of Laplace transform and initial conditions, as the same procedure in the above example when we using equations (9)-(10) for the system (10) and according to (3) we have the following 0^{th} components

$$\begin{aligned} u_0(x, y, t) &:= x + y; v_0(x, y, t) := x + y, \text{ and the recursive relation can be written as follows:} \\ u_1(x, y, t) &:= 2(x + y)t; v_1(x, y, t) := 2(x + y)t; u_2(x, y, t) := 4(x + y)t^2; \\ v_2(x, y, t) &:= 4(x + y)t^2; u_3(x, y, t) := 8(x + y)t^3; v_3(x, y, t) := 8(x + y)t^3; \\ u_4(x, y, t) &:= 16(x + y)t^4; v_4(x, y, t) := 16(x + y)t^4 \\ u_{20}(x, y, t) &:= 1048576(x + y)t^{20}; v_{20}(x, y, t) := 1048576(x + y)t^{20} \end{aligned}$$

Similarly, we can also find other components, and the approximate solution for calculating more 20^{th} . Using (3), the series solutions are therefore given by

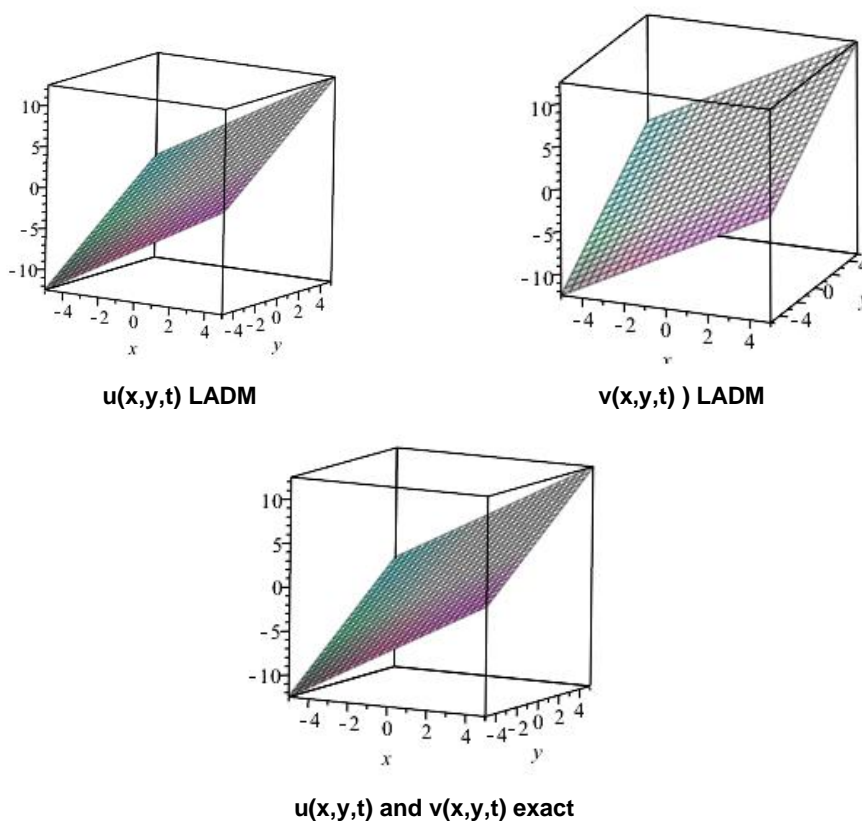
$$\begin{aligned} u(x, y, t) &= v(x, y, t) = (x + y) + 2(x + y)t + 4(x + y)t^2 + 8(x + y)t^3 + 16(x + y)t^4 + \dots \\ &= (x + y)[1 + 2t + 4t^2 + 8t^3 + 16t^4 + \dots] \\ &= (x + y)(1 - 2t)^{-1} = \frac{x + y}{1 - 2t} \end{aligned}$$

that converges to the exact solutions $u(x, y, t) = \frac{x + y}{1 - 2t}$; $v(x, y, t) = \frac{x + y}{1 - 2t}$,

and

Figure 3 ensure that the previous obtained results of system (18) with 30^{th} terms by LADM it converges to the exact solutions. Table 1 show the absolute error between the exact solution and the results obtained from the LADM solution

and comparison with the results obtained by the standard algorithm for Adomian's polynomials ADM and modified ADM solution of system of equation (18) (see [14] and [16]).



Figures 3. show the exact and LADM numerical solutions of $u(x, t)$ and $v(x, t)$ for example 3, $-5 \leq x \leq 5$, $-1 \leq y \leq 1$; $t=0.1$.

Table 1: The numerical results in comparison with the analytical solutions for various values of x , y and t for example 3

t	y	x	Exact(u)=Exact(v)	The absolute error		
				LADM	AMD [16]	Modified AMD[14]
0.1	0.125	0.125	0.312500000	6.55e-016	3.2000e-008	3.0000e-010
		0.785	1.137500000	2.39e-015	1.1600e-007	0.0000e+000
	0.875	0.500	1.250000000	2.62e-015	1.2800e-007	0.0000e+000
		0.125	1.250000000	2.62e-015	1.2800e-007	0.0000e+000
0.2	0.125	0.785	2.075000000	4.35e-015	2.1200e-007	0.0000e+000
		0.500	1.666666667	1.83e-009	4.3690e-005	4.2700e-008
	0.875	0.125	1.516666667	6.67e-009	1.5903e-004	1.5900e-007
		0.785	1.666666667	7.33e-009	1.7476e-004	1.6600e-007
	0.500	0.125	1.666666667	7.33e-009	1.7476e-004	1.6600e-007
		0.785	2.766666667	1.22e-008	2.9010e-004	2.8600e-007
0.3	0.125	0.125	0.625000000	1.37e-005	3.7791e-003	1.4145e-005
		0.785	2.275000000	4.99e-005	1.3756e-002	5.1480e-005
	0.875	0.500	2.500000000	5.48e-005	1.5116e-002	5.6574e-005
		0.125	2.500000000	5.48e-005	1.5116e-002	5.6574e-005
	0.500	0.785	4.150000000	9.10e-005	2.5093e-002	9.3897e-005
		0.125	4.150000000	9.10e-005	2.5093e-002	9.3897e-005



We note from the above results, the absolute error obtained by the proposed algorithm LADM as compared with the absolute error of the standard algorithm for Adomian's polynomials ADM and modified ADM given results more accurate.

Example 4: The mathematical models on many phenomena in applied sciences lead to non-linear PDEs such as the homogeneous form of the system of two dimensional Burger's equations which is proposed as mathematical model of free turbulence [15, 16]

$$\begin{cases} u_t + uu_x + vv_y = \frac{1}{R}(u_{xx} + u_{yy}), \\ v_t + vv_y + uv_x = \frac{1}{R}(v_{xx} + v_{yy}), \end{cases} \quad (19)$$

subject to the initial conditions

$$\begin{cases} u(x, y, 0) = \frac{3}{4} - \frac{1}{4(1 + e^{R(y-x)/8})}, \\ v(x, y, 0) = \frac{3}{4} - \frac{1}{4(1 + e^{R(y-x)/8})}, \end{cases} \quad (20)$$

the exact solution [8]:

$$\begin{cases} u(x, y, t) = \frac{3}{4} - \frac{1}{4(1 + e^{R(-t+4y-4x)/32})}, \\ v(x, y, t) = \frac{3}{4} - \frac{1}{4(1 + e^{R(-t+4y-4x)/32})}, \end{cases}$$

where R is Reynolds number. As the same procedure in the above examples when we use (3)-(6) for the system (19)-(20) and according to (7) –(8), we have the following 0^{th} components:

$$\begin{cases} u_0(x, y, t) = \left(\frac{3}{4} - \frac{1}{4(1 + e^{R(y-x)/8})} \right), \\ v_0(x, y, t) = \left(\frac{3}{4} - \frac{1}{4(1 + e^{R(y-x)/8})} \right), \end{cases} \quad (21)$$

and the recursive relation can be written as follows:

$$\begin{cases} u_{n+1}(x, y, t) = \mathcal{L}^{-1}(\mathcal{L}(F_n - A_n - B_n)), \\ v_{n+1}(x, y, t) = \mathcal{L}^{-1}(\mathcal{L}(G_n - C_n - E_n)), \end{cases} \quad (22)$$

where $n \geq 0$ and the Adomian polynomials A_n, B_n, C_n, E_n, F_n and G_n are defined as follows

$$A_n := \sum_{k=0}^n u_k(x, y, t) \cdot \left(\frac{\partial}{\partial x} (u_{(n-k)}(x, y, t)) \right);$$



$$\begin{aligned}
 B_n &:= \sum_{k=0}^n v_k(x, y, t) \cdot \left(\frac{\partial}{\partial y} (u_{(n-k)}(x, y, t)) \right); \\
 C_n &:= \sum_{k=0}^n v_k(x, y, t) \cdot \left(\frac{\partial}{\partial y} (v_{(n-k)}(x, y, t)) \right); \\
 E_n &:= \sum_{k=0}^n u_k(x, y, t) \cdot \left(\frac{\partial}{\partial x} (v_{(n-k)}(x, y, t)) \right); \\
 F_n &:= \frac{1}{R} \sum_{k=0}^n \left(\frac{\partial^2}{\partial x^2} (u_k(x, y, t)) + \frac{\partial^2}{\partial y^2} (u_k(x, y, t)) \right); \\
 G_n &:= \frac{1}{R} \sum_{k=0}^n \left(\frac{\partial^2}{\partial x^2} (v_k(x, y, t)) + \frac{\partial^2}{\partial y^2} (v_k(x, y, t)) \right).
 \end{aligned} \tag{23}$$

We can give the first few Adomian's polynomials of respectively to get the following components:

$$\begin{aligned}
 A_0 &:= -\frac{1}{128} \frac{\left(2 + 3e^{-\frac{1}{8}R(-y+x)} \right) R e^{-\frac{1}{8}R(-y+x)}}{\left(1 + e^{-\frac{1}{8}R(-y+x)} \right)^3} \\
 A_1 &:= -\frac{1}{4096} \frac{R^2 e^{-\frac{1}{8}R(-y+x)} t \left(3e^{-\frac{1}{4}R(-y+x)} - 2e^{-\frac{1}{8}R(-y+x)} - 2 \right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)} \right)^4} \\
 B_0 &:= \frac{1}{128} \frac{R e^{-\frac{1}{8}R(-y+x)} \left(4 + 3e^{-\frac{1}{8}R(-y+x)} \right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)} \right)^3} \\
 B_1 &:= \frac{1}{4096} \frac{R^2 e^{-\frac{1}{8}R(-y+x)} t \left(3e^{-\frac{1}{4}R(-y+x)} + 2e^{-\frac{1}{8}R(-y+x)} - 4 \right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)} \right)^4} \\
 C_0 &:= -\frac{1}{128} \frac{R e^{-\frac{1}{8}R(-y+x)} \left(4 + 3e^{-\frac{1}{8}R(-y+x)} \right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)} \right)^3} \\
 C_1 &:= -\frac{1}{4096} \frac{R^2 e^{-\frac{1}{8}R(-y+x)} t \left(3e^{-\frac{1}{4}R(-y+x)} + 2e^{-\frac{1}{8}R(-y+x)} - 4 \right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)} \right)^4}
 \end{aligned}$$



$$E_0 := \frac{1}{128} \frac{\left(2 + 3e^{-\frac{1}{8}R(-y+x)}\right) Re^{-\frac{1}{8}R(-y+x)}}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^3}$$

$$E_1 := \frac{1}{4096} \frac{R^2 e^{-\frac{1}{8}R(-y+x)} t \left(3e^{-\frac{1}{4}R(-y+x)} - 2e^{-\frac{1}{8}R(-y+x)} - 2\right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^4}$$

$$F_0 := -\frac{1}{128} \frac{Re^{-\frac{1}{8}R(-y+x)} \left(e^{-\frac{1}{8}R(-y+x)} - 1\right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^3}$$

$F1 :=$

$$\frac{1}{4096} \frac{1}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^4} \left(e^{-\frac{1}{8}R(-y+x)} R \left(-Re^{-\frac{1}{4}R(-y+x)} t \right. \right. \\ \left. \left. + 4Re^{-\frac{1}{8}R(-y+x)} t - 32e^{-\frac{1}{4}R(-y+x)} - Rt + 32 \right) \right)$$

$$G0 := \frac{1}{128} \frac{e^{-\frac{1}{8}R(-y+x)} R \left(-1 + e^{-\frac{1}{8}R(-y+x)} \right)}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^3}$$

$G1 :=$

$$-\frac{1}{4096} \frac{1}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^4} \left(e^{-\frac{1}{8}R(-y+x)} R \left(-Re^{-\frac{1}{4}R(-y+x)} t \right. \right. \\ \left. \left. + 4Re^{-\frac{1}{8}R(-y+x)} t - 32e^{-\frac{1}{4}R(-y+x)} - Rt + 32 \right) \right)$$

and the recursive relation can be written as follows:

$$u_1(x, y, t) := -\frac{1}{128} \frac{Re^{-\frac{1}{8}R(-y+x)} t}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^2}$$

$$v_1(x, y, t) := \frac{1}{128} \frac{Re^{-\frac{1}{8}R(-y+x)} t}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^2}$$

$$u_2(x, y, t) := -\frac{1}{8192} \frac{Re^{-\frac{1}{8}R(-y+x)} \left(e^{-\frac{1}{8}R(-y+x)} - 1\right) t(Rt + 64)}{\left(1 + e^{-\frac{1}{8}R(-y+x)}\right)^3}$$

$$v_2(x, y, t) := \frac{1}{8192} \frac{R e^{-\frac{1}{8} R (-y+x)} \left(e^{-\frac{1}{8} R (-y+x)} - 1 \right) t (Rt + 64)}{\left(1 + e^{-\frac{1}{8} R (-y+x)} \right)^3}.$$

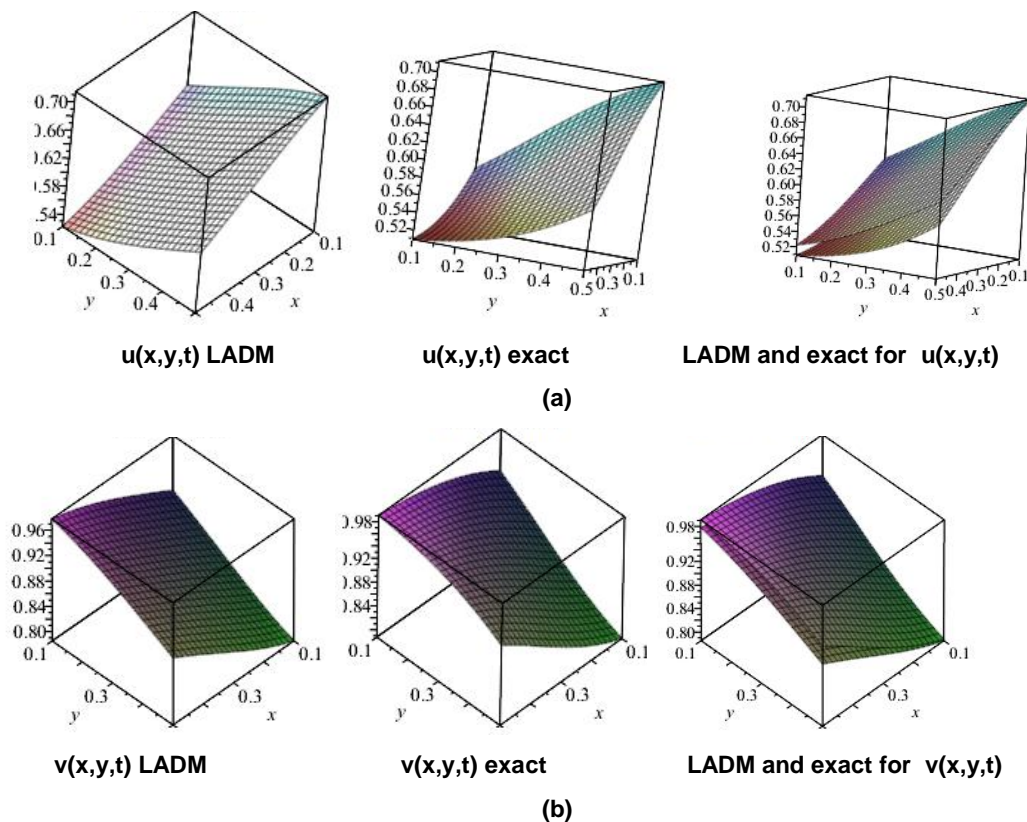
Similarly, we can also find other components, and Using (3), the series solutions are therefore given by

$$u(x, y, t) := -\frac{1}{128} \frac{R e^{-\frac{1}{8} R (-y+x)} t - 96 e^{-\frac{1}{4} R (-y+x)} - 160 e^{-\frac{1}{8} R (-y+x)} - 64}{\left(1 + e^{-\frac{1}{8} R (-y+x)} \right)^2}$$

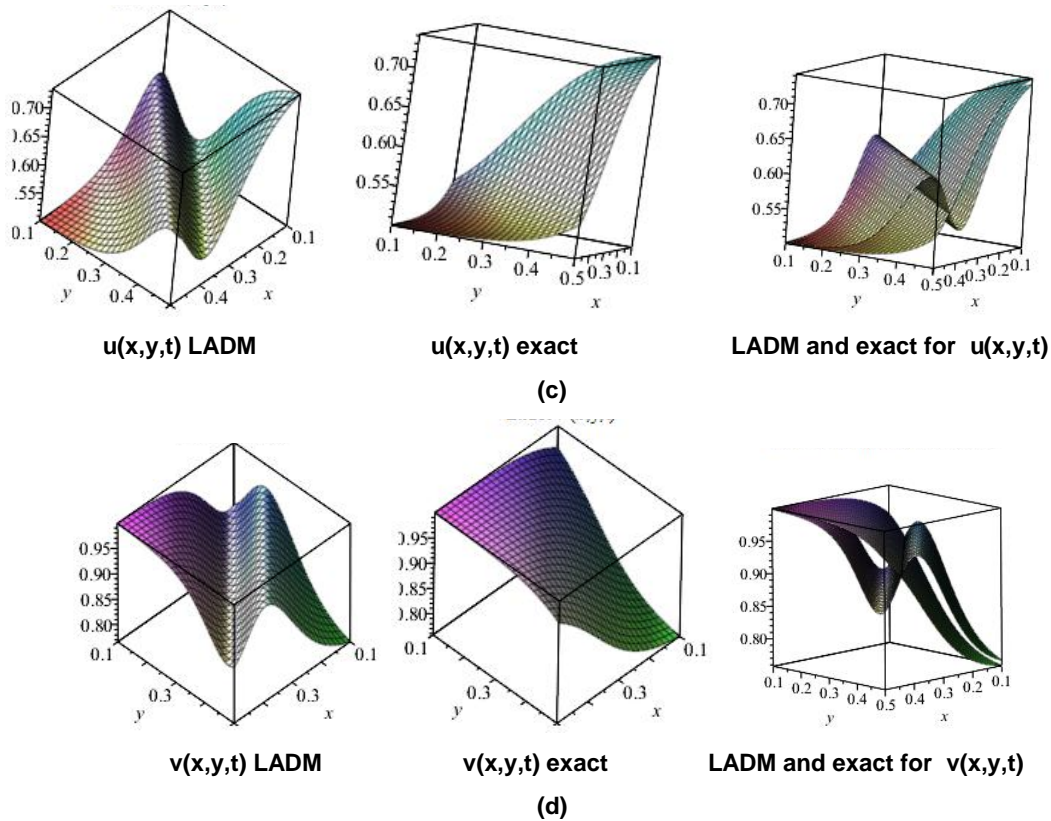
$$v(x, y, t) := \frac{1}{128} \frac{R e^{-\frac{1}{8} R (-y+x)} t + 96 e^{-\frac{1}{4} R (-y+x)} + 224 e^{-\frac{1}{8} R (-y+x)} + 128}{\left(1 + e^{-\frac{1}{8} R (-y+x)} \right)^2}$$

and

Figure 4. (a)-(d) show the exact and numerical solution of system (19) with 3th terms by LADM also table 2 show the absolute error between the exact solution and the results obtained from the LADM solution and comparison with the results obtained by the standard algorithm ADM (see [16]).



Figures 4. (a) and (b) show the exact and LADM of $u(x, y, t)$ and $v(x, y, t)$ for example 4, $0.1 \leq x \leq 0.5$; $0.1 \leq y \leq 0.5$; $t=0.2$; $R=50$



Figures 4. (c) and (d) show the exact and LADM of $u(x,y, t)$ and $v(x, y, t)$ for example 4, $0.1 \leq x \leq 0.5$; $0.1 \leq y \leq 0.5$; $t=0.5$; $R=100$

Table 2: The numerical results in comparison with the analytical solutions for various values of x and t for example 4 when $y = 1$, $R = 50$ and $R=100$ are respectively

t	x	Exact(u) R=50	Exact(v) R=100	The absolute error			
				LADM(u) R=50	AMD(u) [16] R=50	LADM(v) R=100	AMD(v) [16] R=100
0.2	0.1	0.7487736486	0.7499939249	2.73e-004	7.4512e-003	1.88e-006	5.1579e-006
	0.2	0.7477185907	0.7499787972	5.04e-004	1.3845e-003	6.55e-006	1.8001e-005
	0.3	0.7457712717	0.7499260107	9.20e-004	2.5604e-003	2.28e-005	6.2817e-005
	0.4	0.7422137941	0.7497419422	1.65e-003	4.6943e-003	7.96e-005	2.1906e-004
	0.5	0.7358362425	0.7491015994	2.84e-003	8.4719e-003	2.77e-004	7.6234e-004
0.5	0.1	0.7480460278	0.7499844874	6.09e-004	1.7787e-003	1.87e-006	1.8237e-005
	0.2	0.7463741051	0.7499458640	1.28e-003	3.2970e-003	6.53e-006	1.636e-005
	0.3	0.7433101792	0.7498111486	2.07e-003	6.0701e-003	2.28e-005	2.2192e-004
	0.4	0.7377855554	0.7493420815	3.73e-003	1.1040e-002	8.01e-005	7.7219e-004
	0.5	0.7281090401	0.7477185907	6.53e-003	1.9654e-002	2.84e-004	2.6661e-003

The numerical results show that using of LADM by 3- components gives results more accurate than the results using 5- components by standard ADM which is presented in [16].

4. CONCLUSION

LADM is a powerful tool which is capable of handling coupled system of nonlinear of partial differential equations. In this paper the LADM has been successfully applied to find approximate solutions for the homogeneous form of the system of two dimensional Burger's equations without any discretization. The method presents a useful way to develop an analytic



treatment for these systems. The proposed scheme can be applied for system more than two linear and nonlinear partial differential equations.

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