

Global attractor for a class of nonlinear generalized Kirchhoff models

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ABSTRACT

The paper studies the long time behavior of solutions to the initial boundary value problem(IBVP) for a class of Kirchhoff models flow $u_{tt} + \alpha u_t - \beta \Delta u_t - \phi (\|\nabla u\|^2) \Delta u + (1 + \|u\|^2)^{p-1} u = f(x)$. We establish the well-posedness, the existence of the global attractor in natural energy space $(H^2 \cap H_0^1) \times H_0^1$.

Key words: Kirchhoff models; well-posedness; Global attractor.

1 Introduction

In this paper,we are concerned with the existence of global attractor for the following nonlinear plate equation referred to as Kirchhoff models:

$$u_{tt} + \alpha u_{t} - \beta \Delta u_{t} - \phi (\|\nabla u\|^{2}) \Delta u + (1 + |u|^{2})^{p-1} u = f(x) \quad in \quad \Omega \times \mathbb{R}^{+},$$
(1.1)

$$u(x,0) = u_0(x); u_t(x,0) = u_1(x), \quad x \in \Omega,$$
 (1.2)

$$u(x,t)|_{\partial\Omega} = 0, \quad (x) \in \Omega. \tag{1.3}$$

Where Ω is a bounded domain in \mathbb{R}^N , $p \ge 1$, and α, β are positive constants, and the assumptions on $\phi(\|\nabla u\|^2)$ will be specified later.

Global attractor is a basic concept in the study of the asymptotic behavior of solutions for nonlinear evolution equations with various dissipation. From the physical point of view, the global attractor of the dissipative equation(1.1)represents the permanent regime that can be observed when the excitation starts from any point in natural energy space, and its dimension represents the number of degree of freedom of the related turbulent phenomenon and thus the level of complexity concerning the flow. All the information concerning the attractor and its dimension frim the qualitative nature to the quantitative nature then yield valuable information concerning the flows that this physical system can generate. On the physical and numerical simulations[1].

Many authors have focused on the Kirchhoff equations, Igor Chueshov[2]studied the long-time dynamics of Kirchhoff wave models with strong nonlinear damping:

$$\partial_{tt} u - \sigma(\|\nabla u\|^2) \Delta \partial_t u - \phi(\|\nabla u\|^2) \Delta u + f(u) = h(x). \tag{1.4}$$

Tokio Matsuyama and Ryo Ikehata[3] proved on global solutions and energy decay for the wave equations of Kirchhoff type with nonlinear damping terms:

$$u_{tt} - M(\|\nabla u(t)\|_{2}^{2}) \Delta u + \delta |u_{t}|^{p-1} u_{t} = \mu |u|^{q-1} u,$$
(1.5)

with clamped boundary condition

$$u(x,t)|_{\partial\Omega} = 0, \quad t \ge 0,$$
 (1.6)

and M(s) is a positive C^1 - class function for $s \ge 0$ satisfying $M(s) \ge m_0 > 0$ with a constant m_0 , and $\delta > 0$, $\mu \in \mathbf{R}$ are given constants.

Recently, Cheng Jian ling and Yang Zhijian [4] studies the long time behavior of the Kirchhoff type equation with strong damping:

$$u_{tt} - M(\|\nabla u(t)\|_{2}^{2})\Delta u - \Delta u_{t} + g(x, u) + h(u_{t}) = f(x),$$
(1.7)

where $M(s)=1+s^{rac{m}{2}}$, $m\geq 1$. $\Omega\in \mathbf{R}^N$ is a bounded domain with smooth boundary $\partial\Omega$.



Yang Zhijian[5] also studied the longtime behavior of the Kirchhoff type equation with strong damping on \mathbf{R}^N :

$$u_{tt} - M(\|\nabla u\|^2) \Delta u - \Delta u_t + u + u_t + g(x, u) = f(x), \tag{1.8}$$

where $M(s)=1+s^{\frac{m}{2}}$, $m\geq 1$. $\Omega\in \mathbf{R}^N$ is a bounded domain with smooth boundary $\partial\Omega$, f(x) is an external force term. It shows that the related continuous semigroup S(t) possesses a global attractor which is connected and has finite fractal and Hausdorff dimension.

Zhijian Yang and Pengyan Ding[6] studies the longtime dynamics of the Kirchhoff equation with strong damping and critical nonlinearity on \mathbf{R}^N :

$$u_{tt} - \Delta u_{t} - M(\|\nabla u\|^{2}) \Delta u + u_{t} + g(x, u) = f(x), \tag{1.9}$$

where $M \in C^1(\mathbf{R}^+)$, $M'(s) \ge 0$, $M(0) = M_0 > 0$. They established the well-posedness, the existence of the global and exponential attractors in natural energy space $\mathbf{H} = H^1(\mathbf{R}^N) \times L^2(\mathbf{R}^N)$ in critical nonlinearity case.

Claudianor O.Alves and Giovany M.Figueiredo[7]proved the existence of positive solutions for the following class of nonlocal problem:

$$M(\int_{\mathbb{R}^{N}} |\nabla u|^{2} dx + \int_{\mathbb{R}^{N}} V(x) |u|^{2} dx) [-\Delta u + V(x)u] = \lambda f(u) + \gamma u^{\tau},$$
(1.10)

where $\tau=5$ for N=3 and $\tau\in(1,+\infty)$ for N=1,2. λ is a positive parameter and $\gamma\in\{0,1\}$. For more related results, we refer the reader to [8]-[11]. The paper is arranged as follows. In Sec.2, some notations and the main results are stated. In Sec.3, the global existence of solutions to problem (1.1)-(1.3) is established in space $L^{\infty}(0,+\infty;H^1_0\cap L^{2p})\times (L^{\infty}(0,+\infty;L^2)\cap L^2(0,T;H^1_0))$ and $L^{\infty}(0,+\infty;V_2)\times (L^{\infty}(0,+\infty;H^1_0)\cap L^2(0,T;V_2))$. In Sec.4, the existence of global attractor for the dynamical system associated with problem (1.1)-(1.3) is discussed in phase space X_1 .

2 Statement of main results

For brevity, we use the follow abbreviation:

$$\begin{split} L^p &= L^p(\Omega), \quad W^{k,p} = W^{k,p}(\Omega), \quad H^k = W^{k,2}, \quad H = L^2, \quad \left\| \cdot \right\| = \left\| \cdot \right\|_{L^2}, \\ \left\| \cdot \right\|_p &= \left\| \cdot \right\|_{L^p}, \quad V_2 = H^2 \cap H^1_0, \quad V_{2'} = V_{-2}, \quad X_1 = V_2 \times H^1_0, \end{split}$$

with $p \ge 1$. We denote the dual of $W_0^{1,p}$ by $W^{-1,p'}$, with $p' = \frac{p}{p-1}$. And where H^k are the L^2 -based Sobolev

spaces and H_0^k are the completion of $C_0^\infty(\Omega)$ in H^k for k>0. The notation (\cdot,\cdot) for the H – inner product will also be used for the notation of duality pairing between dual spaces.

We define the operator $A: V_2 \rightarrow V_{2'}$,

$$(Au, v) = (\Delta u, \Delta v), \quad for \quad u, v \in V_2.$$

Then, the operators $A^s(s \in R)$ are strictly positive and the spaces $V_s = D(A^{\frac{s}{4}})$ are Hilbert spaces with the scalar products and the norms

$$(u,v)_s = (A^{\frac{s}{4}}u, A^{\frac{s}{4}}v), \quad for \|u\|_{V_s} = \|A^{\frac{s}{4}}u\|,$$

respectively. Obviously,



$$\|u\|_{V_2} = \|A^{\frac{1}{2}}u\| = \|\Delta u\|, \quad \|u\|_{V_1} = \|A^{\frac{1}{4}}u\| = \|\nabla u\|.$$

Now, we state the main results of the paper.

Theorem 2.1. Assume that $(H_1) \ \phi \in C^1(\mathbf{R}^+), \ \phi'(s) \ge 0, \ \phi(0) = \phi_0 \ge 1$,

 $(\boldsymbol{H}_2) \quad f \in \boldsymbol{H}^{-1} \text{, } (u_0, u_1) \in \boldsymbol{H}_0^1 \times \boldsymbol{H} \text{ , } p \geq 1 \text{ . Then the solution } (u, v) \text{ of the problem(1.1)-(1.3) satisfies } \boldsymbol{H}_2 = \boldsymbol{H}_1 \times \boldsymbol{H}_2 + \boldsymbol{H}_2 \times \boldsymbol{H}_3 + \boldsymbol{H}_3 \times \boldsymbol{H}_4 + \boldsymbol{H}_4 \times \boldsymbol{H}_4 +$

$$H_1(t) \le H_1(0)e^{-k_1t} + \frac{C_1}{k_1}(1 - e^{-k_1t}), \tag{2.1}$$

$$\beta \int_0^T \|\nabla v\|^2 ds \le H_1(0) + \int_0^T C_1 ds. \tag{2.2}$$

Where

$$v = u_t + \varepsilon u$$
, $0 < \varepsilon \le \min\{\frac{\alpha}{4}, \frac{\lambda_1}{2\alpha}, \frac{1}{2\beta}\}$, and

$$\begin{split} H_1(t) &= \left\| v \right\|^2 + \int_0^{\left\| \nabla u \right\|^2} \! \phi(s) ds - \beta \varepsilon \left\| \nabla u \right\|^2 + \frac{1}{p} \int_{\Omega} (1 + \left| u \right|^2)^p \, dx \,, \quad \text{then problem} \quad (1.1) - (1.3) \quad \text{admits a solution} \\ u &\in L^{\infty}(0, +\infty; H_0^1 \cap L^{2p}) \,, \quad v \in L^{\infty}(0, +\infty; L^2) \cap L^2(0, T; H_0^1) \,. \end{split}$$

Remark 2.1 In addition to the assumptions of Theorem 2.1, we know that $\phi(s)$ and $\phi'(s)$ are bounded.

Theorem 2.2. In addition to the assumptions of Theorem 2.1, assume that (H_3)

$$\begin{cases} 1 \le p < +\infty & N = 1, 2, \\ 1 \le p \le \frac{N-1}{N-2} & N \ge 3, \end{cases}$$

 $(H_4) \ \ f \in H$, $(u_0,u_1) \in V_2 \times H^1_0$. Then the solution (u,v) of the problem(1.1)-(1.3) satisfies

$$H_3(t) \le H_3(0)e^{-\delta} + \frac{C_5}{\delta}(1 - e^{-\delta}),$$
 (2.3)

$$\beta \int_{0}^{T} \|\Delta u_{t}\|^{2} ds \le H_{3}(0) + \int_{0}^{T} C_{5} ds. \tag{2.4}$$

Then problem (1.1)-(1.3) admits a unique solution $u\in L^\infty(0,+\infty;V_2)$, $u_t\in L^\infty(0,+\infty;H_0^1)\cap L^2(0,T;V_2)$.

Remark 2.2 We denote the solution in Theorem 2.2 by $S(t)(u_0,u_1)=(u(t),u_t(t))$. Then S(t) composes a continuous demigroup in X_1 .

Theorem 2.3 In addition to the assumptions of Theorem 2.2, then the continuous semigroup S(t) defined in Remark 2.1 possesses in X_1 a global attractor which is connected.

3 Global existence of solutions

We first prepare the following well known lemmas which will be needed later.

Lemma 3.1(Sobolev-Poincare) [2][11]. If either $1 \le p < +\infty(N=1,2)$ or $1 \le p \le \frac{N-1}{N-2}(N \ge 3)$, then there is a constant $C(\Omega,4p-2)$ such that

$$||u||_{4p-2} \le C(\Omega, 4p-2)||\nabla u||, \quad for \quad u \in H_0^1(\Omega).$$



In other words,

$$C(\Omega, 4p - 2) = \sup\{\frac{\|u\|_{4p-2}}{\|\nabla u\|} \mid u \in H_0^1(\Omega), u \neq 0\}$$

is positive and finite.

Lemma 3.2(Gronwall's inequality) [11]. Let H(t) be a non-negative absolutely continuous function on $[0,\infty)$ which satisfies the differential inequality

$$\frac{dH}{dt} + kH \le C, \quad t \ge 0,$$

where k > 0 and $C \ge 0$ are constants. Then

$$H(t) \leq R_2, \quad t \geq T(H_0),$$

where $T(\boldsymbol{H}_0)$ is a constant depending on $\boldsymbol{H}_0 = \boldsymbol{H}(0)$.

Proof of Theorem 2.1

Proof. Let $v = u_t + \varepsilon u$, $0 < \varepsilon \le \min\{\frac{\alpha}{4}, \frac{\lambda_1}{2\alpha}, \frac{1}{2\beta}\}$, then v satisfies

$$v_{t} + (\alpha - \varepsilon)v + (\varepsilon^{2} - \alpha\varepsilon)u - \beta\Delta v + \beta\varepsilon\Delta u - \phi(\|\nabla u\|^{2})\Delta u + (1 + |u|^{2})^{p-1}u = f(x).$$
 (3.1)

Taking H – inner product by v in (3.1), we have

$$\frac{1}{2} \frac{d}{dt} \|v\|^{2} + (\alpha - \varepsilon) \|v\|^{2} + (\varepsilon^{2} - \alpha \varepsilon)(u, v) + \beta \|\nabla v\|^{2} + \beta \varepsilon (\Delta u, v) - (\phi(\|\nabla u\|^{2}) \Delta u, v) + ((1 + |u|^{2})^{p-1} u, v) = (f, v).$$
(3.2)

By using Holder's inequality, Young's inequality and Poincare's inequality, we deal with the terms in (3.2) one by one as follow

$$(\alpha - \varepsilon) \|v\|^2 \ge \frac{3\alpha}{4} \|v\|^2 \tag{3.3}$$

$$(\varepsilon^{2} - \alpha \varepsilon)(u, v) \ge \frac{\varepsilon^{2} - \alpha \varepsilon}{\sqrt{\lambda_{1}}} \|\nabla u\| \|v\|$$

$$\geq -\frac{\varepsilon \alpha^{2}}{\lambda_{1}} \|v\|^{2} - \frac{\varepsilon}{4} \|\nabla u\|^{2}$$

$$\geq -\frac{\varepsilon}{4} \|\nabla u\|^2 - \frac{\alpha}{2} \|v\|^2,\tag{3.4}$$

and

$$\beta \varepsilon (\Delta u, v) = -\frac{\beta \varepsilon}{2} \frac{d}{dt} \|\nabla u\|^2 - \beta \varepsilon^2 \|\nabla u\|^2, \tag{3.5}$$

$$-(\phi(\|\nabla u\|^{2})\Delta u, v) = \phi(\|\nabla u\|^{2})(\nabla u, \nabla v) = \frac{1}{2}\frac{d}{dt}(\int_{0}^{\|\nabla u\|^{2}}\phi(s)ds) + \varepsilon\phi(\|\nabla u\|^{2})\|\nabla u\|^{2}$$

$$\geq \frac{1}{2} \frac{d}{dt} \left(\int_0^{\|\nabla u\|^2} \phi(s) ds \right) + \varepsilon \left(\int_0^{\|\nabla u\|^2} \phi(s) ds \right) \tag{3.6}$$



$$((1+|u|^{2})^{p-1}u,v) = \frac{1}{2p}\frac{d}{dt}(\int_{\Omega}(1+|u|^{2})^{p}dx) + \varepsilon\int_{\Omega}(1+|u|^{2})^{p-1}|u|^{2}dx$$

$$\geq \frac{1}{2p}\frac{d}{dt}(\int_{\Omega}(1+|u|^{2})^{p}dx) + \frac{\varepsilon}{p}\int_{\Omega}(1+|u|^{2})^{p}dx - \frac{\Omega}{p\varepsilon}.$$
(3.7)

By (3.3)-(3.7), it follows from that

$$\frac{d}{dt} [\|v\|^{2} + \int_{0}^{\|\nabla u\|^{2}} \phi(s) ds - \beta \varepsilon \|\nabla u\|^{2} + \frac{1}{p} \int_{\Omega} (1 + |u|^{2})^{p} dx] + \frac{\alpha}{2} \|v\|^{2} + \varepsilon (2 \int_{0}^{\|\nabla u\|^{2}} \phi(s) ds - 2(\beta \varepsilon + \frac{1}{4}) \|\nabla u\|^{2}) + \frac{2\varepsilon}{p} \int_{\Omega} (1 + |u|^{2})^{p} dx + \beta \|\nabla v\|^{2} \le \frac{1}{\beta} \|\nabla^{-1} f\|^{2} + \frac{2\Omega}{p\varepsilon}. \tag{3.8}$$

Because of $0<\varepsilon\leq\frac{1}{2\beta}$ and (H_1) , we can get

$$2\int_{0}^{\|\nabla u\|^{2}} \phi(s)ds - (2\beta\varepsilon + \frac{\varepsilon}{2})\|\nabla u\|^{2} \ge \int_{0}^{\|\nabla u\|^{2}} \phi(s)ds - \beta\varepsilon\|\nabla u\|^{2}. \tag{3.9}$$

Substituting(3.9) into (3.8) get

$$\frac{d}{dt} \left[\left\| v \right\|^{2} + \int_{0}^{\left\| \nabla u \right\|^{2}} \phi(s) ds - \beta \varepsilon \left\| \nabla u \right\|^{2} + \frac{1}{p} \int_{\Omega} (1 + \left| u \right|^{2})^{p} dx \right] + \frac{\alpha}{2} \left\| v \right\|^{2} + \varepsilon \left(\int_{0}^{\left\| \nabla u \right\|^{2}} \phi(s) ds - \beta \varepsilon \left\| \nabla u \right\|^{2} \right) + \frac{\varepsilon}{p} \int_{\Omega} (1 + \left| u \right|^{2})^{p} dx + \beta \left\| \nabla v \right\|^{2} \le \frac{1}{\beta} \left\| \nabla^{-1} f \right\|^{2} + \frac{2\Omega}{p\varepsilon}. \tag{3.10}$$

Taking $k_1 = \min\{\frac{\alpha}{2}, \varepsilon\} = \varepsilon$, then

$$\frac{d}{dt}H_{1}(t) + k_{1}H_{1}(t) + \beta \|\nabla v\|^{2} \le \frac{1}{\beta} \|\nabla^{-1}f\|^{2} + \frac{2\Omega}{p\varepsilon} := C_{1}, \tag{3.11}$$

$$H_1(t) \le H_1(0)e^{-k_1t} + \frac{C_1}{k_1}(1 - e^{-k_1t}),$$
 (3.12)

$$\beta \int_0^T \|\nabla v\|^2 ds \le H_1(0) + \int_0^T C_1 ds. \tag{3.13}$$

Proof of Theorem 2.2

Proof. Taking H – inner product by $-\Delta u$, $-\Delta u$, in(1.1), we get

$$\frac{1}{2} \frac{d}{dt} [\alpha \|\nabla u\|^{2} + 2(u_{t}, -\Delta u) + \beta \|\Delta u\|^{2}] + \phi(\|\nabla u\|^{2}) \|\Delta u\|^{2}
= ((1 + |u|^{2})^{p-1} u, \Delta u) + (f, -\Delta u),$$
(3.14)



$$\frac{1}{2} \frac{d}{dt} \|\nabla u_t\|^2 + \alpha \|\nabla u_t\|^2 + \beta \|\Delta u_t\|^2 = \phi(\|\nabla u\|^2)(\Delta u, \Delta u_t)
+ (1 + |u|^2)^{p-1} u, \Delta u_t) + (f, -\Delta u_t).$$
(3.15)

We have

$$|((1+|u|^{2})^{p-1}u,\Delta u)| \leq \begin{cases} |(2^{p-1}u,\Delta u)| & |u| < 1, \\ |((2^{p-1}|u|^{2p-2}u,\Delta u)| & |u| \geq 1, \end{cases}$$
(3.16)

where

$$|(2^{p-1}u, \Delta u)| \le \frac{1}{2} ||\Delta u||^2 + 2^{2p-1} ||u||^2,$$
 (3.17)

$$|\left((2^{p-1} \mid u \mid^{2p-2} u, \Delta u) \mid \leq 2^{p-1} \|u\|_{4p-2}^{2p-1} \|\Delta u\| \leq C_2(\Omega, 4p-2) \cdot 2^{p-1} \|\nabla u\|^{2p-1} \|\Delta u\|$$

$$\leq \frac{1}{8} \|\Delta u\|^2 + 2^{2p-1} C_2 \cdot \|\nabla u\|^{4p-2}, \tag{3.18}$$

then

$$\left| \left((1 + |u|^2)^{p-1} u, \Delta u \right) \right| \le \frac{1}{4} \left\| \Delta u \right\|^2 + 2^{2p-1} \left\| u \right\|^2 + 2^{2p-1} C_2 \cdot \left\| \nabla u \right\|^{4p-2}, \tag{3.19}$$

$$|(f, -\Delta u)| \le \frac{1}{4} ||\Delta u||^2 + ||f||^2.$$
 (3.20)

Also, we have

$$\left| \left((1 + |u|^2)^{p-1} u, \Delta u_t \right) \right| \le \frac{\beta}{4} \left\| \Delta u_t \right\|^2 + \frac{2^{2p-1}}{\beta} \left\| u \right\|^2 + \frac{2^{2p-1}}{\beta} C_2^2 \cdot \left\| \nabla u \right\|^{4p-2}, \tag{3.21}$$

$$\phi(\|\nabla u\|^2) |(\Delta u, \Delta u_t)| \le \frac{\beta}{8} \|\Delta u_t\|^2 + 2 \frac{\phi^2(\|\nabla u\|^2)}{\beta} \|\Delta u\|^2,$$
 (3.22)

$$|(f, -\Delta u_t)| \le \frac{\beta}{8} ||\Delta u_t||^2 + \frac{2}{\beta} ||f||^2.$$
 (3.23)

Substituting (3.19), (3.20) into (3.14), we receive

$$\frac{d}{dt} \left[\alpha \| \nabla u \|^{2} + 2(u_{t}, -\Delta u) + \beta \| \Delta u \|^{2} \right] + \| \Delta u \|^{2}
\leq 2^{2p} \| u \|^{2} + 2^{2p} \cdot C_{2}^{2} \| \nabla u \|^{4p-2} + 2 \| f \|^{2} := C_{3}.$$
(3.24)

Substituting (3.21)-(3.23) into (3.15), we receive

$$\frac{d}{dt} \|\nabla u_t\|^2 + 2\alpha \|\nabla u_t\|^2 + \beta \|\Delta u_t\|^2 \le \frac{4\phi^2(\|\nabla u\|^2)}{\beta} \|\Delta u\|^2 + C_4, \tag{3.25}$$

where
$$C_4 = \frac{2^{2p}}{\beta} \|u\|^2 + C_2^2 \frac{2^{2p}}{\beta} \|\nabla u\|^{4p-2} + \frac{4}{\beta} \|f\|^2$$
.

Let
$$K_1 = \frac{4\phi^2(\left\|\nabla u\right\|^2)}{\beta}$$
 , $K_2 = K_1 + 1$, $(3.24) \times K_2 + (3.25)$, we have



$$\frac{d}{dt} \left[K_2(\alpha \|\nabla u\|^2 + 2(u_t, -\Delta u) + \beta \|\Delta u\|^2) + \|\nabla u_t\|^2 \right] + \|\Delta u\|^2 + 2\alpha \|\nabla u_t\|^2
+ \beta \|\Delta u_t\|^2 \le K_2 C_3 + C_4.$$
(3.26)

Taking H – inner product by u_t in (1.1), we have

$$\frac{d}{dt}H_2 + \alpha \|u_t\|^2 + 2\beta \|\nabla u_t\|^2 \le \frac{\|f\|^2}{2\alpha},\tag{3.27}$$

Let $K_3 = \frac{2K_2}{\beta} + 1$, $(3.27) \times K_3 + (3.26)$, we get

$$\frac{d}{dt}H_{3} + \|\Delta u\|^{2} + 2\alpha\|\nabla u_{t}\|^{2} + \beta\|\Delta u_{t}\|^{2} \le K_{2}C_{3} + C_{4} + K_{3}\frac{\|f\|^{2}}{2\alpha},$$
(3.28)

where

$$\frac{K_{2}\beta\|\Delta u\|^{2}}{2} + \|\nabla u_{t}\|^{2} \leq H_{3}$$

$$= K_{2}(\alpha\|\nabla u\|^{2} + 2(u_{t}, -\Delta u) + \beta\|\Delta u\|^{2}) + \|\nabla u_{t}\|^{2} + K_{3}H_{2}$$

$$\leq \frac{1}{\delta}(\|\Delta u\|^{2} + \alpha\|\nabla u_{t}\|^{2}) + K_{3}H_{2},$$
(3.29)

where δ is a small positive constants. Now we have

$$\frac{d}{dt}H_3 + \delta H_3 + \beta \|\Delta u_t\|^2 \le K_2 C_3 + C_4 + K_3 \frac{\|f\|^2}{2\alpha} + \delta K_3 H_2 := C_5.$$
(3.30)

Hence, according to Gronwall's inequality and integrating (3.30) over (0,T), we get

$$H_3(t) \le H_3(0)e^{-\delta t} + \frac{C_5}{\delta}(1 - e^{-\delta t}),$$
 (3.31)

$$\beta \int_0^T \left\| \Delta u_t \right\|^2 ds \le H_3(0) + \int_0^T C_5 ds. \tag{3.32}$$

We know that u is the solution of problem (1.1)-(1.3), with $u \in L^{\infty}(0,+\infty;V_2)$, $u_t \in L^{\infty}(0,+\infty;H_0^1) \cap L^2(0,T;V_2)$. The uniqueness is standard; let u(t) and v(t) be two solutions, then w(t) = u(t) - v(t) satisfies

$$w_{tt} + \alpha w_{t} - \beta \Delta w_{t} - (\phi (\|\nabla u\|^{2}) \Delta u - \phi (\|\nabla v\|^{2}) \Delta v) +$$

$$(1 + |u|^{2})^{p-1} u - (1 + |v|^{2})^{p-1} v = 0,$$
(3.33)

with w=0 on $[0,+\infty)\times\partial\Omega$ and $w(0)=w_{_t}(0)=0$ in Ω . Taking the H - inner product of (3.33) with $w_{_t}$, one can find that

$$\frac{1}{2}\frac{d}{dt}[\|w_{t}\|^{2} + \phi(\|\nabla u\|^{2})\|\nabla w\|^{2}] + \alpha\|w_{t}\|^{2} + \beta\|\nabla w_{t}\|^{2}$$



$$= \phi'(\|\nabla u\|^2)(\nabla u, \nabla w_t)\|\nabla w\|^2 + (\phi(\|\nabla u\|^2) - \phi(\|\nabla v\|^2))(\Delta v, w_t)$$

$$-((1+|u|^2)^{p-1}u - (1+|v|^2)^{p-1}v, w_t). \tag{3.34}$$

Here, we note that the first and second terms in the right-hand side of (3.34) are bounded by

$$\phi'(\|\nabla u\|^2)(\nabla u, \nabla w_t)\|\nabla w\|^2 \le C_6 \|\nabla w\|^2$$
(3.35)

$$(\phi(\|\nabla u\|^2) - \phi(\|\nabla v\|^2))(\Delta v, w_t) \le C_7 \|\nabla v\| \|w_t\|, \tag{3.36}$$

respectively. Making use of

$$\begin{split} & \left\| (1+|u|^{2})^{p-1}u - (1+|v|^{2})^{p-1}v \right\| \leq \left\| (1+|u|^{2})^{p-1}w + ((1+|u|^{2})^{p-1} - (1+|v|^{2})^{p-1})v \right\| \\ & \leq \left\| (1+|u|^{2})^{p-1}w + (p-1)(1+|\zeta|^{2})^{p-2}2\zeta v w \right\| \\ & \leq \left\| (1+|u|^{2})^{p-1}w + 2(p-1)(1+|\rho|^{2})^{p-1}w \right\| \\ & \leq \left\| (1+|u|^{2})^{p-1}w + 2(p-1)P(1+|\rho|^{2})^{p-1}w \right\| \\ & \leq \left\| (1+|u|^{2})^{p-1}w + 2(p-1)P(1+|\rho|^{2})^{p-1}w \right\| \\ & \leq (2p-1)2^{p-1} \|w\| + 2^{p-1} \|u\|_{4p-2}^{2p-2} \|w\|_{4p-2} \\ & + (2p-2)2^{p-1} \|\rho\|_{4p-2}^{2p-2} \|w\|_{4p-2} \\ & \leq (2p-1)2^{p-1} \|w\| + 2^{p-1}C_{8} \|\nabla u\|^{2p-2} \|\nabla w\| \\ & + (2p-2)2^{p-1}C_{9} \|\nabla \rho\|^{2p-2} \|\nabla w\|, \end{split} \tag{3.37}$$

where $\zeta = \theta u + (1-\theta)v$, $0 \le \theta \le 1$, $\rho = \max\{\zeta, v\}$. One can also find that the last term in the right-hand side of (3.34) is bounded by

$$((1+|u|^{2})^{p-1}u - (1+|v|^{2})^{p-1}v, w_{t}) \le C_{10} \|\nabla w\| \|w_{t}\|.$$
(3.38)

Hence, integrating (3.34) over (0,t), we get

$$\|w_t\|^2 + \phi(\|\nabla u\|^2)\|\nabla w\|^2 \le C_{11} \int_0^t (\|w_t\|^2 + \|\nabla w\|^2) ds, \tag{3.39}$$

which, by Gronwall's inequality, implies $w \equiv 0$. This completes the proof of Theorem 2.2.

4 Bounded absorbing sets and Global attractor in X_1

Lemma 4.1 [8][11] The continuous semigroup S(t) defined on a Banach space X has a global attractor which is connected when the following conditions are satisfied

- 1) There exists a bounded absorbing set $B \subset X$ such that for any bounded set $B_0 \subset X$, $dist(S(t)B_0,B) \to 0$ as $t \to +\infty$.
- 2) S(t) can be decomposed as S(t) = P(t) + U(t), where P(t) is a continuous map from X to itself with the property that, for any bounded set $B_0 \subset X$,

$$\sup_{\theta \in B_0} \|P(t)\theta\|_X \to 0, \quad t \to \infty, \tag{4.1}$$

and U(t) is precompact for $t > T_0$ for some T_0 .

Proof of Theorem 2.3



Proof. According to Theorem 2.2, we get

$$\|\Delta u\|^2 + \|\nabla u_t\|^2 \le CR_2, \quad t \ge T(\|(u_0, u_1)\|_{X_1}).$$
 (4.2)

(4.2) implies that the ball $B(CR_2)$ centered at zero with a radius $\sqrt{CR_2}$ in X_1 is an absorbing set of S(t). Moreover, integrating (3.30) over (t,t+1), respectively, and exploiting (4.2), we have

$$\int_{t}^{t+1} \|\Delta u_{t}(s)\|^{2} ds \le C(\|(u_{0}, u_{1})\|_{X_{1}}, \Omega, \|f\|), \quad t > 0.$$
(4.3)

Decomposition of S(t): Let R>0 be given with $\left\|(u_0,u_1)\right\|_{X_1}\leq R$, we know from (3.31) and (4.2) that

$$\|(u(t), u_t(t))\|_{X_1} \le C_{12}, \quad t \ge 0,$$
 (4.4)

where and in the sequel

$$C_{12} = \begin{cases} C(H_3(0)) & 0 \le t \le T(R), \\ \sqrt{CR_2} & t > T(R). \end{cases}$$
 (4.5)

Let us write now u = v + w, where

$$w_{tt} + \alpha w_t - \beta \Delta w_t - \phi_0 \Delta w = 0, \quad w(0) = u_0, \quad w_t(0) = u_1,$$
 (4.6)

$$v_{tt} + \alpha v_{t} - \beta \Delta v_{t} - \phi_{0} \Delta v = f + \phi (\|\nabla u\|^{2}) \Delta u - (1 + \|u\|^{2})^{p-1} u := \varphi,$$

$$v(0) = 0, \quad v_{\star}(0) = 0.$$
 (4.7)

Lemma 4.2 If $(u_0, u_1) \in B$, (w, w_t) is the solution of (4.6), then

$$\|q\|^2 + \|\nabla w\|^2 \le R(t), \quad t \ge 0,$$
 (4.8)

and

$$R(t) \to 0$$
, as $t \to +\infty$. (4.9)

Where $q = w_t + \varepsilon w$, $0 < \varepsilon \le \min\{\frac{\alpha}{4}, \frac{\lambda_1}{2\alpha}, \frac{1}{2\beta}\}$.

Proof. Let $q = w_t + \varepsilon w$, $0 < \varepsilon \le \min\{\frac{\alpha}{4}, \frac{\lambda_1}{2\alpha}, \frac{1}{2\beta}\}$, then q satisfies

$$q_t + (\alpha - \varepsilon)q + (\varepsilon^2 - \alpha\varepsilon)w - \beta\Delta q - (\phi_0 - \beta\varepsilon)\Delta w = 0. \tag{4.10}$$

Taking H – inner product by q in (4.10), we have

$$\frac{1}{2}\frac{d}{dt}\|q\|^2 + (\alpha - \varepsilon)\|q\|^2 + (\varepsilon^2 - \alpha\varepsilon)(w, q) + \beta\|\nabla q\| + (\phi_0 - \beta\varepsilon)(\nabla w, \nabla q) = 0. \tag{4.11}$$

It is clear that

$$(\alpha - \varepsilon) \|q\|^2 \ge \frac{3\alpha}{4} \|q\|^2, \quad (\varepsilon^2 - \alpha \varepsilon)(w, q) \ge -\frac{\varepsilon}{4} \|\nabla w\|^2 - \frac{\alpha}{2} \|q\|^2, \tag{4.12}$$

$$(\nabla w, \nabla q) = \frac{1}{2} \frac{d}{dt} \|\nabla w\|^2 + \varepsilon \|\nabla w\|^2. \tag{4.13}$$

So



$$\frac{d}{dt} [\|q\|^2 + (\phi_0 - \beta \varepsilon) \|\nabla w\|^2] + \frac{\alpha}{2} \|q\|^2 + \varepsilon (2\phi_0 - 2\beta \varepsilon - \frac{1}{2}) \|\nabla w\|^2 + 2\beta \|\nabla q\|^2 \le 0. \tag{4.14}$$

Because of $0<\varepsilon\leq \frac{1}{2\beta}$, we get $2\phi_0-2\beta\varepsilon-\frac{1}{2}\geq \phi_0-\beta\varepsilon$,

then by (H_1) and gronwall's inequality,

$$\frac{1}{2}(\|q\|^2 + \|\nabla w\|^2) \le \|q\|^2 + (\phi_0 - \beta \varepsilon)\|\nabla w\|^2 \le (\|q_0\|^2 + (\phi_0 - \beta \varepsilon)\|\nabla w_0\|^2)e^{-\varepsilon t}.$$
(4.15)

Lemma 4.2 is proven.

Lemma 4.3 If $(u_0,u_1)\in B$, (v,v_t) is the solution of (4.7), then it exists compact set $N(T)\subset X_1$ and

$$(v, v_t) \in N(T). \tag{4.16}$$

Proof. Applying $A^{\sigma_1}(0 < \sigma_1 = \frac{1}{2})$ to both sides of (4.7), we have

$$\xi_{tt} + \alpha \xi_t - \beta \Delta \xi_t - \phi_0 \Delta \xi = A^{\sigma_1} \varphi, \quad \xi(0) = 0, \quad \xi_t(0) = 0, \tag{4.17}$$

where $\xi = A^{\sigma_1} v$. Let $\eta = \xi_{\scriptscriptstyle t} + \mathcal{E}\!\!\xi$, then

$$\eta_{t} + (\alpha - \varepsilon)\eta + (\varepsilon^{2} - \alpha\varepsilon)\xi - \beta\Delta\eta - (\phi_{0} - \beta\varepsilon)\Delta\xi = A^{\sigma_{1}}\varphi. \tag{4.18}$$

Taking H-inner product by η in (4.18), we obtain

$$\frac{1}{2}\frac{d}{dt}\left[\left\|\eta\right\|^{2} + (\phi_{0} - \beta\varepsilon)\left\|\nabla\xi\right\|^{2}\right] + \frac{\alpha}{4}\left\|\eta\right\|^{2} + \varepsilon(\phi_{0} - \beta\varepsilon - \frac{1}{4})\left\|\nabla\xi\right\|^{2} + \beta\left\|\nabla\eta\right\|^{2} \le (A^{\sigma_{1}}\varphi, \eta). \tag{4.19}$$

By the same argument of Theorem 2.2 we can obtain

$$|(A^{\sigma_1}f,\eta)| \le C_{13}(||f||,\beta) + \frac{\beta}{8} ||\nabla \eta||^2,$$
 (4.20)

$$|(A^{\sigma_1}(\phi(\|\nabla u\|^2)\Delta u), \eta)| \le C_{14}(\phi(\|\nabla u\|^2)\Delta u, \|\Delta u\|, \beta) + \frac{\beta}{8} \|\nabla \eta\|^2, \tag{4.21}$$

$$|(A^{\sigma_1}((1+|u|^2)^{p-1}u),\eta)| \le C_{15}(\|\nabla u\|,\beta) + \frac{\beta}{4}\|\nabla \eta\|^2.$$
(4.22)

It follows from (4.19)-(4.22) that

$$\frac{d}{dt} [\|\boldsymbol{\eta}\|^{2} + (\phi_{0} - \beta \varepsilon)\|\nabla \boldsymbol{\xi}\|^{2}] + \frac{\alpha}{2} \|\boldsymbol{\eta}\|^{2} + \varepsilon(\phi_{0} - \beta \varepsilon)\|\nabla \boldsymbol{\xi}\|^{2} + \beta\|\nabla \boldsymbol{\eta}\|^{2}$$

$$\leq C_{16}(\|\boldsymbol{f}\|, \phi(\|\nabla \boldsymbol{u}\|^{2}, \|\Delta \boldsymbol{u}\|, \beta). \tag{4.23}$$

Then

$$H_4(t) \le H_4(0)e^{-\alpha} + \frac{C_{16}}{\varepsilon}(1 - e^{-\alpha}),$$
 (4.24)

where $H_4=\left\|\eta\right\|^2+(\phi_0-etaarepsilon)\left\|
ablaarepsilon
ight\|^2$. Since $H_4(0)=0$, (4.24) means

$$H_4(t) \le \frac{C_{16}}{\varepsilon} (1 - e^{-\varepsilon t}), \quad t \ge 0,$$
 (4.25)

which implies

$$\left\| \xi_{t} + \xi \xi \right\|^{2} + \left\| \nabla \xi \right\|^{2} \le C_{17}, \tag{4.26}$$



$$\left\| (v, v_t) \right\|_{V_{2+4\sigma_1} \times V_{4\sigma_1}}^2 \le C_{17}, \quad t > 0, \tag{4.27}$$

for $\xi = A^{\sigma_1} v$.

Since $V_{2+4\sigma_1} \times V_{4\sigma_1} \mapsto X_1$ is compact embedded, which means that the bounded set in $V_{2+4\sigma_1} \times V_{4\sigma_1}$ is the compact set in X_1 .

Lemma 4.3 is proved.

Define

$$P(t)(u_0, u_1) = (w(t), w_t(t)), \quad U(t)(u_0, u_1) = (v(t), v_t(t)). \tag{4.28}$$

Obviously, S(t)=P(t)+U(t). Lemma 3.1 shows that for any $(u_0,u_1)\in B_0\subset X_1$, the map $P(t):X_1\to X_1$ is continuous and satisfies(4.1). Moreover, Lemma 3.2 shows that the map U(t) is precompact for $t\geq 0$ for $V_{2+4\sigma_1}\times V_{4\sigma_1}\mapsto X_1$. So S(t) has in X_1 a global attractor A which is connected.

This completes the proof of Theorem 2.3.

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References

- Zhijian Yang; Na Feng and To Fu Ma, "Global attractor for the generalized double dispersion," Nonlinear Analysis 115(2015)103-116.
- 2. Igor Chueshov,"Long-time dynamics of Kirchhoff wave models with strong nonlinear damping,"Journal of Differential Equations 252(2012)1229-1262.
- 3. Tokio Matsuyama and Ryo Ikehata,"On global solutions and energy decay for the wave equations of Kirchhoff type with nonlinear damping terms, "Journal of Mathematical Analysis and Applications 204,729-753115(1996).
- 4. Cheng Jian ling and Yang Zhijian, "Asymptotic behavior of the Kirchhoff type equation, "Acta Mathematica Scientia 2011,31A(4):1008-1021.
- 5. Yang Zhijian,"Longtime behavior of the Kirchhoff type equation with strong damping on \mathbb{R}^N ,"Journal of Differential Equations 242(2007)269-286.
- 6. Zhijian Yang and Pengyan Ding,"Longtime dynamics of the Kirchhoff equation with strong damping and critical nonlinearity on \mathbb{R}^N , "J.Math.Anal.Appl.435(2016)1826-1851.
- 7. Claudianor O.Alves and Giovany M.Figueiredo, "Nonlinear perturbations of a periodic Kirchhoff equation in \mathbb{R}^N , "Nonlinear Analysis 75(2012)2750"C2759.
- 8. Yang Zhijian and Jin Baoxia, "Global attractor for a class of Kirchhoff models, "Journal of Mathematical Physics 50,032701(2009).
- 9. Penghui Lv; Ruijin Lou and Guoguang Lin, "Global attractor for a class of nonlinear generalized Kirchhoff-Boussinesq model, "International Journal of Modern Nonlinear Theory and Application, 2016, 5,82-92.
- 10. Ruijin Lou; Penghui Lv and Guoguang Lin, "Global Attractors for a Class of Generalized Nonlinear Kirchhoff-Sine-Gordon Equation, "International Journal of Modern Nonlinear Theory and Application, 2016, 5, 73-81.
- 11. Lin Guoguang, Nonlinear evolution equation, Yunnan University Press, 2011.