INTERPLAY OF WIENER-HOPF AND HANKEL OPERATORS WITH ALMOST PERIODIC FOURIER SYMBOLS ON STANDARD AND VARIABLE EXPONENT LEBESGUE SPACES[†]

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ABSTRACT. Wiener-Hopf plus Hankel and Wiener-Hopf minus Hankel operators in both frameworks of standard and variable exponent Lebesgue spaces are considered in this paper. The main aim is to describe certain dependencies between the Fredholm property of some Wiener-Hopf operators acting between variable exponent Lebesgue spaces and the invertibility of Wiener-Hopf plus and minus Hankel operators on all the standard Lebesgue spaces. Different types of Fourier symbols will be used but special focus will be considered on the Wiener subclass of almost periodic matrix functions. In the first part of the paper we will give a survey of investigations on related results. This will be useful at the end of the paper to derive the above mentioned dependencies between the operators under study.

1. INTRODUCTION

Although being obviously more general than the standard Lebesgue spaces, variable exponent Lebesgue spaces preserve many of the properties of standard Lebesgue spaces. The most standard properties of variable exponent Lebesgue spaces may be seen in [12]. Namely, they are Banach spaces [12, Theorem 2.5], the Hölder inequality holds [12, Theorem 2.1], it is known a characterization for them to be reflexive [12, Corollary 2.7], and it is also known a condition under which continuous functions are dense [12, Theorem 2.11]. Anyway, variable exponent Lebesgue spaces appeared in the literature for the first time already in a 1931 article by W. Orlicz [18]. This allowed the consideration of the so-called modular spaces which were systematically studied by H. Nakano [15, 16].

Following the development in variable exponent Lebesgue spaces (see the monographs [8, 12]), the research on operators defined between these spaces have recently increased significantly (cf., e.g., [6, 7]). Obviously, this facilitates the possibility to better understand applied problems e.g. in elasticity, fluid dynamics,

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calculus of variations and differential equations (cf., e.g., [19] for specific applications in electrorheological fluids). One of the difficulties in the development of the operator theory in the $L^{p(\cdot)}$ spaces is that the convolution operators are, in general, not bounded in these spaces. Some *invertibility properties* of Wiener-Hopf and Hankel operators – being well-known for standard Lebesgue spaces (cf., e.g., [1, 3, 4, 17]) – are nowadays interestingly analysed in the framework of variable exponent Lebesgue spaces.

In the present paper we will be devoted to analyse consequences and dependencies of the Fredholm and invertibility properties of Wiener-Hopf, Wiener-Hopf plus Hankel and Wiener-Hopf minus Hankel operators (in short, *Wiener-Hopf-Hankel operators*) on variable exponent Lebesgue spaces and also on standard Lebesgue spaces.

In view to define these operators in a formal way, we will start by the definition of the variable exponent Lebesgue spaces. Let $p : \mathbb{R} \to [1, \infty]$ be measurable a.e. finite function. We denote by $L^{p(\cdot)}(\mathbb{R})$ the set of all complex-valued functions fon \mathbb{R} such that

$$I_{p(\cdot)}\left(\frac{f}{\lambda}\right) := \int_{\mathbb{R}} \left|\frac{f(x)}{\lambda}\right|^{p(x)} dx < \infty$$

for some $\lambda > 0$. This set becomes a Banach space when equipped with the norm

$$||f||_{p(\cdot)} := \inf\{\lambda > 0 : I_{p(\cdot)}(f/\lambda) \le 1\}.$$

The space $L^{p(\cdot)}(\mathbb{R})$ is precisely what we are referring as the variable exponent Lebesgue space.

We will suppose that

$$1 < p_{-} := \operatorname{ess}_{x \in \mathbb{R}} \inf p(x) \le \operatorname{ess}_{x \in \mathbb{R}} \sup p(x) =: p_{+} < \infty.$$
(1.1)

Under these conditions, the space $L^{p(\cdot)}(\mathbb{R})$ is separable and reflexive, and its dual space is isomorphic to $L^{p'(\cdot)}(\mathbb{R})$, where $p'(\cdot)$ is the conjugate exponent function defined by

$$\frac{1}{p(x)} + \frac{1}{p'(x)} = 1 \ (x \in \mathbb{R}).$$

Additionally, with condition (1.1) we have that $\|\phi I\|_{\mathcal{L}(L^{p(\cdot)}(\mathbb{R}))} \leq \|\phi\|_{L^{\infty}(\mathbb{R})}$ for a function $\phi \in L^{\infty}(\mathbb{R})$.

Moreover, $L^{p(\cdot)}(\mathbb{R}_+)$ denotes the variable Lebesgue space of complex-valued functions on the positive half-line $\mathbb{R}_+ = (0, +\infty)$, and the subspace of $L^{p(\cdot)}(\mathbb{R})$ formed by all functions supported in the closure of \mathbb{R}_+ is being denoted by $L^{p(\cdot)}_+(\mathbb{R})$.

Given a Banach space X, we denote by X^N the Banach space of all columns of height N with components in X. The norm in X^N is defined by

$$\|(x_1,\ldots,x_N)^{\mathsf{T}}\|_{X^N} = \left(\sum_{j=1}^N \|x_j\|_X^2\right)^{1/2}.$$

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Given a subalgebra B of $L^{\infty}(\mathbb{R})$, we denote by $B^{N \times N}$ the algebra of all $N \times N$ matrices Φ with entries in B. Moreover, we equip $B^{N \times N}$ with the norm

$$\|\Phi\|_{B^{N\times N}} = \|(\Phi_{jk})_{j,k=1}^N\|_{B^{N\times N}} = \left(\sum_{j,k=1}^N \|\Phi_{jk}\|_B^2\right)^{1/2}$$

We are now in condition to identify in a mathematical way the main objects of this work. We will consider Wiener-Hopf plus and minus Hankel operators acting between Lebesgue spaces with variable exponent $p(\cdot)$ (and also with fixed standard exponent $p(\cdot) = p \in (1, +\infty)$), denoted by

$$W_{\Phi} \pm H_{\Phi} : [L_{+}^{p(\cdot)}(\mathbb{R})]^{N} \to [L^{p(\cdot)}(\mathbb{R}_{+})]^{N}$$

with W_{Φ} and H_{Φ} being Wiener-Hopf and Hankel operators defined by

$$W_{\Phi} = r_{+} \mathcal{F}^{-1} \Phi \cdot \mathcal{F}, \qquad H_{\Phi} = r_{+} \mathcal{F}^{-1} \Phi \cdot \mathcal{F}J, \qquad (1.2)$$

respectively. Here, r_+ represents the operator of restriction from $[L^{p(\cdot)}(\mathbb{R})]^N$ into $[L^{p(\cdot)}(\mathbb{R}_+)]^N$, \mathcal{F}^{-1} denotes the inverse of the Fourier transformation \mathcal{F} , Φ is the socalled Fourier symbol, and $J : [L^{p(\cdot)}_+(\mathbb{R})]^N \to [L^{p(\cdot)}(\mathbb{R})]^N$ is the reflection operator given by the rule $J\varphi(x) = \widetilde{\varphi}(x) = \varphi(-x)$ which throughout the paper will be always considered for even functions $p(\cdot)$ (so that J will therefore be a bounded operator in those variable exponent Lebesgue spaces).

One of the goals of this work is to exhibit conditions under which Wiener-Hopf plus and minus Hankel operators are both invertible in the full scale of standard Lebesgue spaces, in the case when the Fourier symbol Φ belongs to the so-called Wiener subclass of the almost periodic matrix functions (which we will define later on in detail).

2. Preliminaries

The boundedness of a wide variety of operators (and in particular of Wiener-Hopf and Hankel operators) follows from the boundedness of the maximal operator on variable exponent Lebesgue spaces.

Given $f \in L^1_{\text{loc}}(\mathbb{R})$, the Hardy-Littlewood maximal operator M is defined by

$$(Mf)(x) := \sup_{x \in \Omega} \frac{1}{|\Omega|} \int_{\Omega} |f(y)| dy,$$

where the supremum is taken over all intervals $\Omega \subset \mathbb{R}$ containing x, and the Cauchy singular integral operator S is defined by

$$(Sf)(x) := \frac{1}{\pi i} \int_{\mathbb{R}} \frac{f(\tau)}{\tau - x} d\tau$$

where the integral is understood in the principal value sense.

Theorem 2.1. (cf. e.g. [9, Theorem 2.1.]) Let $p : \mathbb{R} \to [1, \infty]$ be a measurable function satisfying (1.1). If the Hardy-Littlewood maximal operator M is bounded on $L^{p(\cdot)}(\mathbb{R})$, then the Cauchy singular integral operator S is bounded on $L^{p(\cdot)}(\mathbb{R})$.

The following result states a sufficient condition on $p(\cdot)$ for M to be bounded on $L^{p(\cdot)}(\mathbb{R})$. **Theorem 2.2.** (cf. e.g. [14, Theorem 2.5.]) Let $p : \mathbb{R} \to [1, \infty]$ satisfy (1.1). In addition, suppose that $p(\cdot)$ satisfies

$$|p(x) - p(y)| \le \frac{C}{-\log|x - y|}, \qquad |x - y| \le \frac{1}{2},$$

and

$$|p(x) - p(y)| \le \frac{C}{\log(e + |x|)}, \qquad |x| \le |y|.$$

Then, the Hardy-Littlewood maximal operator is bounded on $L^{p(\cdot)}(\mathbb{R})$.

Let $\mathcal{P}(\mathbb{R})$ denote the class of exponents $p : \mathbb{R} \to [1, \infty]$ continuous on \mathbb{R} satisfying (1.1), and let $\mathcal{B}(\mathbb{R})$ denote the set of all $p(\cdot) \in \mathcal{P}(\mathbb{R})$ such that M is bounded on $L^{p(\cdot)}(\mathbb{R})$.

A function $\phi \in L^{\infty}(\mathbb{R})$ $(p(\cdot) \in \mathcal{B}(\mathbb{R}))$ is called a *Fourier multiplier* on $L^{p(\cdot)}(\mathbb{R})$ if the operator $\mathcal{F}^{-1}\phi\mathcal{F}$ acting on $L^2(\mathbb{R}) \cap L^{p(\cdot)}(\mathbb{R})$, extends uniquely to a bounded operator on $L^{p(\cdot)}(\mathbb{R})$. The set of all Fourier multipliers on $L^{p(\cdot)}(\mathbb{R})$ is denoted by $\mathcal{M}_{p(\cdot)}$.

It follows that for $\Phi \in [\mathcal{M}_{p(\cdot)}]^{N \times N}$ (with an even $p(\cdot) \in \mathcal{B}(\mathbb{R})$) the Wiener-Hopf and the Hankel operators defined in (1.2), $W_{\Phi} : [L_{+}^{p(\cdot)}(\mathbb{R})]^{N} \to [L^{p(\cdot)}(\mathbb{R}_{+})]^{N}$ and $H_{\Phi} : [L_{+}^{p(\cdot)}(\mathbb{R})]^{N} \to [L^{p(\cdot)}(\mathbb{R}_{+})]^{N}$, are bounded. These are in fact necessary and sufficient conditions for the Wiener-Hopf plus and minus Hankel operators to be bounded in variable exponent Lebesgue spaces. In a sense, these conditions are not explicit. For some more explicit conditions which determine the identification of subsets of $\mathcal{M}_{p(\cdot)}$, we would like to refer the reader to [10, §4.2] where, in particular, a generalization of the classical Mikhlin's theorem on Fourier multipliers (cf. [13]) to variable exponent Lebesgue spaces is presented.

In what follows, we will be also using the projection P and its complementary projection Q (Q = I - P), where

$$P = \mathcal{F}\ell_0 r_+ \mathcal{F}^{-1},$$

with $\ell_0: [L^{p(\cdot)}(\mathbb{R}_+)]^N \to [L^{p(\cdot)}_+(\mathbb{R})]^N$ being the zero extension operator.

Let $\dot{\mathbb{R}} := \mathbb{R} \cup \{\infty\}$. We will denote by $C(\dot{\mathbb{R}})$ the set of all continuous functions φ on the real line for which the two limits

$$\varphi(-\infty) := \lim_{x \to -\infty} \varphi(x), \qquad \varphi(+\infty) := \lim_{x \to +\infty} \varphi(x)$$

exist and coincide. The common value of these two limits will be denoted by $\varphi(\infty)$. Furthermore, $C_0(\dot{\mathbb{R}})$ will represent the collection of all $\varphi \in C(\dot{\mathbb{R}})$ for which $\varphi(\infty) = 0$.

For a continuous function $f : \mathbb{R} \to \mathbb{C}$ and $\mathcal{C} \subset \mathbb{R}$, let

$$\operatorname{osc}(f, \mathcal{C}) := \sup_{t, \tau \in \mathcal{C}} |f(t) - f(\tau)|.$$

We denote by SO the class of *slowly oscillating* functions given by

$$SO := \left\{ f \in C(\mathbb{R}) : \lim_{x \to +\infty} \operatorname{osc}(f, [-2x, -x] \cup [x, 2x]) = 0 \right\} \cap L^{\infty}(\mathbb{R}).$$

SO is a unital C*-algebra of $L^{\infty}(\mathbb{R})$ that contains $C(\mathbb{R})$.

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In what follows, we will use the notation $\mathcal{G}B$ for the group of all invertible elements of a Banach algebra B.

3. Relations between different classes of convolution type operators

In order to relate operators and to transfer certain operator properties between different types of operators, we will also use the notion of equivalence relation and equivalence after extension relation between bounded linear operators. We will now recall these concepts in detail.

Consider two bounded linear operators $T: X_1 \to X_2$ and $S: Y_1 \to Y_2$, acting between Banach spaces. The operators T and S are said to be *equivalent* if there are two boundedly invertible linear operators, $E: Y_2 \to X_2$ and $F: X_1 \to Y_1$, such that

$$T = E \ S \ F. \tag{3.1}$$

It directly follows from (3.1) that if two operators are equivalent, then they belong to the same *invertibility class*. More precisely, one of these operators is invertible, left invertible, right invertible or only generalized invertible, if and only if the other operator enjoys the same property.

We say that T is equivalent after extension to S (and write $T \stackrel{*}{\sim} S$) if there are Banach spaces Z_1 and Z_2 and invertible bounded linear operators E and F such that

$$\begin{bmatrix} T & 0 \\ 0 & I_{Z_1} \end{bmatrix} = E \begin{bmatrix} S & 0 \\ 0 & I_{Z_2} \end{bmatrix} F,$$

where I_{Z_1} and I_{Z_2} represent the identity operators in Z_1 and Z_2 , respectively. In case that Z_1 or Z_2 are the trivial space, we will say that the operators T and S are equivalent after one-sided extension.

If two operators are equivalent after extension, then they belong to the same invertibility class.

The following theorem is well-known (cf. [2]) when considering L^{∞} Fourier symbols in Wiener-Hopf and Hankel operators acting between standard L^{p} -spaces. Here, we present its generalization by including the variable exponent Lebesgue spaces case.

Theorem 3.1. Let $\Phi \in \mathcal{G}[\mathcal{M}_{p(\cdot)}]^{N \times N}$ with an even $p(\cdot) \in \mathcal{B}(\mathbb{R})$ in the case of variable exponent Lebesgue spaces or simply $\Phi \in \mathcal{G}[\mathcal{M}_p]^{N \times N}$ with $p(\cdot) = p \in (1, +\infty)$ in the case of standard Lebesgue spaces. The following two operators are equivalent after extension:

(i)
$$\begin{bmatrix} W_{\Phi} + H_{\Phi} & 0\\ 0 & W_{\Phi} - H_{\phi} \end{bmatrix} : [L_{+}^{p(\cdot)}(\mathbb{R})]^{2N} \to [L^{p(\cdot)}(\mathbb{R}_{+})]^{2N};$$

(ii) $W_{\Phi\widetilde{\Phi^{-1}}} : [L_{+}^{p(\cdot)}(\mathbb{R})]^{N} \to [L^{p(\cdot)}(\mathbb{R}_{+})]^{N}.$

Proof. We start by noticing that the operators

$$W_{\Phi} \pm H_{\Phi} : [L^{p(\cdot)}_{+}(\mathbb{R})]^{N} \to [L^{p(\cdot)}(\mathbb{R}_{+})]^{N}$$

and

$$P\phi P \pm P\phi JP + Q : [L^{p(\cdot)}(\mathbb{R})]^N \to [L^{p(\cdot)}(\mathbb{R})]^N$$

are equivalent after extension. Indeed, due to the identity

$$P\Phi P \pm P\Phi JP + Q = (\Phi P \pm \Phi JP + Q)(I - Q\Phi P \mp Q\Phi JP),$$

where the operator $I - Q\Phi P \mp Q\Phi JP$ is invertible (with inverse being given by $I + Q\Phi P \mp Q\Phi JP$, it follows that the operator in (i) is equivalent after extension to the operator

$$\begin{bmatrix} \Phi P + \Phi J P + Q & 0\\ 0 & \Phi P - \Phi J P + Q \end{bmatrix} : [L^{p(\cdot)}(\mathbb{R})]^{2N} \to [L^{p(\cdot)}(\mathbb{R})]^{2N}.$$

Observing that

$$\frac{1}{2}\begin{bmatrix}I & I\\ J & -J\end{bmatrix}\begin{bmatrix}\Phi P + \Phi J P + Q & 0\\ 0 & \Phi P - \Phi J P + Q\end{bmatrix}\begin{bmatrix}I & J\\ I & -J\end{bmatrix}$$
$$= \begin{bmatrix}\Phi & 0\\ \widetilde{\Phi} & I\end{bmatrix}P + \begin{bmatrix}I & \Phi\\ 0 & \widetilde{\Phi}\end{bmatrix}Q,$$

and multiplying this operator from the left with the invertible operator

$$A := \begin{bmatrix} I & -\Phi \widetilde{\Phi^{-1}} I \\ 0 & \widetilde{\Phi^{-1}} I \end{bmatrix},$$

we obtain the equivalent operator $\Psi P + Q$, where

$$\Psi = \begin{bmatrix} 0 & -\Phi \widetilde{\Phi^{-1}} \\ & \\ 1 & \widetilde{\Phi^{-1}} \end{bmatrix}.$$

Since

$$\Psi P + Q = (P\Psi P + Q)(I + Q\Psi P), \qquad (3.2)$$

where the operator on the right-hand side of (3.2) is invertible (with inverse being given by $I - Q\Psi P$, it follows that the operator in (ii) is equivalent to the operator $P\Psi P + Q$ which in its turn is equivalent after extension to the operator $P\Psi P$.

Additionally,

$$P\begin{bmatrix} 0 & -\Phi\overline{\Phi^{-1}} \\ 1 & \overline{\Phi^{-1}} \end{bmatrix} P = P\begin{bmatrix} \Phi\overline{\Phi^{-1}} & 0 \\ 0 & I \end{bmatrix} P\begin{bmatrix} 0 & -I \\ I & \overline{\Phi^{-1}} \end{bmatrix} P.$$
(3.3)

Since

$$B := P \begin{bmatrix} 0 & -I \\ I & \widetilde{\Phi^{-1}} \end{bmatrix} P$$

is invertible, (3.3) shows an explicit equivalence after (one-sided) extension between $P\Psi P$ and $P\Phi \widetilde{\Phi^{-1}}P$. Due to the fact that $P\Phi \widetilde{\Phi^{-1}}P$ is equivalent to $W_{\Phi \widetilde{\Phi^{-1}}}$, altogether, it follows that the operators in (i) and (ii) are equivalent after extension.

Finally, for the case of variable exponent Lebesgue spaces, we stress that all the just used operators are bounded in the variable exponent Lebesgue spaces due to the hypothesis $\Phi \in \mathcal{G}[\mathcal{M}_{p(\cdot)}]^{N \times N}$ and the consideration of an even $p(\cdot) \in \mathcal{B}(\mathbb{R})$.

4. MATRIX-VALUED AP SYMBOLS

The smallest closed subalgebra of $L^{\infty}(\mathbb{R})$ that contains all the functions e_{λ} ($\lambda \in \mathbb{R}$), where $e_{\lambda}(x) = e^{i\lambda x}$, $x \in \mathbb{R}$, is denoted by AP and called the algebra of *almost periodic functions*:

$$AP := alg_{L^{\infty}(\mathbb{R})} \{ e_{\lambda} : \lambda \in \mathbb{R} \}.$$

In addition, we use the notation

$$AP_{+} := alg_{L^{\infty}(\mathbb{R})} \{ e_{\lambda} : \lambda \ge 0 \}, \qquad AP_{-} := alg_{L^{\infty}(\mathbb{R})} \{ e_{\lambda} : \lambda \le 0 \}$$

for these two subclasses of AP (which are still closed subalgebras of $L^{\infty}(\mathbb{R})$).

We will likewise make use of the Wiener subclass of AP (denoted by APW) formed by all those elements from AP which allow a representation by an absolutely convergent series. Therefore, APW is precisely the (proper) subclass of all functions $\varphi \in AP$ which can be written in an absolutely convergent series of the form:

$$\varphi = \sum_j \varphi_j \; e_{\lambda_j} \; ,$$

where $\lambda_j \in \mathbb{R}$ and $\sum_j |\varphi_j| < \infty$.

We recall that all AP functions have a well-known mean value. The existence of such a number is provided in the next standard result from the theory of almost periodic functions (cf., e.g., [5]).

Proposition 4.1. Let $A \subset (0, \infty)$ be an unbounded set and let

$$\{I_{\alpha}\}_{\alpha\in A} = \{(x_{\alpha}, y_{\alpha})\}_{\alpha\in A}$$

be a family of intervals $I_{\alpha} \subset \mathbb{R}$ such that $|I_{\alpha}| = y_{\alpha} - x_{\alpha} \to \infty$ as $\alpha \to \infty$. If $\varphi \in AP$, then the limit

$$M(\varphi) := \lim_{\alpha \to \infty} \frac{1}{|I_{\alpha}|} \int_{I_{\alpha}} \varphi(x) \, dx$$

exists, is finite, and is independent of the particular choice of the family $\{I_{\alpha}\}$.

For any $\varphi \in AP$, the number that has been just introduced $M(\varphi)$ is called the *Bohr mean value* or simply the *mean value* of φ . In the matrix case the *mean value* is defined entry-wise.

Let $APW^ (APW^+)$ be the set of all functions $\psi \in APW$ such that $\Omega(\psi) \subset (-\infty, 0]$ $(\Omega(\psi) \subset [0, +\infty)$, respectively). Here $\Omega(\psi) = \{\lambda \in \mathbb{R} : M(\psi e_{-\lambda}) \neq 0\}$ is the *Bohr-Fourier spectrum* of ψ . Naturally, $APW^- \subset AP^-$ and $APW^+ \subset AP^+$.

Let us now recall the so-called right AP and APW factorizations and some other notions associated with them.

Definition 4.2. A matrix function $\Phi \in \mathcal{G}AP^{N \times N}$ is said to admit a *right* AP *factorization* if it can be represented in the form

$$\Phi(x) = \Phi_{-}(x)D(x)\Phi_{+}(x)$$

for all $x \in \mathbb{R}$, with

$$\Phi_{-} \in \mathcal{G}AP_{-}^{N \times N}, \quad \Phi_{+} \in \mathcal{G}AP_{+}^{N \times N}, \tag{4.1}$$

and D is a diagonal matrix of the form $D(x) = \text{diag}[e^{i\lambda_1 x}, \dots, e^{i\lambda_N x}], \lambda_j \in \mathbb{R}$. The numbers λ_j are called the *right AP indices* of the factorization. A right *AP* factorization with $D = I_{N \times N}$ is referred to be a canonical right *AP* factorization. In the case when (4.1) is substituted by

$$\Phi_{-} \in \mathcal{G}APW_{-}^{N \times N}, \quad \Phi_{+} \in \mathcal{G}APW_{+}^{N \times N},$$

we will be in the presence of a right APW factorization which will also turn to be a canonical one in case of $D = I_{N \times N}$.

The vector containing the right AP indices will be denoted by $k(\Phi)$, i.e., in the above case $k(\Phi) := (\lambda_1, \ldots, \lambda_N)$. If we consider the case with equal *right* AP indices $(k(\Phi) := (\lambda_1, \lambda_1, \ldots, \lambda_1))$, then the matrix $\mathbf{d}(\Phi) := M(\Phi_-)M(\Phi_+)$ is independent of the particular choice of the right AP factorization. In this case, this matrix $\mathbf{d}(\Phi)$ is called the *geometric mean* of Φ .

The just presented factorizations allow the possibility to find conditions which induce the invertibility or, at least, the Fredholm property of Wiener-Hopf operators in standard Lebesgue spaces. Anyway, for $APW^{N\times N}$ even stronger results appear as it is the case of the following Kurbatov Theorem (cf. [11]).

Theorem 4.3. If $\Phi \in APW^{N \times N}$, then $W_{\Phi} : [L^p_+(\mathbb{R})]^N \to [L^p(\mathbb{R}_+)]^N$ is invertible for some $p \in (1, +\infty)$ if and only if $W_{\Phi} : [L^p_+(\mathbb{R})]^N \to [L^p(\mathbb{R}_+)]^N$ is invertible for all $p \in (1, +\infty)$.

5. Matrix-valued SAP symbols

The C^{*}-algebra SAP of all semi-almost periodic functions on \mathbb{R} is the smallest closed subalgebra of $L^{\infty}(\mathbb{R})$ that contains AP and $C(\overline{\mathbb{R}})$:

$$SAP = alg_{L^{\infty}(\mathbb{R})} \{AP, C(\overline{\mathbb{R}})\}.$$

In addition, it is possible to interpret the SAP functions in a different form due to the following characterization of D. Sarason [20].

Theorem 5.1. Let $u \in C(\overline{\mathbb{R}})$ be any function for which $u(-\infty) = 0$ and $u(+\infty) = 1$. 1. If $\varphi \in SAP$, then there is $\varphi_{\ell}, \varphi_r \in AP$ and $\varphi_0 \in C_0(\mathbb{R})$ such that

$$\varphi = (1 - u)\varphi_{\ell} + u\varphi_r + \varphi_0.$$

The functions $\varphi_{\ell}, \varphi_r$ are uniquely determined by φ , and independent of the particular choice of u. The maps

$$\varphi \mapsto \varphi_{\ell}, \qquad \varphi \mapsto \varphi_r$$

are C^* -algebra homomorphisms of SAP onto AP.

The last theorem is also valid in the matrix case. In the next section we will use the following result which allows a certain connection between the Fredholm property of some classes of singular integral operators on variable exponent Lebesgue spaces and consequent invertible operators on specific standard Lebesgue spaces. **Theorem 5.2.** ([9, Theorem 1.2]) Let $\Phi \in SAP^{N \times N}$ and $p(\cdot) \in \mathcal{B}(\mathbb{R}) \cap SO$. If the singular integral operator $\Phi P + Q$ is Fredholm on the variable Lebesgue space $[L^{p(\cdot)}(\mathbb{R})]^N$, then:

(a) there is an exponent q_r lying in the segment

$$\left[\lim_{x \to +\infty} \inf p(x), \lim_{x \to +\infty} \sup p(x)\right]$$

such that $\Phi_r P + Q$ is invertible on the standard Lebesgue space $[L^{q_r}(\mathbb{R})]^N$; (b) there is an exponent q_ℓ lying in the segment

$$\left[\lim_{x \to -\infty} \inf p(x), \lim_{x \to -\infty} \sup p(x)\right]$$

such that $\Phi_{\ell}P + Q$ is invertible on the standard Lebesgue space $[L^{q_{\ell}}(\mathbb{R})]^N$.

6. Wiener-Hopf and Wiener-Hopf-Hankel operators with APWFourier symbols in different Lebesgue spaces

Making use of the previous results and, in particular, the equivalence after extension relation presented in Theorem 3.1, we are in position to describe an interplay between the Fredholm property of a Wiener-Hopf operator acting on variable exponent Lebesgue spaces and the invertibility of some associated Wiener-Hopf plus and minus Hankel operators on standard Lebesgue spaces in all the possible range of the p Lebesgue constant exponents.

Theorem 6.1. Let $\Phi \in \mathcal{G}APW^{N \times N}$ and $p(\cdot) \in \mathcal{B}(\mathbb{R}) \cap SO$. If the Wiener-Hopf operator acting between variable exponent Lebesgue spaces

$$W_{\Phi\widetilde{\Phi^{-1}}}: [L^{p(\cdot)}_+(\mathbb{R})]^N \to [L^{p(\cdot)}(\mathbb{R}_+)]^N$$

has the Fredholm property, then the Wiener-Hopf plus Hankel and the Wiener-Hopf minus Hankel operators

$$W_{\Phi} \pm H_{\Phi} : [L^p_+(\mathbb{R})]^N \to [L^p(\mathbb{R}_+)]^N$$

are both invertible operators for all $p \in (1, +\infty)$.

Proof. Having in mind that $\Phi \in \mathcal{G}APW^{N \times N}$, we also have $\Phi \Phi^{-1} \in APW^{N \times N}$ and, moreover,

$$(\Phi \widetilde{\Phi^{-1}})_{\ell} = (\Phi \widetilde{\Phi^{-1}})_r = \Phi \widetilde{\Phi^{-1}}.$$

Therefore, from Theorem 5.2, we obtain that there is an exponent

$$q \in [p_-, p_+]$$

(cf. (1.1)) such that the singular integral operator $\Phi \widetilde{\Phi^{-1}}P + Q$ is invertible on the standard Lebesgue space $[L^q(\mathbb{R})]^N$.

However,

$$\Phi \widetilde{\Phi^{-1}} P + Q : [L^q(\mathbb{R})]^N \to [L^q(\mathbb{R})]^N$$

is equivalent after extension to $W_{\Phi\widetilde{\Phi^{-1}}}: [L^q_+(\mathbb{R})]^N \to [L^q(\mathbb{R}_+)]^N$. Thus, $W_{\Phi\widetilde{\Phi^{-1}}}$ is also invertible on the standard Lebesgue space $[L^q_+(\mathbb{R})]^N$.

Therefore, due to the Kurbatov Theorem,

$$W_{\Phi\widetilde{\Phi^{-1}}}: [L^p_+(\mathbb{R})]^N \to [L^p(\mathbb{R}_+)]^N$$

is an invertible operator for all $p \in (1, +\infty)$. Using the equivalence after extension relation, presented in Theorem 3.1, we obtain that the Wiener-Hopf plus and minus Hankel operators $W_{\Phi} + H_{\Phi} : [L^p_+(\mathbb{R})]^N \to [L^p(\mathbb{R}_+)]^N$ and $W_{\Phi} - H_{\Phi} : [L^p_+(\mathbb{R})]^N \to [L^p(\mathbb{R}_+)]^N$ are both invertible operators for all $p \in (1, +\infty)$. \Box

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