

We are IntechOpen, the world's leading publisher of Open Access books Built by scientists, for scientists

4,800

Open access books available

122,000

International authors and editors

135M

Downloads

Our authors are among the

154

Countries delivered to

TOP 1%

most cited scientists

12.2%

Contributors from top 500 universities



WEB OF SCIENCE™

Selection of our books indexed in the Book Citation Index
in Web of Science™ Core Collection (BKCI)

Interested in publishing with us?
Contact book.department@intechopen.com

Numbers displayed above are based on latest data collected.

For more information visit www.intechopen.com



Ferroelectric Domain Imaging Multiferroic Films Using Piezoresponse Force Microscopy

Hongyang Zhao, Hideo Kimura, Qiwen Yao,
Lei Guo, Zhenxiang Cheng and Xiaolin Wang

Additional information is available at the end of the chapter

<http://dx.doi.org/10.5772/52519>

1. Introduction

Recently, multiferroic materials with the magnetoelectric coupling of ferroelectric (or antiferroelectric) properties and ferromagnetic (or antiferromagnetic) properties have attracted a lot of attention.[1-4] Among them, BiFeO_3 (BFO) and YMnO_3 has been intensively studied. For such ABO_3 perovskite structured ferroelectric materials, they usually show antiferromagnetic order because the same B site magnetic element except BiMnO_3 is ferromagnet. While for the $\text{A}_2\text{BB}'\text{O}_6$ double perovskite oxides, the combination between B and B' give rise to a ferromagnetic coupling. They are also expected to be multiferroic materials. Several bismuth-based double perovskite oxides ($\text{BiBB}'\text{O}_6$) have aroused great interest like $\text{Bi}_2\text{NiMnO}_6$, $\text{La}_2\text{NiMnO}_6$, $\text{BiFeO}_3\text{-BiCrO}_3$. But as we know, few researches are focused on $\text{Bi}_2\text{FeMnO}_6$. We believe it is particular interesting to investigate $\text{Bi}_2\text{FeMnO}_6$ (BFM). The origin of the ferromagnetism in these compounds has been discussed in many reports. According to Goodenough-Kanamori's (GK) rules, many ferromagnets have been designed through the coupling of two B site ions with and without e_g electrons. In BFM, the 180 degree $-\text{Fe}^{3+}\text{-O-Mn}^{3+}$ bonds is quasistatic, partly because the strong Jahn-Teller uniaxial strain in an octahedral site. Because the complication of the double perovskite system, there are still some questions about the violation of GK rules in some cases and the origin of the ferromagnetism or antiferromagnetism. The other problem is the bad ferroelectric properties. In order to characterization of their ferroelectric/piezoelectric properties, scanning probe microscopy (SPM) techniques were used.

Multiferroic materials can be classified into two categories: one is single-phase materials; the other is multilayer or composite hetero-structures that contain more than one ferroic phases [5]. The most desirable multiferroic material is the intrinsic single-phase material, which is

still rarely produced although significant advancements had been made recently. Therefore, it is essential to broaden the searching field for new candidates of multiferroics. This work focuses on both types multiferroic materials: Single-phase $\text{Bi}_2\text{FeMnO}_6$ and multilayered $\text{YMnO}_3/\text{SnTiO}_3$. Scanning probe microscopy (SPM) techniques were used for ferroelectric domain imaging of these multiferroic materials.

As we know, most piezoelectric/ferroelectric materials with good performances are based on the perovskite-type oxides of ABO_3 , among them, the most extensively studied ones should be PbTiO_3 and $\text{Pb}(\text{ZrTi})\text{O}_3$ based materials due to their high dielectric constants and good piezoelectric/ferroelectric properties. However, these materials containing lead which leads to environmental problems. Thus, it is desirable to develop new lead-free piezoelectric materials to replace PZT based piezoelectrics for environmental protection. Through first principle calculations, SnTiO_3 containing Sn^{2+} ions was estimated to have excellent ferroelectric properties [6, 7]. The calculated results indicated that the spontaneous polarization and piezoelectric coefficients of SnTiO_3 is comparable with those of PbTiO_3 . Moreover, the most stable structure is tetragonal perovskite with $a=b=3.80\text{\AA}$ and $c=4.09\text{\AA}$ [6, 7]. The metastable SnTiO_3 phase is very difficult to obtain, a good way to stabilize the SnTiO_3 phase is to mix it with other compounds. Several studies have been focused on the synthesis of Sn-doped ceramics in BaTiO_3 system, such as $(\text{Ba}_{0.6}\text{Sr}_{0.4})(\text{Ti}_{1-x}\text{Sn}_x)\text{O}_3$ [8], $(\text{Ba}_{1-x-y}\text{Ca}_x\text{Sn}_y)\text{TiO}_3$ [9]. However, it is difficult to obtain Sn^{2+} by traditional bulk synthesis techniques. On the other hand, it is easier to fabricate SnTiO_3 using pulsed laser deposition method (PLD). Our interest was to find out whether the SnTiO_3 phase can be stabilized through the layer-by-layer deposition using PLD, and whether it is possible to obtain such metastable materials in non-equilibrium conditions.

The hexagonal manganite YMnO_3 , which shows an antiferromagnetic transition at $T_N = 75\text{ K}$, and a ferroelectric transition $T_C = 913\text{ K}$, is one of the rare existing multiferroics [10-12]. It was chosen as the basic composition to sandwich SnTiO_{3+x} phase and to form YST multilayer film. In this work, we reported the effect of substrate orientation and layer numbers of YMnO_3 and SnTiO_{3+x} on the piezoelectric/ferroelectric and magnetic properties in the designed layered YST system. It was hoped that our method would serve as a model system to introduce the use of PLD techniques in the growth of multilayer multiferroic materials and metastable materials. Such techniques could be promising for device design and the searching in fabricating new materials.

2. SPM system

SPM has emerged as a powerful tool for high-resolution characterization of ferroelectrics for the first time providing an opportunity for non-destructive visualization of ferroelectric domain structures at the nanoscale. [13] These nanostructures can be used for studying the intrinsic size effects in ferroelectrics as well as for addressing such technologically important issues as processing damage, interfacial strain, grain size, aspect ratio effect, edge effect, domain pinning and imprint. [14] The system composes of a conducting tip in con-

tact with the dielectric surface on a conductive substrate can be considered as a capacitor. SPM is a well-established field offering multiple opportunities for new discoveries and breakthroughs. [15-17]

Piezoelectric materials provide an additional response to the applied ac electric field due to the converse piezoelectric effect: $\Delta l = d_{33}V$, where Δl is the displacement, d_{33} is the effective longitudinal piezoelectric coefficient. [18, 19] It follows that both electrostatic and piezoelectric signals are linear with the applied voltage and thus contribute to the measured PFM response. These measurements are referred to as out-of-plane (OP) measurement. Measurements of local hysteresis loops are of great importance in inhomogeneous or polycrystalline ferroelectrics because they are able to quantify polarization switching on a scale significantly smaller than the grain size or inhomogeneity variation. Macroscopically, the switching occurs via the nucleation and growth of a large number of reverse domains in the situation where the applied electric field is uniform. Therefore, the d_{33} hysteresis reflects the switching averaged over the entire sample under the electrode. [20, 21] In the PFM experimental conditions, the electric field is strongly localized and inhomogeneous; therefore, the polarization switching starts with the nucleation of a single domain just under the tip.

3. Material designation and characterization

Multiferroic materials have been attracting considerable attention especially in the last ten years. [22-25] One of the most appealing aspects of multiferroics is their magnetoelectric coupling. Among them, the most studied is the perovskite BiFeO_3 (BFO) with room temperature multiferroic properties and YMnO_3 . Based on BFO and YMnO_3 , we designed new multiferroics and fabricated the films using pulsed laser deposition (PLD) method. The material system includes double perovskite multiferroic $\text{Bi}_2\text{FeMnO}_6$ (BFM) and multilayered $\text{YMnO}_3/\text{SnTiO}_{3+x}$ (YST). Compared to BFO, the magnetic properties of BFM and YST were greatly improved. However, until now there is no report about their ferroelectric properties because the difficulty of obtaining well-shaped polarization hysteresis loops. It is important to study ferroelectric properties because the possible coupling between ferroelectric and antiferromagnetic domains and its lead-free nature. As is well known, the ferroelectric property is mainly determined by the domain structures and domain wall motions. The most exciting recent developments in the field of multiferroics and the most promise for future discoveries are in interfacial phenomena. [26] The interfaces include those that emerge spontaneously and those that are artificially engineered. Therefore, the domain structure and polarization switching were studied in these three new multiferroic films using piezoresponse force microscopy (PFM).

The emerging technique of PFM is proved to be a powerful tool to study piezoelectric and ferroelectric materials in such cases and extensive contributions have been published. In PFM, the tip contacts with the sample surface and the deformation (expansion or contraction of the sample) is detected as a tip deflection. The local piezoresponse hysteresis loop and information on local ferroelectric behavior can be obtained because the strong coupling

between polarization and electromechanical response in ferroelectric materials. In the present study, we attempt to use PFM to study the ferroelectric/piezoelectric properties in films of BFM and YST. PFM response was measured with a conducting tip (Rh-coated Si cantilever, $k \sim 1.6 \text{ N m}^{-1}$) by an SII Nanotechnology E-sweep AFM. PFM responses were measured as a function of applied DC bias (V_{dc}) with a small ac voltage applied to the bottom electrode (substrate) in the contact mode, and the resulting piezoelectric deformations transmitted to the cantilever were detected from the global deflection signal using a lock-in amplifier.

3.1. Characterization of BFM film

Magnetism and ferroelectricity exclude each other in single phase multiferroics. It is difficult for designing multiferroics with good magnetic and ferroelectric properties. Our interest is to design new candidate multiferroics based on BiFeO_3 . BiFeO_3 is a well-known multiferroic material with antiferromagnetic Neel temperature of 643K, which can be synthesized in a moderate condition. [22, 23, 25] In contrast, BiMnO_3 is ferromagnetic with $T_c = 110\text{K}$ and it needs high-pressure synthesis. [27, 28] The possible magnetoelectric coupling has motivated a lot of studies on the ferroelectric and antiferromagnetic domains. The ferroelectric domain structures of BFO in the form of ceramics, single crystals and thin films have been intensively studied using PFM, TEM and other techniques.

Single phase BFM ceramics could be synthesized by conventional solid state method as the target. For BFM ceramics, the starting materials of Bi_2O_3 , Fe_3O_4 , MnCO_3 were weighed according to the molecular mole ratio with 10 mol% extra Bi_2O_3 . They were mixed, pressed into pellets and sintered at $800 \text{ }^\circ\text{C}$ for 3 h. Then the ceramics were crushed, ground, pressed into pellets and sintered again at $880 \text{ }^\circ\text{C}$ for 1 h. BFM films were deposited on SrTiO_3 (STO) substrate by pulsed laser deposition (PLD) method at 650°C with $500 \sim 600\text{mTorr}$ dynamic oxygen. [29, 30]

The structure of BFM was calculated [31] and it is connected with the magnetic configurations as shown in Figure 1. It has three possible space groups of $\text{Pm}3\text{m}$, $\text{R}3$ and $\text{C}2$ and the magnetic configurations were presumed for each structure symmetry to be G-type antiferromagnetic (G-AFM) and ferromagnetic (FM) structures. The most stable structure of BFM is monoclinic with $\text{C}2$ space group. Mn tends to show 3+ valence which will induce a large distortion because it is Jahn-Teller ion. The valence of Mn and Fe has been studied in the former work. [30, 32]

Figure 2 shows the results of the PFM images of BFM film. Several features could be observed: firstly, the obvious contrast could be seen and the grains in PFM and topography are correspondence; secondly, the existence of contrasts on both OP and IP indicates multiple orientations of domains; thirdly, the IP contrast is not so clear as OP contrast, that is to say, the suppression of the in-plane response for heterostructures suggesting a constrained ferroelectric domain-orientation along the OP direction. In the former paper, the ferroelectric domain switching and typical butterfly loops were observed. [32]

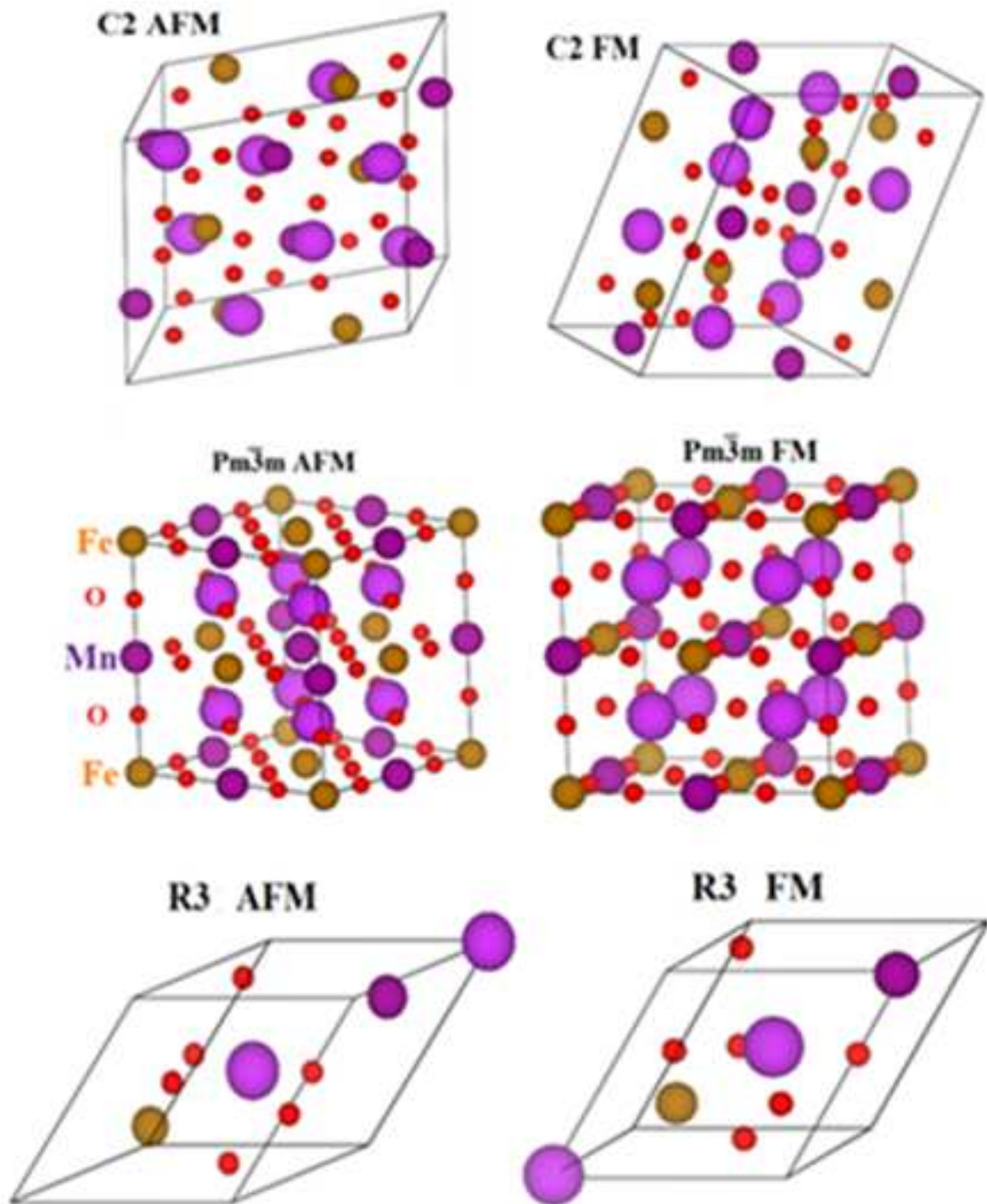


Figure 1. Calculated six structures of BFM.

The domains in BFO films have been analyzed in detail in Ref [19]: the bigger spontaneous ferroelectric domains were observed in BFO than in other ferroelectrics without multiferroic properties; the domains were irregular but the the model was predicted and consistent with the experimental results. According to the present results of BFM, further study is needed to obtain the domain morphology and do the calculation.

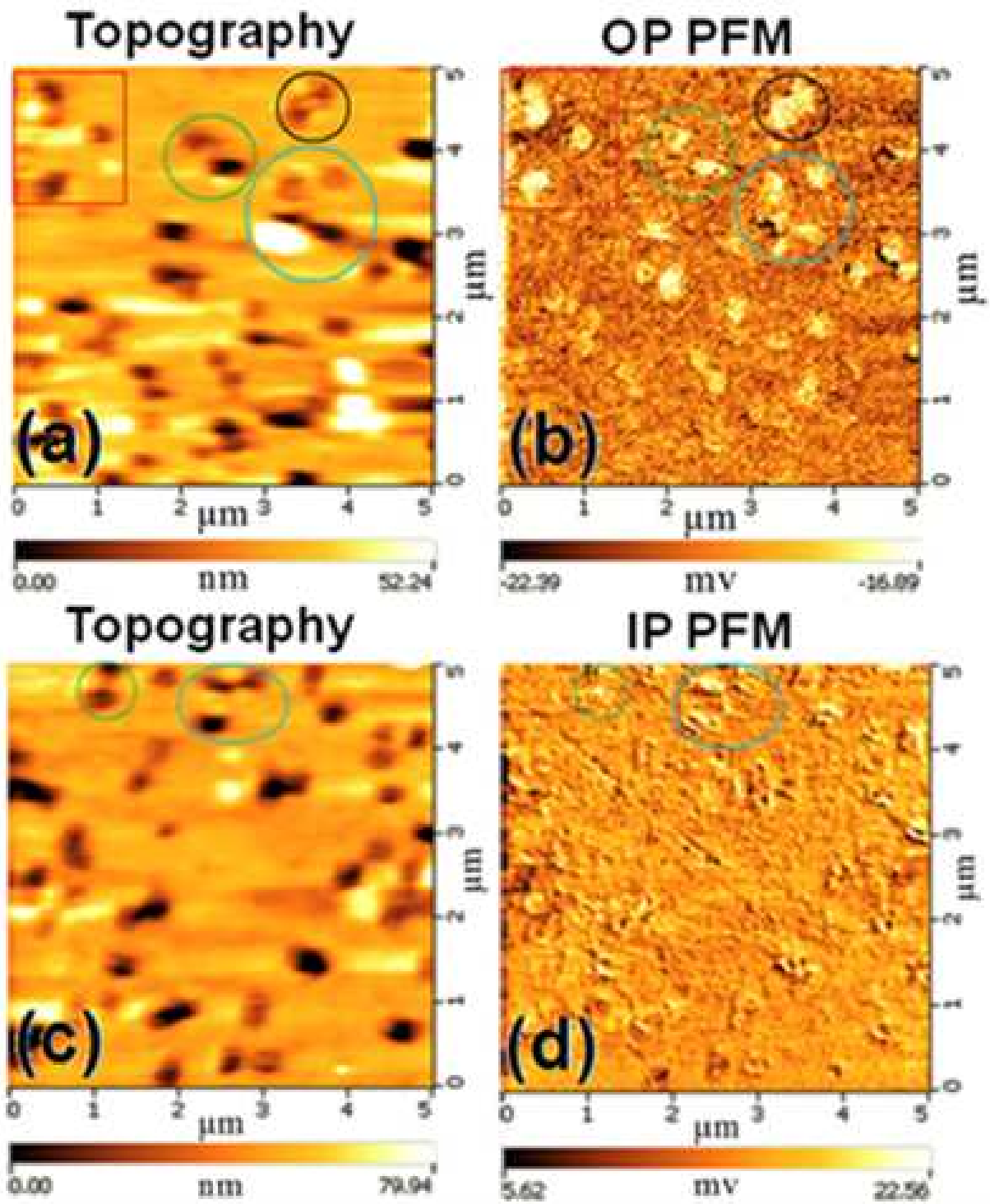


Figure 2. PFM images of BFM film

3.2. Characterization of YST film

Polycrystalline YMnO_3 and TiO_2 ceramics were synthesized by conventional solid state method as the targets. On addition the SnO and SnO_2 commercial targets were also used at the same time in our process. The deposition of these films includes the following major steps. YST films were deposited on (100) and (110) Nb: SrTiO_3 (STO) substrate using a pulsed laser deposition (PLD) system at 700°C with $10^{-1} \sim 10^{-5}$ Torr dynamic oxygen. The targets were alternately switched constantly and the films were obtained in a layer-by-layer growth mode. After deposition, the films were annealed in the same condition for 15 minutes at 700°C and then cooled to room temperature. [33]

In YST multilayered films, one layer is defined to be comprised of two sub-layers: (1) YMnO_3 and (2) SnTiO_{3+x} . The films deposited on (100), (111) and (110) STO are expressed as YST100, YST111 and YST110, respectively. The film YST110-4 denotes the films deposited on (110) STO with four layers, as shown clearly in Figure 3. The total deposition time for YMnO_3 and SnTiO_{3+x} is the same of 40 and 20 minutes, respectively.

XRD patterns for the five films of YST100-2, YST100-4 were shown in Figure 4. The peaks were identified using the XRD results of SnTiO_3 (\odot), YMnO_3 (Y) (shown in Figure 9, FeTiO_3 (∇) [34, 35]. SnTiO_3 is metastable and it showed two combined phases, one is FeTiO_3 , and the other is the good ferroelectric phase which has tetragonal structure. XRD patterns of SnTiO_3 (\odot) were shown in Figure 5 in supplementary materials, which is obtained using the calculated data[6, 7]. The symbol of “?” represents the phase which we cannot identify so far, it could be a peak from the (111) TiO_2 phase.

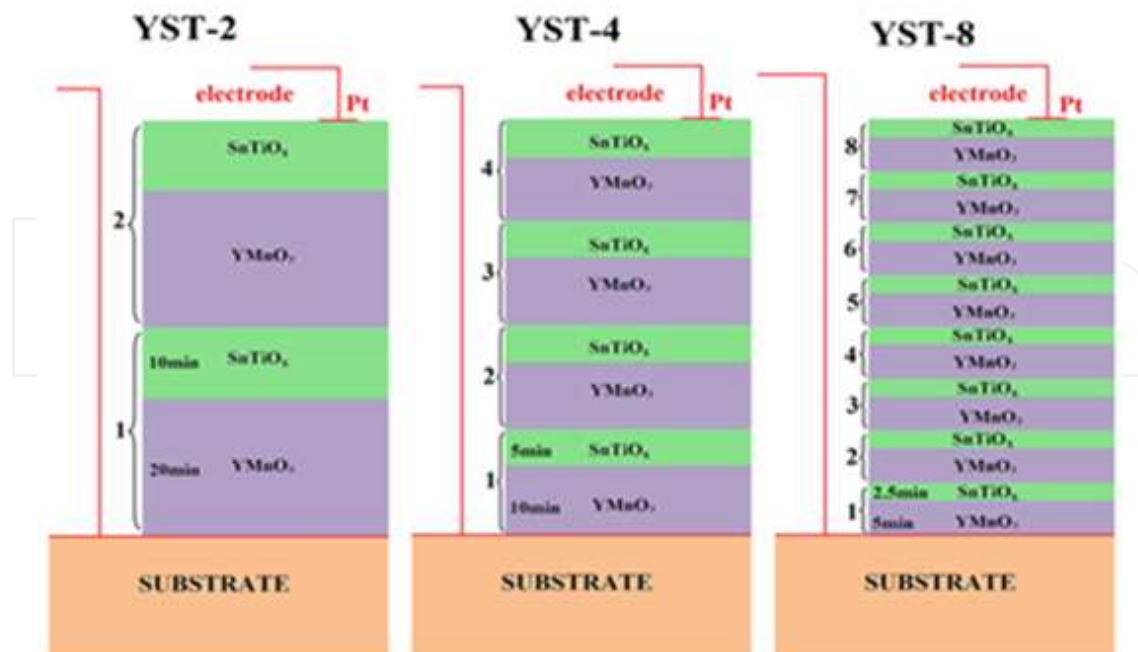


Figure 3. Schematic explanation of films with two, four and eight layers, each layer contains two sublayers of YMnO_3 and SnTiO_3 .

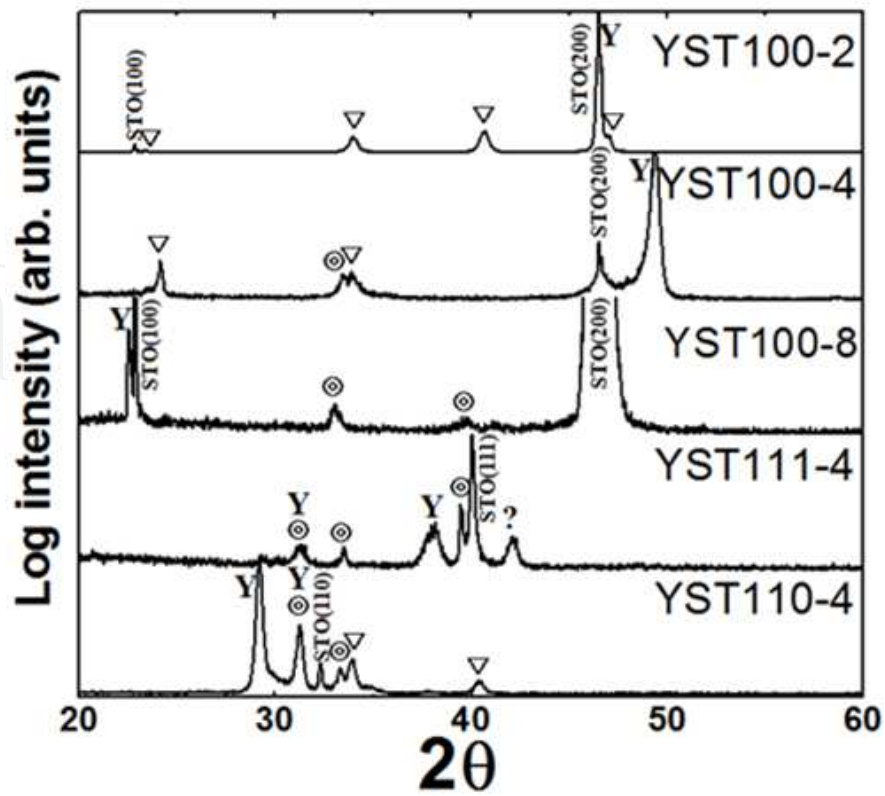


Figure 4. XRD patterns for the five films of YST100-2, YST100-4. [31]

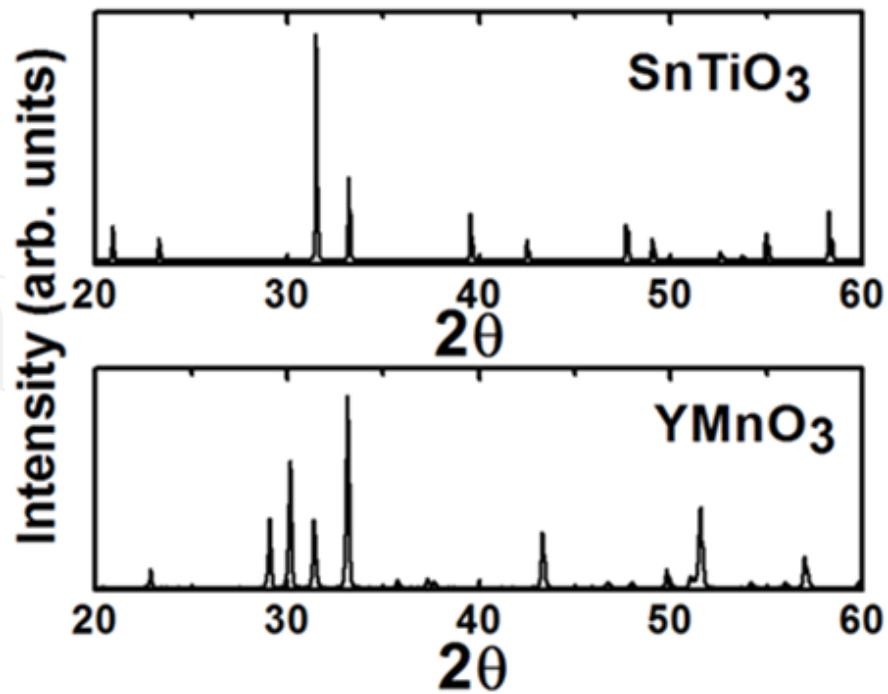


Figure 5. XRD patterns of calculated SnTiO₃ and YMnO₃ targets. [33]

Many reports about the YMnO_3 films can be found but most of them without ferroelectric characterization provided. This is mainly ascribed to the difficulty in obtaining a satisfactory ferroelectric measurement result. Figure 6 shows the electrical polarization hysteresis loops (P-E loops) of YST110-4 film. The ferroelectric type hysteresis loop was observed. It is obvious that the P-E loop is significantly different for different layer-number films, as well as for different substrate orientations. For the same STO orientation, the P-E loops were improved as the layer number increase, and the film fabricated on (110) STO shows improved properties compared to the film on (100) STO with the same layer number. As for the same deposition time for the four samples, it is found that more layers of SnTiO_3 in the YST system shows much better ferroelectric properties. It is suggested that the SnTiO_3 phase can exist only in ultra-thin thickness and can be stabilized by YMnO_3 sub-layers. That is to say, SnTiO_3 is indeed a ferroelectric materials but it is a great challenge [36] to obtain stable single phase SnTiO_3 films that is with good ferroelectric performance (it will be tried in our future works). Although the electric charge in the interface may affect its ferroelectric properties, the improved properties of YST110-4 proved that its effect was not significant. While the P-E loops were not improved in YST111-4, all the observations showed the anisotropy of the films.

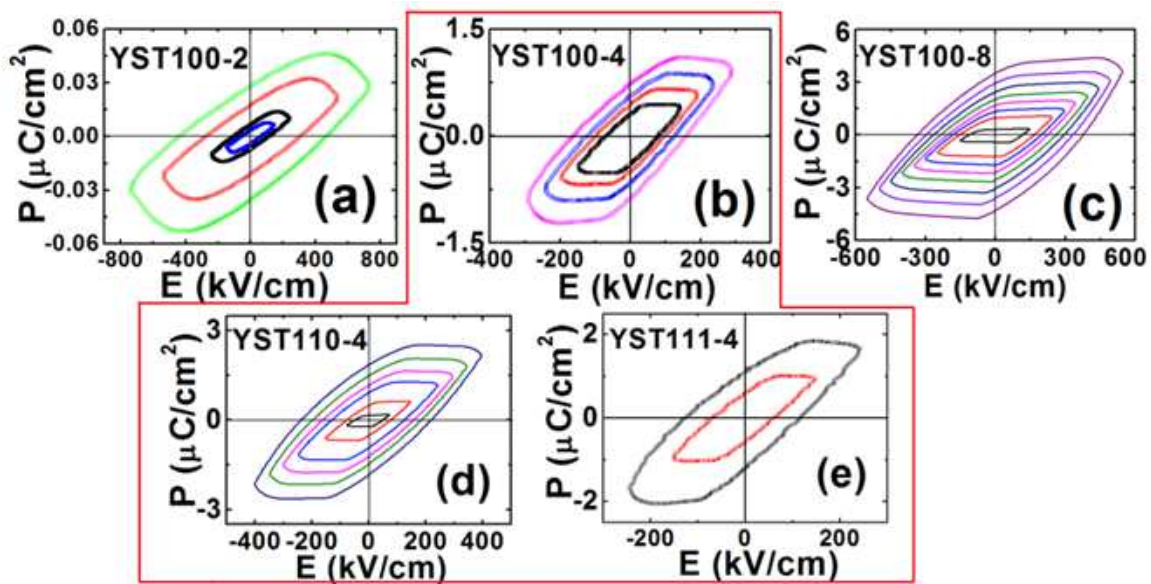


Figure 6. Ferroelectric polarization loop of YST films. [31]

Although well developed P-E loops were observed for YST100-8 and YST110-4, the loops still show the rounded features, indicating that contributions to the hysteresis loop from movable charges was significant. In addition, the P-E loops of the four films have revealed that the polarization is rather small (just a few $\mu\text{C}/\text{cm}^2$) and they do not exhibit saturation. The fatigue properties were shown in Figure 7 for YST110-4. The polarization measurement was at 21V, the switching voltage was 14V and the frequency was 5k Hz. The remnant polarization decreased about 40% after 10^7 read/write cycles for YST110-4. From all the comparisons, YST110-4 sample shows the comparable ferroelectric properties with YST100-8.

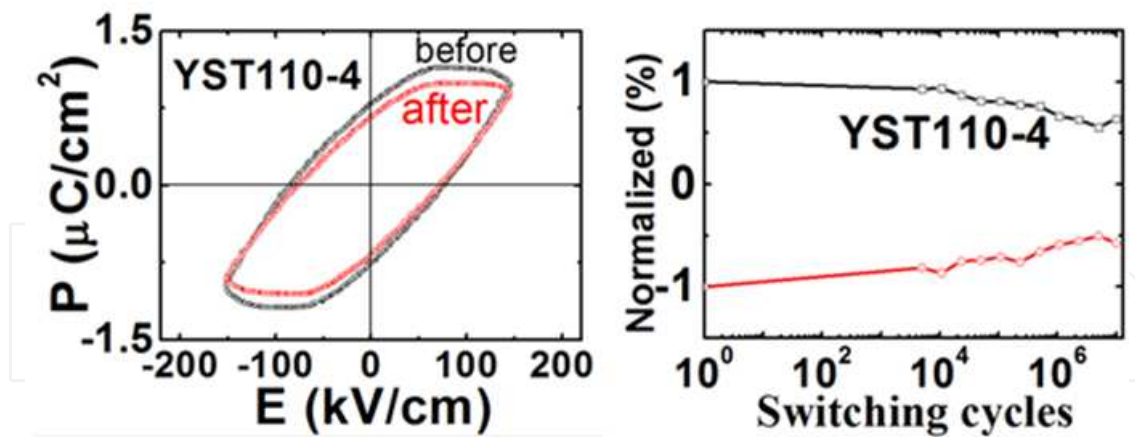


Figure 7. Fatigue measurement of YST110-4 film. [31]

Figure 8 shows TEM results of the YST110-4 film. It clearly shows four layers and one layer which is close to the substrate seems have some reaction with the substrate. The other three layers are homogeneous and we could see the domain structures through this image. The surface of the film was studied using AFM as shown in Figure 9. As expected, Figure 10 displays contrast over the polarized square after poled by positive and negative 10 V voltage, due to the different phases of the PFM response for the up and down domains. The obvious change of the contrast in YST110-4 film confirms that the polarization reversal is indeed possible and that the film is ferroelectric at room temperature.



Figure 8. TEM image of YST110-4 film.

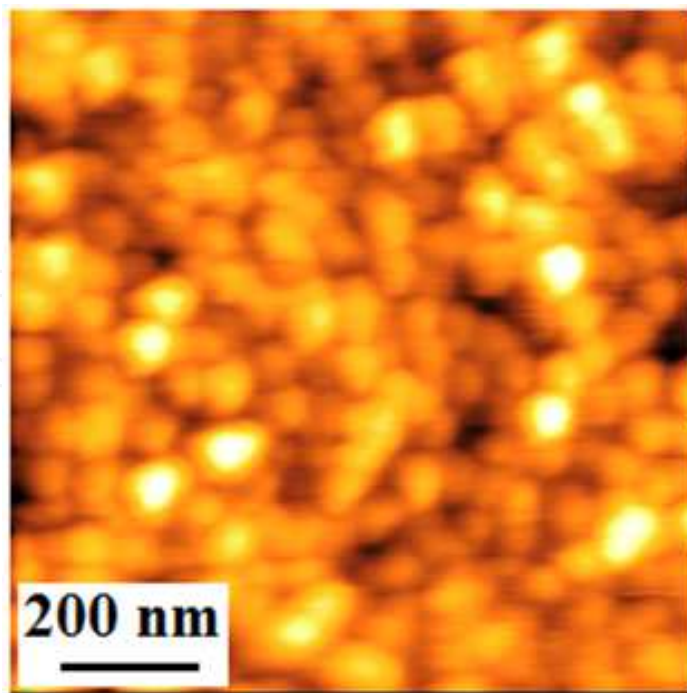


Figure 9. AFM topography image of YST110-4 film.

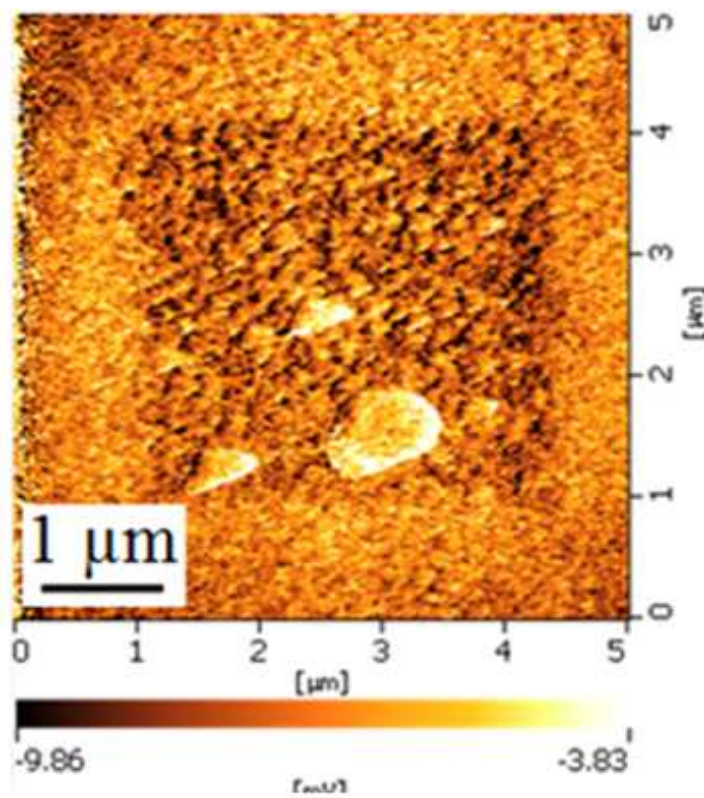


Figure 10. PFM image of YST110-4 film after poled using ± 10 V voltage

As seen in Figure 11, when a direct current voltage up to 10 V was applied on the samples, the sample exhibited “butterfly” loop. The loops were not symmetrical due to the asymmetry of the upper (tip) and bottom (substrate) electrodes. In addition, the substrate may also affect the d_{33} . According to the equation $d_{33} = \Delta l / V$, where Δl is the displacement, the effective d_{33} could be calculated. At the voltage of 10 V, both samples show the maximum effective d_{33} of about 6.21 pm/V for YST110-4 sample.

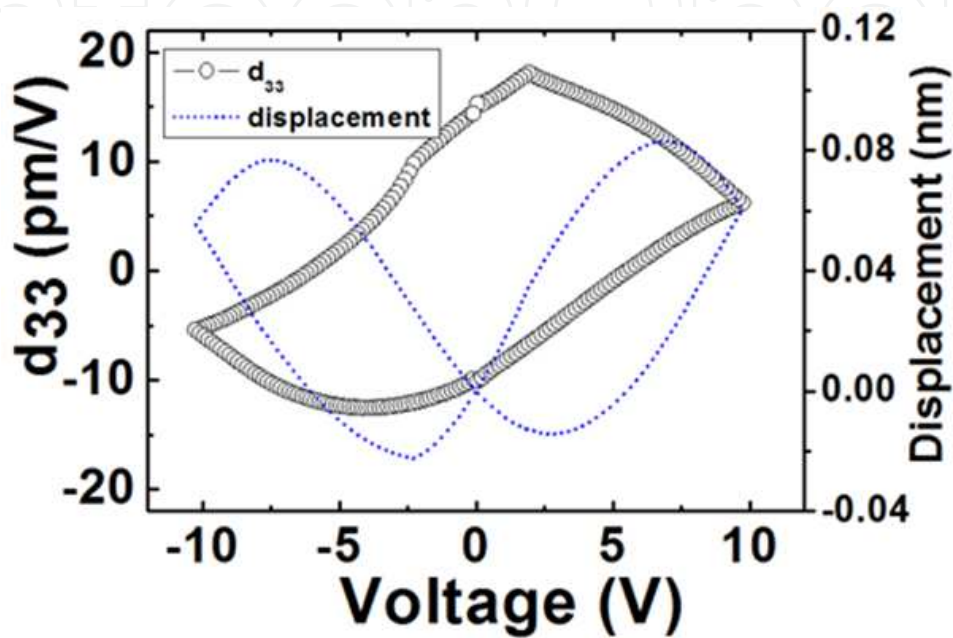


Figure 11. Butterfly and local piezoresponse hysteresis loops of YST110-4 film.

4. Conclusions

The novel multiferroic films of BFM and multilayered YST were successfully produced using PLD method. Both of them were characterized using PFM method. The special or excellent properties can often be found in the metastable materials. Through calculation, six structures are presumed in BFM with two magnetic configurations of G-AFM and FM. In YST films, the metastable SnTiO_3 phase was obtained. The improved ferroelectric properties were observed through increasing the layer numbers and more SnTiO_3 phase was stabilized, moreover, the change in contrast after bias is applied indicates a change in polarization direction and hence ferroelectric switching. Although it is a great challenge in obtaining the SnTiO_3 thin films, we believe that through the optimization of fabrication process and conditions, the single phase SnTiO_3 or multilayer films, or composite materials containing SnTiO_3 could become a new generation lead-free piezoelectric/ferroelectric material. For a thorough understanding of the mechanisms of the films, knowledge of the domain structures is a prerequisite, which is of crucial importance to increase and tune the functionality of multiferroic films.

Acknowledgements

The authors gratefully acknowledge Dr. Minoru Osada, Dr. Kazuya Terabe in NIMS and Prof. H. R. Zeng in Shanghai Institute of Ceramics for their experimental help and discussions. Part of this work was supported by JST ALCA and MEXT GRENE.

Author details

Hongyang Zhao^{1,2*}, Hideo Kimura¹, Qiwen Yao¹, Lei Guo¹, Zhenxiang Cheng³ and Xiaolin Wang³

*Address all correspondence to: zhao.hongyang@nims.go.jp

1 National Institute for Materials Science, Sengen 1-2-1, Tsukuba, Japan

2 Shanghai Institute of Ceramics, Chinese Academy of Sciences, Shanghai, China

3 Institute for Superconducting and Electronics Materials, University of Wollongong, Innovation Campus, Fairy Meadow, Australia

References

- [1] Hill N. A., Why are there so few magnetic ferroelectrics? *J. Phys. Chem. B* 2000; 104: 6694.
- [2] Fiebig M., T. Lottermoser, D. Frohlich, A. V. Goltsev, R. V. Pisarev, Observation of coupled magnetic and electric domains. *Nature* 2002; 419: 818.
- [3] Eerenstein W., Mathur N. D. and Scott J.F., Multiferroic and magnetoelectric materials. *Nature* 2006; 442: 759-765.
- [4] Singh M. K., Prellier W., Singh M. P., Katiyar R. S., Scott J. F., Spin-glass transition in single-crystal BiFeO₃. *Phys. Rev. B* 2008; 77: 144403.
- [5] J. Das, Y.Y. Song, M.Z. Wu, Electric-field control of ferromagnetic resonance in monolithic BaFe₁₂O₁₉-Ba_{0.5}Sr_{0.5}TiO₃ heterostructures. *J. Appl. Phys.* 2010; 108: 043911.
- [6] Konishi Y., Ohasawa M., Yonezawa Y., Tanimura Y., Chikyow T., Wakisaka T., Koinuma H., Miyamoto A., Kubo M., Sasata K., *Mater. Res. Soc. Symp. Proc.* 2003; 748 U3.13.1.
- [7] Uratani Y., Shishidou T., Oguchi T., First-Principles Study of Lead-Free Piezoelectric SnTiO₃. *Jpn. J. Appl. Phys.* 2008; 47 (9), 7735-7739.

- [8] Ha J. Y., Lin L. W., Jeong D. Y., Yoon S. J., Choi J. W., Improved Figure of Merit of (Ba,Sr)TiO₃-Based Ceramics by Sn Substitution. *Jpn. J. Appl. Phys.* 2009; 48 011402.
- [9] Suzuki S., Takeda T., Ando A., Takagi H., Ferroelectric phase transition in Sn²⁺ ions doped (Ba,Ca)TiO₃ ceramics. *Appl. Phys. Lett.* 2010; 96: 132903.
- [10] Marin L.W., Crane S.P., Chu Y-H., Holcomb M.B., Gajek M., Huijben M., Yang C-H., Balke N., Ramesh R., Multiferroics and magnetoelectrics: thin films and nanostructures. *J. Phys.: Condens. Matter* 2008; 20: 434220.
- [11] Yakel H.L., Koehler W.C., Bertaut E.F., Forrat E.F., On the crystal structure of the manganese(III) trioxides of the heavy lanthanides and yttrium. *Acta Crystallogr.* 1963; 16: 957-962.
- [12] Bertaut E. F., Pauthenet R., Mercier M., Proprietes magnetiques et structures du manganite d'yttrium. *Phys. Lett.* 1963; 7: 110-111.
- [13] Gruverman A, Kholkin A, Nanoscale Ferroelectrics: Processing, Characterization and Future Trends. *Rep. Prog. Phys.* 2006; 69: 2443-2474.
- [14] Catalan G., Noheda B, McAneney J, Sinnamon L J, Gregg J M Strain gradients in epitaxial ferroelectrics. *Phys. Rev. B* 2005; 72: 0201102.
- [15] Muralt P. J. *Micromech. Ferroelectric thin films for micro-sensors and actuators* *Micromech. Eng.* 2000;10: 136.
- [16] Terabe K, Nakamura M, Takekawa S, Kitamura K, Higuchi S, Gotoh Y, Cho Y Microscale to nanoscale ferroelectric domain and surface engineering of a near-stoichiometric LiNbO₃ crystal. *Appl. Phys. Lett.* 2003; 82: 433-435.
- [17] Fong D D, Stephenson G B, Streigger S K, Eastman J A, Auciello O, Fuoss P H, Thompson C Ferroelectricity in ultrathin perovskite films *Science*. 2004; 304: 1650.
- [18] Kholkin A L, Bdikin I K, Shvartsman V V, Orlova A, Kiselev D, Bogomolov V, *MRS Proc. E* 2005; 838: O7.6.
- [19] Catalan G., Bea H., Fusil S., Bibes M., Paruch P., Barthelemy A., Scott J. F., Fractal Dimension and Size Scaling of Domains in Thin Films of Multiferroic BiFeO₃. *Phys. Rev. Lett.* 2008; 100 027602.
- [20] Kalinin S V, Gruverman A and Bonnell D A Quantitative analysis of nanoscale switching in SrBi₂Ta₂O₉ thin films by piezoresponse force microscopy *Appl. Phys. Lett.* 2004; 85: 795-797
- [21] Tybell T, Paruch P, Giamarchi T, Triscone J-M Domain wall creep in epitaxial ferroelectric Pb(Zr_{0.2}Ti_{0.8})O₃ thin films 2002 *Phys. Rev. Lett.* 89 097601
- [22] Catalan G., Scott J. F., *Adv. Mater. Physics and Applications of Bismuth Ferrite.* 2009; 21, 2643.
- [23] Wang J., Neaton B. J., Zheng H., Nagarajan V., Ogale S. B., Liu B., Viehland D., Vaithyanathan V., Schlom D. G., Waghmare U. V., Spaldin N. A., Rabe K. M., Wutting

- M., Ramesh R., Epitaxial BiFeO₃ Multiferroic Thin Film Heterostructures. *Science* 2003; 299: 1719-1722.
- [24] Kimura T., Goto T., Shintani H., Ishizaka K., Arima T., Tokura Y., Magnetic Control of Ferroelectric Polarization. *Nature* 2003; 55: 426.
- [25] Cheng Z. X., Wang X. L., Dou S. X., Ozawa K., Kimura H., Improved ferroelectric properties in multiferroic BiFeO₃ thin films through La and Nbcodoping. *Phys. Rev. B* 2008; 77: 092101.
- [26] Spaldin N. A., Cheong S. W., Ramesh R., Multiferroics: Past, present, and future. *Physics Today* 2010; 63: 38-43.
- [27] T. Atou, H. Chiba, K. Ohoyama, Y. Yamaguchi, Y. Syono, Structure determination of ferromagnetic perovskite BiMnO₃. *J. Solid State Chem.* 1999: 145(2) 639-642.
- [28] T. Kimura, S. Kawamoto, I. Yamada, M. Azuma, M. Takano, Y. Tokura, Magnetocapacitance effect in multiferroic BiMnO₃, *Phys. Rev. B* 2003: 67 (R), 180401-180404.
- [29] Zhao, H.Y.; Kimura, H.; Cheng, Z.X.; Wang, X.L. & Nishida, T. Room temperature multiferroic properties of Nd:BiFeO₃/Bi₂FeMnO₆ bilayered films. *Appl. Phys. Lett.* 2009; 95(23)232904
- [30] Zhao, H.Y.; Kimura, H.; Cheng, Z.X.; Wang, X.L.; Ozawa, K. & Nishida, T. Magnetic characterization of Bi₂FeMnO₆ film grown on (100) SrTiO₃ substrate *Phys. Status Solidi RRL* 2010; 4(11) 314.
- [31] Bi, L.; Taussig, A.R.; Kim, H-S.; Wang, L.; Dionne, G.F.; Bono, D.; Persson, K.; Ceder, G. & Ross, C.A. Structural, magnetic, and optical properties of BiFeO₃ and Bi₂FeMnO₆ epitaxial thin films: An experimental and first-principles study. *Phys. Rev. B* 2008; 78 (10) 104106.
- [32] Zhao H Y, Kimura H.; Cheng Z.X.; Wang X.L., New multiferroic materials: Bi₂FeMnO₆. *Ferroelectrics-Material Aspects. InTech*; 2011.
- [33] Zhao H Y, Kimura H.; Cheng Z.X.; Wang X.L., Yao Q W, Osada M, Li B W, Nishida T A new multiferroic heterostructure of YMnO₃/SnTiO_{3+x}. *Scrip. Mater.* 2011; 65: 618-621
- [34] Harrison R.J., Redfern S.A.T., Smith R.I., Thermodynamics of the R $\bar{3}$ to R $\bar{3}$ phase transition in the ilmenite-hematite solid solution. *Am. Mineral.* 2000; 85: 1694-1705.
- [35] Leinenweber K., Utsmi W., Tsuchida Y., Yagi T., Kuita K., Unquenchable High-Pressure Perovskite Polymorphs Of MnSnO₃ And FeTiO₃. *Phys. Chem. Miner.* 1991; 18: 244.
- [36] Venkatesan S., Daumont C., Kooi B.J., Noheda B., J Hossoneff Th.M. De, Nanoscale domain evolution in thin films of multiferroic TbMnO₃. *Phys. Rev. B* 2009; 80: 214111.

